
The Evaluation of the Leading Hadronic Contribution to a_μ
Consolidation of MUonE Experiment and Recent Developments in Low Energy e^+e^- Data

MCMULE for MUonE

Marco Rocco for MCMULE

Paul Scherrer Institut

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- ✿ higher-order predictions and comparison with precision experiments
- ✿ focus on $2 \rightarrow 2$ low-energy QED+ scattering processes
- ✿ **input:** matrix elements by us or others (at NNLO + first visits at N3LO)
- ✿ **output:** physical cross section for any physical observable at fixed order
- ✿ at present an integrator, generator features under testing

MCMULE

Monte Carlo for MUons and other LEptons

code → <https://mule-tools.gitlab.io/>

docs → <https://mcmule.readthedocs.io/>





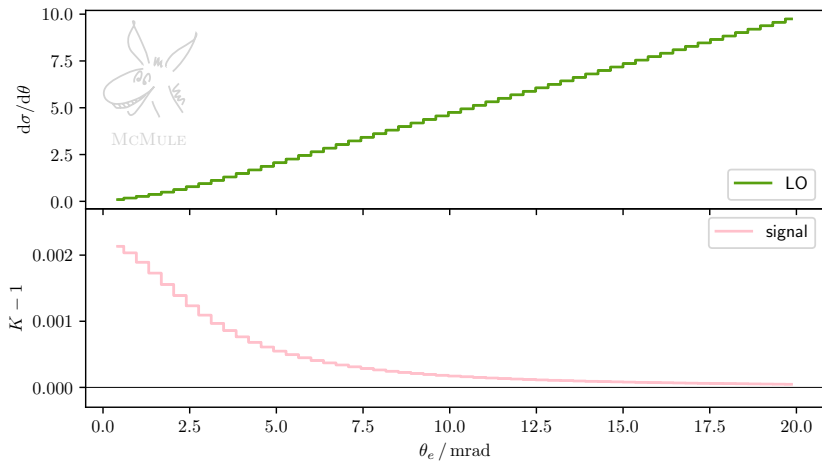
phenomenology

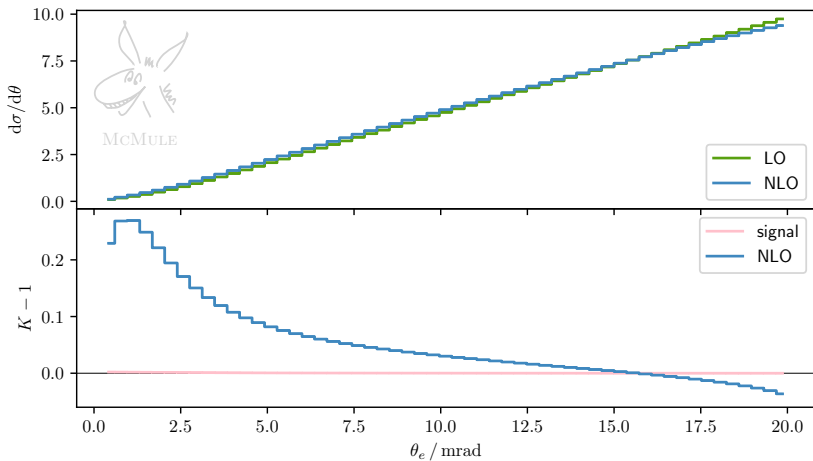
$$\mu e \rightarrow \mu e \quad @ \text{ NNLO}$$

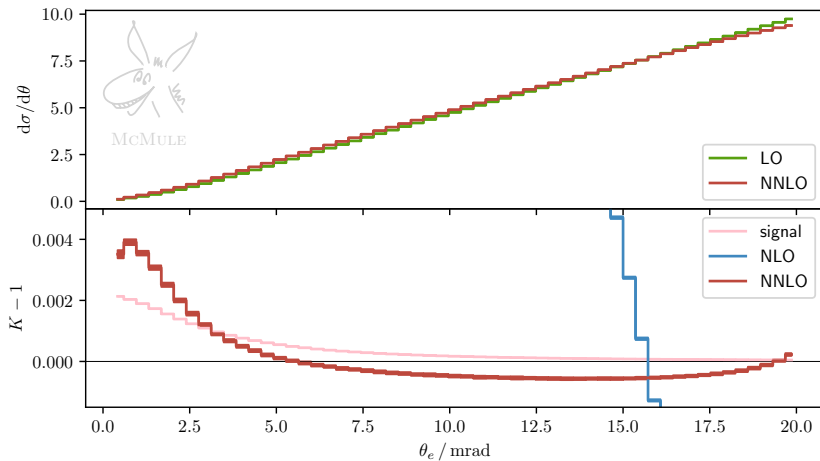
- kinematical setup mimics MUonE:

$$S1 \quad :: \quad E_{\mu,i} = 160 \text{ GeV} \quad E_{e,f} > 1 \text{ GeV} \quad \theta_{\mu,f} > 0.3 \text{ mrad}$$

- results for different kinematical scenarios and any IR safe observable
- no mass is neglected



