

# Recent progress of chiral three-nucleon forces and applications to nuclei and matter

Kai Hebeler

Mainz, June 24, 2024



MITP  
TOPICAL  
WORKSHOP

Uncertainty Quantification in Nuclear Physics  
June 24 – 28, 2024

<https://indico.mitp.uni-mainz.de/event/357>

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The banner features a red triangle on the left with the text 'MITP TOPICAL WORKSHOP'. The central part shows a scenic view of Mainz, Germany, with a large cathedral and a river. On the right, there is a stylized graphic of overlapping green and blue waves. The text 'Uncertainty Quantification in Nuclear Physics' and 'June 24 – 28, 2024' is positioned above the wave graphic. A small globe icon and a red arrow point to the URL 'https://indico.mitp.uni-mainz.de/event/357'. The MITP logo is located in the bottom left corner of the banner.



# Outline

## Uncertainties related to chiral three-nucleon interactions:

- I. Chiral expansion and different regularization schemes
- II. Fixing of low-energy couplings + SRG evolution
- III. Inclusion of 3NFs in many-body calculations

# 3NFs in different regularization schemes

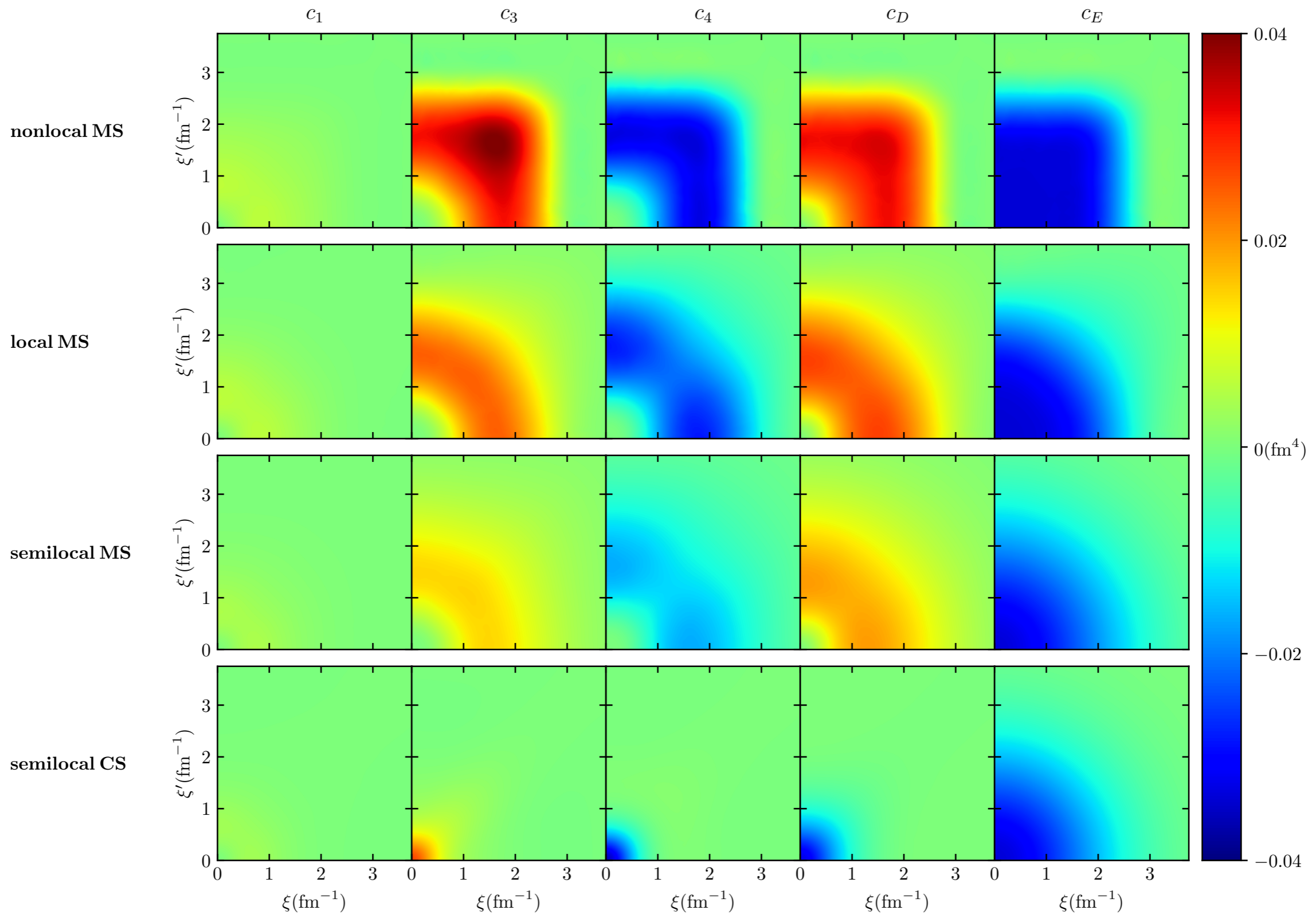
	momentum space	coordinate space
<b>nonlocal</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b><u>nonlocal MS</u></b>  $f_{\Lambda}^{\text{long}}(\mathbf{p}, \mathbf{q}) = \exp\left[-((\mathbf{p}^2 + 3/4 \mathbf{q}^2)/\Lambda^2)^n\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{p}, \mathbf{q}) = f_{\Lambda}^{\text{long}}(\mathbf{p}, \mathbf{q}) = f_R(\mathbf{p}, \mathbf{q})$  $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}}   \mathbf{p} \mathbf{q} \rangle = f_R(\mathbf{p}', \mathbf{q}') \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle f_R(\mathbf{p}, \mathbf{q})$	
<b>local</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b><u>local MS</u></b>  $f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = \exp\left[-(\mathbf{Q}_i^2/\Lambda^2)^2\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{Q}_i) = f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = f_{\Lambda}(\mathbf{Q}_i)$  $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}}   \mathbf{p} \mathbf{q} \rangle = \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle \prod_i f_R(\mathbf{Q}_i)$	<b><u>local CS</u></b>  $f_R^{\text{long}}(\mathbf{r}) = 1 - \exp\left[-(r^2/R^2)^n\right]$ $f_R^{\text{short}}(\mathbf{r}) = \exp\left[-(r^2/R^2)^n\right]$  $V_{3N}^{\pi, \text{reg}}(\mathbf{r}_{ij}) = f_R^{\text{long}}(\mathbf{r}_{ij}) V_{3N}^{\pi}(\mathbf{r}_{ij})$ $\delta^{\text{reg}}(\mathbf{r}_{ij}) = \alpha_n f_R^{\text{short}}(\mathbf{r}_{ij})$
<b>semilocal</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b><u>semilocal MS</u></b>  $f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = \exp\left[-(\mathbf{Q}_i^2 + m_{\pi}^2)/\Lambda^2\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{p}) = \exp\left[-\mathbf{p}^2/\Lambda^2\right]$  $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \pi}   \mathbf{p} \mathbf{q} \rangle = f_R^{\text{long}}(\mathbf{Q}_i) \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \delta}   \mathbf{p} \mathbf{q} \rangle = f_{\Lambda}^{\text{short}}(\mathbf{p}'_{\delta}) \langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\delta}   \mathbf{p} \mathbf{q} \rangle f_{\Lambda}^{\text{short}}(\mathbf{p}_{\delta})$	<b><u>semilocal CS</u></b>  $f_R^{\text{long}}(\mathbf{r}) = (1 - \exp\left[-r^2/R^2\right])^n$ $f_{\Lambda}^{\text{short}}(\mathbf{p}) = \exp\left[-\mathbf{p}^2/\Lambda^2\right]$  $V_{3N}^{\pi, \text{reg}}(\mathbf{r}_{ij}) = f_R^{\text{long}}(\mathbf{r}_{ij}) V_{3N}^{\pi}(\mathbf{r}_{ij})$ $\delta(\mathbf{r}_{ij}) \xrightarrow{FT} V_{3N}^{\delta}$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \delta}   \mathbf{p} \mathbf{q} \rangle = f_{\Lambda}^{\text{short}}(\mathbf{p}'_{\delta}) \langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\delta}   \mathbf{p} \mathbf{q} \rangle f_{\Lambda}^{\text{short}}(\mathbf{p}_{\delta})$

# 3NFs in different regularization schemes

	momentum space	coordinate space
<b>nonlocal</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b>nonlocal MS</b> <span style="float: right;">N<sup>2</sup>LO N<sup>3</sup>LO</span> $f_{\Lambda}^{\text{long}}(\mathbf{p}, \mathbf{q}) = \exp\left[-((\mathbf{p}^2 + 3/4 \mathbf{q}^2)/\Lambda^2)^n\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{p}, \mathbf{q}) = f_{\Lambda}^{\text{long}}(\mathbf{p}, \mathbf{q}) = f_R(\mathbf{p}, \mathbf{q})$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}}   \mathbf{p} \mathbf{q} \rangle = f_R(\mathbf{p}', \mathbf{q}') \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle f_R(\mathbf{p}, \mathbf{q})$	Violation of chiral symmetry at N3LO? Non-renormalizable? <i>Epelbaum, Krebs</i>
<b>local</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b>local MS</b> <span style="float: right;">N<sup>2</sup>LO N<sup>3</sup>LO</span> $f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = \exp\left[-(\mathbf{Q}_i^2/\Lambda^2)^2\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{Q}_i) = f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = f_{\Lambda}(\mathbf{Q}_i)$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}}   \mathbf{p} \mathbf{q} \rangle = \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle \prod_i f_R(\mathbf{Q}_i)$	<b>local CS</b> <span style="float: right;">N<sup>2</sup>LO N<sup>3</sup>LO</span> $f_R^{\text{long}}(\mathbf{r}) = 1 - \exp\left[-(r^2/R^2)^n\right]$ $f_R^{\text{short}}(\mathbf{r}) = \exp\left[-(r^2/R^2)^n\right]$ $V_{3N}^{\pi, \text{reg}}(\mathbf{r}_{ij}) = f_R^{\text{long}}(\mathbf{r}_{ij}) V_{3N}^{\pi}(\mathbf{r}_{ij})$ $\delta^{\text{reg}}(\mathbf{r}_{ij}) = \alpha_n f_R^{\text{short}}(\mathbf{r}_{ij})$ Development of large cutoff interactions for QMC calculations in progress <i>Tews et al.</i> → Reduction of cutoff artifacts.
<b>semilocal</b> <i>regulators:</i> long-range short-range  <i>regularization:</i>	<b>semilocal MS</b> <span style="float: right;">N<sup>2</sup>LO N<sup>3</sup>LO</span> $f_{\Lambda}^{\text{long}}(\mathbf{Q}_i) = \exp\left[-(\mathbf{Q}_i^2 + m_{\pi}^2)/\Lambda^2\right]$ $f_{\Lambda}^{\text{short}}(\mathbf{p}) = \exp\left[-\mathbf{p}^2/\Lambda^2\right]$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \pi}   \mathbf{p} \mathbf{q} \rangle = f_R^{\text{long}}(\mathbf{Q}_i) \langle \mathbf{p}' \mathbf{q}'   V_{3N}   \mathbf{p} \mathbf{q} \rangle$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \delta}   \mathbf{p} \mathbf{q} \rangle = f_{\Lambda}^{\text{short}}(\mathbf{p}'_{\delta}) \langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\delta}   \mathbf{p} \mathbf{q} \rangle f_{\Lambda}^{\text{short}}(\mathbf{p}_{\delta})$ Formal derivation 'basically' finished. Implementation in progress... (hard!)	<b>semilocal CS</b> <span style="float: right;">N<sup>2</sup>LO N<sup>3</sup>LO</span> $f_R^{\text{long}}(\mathbf{r}) = (1 - \exp\left[-r^2/R^2\right])^n$ $f_{\Lambda}^{\text{short}}(\mathbf{p}) = \exp\left[-\mathbf{p}^2/\Lambda^2\right]$ $V_{3N}^{\pi, \text{reg}}(\mathbf{r}_{ij}) = f_R^{\text{long}}(\mathbf{r}_{ij}) V_{3N}^{\pi}(\mathbf{r}_{ij})$ $\delta(\mathbf{r}_{ij}) \xrightarrow{FT} V_{3N}^{\delta}$ $\langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\text{reg}, \delta}   \mathbf{p} \mathbf{q} \rangle = f_{\Lambda}^{\text{short}}(\mathbf{p}'_{\delta}) \langle \mathbf{p}' \mathbf{q}'   V_{3N}^{\delta}   \mathbf{p} \mathbf{q} \rangle f_{\Lambda}^{\text{short}}(\mathbf{p}_{\delta})$ Derivation of consistently regularised currents at this order very hard, switched to SMS



