

Dispersion relations and radiative corrections for the two-pion channel

Peter Stoffer

Physik-Institut, University of Zurich, and Paul Scherrer Institut

in collaboration with G. Colangelo and M. Hoferichter,
JHEP **02** (2019) 006; arXiv:2308.04217 [hep-ph]

with G. Colangelo, M. Hoferichter, and B. Kubis,
JHEP **10** (2022) 032

with **J. Lüdtk**e and M. Procura,
JHEP **04** (2023) 125

and with **N. Gerasis**, **E. Kaziukėnas**, **J.-N. Toelstede**, and **T. Leplumey**
work in progress

MUonE Workshop, MITP Mainz, Germany, June 5, 2024



University of
Zurich^{UZH}

PAUL SCHERRER INSTITUT



Swiss National
Science Foundation

- 1 Introduction
- 2 Dispersive analysis of pion VFF
- 3 Zeros in the form factor
- 4 Structure-dependent radiative corrections
- 5 Summary

- 1 Introduction
- 2 Dispersive analysis of pion VFF
- 3 Zeros in the form factor
- 4 Structure-dependent radiative corrections
- 5 Summary

Two-pion contribution to HVP

- $\pi\pi$ contribution amounts to more than 70% of HVP contribution
- dominant source of HVP uncertainty
- can be expressed in terms of **pion vector form factor** \Rightarrow constraints from **analyticity and unitarity**

A multitude of puzzles in HVP

- tension between BMWc **lattice-QCD** and dispersive evaluations based on older e^+e^- **cross sections**
- discrepancy between **CMD-3** and all previous e^+e^- experiments
- ongoing scrutiny of both lattice and dispersive evaluations
- role of **radiative corrections**?

- 1 Introduction
- 2 Dispersive analysis of pion VFF**
- 3 Zeros in the form factor
- 4 Structure-dependent radiative corrections
- 5 Summary

Unitarity and analyticity

implications of unitarity (two-pion intermediate states):

- ① $\pi\pi$ contribution to HVP—pion vector form factor (VFF)
- ② pion VFF— $\pi\pi$ scattering
- ③ $\pi\pi$ scattering— $\pi\pi$ scattering

$\sigma(e^+e^- \rightarrow \pi^+\pi^-) \propto |F_\pi^V(s)|^2$

analyticity \Rightarrow dispersion relation for HVP contribution

Unitarity and analyticity

implications of unitarity (two-pion intermediate states):

- ① $\pi\pi$ contribution to HVP—pion vector form factor (VFF)
- ② pion VFF— $\pi\pi$ scattering
- ③ $\pi\pi$ scattering— $\pi\pi$ scattering

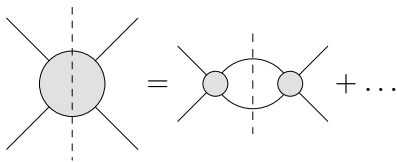
$$F_{\pi}^V(s) = |F_{\pi}^V(s)| e^{i\delta_1^0(s) + \dots}$$

analyticity \Rightarrow dispersion relation for pion VFF

Unitarity and analyticity

implications of unitarity (two-pion intermediate states):

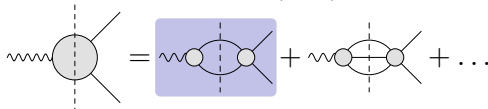
- 1 $\pi\pi$ contribution to HVP—pion vector form factor (VFF)
- 2 pion VFF— $\pi\pi$ scattering
- 3 $\pi\pi$ scattering— $\pi\pi$ scattering



analyticity, crossing, PW expansion \Rightarrow Roy equations

Dispersive representation of pion VFF

→ Colangelo, Hoferichter, Stoffer, JHEP **02** (2019) 006



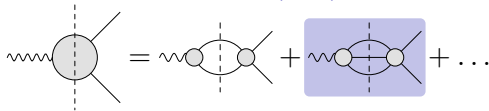
$$F_{\pi}^V(s) = \Omega_1^1(s) \times G_{\omega}(s) \times G_{\text{in}}^N(s)$$

