

Color-Dual EFT

MITP 2022

Based on ongoing work with
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Suna Zekioglu



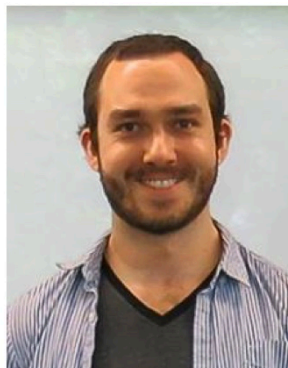
Nic Pavao



Aslan Seifi



Ingrid Vazquez-Holm
Bogdan Stolica

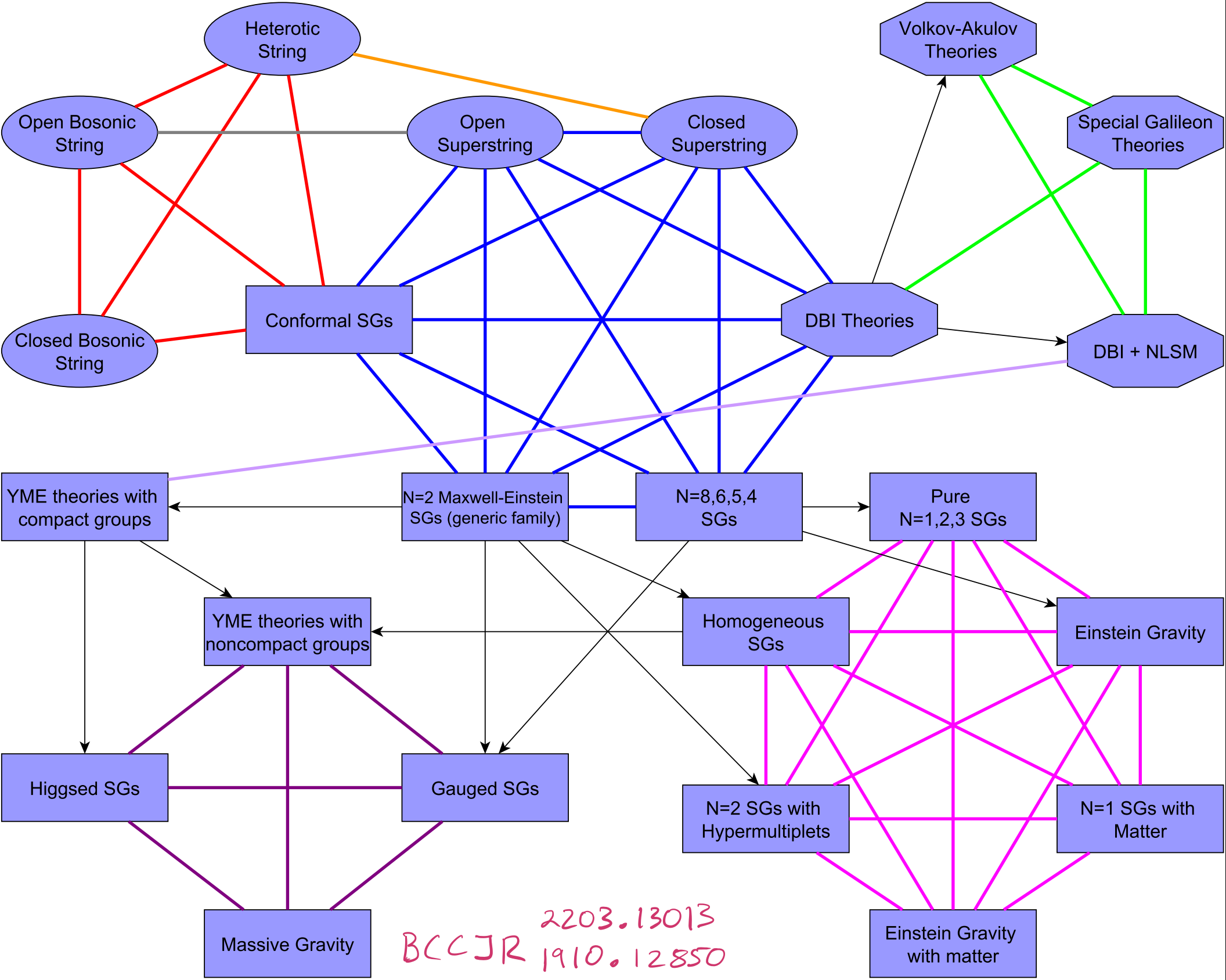


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James Mangan



Matt Lewandowski
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INTRODUCTION

double copy invites us to construct
+ combine our favorite theories using
building blocks

$$\text{SYM} = \text{color} \otimes \text{vector} + \text{susy}$$

$$\text{NLSM} = \text{color} \otimes \pi$$

$$\text{SDBIVA} = \pi \otimes \text{vector} + \text{susy}$$

$$\text{SG} = \text{vector} + \text{susy} \otimes \text{vector} + \text{susy}$$

$$\otimes \equiv \left(\underset{\text{(tree level)}}{\text{KLT}} = \underset{\text{(tree level)}}{\text{BCJ}} = \text{CHY} \right)$$

This talk will largely be at Tree Level
but borrowing insight & approaches from
Loop level

Will refine language around double

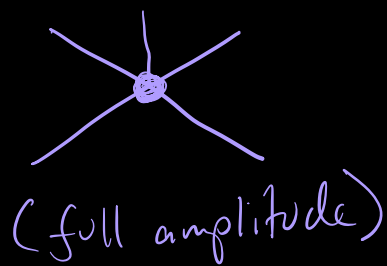
copy via Something New: alg. composition

JJMC, ZR, ZY, SŽ

1910.12850

The diagram shows an equation between two Feynman diagrams. On the left, a wavy line enters a circle containing a double slash $//h$, with an arrow pointing upwards from the top of the circle. A wavy line exits the right side of the circle. This is followed by an equals sign. On the right, there are two identical Feynman diagrams multiplied together, indicated by a central dot. Each diagram consists of a wavy line entering a circle with $//h$ and an upward arrow, and a wavy line exiting the right side.

Composition $\neq$ Double copy



$=$

(adjoint)

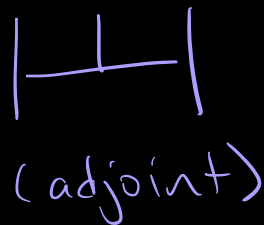


Double
copy

(adjoint)



v.s.



$=$



(adjoint composition)

E.g.