# Illustration of 3NFs in different regularization schemes



$$\xi^2 = p^2 + 3/4q^2 \quad \tan \theta = p/(\sqrt{3}/2q) = \frac{\pi}{4}$$

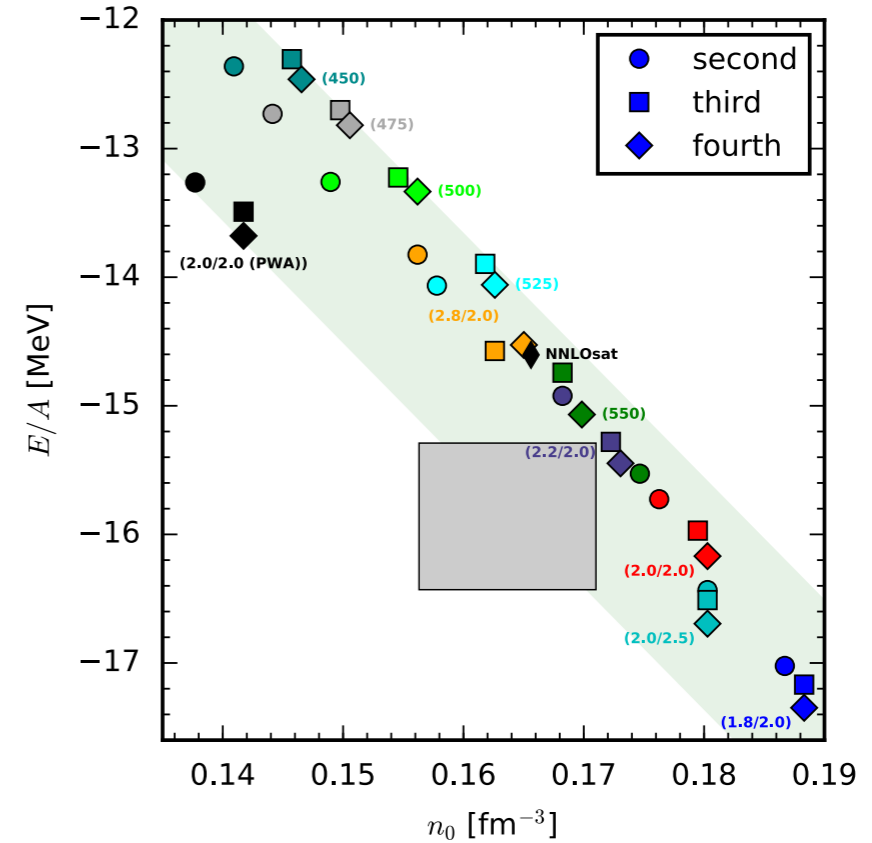
## Uncertainty II: Fitting of low-energy couplings + SRG evolution

- choice of observables
- separate or simultaneous fits of NN and 3N LECs?
- using few-body and/or many-body observables?
- using bare or low-resolution interactions for fits?
- low resolution interactions:
  - ❖ evolve NN interactions to lower scales via the RG
  - ❖ fit the 3N LECs at a low cutoff scale ( $\Lambda_{3N} \sim 2 \text{ fm}^{-1}$ )
  - ❖ e.g., use  ${}^3\text{H}$  binding energy and  ${}^4\text{He}$  radius to fix  $c_D$  and  $c_E$

$\Lambda$ or $\lambda/\Lambda_{3NF}$ [ $\text{fm}^{-1}$ ]	$V_{\text{low } k}$		SRG	
	$c_D$	$c_E$	$c_D$	$c_E$
1.8/2.0 (EM $c_i$ 's)	+1.621	-0.143	+1.264	-0.120
2.0/2.0 (EM $c_i$ 's)	+1.705	-0.109	+1.271	-0.131
2.0/2.5 (EM $c_i$ 's)	+0.230	-0.538	-0.292	-0.592
2.2/2.0 (EM $c_i$ 's)	+1.575	-0.102	+1.214	-0.137
2.8/2.0 (EM $c_i$ 's)	+1.463	-0.029	+1.278	-0.078
2.0/2.0 (EGM $c_i$ 's)	-4.381	-1.126	-4.828	-1.152
2.0/2.0 (PWA $c_i$ 's)	-2.632	-0.677	-3.007	-0.686

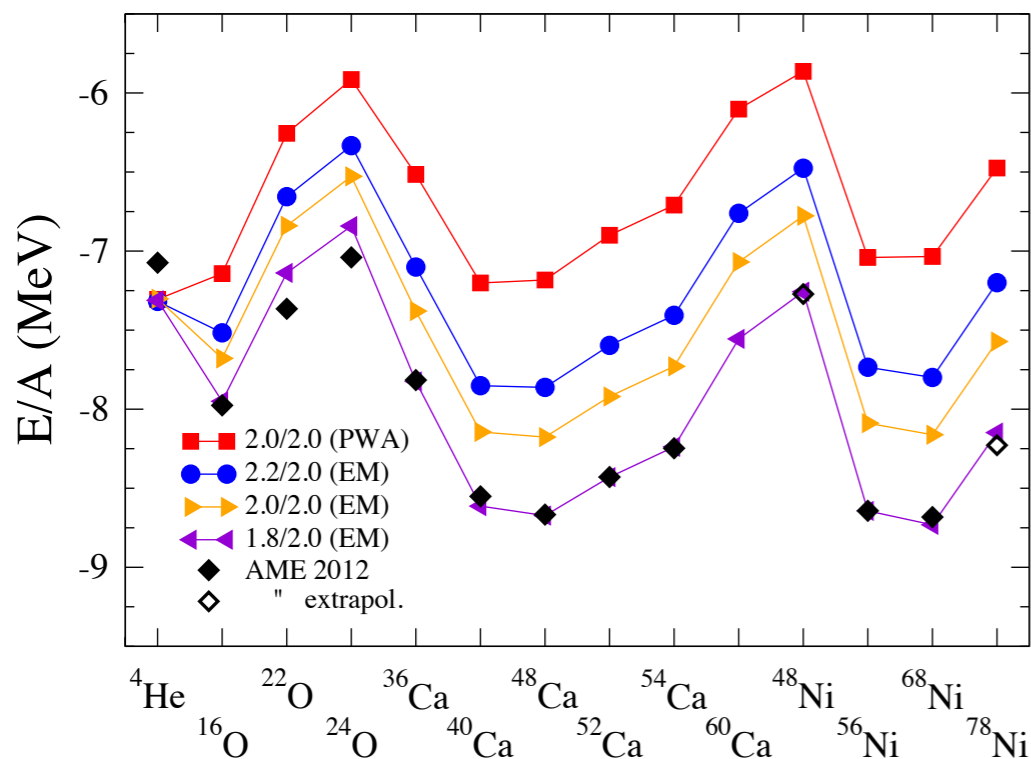
# Low-resolution fits versus consistent NN+3N evolution

NN SRG evolution + 3N fits					
$\lambda_{\text{SRG}} \text{ (fm}^{-1}\text{)}$	$\Lambda_{3\text{NF}} \text{ (fm}^{-1}\text{)}$	$c_D$	$c_E$	$r_{3\text{H}} \text{ (fm)}$	$E_{4\text{He}} \text{ (MeV)}$
$\infty$	2.0	+1.5	0.114	1.601	-28.64(4)
2.8	2.0 [116]	+1.278	-0.078	1.604	-28.75(2)
2.6	2.0	+1.26	-0.099	1.605	-28.77(2)
2.4	2.0	+1.265	-0.115	1.606	-28.80(2)
2.2	2.0 [116]	+1.214	-0.137	1.608	-28.86(2)
2.0	2.0 [116]	+1.271	-0.131	1.612	-28.95(2)
1.8	2.0 [116]	+1.264	-0.120	1.617	-29.11(2)
1.6	2.0	+1.25	-0.075	1.626	-29.42(2)
$\infty$	2.5	-1.45	-0.633	1.604	-28.65(4)
2.8	2.5	-1.35	-0.735	1.606	-28.84(2)
2.6	2.5	-1.2	-0.75	1.606	-28.85(2)
2.4	2.5	-1.0	-0.725	1.607	-28.89(2)
2.2	2.5	-0.7	-0.675	1.609	-28.95(2)
2.0	2.5 [116]	-0.292	-0.592	1.612	-29.05(2)
1.8	2.5	0.05	-0.503	1.617	-29.21(2)
1.6	2.5	0.55	-0.353	1.626	-29.48(2)

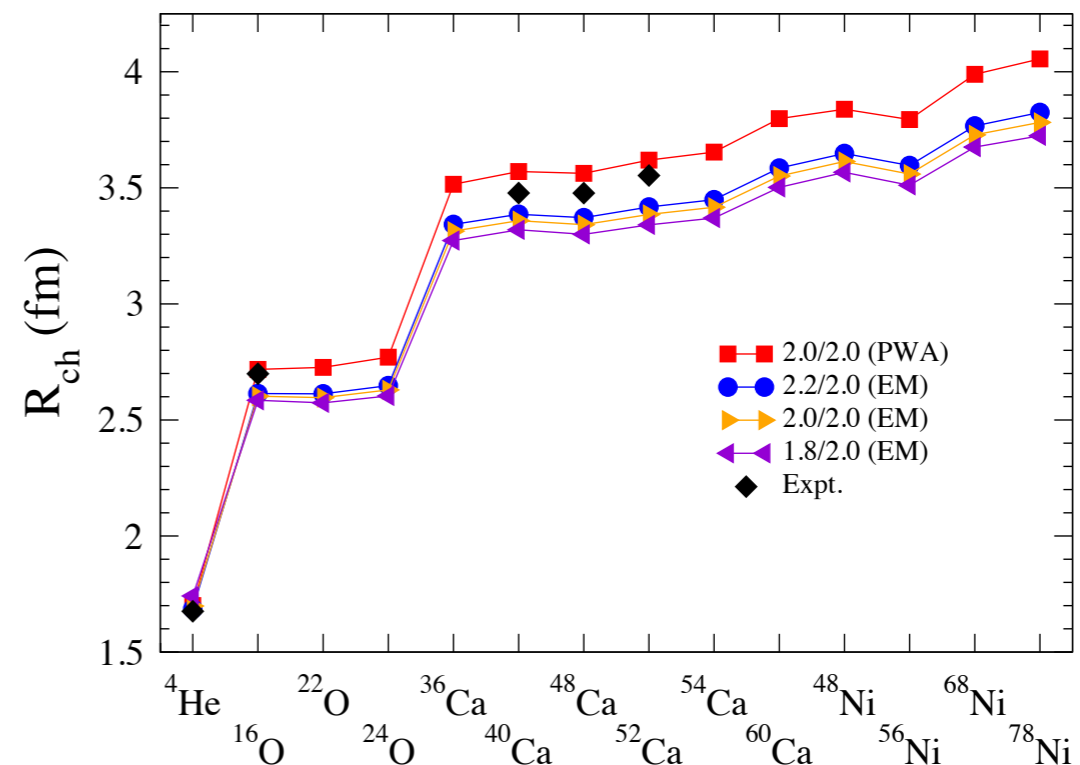


KH, Bogner, Furnstahl, Nogga, Schwenk, PRC 83, 031301 (2011)  
 KH, Phys. Rept. 890, 1 (2021)

Drischler, KH, Schwenk, PRL 122, 042501 (2019)

