

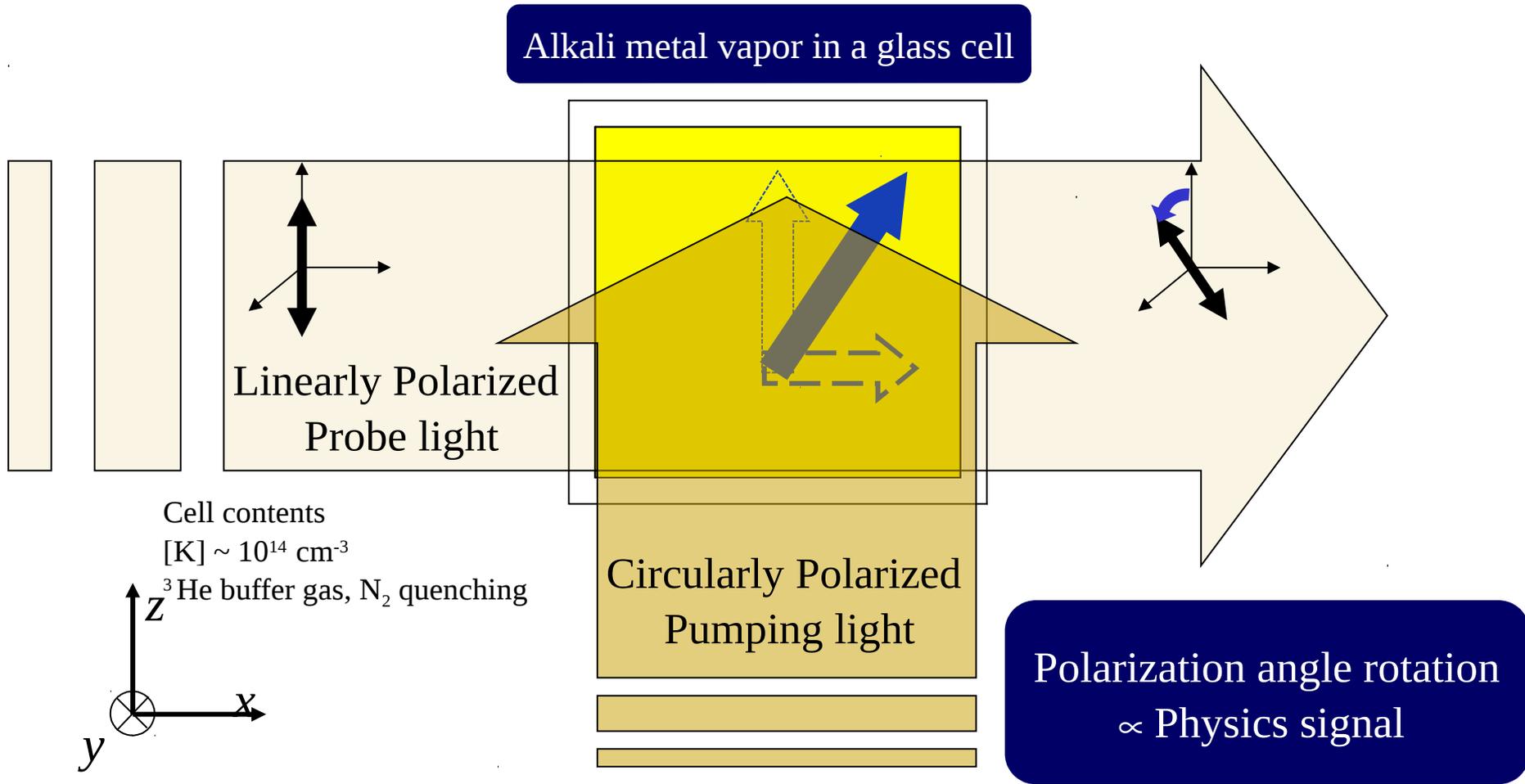
Precision measurements with  
atomic co-magnetometer at the  
South Pole

**Michael Romalis**  
**Princeton University**

# Outline

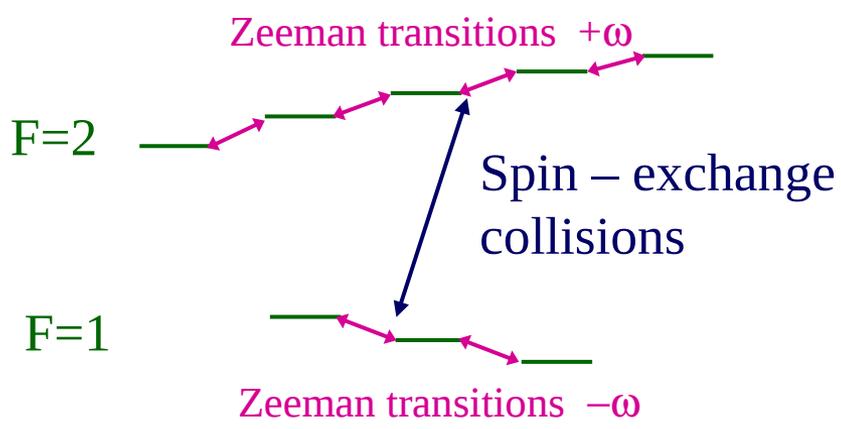
- Alkali metal - noble gas co-magnetometer
- Rotating co-magnetometer at the South Pole
- New Physics Constraints
  - ⇒ Lorentz violation
  - ⇒ Long-range spin-dependent forces
  - ⇒ Slowly oscillating fields
- Current experiments
  - ⇒ Search for spin-mass interaction on 20 cm scale
  - ⇒ Search for spin-spin interactions

# Operation of Atomic Co-Magnetometer

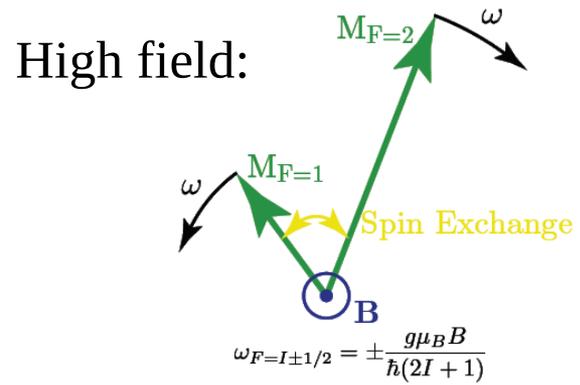
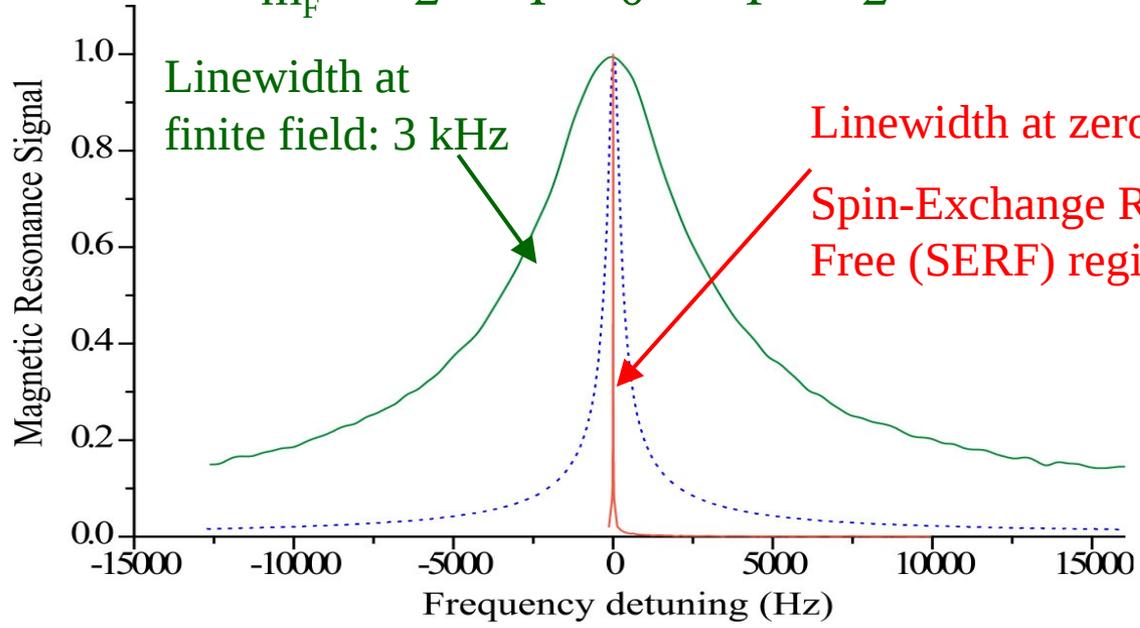


# Elimination of spin-exchange broadening at zero field

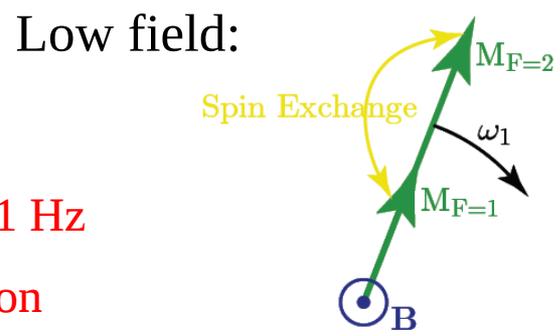
Ground state Zeeman and hyperfine levels in K



$m_F = -2 \quad -1 \quad 0 \quad 1 \quad 2$



$$\omega_{F=I \pm 1/2} = \pm \frac{g\mu_B B}{\hbar(2I+1)}$$



$$\omega_1 = \frac{3(2I+1)}{3+4I(I+1)}\omega = \frac{2}{3}\omega_{F=I \pm 1/2}$$

W. Happer and H. Tang, PRL 31, 273 (1973); J. C. Allred, R. N. Lyman, T. W. Kornack, and MVR, PRL. 89, 130801 (2002)

# K-<sup>3</sup>He Co-magnetometer

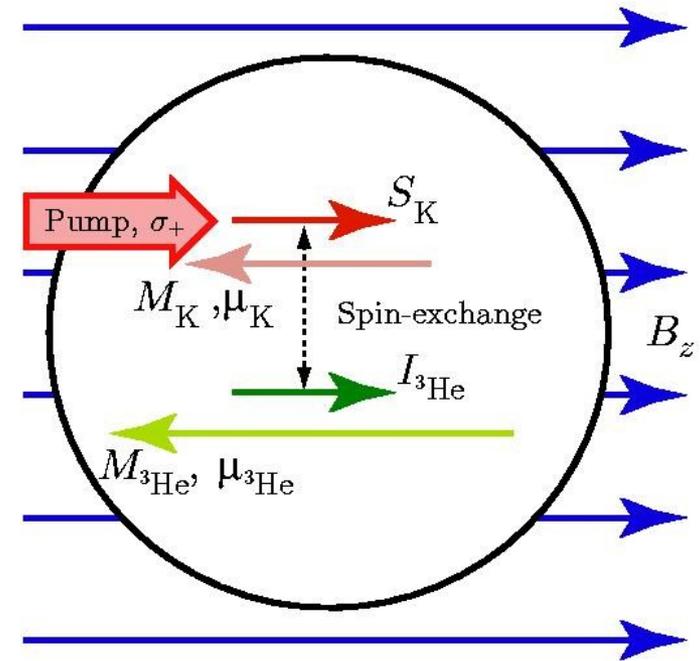
1. Optically pump potassium atoms at high density ( $10^{13}$ - $10^{14}$ /cm<sup>3</sup>)
2. <sup>3</sup>He nuclear spins are polarized by spin-exchange collisions with K vapor
3. Polarized <sup>3</sup>He creates a magnetic field felt by K atoms

$$B_K = \frac{8\pi}{3} \kappa_0 M_{\text{He}}$$

4. Apply external magnetic field  $B_z$  to cancel field  $B_K$   
⇒ K magnetometer operates near zero magnetic field

5. At zero field and high alkali density K-K spin-exchange relaxation is suppressed
6. Obtain high sensitivity of K to magnetic fields in spin-exchange relaxation free (SERF) regime

