

# GRAVITATIONAL WAVES FROM A DARK ~~SU(3)~~ SU(N) PHASE TRANSITION

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Based on work with:  
[Enrico Morgante](#)  
[Nicklas Ramberg](#)



The Future of Fundamental Composite  
Dynamics: Colliders, Cosmology, Tools

MITP Mainz

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# Overview

Motivation for confining/QCD-like dark sectors

The confinement/deconfinement phase transition (PT)

- ▶ Gravitational waves from a first order PT
- ▶ The NANOGrav GW hint for a low scale PT

Towards quantitative predictions of GWs in strongly coupled PTs

With Enrico Morgante, Nicklas Ramberg, 2210.11821

# Why should you care about dark SU(N)

## Top down perspective:

- ▶ Many string compactifications contain hidden sectors with new gauge symmetries
- ▶ Straightforward extension of SM

## Useful for model building

- ▶ LHC-safe solutions to the hierarchy problem (twin Higgs...)
- ▶ Axion models and composite axions, relaxion

## Interesting (&new) phenomenology

- ▶ DM with self-interactions, SIMP mechanism
- ▶ Unique collider signatures...

# Composite/QCD-like DM

Alternative to elementary WIMP models

Phenomenologically viable, “generic” possibility in presence of hidden sectors

Some nice features:

- ▶ DM stability, mass scale
- ▶ Self interactions, unique collider pheno
- ▶ Natural implementation of SIMP mechanism (3- $\rightarrow$ 2 annihilation)
- ▶ Glueball dark matter

e.g. Bai, PS, 2014  
PS, Stolarski, Weiler, JHEP 2015

Hochberg, Kuflik, Murayama,  
Volansky, Wacker, 2014

...

e.g. Soni, Zhang, 2016  
Asadi, Kramer, Kuflik, Slatyer, Smirnov, 2022  
Carenza, Pasechnik, Salinas, Wang, 2022

# Models I'm interested here

Nonabelian  $SU(N)$  dark sector, confinement scale  $\Lambda_d$

$n_f$  light/massless **dark quarks**

$$n_f = 0$$

Glueball DM

PT from center  
symmetry restoration

$$n_f > 0$$

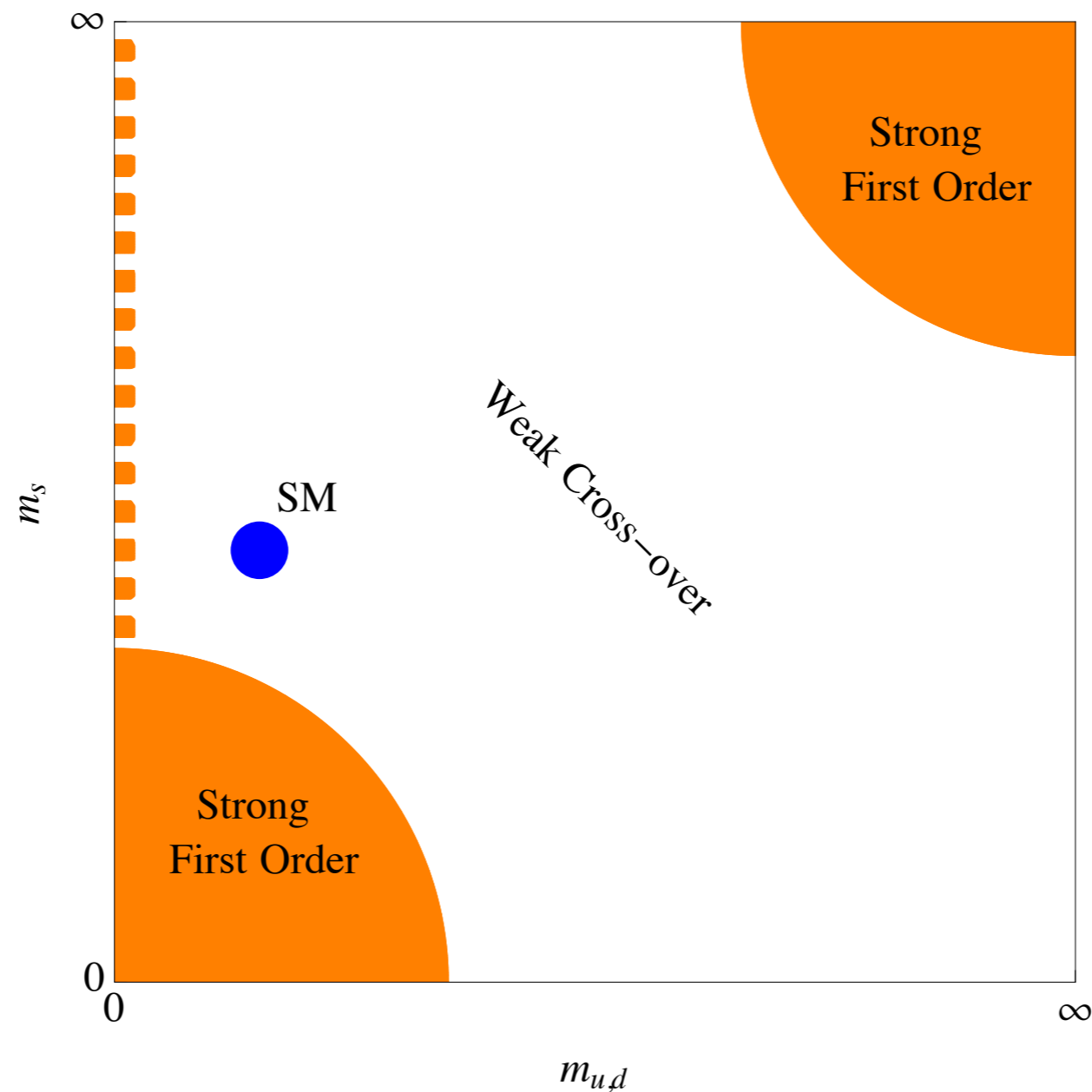
Dark Baryons  
or Dark Pions

Chiral Symmetry Breaking

A new PT is a robust prediction of these scenarios

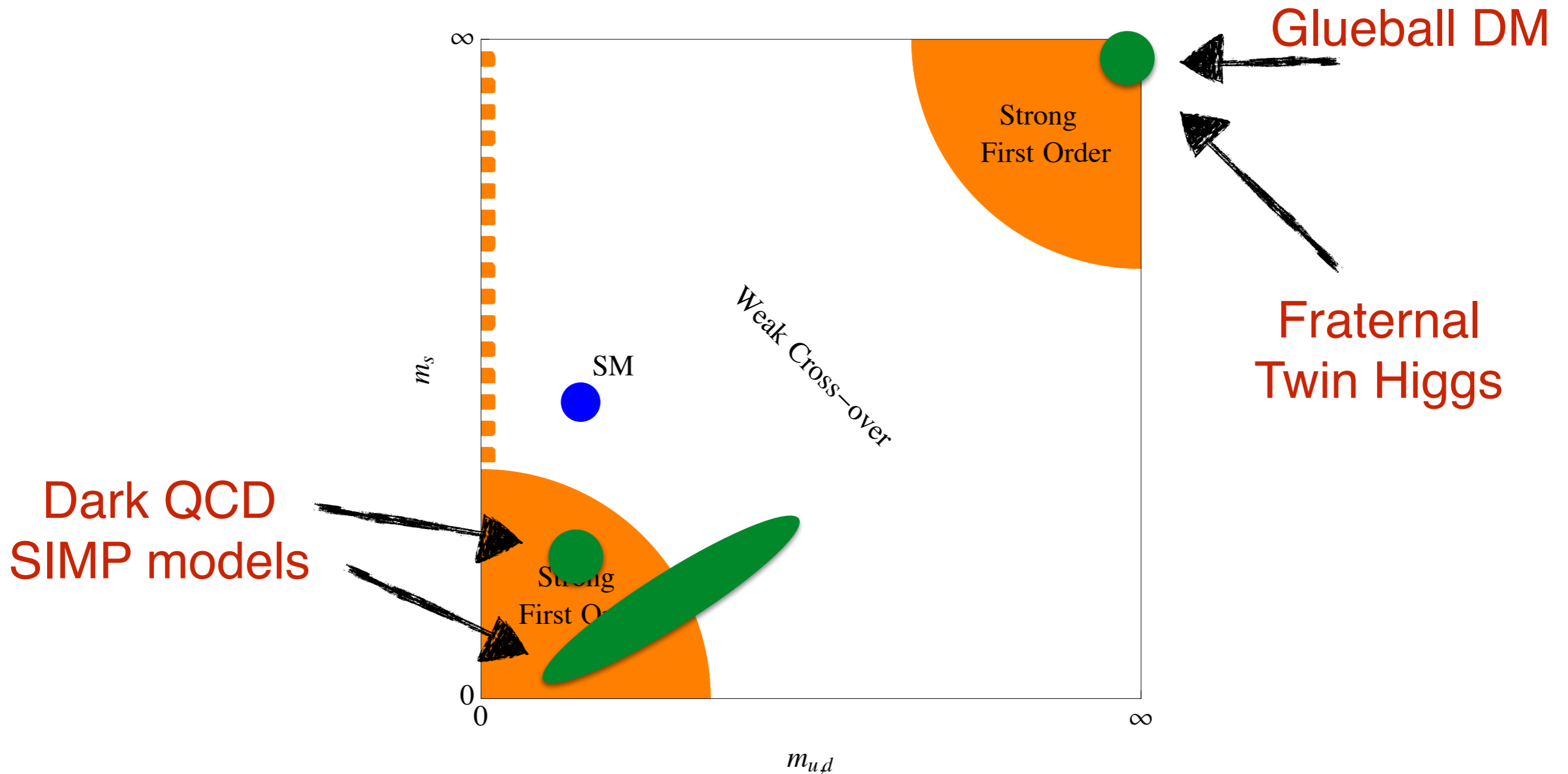
# What is the nature of the PT?

QCD phase diagram



PS, 2016

# Phase Diagram II



PS, 2016

# SU(N) - PT

Consider.  $SU(N_d)$  with  $n_f$  massless flavours

PT is first order for

▶  $N_d \geq 3$  ,  $n_f = 0$

Svetitsky, Yaffe, 1982  
M. Panero, 2009

▶  $N_d \geq 3$  ,  $3 \leq n_f < 4N_d$

Pisarski, Wilczek, 1983

Not for:

▶  $n_f = 1$  (no global symmetry, no PT)

▶  $n_f = 2$  (not yet known)

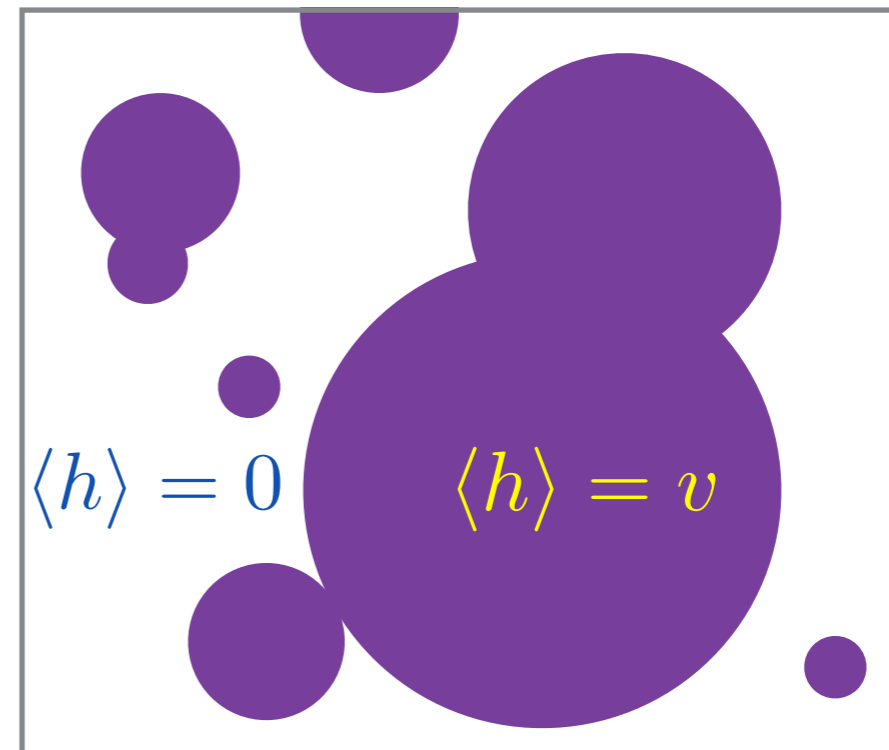
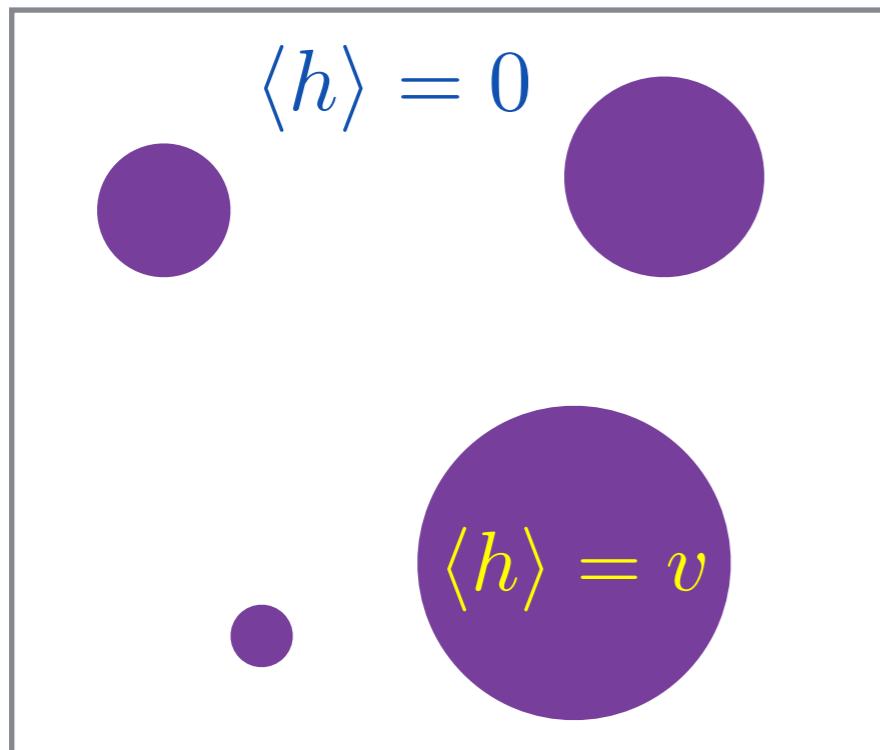
Note: Nature of the PT does not depend on arbitrary model parameters



