

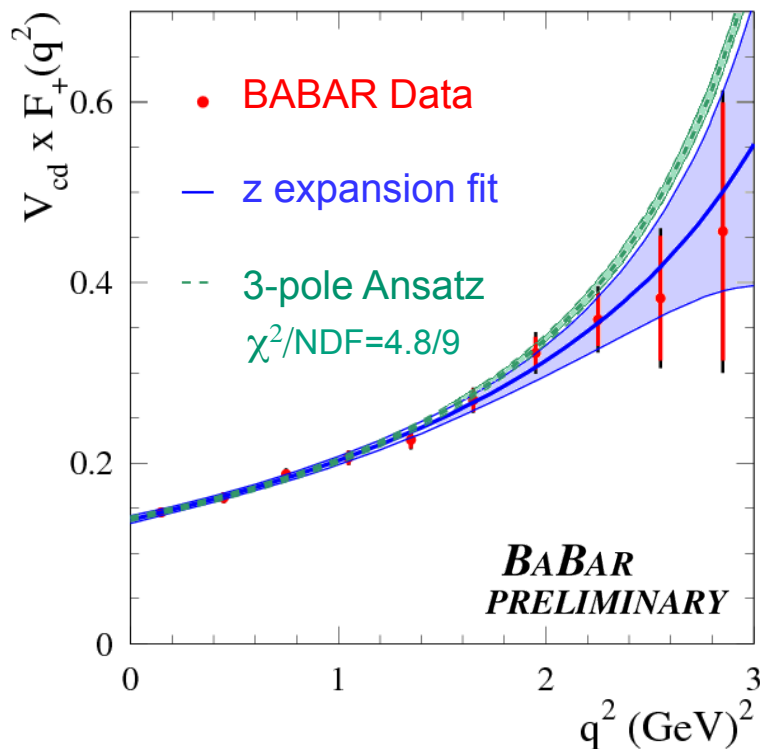
3-Pole Approximation to $D^0 \rightarrow \pi^- e^+ \nu$ Form Factors

Becirevic et al. arXiv:1407.1019

- As recently proposed, the addition of a 3rd effective pole greatly improves the prediction

$$f_+(q^2) = \frac{\gamma_0}{m_{D^*}^2 - q^2} + \frac{\gamma_1}{m_{D^{*'}}^2 - q^2} + \frac{\gamma_{\text{eff}}}{m_{\text{eff}}^2 - q^2}$$

$$\text{with } \gamma_n = \text{Res}_{q^2=m_{D_n^*}^2} f_+^{D\pi}(q^2)$$



Constraining the residues γ_0 and γ_1 to expectation,

$$\gamma_0 = 4.17 \pm 0.13 \text{ GeV}^2 \quad \gamma_1 = -1.1 \pm 0.4 \text{ GeV}^2$$

based on various LQCD calculations

with $m_{D^{*+}} = 2,010 \pm 0.13 \text{ MeV}$ $m_{D^{*'}} = 2,610 \pm 4 \text{ MeV}$

and imposing $\gamma_{\text{eff}} + \gamma_0 + \gamma_1 = 0$

results in

$$m_{\text{eff}} = (3.6 \pm 0.3) \text{ GeV}$$

a value larger than the mass of the next $J^P=1^-$ state of 3.1 GeV predicted by quark models.

3-pole ansatz fits the data very well up to $\sim 2 \text{ GeV}^2$!

3-Pole Fit to $B \rightarrow \pi l^+ \nu$ Decay Rate

- Following suggestions by theorists, we applied the same ansatz to B decays

$$f_+^{B\pi}(q^2) = \frac{\beta_0}{m_{B^*}^2 - q^2} + \frac{\beta_1}{m_{B^{*'}}^2 - q^2} + \frac{\beta_{\text{eff}}}{m_{\text{eff}}^2 - q^2}$$