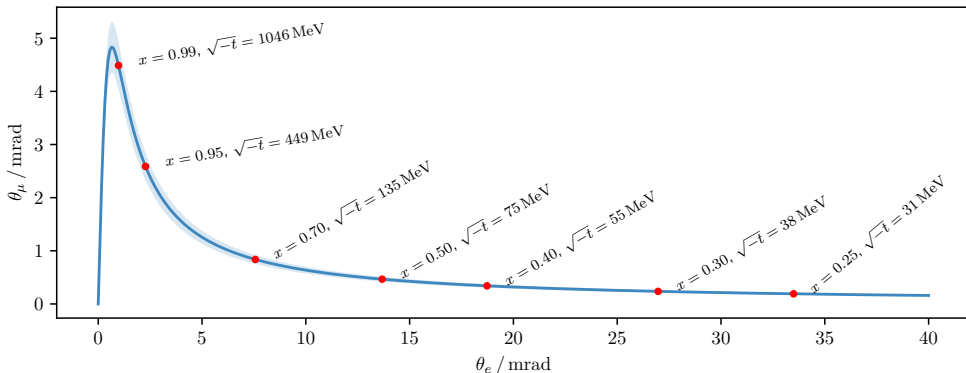




how to handle hard radiation?

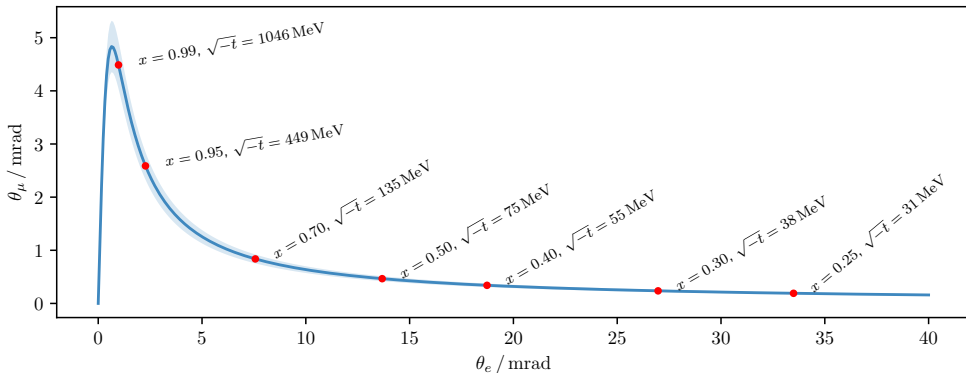
- elasticity veto! $\rightarrow 0.9 < \frac{\theta_{\mu,f}}{\theta_{\mu,f}^{el}} < 1.1$ (S2)

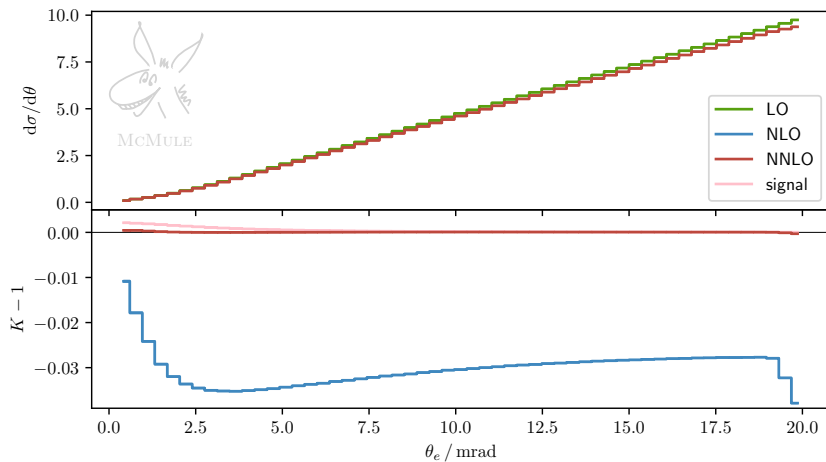
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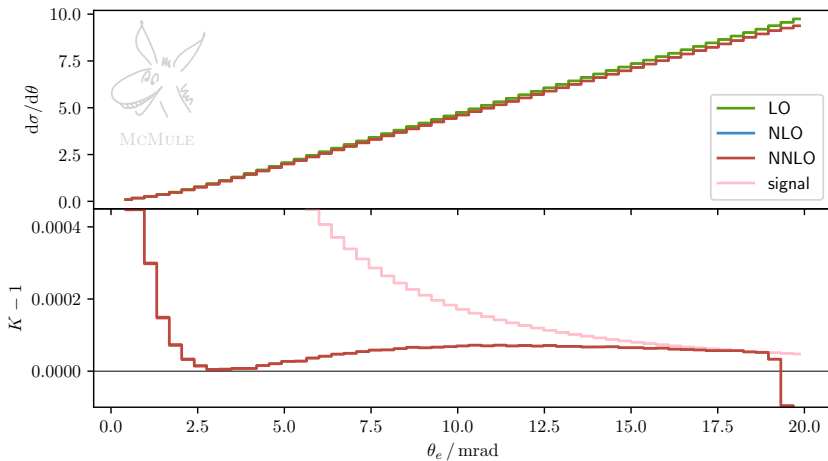


how to handle hard radiation?