$\overline{LL}$

PIONS

$=$

$(D\phi)^2$

$\uparrow$

covariantized



$(D\phi)^2$

$\uparrow$

free scalar

Summary:

- composition lets us build ε 1910.12850
2104.08370

classify EFT operators for
gauge & gravity theories

with SMALL set of Building Blocks

- $N=4$ SG may have a 2203.03592

non-local color-dual fate:

- Berkovits Witten CSG
- Heterotic String

EFT

- Higher derivative operators (or counterterms)

let you parameterize your ignorance of new physics
with available fields

$$\mathcal{L} = -\partial^2 \phi + g \phi^3 + c_2 \frac{(\partial^2 \phi^2)}{\Lambda^2} + c_4 \frac{\partial^4 \phi^4}{\Lambda^4} + \dots$$

Goal :

$$\partial^n F^m, \partial^n R^m$$

How we calculate at loops

- Not Feynman diagrams

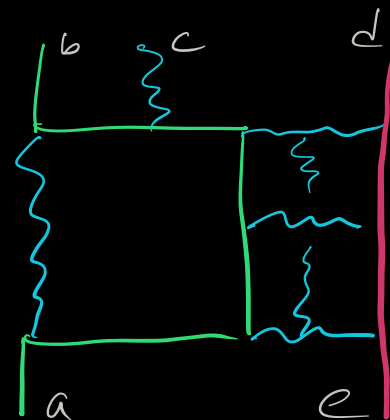
- off shell unweildy
- factorial # of diagrams

- But Graph Organized

- unitarity methods efficient for
VERIFICATION & Cut construction
- Virtual integration wants local reps.

(Gauge) Loop Integrand Map

Mapping from labeled topology



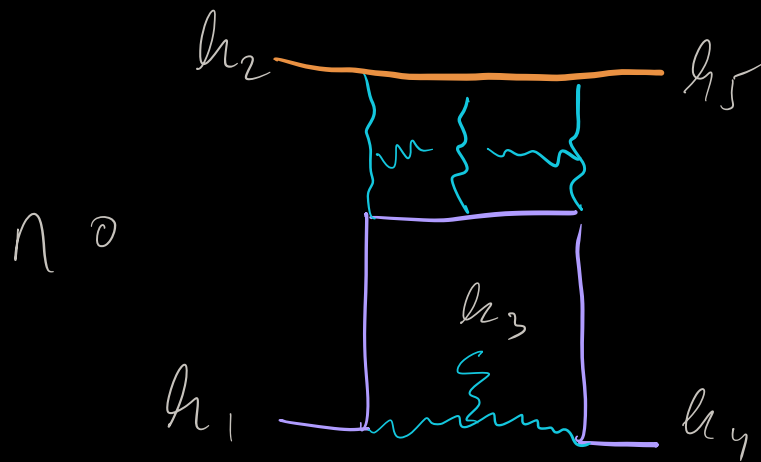
to color weights $\left(\text{wavy line} f^{abc} \frac{i}{\xi_a} T_{ij}^a \right) \subset \mathbb{H}$

propagator weights $\left(\xrightarrow{\ell} = \frac{1}{\ell^2 - m^2} \right)$

numerator weights $(k_a \dots k_e, \epsilon_c, l_1, \dots, l_4)$

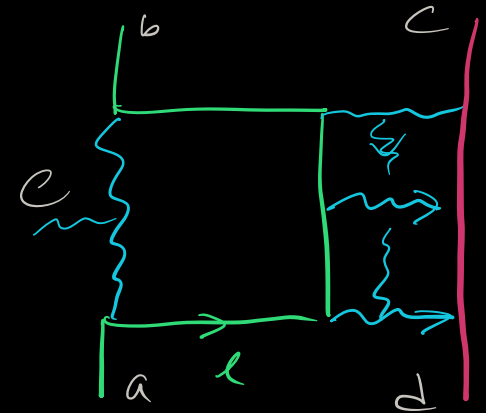
Weights should obey

symmetries of topology



?

$a \rightarrow h_1$		$a \rightarrow h_4$
$b \rightarrow h_2$	or	$b \rightarrow h_5$
$c \rightarrow h_3$		$c \rightarrow h_3$



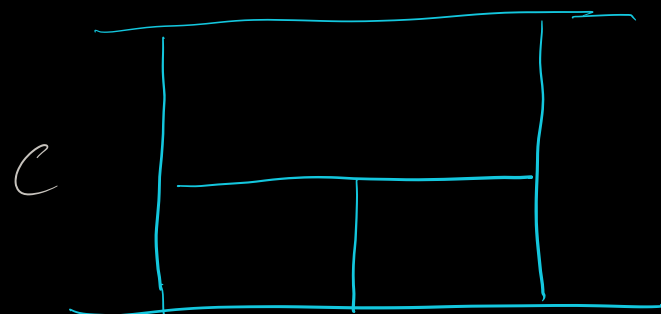
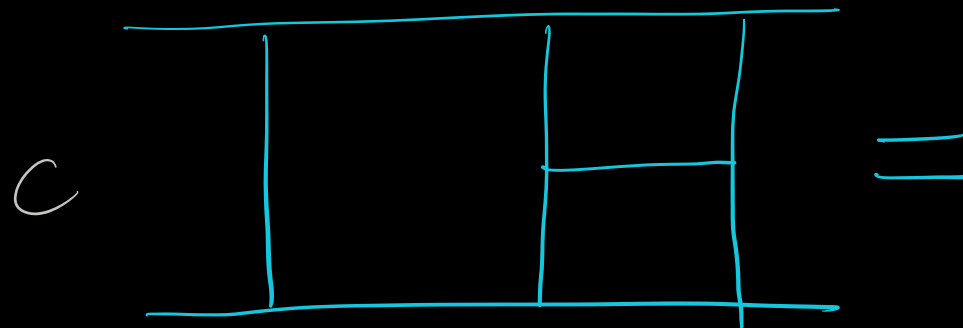
ϵ : unitarity cuts.

Recipe for Calculation

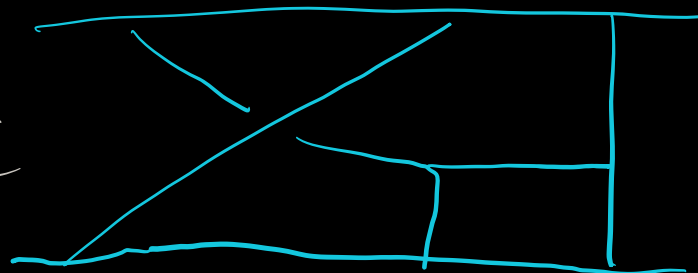
- Give Ansätze to topologies
- Fix on Symmetries & Cuts
- Duality between color & kinematics
can help!