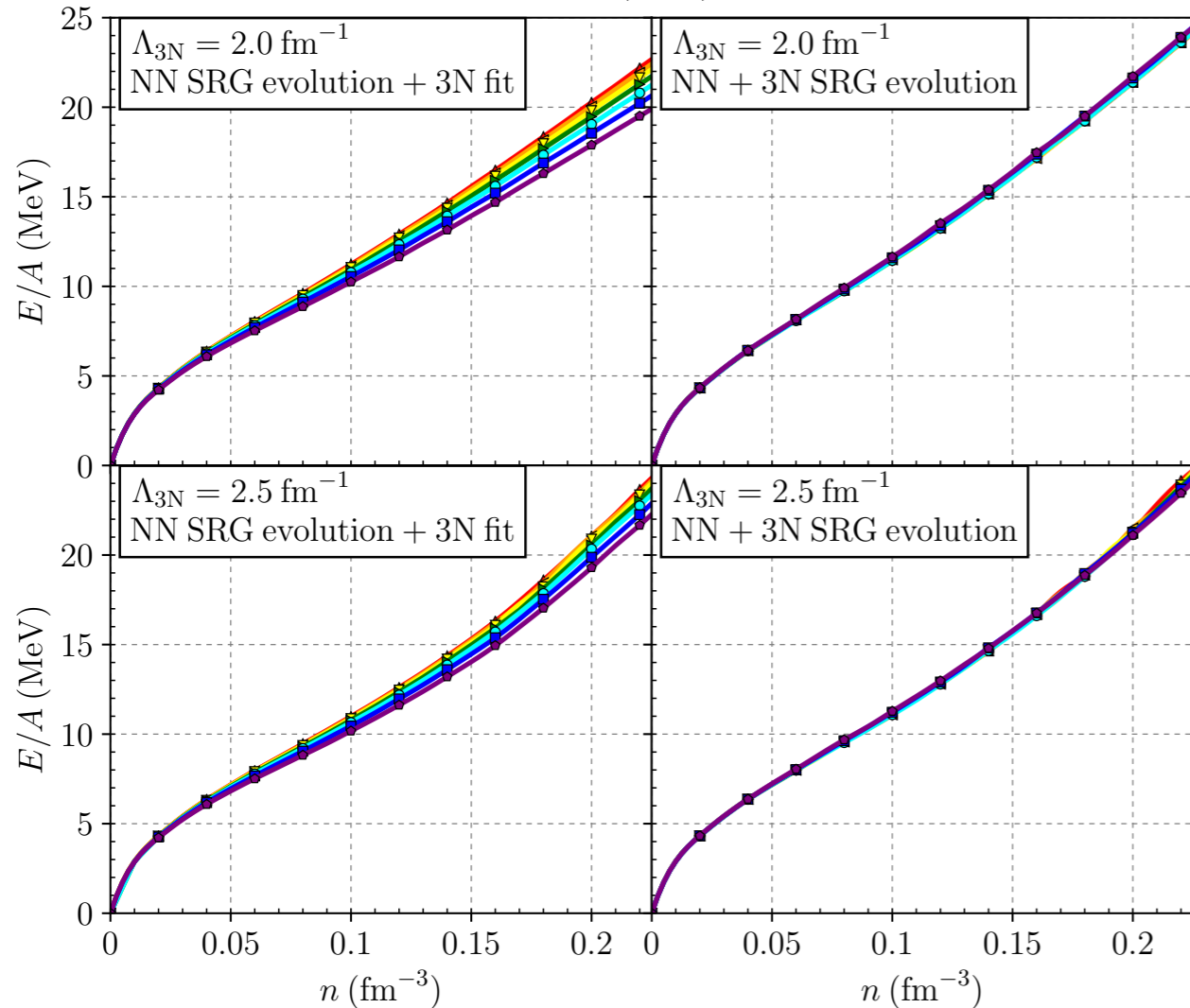
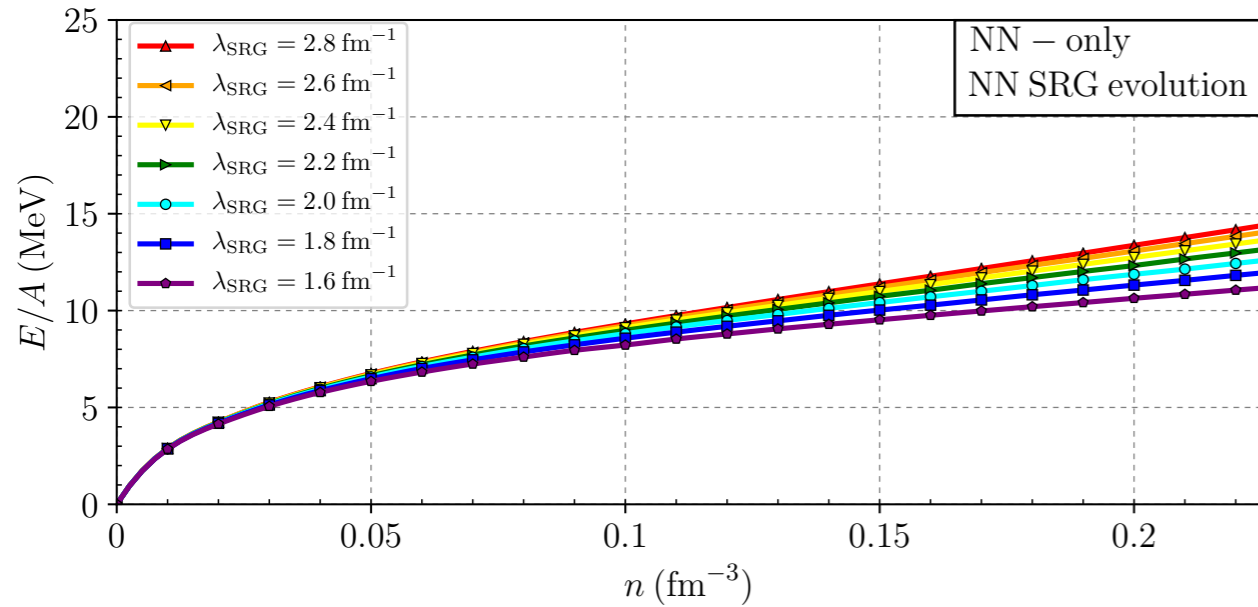


Simonis, Stroberg, KH, Holt, Schwenk, PRC 96, 014303 (2017)



# Low-resolution fits versus consistent NN+3N evolution

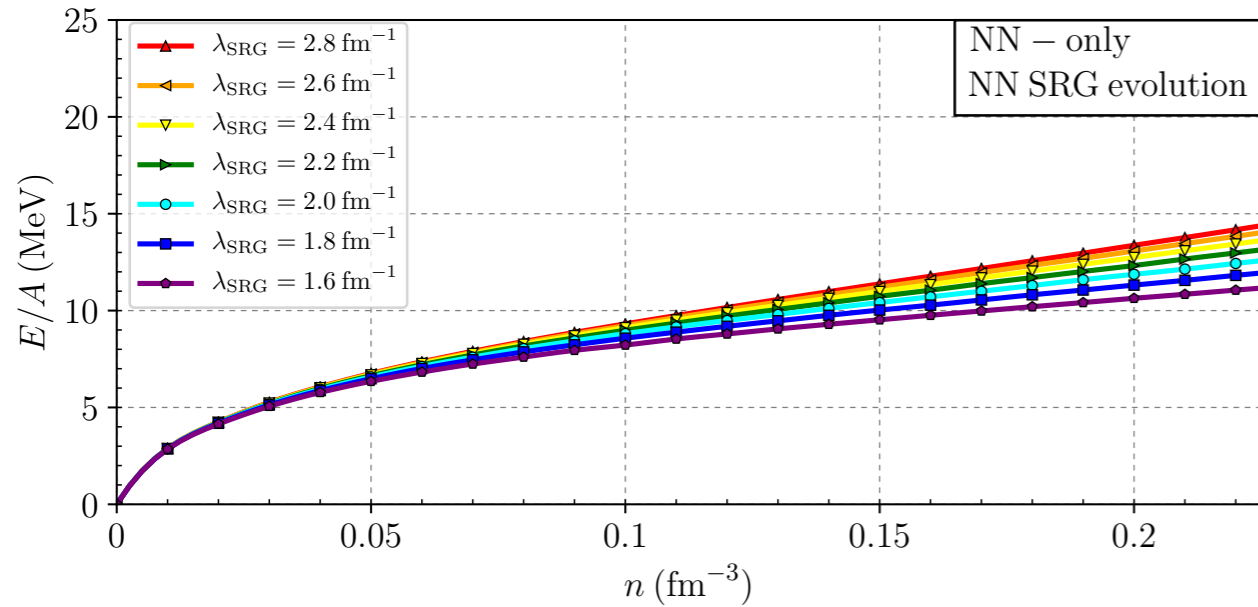
pure neutron matter (PNM)



