Status of MUonE theory

F. Piccinini



INFN, Sezione di Pavia (Italy)

The Evaluation of the Leading Contribution to the Muon g-2 Mainz, 3-7 June 2024



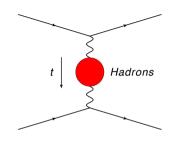
Fulvio Piccinini (INFN, Pavia) Status of MUonE theory 1/26



- G. Abbiendi, C.M. Carloni Calame, U. Marconi, C. Matteuzzi, G. Montagna, O. Nicrosini, M. Passera, F. Piccinini, R. Tenchini, L. Trentadue, G. Venanzoni,
 Measuring the leading hadronic contribution to the muon g-2 via μe scattering
 Eur. Phys. J. C 77 (2017) no.3, 139 arXiv:1609.08987 [hep-ph]
- ★ C. M. Carloni Calame, M. Passera, L. Trentadue and G. Venanzoni, A new approach to evaluate the leading hadronic corrections to the muon g-2 Phys. Lett. B 746 (2015) 325 - arXiv:1504.02228 [hep-ph]

$$a_{\mu}^{\mathrm{HLO}} = \frac{\alpha}{\pi} \int_{0}^{1} dx (1-x) \Delta \alpha_{\mathrm{had}}[t(x)]$$

$$t(x) = \frac{x^{2} m_{\mu}^{2}}{x-1} < 0$$



e.a. Lautrup, Peterman, De Rafael, Phys. Rept. 3 (1972) 193

- The hadronic VP correction to the running of α enters
- $\Delta \alpha_{\rm had}(t)$ can be directly measured in a (single) experiment involving a space-like scattering process and $\mathbf{a}_{\mu}^{\mathrm{HLO}}$ obtained through numerical integration Carloni Calame, Passera, Trentadue, Venanzoni PLB 746 (2015) 325

* A data-driven evaluation of a_{μ}^{HLO} , but with space-like data

Kernel functions for $\mathbf{a}_{\mu}^{\mathbf{HVP}}$

• LO:
$$\frac{\alpha}{\pi}(1-x)$$



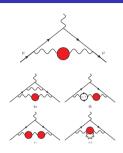
Kernel functions for $\mathbf{a}_{\mu}^{\mathbf{HVP}}$

- LO: $\frac{\alpha}{\pi}(1-x)$
- NLO

E. Balzani, S. Laporta, M. Passera, Phys. Lett. B834 (2022) 137462

A.V. Nesterenko, J. Phys. G49 (2022) 5, 055001;

J. Phys. G50 (2022) 2, 029401



Kernel functions for $\mathbf{a}_{\mu}^{\mathbf{HVP}}$

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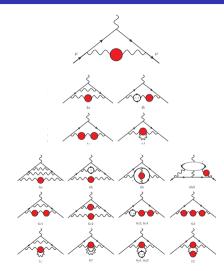
E. Balzani, S. Laporta, M. Passera, Phys. Lett. B834 (2022) 137462

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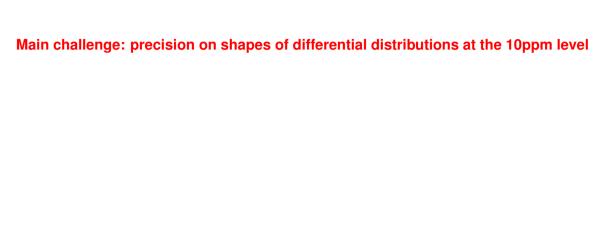
J. Phys. G50 (2022) 2, 029401

NNLO

E. Balzani, S. Laporta, M. Passera, Phys. Lett. B834 (2022) 137462



⇒ talk by S. Laporta



Radiative corrections to the signal

Radiative corrections to the signal

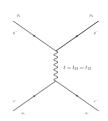
Predictions for Background processes

Radiative corrections to the signal

Predictions for Background processes

High precision Monte Carlo simulation tools required

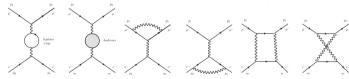
First step towards precision: QED NLO and MC (2018)



analytical expression for tree level

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2}{\lambda(s, m_{\mu}^2, m_e^2)} \left[\frac{(s - m_{\mu}^2 - m_e^2)^2}{t^2} + \frac{s}{t} + \frac{1}{2} \right]$$

