#### using the cosmological constant to learn about spacetime

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► Work in progress.

#### Outline

► The cosmologícal constant problem.

► The role of the gravitational symmetries.

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Weyl transverse gravity.

The cosmological constant problem

Failure of the decoupling principle behind effective field theory.

Some quantities are extremely sensitive to high-energy physics.

Relevant, but non-natural operator:

$$\Lambda \sqrt{-g}$$

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 Value of the cosmological constant: both ultravioled and infrared aspects; quantum gravity.

## Our solar system

Natural theoretical prediction:

$$|\Lambda_{\rm vac}|\sim$$
 10 $^8~{
m GeV}^4$ 

Measured perihelion precession of Mercury:

 $\Delta arphi =$  574.10  $\pm$  0.65 arc-seconds per century



## The quantum vacuum

Scattering amplitudes and vacuum bubbles:



Vacuum bubbles couple to a dynamical volume form

$$\boldsymbol{\epsilon} = \sqrt{|g|} \, dx^{1} \wedge \dots \wedge dx^{n}$$

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# Consistent low-energy effective description

No modifications of the classical physics.

- Respect classical solutions of general relativity and the particle spectrum of the standard model.
- Different semiclassical and quantum physics in the presence of gravitational fields.
- Cosmological constant as mysterious as (but no more than) any other parameter: gravitational constant, electron charge, ...

#### Fírst steps

Shift symmetry

$$\mathcal{L} \longrightarrow \mathcal{L} + c_o$$

▶ Classically: Fixed volume element  $\omega$ ; corresponding contribution

Radiative corrections: we need a symmetry that forbids the term

$$\int {m \epsilon} \, \Lambda = \int {
m d}^4 { imes} \sqrt{- { ilde { extsf{9}}}} \, \Lambda$$

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Scale transformations of the gravitational field:

$$g_{ab} \longrightarrow \zeta^2 g_{ab} \qquad \zeta \in \mathbb{R}$$

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## The gravitational action

 The Einstein-Hilbert action is not invariant under scale transformations.

 Second-order field theory: longitudinal diffeomorphisms should be broken.

 Keeping the counting of degrees of freedom: local scale transformations

$$g_{ab} \longrightarrow \zeta^2(x) g_{ab}$$

Dynamical conformal structures, parametrized by tensor densities

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Weyl transverse gravity

$$\mathcal{A} := \frac{1}{2\kappa} \int_{\mathcal{M}} \boldsymbol{\omega} \, \mathbb{R}[|\boldsymbol{\omega}|^{1/n} |\boldsymbol{g}|^{-1/n} \boldsymbol{g}_{ab}]$$

- Theory of dynamical conformal structures on n-dimensional manifold M.
- Invariant under transverse diffeomorphisms and Weyl transformations:

$$\delta_{\xi,\varphi}g_{ab} = \mathcal{L}_{\xi}g_{ab} + \varphi g_{ab} \qquad \nabla_a \xi^a = 0$$

► Dynamical volume element √-g forbidden by symmetries.

Matter couples to the composite field

 $|\omega|^{1/n}|g|^{-1/n}g_{ab}$ 

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## From gravity to gravitons

• Expand around flat spacetime  $\eta_{ab}$ :

$$g_{ab} = \eta_{ab} + \lambda h_{ab}$$

At the lowest order: invariant under

$${\mathfrak h}'^{ab}={\mathfrak h}^{ab}+\eta^{ac}\partial_c\xi^b+\eta^{bc}\partial_c\xi^a+\phi\eta^{ab}\qquad\qquad\partial_a\xi^a=o$$

- On-shell equivalence to Fierz-Pauli theory (Izawa1995, Alvarez2005).
- Higher orders describe the interaction vertices of gravitons.
- Nonlinear theory of a spin-2 particle.

## classical theory

• In the gauge  $g = \omega$ , one recovers the traceless Einstein equations

$$R_{ab} - \frac{1}{4}Rg_{ab} = \kappa \left( T_{ab} - \frac{1}{4}Tg_{ab} \right)$$

- These equations are equivalent to Einstein field equations in the same gauge (Ellis2010).
- $\blacktriangleright$  The cosmological constant  $\Lambda$  is a constant of integration.
- Shift symmetry  $\mathcal{L} \to \mathcal{L} + C_0$ :

$$T_{ab} \longrightarrow T_{ab} + g_{ab}C_o, \qquad \Lambda \longrightarrow \Lambda - \frac{\pi}{a}C_o$$

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#### Anomalíes

 Generic result: not all symmetries can be preserved in the quantization. Path integral:

 $\int [\mathcal{D}\Psi] \exp(iS[\Psi])$ 

 A symmetry is not anomalous per se, but with respect to other symmetries.