- elasticity veto! $\rightarrow 0.9 < \frac{\theta_{\mu,f}}{\theta_{\mu,f}^{el}} < 1.1$ (S2)
- something else?









a peek in the stable theory

$$\begin{aligned}
 & \int d\Phi_2 \left| \begin{array}{c} \text{tree} \\ \text{tree} \\ \text{tree} \\ \dots \end{array} \right|^2 \\
 & + \int d\Phi_3 \left| \begin{array}{c} \text{one-loop} \\ \text{one-loop} \\ \dots \end{array} \right|^2 \\
 & + \int d\Phi_4 \left| \begin{array}{c} \text{two-loop} \\ \dots \end{array} \right|^2
 \end{aligned}$$

- ① fully-differential PS integration
→ FKS^ℓ
- ② virtual amplitudes with massive particles
→ one-loop: OpenLoops
→ two-loop: massification
- ③ numerical instabilities due to pseudo-singularities
→ next-to-soft stabilisation



local subtraction of infrared divergences

$$\begin{aligned}
 & \int d\Phi_n \left\{ \text{red blob} + \int d\Phi_\gamma \text{red blob with } \zeta \right\} \\
 &= \int d\Phi_n d\Phi_\gamma \left\{ \text{red blob with } \zeta - \text{green blob} \right\} + \int d\Phi_n \left\{ \text{red blob} + \int d\Phi_\gamma \text{green blob} \right\}
 \end{aligned}$$

- ◇ exploits exponentiation of **soft singularities** [VFS 61]
- ◇ works at **all orders** in QED [Engel, Signer, Ulrich 19]




- ◇ singularities are dealt with **locally** → **stable** numerical integration
- ◇ subtraction makes negative-weighted events much more frequent
- ◇ theory error: 0

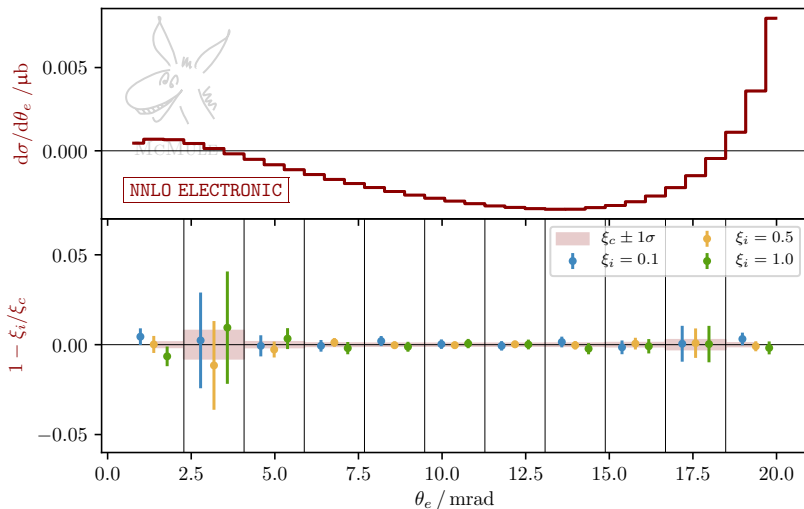


- *photonic* and *fermionic* (\rightarrow hyperspherically [Fael 18]) corrections
- photonic at NNLO are split as

$$d\sigma^{(2)} = \int d\Phi_n \mathcal{M}_n^{(2)} + \int d\Phi_{n+1} \mathcal{M}_{n+1}^{(1)} + \int d\Phi_{n+2} \mathcal{M}_{n+2}^{(0)}$$

- for each part identify *gauge-invariant subsets* based on lepton charges (q for electron, Q for muon)

- ◇ $q^6 Q^2$:: *electronic*

- ◇ $\{q^5 Q^3, q^4 Q^4, q^3 Q^5\}$:: *mixed*

- ◇ $q^2 Q^6$:: *muonic*




full 2-loop amplitude with $M \neq 0$, $m = 0 \rightarrow$ [Bonciani et al. 21]

full 2-loop amplitude with $M \neq 0$, $m \neq 0 \rightarrow$ [??]

- ◇ exploit scale hierarchy $m^2 \ll M^2, Q^2$, expand in $m^2/Q^2 \sim 0$

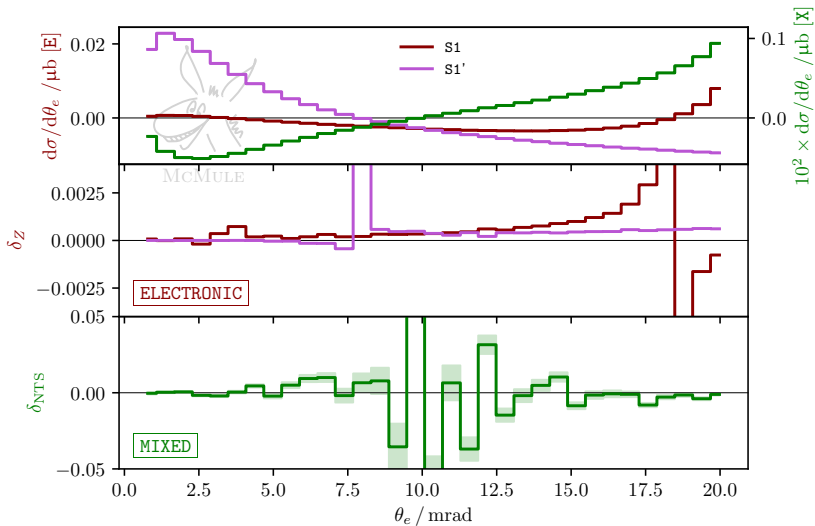
$$\text{Diagram} \sim A \log^2 \frac{m^2}{Q^2} + B \log \frac{m^2}{Q^2} + C + \mathcal{O}\left(\frac{m^2}{Q^2}\right)$$