# First order phase transitions produce GWs

First order PT  $\rightarrow$  Bubbles nucleate, expand

Bubble collisions  $\rightarrow$  Gravitational Waves



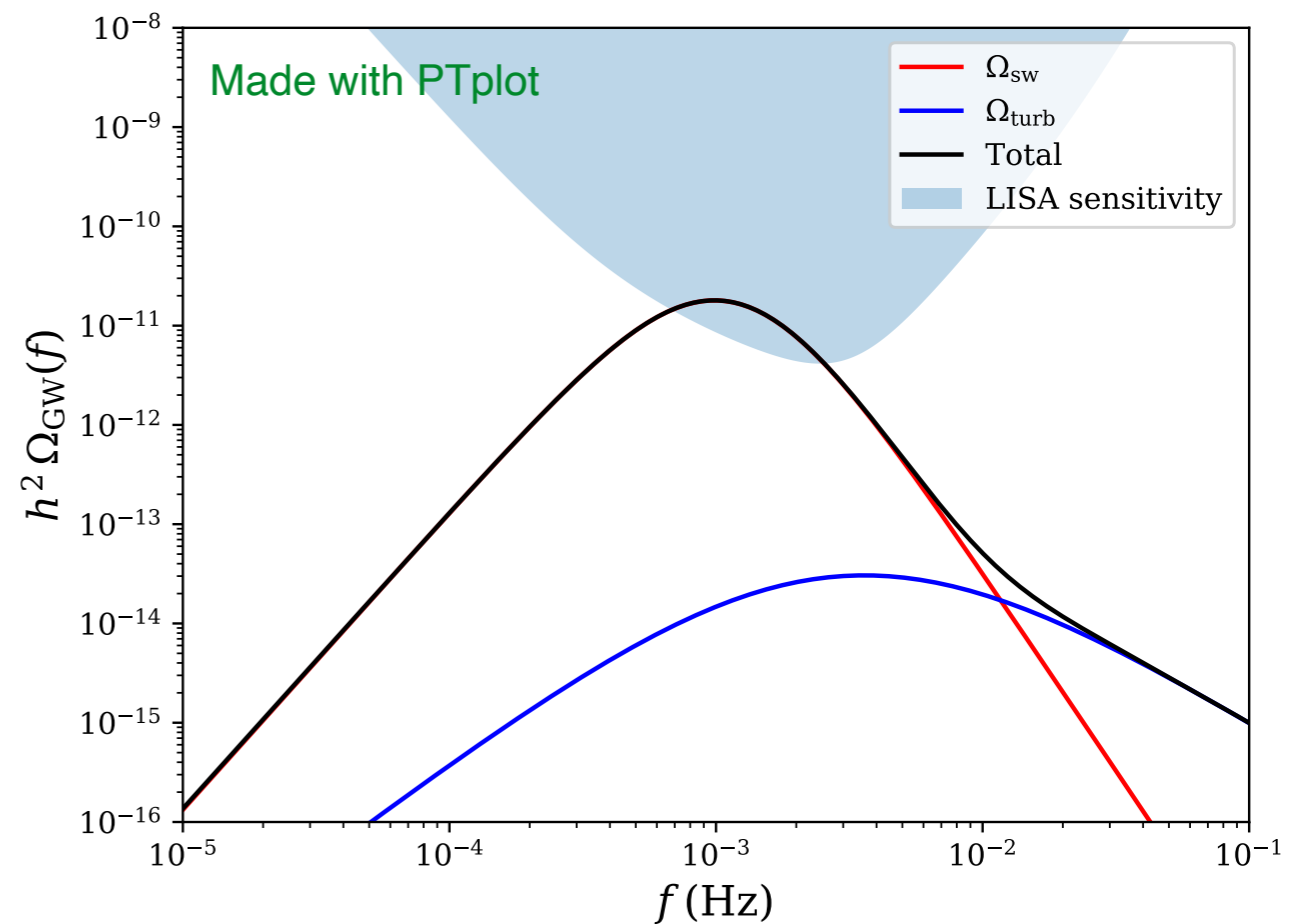
# PT signal

PT characterised by few parameters:

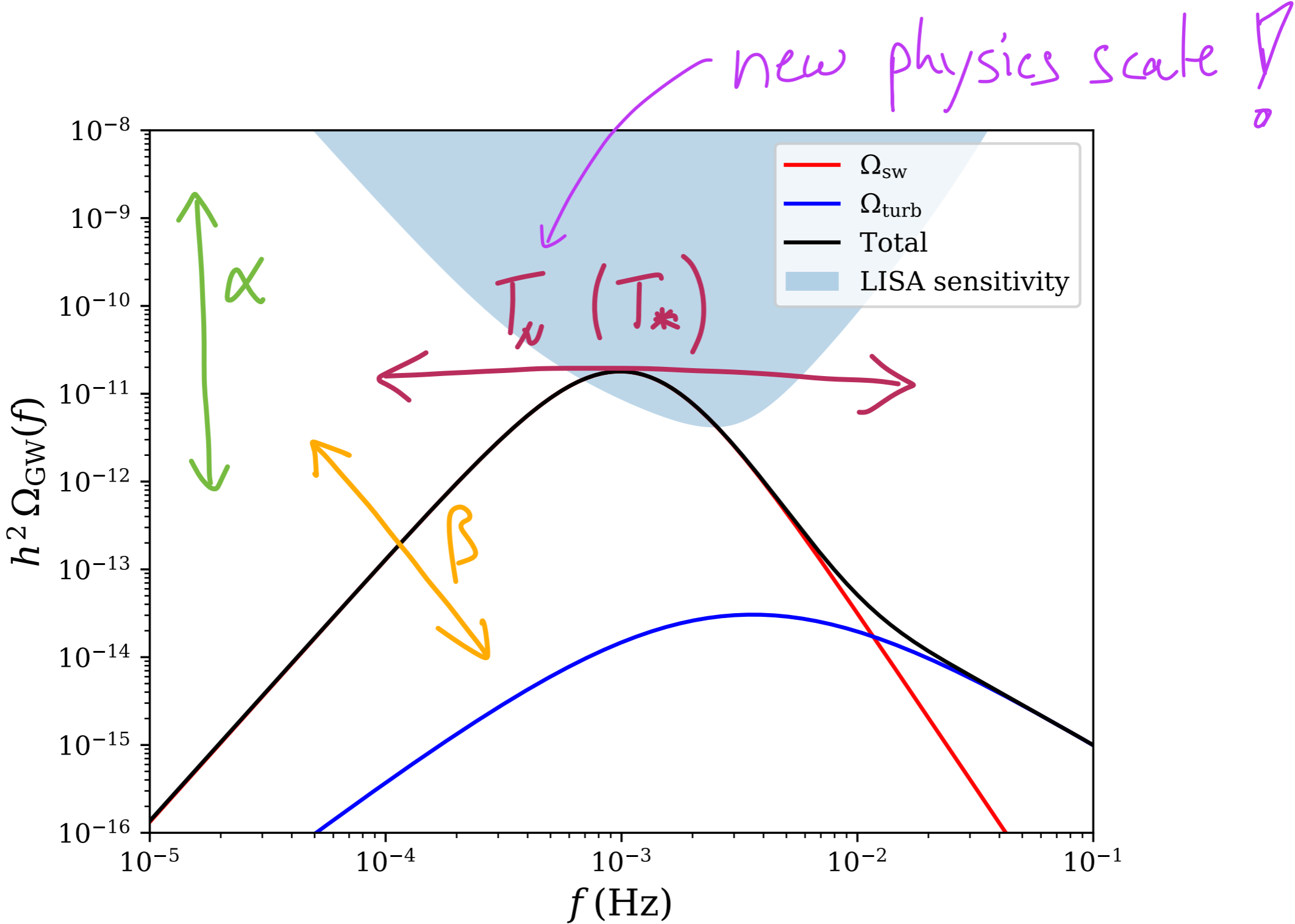
- Latent heat  $\alpha \approx \frac{\Omega_{\text{vacuum}}}{\Omega_{\text{rad}}}$
- Bubble wall velocity  $v$
- Bubble nucleation rate  $\beta$
- PT temperature  $T_*$

More details, see e.g.:

Summary and recommendations:  
1910.13125  
(LISA Cosmology WG)



# Signal properties



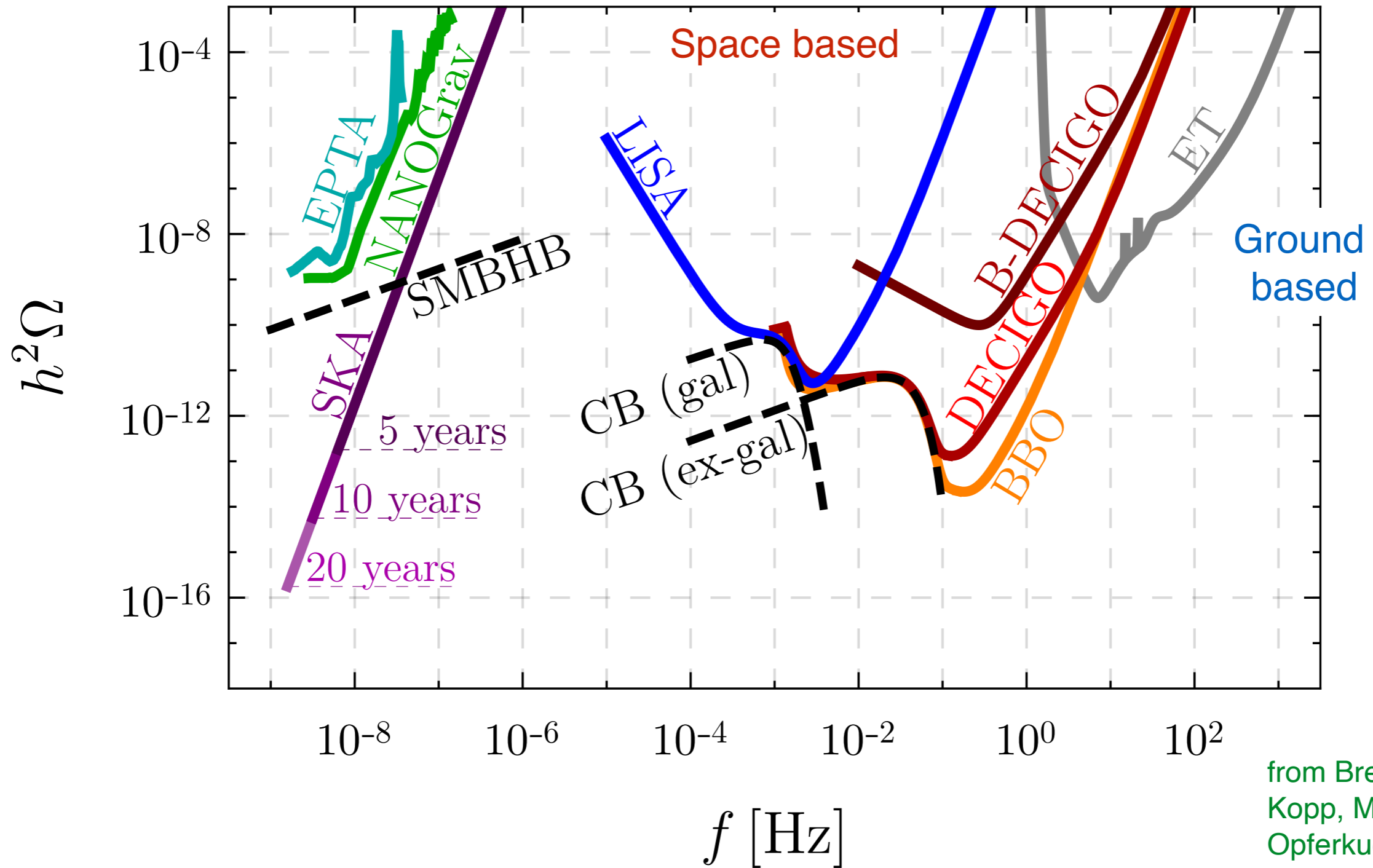
# There are many (planned) experiments

New physics scale

GeV

TeV

PeV



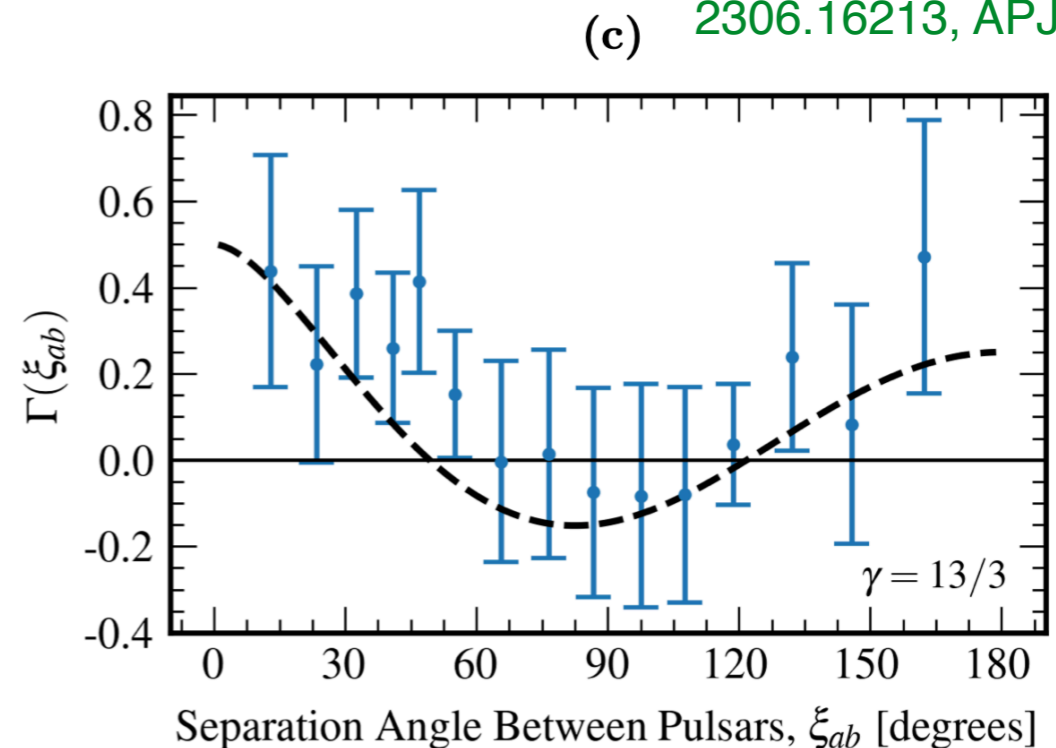
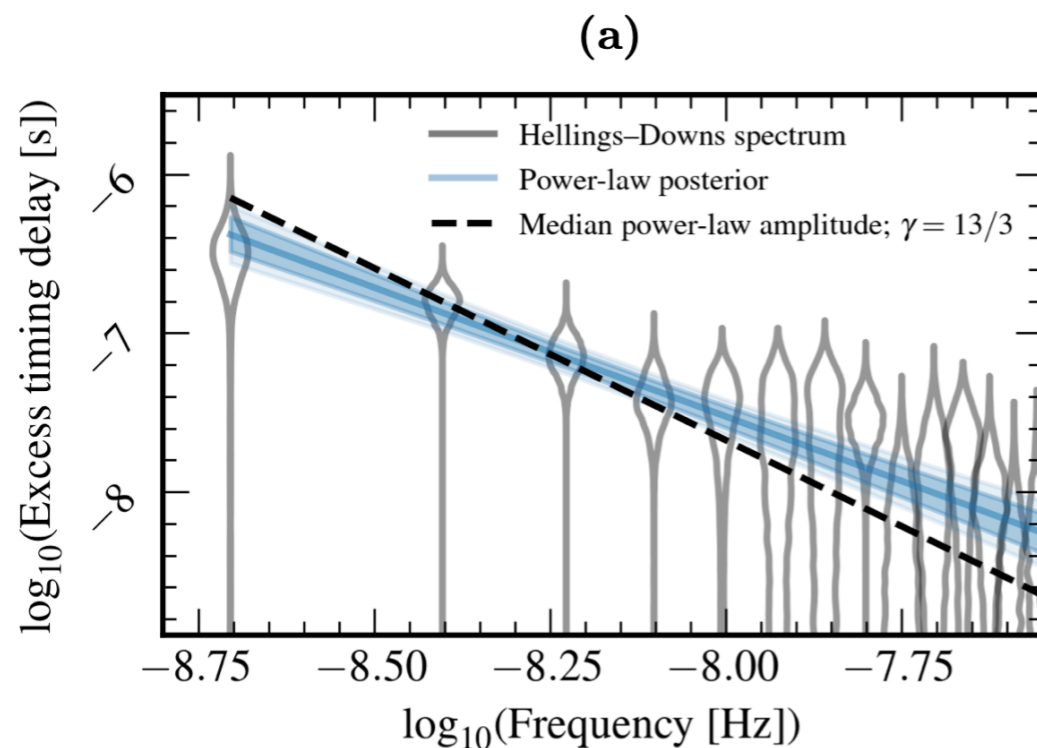
from Breitbach,  
Kopp, Madge,  
Opferkuch, PS  
1811.11175