 A necessary condition is that different symmetries act on the same degrees of freedom.

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Conformal anomaly

Díffeomorphisms:

$$\delta \sqrt{|g|} \propto \nabla_a \xi^a$$

Conformal transformations:

 $\delta \sqrt{|g|} \propto \Omega^4$ 



# Avoiding the anomalies



# Path integral for a scalar field

Inner product:

$$\langle \phi, \phi' 
angle := \int_{\mathcal{M}} d^{\mathbb{D}} \times \sqrt{|\hat{g}|} \, \phi(x) \phi'(x) = \int_{\mathcal{M}} \omega \, \phi \, \phi'$$

Decomposition coefficients:

$$c_{n}:=\langle \phi_{n},\phi
angle =\int_{\mathcal{M}}\omega\,\phi_{n}\,\phi$$

Path integral measure:

$$\prod_{n=0}^{\infty} \frac{dc_n}{\sqrt{2\pi}}$$

Absence of anomalies:

$$\delta c_n = \int_{\mathcal{M}} d^{\mathbb{D}} \mathbf{x} \, \boldsymbol{\phi}_n(\mathbf{x}) \boldsymbol{\phi}(\mathbf{x}) \, \delta \sqrt{|\hat{g}|} = 0$$

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#### Semiclassical theory

 Classical gravitational fields, quantum matter fields: effective action

$$\mathcal{S}_{ ext{gab}} = rac{1}{2}$$
 ln det $(\mathcal{O}_{ ext{gab}})$ 

Heat kernel expansion of the effective action:

$$S_{g} = S_{go} - \frac{1}{32\pi^{2}} \int_{\mathcal{M}} \boldsymbol{\omega} \left\{ \mu^{2} [a_{1}(\hat{g}_{ab}) - a_{1}(\hat{g}_{ab}^{o})] + \iota \boldsymbol{\omega} (\mu^{2}/\boldsymbol{\omega}^{2}) [a_{2}(\hat{g}_{ab}) - a_{2}(\hat{g}_{ab}^{o})] \right\}$$

No term corresponding to a<sub>o</sub>, which in general relativity leads to the renormalization of the cosmological constant.

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• Due to the non-dynamical volume form  $\omega$ .

## Renormalization group

 Renormalization of gravitational couplings; e.g., gravitational constant

$$\frac{1}{\kappa} = \frac{1}{\kappa_0} + C_1 \mu^2 + C_2 \log\left(\frac{\mu}{C_3}\right)$$

There is NO renormalization equation for the cosmological constant.

 Cosmological constant protected by gravitational scale transformations:

$$g_{ab} 
ightarrow \zeta^2 g_{ab}$$

▶ Shift symmetry on the Lagrangian  $\mathcal{L} \rightarrow \mathcal{L} + C_0$ ;  $C_0$  drops off from field equations.

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#### Quantum theory

- Quantum gravitational field.
- Formal definition of the path integral.
- Correlation functions:

$$\langle \Phi(\mathbf{x}_{1})\Phi(\mathbf{x}_{2})...\Phi(\mathbf{x}_{n})\rangle = \frac{\int [\mathcal{D}g][\mathcal{D}\psi]\exp(i\mathbf{S})\Phi(\mathbf{x}_{1})\Phi(\mathbf{x}_{2})...\Phi(\mathbf{x}_{n})}{\int [\mathcal{D}g][\mathcal{D}\psi]\exp(i\mathbf{S})}$$

 Contributions that in general relativity would renormalize the cosmological constant are cancelled out:

$$S_{o} := c_{o} \int_{\mathcal{M}} \boldsymbol{\omega}$$

#### Díscussion

- Metric volume form  $\epsilon$  versus non-dynamical volume form  $\omega$ .
- Essentially equivalent classical field equations. What does it tell us?
- But different semiclassical properties.
- A non-dynamical volume form ω seems less elegant, but provides a better fit to the properties we observe.

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#### Conclusions

 Theory of dynamical causal structures that uses an auxiliary, non-dynamical volume form.

Standard classical solutions for the gravitational field equations.

 Effective description that avoids the (ultraviolet) cosmological constant problem: non-anomalous gravitational scale invariance.

 Being an effective description, it invites looking for completions. New suggestions for quantum gravity?



Thank you for your attention.

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