*Turn most-sensitive atomic magnetometer into a co-magnetometer*



J. C. Allred, R. N. Lyman, T. W. Kornack, and MVR, PRL **89**, 130801 (2002)

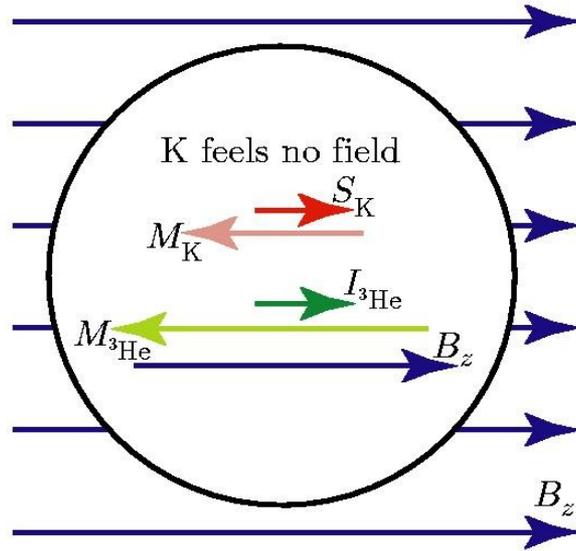
I. K. Kominis, T. W. Kornack, J. C. Allred and MVR, Nature **422**, 596 (2003)

T.W. Kornack and MVR, PRL **89**, 253002 (2002)

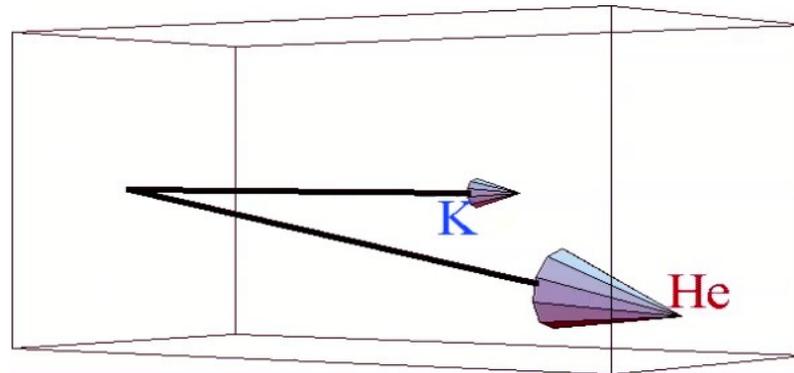
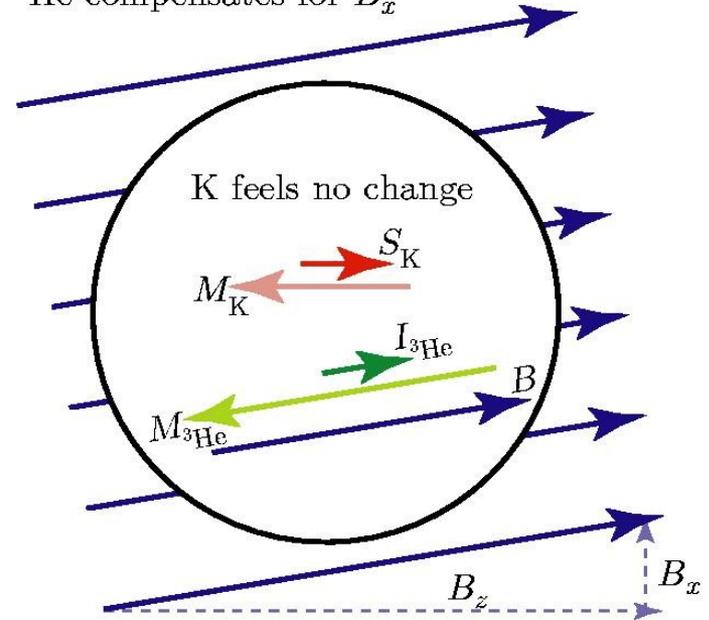
T. W. Kornack, R. K. Ghosh and MVR, PRL **95**, 230801 (2005)

# Magnetic field self-compensation

(a)  $^3\text{He}$  cancels the external field  $B_z$



(b)  $^3\text{He}$  compensates for  $B_x$



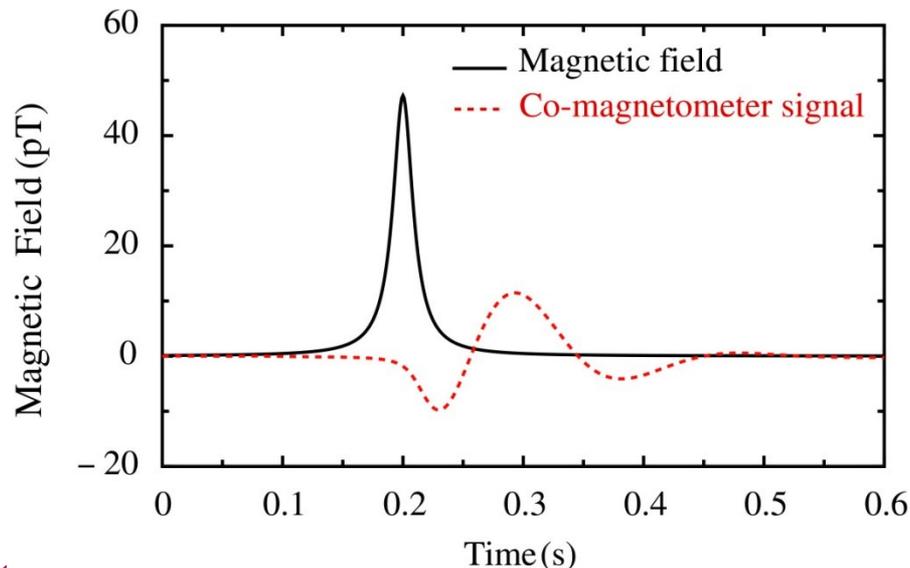
# Response to transient signals

- Fast transient response

- ⇒  $^3\text{He}$  has  $T_2$  of 1000s of seconds

- ⇒ Transient signals decay in 0.3 seconds

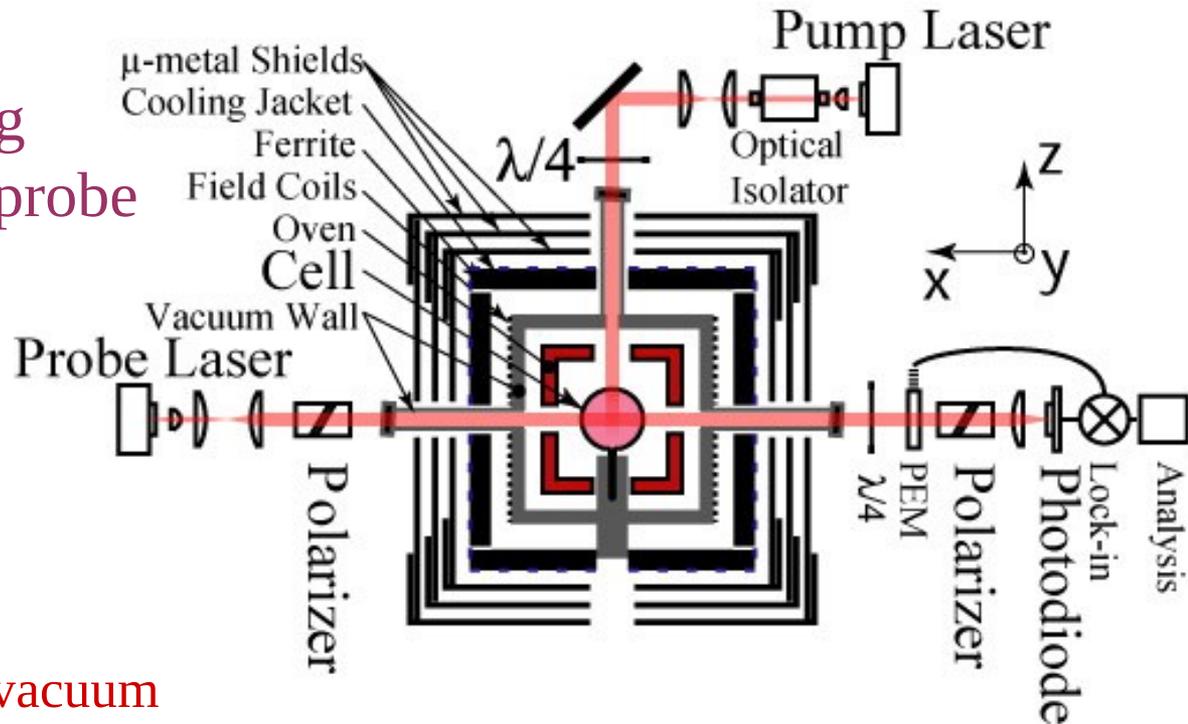
- ⇒ Due to spin-damping coupling to K atoms



- Integral of the signal is proportional to spin rotation angle for arbitrary pulse shape

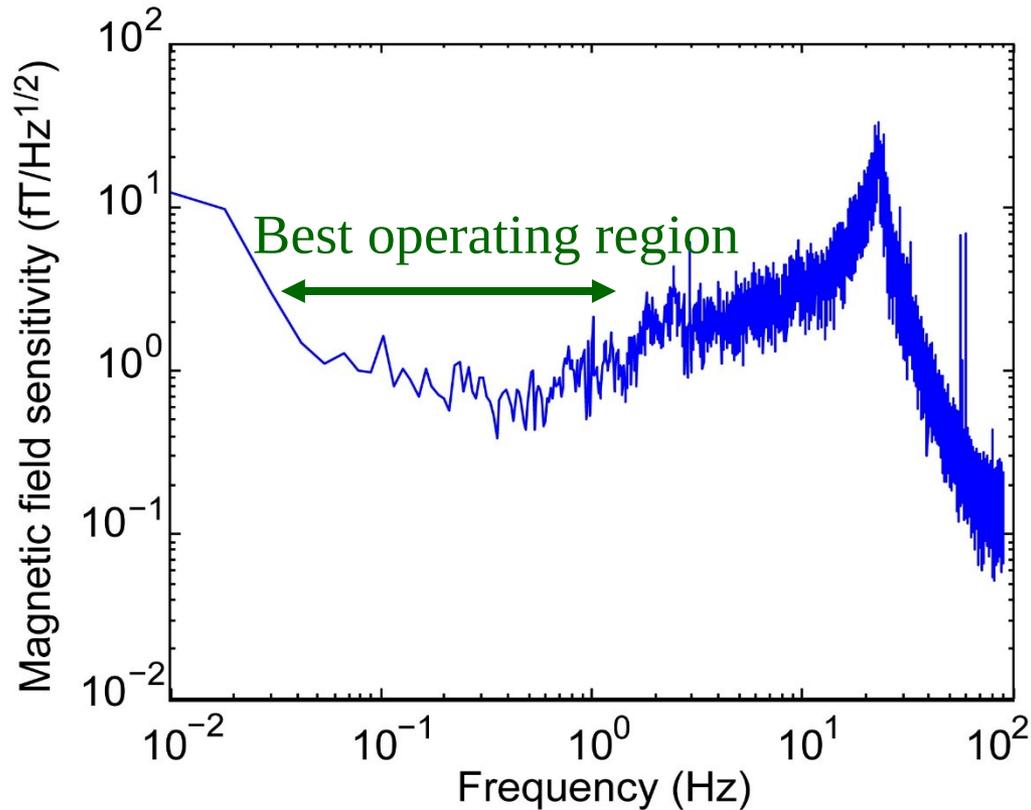
# Co-magnetometer Setup

- Simple pump-probe arrangement
- Measure Faraday rotation of far-detuned probe beam
- Sensitive to spin coupling orthogonal to pump and probe



- Details:
  - ⇒ Ferrite inner-most shield
  - ⇒ 3 layers of  $\mu$ -metal
  - ⇒ Cell and beams in mtorr vacuum
  - ⇒ Polarization modulation of probe beam for polarimetry at  $10^{-7}\text{rad/Hz}^{1/2}$
  - ⇒ Whole apparatus in vacuum at 1 Torr