Some recent  
developments

# Pulsar timing arrays

NANOGrav has observed evidence for a stochastic GW background at nano-Hz frequencies:

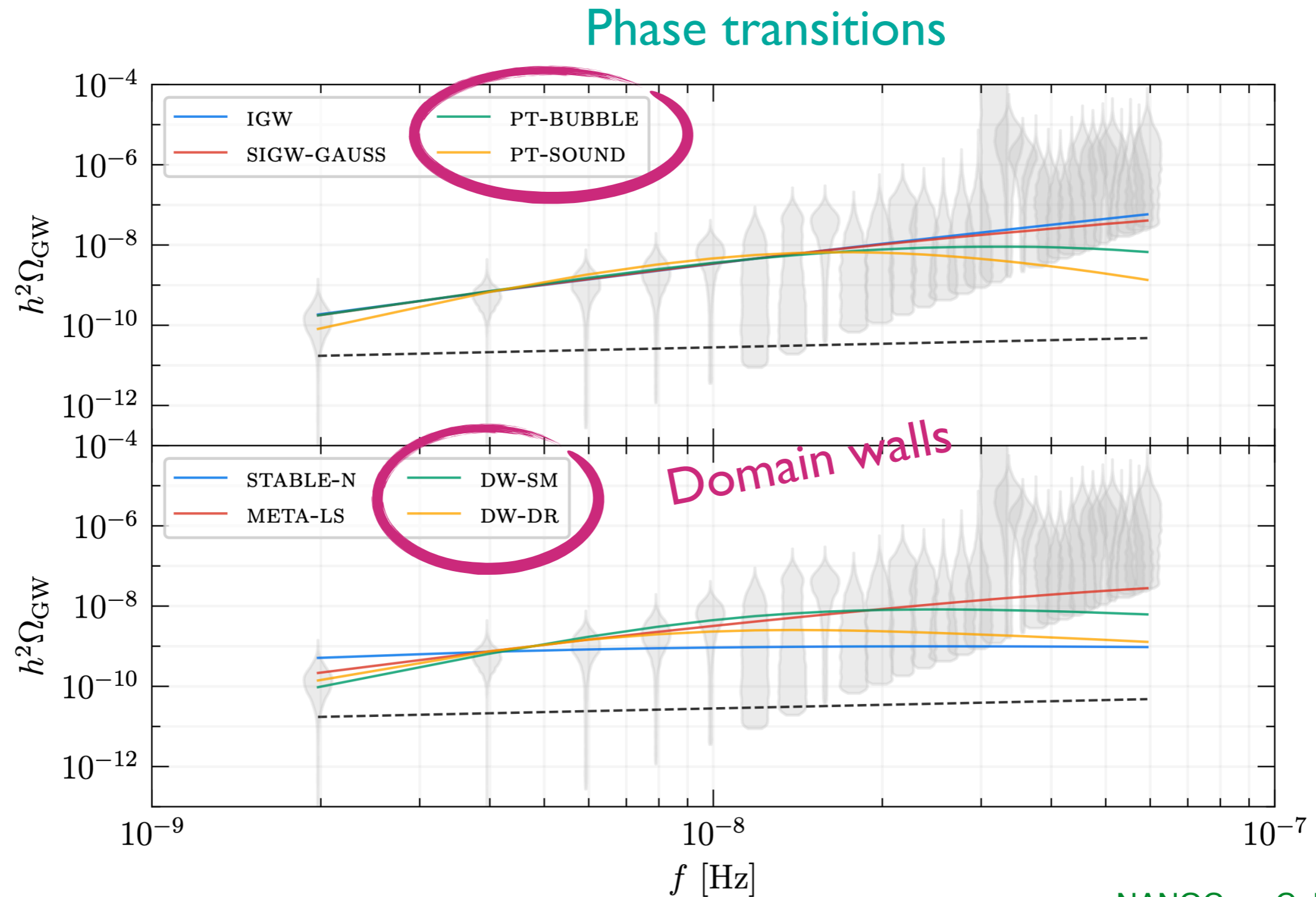
NANOGrav Collaboration,  
2306.16213, APJL 951



Strong evidence for Hellings-Downs correlation

Also supported by new EPTA+InPTA, CPTA data (PPTA less)

# Compatible with primordial GWs from new physics



NANOGrav Collaboration,  
2306.16219, APJL 951

# Thoughts:

This is a very strong signal!

$$\Omega_{\text{GW, today}} \sim 10^{-9}$$

Comparison: The photon density today is  $\Omega_{\gamma} \sim 10^{-5}$ , but photons were in thermal equilibrium in early Universe

Any source that can explain this must:

- ▶ Represent a significant fraction of the total energy density at the time of production,  $T_* \sim (10 - 1000) \text{ MeV}$
- ▶ Be very efficient at converting that energy to GW radiation
- ▶ Then disappear before onset of BBN,  $T \sim 1 \text{ MeV}$

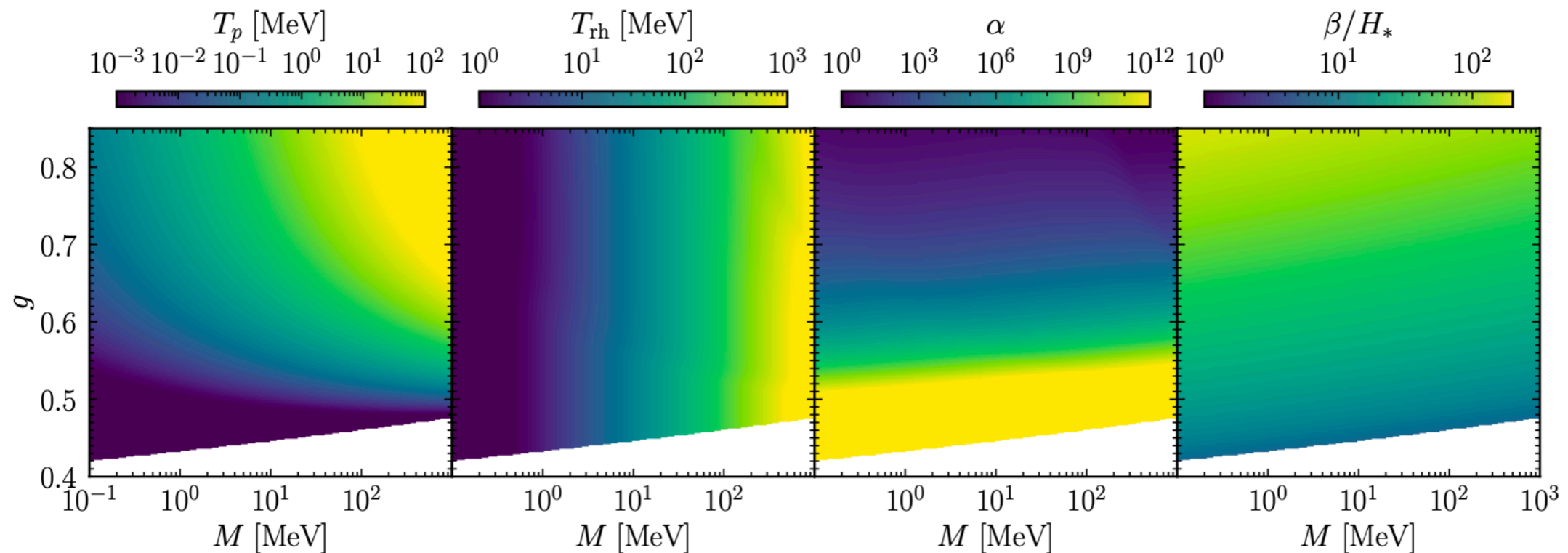


# Supercooled phase transitions

Benchmark model: Coleman-Weinberg model with vanishing tree level potential

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^2 + D_\mu\Phi^\dagger D^\mu\Phi - V(\Phi, T)$$

Two parameter model: Mass scale  $M$  and coupling  $g$



Madge et al,  
[2306.14856](https://arxiv.org/abs/2306.14856)

Signal dominated by colliding bubbles and sound shells

Simulated by Lewicki and Vaskonen, 2208.11697

# Supercooled phase transitions

Madge et al,  
2306.14856

Comparison with  
12 year data

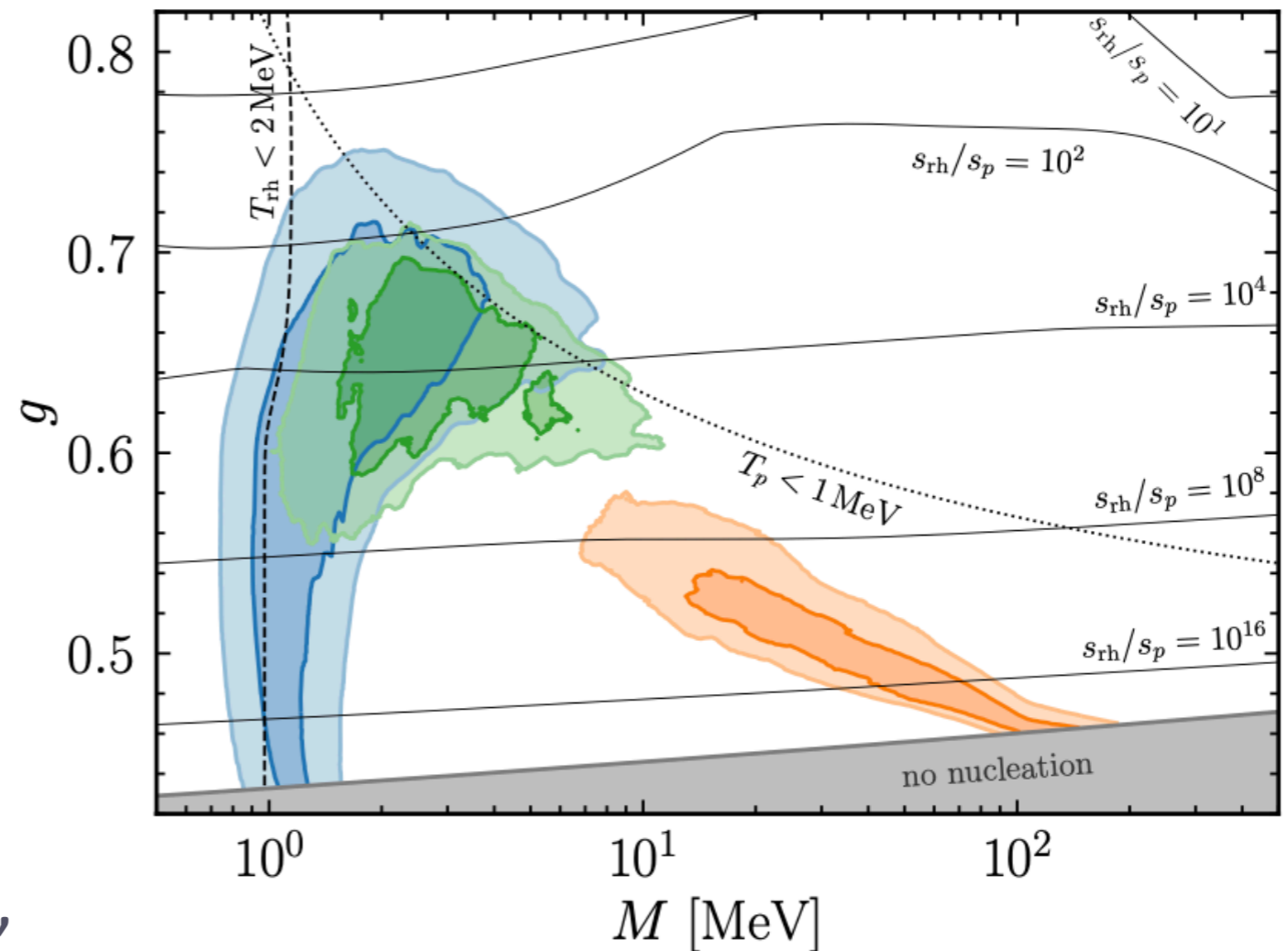
Large supercooling  
and reheating

- ▶ Dilution of baryons,  
dark matter
- ▶ Two BBNs

Pheno: Light scalar  $m_\phi \approx M$ ,  
decay to electrons and photons

Higgs portal not viable, instead

FCC? Or low energy e+e- machine (e.g. MESA in Mainz)



$$\mathcal{L} \supset c_{ee} \frac{|\Phi|^2}{\Lambda^2} LH\bar{e} + c_{\gamma\gamma} \frac{|\Phi|^2}{\Lambda^2} F_{\mu\nu} F^{\mu\nu}$$

Can this be from the  
(dark) QCD phase  
transition?

# Towards quantitative predictions for SU(N) PTs

Strong coupling  $\rightarrow$  Non-perturbative methods required

Lattice?

- ▶ Good for equilibrium thermodynamics (free energy, pressure...)
- ▶ Easier for  $n_f = 0$ , lots of data for  $N_c = 3$
- ▶ No real time dynamics

Holography (AdS/CFT)

- ▶ Allows perturbative calculations
- ▶ Works best for large  $N_c$  and in CFT limit e.g. Hindmarsh et al, Cotrone et al, ...

# Combine both approaches

## Improved holographic QCD

$$\mathcal{S}_5 = -M_P^3 N_c^2 \int d^5x \sqrt{g} \left[ R - \frac{4}{3} (\partial\Phi)^2 + V(\Phi) \right] + 2M_P^3 N_c^2 \int_{\partial M} d^4x \sqrt{h} K$$

- ▶ AdS Einstein-dilaton gravity  $\leftrightarrow$  4D CFT
- ▶ Dilaton potential  $V(\Phi)$
- ▶ Dilaton  $\lambda = \exp \Phi \leftrightarrow$  't Hooft coupling  $\lambda_t = N_c g_{YM}^2$
- ▶ ...
- ▶ Solutions of EOM  $\leftrightarrow$  phases of SU(N)

Gürsoy, Kiritsis, Mazzanti, Nitti  
0707.1324, 0707.1349, 0812.0792, 0903.2859, ...