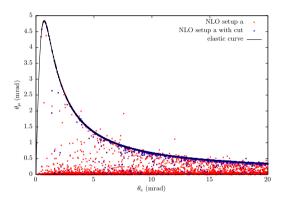
- VP gauge invariant subset of NLO rad. corr.
- factorized over tree-level: $\alpha \rightarrow \alpha(t)$
- ullet QED NLO virtual diagrams and real emission diagrams with exact finite m_e and m_μ effects

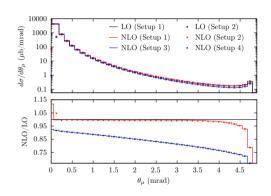


- tree-level Z-exchange important at the 10^{-5} level ($\sim tG_{\mu}/4\pi\alpha\sqrt{2}$ in the Fermi theory)
- SM weak RCs at most at a few 10^{-6} level, negligible

Alacevich et al. JHEP 02 (2019) 155

First realistic description of scattering events

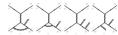




- many points fall out of the $2 \to 2$ correlation curve $\theta_\mu \theta_e$ because of the radiative events
- NLO QED radiative corrections at the % level, enhanced by exclusive event selections

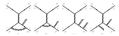
exact calculation of corrections along one lepton line with all finite mass effects





exact calculation of corrections along one lepton line with all finite mass effects





• two independent calculations, with different IR singularities handling procedures (slicing and subtraction)

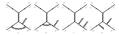
Carloni Calame et al., JHEP 11 (2020) 028,

P. Banerjee, T. Engel, A. Signer, Y. Ulrich, SciPost Phys. 9 (2020) 027

• implemented in Mesmer and McMule, perfect numerical agreement

exact calculation of corrections along one lepton line with all finite mass effects





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Carloni Calame et al., JHEP 11 (2020) 028,

P. Banerjee, T. Engel, A. Signer, Y. Ulrich, SciPost Phys. 9 (2020) 027

- implemented in Mesmer and McMule, perfect numerical agreement
- NNLO with finite mass effects and approximate up-down interference in Mesmer
 - interference of LO $\mu e
 ightarrow \mu e$ amplitude with



NNLO double-virtual amplitudes where at least 2 photons connect the e and μ lines are approximated according to the Yennie-Frautschi-Suura ('61) formalism to catch the IR divergent structure

• complete calculation of the amplitude $f^+f^- o F^+F^-$ with $m_f=0$, $m_F
eq 0$

R. Bonciani et al., PRL 128 (2022)

- complete calculation of the amplitude $f^+f^- o F^+F^-$ with $m_f=0$, $m_F
 eq 0$
- "massification" to recover the leading m_e terms, i.e. neglecting powers of m_e^2/Q^2

T. Engel, C. Gnendiger, A. Signer and Y. Ulrich, JHEP 02 (2019) 118

Y. Ulrich, PoS RADCOR2023 (2024) 077

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FKS^ℓ subtraction scheme

T. Engel, A. Signer, Y. Ulrich, JHEP 01 (2020) 085

- complete calculation of the amplitude $f^+f^- \to F^+F^-$ with $m_f=0$, $m_F \neq 0$
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 Next-to-soft stabilisation, to obtain numerical stability in real-virtual corrections with soft and/or collinear photon configurations

T. Engel, A. Signer, Y. Ulrich, JHEP 04 (2022) 097; T. Engel, JHEP 07 (2023) 177

Fulvio Piccinini (INFN, Pavia) Status of MUonE theory 9/26

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T. Engel, A. Signer, Y. Ulrich, JHEP 04 (2022) 097; T. Engel, JHEP 07 (2023) 177