# Magnetic field sensitivity



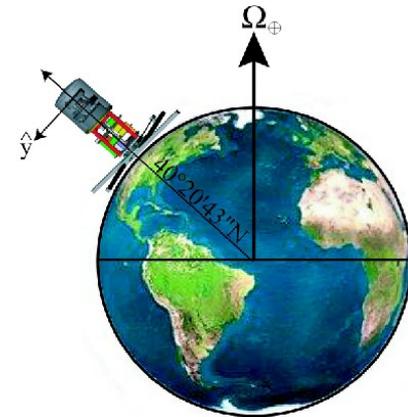
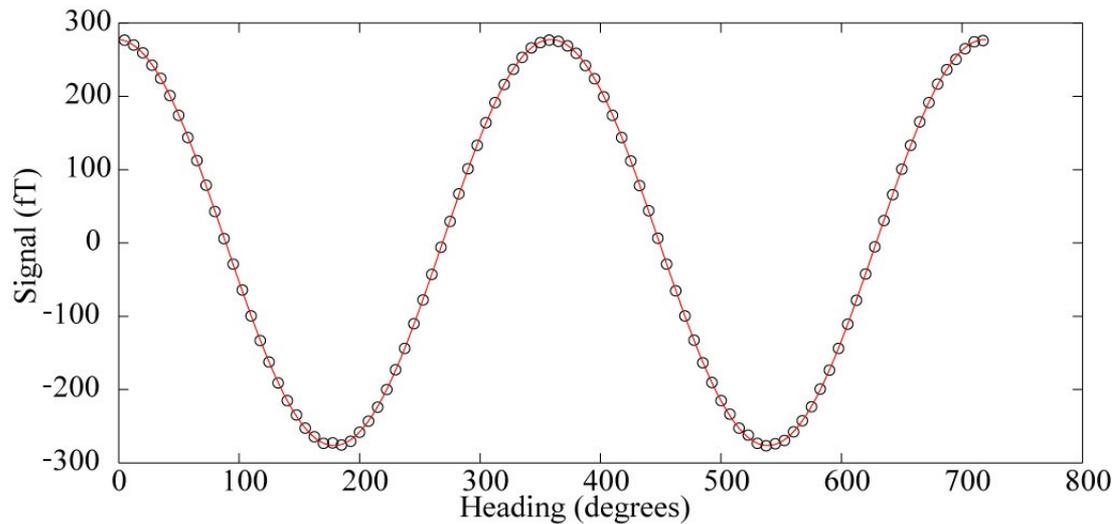
- Sensitivity of  $\sim 1 \text{ fT/Hz}^{1/2}$  for both electron and nuclear interactions  
 $\Rightarrow$  Frequency uncertainty of  $20 \text{ pHz/month}^{1/2} = 10^{-25} \text{ eV}$  for  $^3\text{He}$   
 $20 \text{ nHz/month}^{1/2} = 10^{-22} \text{ eV}$  for electrons
- So search for preferred spatial direction, reverse co-magnetometer orientation every 20 sec to operate in the region of best sensitivity

# Rotating K-<sup>3</sup>He co-magnetometer

- Rotate – stop – measure – rotate  
⇒ Fast transient response crucial
- Record signal as a function of magnetometer orientation

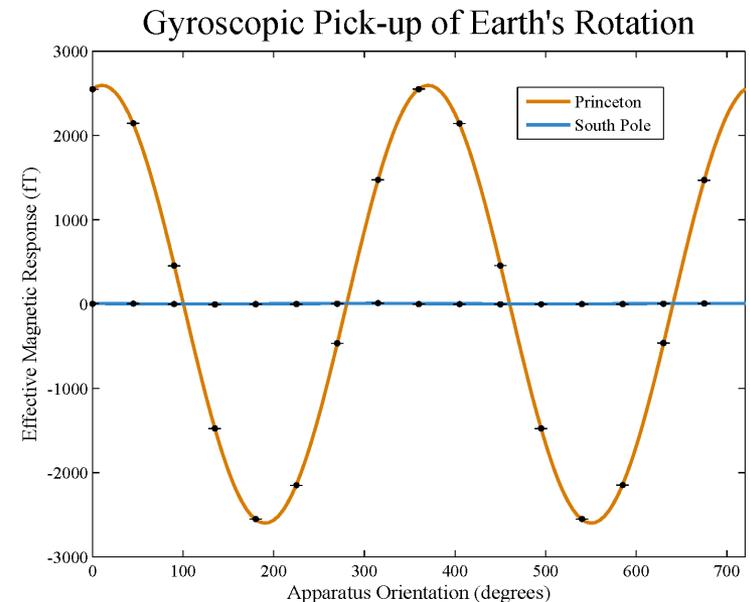
$$b^{eff} = \frac{\Omega}{\gamma}$$

$$S = \frac{P_z \gamma_e \Omega_y}{R} \left( \frac{1}{\gamma_e} - \frac{1}{\gamma_n} \right)$$



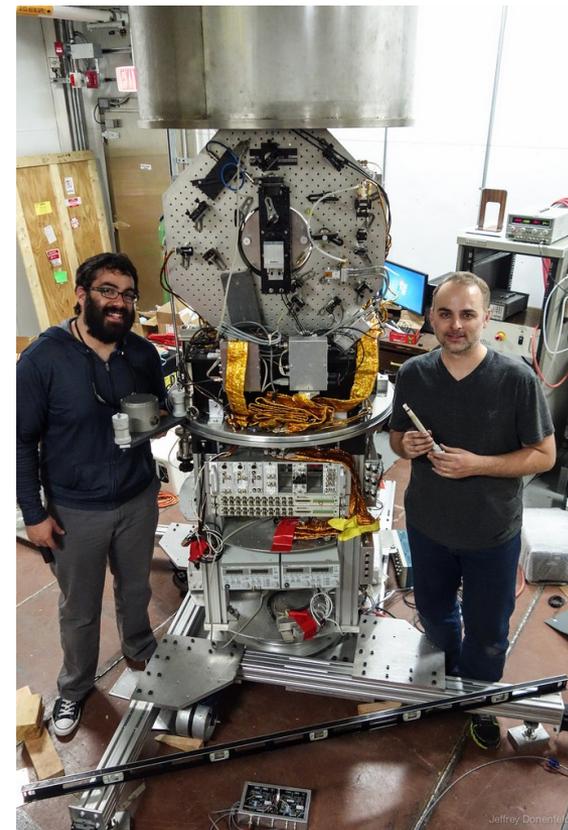
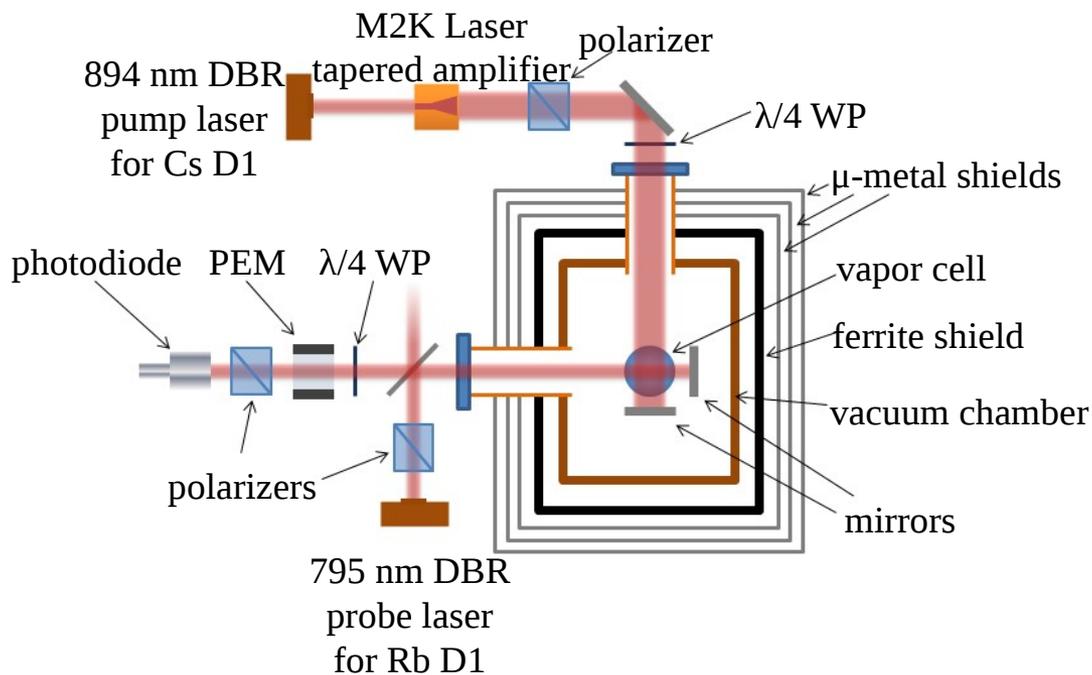
# South Pole

- Most systematic errors are due to two preferred directions in the lab: gravity vector and Earth rotation vector
- If the two vectors are aligned, rotation about that axis will eliminate most systematic errors
- Amundsen-Scott South Pole Station
  - ⇒ Lab location within 200 meters of geographic South Pole



# South Pole Setup

- Use  $^{21}\text{Ne}$  with  $I=3/2$  to look for tensor CPT-even Lorentz-violating effects
- Reliable operation with minimal human intervention:
  - Simple optical setup with DBR diode lasers
  - Whole apparatus in vacuum at 1 Torr
  - Automatic fine-tuning and calibration procedures
  - Remote-controlled mirrors, lasers, etc

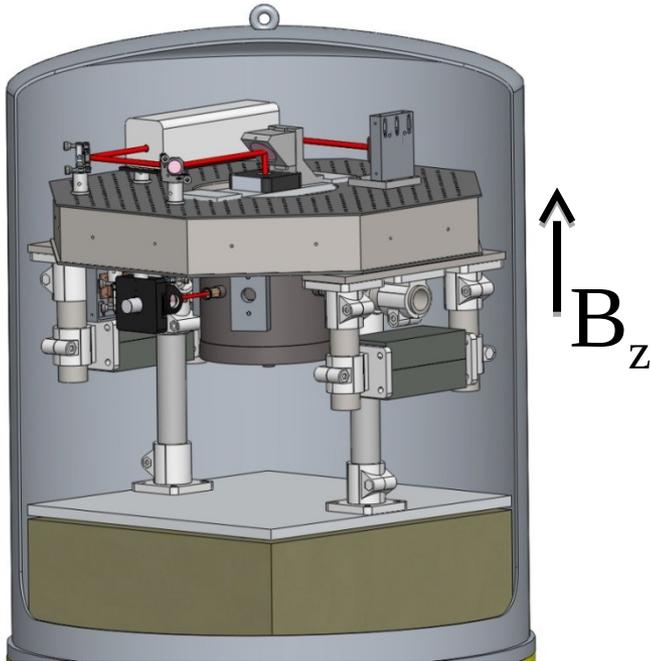


**10 mm**



# Apparatus Orientations

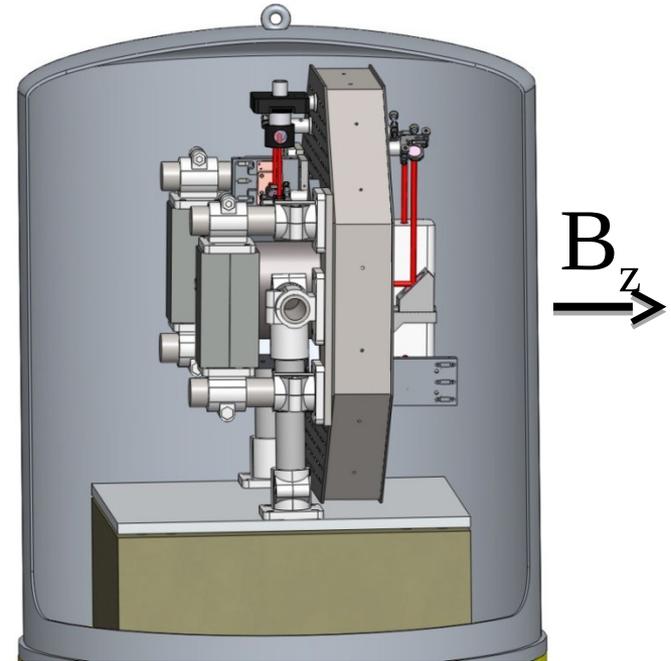
Dipole and quadrupole Lorentz violating coefficients are constrained by operating with the quantization axis in two orthogonal configurations