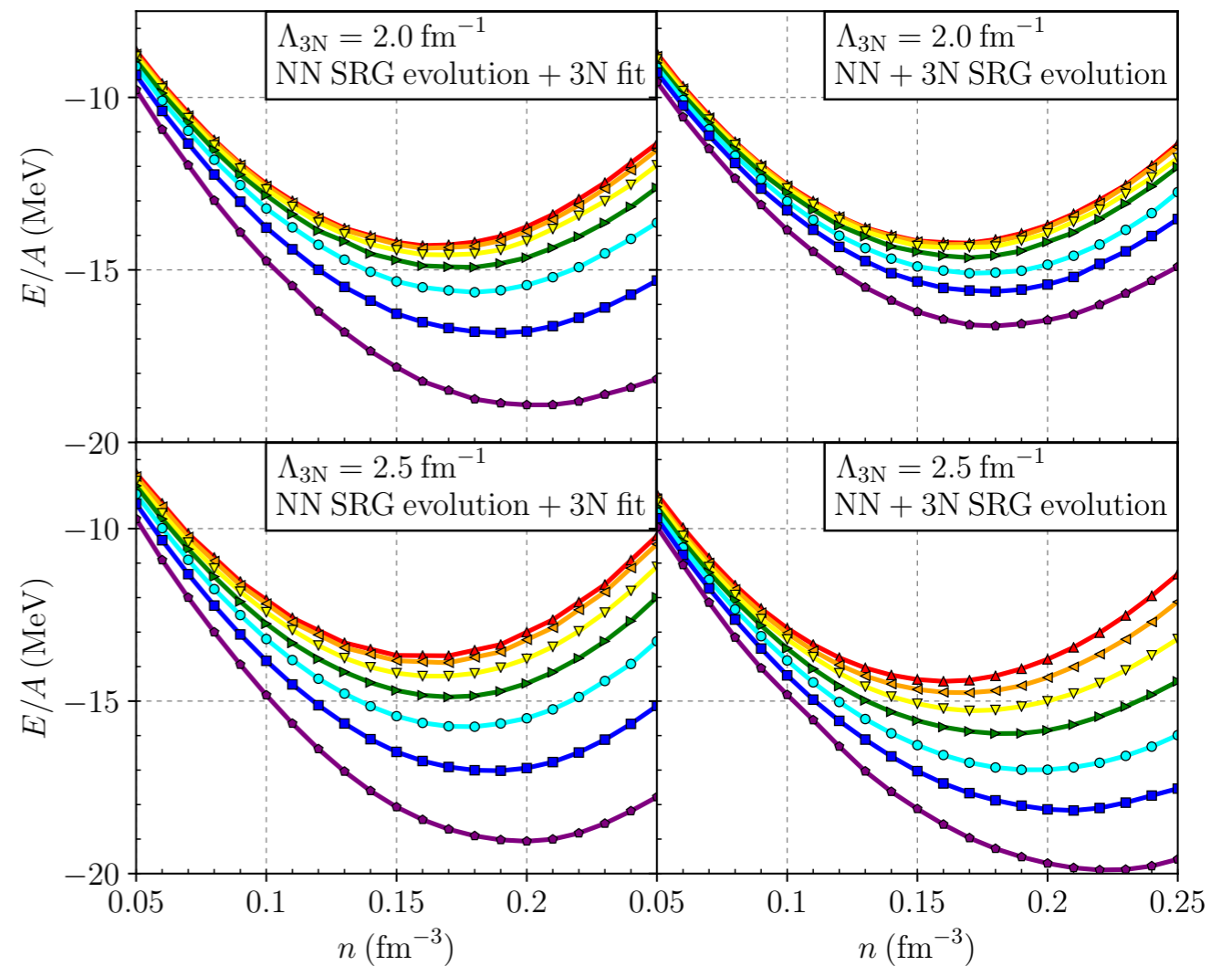
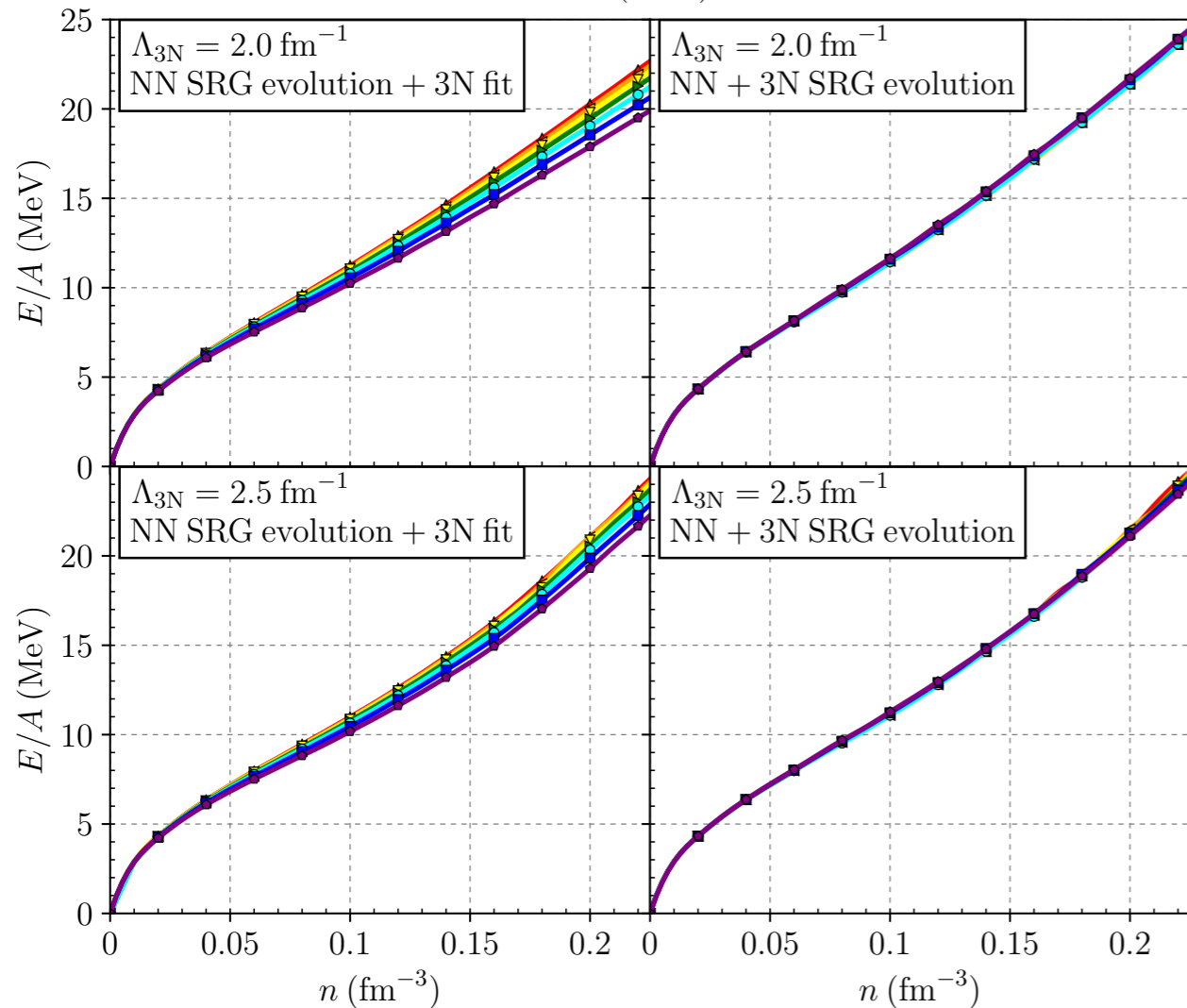
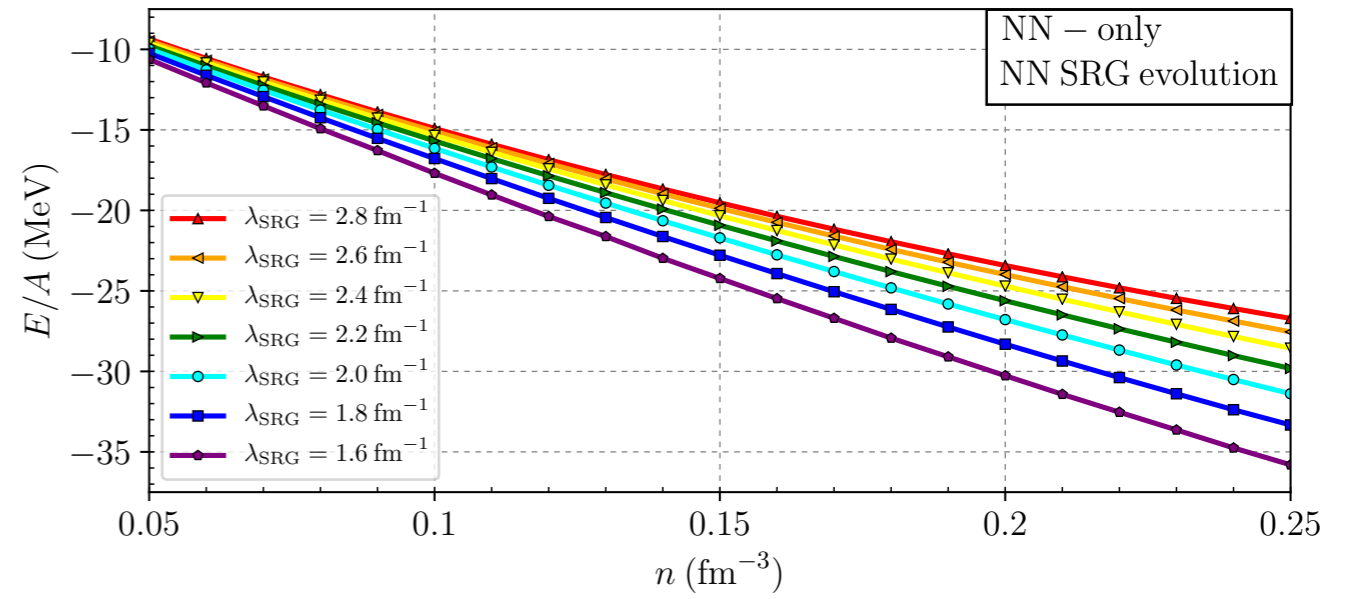
		NN+3N SRG evolution		
$\lambda_{\text{SRG}}$ ( $\text{fm}^{-1}$ )	$\Lambda_{3\text{NF}}$ ( $\text{fm}^{-1}$ )	$E_{3\text{H}}$ (MeV)	$r_{3\text{H}}$ (fm)	$E_{4\text{He}}$ (MeV)
$\infty$	2.0	-8.482	1.601	-28.64(4)
2.8	2.0 [116]	-8.482	1.605	-28.72(2)
2.6	2.0	-8.481	1.606	-28.73(2)
2.4	2.0	-8.481	1.608	-28.73(2)
2.2	2.0 [116]	-8.480	1.611	-28.74(2)
2.0	2.0 [116]	-8.479	1.615	-28.75(2)
1.8	2.0 [116]	-8.478	1.622	-28.76(2)
1.6	2.0	-8.476	1.635	-28.79(2)
<hr/>				
$\infty$	2.5	-8.482	1.604	-28.65(4)
2.8	2.5	-8.482	1.608	-28.75(2)
2.6	2.5	-8.482	1.609	-28.76(2)
2.4	2.5	-8.482	1.610	-28.77(2)
2.2	2.5	-8.481	1.613	-28.77(2)
2.0	2.5 [116]	-8.481	1.617	-28.77(2)
1.8	2.5	-8.480	1.625	-28.77(2)
1.6	2.5	-8.478	1.638	-28.77(2)

# Low-resolution fits versus consistent NN+3N evolution

pure neutron matter (PNM)



symmetric nuclear matter (SNM)

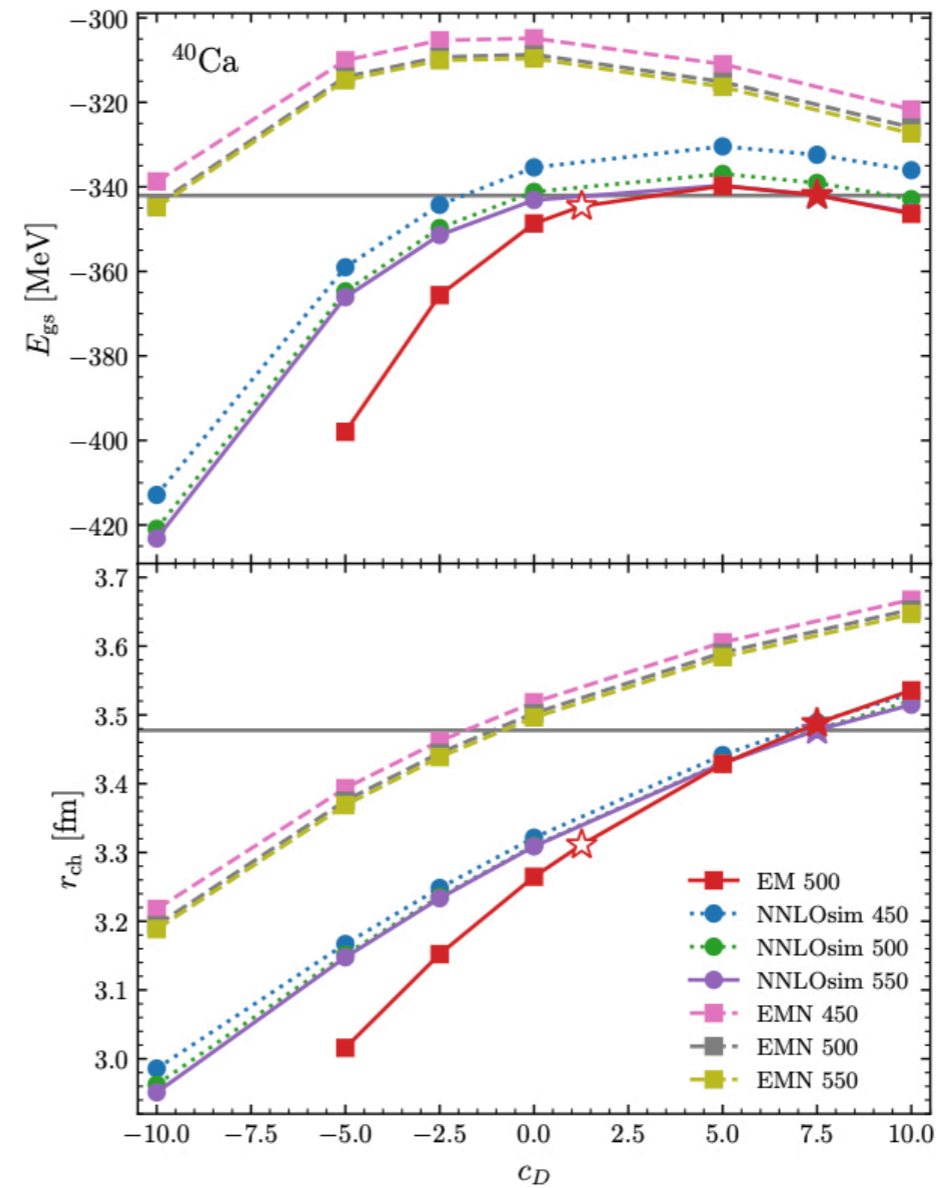
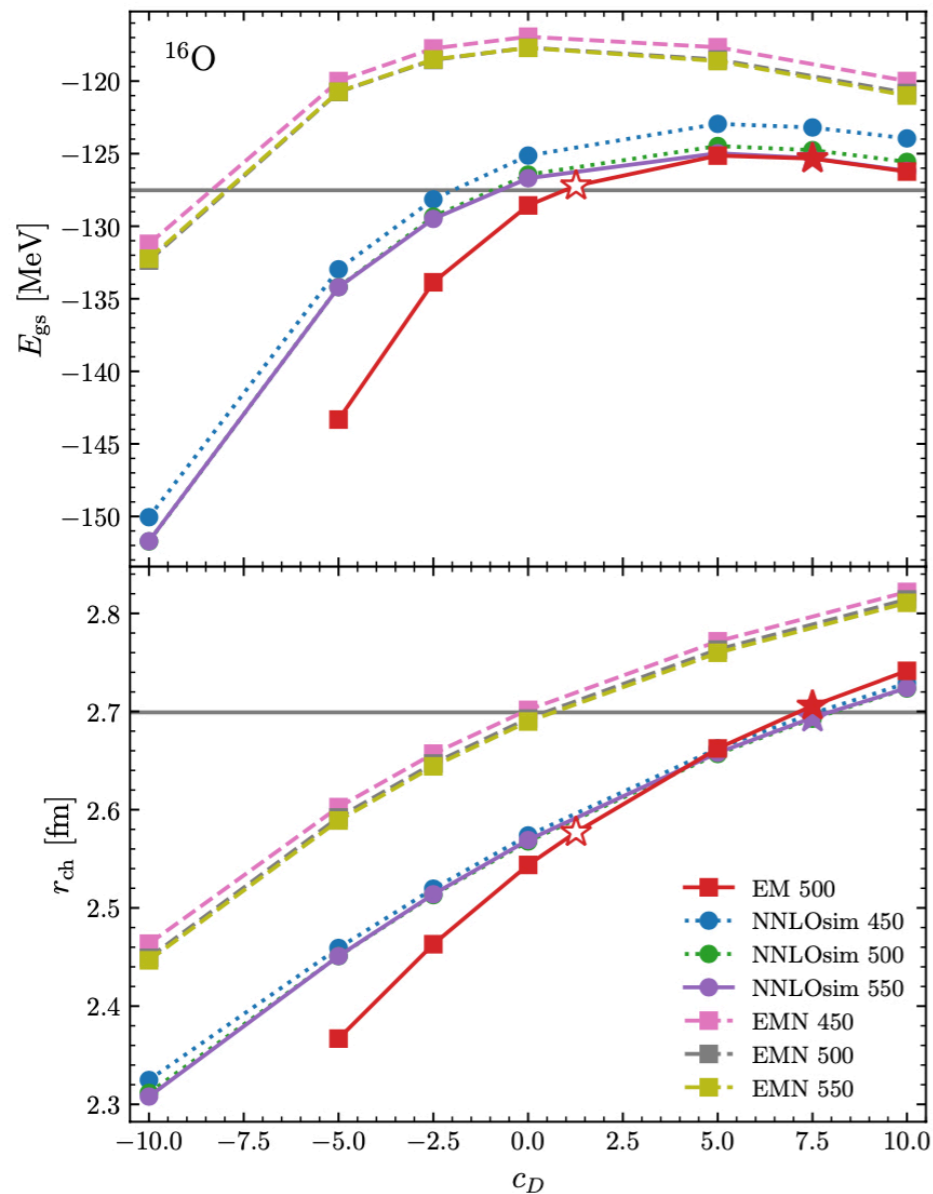