- ◇ massification: $\mathcal{A}_{mM}(m) = \mathcal{S} \times \mathcal{Z} \times \mathcal{Z} \times \mathcal{A}_{mM}(0) + \mathcal{O}(m)$

[Penin 06, Becher, Melnikov 07; Engel, Gnendiger, Signer, Ulrich 18]

- ◇ theory error: $\mathcal{O}(10^{-3})$ @ NNLO $\sim \mathcal{O}(10^{-6})$

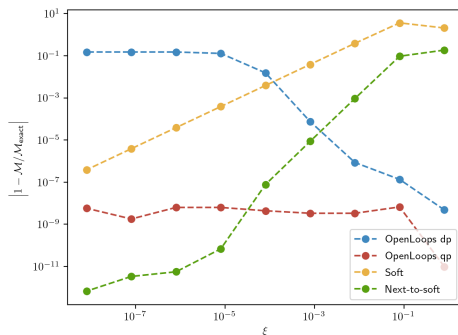




OpenLoops [Buccioni, Pozzorini, Zoller 18, Buccioni et al. 19]

LBK theorem [LBK 58-61, Engel, Signer, Ulrich 21, 2xEngel 23]

$$\begin{array}{c} \text{wavy line} \\ \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \end{array} \xrightarrow{E_\gamma \rightarrow 0} \varepsilon \begin{array}{c} \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \end{array} + (D_{\text{LBK}} + \mathcal{S}) \begin{array}{c} \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \end{array} + \mathcal{O}(E_\gamma^0)$$



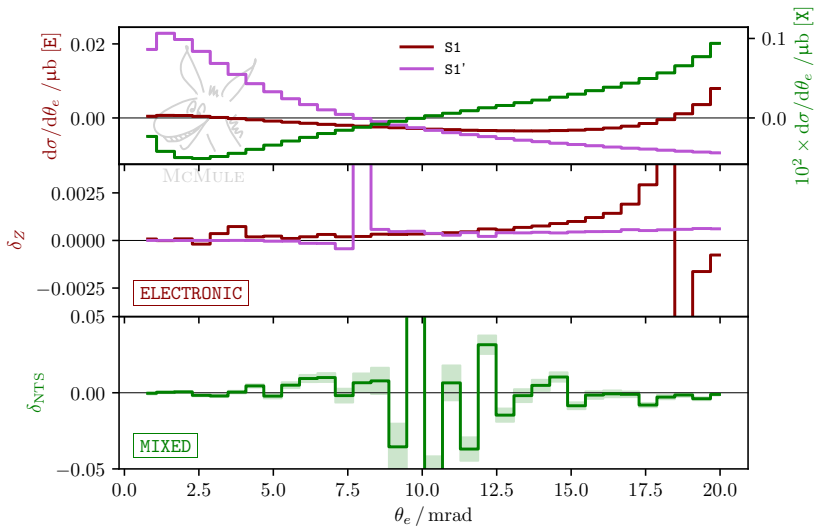
◇ introduce NTS stabilisation [McMule 21, 22]

– if $E_\gamma < E_{\text{NTS}} \sim 10^{-3} \sqrt{s}/2$
switch to the expansion above

– theory error:
 $\mathcal{O}(10^{-2})$ @ NNLO $\sim \mathcal{O}(10^{-5})$

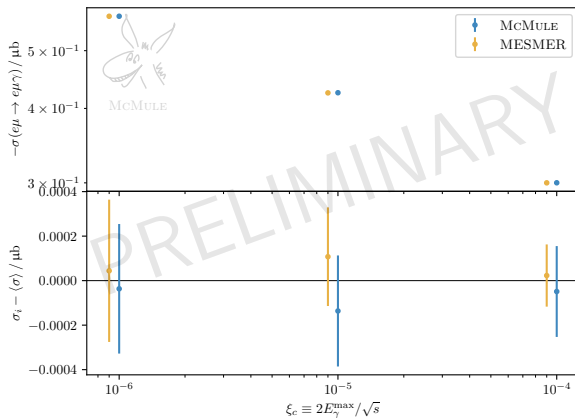
– nice(r) carbon footprint





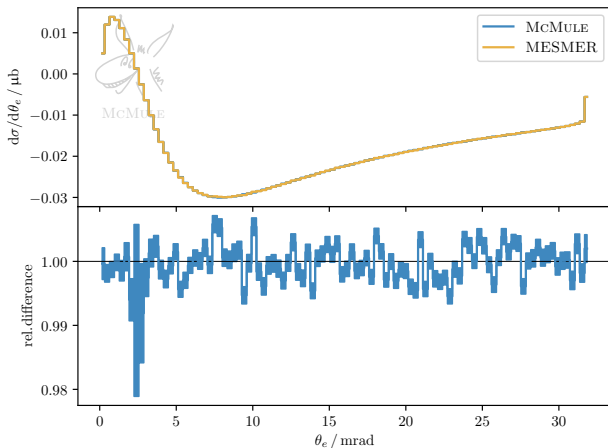
(FKS^ℓ + DIMREG) vs (slicing + m_γ)