# Improved holographic QCD

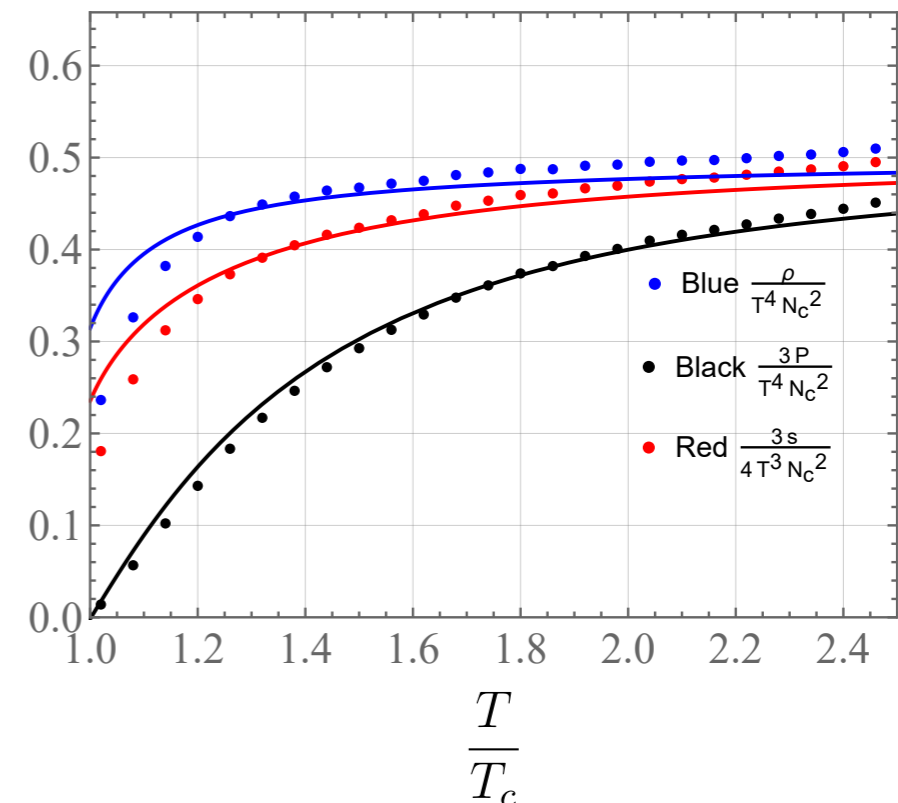
Want this to reproduce SU(N) theories

- ▶ Confinement in IR ( $\lambda \rightarrow \infty$ )
- ▶ Yang Mills beta function in UV ( $\lambda \rightarrow 0$ )

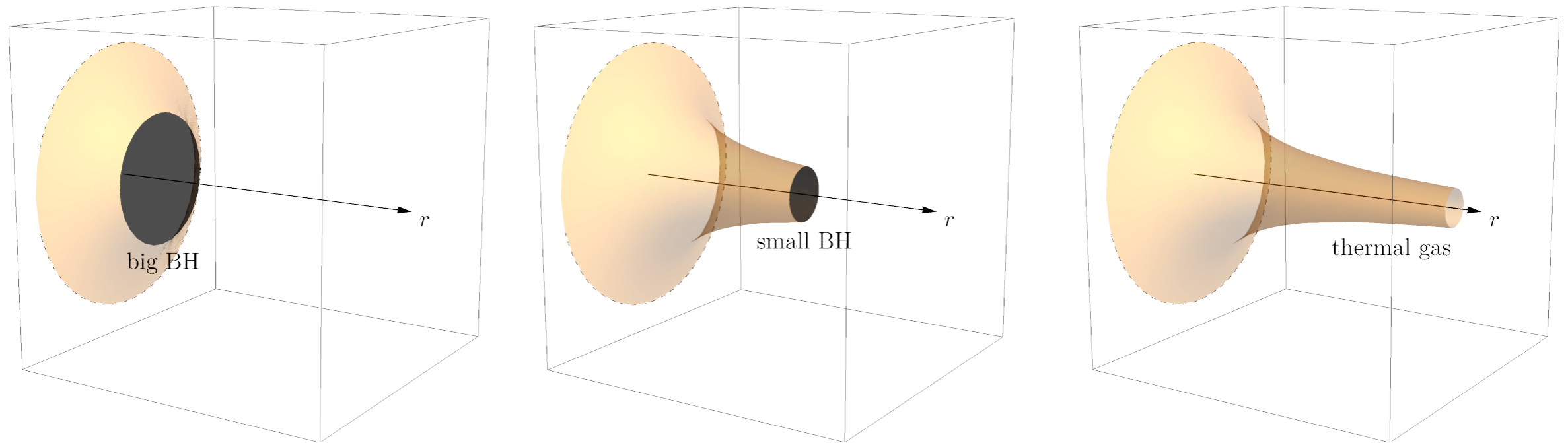
$$V(\lambda) = \frac{12}{\ell^2} \left\{ 1 + V_0 \lambda + V_1 \lambda^{4/3} [\log(1 + V_2 \lambda^{4/3} + V_3 \lambda^2)]^{1/2} \right\}$$

Fix parameters:

- ▶  $V_0, V_2$  to reproduce 2 loop YM running in UV
- ▶  $V_1, V_3$  fit to reproduce SU(3) lattice thermodynamics in IR



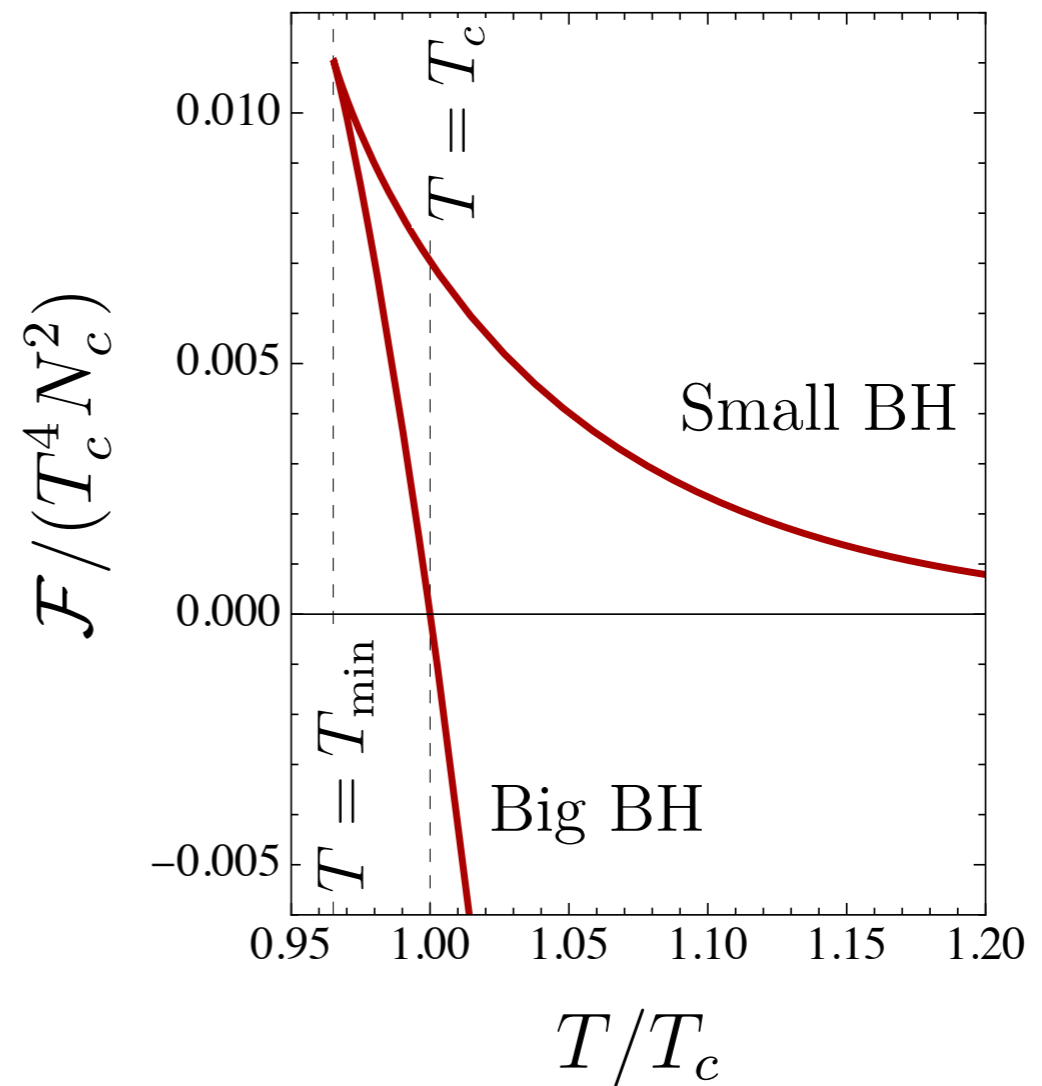
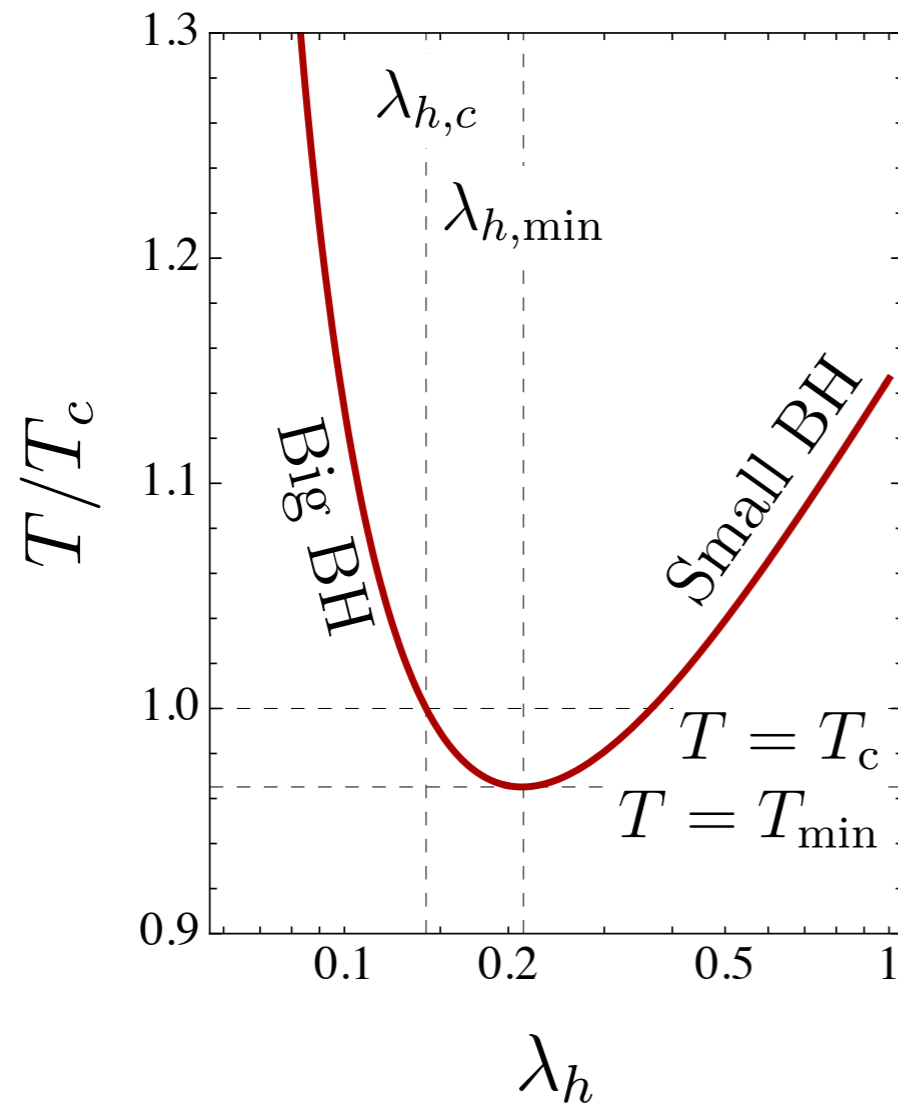
# The phase transition in ihQCD



## Three solutions

- ▶ Big BH: Deconfined phase
- ▶ Small BH: Unstable, saddle point
- ▶ Thermal gas: Confined phase

# The phase transition in ihQCD II



At  $T = T_c$ , deconfined phase becomes meta-stable



# The phase transition in ihQCD III

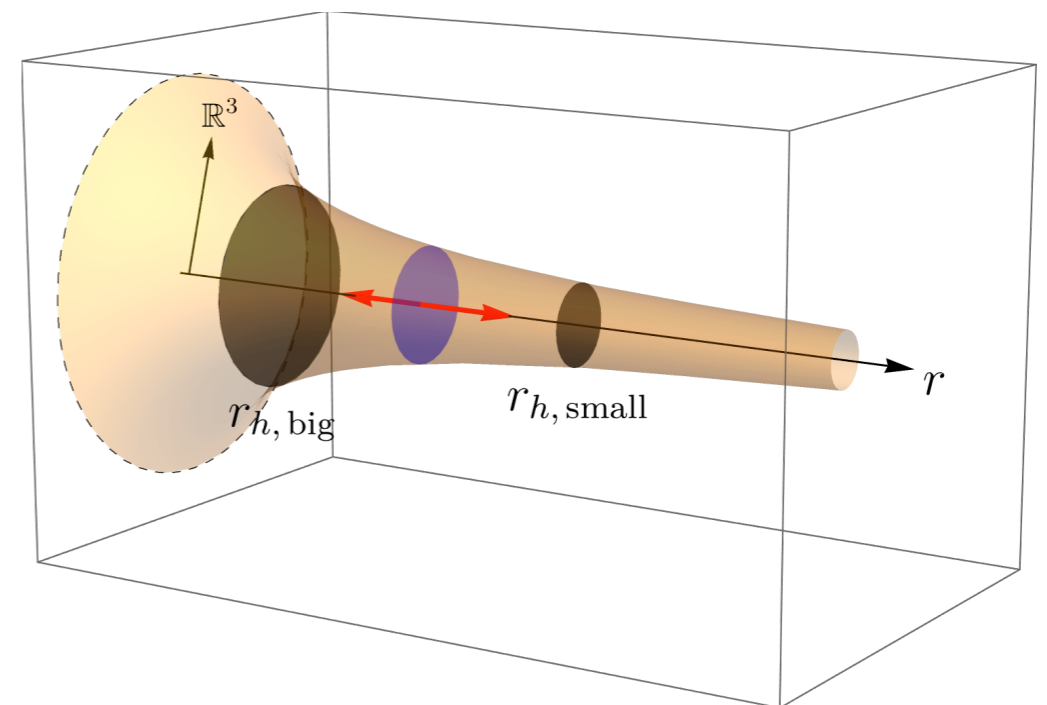
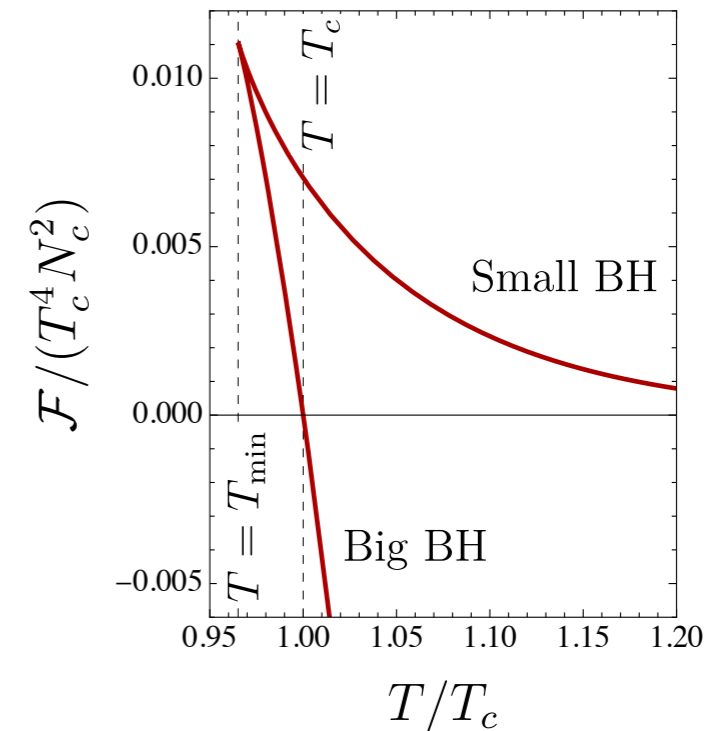
Hawking Page transition, with small BH acting as instanton