- Omnès function with elastic $\pi\pi$ -scattering P -wave phase shift $\delta_1^1(s)$ as input:

$$\Omega_1^1(s) = \exp \left\{ \frac{s}{\pi} \int_{4M_{\pi}^2}^{\infty} ds' \frac{\delta_1^1(s')}{s'(s' - s)} \right\}$$

Dispersive representation of pion VFF

→ Colangelo, Hoferichter, Stoffer, JHEP **02** (2019) 006



$$F_{\pi}^V(s) = \Omega_1^1(s) \times G_{\omega}(s) \times G_{\text{in}}^N(s)$$

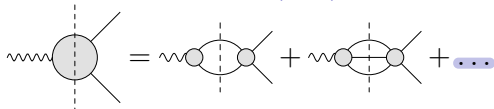
- isospin-breaking 3π intermediate state: negligible apart from ω resonance (ρ - ω interference effect)

$$G_{\omega}(s) = 1 + \frac{s}{\pi} \int_{9M_{\pi}^2}^{\infty} ds' \frac{\text{Im}g_{\omega}(s')}{s'(s'-s)} \left(\frac{1 - \frac{9M_{\pi}^2}{s'}}{1 - \frac{9M_{\pi}^2}{M_{\omega}^2}} \right)^4,$$

$$g_{\omega}(s) = 1 + \epsilon_{\omega} \frac{s}{(M_{\omega} - \frac{i}{2}\Gamma_{\omega})^2 - s}$$

Dispersive representation of pion VFF

→ Colangelo, Hoferichter, Stoffer, JHEP **02** (2019) 006



$$F_{\pi}^V(s) = \Omega_1^1(s) \times G_{\omega}(s) \times G_{\text{in}}^N(s)$$

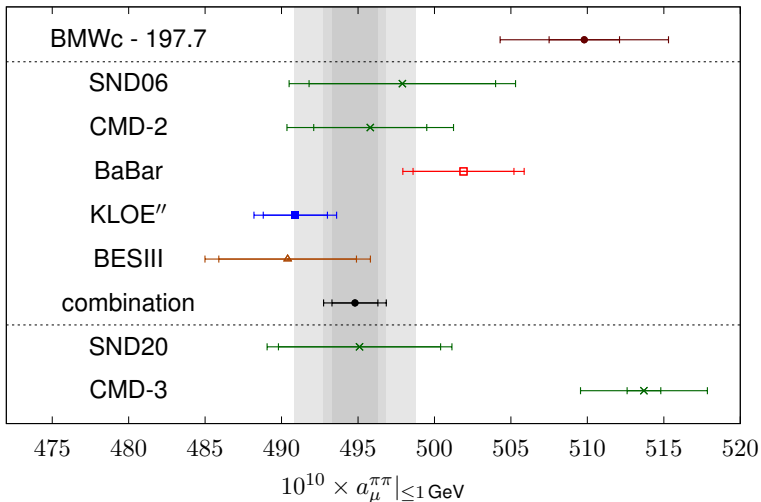
- heavier intermediate states: 4π (mainly $\pi^0\omega$), $\bar{K}K$, ...
- described in terms of a conformal polynomial with cut starting at $\pi^0\omega$ threshold

$$G_{\text{in}}^N(s) = 1 + \sum_{k=1}^N c_k (z^k(s) - z^k(0))$$

- correct P -wave threshold behavior imposed
- potentially leads to **zeros of the form factor**

Result for $a_{\mu}^{\text{HVP},\pi\pi}$ below 1 GeV

→ Colangelo, Hoferichter, Stoffer, JHEP **02** (2019) 006 and 2308.04217 [hep-ph]
 Colangelo, Hoferichter, Kubis, Stoffer, JHEP **10** (2022) 032



CMD-3 vs. all the rest

→ Colangelo, Hoferichter, Stoffer, 2308.04217 [hep-ph]

discrepancy	$a_{\mu}^{\pi\pi} _{[0.60, 0.88] \text{ GeV}}$	$a_{\mu}^{\pi\pi} _{\leq 1 \text{ GeV}}$	int window
SND06	1.8σ	1.7σ	1.7σ
CMD-2	2.3σ	2.0σ	2.1σ
BaBar	3.3σ	2.9σ	3.1σ
KLOE''	5.6σ	4.8σ	5.4σ
BESIII	3.0σ	2.8σ	3.1σ
SND20	2.2σ	2.1σ	2.2σ
Combination	4.2σ (6.1σ)	3.7σ (5.0σ)	3.8σ (5.7σ)