$B_z$  Vertical

1<sup>st</sup> Harmonic:  $c_X, c_Y, \tilde{b}_X, \tilde{b}_Y$

2<sup>nd</sup> Harmonic: none

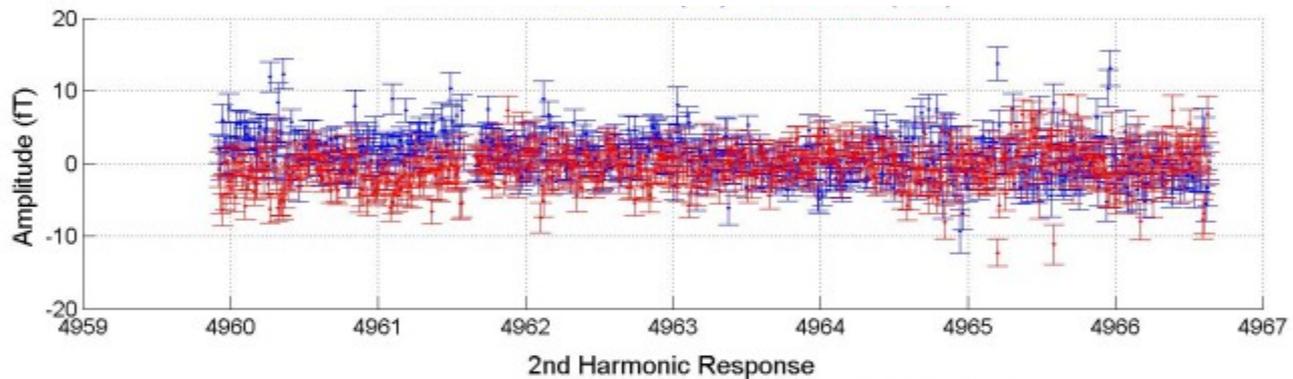
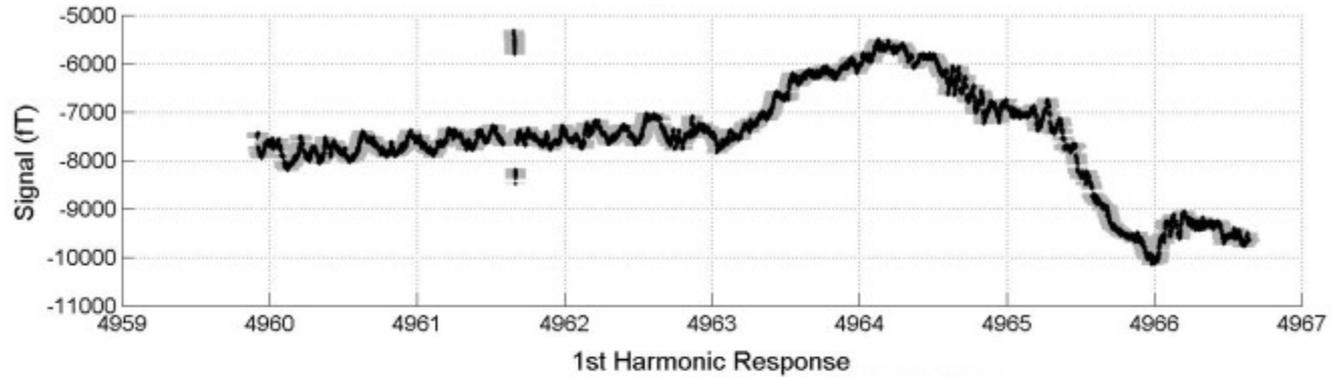


$B_z$  Horizontal

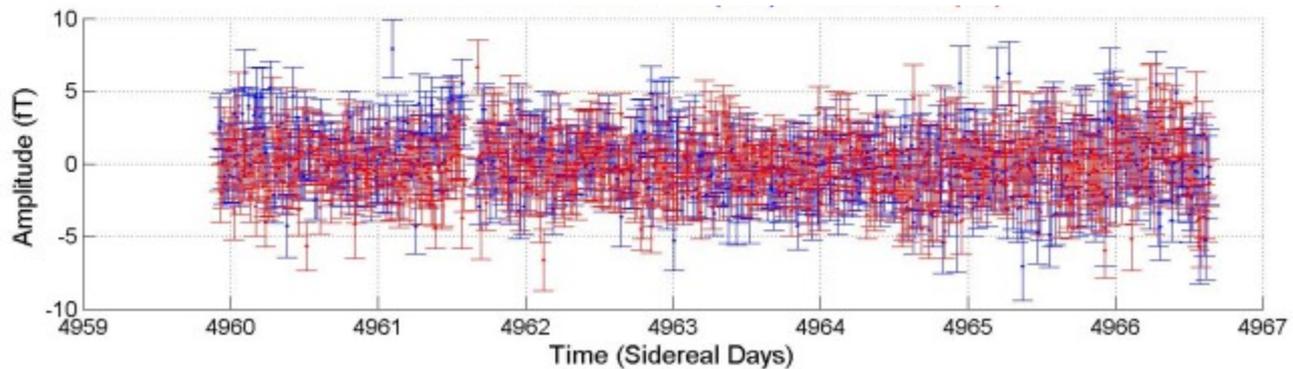
1<sup>st</sup> Harmonic:  $\tilde{b}_X, \tilde{b}_Y$

2<sup>nd</sup> Harmonic:  $c_-, c_Z$

# South pole data sample

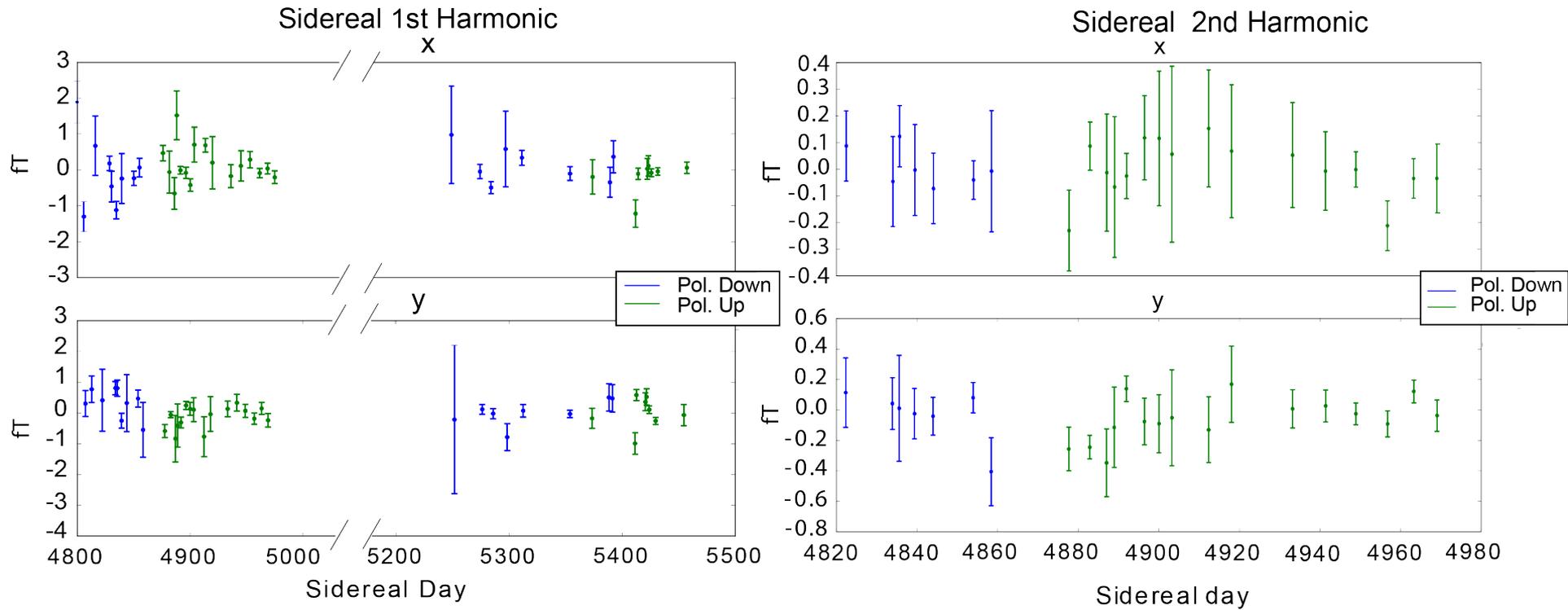


$$\chi^2=1.7$$



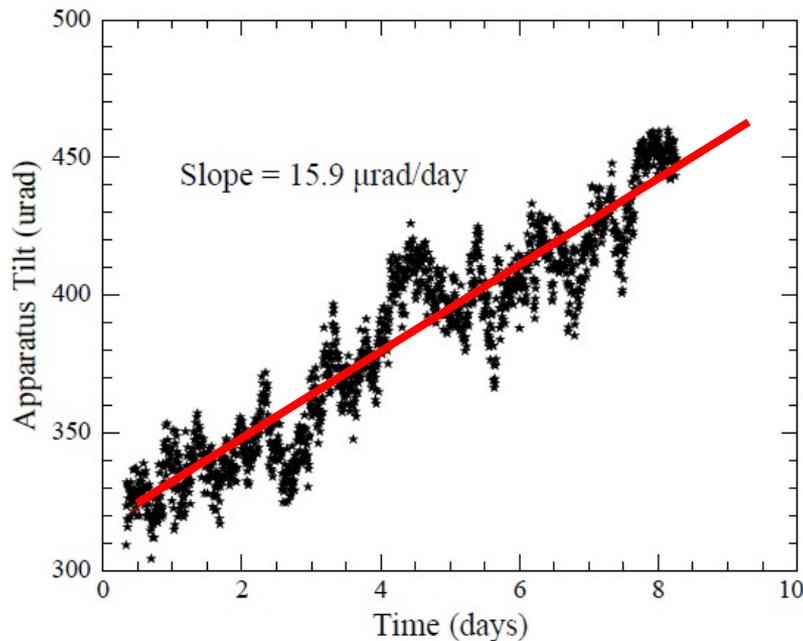
$$\chi^2=1.1$$

# Summary of Lorentz-violation data

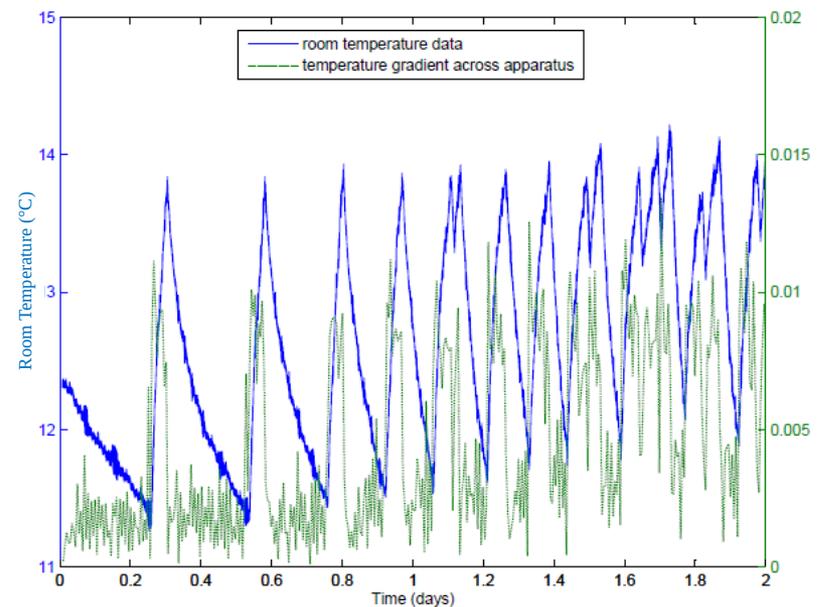


- Two years of data taking
- About 60% duty factor

# Challenges at the Pole



The building's tilt on ice is slowly drifting  
Requires regular automatic leveling



Aggressive temperature cycling  
Temperature gradient across apparatus

## Other challenges:

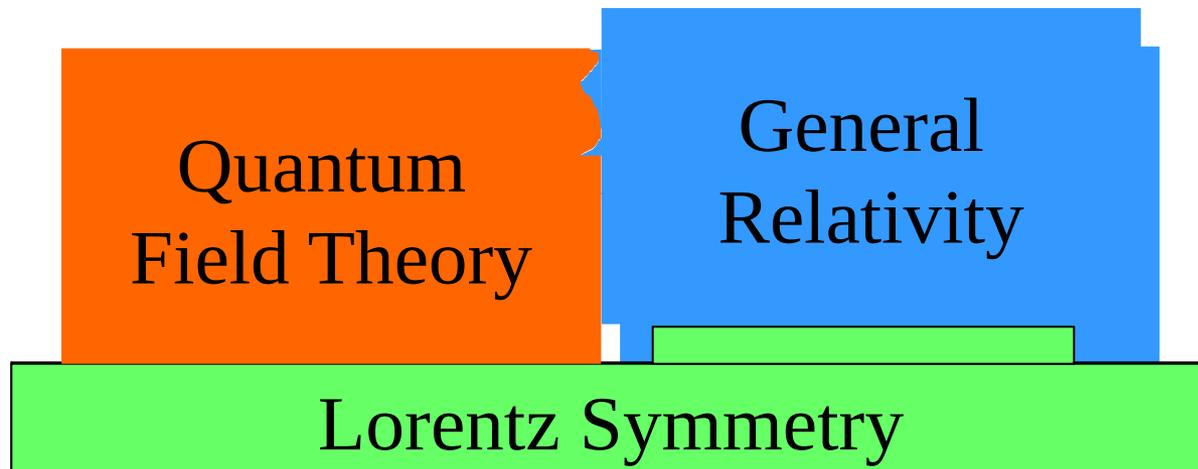
Isolation platform damping failed, probe laser burned out, air-bearing rotation stage got stuck, etc... Need spares for everything.