# New and improved 'magic' interactions

Arthuis, KH, Schwenk, arXiv:2401.06675



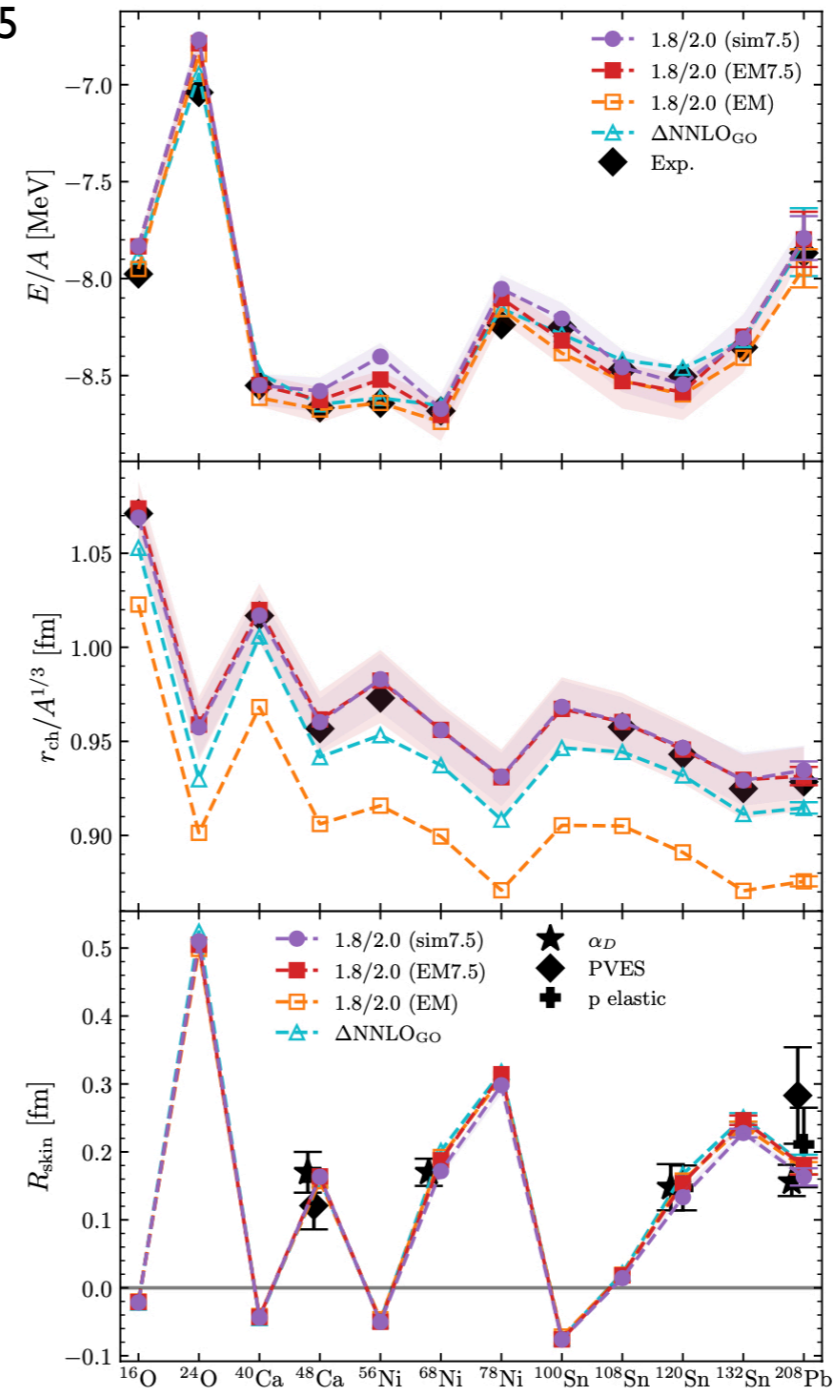
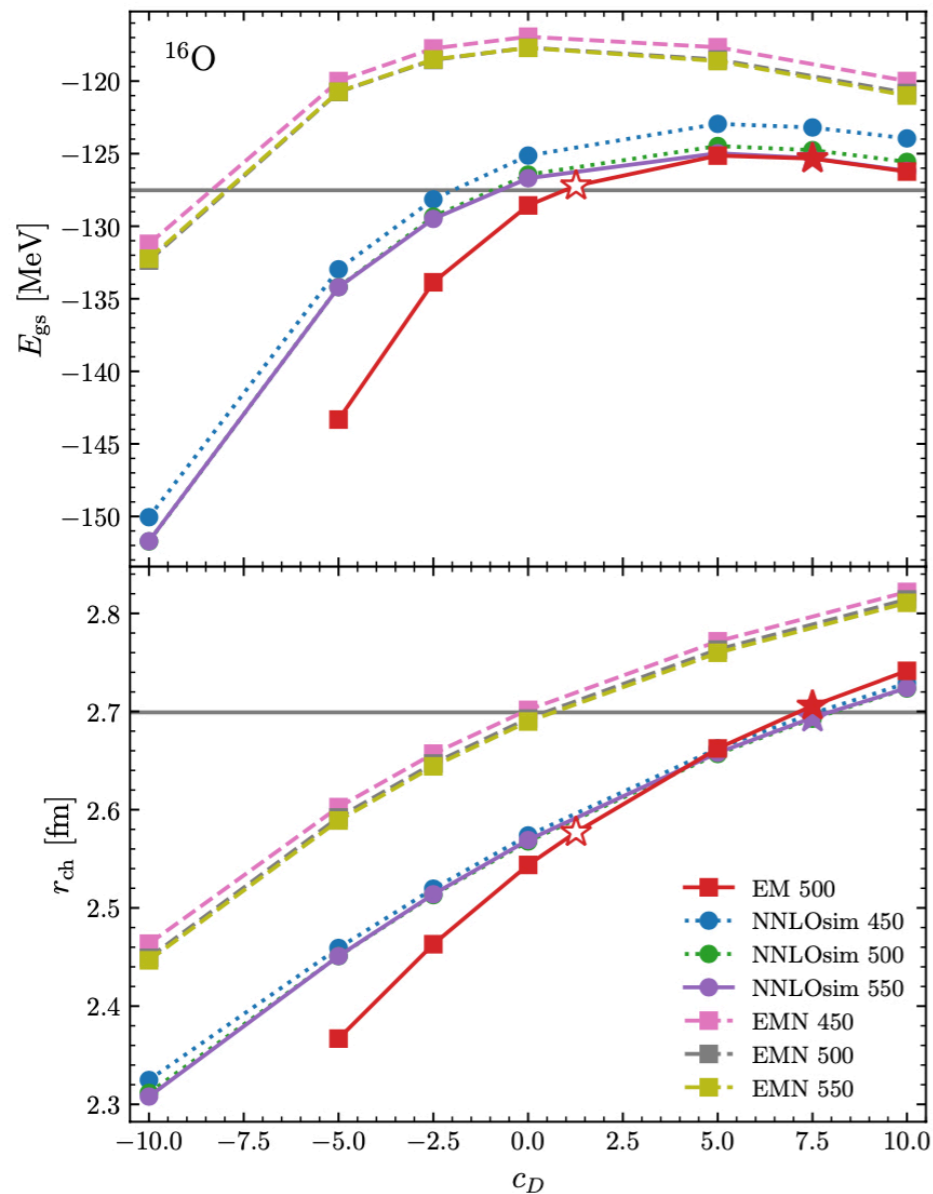
- Fit to  $c_D$  and  $c_E$  to  $^3\text{H}$  energy and **both** binding energies and radius of  $^{16}\text{O}$  using EM500 and NNLO<sub>sim</sub> 550 low-resolution NN interactions
- Simultaneous reproduction of experimental binding energies and charge radii  
*see talk by Achim for more details and results*





# New and improved 'magic' interactions

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- Fit to  $c_D$  and  $c_E$  to  $^3\text{H}$  energy and **both** binding energies and radius of  $^{16}\text{O}$  using EM500 and NNLO<sub>sim</sub> 550 low-resolution NN interactions
- Simultaneous reproduction of experimental binding energies and charge radii  
see *talk by Achim for more details and results*

# Uncertainty III: Inclusion of 3NF in many-body calculations

A general Hamiltonian consisting of kinetic energy, NN and 3N interactions:

$$\hat{H} = \hat{T}_{\text{rel}} + \hat{V}_{\text{NN}} + \hat{V}_{\text{3N}}$$

$$\hat{T}_{\text{rel}} = \sum_{ij} \langle i|T|j\rangle \hat{a}_i^\dagger \hat{a}_j,$$

$$\hat{V}_{\text{NN}} = \frac{1}{(2!)^2} \sum_{ijkl} \langle ij|V_{\text{NN}}^{\text{as}}|kl\rangle \hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_l \hat{a}_k,$$

$$\hat{V}_{\text{3N}} = \frac{1}{(3!)^2} \sum_{ijklmn} \langle ijk|V_{\text{3N}}^{\text{as}}|lmn\rangle \hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_k^\dagger \hat{a}_n \hat{a}_m \hat{a}_l$$

can be exactly rewritten in a normal-ordered form:

$$\Gamma_{\text{HF}}^{(0)} = \sum_i n_i \langle i|T|i\rangle + \frac{1}{2} \sum_{ij} n_i n_j \langle ij|V_{\text{NN}}^{\text{as}}|ij\rangle + \frac{1}{6} \sum_{ijk} n_i n_j n_k \langle ijk|V_{\text{3N}}^{\text{as}}|ijk\rangle,$$

$$\hat{\Gamma}_{\text{HF}}^{(1)} = \sum_{ij} \left[ \langle i|T|j\rangle + \sum_k n_k \langle ik|V_{\text{NN}}^{\text{as}}|jk\rangle + \frac{1}{2} \sum_{kl} n_k n_l \langle ikl|V_{\text{3N}}^{\text{as}}|jkl\rangle \right] N(\hat{a}_i^\dagger \hat{a}_j),$$

$$\hat{\Gamma}_{\text{HF}}^{(2)} = \sum_{ijkl} \left[ \langle ij|V_{\text{NN}}^{\text{as}}|kl\rangle + \sum_m n_m \langle ijm|V_{\text{3N}}^{\text{as}}|klm\rangle \right] N(\hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_l \hat{a}_k),$$

$$\hat{\Gamma}_{\text{HF}}^{(3)} = \sum_{ijklmn} \langle ijk|V_{\text{3N}}^{\text{as}}|lmn\rangle N(\hat{a}_i^\dagger \hat{a}_j^\dagger \hat{a}_k^\dagger \hat{a}_n \hat{a}_m \hat{a}_l).$$

$$\hat{H} = \Gamma_{\text{HF}}^{(0)} + \hat{\Gamma}_{\text{HF}}^{(1)} + \hat{\Gamma}_{\text{HF}}^{(2)} + \hat{\Gamma}_{\text{HF}}^{(3)}$$

# Traditional normal ordering framework for 3N interactions

## 1. transformation to Jacobi HO basis plus antisymmetrization

$$\langle p'q'\alpha' | V_{3N}^{(i),\text{reg}} | pq\alpha \rangle \rightarrow \langle N'n'\alpha' | V_{3N}^{(\text{as},\text{reg})} | Nn\alpha \rangle$$