(discrepancies in brackets exclude systematic effect due to BaBar–KLOE tension)

- p -value of fit to CMD-3: 20%
- $\pi\pi$ phase shifts reasonable, main effect in conformal polynomial
- effect on charge radius as expected for rather uniform cross-section shift

- 1 Introduction
- 2 Dispersive analysis of pion VFF
- 3 Zeros in the form factor**
- 4 Structure-dependent radiative corrections
- 5 Summary

Zeros in the pion VFF?

- presence of zeros can in principle be tested with modulus sum rule \rightarrow H. Leutwyler, arXiv:hep-ph/0212324

$$\psi(s) := \frac{1}{(s_0 - s)^{3/2}} \log \frac{F_\pi^V(s)}{F_\pi^V(s_0)}, \quad s_0 = 4M_\pi^2$$

\Rightarrow check if $\psi(s)$ fulfills unsubtracted dispersion relation

$$\psi(s) \stackrel{?}{=} \frac{1}{\pi} \int_{s_0}^{\infty} ds' \frac{\text{Im}\psi(s')}{s' - s}, \quad \text{Im}\psi(s) = -\frac{1}{(s - s_0)^{3/2}} \log \left| \frac{F_\pi^V(s)}{F_\pi^V(s_0)} \right|$$

- only need **modulus of form factor** \Rightarrow experiment

Zeros in the pion VFF?

- no zeros possible in the region of validity of χ PT
→ H. Leutwyler, arXiv:hep-ph/0212324
- zeros in low-energy region excluded via unitarity/analyticity
→ B. Ananthanarayan, I. Caprini, I. Sentitemsu Imsong, PRD **83** (2011) 096002
- zeros excluded at large values of $|s|$ from asymptotic behavior → G. P. Lepage, S. J. Brodsky, PLB **87** (1979) 359
- use VFF parametrization to **test presence of zeros**:
 - fits lead to $G_{\text{in}}^N(s)$ free of zeros for $N \leq 4$
 - for $N > 4$, zeros show up, accompanied by **fit instabilities**
 - zeros for $N > 4$ source of **main systematic uncertainty** in our representation

Constrained fits without zeros

→ work in progress with **Thomas Leplumey** (ETH master student)

- impose absence of zeros, either via explicit parametrization, or **sum-rule constraint**

$$\log G_{\text{in}}^N(s_{\text{in}}) = \frac{1}{\pi} \int_{s_{\text{in}}}^{\infty} \frac{ds'}{s'} \frac{s_{\text{in}}^{3/2}}{(s' - s_{\text{in}})^{3/2}} \log \left| \frac{G_{\text{in}}^N(s')}{G_{\text{in}}^N(s_{\text{in}})} \right|$$

- observe stabilization of fits for larger $N \Rightarrow$ **main source of uncertainty eliminated**

Constrained fits without zeros

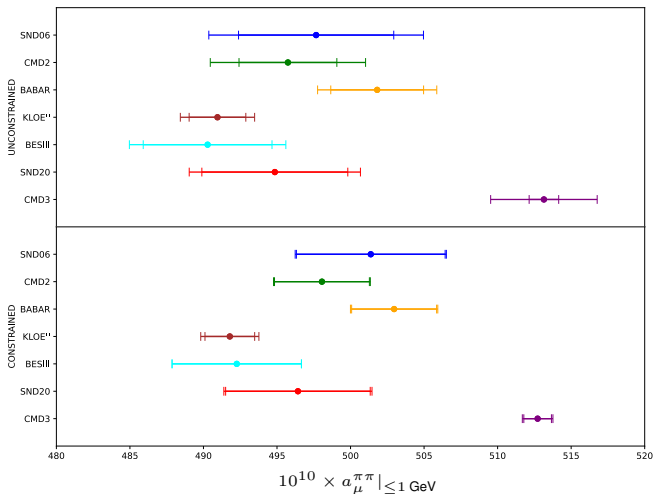
→ work in progress with **Thomas Leplumey** (ETH master student)