$e\mu \rightarrow e\mu\gamma$ @ NLO with $\xi_c = \omega_s = 10^{-\{6,5,4\}}$ (MESMER as in [Carloni et al. 20])



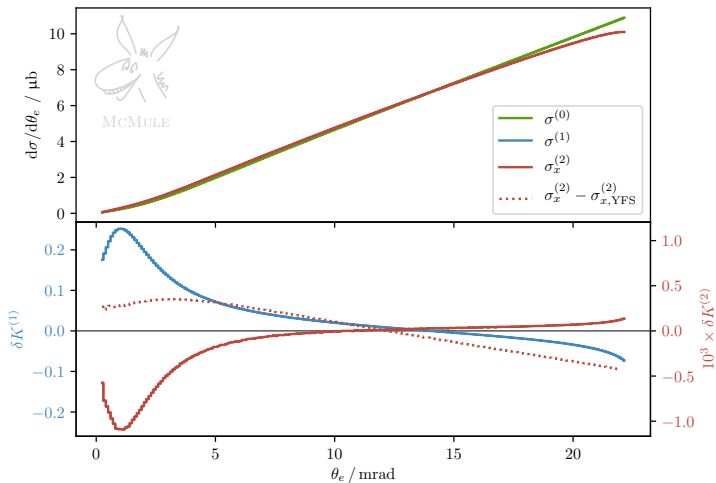
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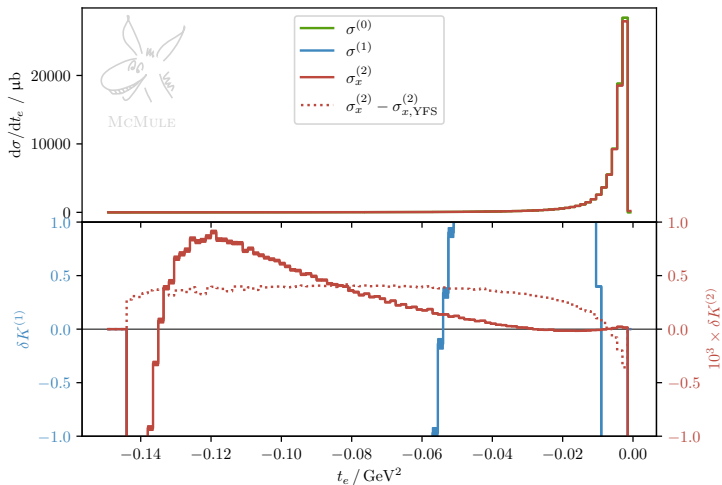


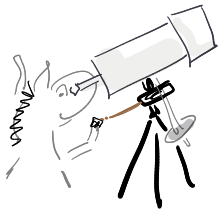
(massification) vs (YFS approximation)

- McMULE + massified VV ($m_e = 0$, [Mastrolia et al 18-22]):
 $-2.50(1) \cdot 10^{-2} \mu\text{b}$
- McMULE + YFS approximation for VV (courtesy of Carlo):
 $6.45(1) \cdot 10^{-2} \mu\text{b}$
- MESMER:
 $6.47(3) \cdot 10^{-2} \mu\text{b}$
– error estimate (LO units): $\left(\frac{\alpha}{\pi}\right)^2 \log \frac{m_e^2}{m_\mu^2} \sim 6 \cdot 10^{-4}$
[LO = $1.214(1) \cdot 10^2 \mu\text{b}$]