- · with the above ingredients
 - ullet NNLO calculation neglecting terms of $\mathcal{O}(m_e^2/Q^2)$ in <code>McMule</code>

A. Broggio et al., JHEP 01 (2023) 112

⇒ talk by M. Rocco

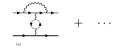
Fulvio Piccinini (INFN, Pavia) Status of MUonE theory 9/26

NNLO virtual leptonic pairs (vacuum polarization insertion) (2021)

- any lepton (and hadron) in the VP blobs
- interfered with $\mu e \rightarrow \mu e$ or $\mu e \rightarrow \mu e \gamma$ amplitudes









• interfered with $\mu e
ightarrow \mu e$ amplitude



2-loop integral evaluated with dispersion relation techniques in Mesmer

used e.g. in the past for Bhabha: Actis et al., Phys. Rev. Lett. 100 (2008) 131602; Carloni Calame et al., JHEP 07 (2011) 126

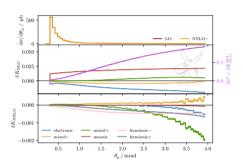
$$\frac{g_{\mu\nu}}{q^2+i\epsilon} \rightarrow g_{\mu\nu} \frac{\alpha}{3\pi} \int_{4m_\ell^2}^{\infty} \frac{dz}{z} \frac{R_\ell(z)}{q^2-z+i\epsilon} = g_{\mu\nu} \frac{\alpha}{3\pi} \int_{4m_\ell^2}^{\infty} \frac{dz}{z} \frac{1}{q^2-z+i\epsilon} \left(1 + \frac{4m_\ell^2}{2z}\right) \sqrt{1 - \frac{4m_\ell^2}{z}}$$

• 2-loop integral evaluated (also) with hyperspherical method in McMule

M. Fael, JHEP02 (2019) 027

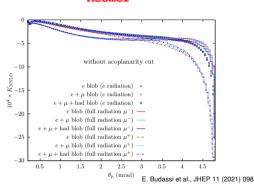
NNLO order of magnitude

McMule



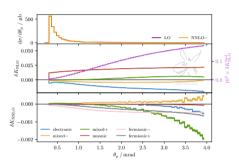
A. Broggio et al., JHEP 01 (2023) 112

Mesmer



NNLO order of magnitude

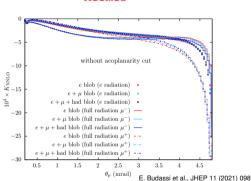
McMule



A. Broggio et al., JHEP 01 (2023) 112

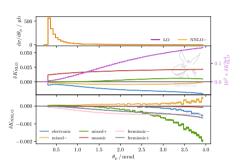
• NNLO corrections at the $10^{-4}-10^{-3}$ level

Mesmer

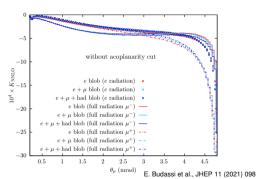


NNLO order of magnitude

McMule







A. Broggio et al., JHEP 01 (2023) 112

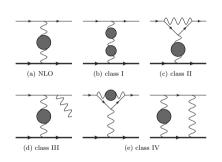
- NNLO corrections at the $10^{-4}-10^{-3}$ level
- eventually fixed order calculations need to be matched to resummation of higher order corrections, through PS techniques (e.g. BaBayaga) or YFS techniques (e.g. KKMC/SHERPA)

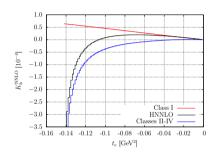
⇒ talk by A. Price

11/26

NNLO hadronic contributions (2019)

using the dispersion relation approach





Fael, Passera, Phys. Rev. Lett. 122 (2019) 192001

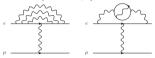
- corrections of the order of 10⁻⁴
- hyperspherical integration method to calculate hadronic NNLO corrections, where the hadronic vacuum polarization is employed in the space-like region (used in McMule)

 M. Fael, JHEP02 (2019) 027

Towards N³LO on the electron line

Y. Ulrich, N³LO kick-off workstop/thinkstart, Durham, 3-5 August 2022

• All virtual (three loops)