To compute bounce action, need effective action (or free energy) along the full path

Interpolate between big and small BH solutions

- ▶ Do some hard work...
- ▶ Win :)

Morgante, Ramberg, PS, 2210.11821



# Effective potential and bounce action

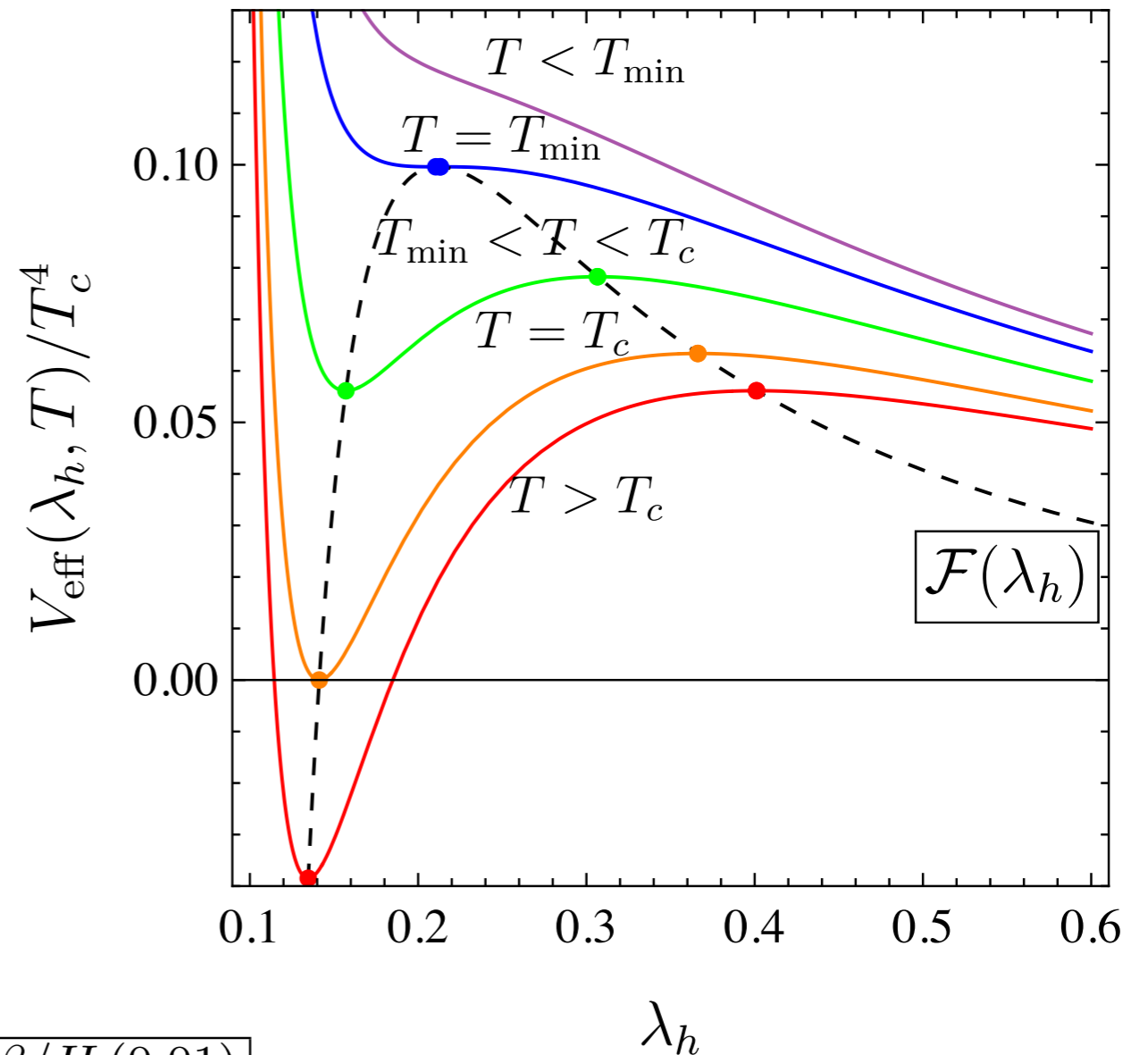
## Bounce action

$$S_{\text{eff}} = \frac{4\pi}{T} \int d\rho \rho^2 \left[ c \frac{N_c^2}{16\pi^2} (\partial_r \lambda_h(r))^2 + V_{\text{eff}}(\lambda_h(r)) \right]$$

## Tunneling decay rate

$$\Gamma = T^4 \left( \frac{S_B}{2\pi} \right)^{3/2} e^{-S_B}$$

Allows us to compute  
 $\alpha$  and  $\beta$



	$\alpha$	$\beta/H (v_w = 1)$	$\beta/H (0.1)$	$\beta/H (0.01)$
$T_c = 50 \text{ MeV}$	0.343	$9.0 \times 10^4$	$8.6 \times 10^4$	$8.2 \times 10^4$
100 GeV	0.343	$6.8 \times 10^4$	$6.4 \times 10^4$	$6.1 \times 10^4$

Morgante, Ramberg, PS, 2210.11821

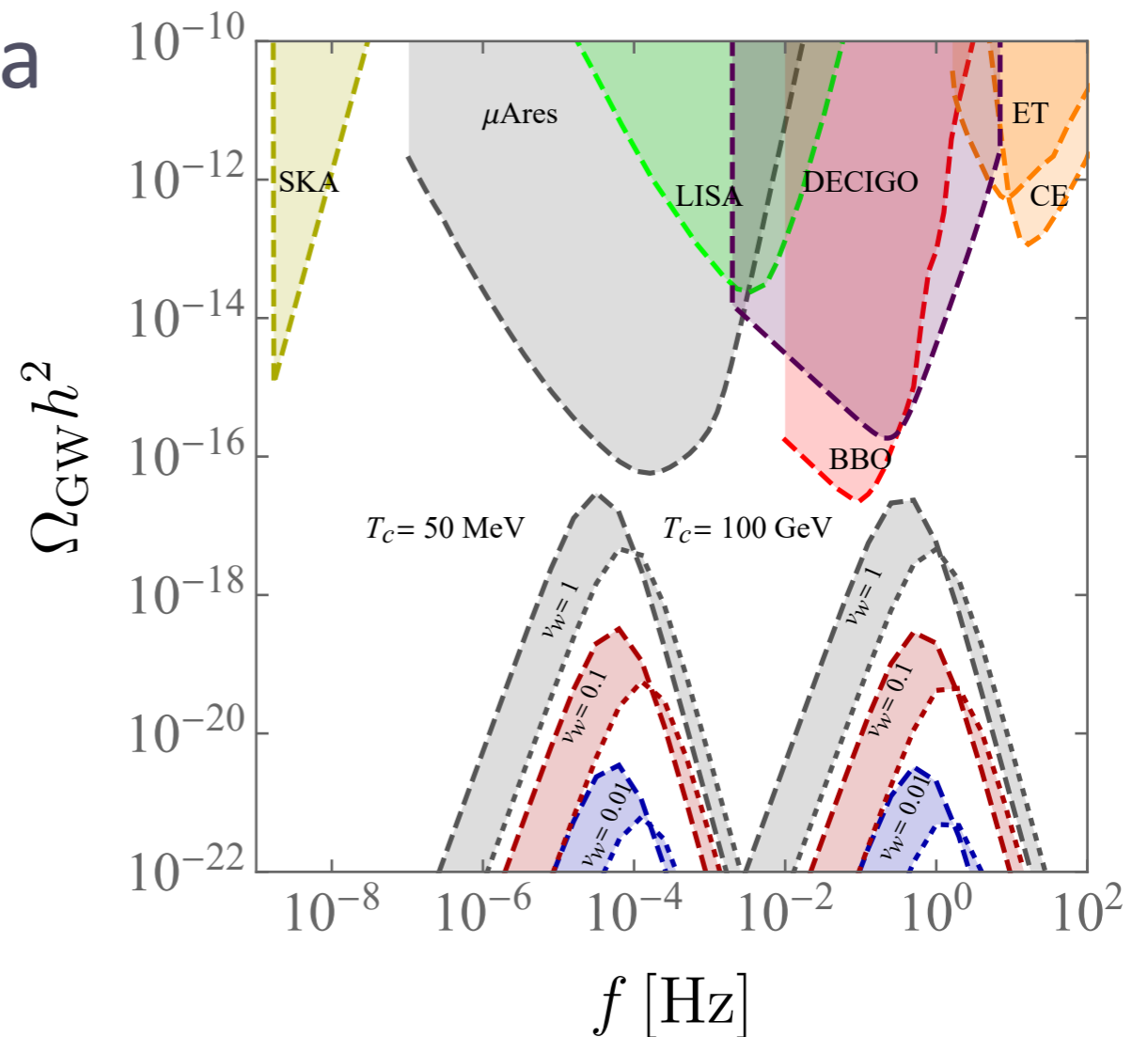
# GW spectrum

First prediction for GW spectra of QCD-like dark sectors from holography

- ▶ for  $N_c = 3, n_f = 0$
- ▶ Some work remains (wall velocity)
- ▶ Larger signal possible for larger  $N_c, n_f$
- ▶ Agrees with estimates based on effective theories and lattice data

(e.g. Halverson+ 2012.04071, Huang+ 2012.11614, March-Russell+ 1505.07109)

Morgante, Ramberg, PS, 2210.11821



# Work in progress

## Bounce action

$$\mathcal{S}_{\text{eff}} = \frac{4\pi}{T} \int d\rho \rho^2 \left[ c \frac{N_c^2}{16\pi^2} (\partial_r \lambda_h(r))^2 + V_{\text{eff}}(\lambda_h(r)) \right]$$

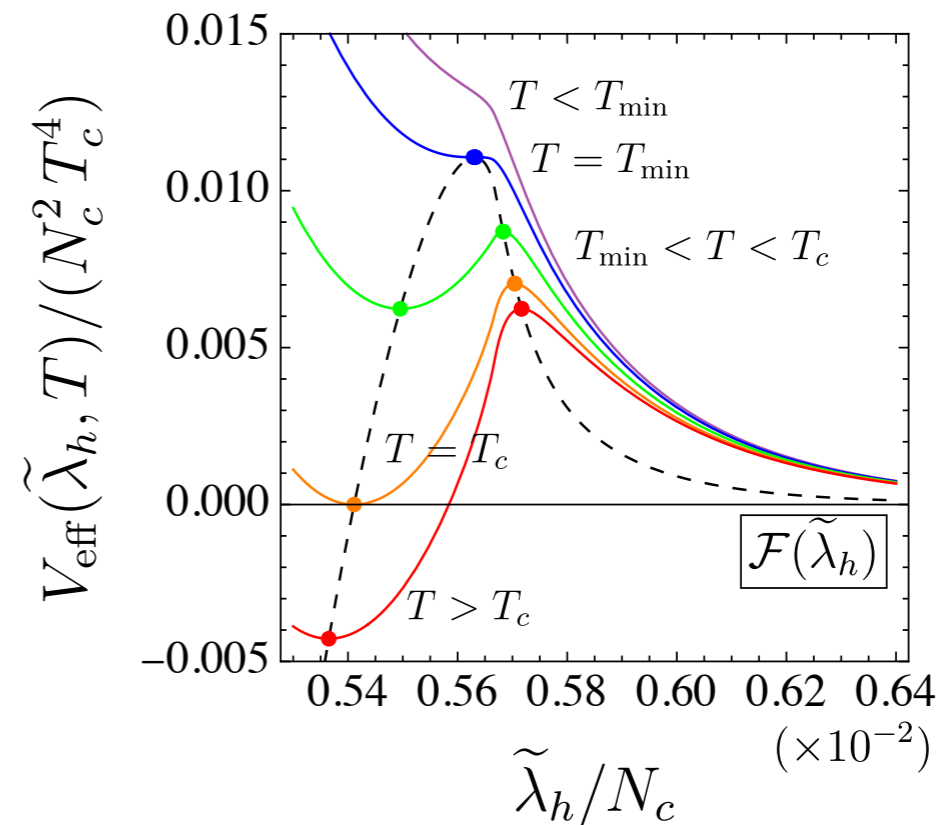
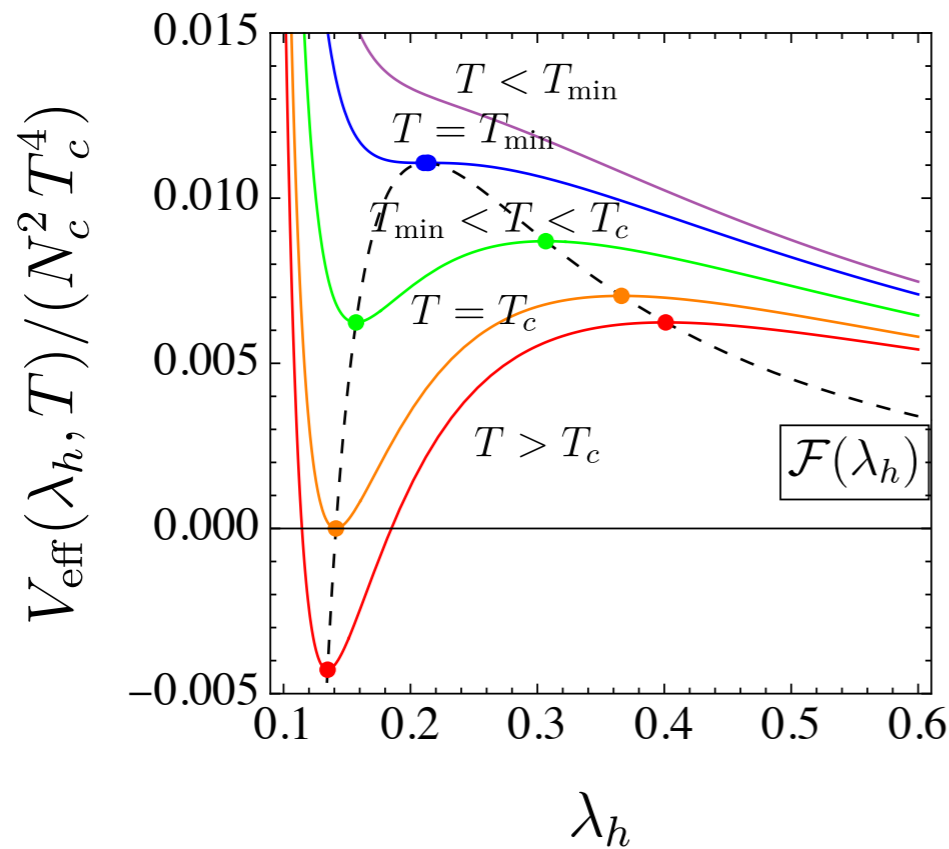
Normalisation “c” of kinetic term unknown