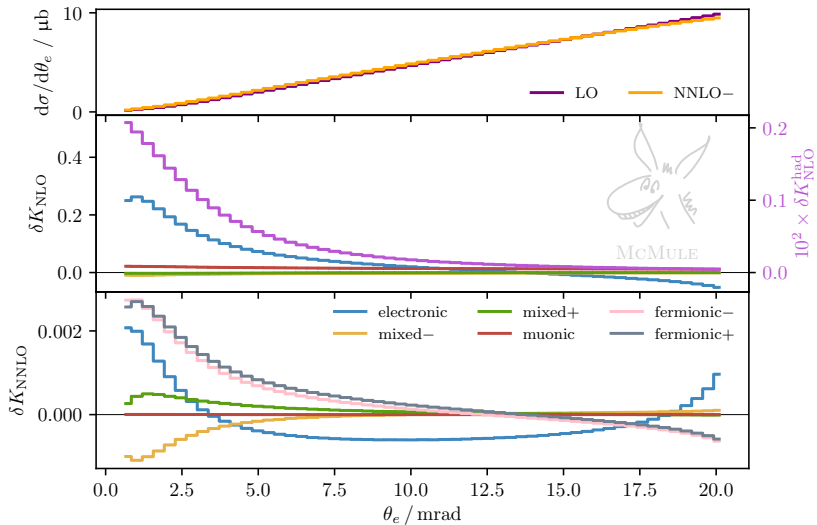
distributions: θ_e 

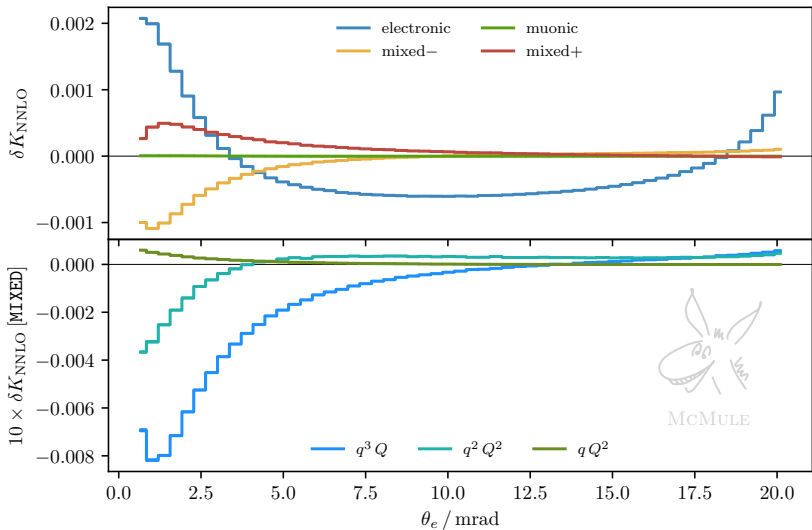
distributions:





looking for another N





Approximate N^3LO

$$d\sigma^{(3)} = \int d\Phi_n \overset{vvv}{\mathcal{M}}_n^{(3)} + \int d\Phi_{n+1} \overset{vv}{\mathcal{M}}_{n+1}^{(2)} + \int d\Phi_{n+2} \overset{rvv}{\mathcal{M}}_{n+2}^{(1)} + \int d\Phi_{n+3} \overset{rrv}{\mathcal{M}}_{n+3}^{(0)}$$

- IR subtraction with FKS^2
- $\mathcal{M}_{n+2}^{(1)}$ & $\mathcal{M}_{n+3}^{(0)}$ via OpenLoops + NTS stabilisation
- $\mathcal{M}_n^{(3)}$ via massification (if $m=0$ ✓) or eikonal or ?
- $\mathcal{M}_{n+1}^{(2)}$ via

{	NTS approximation
	collinear approximation ← \sum 2-loop massive splitting amplitude
	massification

[Tim Engel]

MCMULE is predominantly an integrator

- we can calculate $\sigma = \int d\Phi \mathcal{M} S(\{p_i\})$
- **measurement function** S can be implemented numerically \sim cuts, histograms
- event generators produce events (more or less) distributed according to $w = d\Phi \mathcal{M}$
- trivial solution: dump every event $\{p_i\}$ and weight w to file (“garden hose approach”)



[Yannick Ulrich]

minimise $\{p_i\}$ to propagate through the expensive detector simulation

- the w can be negative beyond LO and span many orders of magnitude
 - clever sampling can help but not fully solve the problem
 - if $r \times N$ of N weights are negative, we need $\propto 1/(1 - 2r)^2$ events
- ⇒ reduce r as much as possible by cancelling negative weights as early as possible

[Yannick Ulrich]

... at NLO for simplicity

$$\sigma_{\text{NLO}} = \int \text{tree} + \frac{\alpha}{4\pi} \int \text{loop} + \frac{\alpha}{4\pi} \int \text{loop}^{\text{wavy}}$$

- slicing: fairly few negative weights **but** numerically construct $\log \omega_c$

$$= \int \underbrace{\left(\text{tree} + \frac{\alpha}{4\pi} \text{loop} + \frac{\alpha}{4\pi} \int_1 \text{loop}^{\text{green}} \right)}_{\text{mostly } > 0} + \frac{\alpha}{4\pi} \int_{\omega > \omega_c} \underbrace{\text{loop}^{\text{wavy}}}_{> 0}$$

- subtraction: easier integration **but** lots and lots of negative weights ($\mathcal{O}(5\%)$ at NLO, more at NNLO)

$$= \int \underbrace{\left(\text{tree} + \frac{\alpha}{4\pi} \text{loop} + \frac{\alpha}{4\pi} \int_1 \text{loop}^{\text{green}} \right)}_{\text{mostly } > 0} + \frac{\alpha}{4\pi} \int \underbrace{\left(\text{loop}^{\text{wavy}} - \text{loop}^{\text{green}} \right)}_{\text{whatever}}$$

[Yannick Ulrich]

two observations

- ① cross section $\sigma = \int_{\mathcal{C}} d\sigma > 0$, irregardless of the size of integration region \mathcal{C}
- ② experiments have a finite resolution
(we already knew that because we can't see soft photons)

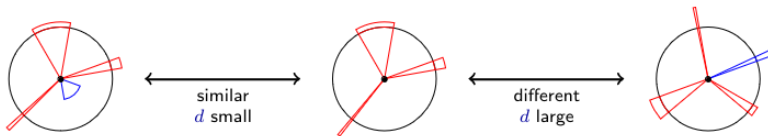
algorithm to remove negative weights [Andersen, Maier 21]