**First atomic physics experiment operated at the South Pole**  
**First experiment to take advantage of geographic pole location**

# Tests of Lorentz symmetry

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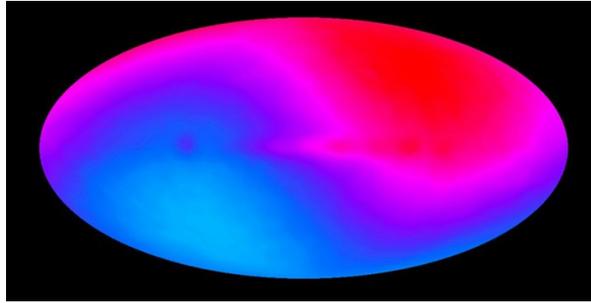
- Lorentz symmetry is at the foundation of two very successful but mutually incompatible theories:
  - ⇒ General Relativity
  - ⇒ Quantum Field Theory
- One approach for resolving this problem is to modify Lorentz symmetry



# Is the space really isotropic?

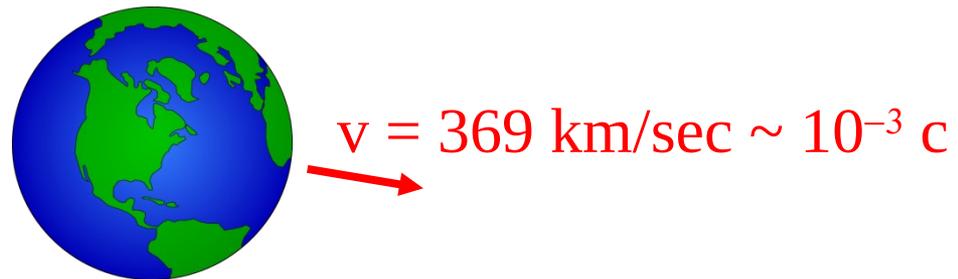
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- Cosmic Microwave Background Radiation Map



⇒ The universe appears warmer on one side!

- Well, we are actually moving relative to CMB rest frame



- ⇒ Space and time vector components mix by Lorentz transformation
- ⇒ A test of spatial isotropy becomes a true test of Lorentz invariance (i.e. equivalence of space and time)

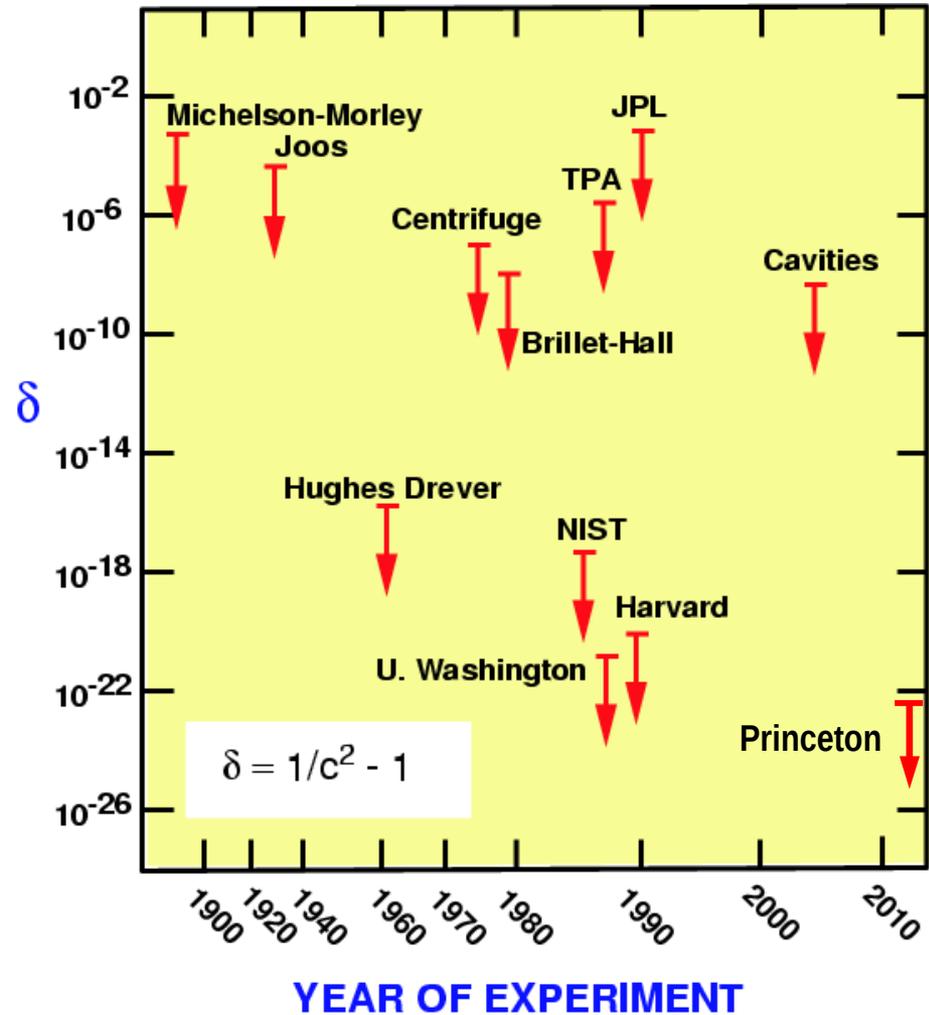
# Local Lorentz Invariance

- Is the speed of light (photons) rotationally invariant in our moving frame?

⇒ First established by Michelson-Morley experiment as a foundation of Special Relativity

- Is the speed of “light” as it enters into particle Lorentz transformation rotationally invariant in the moving frame?

⇒ Best constrained by Hughes-Drever experiments due to finite kinetic energy of nucleons



From Clifford M. Will,  
*Living Rev. Relativity* **9**, (2006)

# Parametrization of Lorentz violation

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$$L = - \bar{\psi} (m + a_{\mu} \gamma^{\mu} + b_{\mu} \gamma_5 \gamma^{\mu}) \psi + \frac{i}{2} \bar{\psi} (\gamma_{\nu} + c_{\mu\nu} \gamma^{\mu} + d_{\mu\nu} \gamma_5 \gamma^{\mu}) \vec{\partial}^{\nu} \psi$$

$a, b$  - CPT-odd  
 $c, d$  - CPT-even  
Alan Kostelecky

⇒  $a_{\mu}, b_{\mu}, c_{\mu\nu}, d_{\mu\nu}$  are vector fields in space with non-zero expectation value

⇒ Vector and tensor analogues to the scalar Higgs vacuum expectation value

- Maximum attainable particle velocity

$$v_{MAX} = c(1 - c_{00} - c_{0j} \hat{v}_j - c_{jk} \hat{v}_j \hat{v}_k)$$

⇒ Implications for ultra-high energy cosmic rays, Cherenkov radiation, etc

⇒ Many laboratory limits (optical cavities, cold atoms, etc)

- Something special needs to happen when particle momentum reaches Planck scale

⇒ Doubly-special relativity

⇒ Horava-Lifshitz gravity

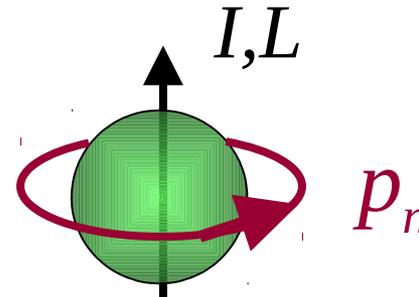
⇒ Your favorite recent theory

# Search for CPT-even Lorentz violation with nuclear spin

- Need nuclei with orbital angular momentum and total spin  $> 1/2$
- Quadrupole energy shift due to angular momentum of the valence nucleon:

$$E_Q \sim (c_{11} + c_{22} - 2c_{33}) \langle p_x^2 + p_y^2 - 2p_z^2 \rangle$$

$$p_x^2 + p_y^2 - 2p_z^2 > 0$$



- Previously has been searched for in experiments using  $^{201}\text{Hg}$  and  $^{21}\text{Ne}$  with sensitivity of about  $0.5 \mu\text{Hz}$

$$\Delta E(t) = E_0 + E_{1X} \cos \Omega t + E_{1Y} \sin \Omega t + E_{2X} \cos 2\Omega t + E_{2Y} \sin 2\Omega t$$

Sidereal Variation                      Semi-sidereal Variation

$$c_{\mu\nu} = \begin{pmatrix} c_{TT} & c_{TX} & c_{TY} & c_{TZ} \\ c_{XT} & c_{XX} & c_{XY} & c_{XZ} \\ c_{YT} & c_{YX} & c_{YY} & c_{YZ} \\ c_{ZT} & c_{ZX} & c_{ZY} & c_{ZZ} \end{pmatrix}$$