- **marginal impact** on χ^2/dof of fit, ω mass and mixing parameter, **central values** of $\pi\pi$ phase
- systematic **uncertainties much reduced** for $\pi\pi$ -phase δ_1^1 , $a_{\mu}^{\pi\pi}$, and pion charge radius $\langle r_{\pi}^2 \rangle$
- fits now lead to results for $\langle r_{\pi}^2 \rangle$ that could be used to discriminate between experiments \Rightarrow opportunity for **independent lattice-QCD checks**

→ Colangelo, Hoferichter, Stoffer, PLB **814** (2021) 136073

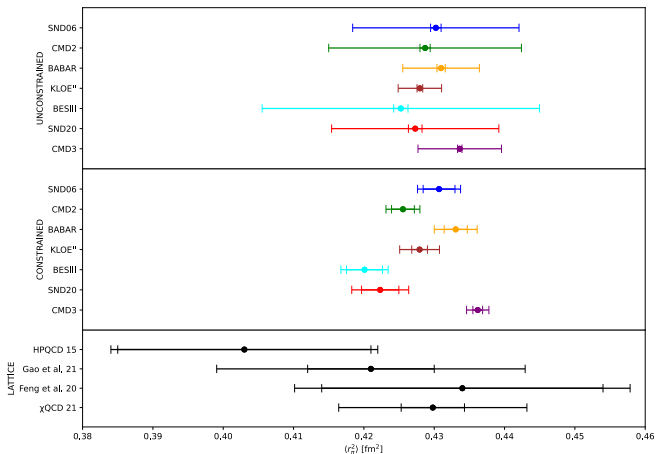
Constrained fits without zeros: $a_{\mu}^{\pi\pi}$

→ work in progress with **Thomas Leplumey** (ETH master student)



Constrained fits without zeros: pion charge radius $\langle r_\pi^2 \rangle$

→ work in progress with **Thomas Leplumey** (ETH master student)



- 1 Introduction
- 2 Dispersive analysis of pion VFF
- 3 Zeros in the form factor
- 4 Structure-dependent radiative corrections**
- 5 Summary

How can we resolve the discrepancies?

- CMD-3 vs. CMD-2 (and SND): experimental issue?
- SND20: **incompatible** with unitarity/analyticity constraints
(p -value: 3.8×10^{-3})
- CMD-3 vs. radiative-return experiments:
model-dependent theory input—how reliable are uncertainties?

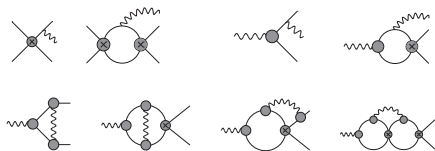
How can we resolve the discrepancies?

- certainly **no conceptual problem** with dispersive approach per se
- dispersive approach relies on **data input**
- but experiments require **theory input**
 - ⇒ try to **reduce model dependence** in that theory input
 - ⇒ need **more dispersion theory**, not less!

Dispersive approach to isospin corrections in

$\pi\pi$ scattering and $F_\pi^V \rightarrow$ talk by G. Colangelo at Zurich WorkStop 2023

\rightarrow G. Colangelo, M. Cottini, J. Monnard, J. Ruiz de Elvira, work in progress

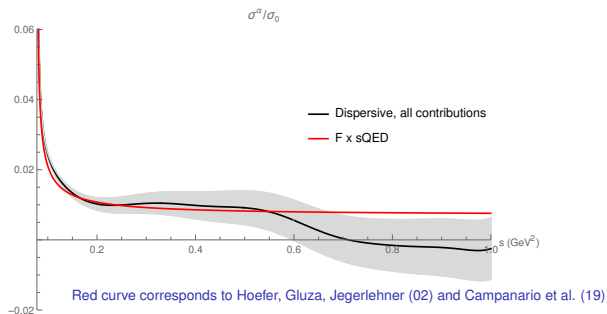


- pion-mass difference in Roy equations
- photonic corrections (real + virtual) to $\pi\pi$ scattering and pion vector form factor