## 2. transformation to single particle basis

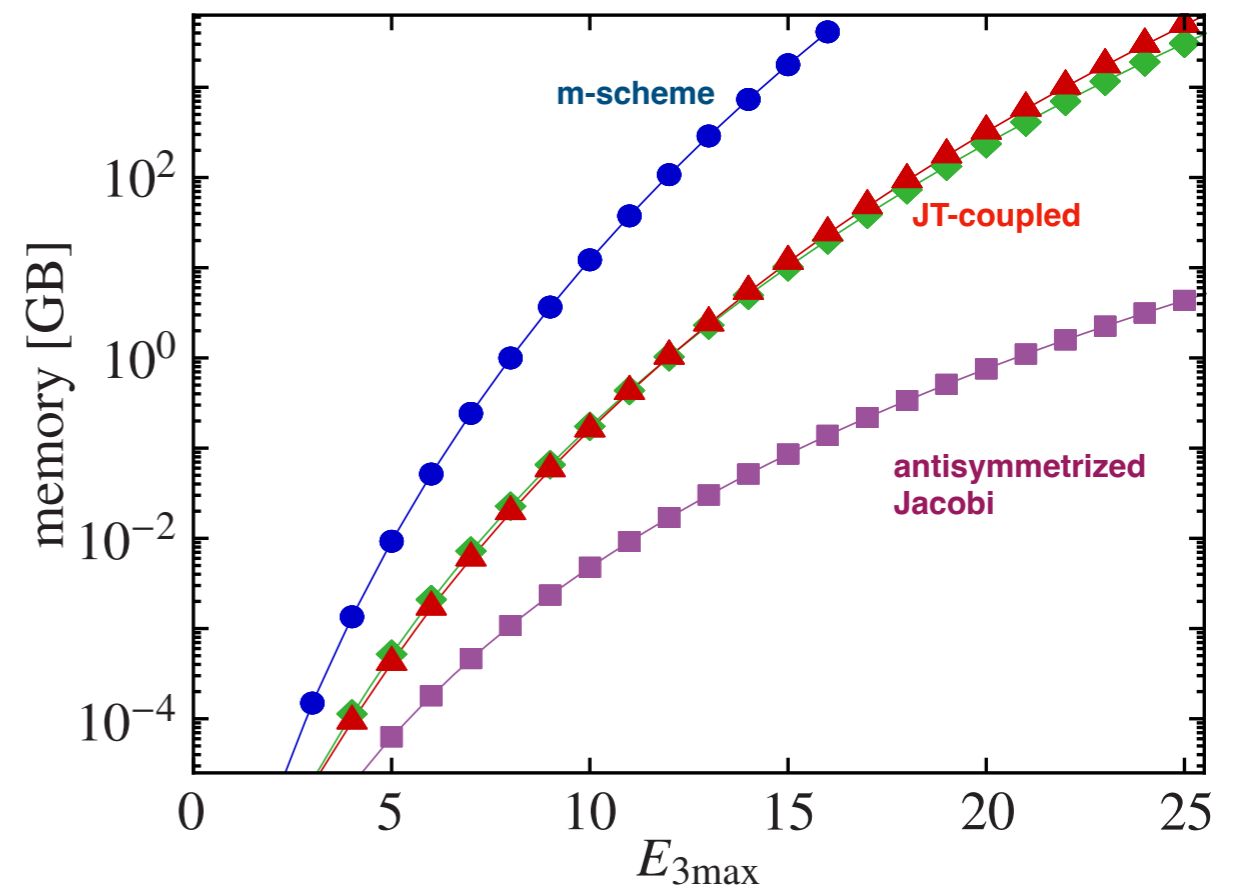
$$\langle N'n'\alpha' | V_{3N}^{\text{as},\text{reg}} | Nn\alpha \rangle \rightarrow \langle 1'2'3' | V_{3N}^{\text{as},\text{reg}} | 123 \rangle$$

## 3. Normal ordering with respect to some reference state

$$\langle 1'2' | \bar{V} | 12 \rangle = \sum_3 \bar{n}_3 \langle 1'2'3 | V_{3N}^{\text{as}} | 123 \rangle$$

- severe memory limitations for handling of single-particle matrix elements with increasing  $E_{3\text{max}}$
- Significant optimisations possible when storing only those matrix elements needed for normal ordering

Miyagi at al., PRC 105, 1 (2022)



Roth at al., PRC 90 024325 (2014)

# Novel normal ordering framework for 3N interactions

1. Use momentum space and expand reference state in HO basis:

$$\begin{aligned}\langle \mathbf{k}'_1 \mathbf{k}'_2 | \bar{V} | \mathbf{k}_1 \mathbf{k}_2 \rangle &= \sum_{n_3 l_3 m_3} \bar{n}_3 \langle \mathbf{k}'_1 \mathbf{k}'_2 \gamma_3 | V_{3N}^{\text{as}} | \mathbf{k}_1 \mathbf{k}_2 \gamma_3 \rangle \\ &= \int d\mathbf{k}_3 d\mathbf{k}'_3 \langle \mathbf{k}'_1 \mathbf{k}'_2 \mathbf{k}'_3 | V_{3N}^{\text{as}} | \mathbf{k}_1 \mathbf{k}_2 \mathbf{k}_3 \rangle \sum_{n_3 l_3 m_3} \bar{n}_3 \langle \gamma_3 | \mathbf{k}'_3 \rangle \langle \mathbf{k}_3 | \gamma_3 \rangle\end{aligned}$$

2. Rewrite interaction in Jacobi momentum basis:

$$\langle \mathbf{p}' \mathbf{P}' | \bar{V} | \mathbf{p} \mathbf{P} \rangle = \int d\mathbf{k}_3 d\mathbf{k}'_3 \langle \mathbf{p}' \mathbf{q}' | V_{3N}^{\text{as}} | \mathbf{p} \mathbf{q} \rangle \delta(\mathbf{P} + \mathbf{k}_3 - \mathbf{P}' - \mathbf{k}'_3) \sum_{n_3 l_3 m_3} \bar{n}_3 \langle \gamma_3 | \mathbf{k}'_3 \rangle \langle \mathbf{k}_3 | \gamma_3 \rangle$$

3. Decomposition in Jacobi partial wave momentum states:

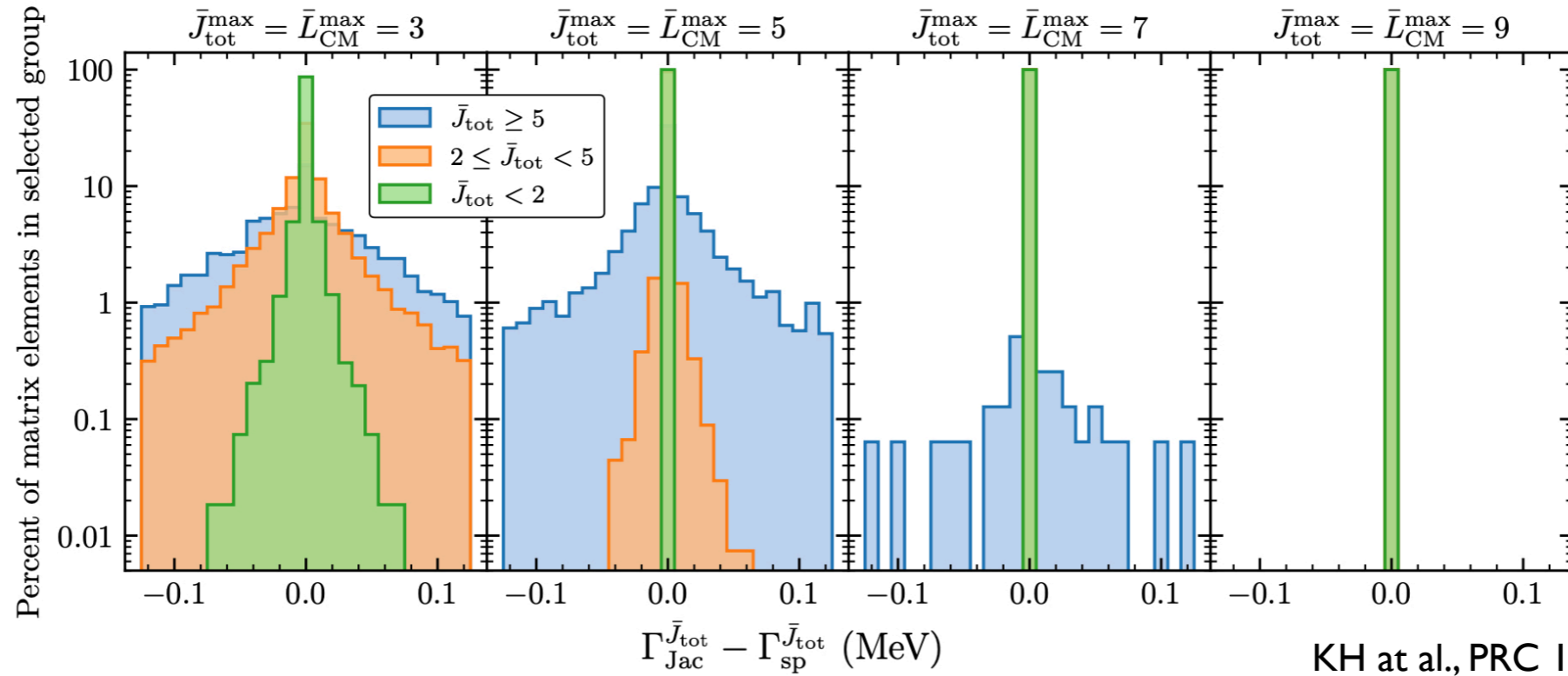
$$\begin{aligned}\langle p' P' L' M' L'_{cm} M'_{cm} | \bar{V} | p P L M L_{cm} M_{cm} \rangle \\ = \int d\hat{\mathbf{p}} d\hat{\mathbf{P}} d\hat{\mathbf{p}}' d\hat{\mathbf{P}}' Y_{L'_{cm} M'_{cm}}^*(\hat{\mathbf{P}}') Y_{L' M'}^*(\hat{\mathbf{p}}') \langle \mathbf{p}' \mathbf{P}' | \bar{V} | \mathbf{p} \mathbf{P} \rangle Y_{L_{cm} M_{cm}}(\hat{\mathbf{P}}) Y_{LM}(\hat{\mathbf{p}})\end{aligned}$$

4. transform matrix elements to Jacobi HO basis

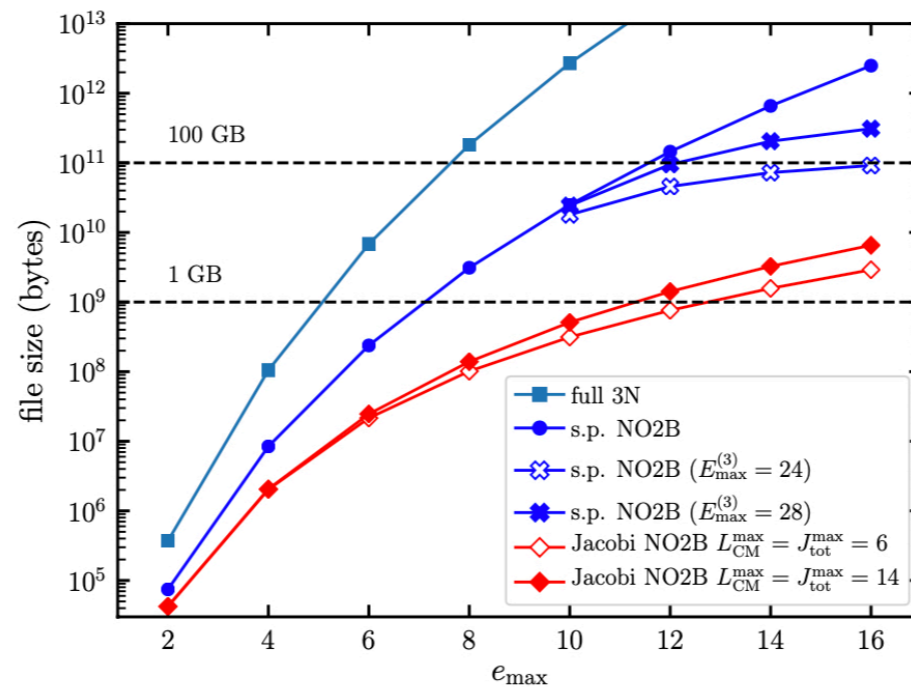
$$\begin{aligned}\langle p' P' L' M' L'_{cm} M'_{cm} | \bar{V} | p P L M L_{cm} M_{cm} \rangle \\ \rightarrow \langle n'_p N'_P L' M' L'_{cm} M'_{cm} | \bar{V} | n_p N_P L M L_{cm} M_{cm} \rangle\end{aligned}$$