- pick an event with $w_i < 0$
- find nearby events until $\sum_{i \in \mathcal{C}} w_i > 0$
- if \mathcal{C} gets too big (events become resolvable), abort (or add more events)
- else $w_i \rightarrow \frac{\sum_{j \in \mathcal{C}} w_j}{\sum_{j \in \mathcal{C}} |w_j|} w_i$

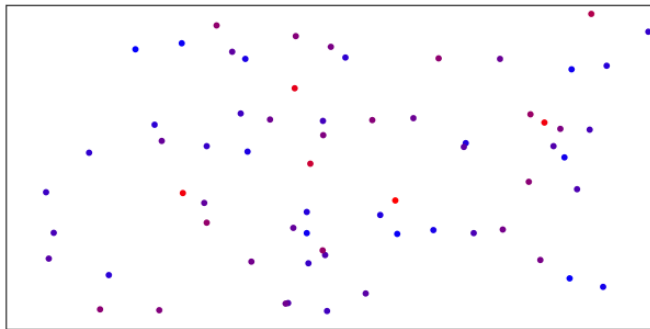
we can remove negative weights without biasing physical observables!

we need to define a **metric** in event space $d(e_1, e_2) \geq 0$

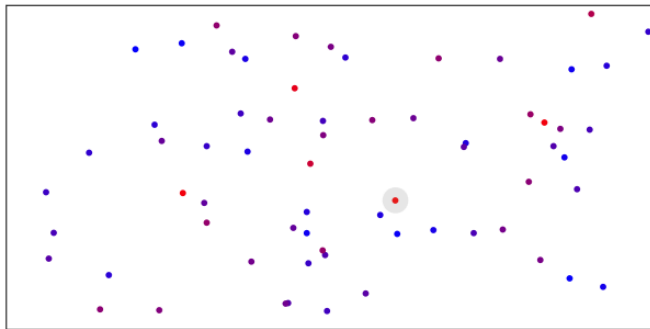
- doesn't really matter how we do this as long as IR safe
(events with soft photons are near each other)
- ideally: events that look similar are closer to each other than those that don't
- MUonE example: $d(e_1, e_2) = \sqrt{|\theta_1^e - \theta_2^e|^2 + |\theta_1^\mu - \theta_2^\mu|^2}$
- can add ϕ and/or energy information, depending on analysis



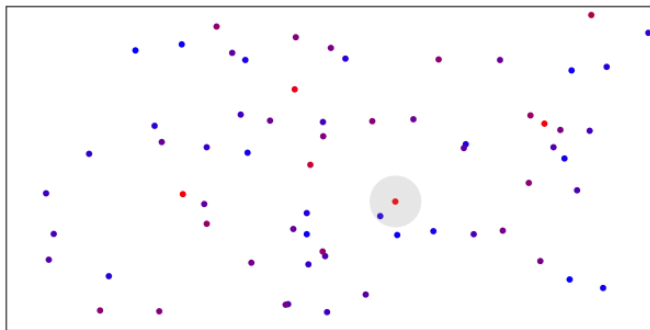
[Yannick Ulrich]



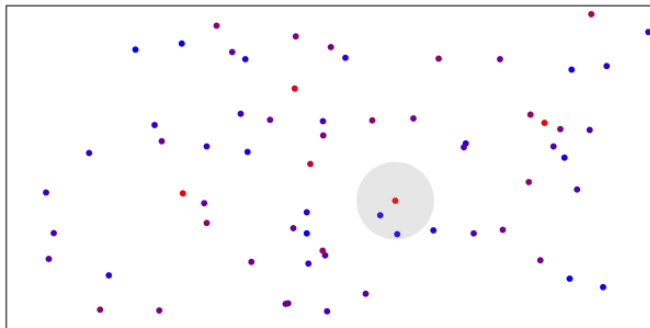
[Yannick Ulrich]