Dispersive approach to isospin corrections in

$\pi\pi$ scattering and $F_\pi^V \rightarrow$ talk by G. Colangelo at Zurich WorkStop 2023

\rightarrow G. Colangelo, M. Cottini, J. Monnard, J. Ruiz de Elvira, work in progress



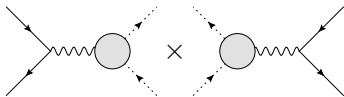
\rightarrow J. Monnard, PhD thesis (2021)

\Rightarrow no dramatic effects found

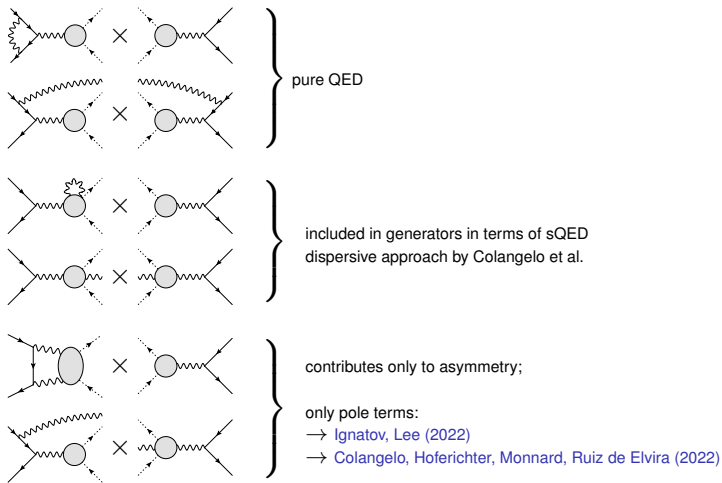
Comparison of MCs → talks by A. Signer, Y. Ulrich

- radiative-return experiments: PHOKHARA
 - FSR from pointlike pions
 - boxes, pentagons with vector form factor *outside* loop integral
- direct scan experiments: MCGPJ
 - FSR from pointlike pions
 - box diagrams for asymmetry with vector form factor inside loop

Direct scan experiments: LO



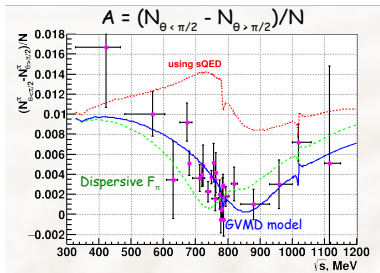
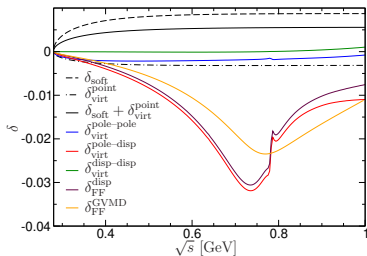
Direct scan experiments: NLO



Forward-backward asymmetry



→ talks by G. Colangelo and F. Ignatov at Zurich WorkStop 2023



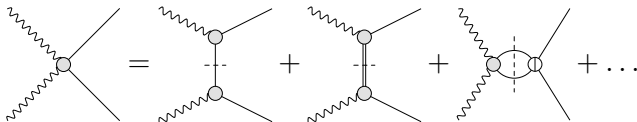
→ Colangelo, Hoferichter, Monnard, Ruiz de Elvira (2022)

→ Ignatov, Lee (2022)

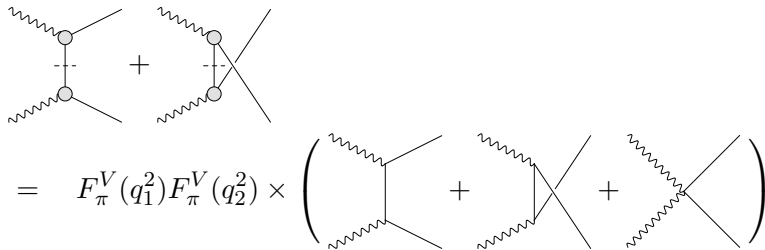
4 Structure-dependent radiative corrections