Can be determined for simpler geometries using prescription of [Bigazzi et al, 2008.02579](#)

Here: Not so easy. But obtained an estimate that has proper normalisation up to a factor 1-10

Morgante, Ramberg, PS, in preparation

# Work in progress



Effective action for canonically normalised field

Unfortunately this seems to further weakens the GW signal

Morgante, Ramberg, PS, in preparation

# Work in progress

Furthermore: Redid fit of dilation potential to newer lattice data and for  $N_c$  up to 8 - no significant  $N_c$  dependence

Extrapolation for large  $N_c$  - signal strengthens, but still unobservable

Exploring different methods to estimate wall velocity

Morgante, Ramberg, PS, in preparation

# Summary

Confining dark sectors are interesting

Predict additional phase transition in the early Universe

- ▶ First order for many scenarios, without tuning of parameters

PTA data hints towards a strong first order PT at the MeV scale, potentially in a dark sector

Holography allows computation of PT observables also at strong coupling

In general: GWs are a new window into the early Universe, with lots of data expected in near (PTAs) and far (LISA, ET) future





# Some more details on improved holographic QCD

Thermal gas Ansatz (confined phase):

$$ds^2 = b_0^2(r)(dr^2 - dt^2 + dx^m dx_m).$$

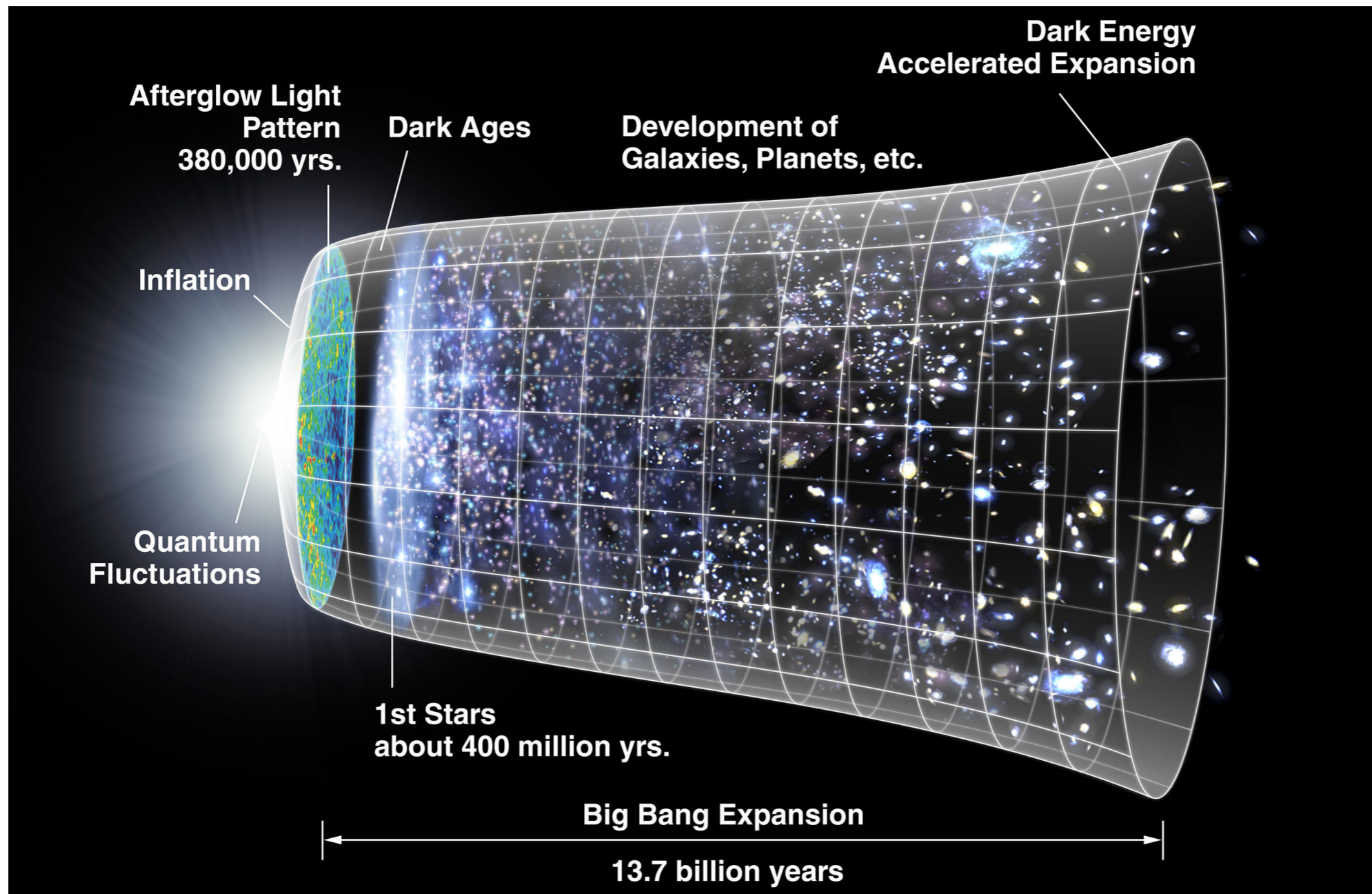
AdS black hole Ansatz (deconfined phase):

$$ds^2 = b^2(r) \left( \frac{dr^2}{f(r)} - f(r)dt^2 + dx^m dx_m \right)$$

Functions  $b_0(r)$ ,  $b(r)$ ,  $f(r)$  determined from EOM

Solutions asymptotic to AdS in UV, with log corrections

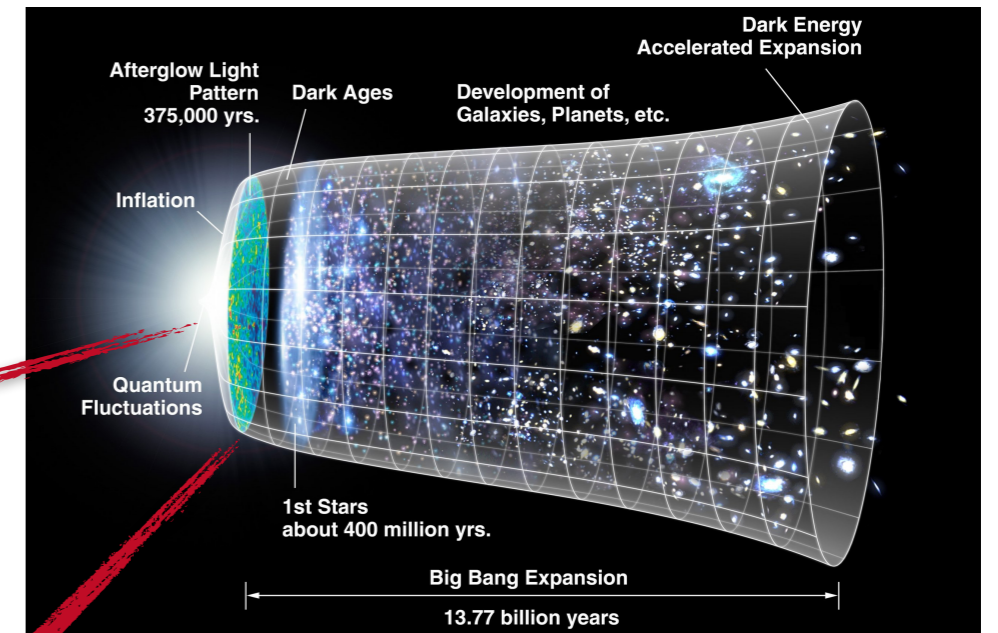
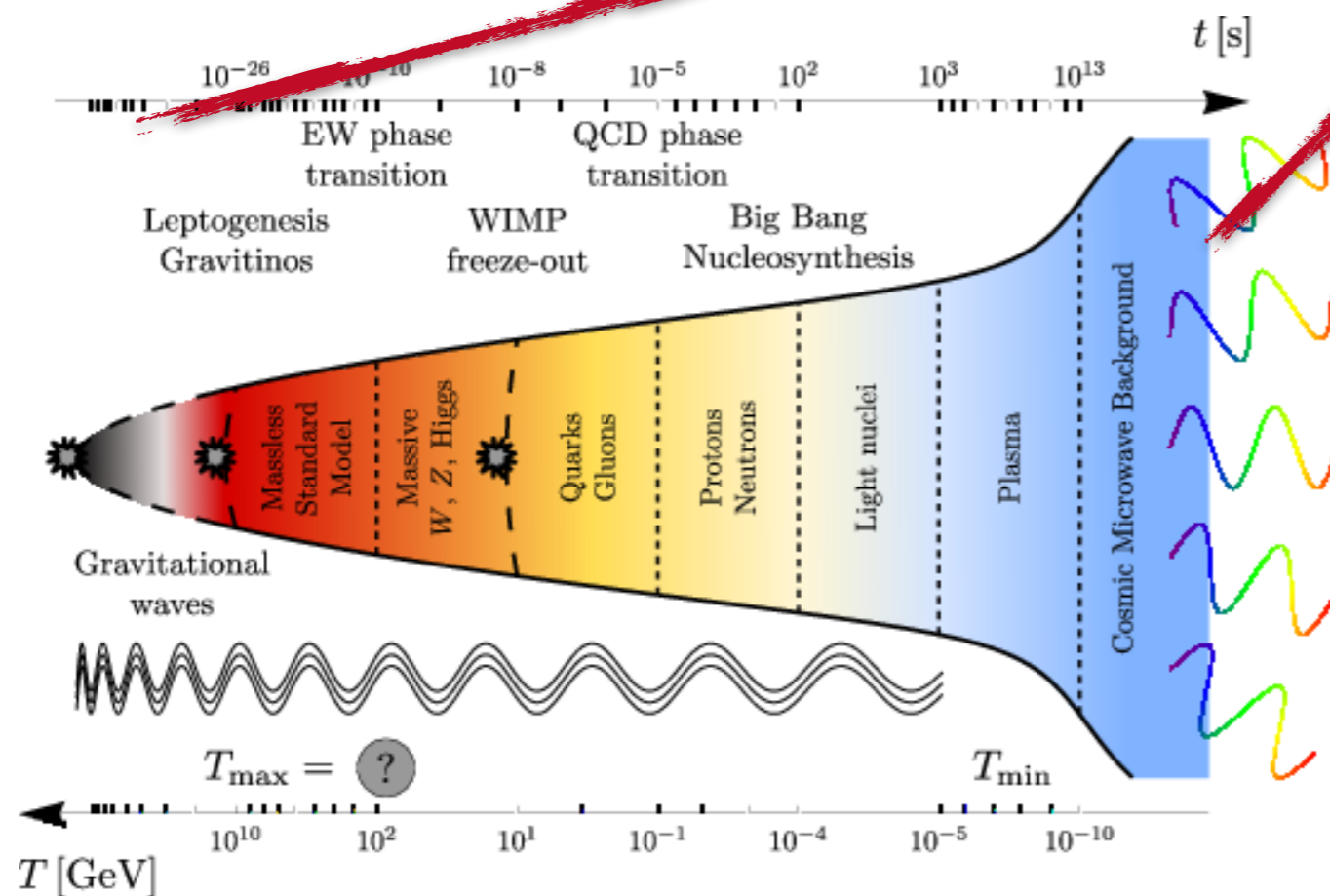
# Thermal History



# Gravitational Waves?

Zoom into interesting region

New window into early universe

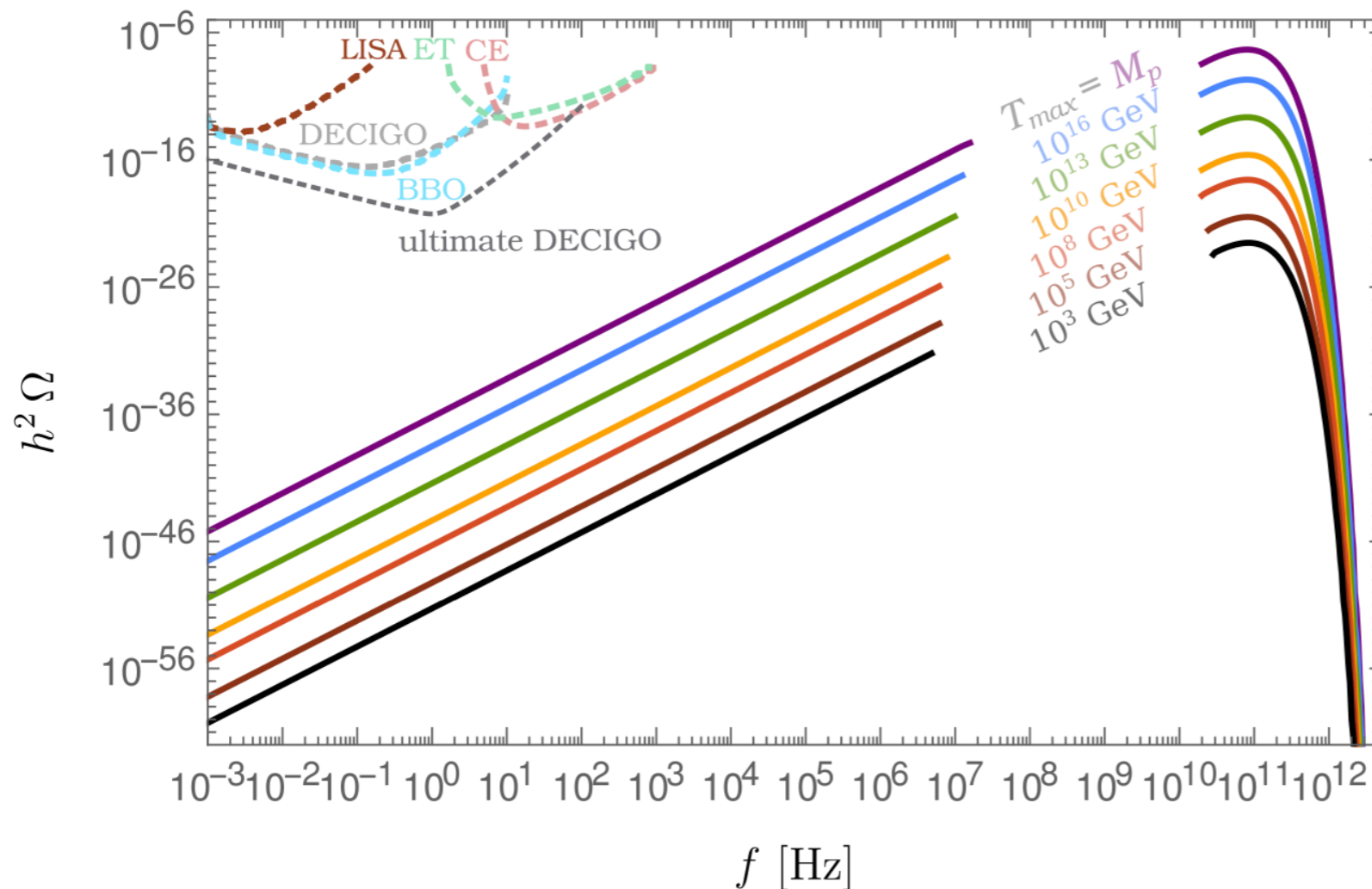


- e.g.
- Electroweak symmetry breaking
  - Baryogenesis
  - Dark matter production