Adjoint-type color weights

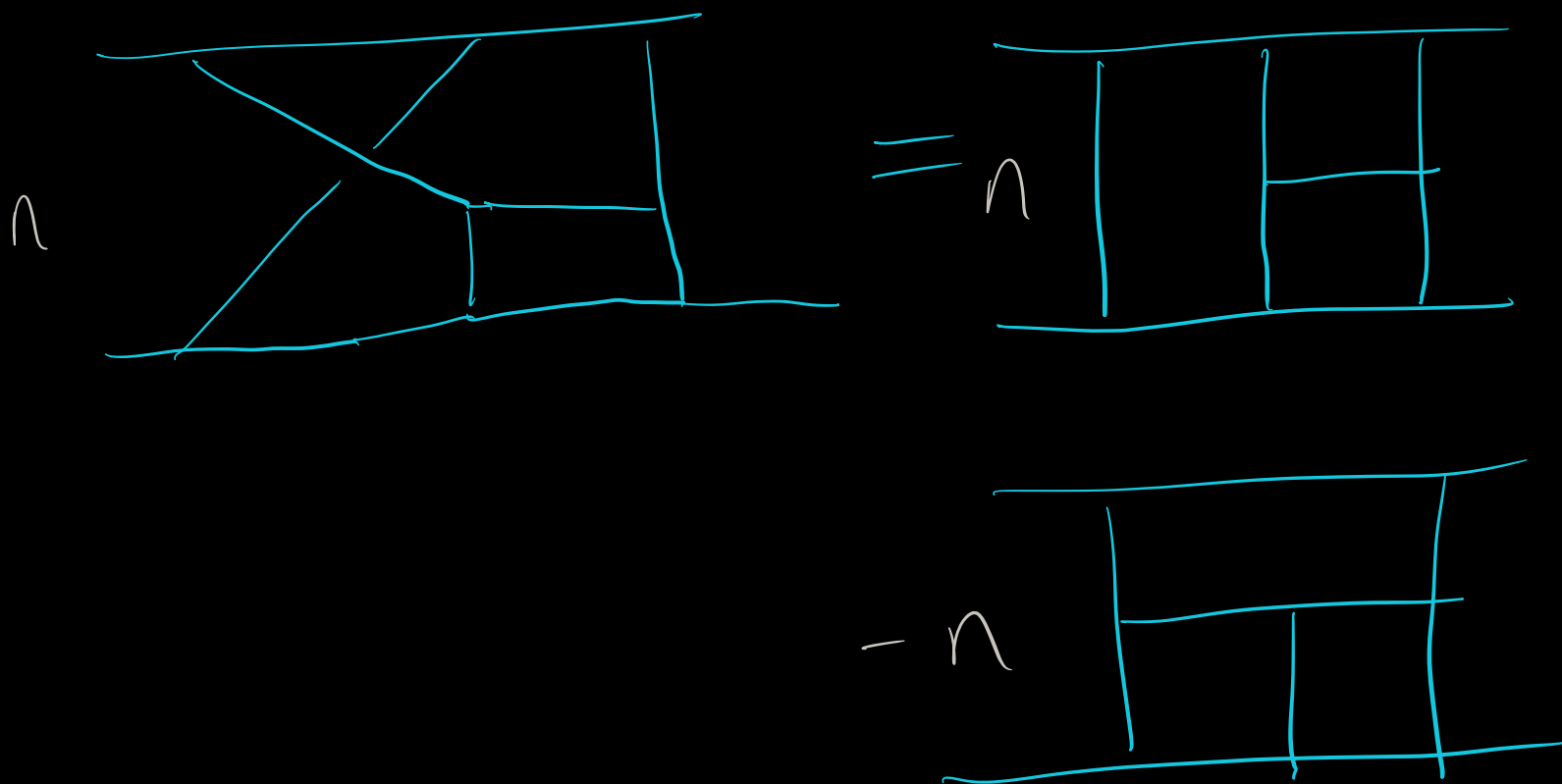
f^{abc}



$+ C$



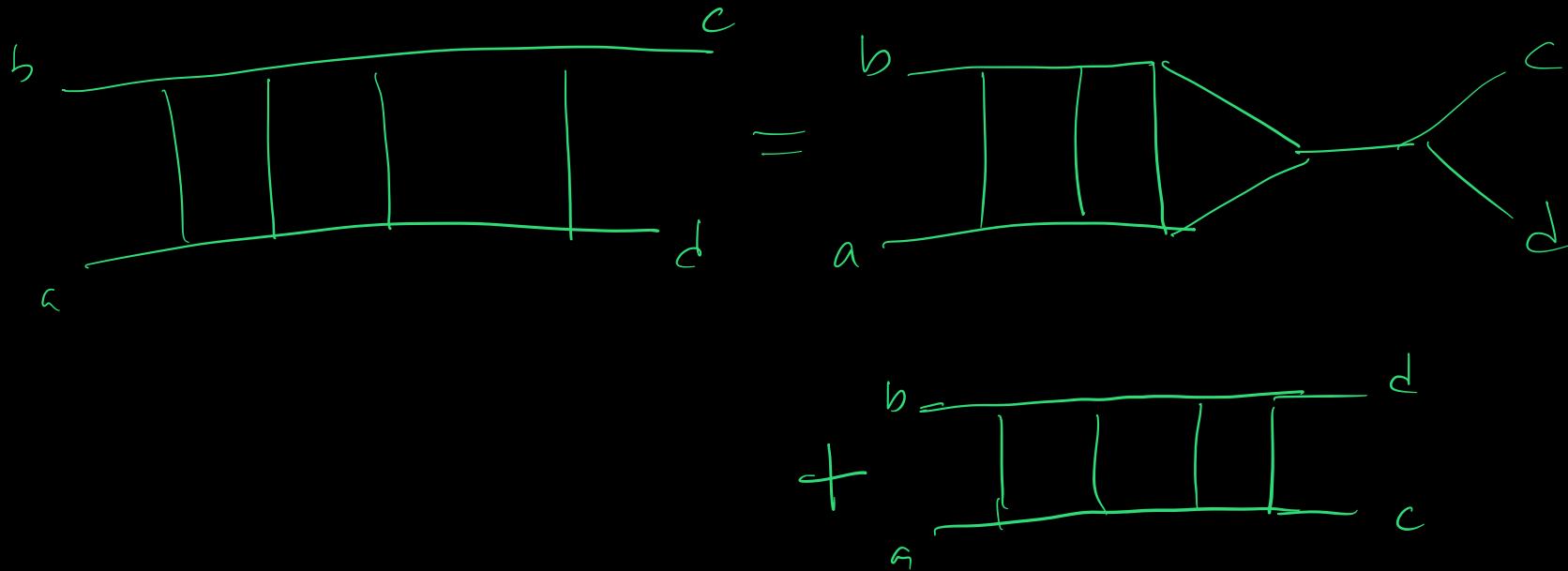
Demand this holds for n as well!