5. transformation to single-particle HO basis via generalized Talmi transformation (taking into account  $L_{cm}$  dependence)

# Novel normal ordering framework for 3N interactions



KH at al., PRC 107, 024310 (2023)

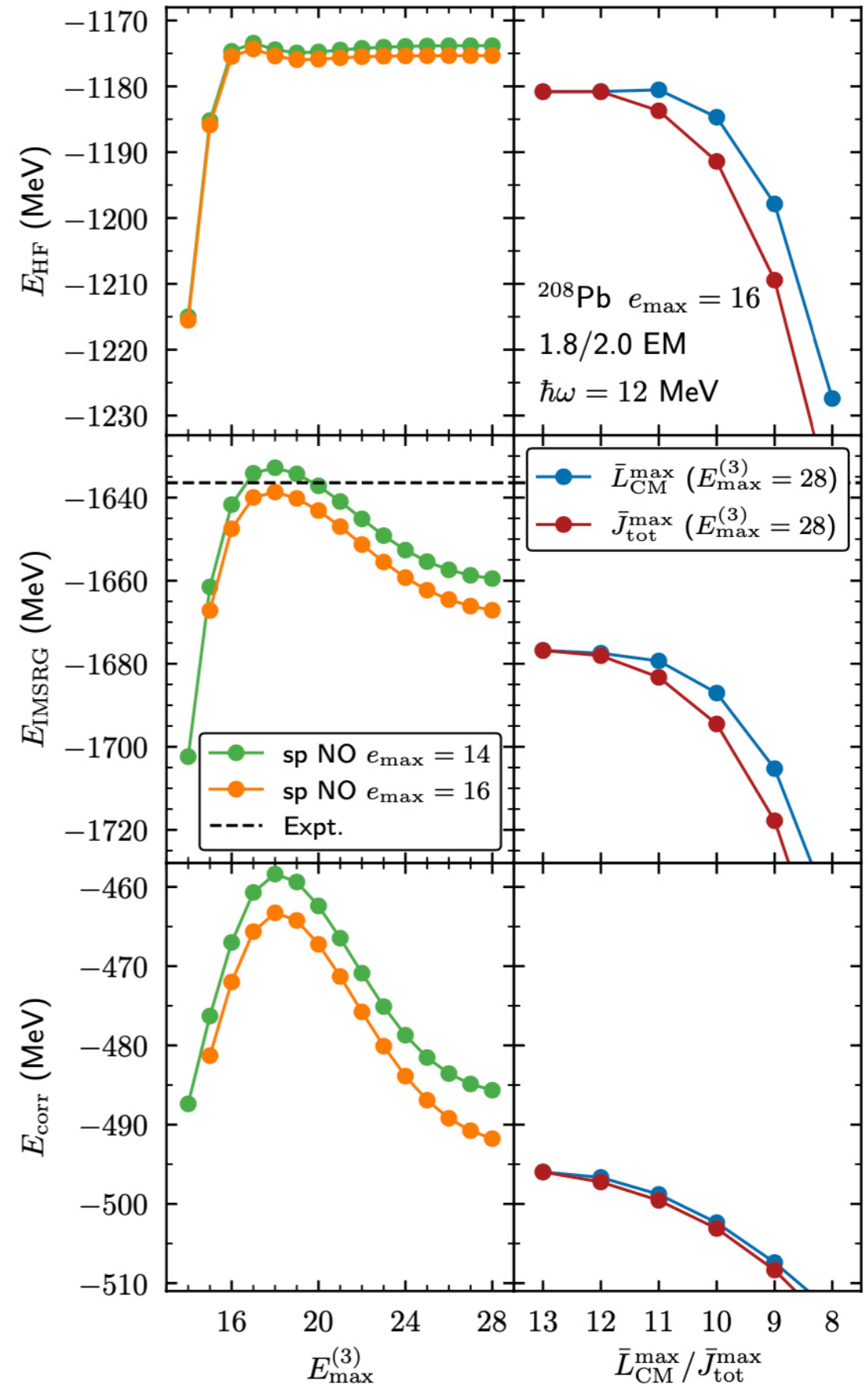
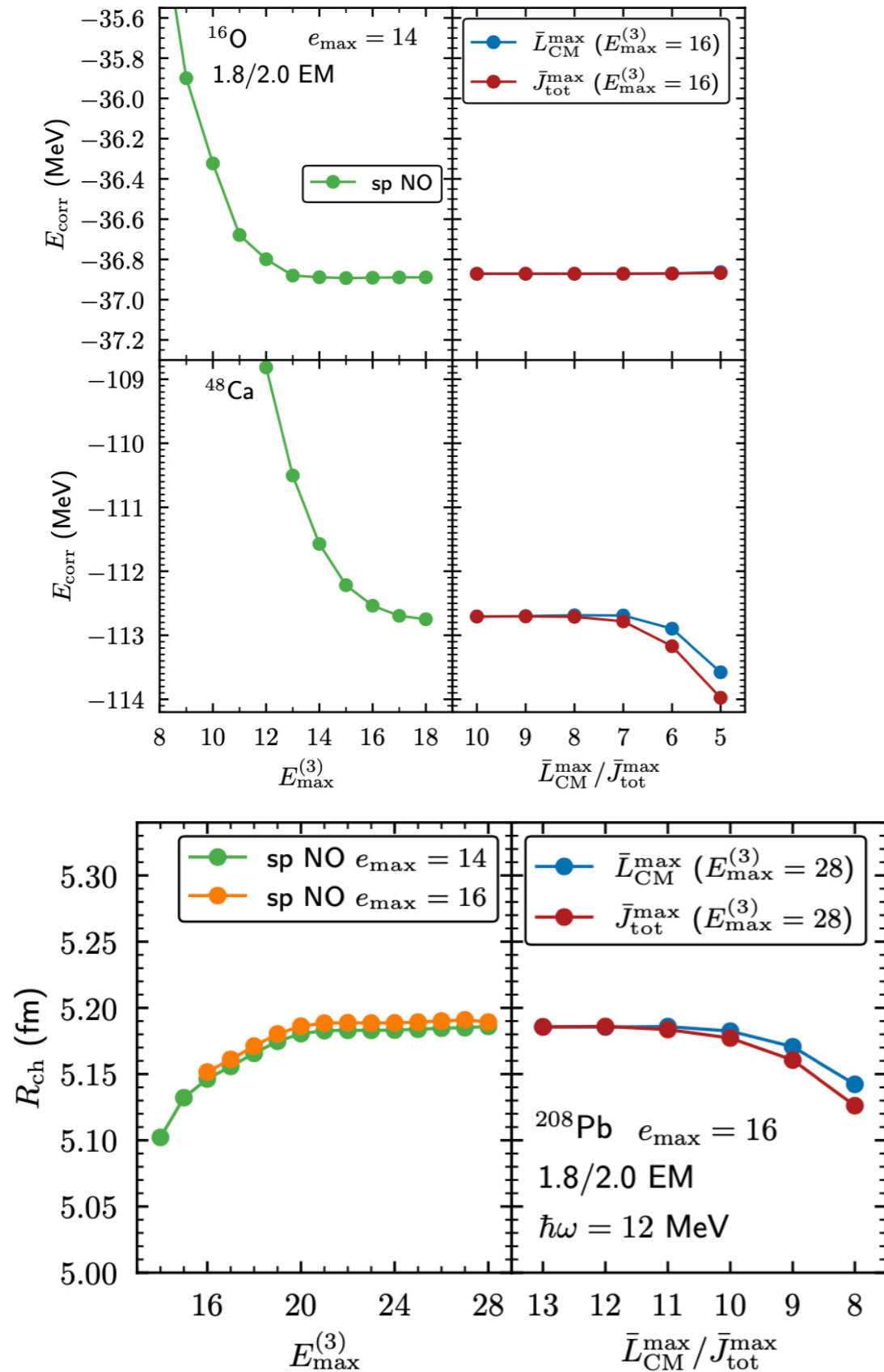


- at no stage single-particle 3N HO matrix elements needed
- $N_{\max}$  can be increased straightforwardly
- basis size and storage space determined by  $J_{\max}$  and  $L_{\text{cm},\max}$ ,



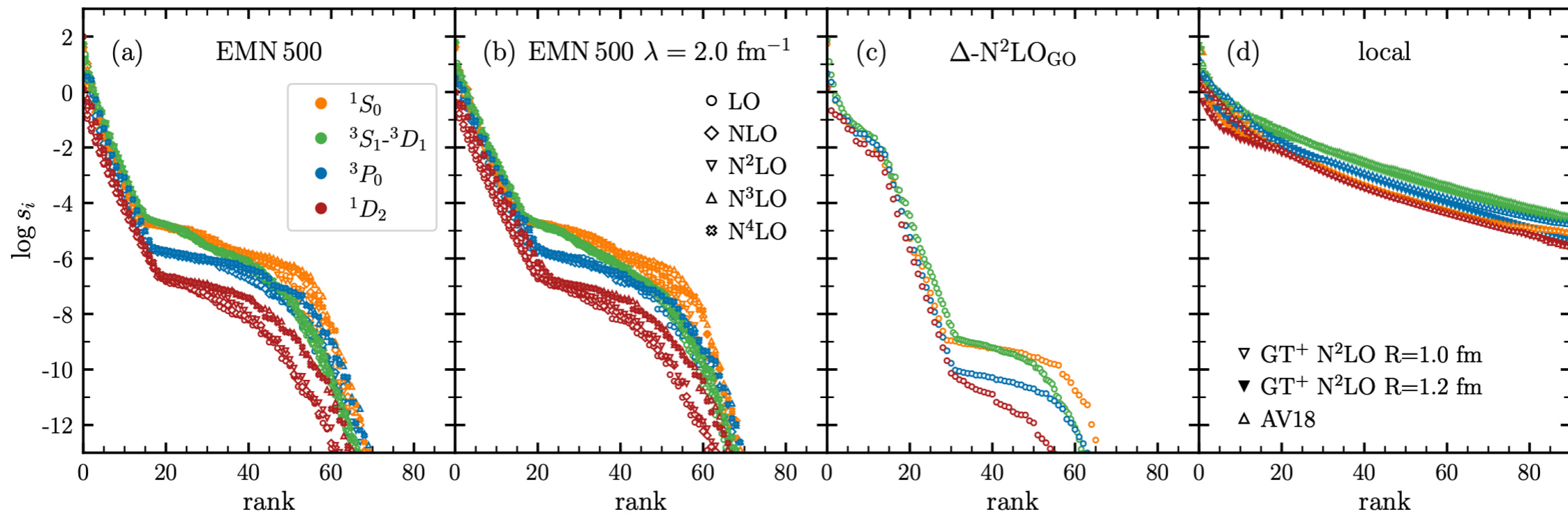
# Novel normal ordering framework for 3N interactions

KH et al., PRC 107, 024310 (2023)