• Single real emission (two loops)



• Double real emission (one loops)

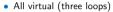


Triple real



M. Fael, MUonE Collaboration Meeting, 16/05/2023, CERN

Y. Ulrich, N³LO kick-off workstop/thinkstart, Durham, 3-5 August 2022





• Single real emission (two loops)



Double real emission (one loops)



Triple real



M. Fael, MUonE Collaboration Meeting, 16/05/2023, CERN

 this contribution will allow improved perturbative predictions and more reliable theoretical uncertainty estimates

Recent progress

the three-loop form factor with finite fermion mass is now available

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev. Lett 128 (2022) 172003

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev.D 106 (2022) 034029

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev.D 107 (2023) 094017

All order subtraction scheme FKS^ℓ availale

T. Engel, A. Signer, Y. Ulrich, JHEP 01 (2020) 085

 very recent generalisation of the LBK theorem to multi-photon emission ⇒ extension of next-to-soft stabilisation to multiple radiation

T. Engel, JHEP 03 (2024) 004

ullet real-virtual-virtual corrections recently recalculated with $m_e o 0$

S. Badger, J. Krys, R. Moodle, S. Zoia, JHEP 11 (2023) 041

V.S. Fadin, R.N. Lee, JHEP 11 (2023) 148



Fixed target experiment \Longrightarrow bound electron effects

very recently estimated

R. Plestid and M.B. Wise, arXiv:2403.12184

Fixed target experiment ⇒ bound electron effects

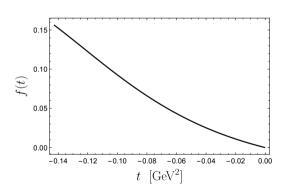
very recently estimated

R. Plestid and M.B. Wise, arXiv:2403.12184

ullet for C

$$\frac{1}{\sigma} \frac{d\sigma}{dt} = \frac{1}{\sigma^0} \frac{d\sigma^0}{dt} \left(1 - Kf(t) \right)$$

• $K = 4.5 \cdot 10^{-4}$, scaling as $1/Z_A$

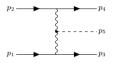


⇒ talk by R. Plestid today

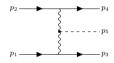
15/26

- pion pair production forbidden kinematically with the available \sqrt{s}

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- single π^0 production possible



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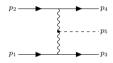


• π^0 production calculated and shown to be well below 10^{-5} w.r.t. $\mu e o \mu e$

E. Budassi et al., PLB 829 (2022) 137138

Backgrounds

- pion pair production forbidden kinematically with the available \sqrt{s}
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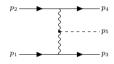
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E. Budassi et al., PLB 829 (2022) 137138

lepton pair production

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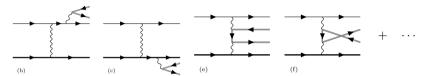
E. Budassi et al., PLB 829 (2022) 137138

lepton pair production

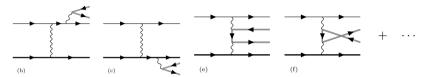
•
$$\mu^{\pm}e^{-} \rightarrow \mu^{\pm}e^{-}\ell^{+}\ell^{-}$$

•
$$\mu^{\pm}N \rightarrow \mu^{\pm}N\ell^{+}\ell^{-}$$

• it also contributes at NNLO accuracy



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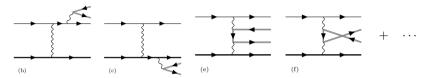


• the emission of an extra electron pair $\mu e \to \mu e \ e^+ e^-$ is potentially a dramatically large background, because of the presence of "peripheral" diagrams which develop powers of collinear logarithms upon integration

G. Racah, Il Nuovo Cimento 14 (1937) 83-113; L.D. Landau, E.M. Lifschitz, Phys. Z. Sowjetunion 6 (1934) 244; H.J. Bhabha, Proc. Roy. Soc. Lond. A152 (1935) 559;