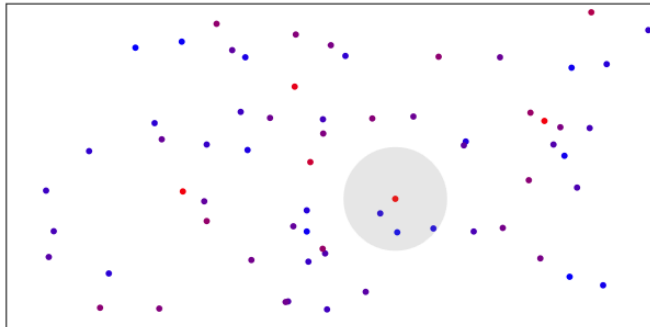
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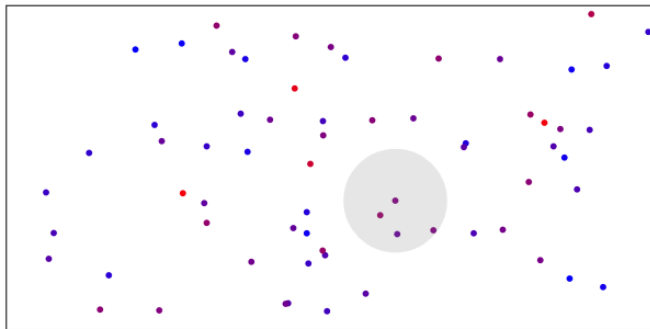
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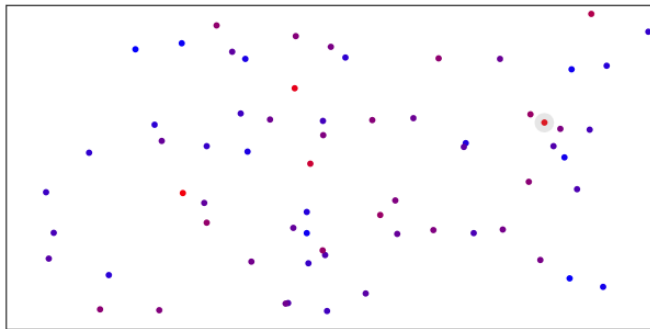
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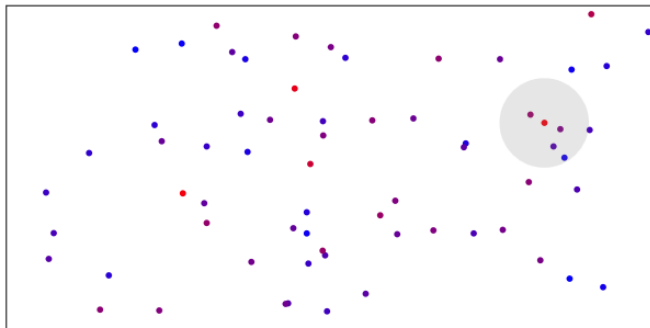
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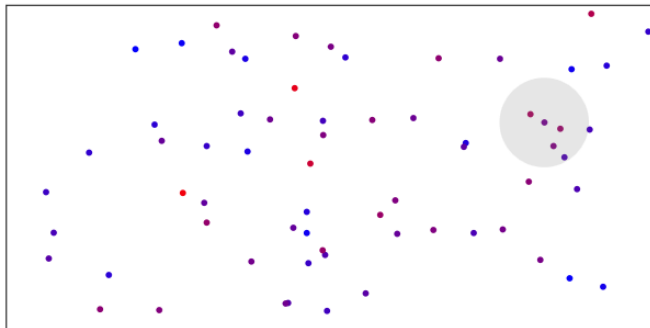
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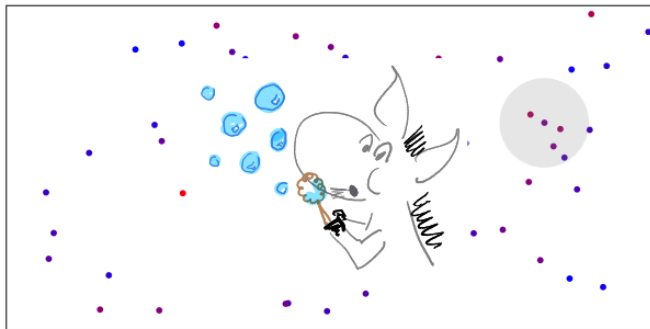
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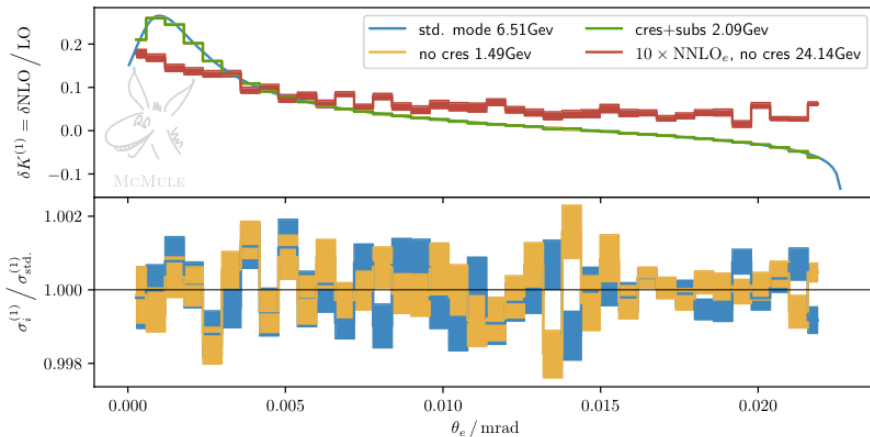
[Yannick Ulrich]



[Yannick Ulrich]



[Yannick Ulrich]



$r \approx 2 \times 10^{-2} \rightarrow 2 \times 10^{-5}$, $w_{\text{min}} / \langle w \rangle = -10^5 \rightarrow -10^{-3}$, similar at NNLO

[Yannick Ulrich]

- ◇ NNLO with different external masses [2212.06481]
- ◇ naive theory error (missing higher orders) $\mathcal{O}(10^{-5})$



- two-masses two-loop to stress test massification?
- use YFS approximation at N^3LO ?
 - feasibility: can it be used for $2 \rightarrow 3$ kinematics?
 - cross-check for RVV computed with MCMULE cut-and-patch?
- deeper insights on collinear hierarchies at N^3LO ?
- resummation (analytic & parton shower) at NLL?
- ...