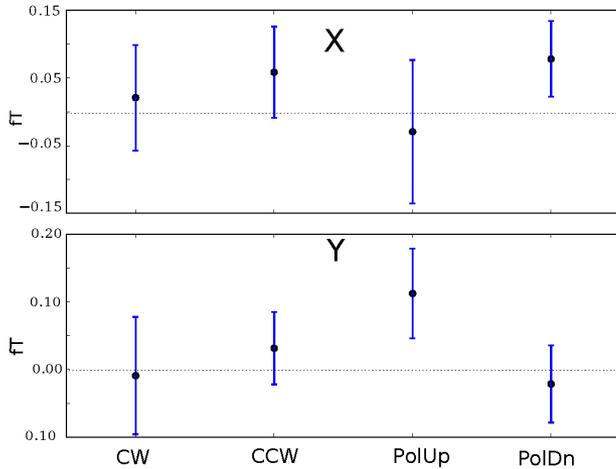
} 2<sup>nd</sup> Harmonic

} 1<sup>st</sup> Harmonic

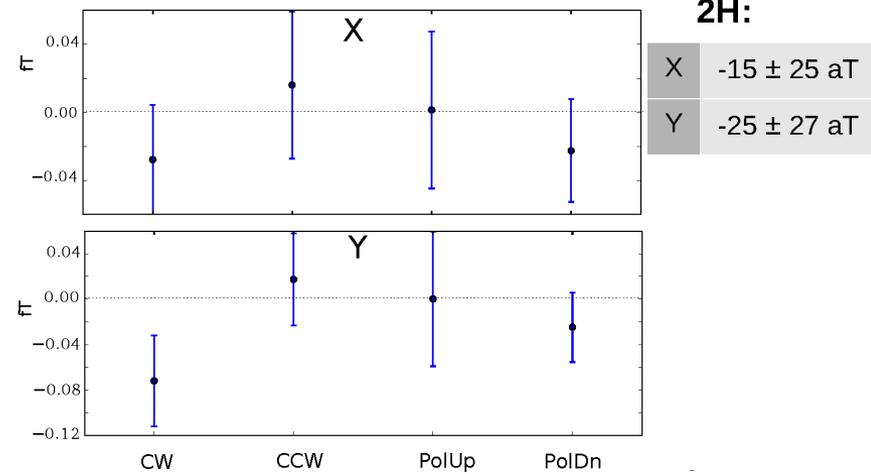
Suppressed by  $v_{\text{Earth}}$

# Preliminary Results

Subsets, 1st Harmonic

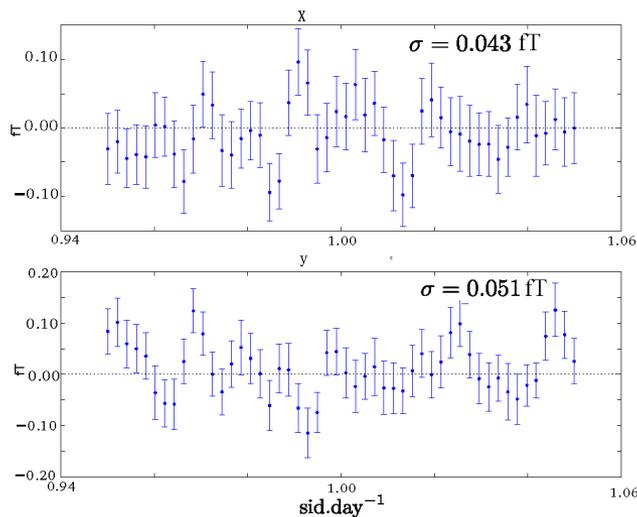


Subsets, 2nd Harmonic

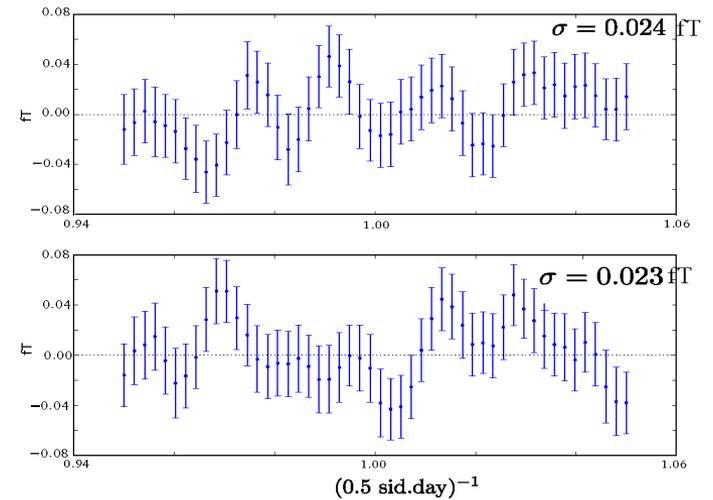


Vary frequency of the fit around sidereal period to independently estimate errors

1H (fit frequency)



2nd Harmonic (vs fit frequency)



# Constraints on SME coefficients

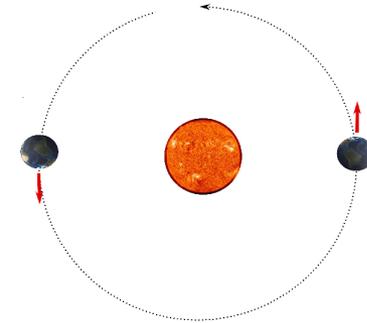
## Spatial

SME coeff	$\bar{\chi}$	$\sigma$	units	Improve
$c_{yz}^n + c_{zy}^n$	2.6	4.5	$\times 10^{-30}$	$\times 10$ [1]
$c_{xz}^n + c_{zx}^n$	2.2	4.8	$\times 10^{-30}$	$\times 10$ [1]
$c_{xy}^n + c_{yx}^n$	-1.9	3.1	$\times 10^{-30}$	$\times 3$ [1]
$c_{xx}^n - c_{yy}^n$	3.0	2.9	$\times 10^{-30}$	$\times 3$ [1]
$\tilde{b}_x$	-1.0	1.5	$\times 10^{-33}$ GeV	$\times 1$ [2]
$\tilde{b}_y$	0.5	1.6	$\times 10^{-33}$ GeV	$\times 1$ [2]

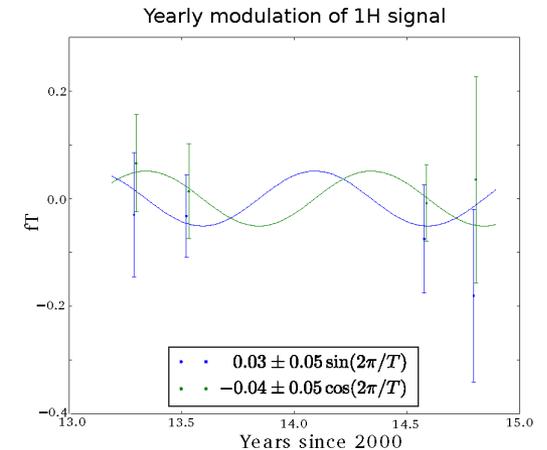
## Boost (suppressed by $v_{\text{earth}}/c$ )

coeff	$\bar{\chi}$	$\sigma$	units	Improve
$c_{tx}^n + c_{xt}^n$	0.1	3.0	$\times 10^{-26}$	-
$c_{ty}^n + c_{yt}^n$	0.4	3.2	$\times 10^{-26}$	-
$c_{tz}^n + c_{zt}^n$	-2.9	4.9	$\times 10^{-26}$	-
$\tilde{b}_t$	-0.25	1.6	$\times 10^{-29}$ GeV	$\times 50$ [3]

## Boost



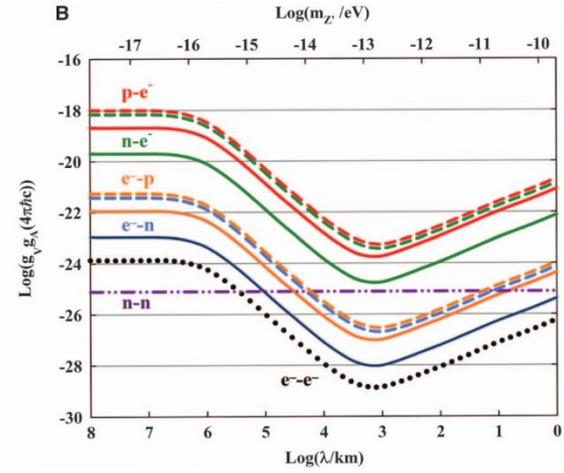
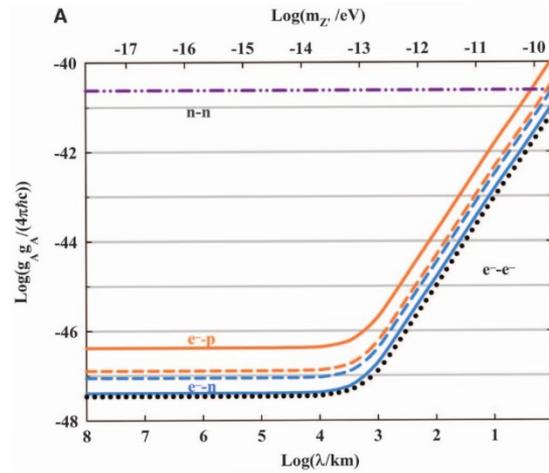
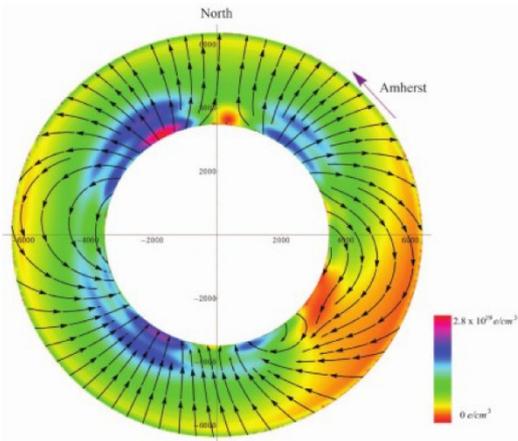
Look for yearly modulation of correlation signal. Suppressed by  $v_{\text{earth}}/c$



- [1] Smiciklas et al., PRL. 107, 171604 (2011)
- [2] Brown et al. PRL 105, 151604 (2010)
- [3] Cane et. al. PRL 93, 230801 (2004)

# Long-range spin-spin interactions with Geo-electrons

## Earth spin-density map



Cause a static lab-frame interaction  
No Earth's rotation background at South Pole

Larry Hunter et al, Science 339, 928 (2013)

Interactions that scale as  $1/r$  or  $1/r^2$   
- Vector bosons mediated,  
- Unparticles

	$\bar{\chi}$	$\sigma$	units	Imprv.
$\frac{g_A^e g_A^n}{4\pi\hbar c}$	-1.2	0.7	$\times 10^{-49}$	$\times 30$ [1]
$\frac{g_V g_A^n}{4\pi\hbar c}$	-5	-2.8	$\times 10^{-25}$	$\times 20$ [1]

# Slowly-modulated signals: light axions, dark photons

Look for modulation of the sidereal frame signal, i.e. amplitude modulation of the correlation signal

$$\delta E = A \sin(\omega t) (\cos(\omega_{sd} t) \hat{X} + \sin(\omega_{sd} t) \hat{Y})$$

For e.g. axion-wind:

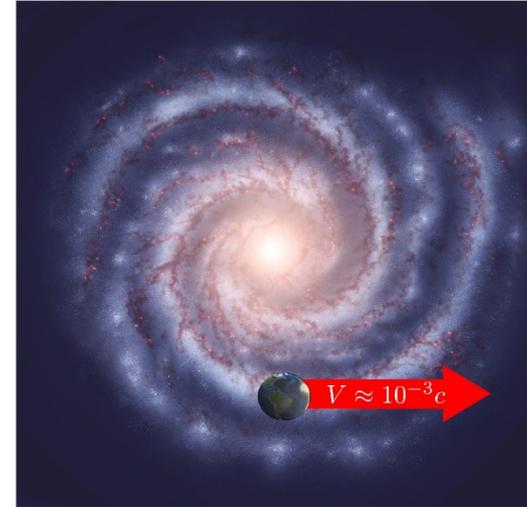
$$\omega = m_a c^2 / \hbar$$

$$A = \frac{(C_N a_0)}{(2 f_a)} (p_a \cdot \sigma_N)$$

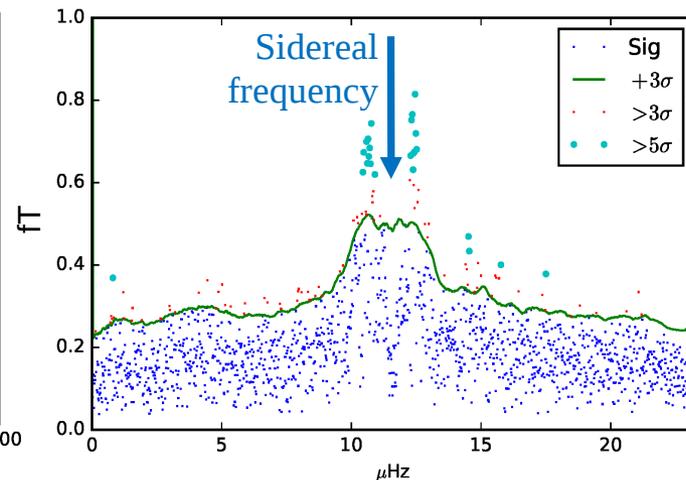
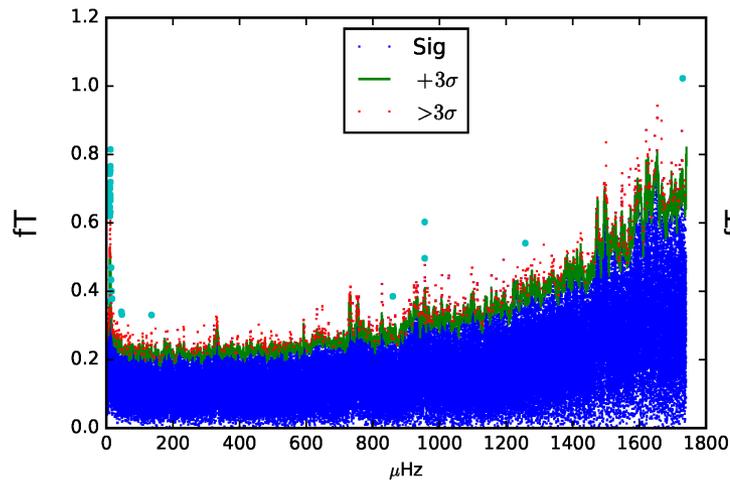
Careful:

Look-elsewhere effects

Interference with sidereal frequency giving rise to slow drifts

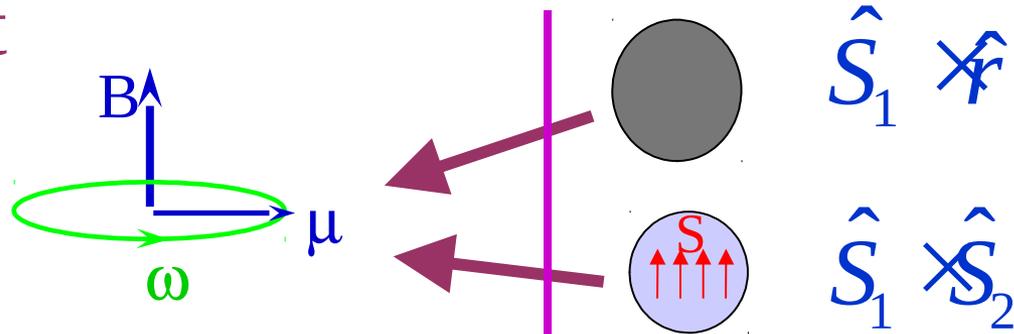


General sensitivity to  $\delta E$  on the order of  $10^{-32}$  GeV in the frequency range 0.1-1500  $\mu\text{Hz}$

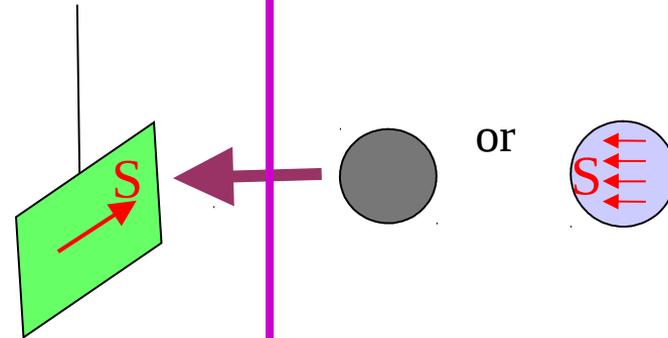


# Searches for spin-dependent forces

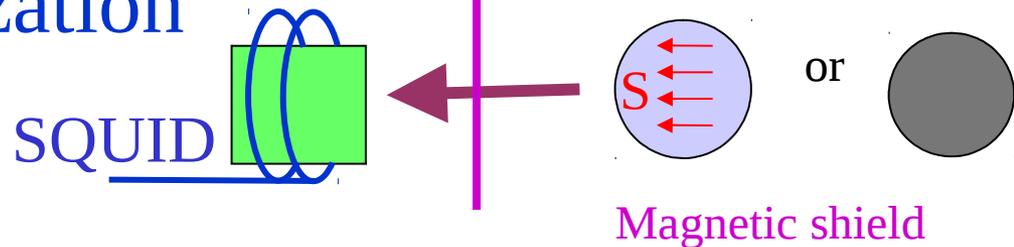
- Frequency shift



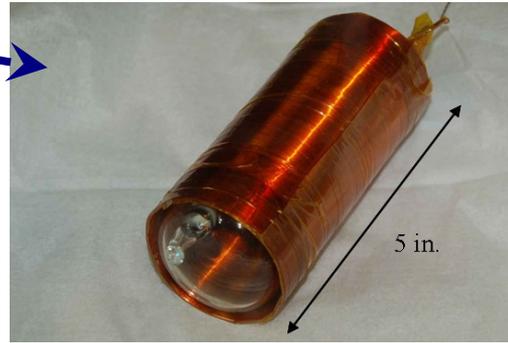
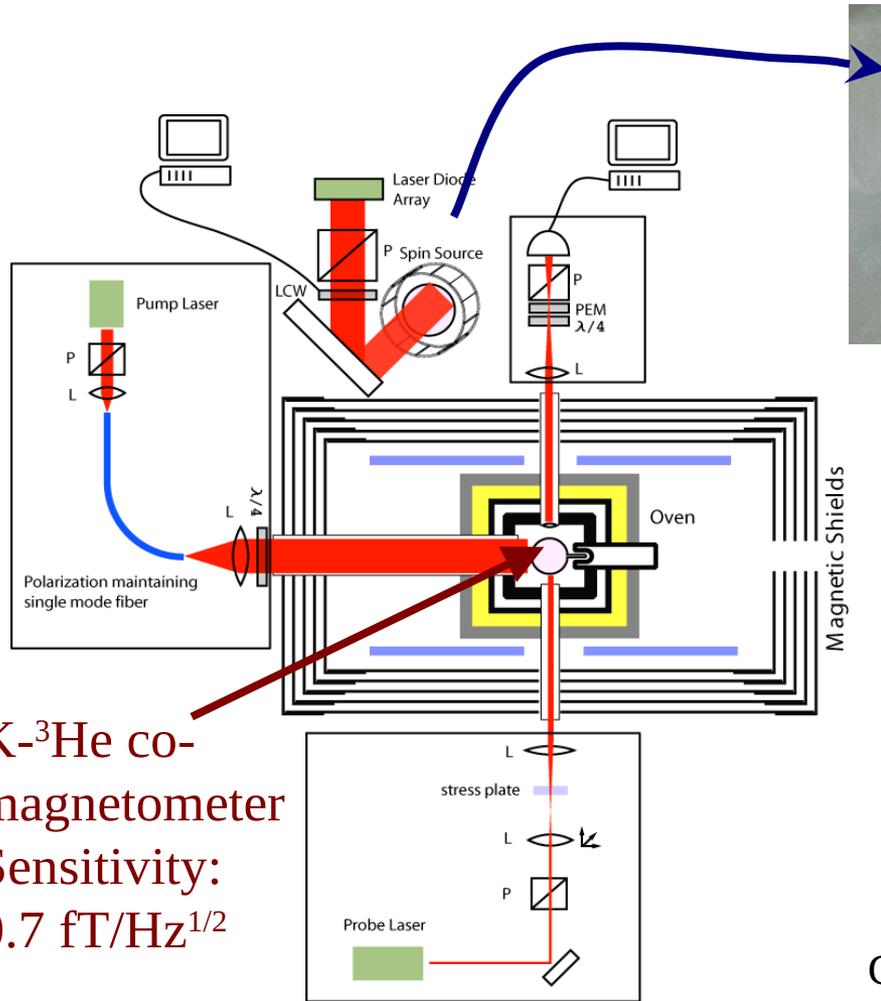
- Acceleration



- Induced magnetization

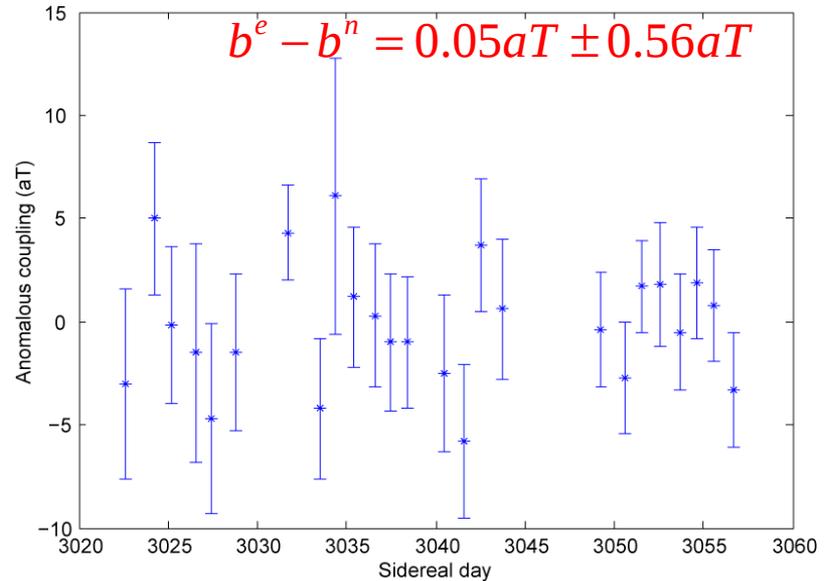


# Search for nuclear spin-dependent forces



Spin Source:  
 $10^{22}$  <sup>3</sup>He spins at 20 atm.  
 Spin direction reversed every 3 sec with Adiabatic Fast Passage