Dispersion relations for $\gamma^* \gamma^* \rightarrow \pi^+ \pi^-$

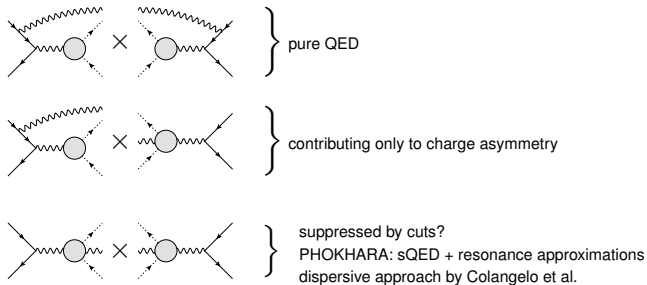
→ Colangelo, Hoferichter, Procura, Stoffer (2015), Hoferichter, Stoffer (2019)



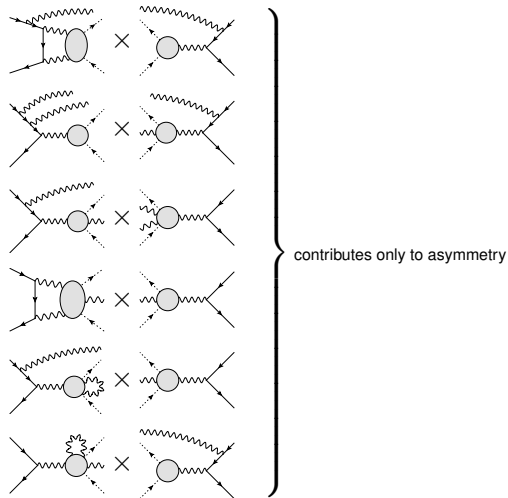
pole term = FsQED



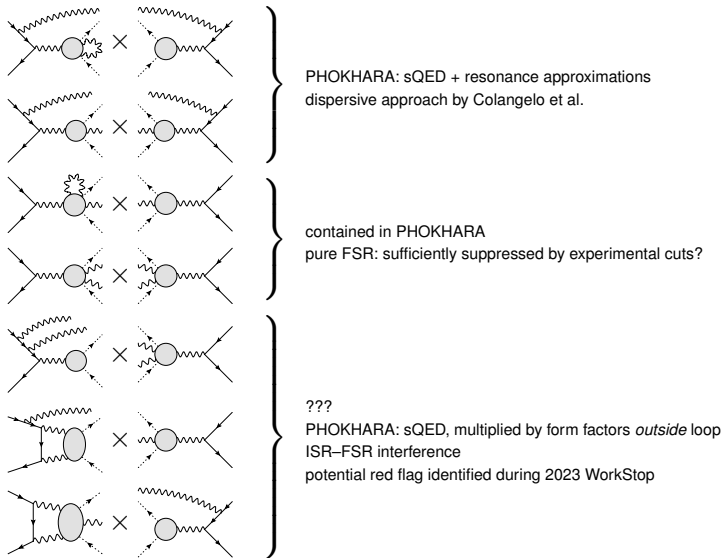
Radiative-return experiments: LO



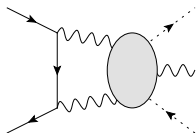
Radiative-return experiments: NLO (omitting pure QED corrections to LO)



Radiative-return experiments: NLO (omitting pure QED corrections to LO)



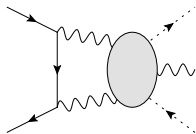
Most difficult sub-process: $\gamma^* \gamma^* \gamma \rightarrow \pi^+ \pi^-$



- PHOKHARA: $s\text{QED} \times F_{\pi}^V(s)$ (s : e^+e^- invariant squared energy)—model prescription, which achieves cancellation of IR singularities
- **not FsQED** (= dispersive pole terms): lesson learnt from asymmetry might raise concerns
- here: **dispersive pole terms expected to be bad approximation**: $\pi\pi$ system in p -wave, ρ resonance in rescattering