R.N. Lee, A.A. Lyubyakin, V.A. Smirnov, Phys. Lett. B 848 (2024) 138408

it also contributes at NNLO accuracy



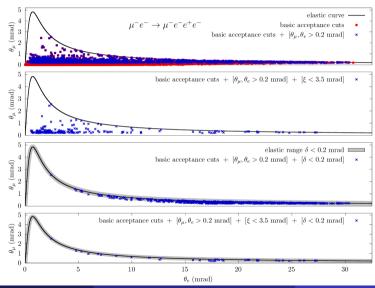
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• $\mu^\pm e^- o \mu^\pm e^- \ell^+ \ell^-$ calculated with finite mass effects and implemented in Mesmer

simulation of $5\cdot 10^5$ points of $\mu^\pm e^- o \mu^\pm e^- \ell^+ \ell^-$



Real pair emission from scattering on nucleus: $\mu^{\pm}N o \mu^{\pm}N\ell^{+}\ell^{-}$

G. Abbiendi et al., Phys. Lett B854 (2024) 138720

• it can mimic the signal if one particle is not reconstructed or two tracks overlap within resolution

Real pair emission from scattering on nucleus: $\mu^{\pm}N o \mu^{\pm}N\ell^{+}\ell^{-}$

G. Abbiendi et al., Phys. Lett B854 (2024) 138720

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ullet GEANT4: "for the process of e^+e^- pair production the muon deflection is neglected"

A.G. Bogdanov et al., IEEE transactions on nuclear science, 53, n. 2, April 2006

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A.G. Bogdanov et al., IEEE transactions on nuclear science, 53, n. 2, April 2006

 \implies a dedicated calculation implemented in the Monte Carlo generator ${ t Mesmer}$

Fulvio Piccinini (INFN, Pavia) Status of MUonE theory 20/26

• approximation: scattering on the external nucleus field

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- finite extension of the nucleus through a form factor

$$F_Z(q) = \frac{1}{Ze} \int_0^\infty dr \, r^2 \rho_Z(r) \frac{\sin(qr)}{qr}$$

- q : momentum transferred to the nucleus
- ρ_Z : nuclear charged density

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- q: momentum transferred to the nucleus
- ρ_Z : nuclear charged density
- different models for charge density

J. Heeck, R. Szafron, Y. Uesaka, PRD 105 (2022) 053006

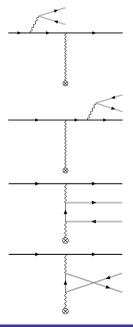
- $F_Z(q) = 1$ (conservative)
- 1 parameter Fermi model (1pF)

$$\rho_Z(r) = \frac{\rho_0}{1 + \exp\frac{r - c}{z}}$$

Fourier Bessel expansion (FB)

$$\rho_Z(r) = \sum_{k=0}^{n} a_k j_0 \left(\frac{k\pi r}{R}\right), \quad r \ge R$$
$$= 0 > R$$

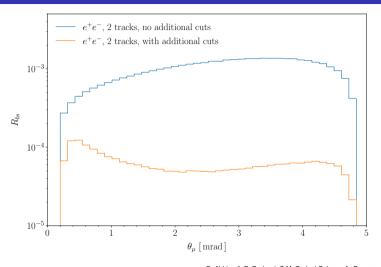
modified-harmonic oscillator model



20/26

Status of MUonE theory

Background/signal ratio



G. Abbiendi, E. Budassi, C.M. Carloni Calame, A. Gurgone, F.P., Phys.Lett.B 854 (2024) 138720

 \Longrightarrow talk by A. Gurgone

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

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• Effects of heavy $(M_{NP}\gg 1~{\rm GeV})$ NP mediators investigated through EFT with dim-6 operators

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

- Effects of heavy $(M_{NP}\gg 1~{\rm GeV})$ NP mediators investigated through EFT with dim-6 operators
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- Effects of heavy $(M_{NP}\gg 1~{\rm GeV})$ NP mediators investigated through EFT with dim-6 operators
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- Effects of **light** $(M_{NP} \le 1 \text{ GeV})$ NP mediators investigated with spin-dependent general models