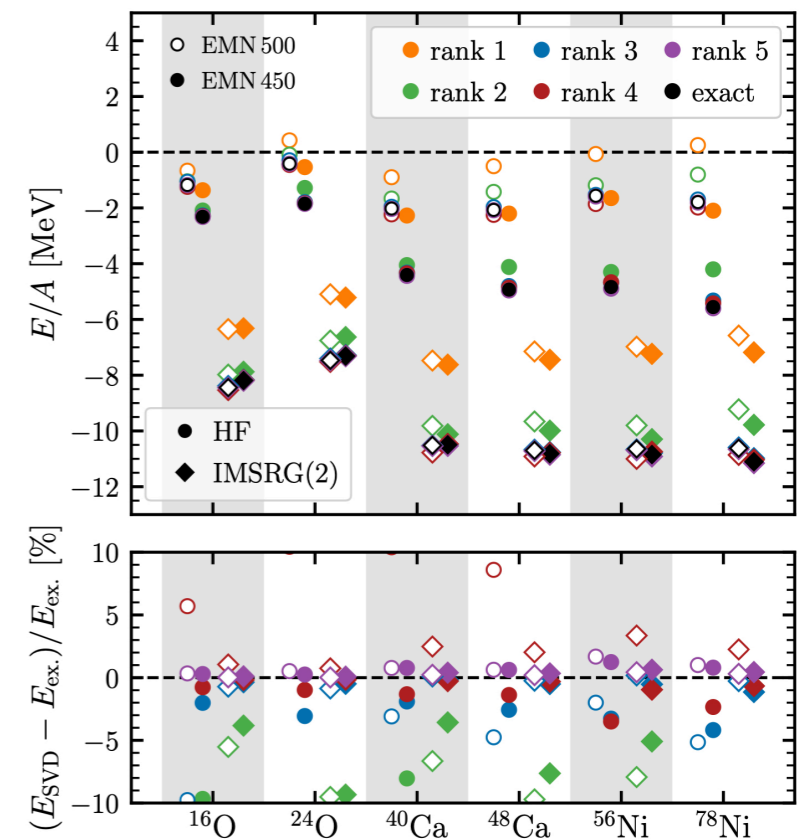
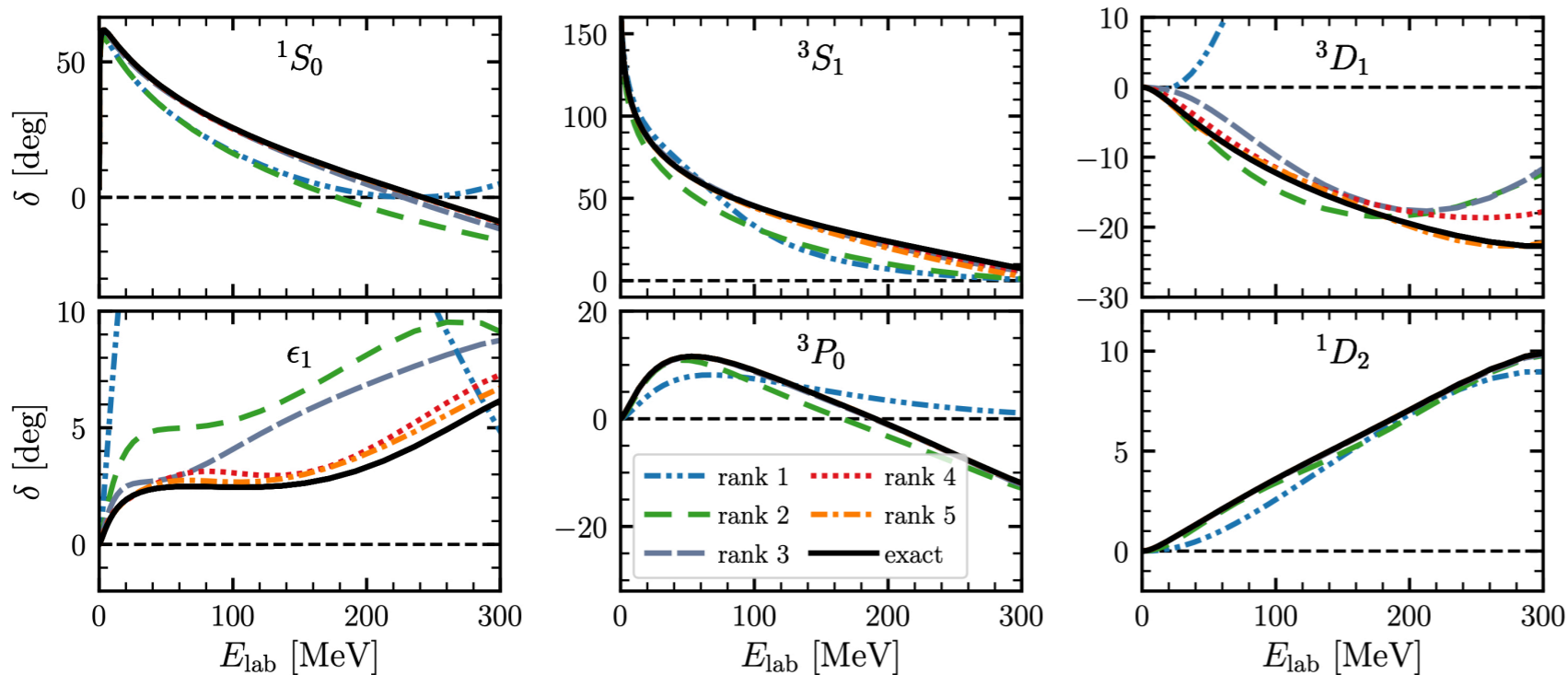




# Singular value decomposition of NN interactions



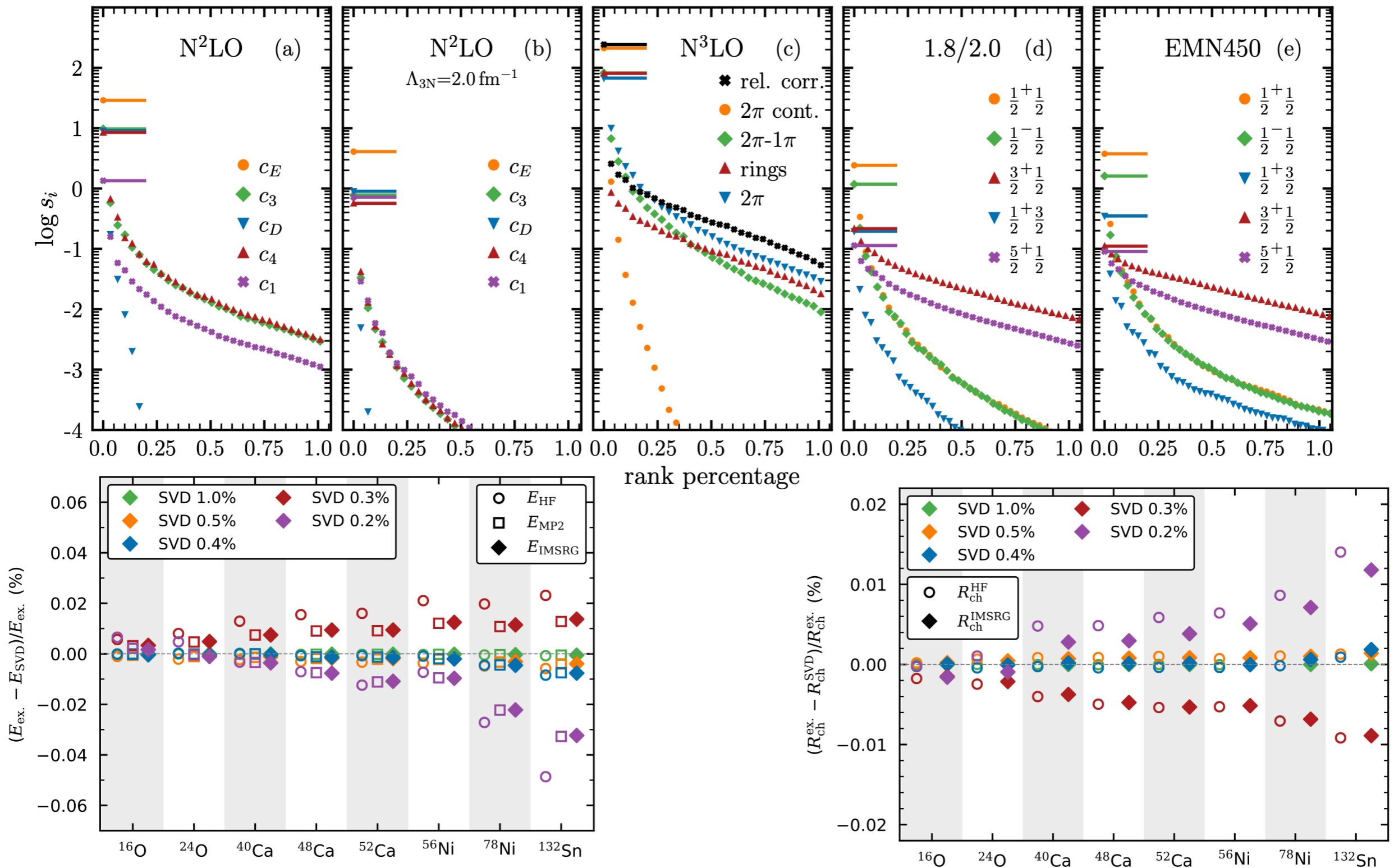
Tichai et al., PLB 821, 136623 (2021)



- excellent agreement with full results for phase shifts and binding energies for very low number of ranks (= number of retained singular values)

# (Randomized) singular value decomposition of 3N interactions

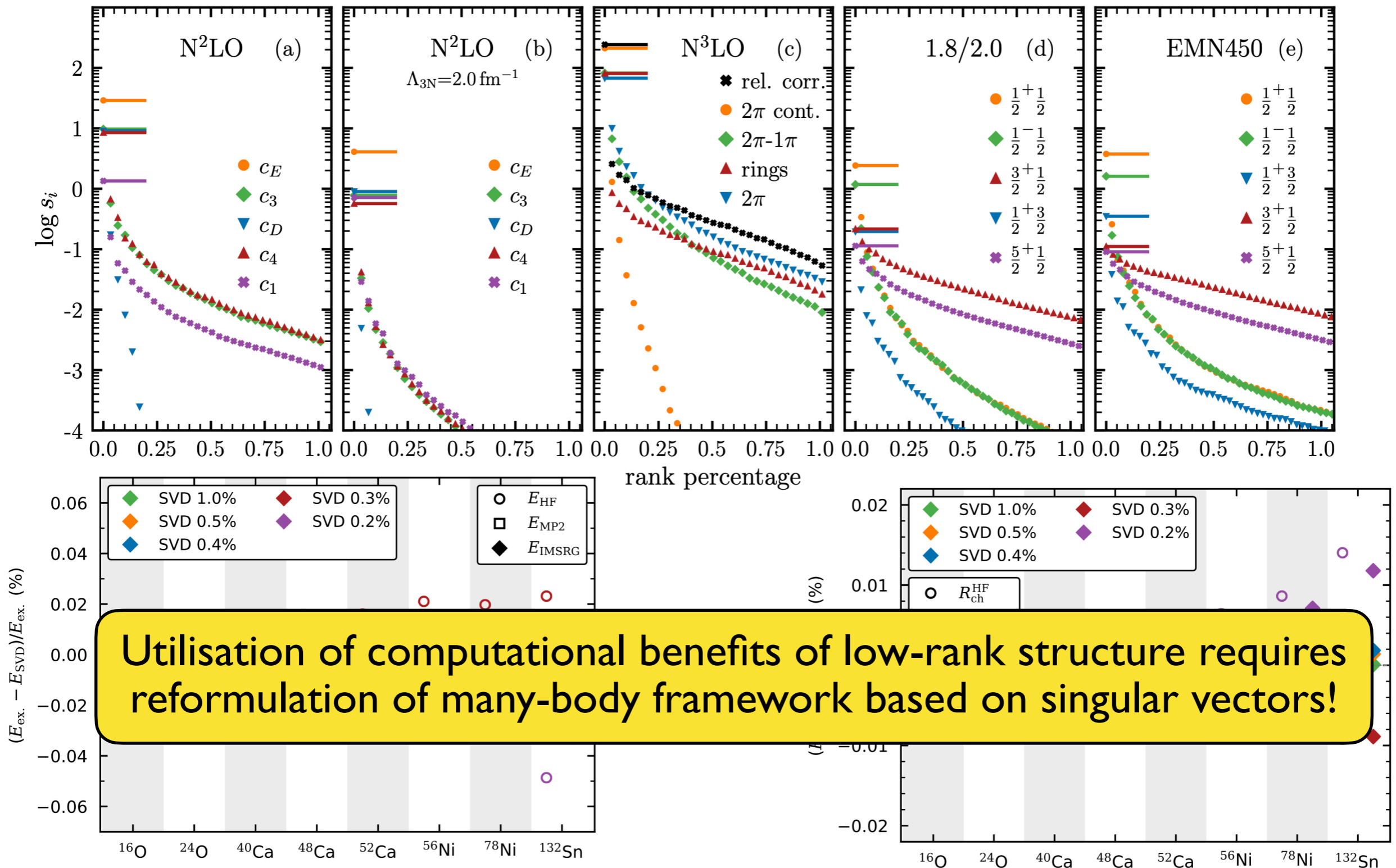
Tichai et al., arXiv:2307.15572



- again good agreement with full results for binding energies and charge radii for very low number of ranks (NN interactions not SVD-decomposed)

# (Randomized) singular value decomposition of 3N interactions

Tichai et al., arXiv:2307.15572



- again good agreement with full results for binding energies and charge radii for very low number of ranks (NN interactions not SVD-decomposed)

Which contributions should a comprehensive uncertainty estimate contain?

- I. Power counting scheme
- II. Chiral expansion
- III. Regularization schemes
- IV. Different fitting strategies for low-energy couplings
- V. Truncation in many-body expansions/SRG evolution
- VI. Basis truncations
- VII....?