All of 3-loop 4pt $N=4$ SYM in one topology!

Bern, Carrasco, Johansson '08
'10

Note feed back constraints:



Jacobi relations are functional at loop level

so Ansätze are radical, (but can get too expensive)

Higher Derivative Operators

- $\mathcal{O}_i$ in $\mathcal{L} \longrightarrow$ h_i in $\mathcal{A}$
- Game is to figure out all distinct ways of sprinkling in some order in MD h into your scattering amplitude.

w/out messing w/ symmetries.

- Easy: scalar permutation invariants

4 pt tree level adjoint type $\eta(1-1)$



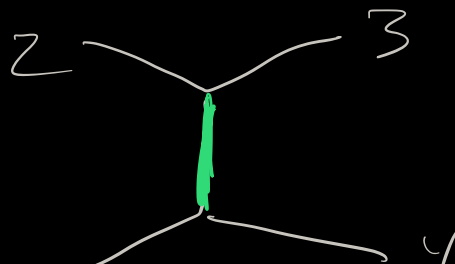
$$s = (k_1 + k_2)^2$$

$$= s_{12}$$

$$= s_{34}$$

$$g(s, t, u)$$

$$= -g(s, u, t)$$



$$t = (k_1 + k_4)^2$$

$$= s_{14} = s_{23}$$

$$g(t, s, u)$$

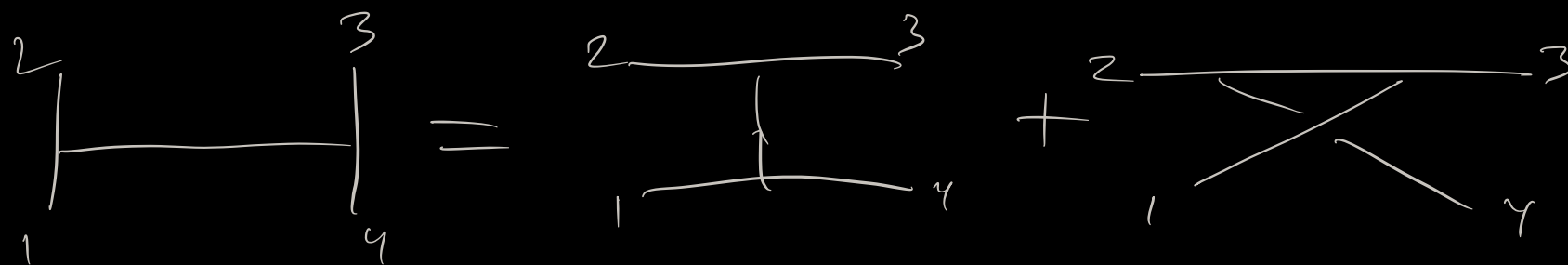


$$u = (k_1 + k_3)^2$$

$$= s_{13} = s_{24}$$

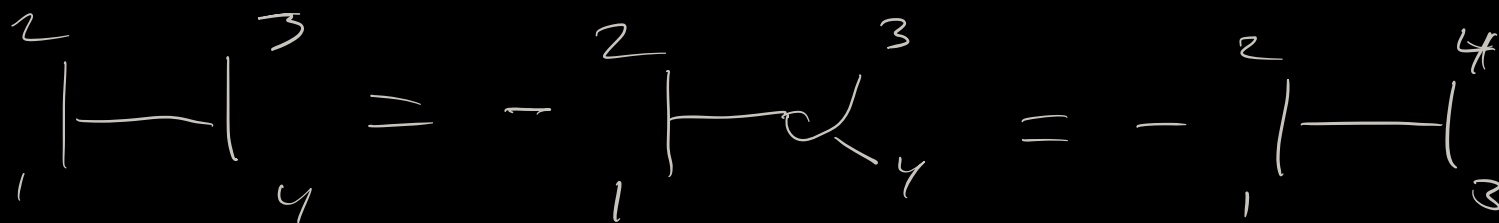
$$g(u, t, s)$$

Adjoint conditions



$$n_a(s, t, u) = n_a(t, s, u) + n_a(u, t, s)$$

ε



$$n_a(s, t, u) = -n_a(s, u, t)$$

Cov free scalar

$$(\downarrow \phi)^2 \sim$$

$$(\partial_m \phi_a) A_b^m \phi_c f^{abc}$$

~~sum~~

Adjoint type!

$$\begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 3 \\ | \\ 4 \end{array} = \begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 3 \\ | \\ 4 \end{array} + \begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 3 \\ | \\ 4 \end{array}$$

$$\underline{n_a(s,t,u)} = \underline{n_a(t,s,u)} + \underline{n_a(u,t,s)}$$

$$\begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 3 \\ | \\ 4 \end{array} = - \begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 3 \\ | \\ 4 \end{array} = - \begin{array}{c} 2 \\ | \\ 1 \end{array} \text{---} \begin{array}{c} 4 \\ | \\ 3 \end{array}$$

$$\underline{n_a(s,t,u)} = - \underline{n_a(s,u,t)}$$

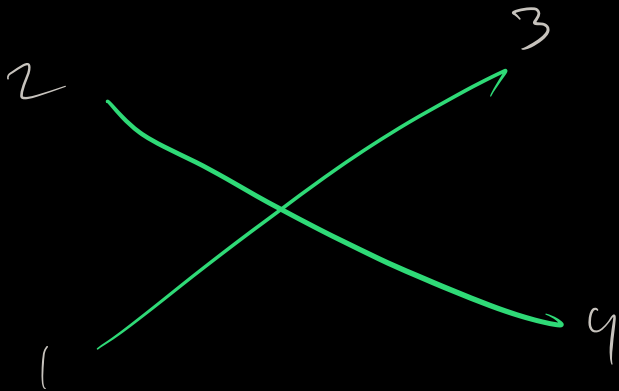
~~sum~~

$$n^{ss}(a,b,c) = b - c$$

$$\underline{(t-u)} = \underline{(s-u)} + \underline{(t-s)}$$

$$\underline{t-u} = - \underline{(u-t)}$$

Permutation Invt Conditions



$$s \notin A^{\text{ym}}(s, t)$$

P_4 doesn't care about leg order
invt under S_4 .