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

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 - spin—0 NP mediators (ALPs)
 - spin-1 NP mediators (Dark Photons, light Z^\prime vector bosons)

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

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HVP determination with MUonE data will be robust against New Physics

Possible New Physics studies with MUonE (in complementary regions to $\Delta \alpha_h$)

- interesting proposals for NP searches at MUonE (new light mediators) in $2 \to 3$ processes
 - invisibly decaying light Z' in $\mu e \to \mu e Z'$

Asai et al., Phys. Rev. D106 (2022) 5

- a relevant background can be $\mu e o \mu e \pi^0$, in addition to $\mu e o \mu e \gamma$
- long-lived mediators with displaced vertex signatures $\mu e \to \mu e A' \to \mu e e^+ e^-$

Galon et al., Phys.Rev.D 107 (2023) 095003

• through scattering off the target nuclei $\mu N \to \mu N X \to \mu N e^+ e^-$

Grilli di Cortona and E. Nardi, Phys. Rev. D105 (2022) L111701

Summary

- Given its precision requirements, MUonE represents a challenge for
 - QFD corrections
 - background calculation
- at present we have two independent Monte Carlo tools, Mesmer and McMule featuring
 - NLO QED corrections
 - NNLO QED corections from single lepton legs
 - YFS inspired approximation to the full NNLO QED in Mesmer
 - full NNLO QED with electron "massification" in McMule
 - pair production in Mesmer
 - $\mu^{\pm}e^{-} \rightarrow \mu^{\pm}e^{-}\ell^{+}\ell^{-}$
 - $\mu^{\pm}N \rightarrow \mu^{\pm}N\ell^{+}\ell^{-}$
- efforts for N³LO started
- work in progress for matching with higher order QED corrections

Theoretical progress, thanks also to past

- MUonE theory workshops
 - Theory Kickoff Workshop, Padova, 4-5 September 2017
 - MITP Workshop, Mainz 19-23 February 2018
 - 2nd Workstop/ThinkStart, Zürich, 4-7 February 2019
 - N³LO kick-off workstop/thinkstart IPPP Durham, 3-5 August 2022
 - MITP Workshop, Mainz 14-18 November 2022
- Five General MUonE Collaboration Meetings

A collection of references on calculation developments

- → Carloni Calame et al., PLB 746 (2015), 325
- → Abbiendi et al., EPJ C77 (2017), 139
- → Mastrolia et al., JHEP 11 (2017) 198
- → Di Vita et al., JHEP 09 (2018) 016
- → Alacevich et al., JHEP 02 (2019) 155
- → Fael and Passera, PRL 122 (2019) 19, 192001
- → Fael, JHEP 02 (2019) 027
- → Engel et al., JHEP 02 (2019) 118
- → Engel et al., JHEP 01 (2020) 085
- → Carloni Calame et al., JHEP 11 (2020) 028
- → Banerjee et al., SciPost Phys. 9 (2020), 027
- → Banerjee et al., EPJC 80 (2020) 6, 591
- → Budassi et al., JHEP 11 (2021) 098
- → Balzani et al., PLB 834 (2022) 137462

- → Bonciani et al., PRL 128 (2022) 2, 022002
- Budassi et al., PLB 829 (2022) 137138
- → Engel et al., JHEP 04 (2022) 097
- → Fael et al., PRL 128 (2022) 172003
- → Fael et al., PRD 106 (2022) 034029
- → Broggio et al., JHEP 01 (2023) 112
- → Fael et al., PRD 107 (2023) 094017
- → Engel, JHEP 07 (2023) 177
- → Badger et al., JHEP 11 (2023) 041
- → Fadin and Lee., JHEP 11 (2023) 148
- → Ahmed et al., JHEP 01 (2024) 010
- → Engel, JHEP 03 (2024) 004
- → Abbiendi et al., PLB 854 (2024) 138720
- → Plestid and Wise, arXiv:2403.12184