K-<sup>3</sup>He co-magnetometer  
 Sensitivity:  
 $0.7 \text{ fT/Hz}^{1/2}$

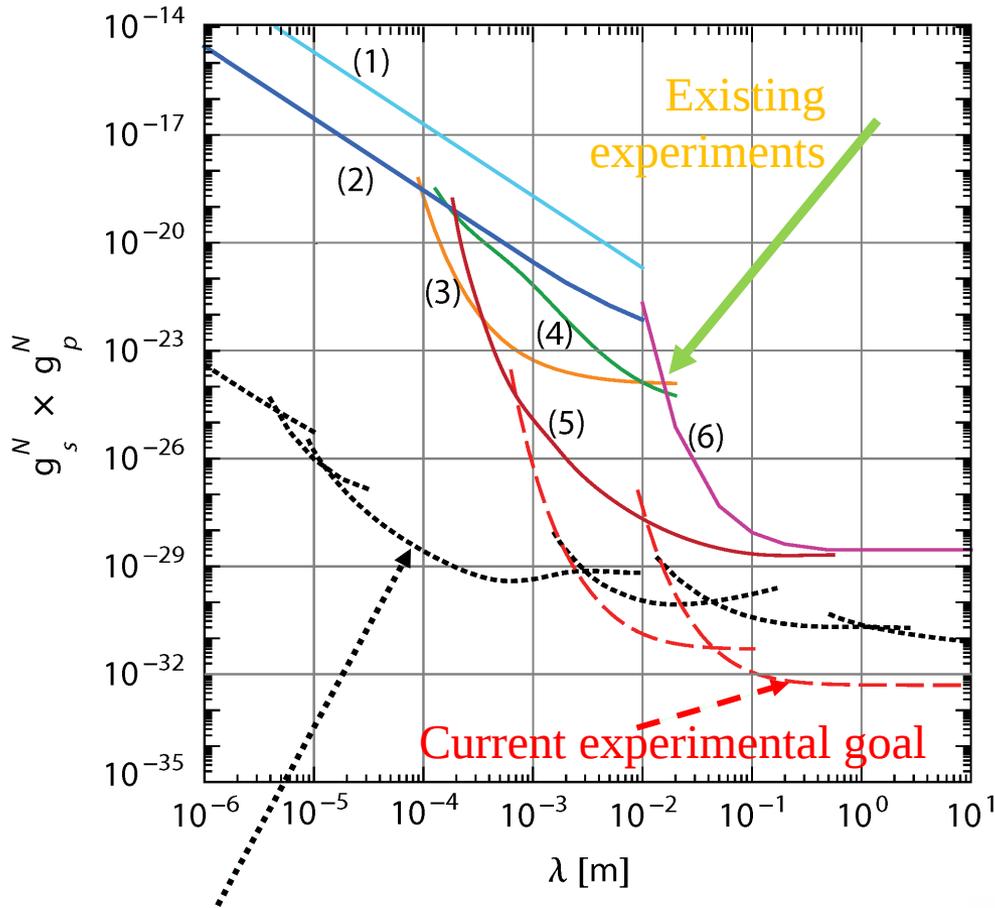


G. Vasilakis, J. M. Brown, T. W. Kornack, MVR, Phys. Rev. Lett. **103**, 261801 (2009)

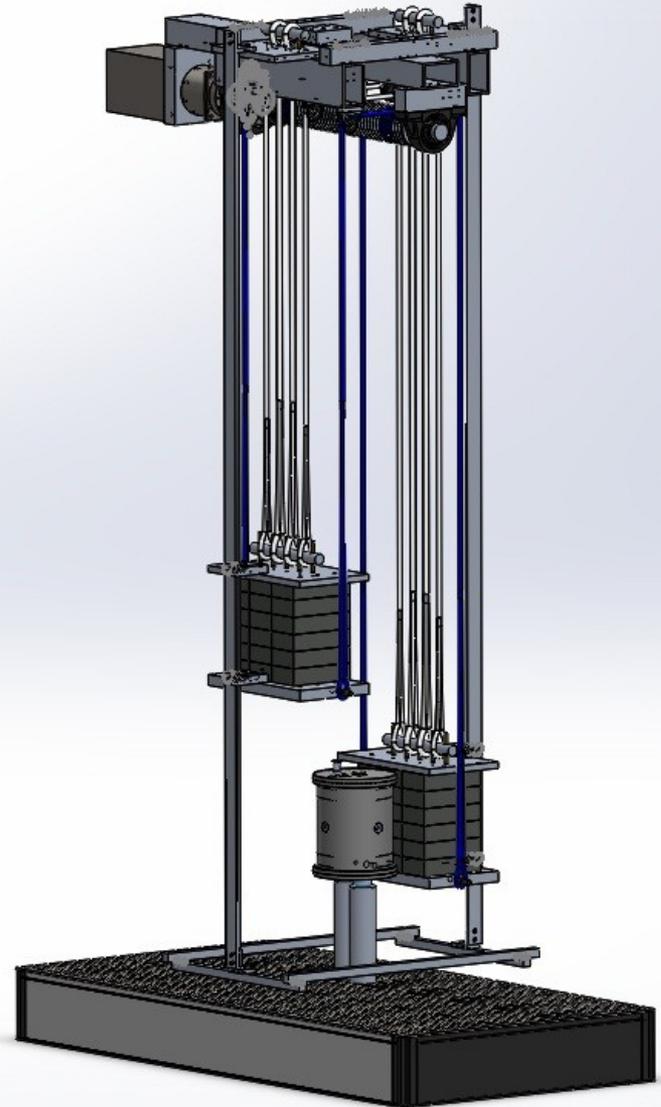
Uncertainty ( $1\sigma$ ) = 18 pHz or  $4.3 \cdot 10^{-35}$  GeV <sup>3</sup>He energy after 1 month  
 Smallest energy shift ever measured

# Spin-mass searches with co-magnetometer

- Will be more sensitive than astrophysical limits



Astrophysical  $\times$  gravitational limits from G. Raffelt  
Phys. Rev. D **86**, 015001 (2012)



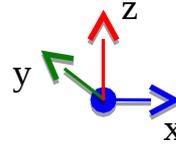
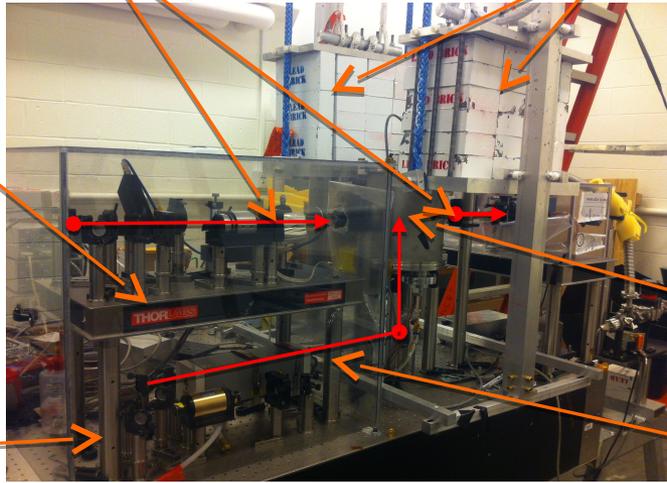
# Movable mass constructed and tested

Probe beam Pb Masses (~ 210 kg each)

Apparatus is suspended from ceiling to reduce mechanical coupling to optical table

Probe Beam Optics

Pump Beam Optics



Glass cell (within vacuum chamber)

Pump beam

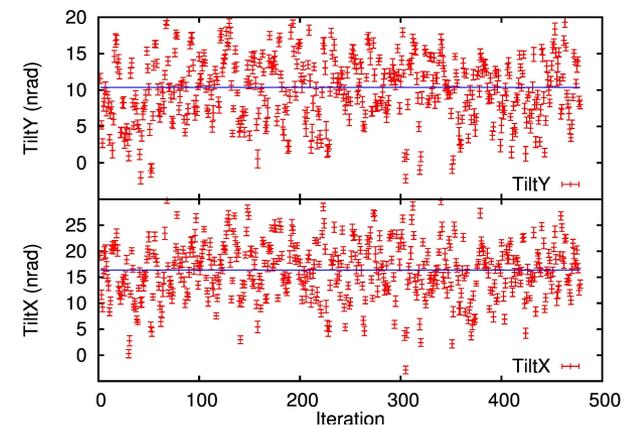
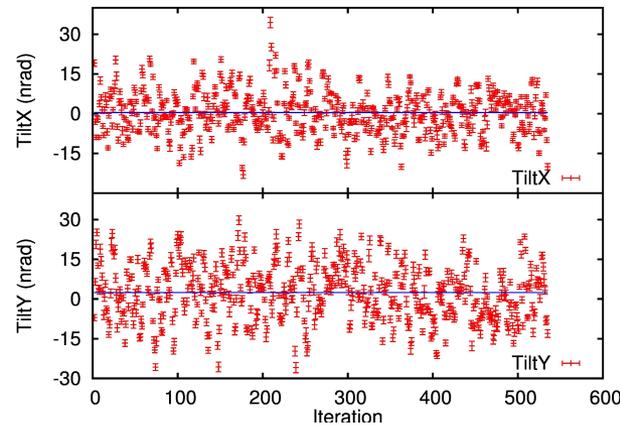
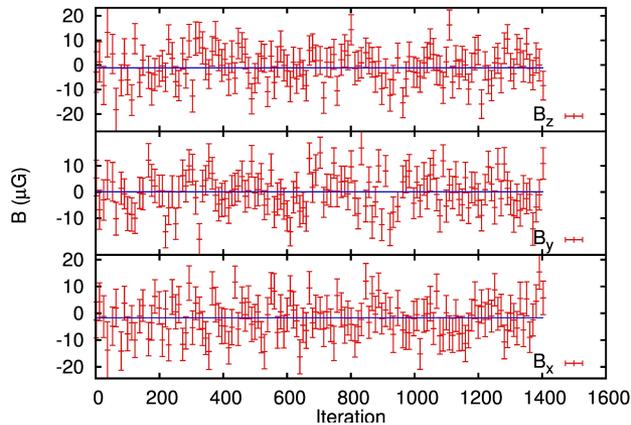


A Yaskawa 4kW servo motor smoothly raise/lower the two 210 kg Pb weights through a distance of 0.5m in 1s.

Magnetic correlation  $\sim 2 \mu\text{G}$

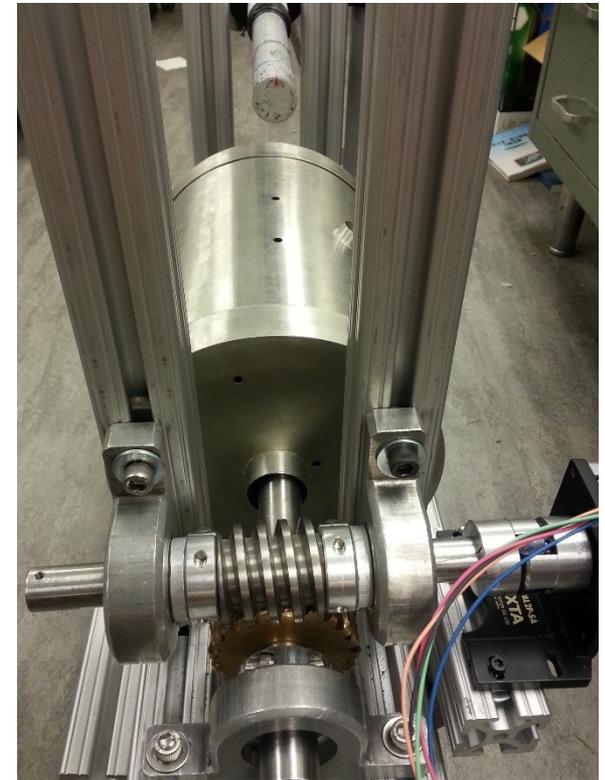
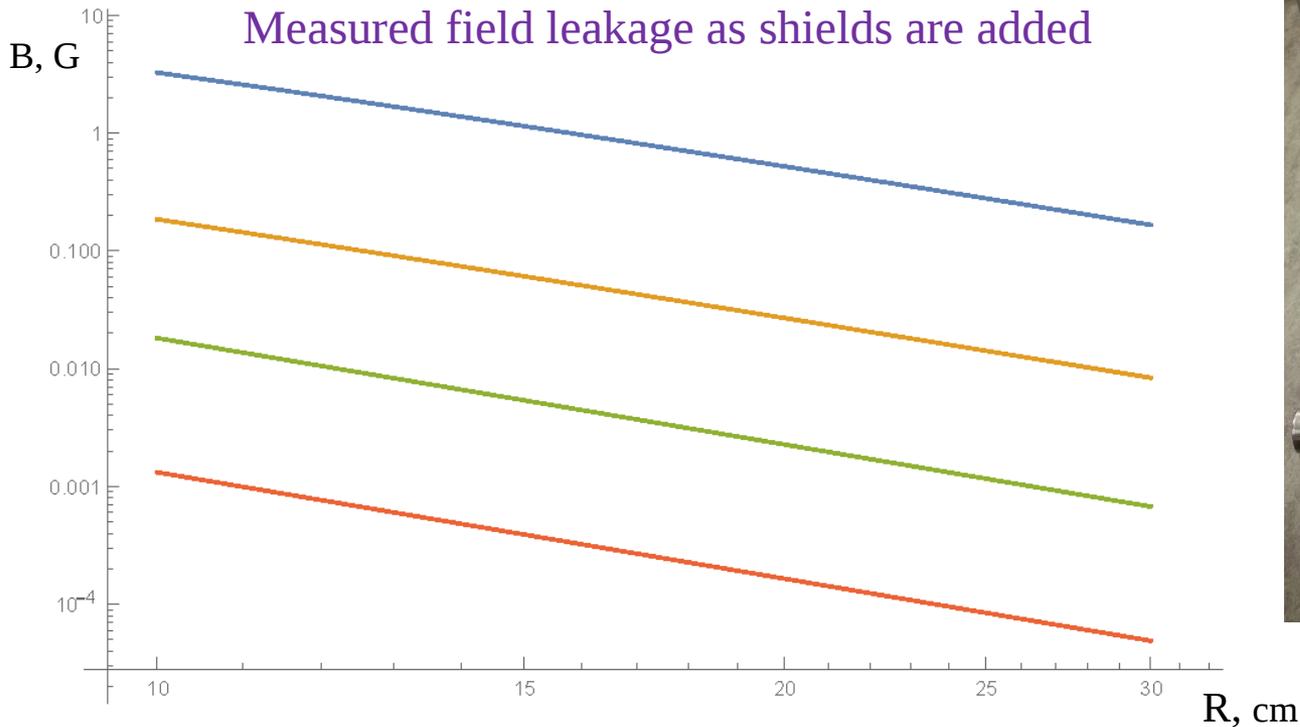
Optical table tilt correlation  $\sim 2 \text{ nrad}$

Gravitational effect on Tiltmeter  $19.4 \pm 0.5 \text{ nrad}$



# Spin-spin long-range force

- Use a permanent magnet spin source with  $10^{24}$  polarized electrons (Eöt-Wash approach)
- Use co-magnetometer as spin sensor
- Limits both e-e and e-n interactions
- Expect  $g_p \sim 10^{-9}$ , better than current laboratory limits but not quite reaching astrophysics limits
- Currently testing magnetic field leakage with 3 shields



# Conclusions

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- Atomic co-magnetometers set the most stringent limits on both CPT-odd and CPT-even Lorentz –violation coefficients
- Set limits on spin-dependent forces at 20 pHz level, the most sensitive energy shift measurements
- Can place limits on oscillating spin couplings in the  $\mu\text{Hz}$ -Hz range
- Search for spin-mass coupling under way, should exceed astrophysical limits.

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