$$c \left(\begin{array}{cc} b & e \\ a & d \end{array} \right) = d^{abcd}$$

Note, no linear perm int for massless

$$s + t + u = 0$$

But can compose adjoint g_a to get
perm int.

$$N^{\text{adj}} \oplus \tilde{N}^{\text{adj}} = N_s \tilde{N}_s + N_t \tilde{N}_t + N_u \tilde{N}_u$$

$$P_2 = N^{ss} \oplus N^{ss} = s^2 + t^2 + u^2 = \sigma_2$$

Note also, trivially

$$N_s = \tilde{N}^{\text{adj}} \otimes_a P = P \otimes_a \tilde{N}^{\text{adj}} = P \tilde{N}_s^{\text{adj}}$$

Adjoint Comp at 4pt !

$$\eta_s = \tilde{\eta}^{\text{adj}} \textcircled{a}_s \check{\eta}^{\text{adj}} = \tilde{\eta}_t \check{\eta}_t - \tilde{\eta}_u \check{\eta}_u$$

$n_t, \dot{,} n_u$ just come functionally by relating.

$$\overline{\Pi}_s = n^{ss} \textcircled{a}_s n^{ss} \propto s(t-u)$$

$$P_3 = \underbrace{n^{ss}}_1 \textcircled{\Phi} \underbrace{\overline{\Pi}}_2 \propto st u = \mathcal{O}_3$$

$$P_4 = \overline{\Pi} \textcircled{\Phi} \overline{\Pi} \propto \mathcal{O}_2 \textcircled{\Phi} \mathcal{O}_2 = (s^2 + t^2 + u^2)^2$$

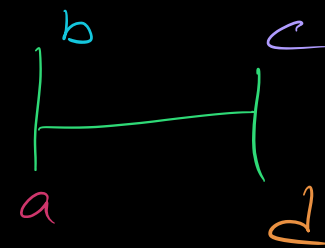
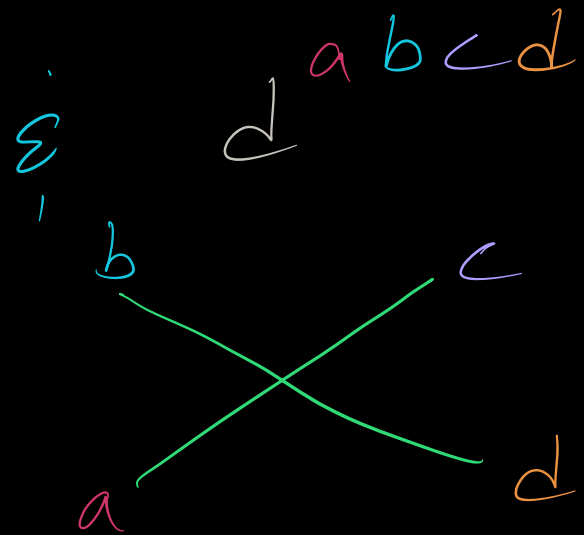
Simple Covariantized Scalar

spans all scalar
adjoint type BB
under composition.

- Climbs a ladder MD by MD
- Closes after 2nd rung under perin invariance.

$$n_s^{MD} = \sum_{X,Y}^{3Y+2X=MD-1} C_{XY}^{ss} G_2^X G_3^Y n_s^{ss} + \sum_{X,Y}^{3Y+2X=MD-2} C_{XY}^{\pi} G_2^X G_3^Y \pi_s$$

Can add color: $c_s^{ads} = f^{ab} e f^{cd}$



to $n's$ using
composition in
relevant ways.

eg:

$$\overline{\Pi}_s^{abcd} = d^{abcd} \overline{\Pi}_s$$

Note: $\text{tr}(F^4) = d^{abcd} \text{st} A^{ym}(1234)$

$= \prod^{abcd} \otimes \text{vector}$

It's ordered amplitudes don't satisfy

$(n-3)!$ BCJ amplitudes relations

(they're perm inv!)

Rather it is a double copy

C. S.
Shruti
Paranjape's
TALK!

This allows us to recover
Open & Closed Superstrings
at Tree level as

FIELD THEORY Double Copies

from scalar perm invariants ($\underline{g_2}, \underline{g_3}$)
scalar adjoint numerators ($\underline{n^{ss}}, \underline{\Pi}$)
single trace color weights ($\underline{f^{abc}}, \underline{d^{abcd}}$)
& sym numerator ($\underline{n^{sym}}$)

$$OSS = \underline{Z} \otimes sYM$$

$$CSS = sYM \otimes sYM^{(SV)}$$

$$\begin{aligned} \underline{Z}_s = & \left(\underline{C^{adj}} \otimes n^{ss} \right)_s 6_2^x 6_3^y C_{x,y}^{ss} \\ & + \left(\underline{C^{adj}} \otimes \Pi \right)_s 6_2^x 6_3^y C_{x,y}^{\Pi} \\ & + \left(\underline{d^{abcd}} \Pi \right)_s 6_2^x 6_3^y C_{x,y}^d \end{aligned}$$

$$sYM_s^{(SV)} = n_s^{YM} 6_2^x 6_3^y C_{x,y}^{SV}$$

Composition Very Powerful

• Could verify up to insanely high MD
 $(\alpha')^{40} \sim (s_{ij})^{40}$ before finding the right way

→ express this in terms of OSS

directly in closed $\Gamma(\)$ form at 4pt.

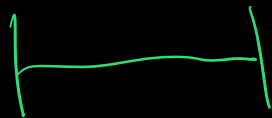
• Check out PRL for example

$\hat{\epsilon}_i$ proofs

Why no $(C^{SS} @ n^{YM})$ or
 $(C^{\pi} @ n^{YM})$?

Resulting numerators satisfy adjoint
constraints but NOT Gauge Inv.

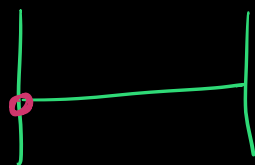
What about eg other tensor structures?
turns out $\exists$ 8 spanning building blocks
at 4pts.