Most difficult sub-process: $\gamma^* \gamma^* \gamma \rightarrow \pi^+ \pi^-$

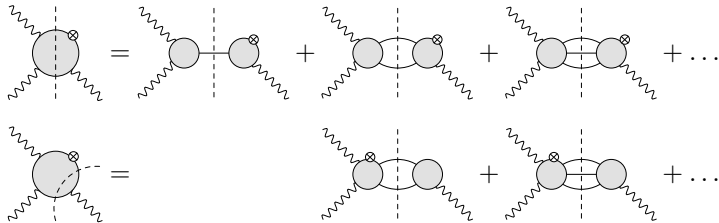
→ work in progress with **Emilis Kaziukėnas**, **Nikolas Geralis** (ETH master students), **J.-N. Toelstede**



- goal: dispersive treatment of $\gamma^* \gamma^* \gamma \rightarrow \pi \pi$
- synergies with **dispersive approach to HLbL** in triangle kinematics → talk by **M. Hoferichter**
- warm-up: $\gamma^* \gamma^* \gamma \rightarrow \pi \pi$ at NLO in $SU(2)$ χ PT computed
→ work in progress with **Emilis Kaziukėnas**
⇒ useful to understand **analytic structure** and for fixing **subtraction constants**

HLbL in triangle kinematics

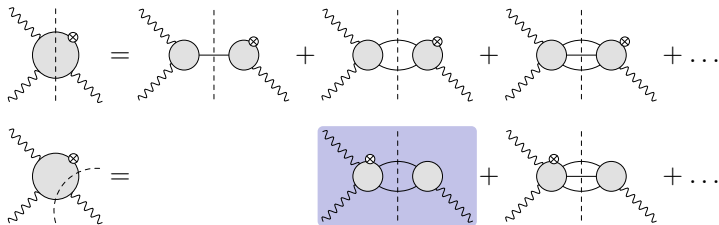
→ Lütke, Procura, Stoffer, JHEP **04** (2023) 125



- same sub-process
- for HLbL: only soft-photon limit required
- beyond soft limit: **ambiguities in tensor decomposition**
need to be addressed for $e^+e^- \rightarrow \pi^+\pi^-\gamma$
- dispersive definition of **pole terms** non-trivial
→ work in progress with **Emilis Kaziukėnas**

HLbL in triangle kinematics

→ Lüdtke, Procura, Stoffer, JHEP **04** (2023) 125



- same sub-process
- for HLbL: only soft-photon limit required
- beyond soft limit: **ambiguities in tensor decomposition**
need to be addressed for $e^+e^- \rightarrow \pi^+\pi^-\gamma$
- dispersive definition of **pole terms** non-trivial
→ work in progress with **Emilis Kaziukėnas**

- 1 Introduction
- 2 Dispersive analysis of pion VFF
- 3 Zeros in the form factor
- 4 Structure-dependent radiative corrections
- 5 Summary**

Summary

- systematic uncertainties in pion VFF drastically reduced if **zeros are excluded**
- **data do not prefer zeros**: presence of zeros in fit connected with instabilities
- reduce model dependence in radiative corrections: **rely on dispersion theory**
- $\gamma^*\gamma^* \rightarrow \pi\pi$ sub-process: well understood **pion-pole terms**, rescattering could be included
- $\gamma^*\gamma^*\gamma \rightarrow \pi\pi$ sub-process: **very difficult**, but synergies with new dispersive approach to HLbL; pole terms not enough

Backup

Tension with lattice QCD

→ Colangelo, Hoferichter, Stoffer, PLB **814** (2021) 136073

- force a different HVP contribution in VFF fits by including “lattice” datum with tiny uncertainty
- three different scenarios:
 - “low-energy” physics: $\pi\pi$ phase shifts
 - “high-energy” physics: inelastic effects, C_k
 - all parameters free
- study effects on pion charge radius, hadronic running of $\alpha_{\text{QED}}^{\text{eff}}$, phase shifts, cross sections