constraining the residua β_0 and β_1 to expectation,

$$\beta_0 = 24.9 \pm 4.2 \text{ GeV}^2 \quad \beta_1 = -8 \pm 2 \text{ GeV}^2$$

based on various LQCD calculations

and imposing $\beta_{\text{eff}} + \beta_0 + \beta_1 = 0$

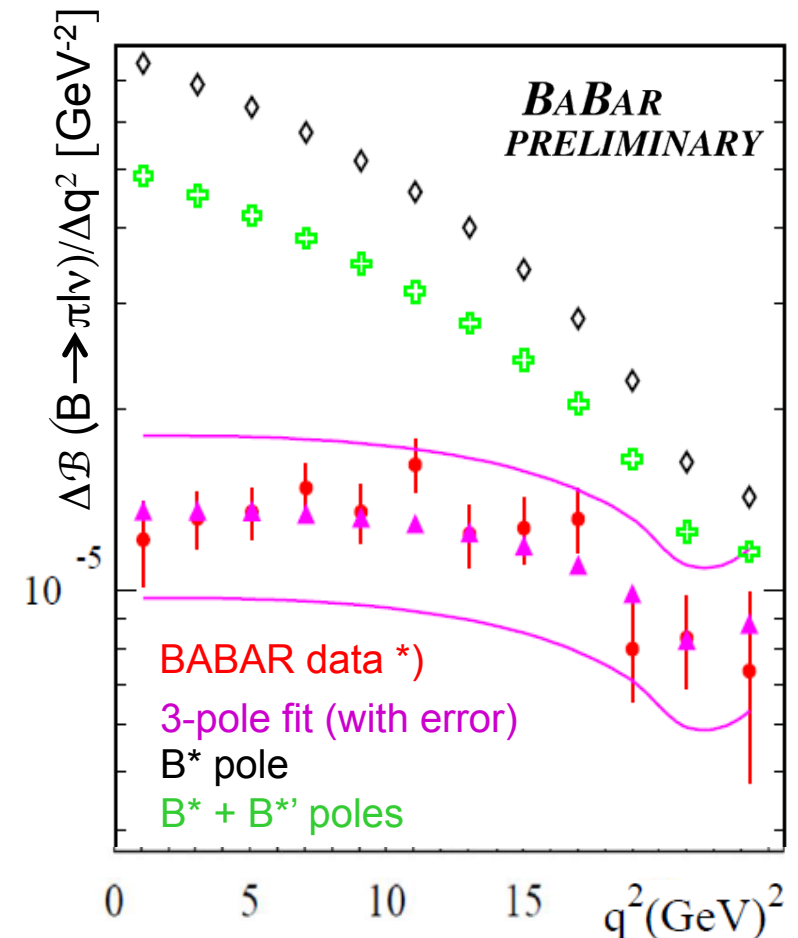
with $m_{B^*} = 5.325 \text{ GeV}$ $m_{B^{*'}} = 5.491 \text{ GeV}$

we obtain

$$m_{\text{eff}} = (7.4 \pm 0.4) \text{ GeV}$$

$$|V_{ub}| = (2.6 \pm 0.2_{\text{exp}} \pm 0.4_{\text{LQCD}})$$

Uncertainties (15%) dominated by uncertainty in coupling constants $g_{B^*B\pi}$, taken from LQCD!



*) BABAR PRD, 092004 (2012)

Extrapolation of FF from $D^0 \rightarrow \pi^- e^+ \nu$ to $B^0 \rightarrow \pi^- l^+ \nu$

- It has been suggested years ago, that dynamics of D and B meson decays should be closely related.
- Specifically, for $w_B = w_D$ kinematic factors cancel

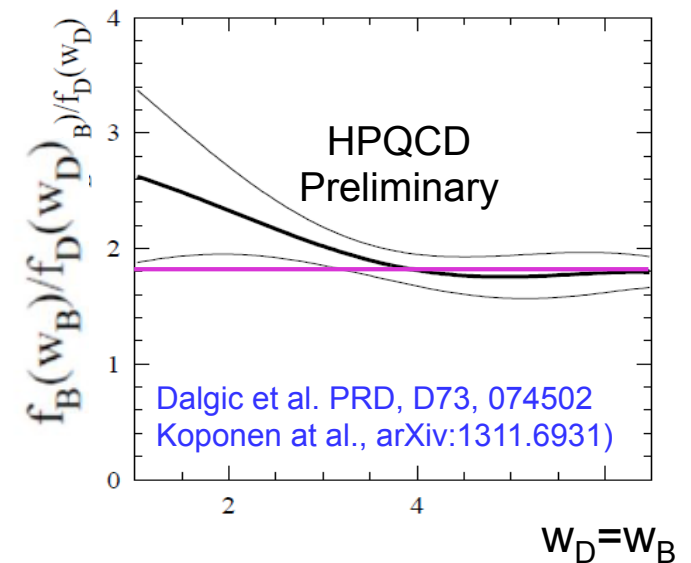
$$\frac{dB^B}{dw} = \frac{dB^D}{dw} \Big|_{\text{meas.}} \frac{m_B \tau_B}{m_D \tau_D} \left(\frac{|V_{ub}|}{|V_{cd}|} \right)^2 R_{BD}^2$$