γ_m

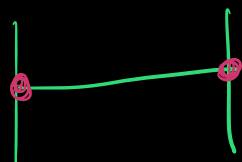
Vector Building
Blocks

$\text{tr} F^3$



α' above γ_m + 2 more

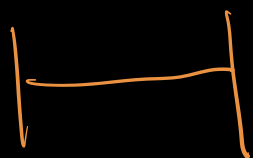
$(\text{tr} F^2)^2 / G_3$



$2 \text{tr} F^3 + \text{tr}(F^4)$

α'^2 above γ_m

+ 2 more

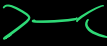


$\partial^2 F^4$

All other local color dual vectors spanned
by these ϵ_i perm invariants.

Eg Bosonic String: $Z \oplus (DF)^2 + YM$

$$\frac{(DF)^2 + YM}{N} = N_{YM} + \frac{\left[\alpha' (n^F)^3 + \alpha' (n^{2F^3 + F^4}) + \alpha' (n^{\partial^2 F^4}) + \alpha' \hbar (DF)^4 \right]}{1 - \alpha'^2 G_2 - \alpha'^3 G_3}$$

Diagrams above the terms in the numerator:
    

Fine

But can these be algebraically constructed?

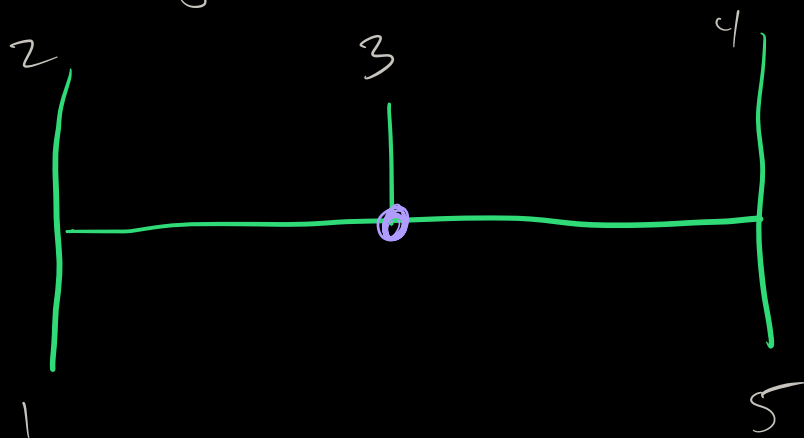
= STAY TUNED = JMC,
LR, SŽ

Is this just a feature of 4pt?

No! 5pt even richer:

a fascinating surprise

No adjoint linear building block

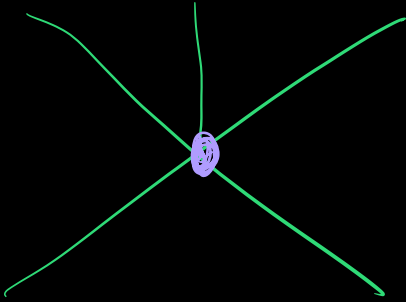


Relaxed Instead

$$(h_1 - h_2) \cdot (h_5 - h_4)$$

JJMC, LR, SŽ
2104.08370

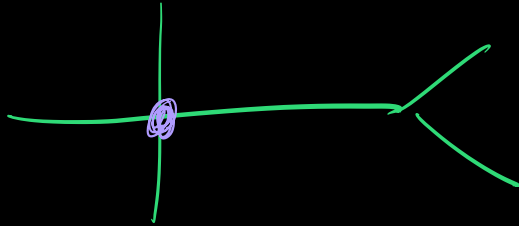
5 pt Algebraic Structures



P_5

d^5

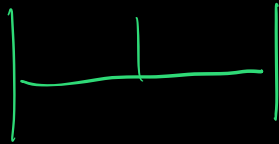
$\textcircled{P}$



Hybrid

$d^4 f^3$

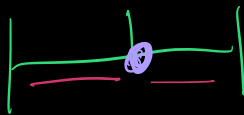
$\textcircled{h}$



Adj

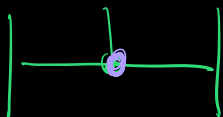
$f^3 f^3 f^3$

$\textcircled{a}$



Relaxed Adj

$\textcircled{r}$



Sandwich

$f^3 d^3 f^3$

$\textcircled{s}$

Lesson from five points.

- Odd multiplicity EFT

linear building block = Relaxed

- Even multiplicity = Adjoint



of vector adjoint color-dual BB?

Open Question

- even SpT exhausting via
ansatz, motivating constructive
approach!

STAY TUNED!

Color-dual Fate of $N=4$ SG

JJMC, ML, NP 2203.03592

- Diverges in 4D at 4-loops

Bern, Davies, Dixon
Smirnov, Smirnov '13

- Seems tied to anomalous behavior
at 1-loop

JJMC, Kallosh
Roiban, Tseytlin '13

- Counterterm related to $\text{tr } F^3$

— does resolve 1-loop anomaly behavior.

— doesn't mess up UV at 1 & 2 loops.

'17 Bern, Edison, Kosower, Parra-Martinez

'17 Bern, Parra-Martinez, Roiban

'19 Bern, Kosower,
Parra-Martinez

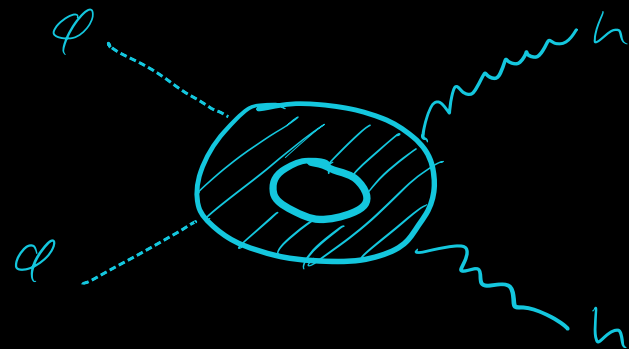
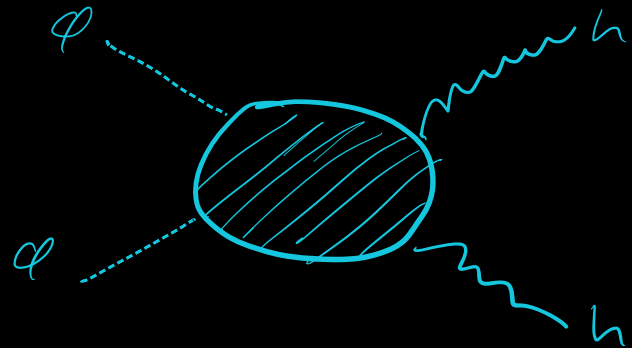
ANOMALOUS BEHAVIOR