# Standard model

## The hot early Universe sources GWs!

- ▶ Classical picture: thermal fluctuations source tensor fluctuations
- ▶ Quantum picture: gluon + gluon  $\rightarrow$  graviton



From Ringwald,  
Schütte-Engel, Tamarit, 2020

Original computations:  
Ghiglieri, Laine, 2015  
Ghiglieri, Jackson, Laine,  
Zhu, 2020

# Dark QCD

Models I'm interested in here:

Nonabelian  $SU(N)$  dark sector, confinement scale  $\Lambda_d$

$n_f$  light/massless flavours

$$n_f = 0$$

Glueball DM

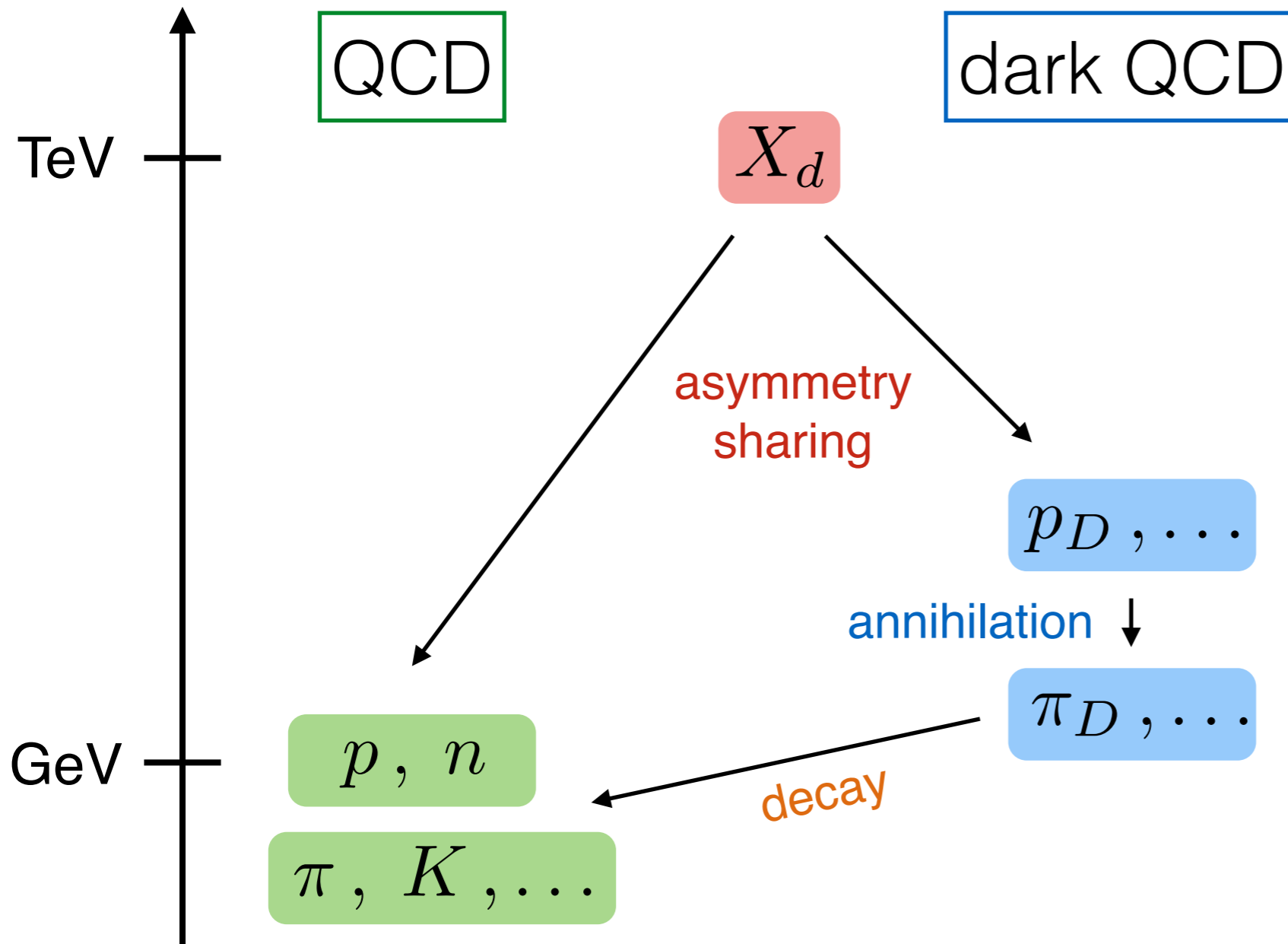
PT from center  
symmetry restoration

$$n_f > 0$$

Dark Baryons  
or Dark Pions

Chiral Symmetry Breaking

# Composite DM / Hidden Sector



- SU(N) dark sector with neutral "dark quarks"
- Confinement scale  $\Lambda_{\text{darkQCD}}$
- DM is composite "dark proton"

Bai, PS, PRD 89, 2014  
PS, Stolarski, Weiler, JHEP 2015

many other works!

Similar setup e.g.: Blennow et al; Cohen et al; Frandsen et al;  
Hidden Valleys: Strassler, Zurek;...

# New physics mass scales for PTA

## Phase transition

►  $T_* = M_{\text{NP}} \sim \text{MeV} - \text{GeV}$

## Audible axions:

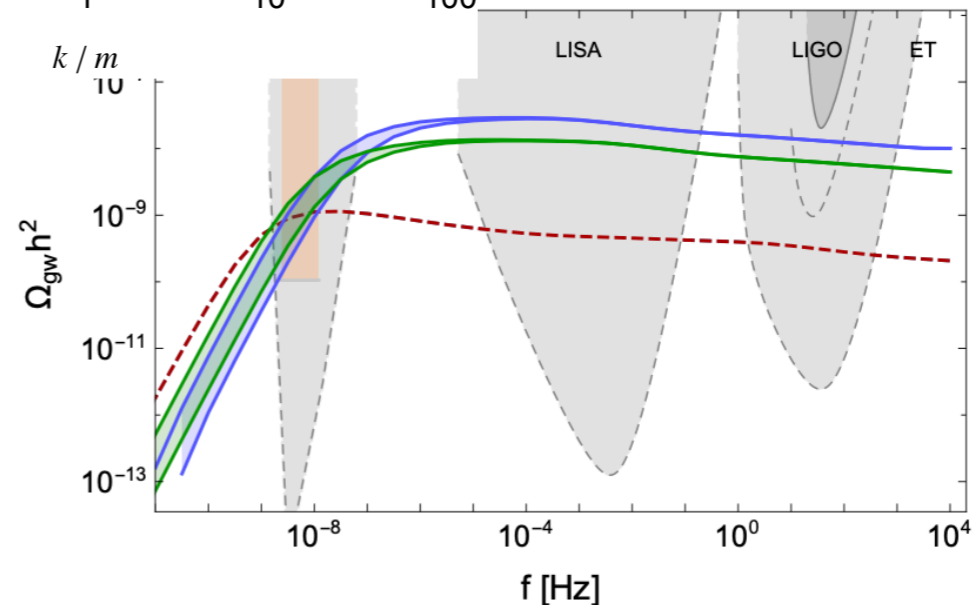
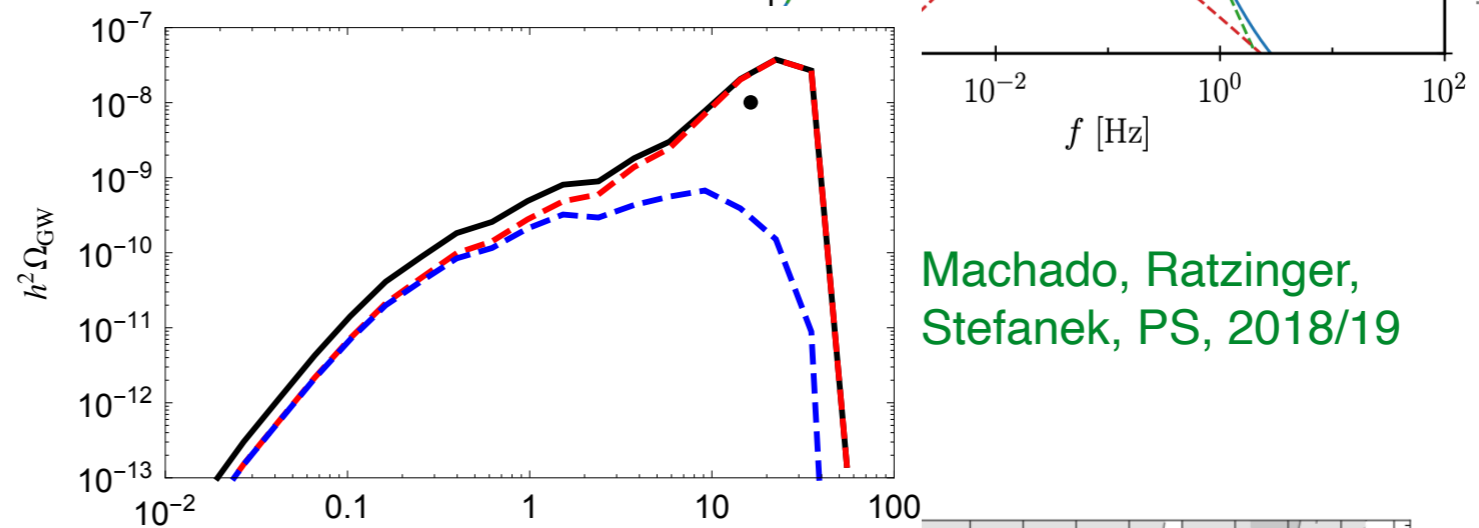
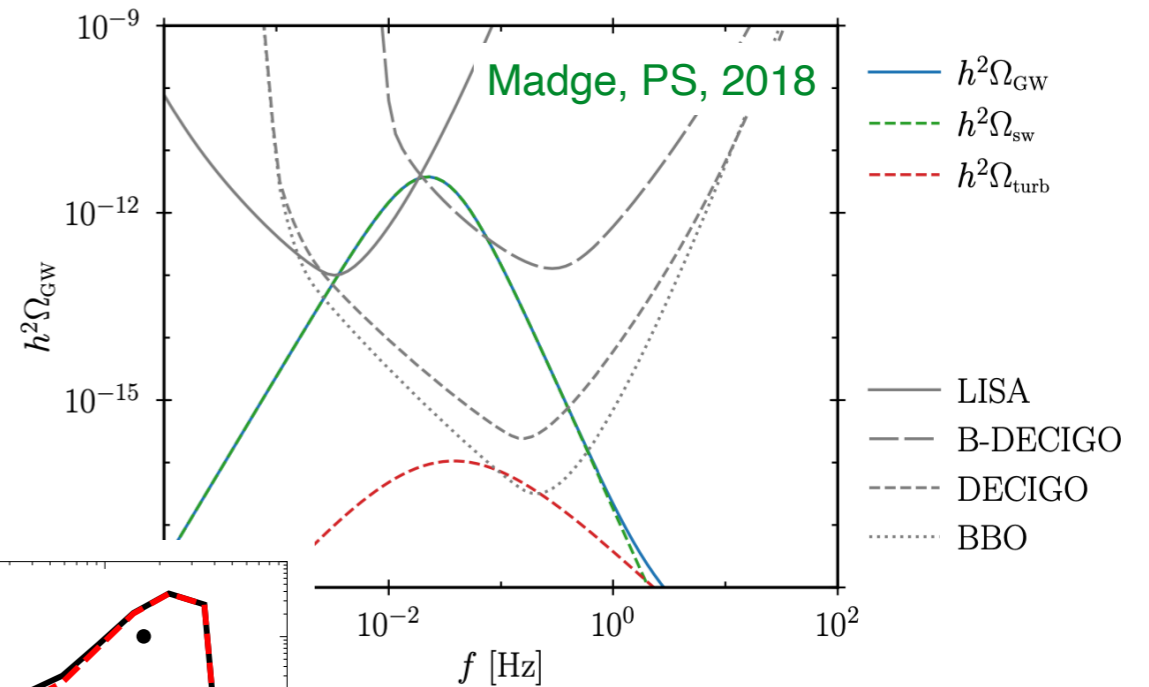
►  $T_*^2 / M_{\text{Pl}} = M_{\text{NP}}$

→  $M_{\text{NP}} \sim 10^{-14} \text{ eV}$

## Cosmic strings/domain walls

►  $T_* \sim \Gamma_{\text{decay}}$

$M_{\text{NP}} \gg \text{MeV}$



# Benchmark models with large signals

Madge et al,  
[2306.14856](#)

## Global cosmic strings

- ▶ In trouble with  $N_{\text{eff}}$

## Audible axions

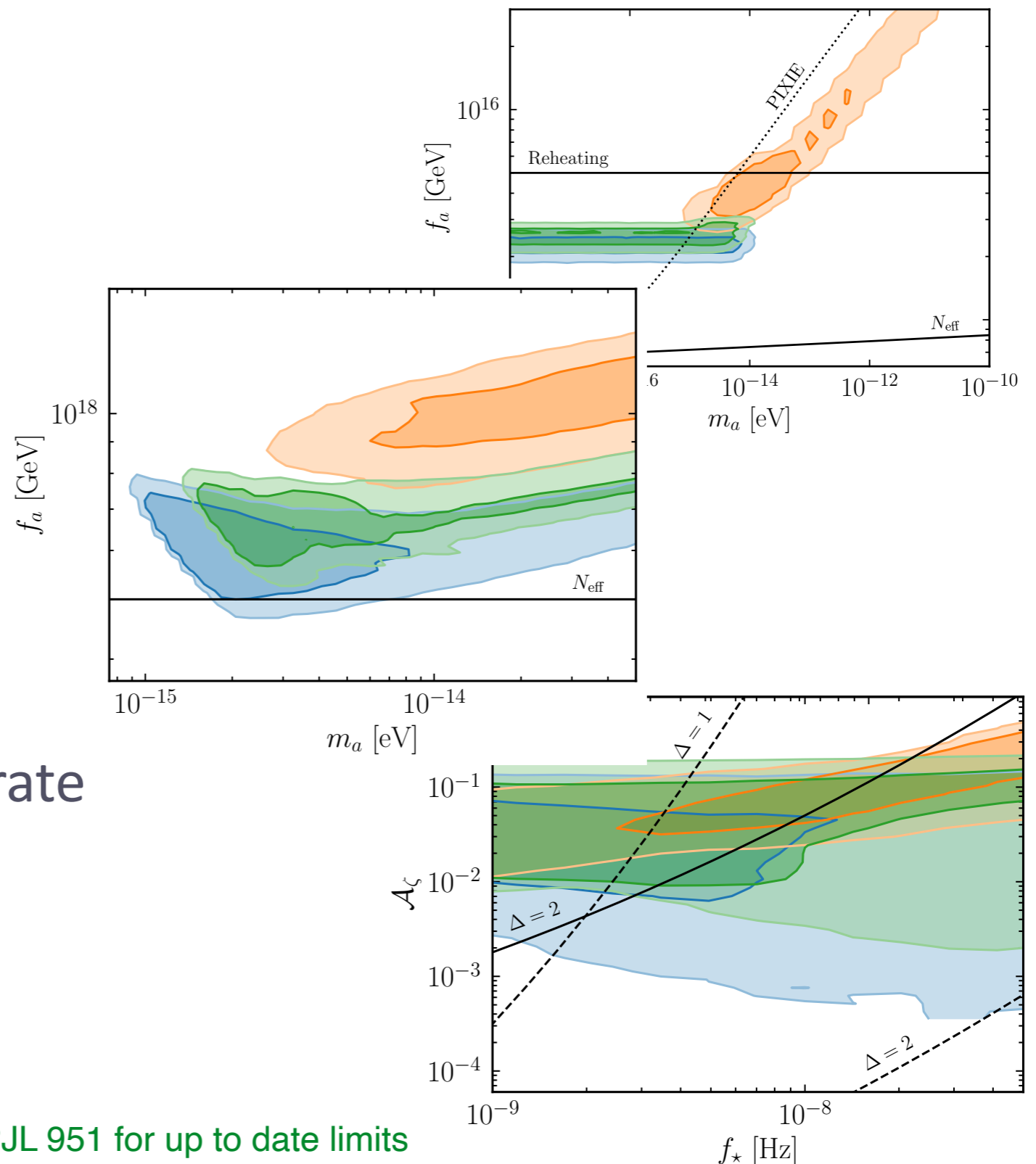
- ▶ Barely consistent with  $N_{\text{eff}}$