with $w_{B,D} = \frac{M_{B,D}^2 + m_\pi^2 - q^2}{2M_{B,D}m_\pi}$

Kinematically, there is a common range for

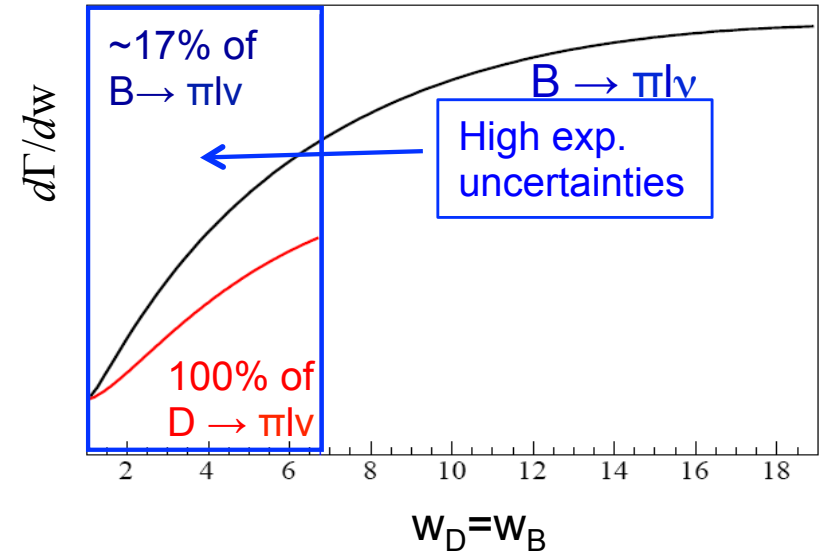
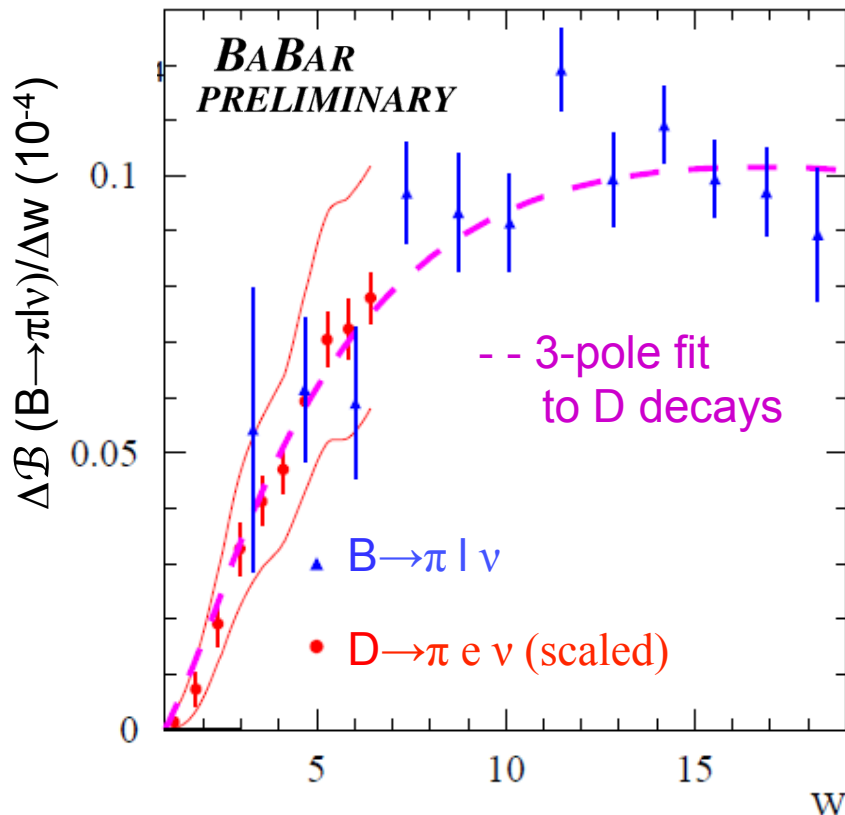
$$w_B = w_D \{ 1.0 - 6.7 \} \quad q_{\text{max}}^2 \longrightarrow w=1$$

- Future LQCD calculation might determine the FF ratio R_{BD} with high precision.
- Preliminary calculation by HPQCD indicate small variation of the ratio of individual FF R_{BD} as function of w ; average for $w > 4$: $\langle R_{BD} \rangle = 1.8 \pm 0.2$



Extrapolation of FF from $D^0 \rightarrow \pi^- e^+ \nu$ to $B^0 \rightarrow \pi^+ l^+ \nu$

Derive $\mathcal{B}(B^0 \rightarrow \pi^+ l^+ \nu)$ from $D^0 \rightarrow \pi^- e^+ \nu$ data with $R_{BD} = 1.8 \pm 0.2$ and adjust $|V_{ub}|$ to 3.65×10^{-3}



- Agreement very good, may not be a surprise: Translation based on same LQCD calculation for B s.l. decays, used to extract $|V_{ub}|$
- Alternatively, use 3-pole FF fit to fit to $\mathcal{B}(D^0 \rightarrow \pi^+ l^+ \nu)$ and extrapolate to unphysical region, obtain independent result for $|V_{ub}|$

$$|V_{ub}| = (3.65 \pm 0.18_{\text{exp}} \pm 0.40_{\text{RD}}) \times 10^{-3}$$

Need improvement in LQCD calculation!