Marcus '85

$$A_{\text{tree}}^{SG}(\varphi\varphi h^+h^+) = 0$$

$N=4$

$$A_{\text{loop}}^{SG}(\varphi\varphi h^+h^+) \neq 0$$



Can understand behavior via double-copy

half max
SG = maximal SYM $\otimes$ pure YM

ANOMALOUS BEHAVIOR

$$A^{sb}(\varphi\varphi h^+h^+) =$$

$$g^+ \otimes g^+ = h^+$$

$$g^+ \otimes g^- = \varphi$$

$$g^- \otimes g^- = h^-$$

$$\underbrace{A^{SYM}(g^-g^-g^+g^+)}_{\text{B/C SUSY: ONLY VALID HEL STRUCTURE TREE OR LOOP}} \otimes \underbrace{A^{YM}(g^+g^+g^+g^+)}_{\text{ZERO at tree level (secretly SUSY at tree) NON-ZERO AT 1 LOOP!}}$$

B/C SUSY: ONLY
VALID HEL STRUCTURE
TREE OR LOOP

ZERO at tree level
(secretly SUSY
at tree)
NON-ZERO AT 1 LOOP!

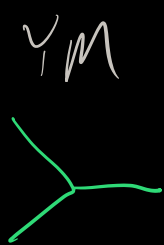
To cancel anomaly need COUNTERTERM

that admits 4 g^+ even at tree level $\underbrace{\text{tr}(F^3)}$

So is $\text{tr} F^2 + \alpha' \text{tr} F^3$

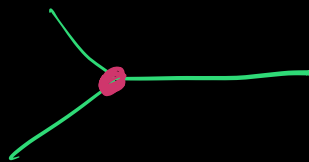
Double Copy Consistent?

(ie color dual $\hat{=}$ factorizing for
all multiplicity at tree level.)



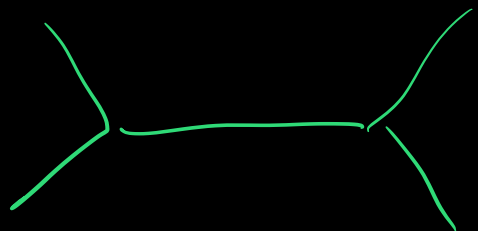
+ α'

F^3

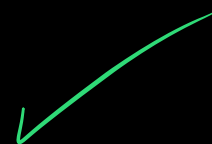


3pt sure!

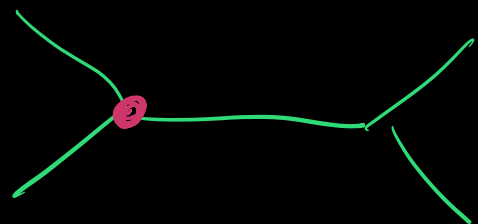
4pt



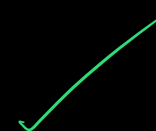
(YM)



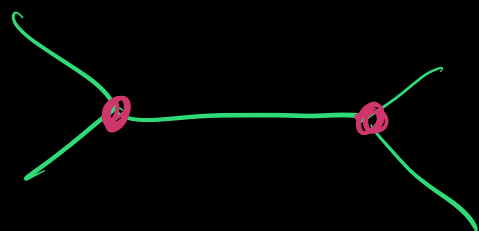
α'



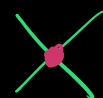
(single insertion)



α'^2



(double insertion)



Needs 4pt $g(F^4)$ et. OK

So is $\text{tr} F^2 + \alpha' \text{tr} F^3 + \dots$

Double Copy Consistent

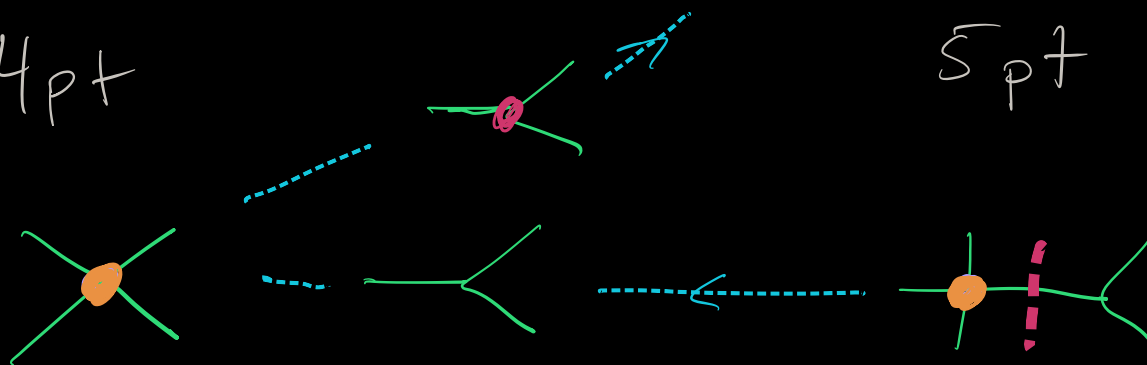
with a finite # of CT?