## Scalar induced GWs

- ▶ OK, but difficult to generate

## Domain walls

- ▶ Better :)



See also NANOGrav Collaboration, [2306.16219](#), APJL 951 for up to date limits



# Axion/ALP domain walls

Domain walls appear when discrete symmetries are spontaneously broken to degenerate ground states

Long lasting GW source, until DWs annihilate, before dominating the Universe ideally

Review:  
Saikawa,  
[1703.02576](#)

Axion DW:  $U(1)_{\text{PQ}} \rightarrow Z_N$

Surface tension  $\sigma = 8m_a f_a^2$

Annihilation triggered by QCD instantons

$$T_{\text{ann}} \sim 1 \text{ GeV} \left( \frac{g_*(T_{\text{ann}})}{80} \right)^{-\frac{1}{4}} \left( \frac{\Lambda_{\text{QCD}}}{400 \text{ MeV}} \right)^2 \left( \frac{10^7 \text{ GeV}}{f_a} \right) \sqrt{\frac{10 \text{ GeV}}{m_a}}$$

Madge et al,  
[2306.14856](#)

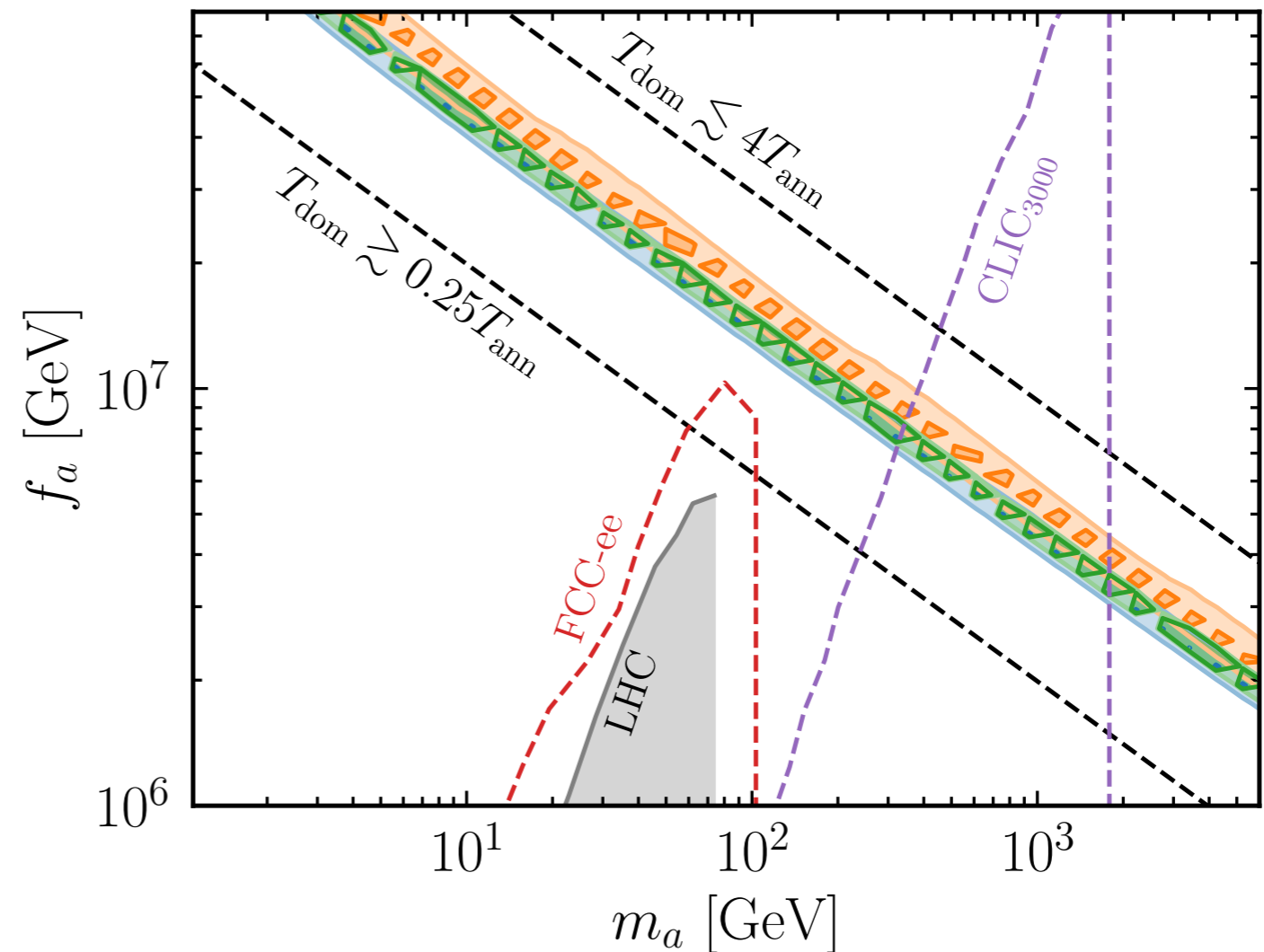
# Axion/ALP domain walls

Madge et al,  
2306.14856

Concrete model:  
Aligned/clockwork  
Axions [Higaki et al, 1606.05552](#)

Heavy axion  
“partners” at weak  
scale

In reach of future  
colliders [Bauer et al, 1808.10323](#)



- ▶ Maybe room for improvement (FCC-hh?)

# Invisibly decaying DWs

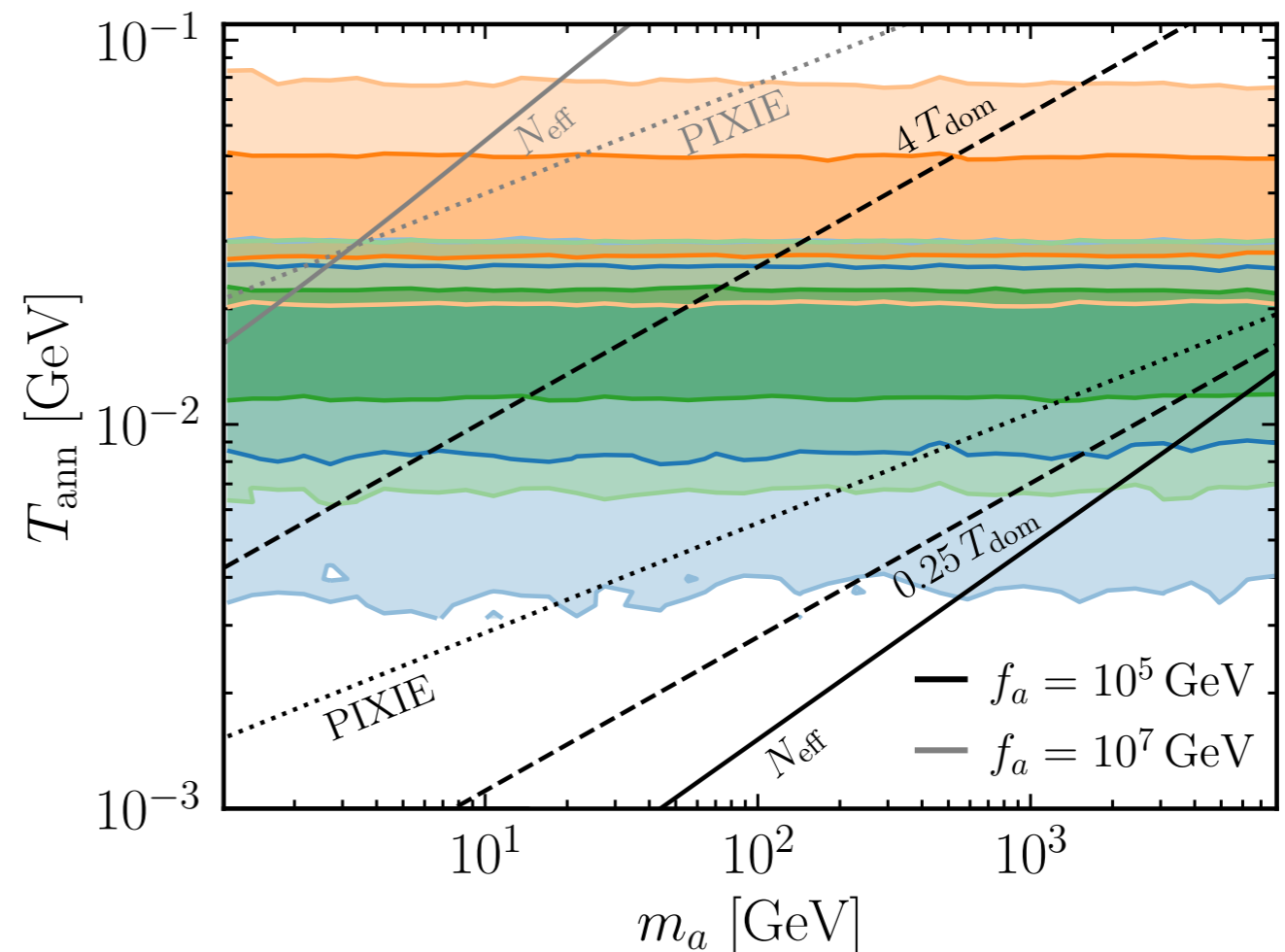
Madge et al,  
[2306.14856](https://arxiv.org/abs/2306.14856)

DWs annihilate to  
dark radiation

- ▶  $N_{\text{eff}}$  ok mostly

Dark sector anisotropies  
induce CMB spectral  
distortions

- ▶ In reach of future  
experiments (PIXIE)



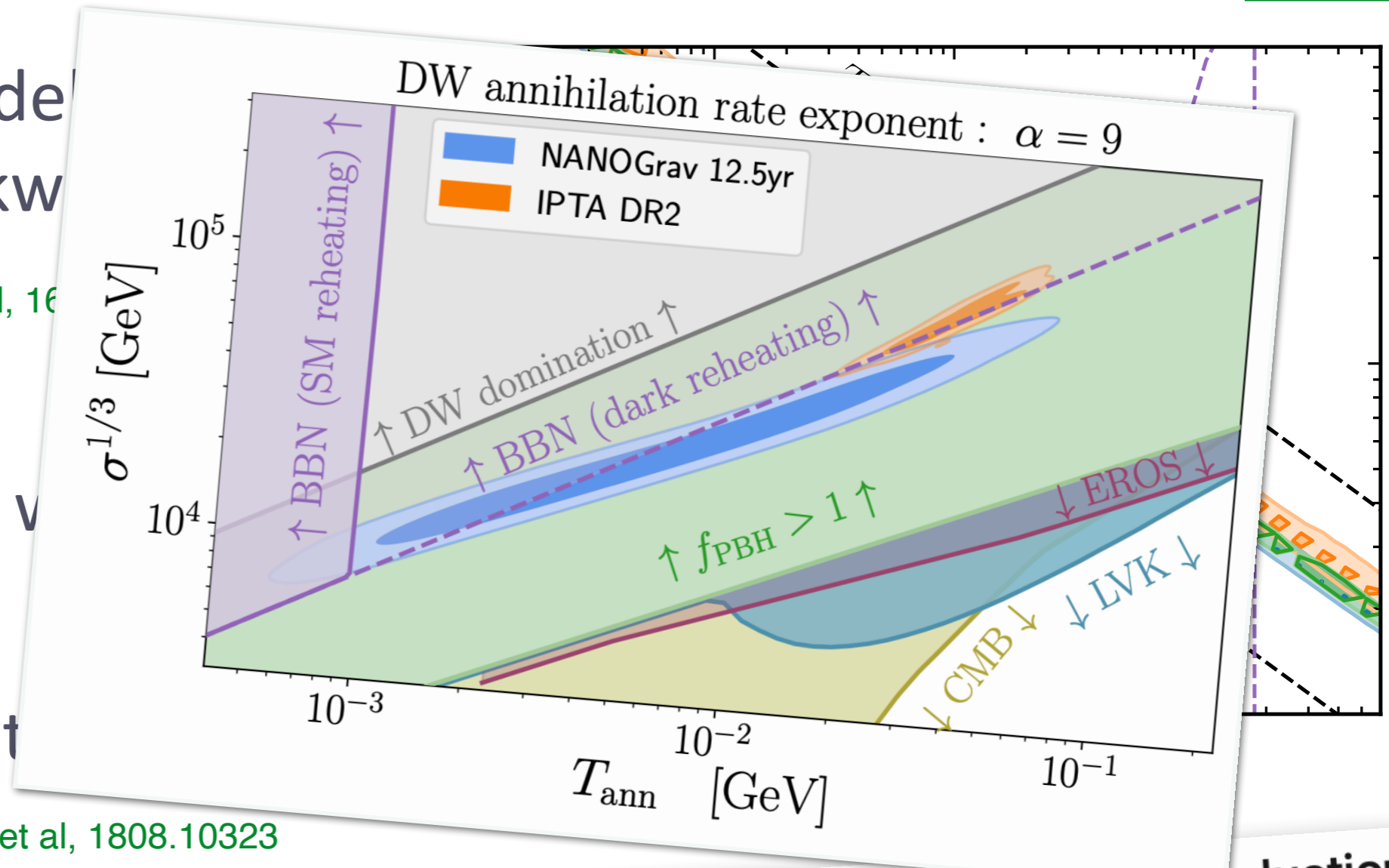
# Axion/ALP domain walls

Madge et al,  
2306.14856

Concrete model  
Aligned/clockwise  
Axions [Higaki et al, 1608.07441](#)

Heavy axion  
“partners” at v  
scale

In reach of future  
colliders [Bauer et al, 1808.10323](#)



► Maybe...  
**Domain wall interpretation of the PTA signal confronting black hole overproduction**

Yann Gouttenoire (Tel Aviv U.), Edoardo Vitagliano (Hebrew U.)

Jun 30, 2023

# One more: Primordial black holes

pBH mergers can explain the data

Clustering needed to evade most stringent bounds

Expect larger anisotropies than from primordial sources