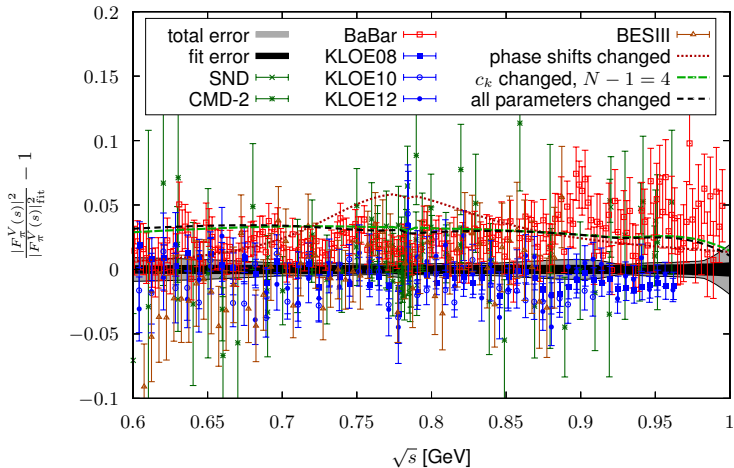
Modifying $a_{\mu}^{\pi\pi} |_{\leq 1 \text{ GeV}}$

→ Colangelo, Hoferichter, Stoffer, PLB **814** (2021) 136073

- “low-energy” scenario requires large local changes in the cross section in the ρ region
- “high-energy” scenario has an impact on **pion charge radius** and the space-like VFF \Rightarrow chance for independent lattice-QCD checks

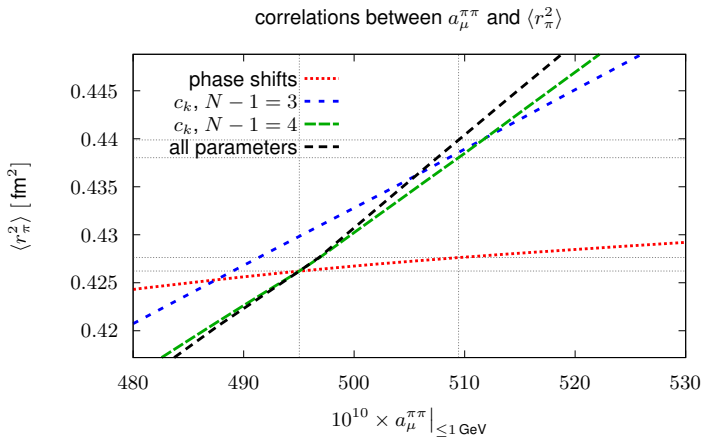
Modifying $a_{\mu}^{\pi\pi} |_{\leq 1 \text{ GeV}}$

→ Colangelo, Hoferichter, Stoffer, PLB **814** (2021) 136073

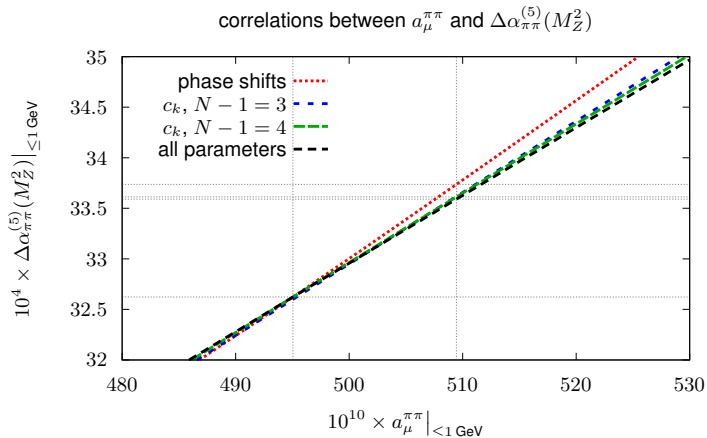


Modifying $a_\mu^{\pi\pi} |_{\leq 1 \text{ GeV}}$

→ Colangelo, Hoferichter, Stoffer, PLB **814** (2021) 136073



Modifying $a_\mu^{\pi\pi} |_{\leq 1 \text{ GeV}}$



Modifying $a_{\mu}^{\pi\pi} |_{\leq 1 \text{ GeV}}$ 