Apparently not

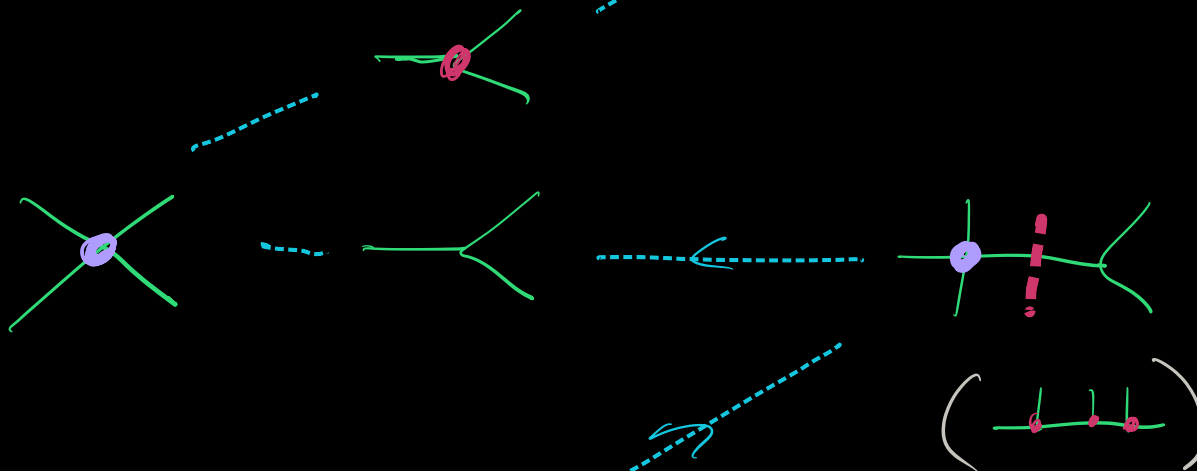
4pt

5pt

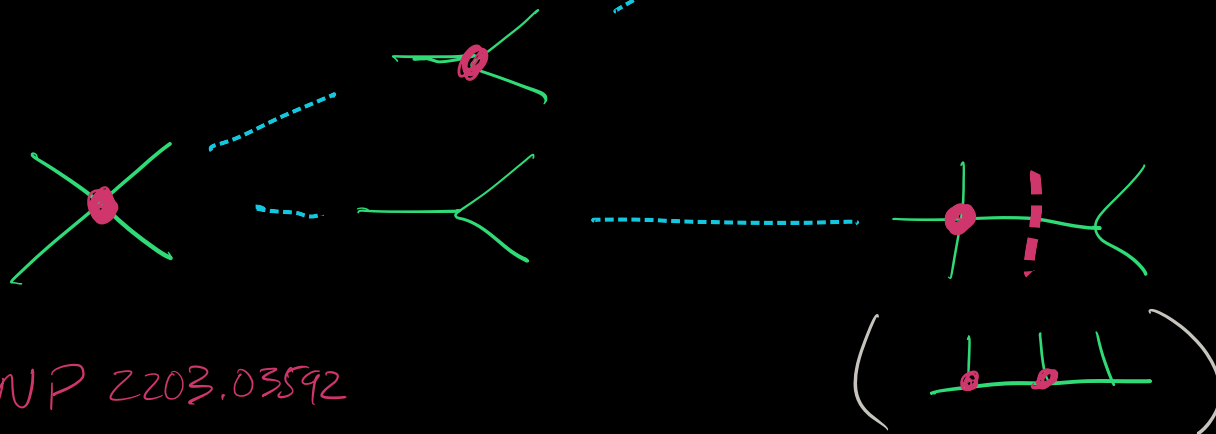
$(\alpha')^4$



$(\alpha')^3$



$(\alpha')^2$



JJMC, ML, NP 2203.03592

Requiring $C/K \in$

Factorization even

Between 4 $\in$ 5 point

INDUCES A TOWER

ALL THE WAY TO THE

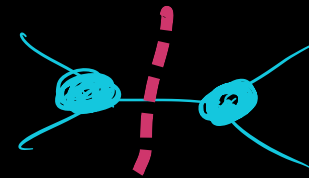
UV

How can you tell?

4pt Ansatz:

$$\eta^{4pt} = \sum_{i \in I}^8 \sum_{x,y}^{2x+3y < \max} C_{xy}^i G_2^x G_3^y \eta_i$$

fix on cuts to 3 points

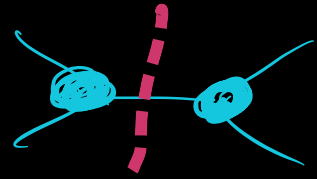


Everything is either

precluded, set to α' or 1 (YM coupling),

or free (could be set to $\emptyset$)

$$\mathcal{L} = \mathcal{L}^{\text{YM}} + \alpha' \mathcal{L}^{\text{YM}+F^3} + \alpha'^2 \mathcal{L}^{(F^3)^2+F^4} +$$



$$\alpha'^3 \left(\mathcal{C}_3 \left(\mathcal{L}^{\text{D}^2 F^4} + \mathcal{G}_2 \mathcal{L}^{\text{YM}+F^3} \right) + \right.$$

$$\left. \tilde{\mathcal{C}}_3 \left(\mathcal{G}_3 \mathcal{L}^{\text{YM}} \right) \right) +$$

$$\alpha'^4 \left[\mathcal{C}_4 \left(\mathcal{L}^{(\text{D}^2 F)^2} \right) + \mathcal{G}_2 \mathcal{L}^{(F^3)^2+F^4} \right. \\ \left. + \check{\mathcal{C}}_4 \mathcal{L}^{(\text{D}^2 F^2)^2} + \tilde{\mathcal{C}}_4 \mathcal{C}_3 \mathcal{L}^{\text{YM}+F^3} \right]$$

$\mathcal{C}$ unconstrained - could be zero.

UNLESS ladder forces particular values.

which it does.

(feedback on
5 pt c/k i
factorization)

$$C_3 = 1$$

$$\tilde{C}_4 = 1 + \tilde{C}_3$$

$$C_4 = \tilde{C}_4 = 1$$

$\Rightarrow$

$$n = n^{(DF)^2 + 4m} + \mathcal{O}(\alpha'^5) \\ + \tilde{C}_3 \alpha'^3 G_3 (n^{4m} + \alpha' n^{F^3} + \dots)$$

$$N=4 \text{ SG} + \text{c.t.} = N=4 \text{ SYM} \oplus \left[\begin{array}{l} \text{YM} + \text{tr}(F^3) \\ + \dots \end{array} \right] \equiv$$

$$\left[\begin{array}{l} \text{① } F^2 + \text{YM} \\ + \dots \end{array} \right]$$

= Berkovitz Witten SG

of Heterotic String

Calculated through α'^4

Proof & Loop level results await!

(A few words about DF^2)

$$DF^2 \oplus SYM = CSG$$

Johanson, Nohle
117

$$(DF^2 + YM) \oplus SYM = \text{Weyl-Einstein}$$

Johansson, Mogull,
Teng '18

$$(DF^2 + YM) \oplus \mathbb{Z} = \text{Bosonic String}$$

$$(DF^2 + YM)^{SV} \oplus SYM = \text{gravitons in heterotic string}$$

$$(X)_c^{SV} = X_a \oplus^{ab} SV(\mathbb{Z}_{bc})$$

Azvedo, Chiodaroli,
Johansson, Schlotterer
118

Next Steps

- Composition to build vector BB?
- Proof that $\text{tr} F^3$ closes to $(\text{DF})^2 + \text{YM}$
- Effect of tower on explicit
loop-level calculation in $N=4 \text{ SG} + \dots$
- Loop-level composition?
- Lots to Do!

Summary

- composition lets us build & classify EFT operators for gauge & gravity theories with SMALL set of Building Blocks
- $N=4$ SG may have a non-local color-dual fate :
 - Berkovits-Witten Conformal SG
 - Heterotic String