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# Effective Field Theory for Halo Nuclei: Lecture 3

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Ohio University



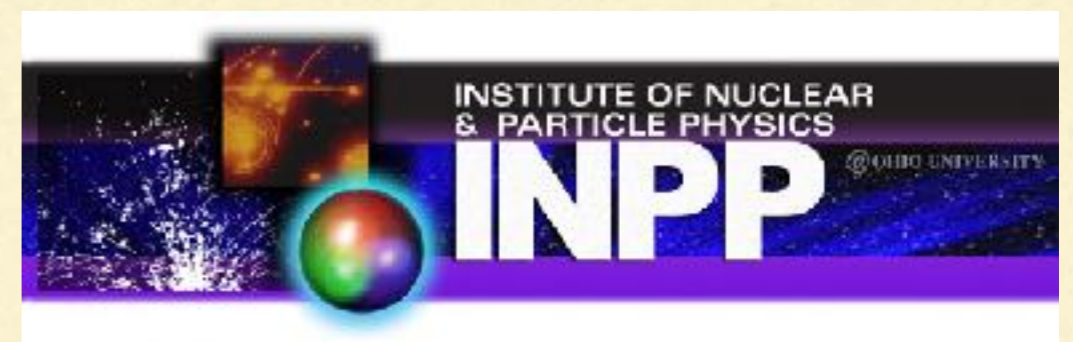
OHIO  
UNIVERSITY

Collaborators: Bijaya Acharya

Pierre Capel

Hans-Werner Hammer

Chen Ji



**RESEARCH SUPPORTED BY THE US DEPARTMENT OF ENERGY**

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# The story so far

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- Halo EFT seeks to describe halo nuclei for  $k \sim 1/R_{\text{halo}} \ll 1/R_{\text{core}}$ , i.e.,  $kr_0 \ll 1$ .
- We are interested in shallow bound states (or virtual bound states), states for which  $\gamma r_0 \ll 1 \Leftrightarrow r_0/a_0 \ll 1$

$$t_0^{NLO}(E) = -\frac{2\pi}{m_R} \frac{1}{-\frac{1}{a_0} + \frac{1}{2}r_0k^2 - ik}$$

- Predicts a shallow bound state with co-ordinate space wave

function  $\psi(\mathbf{r}) = \frac{A}{\sqrt{4\pi}} \frac{e^{-\gamma r}}{r}$

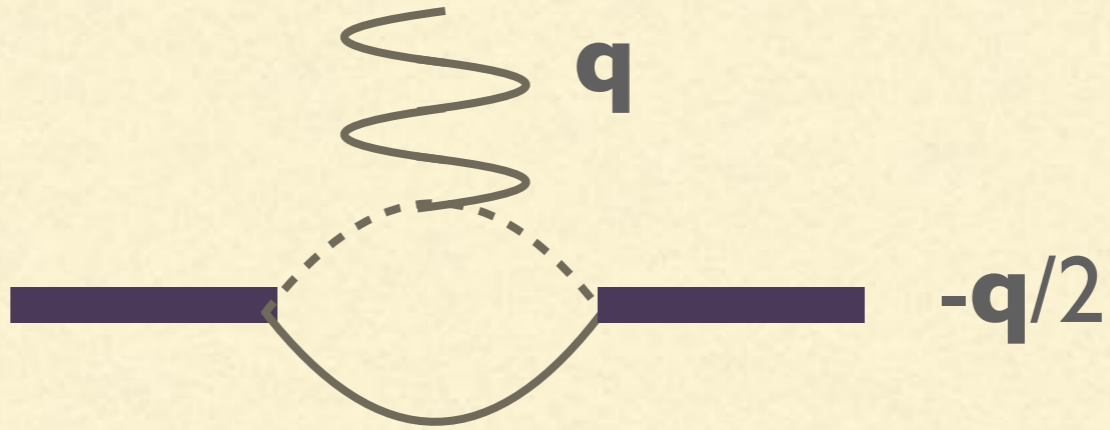
- $A^2 = \frac{2\gamma}{1 - \gamma r}$
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# Electric form factor

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$$-i|e|Z_c G_c^{NLO}(|\mathbf{q}|) =$$

 $\mathbf{q}/2$ 

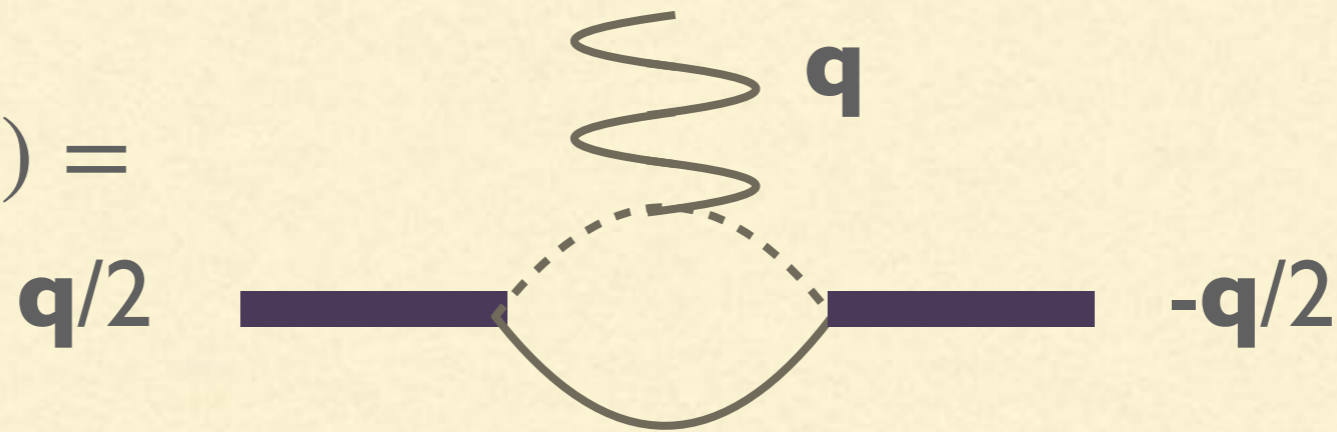


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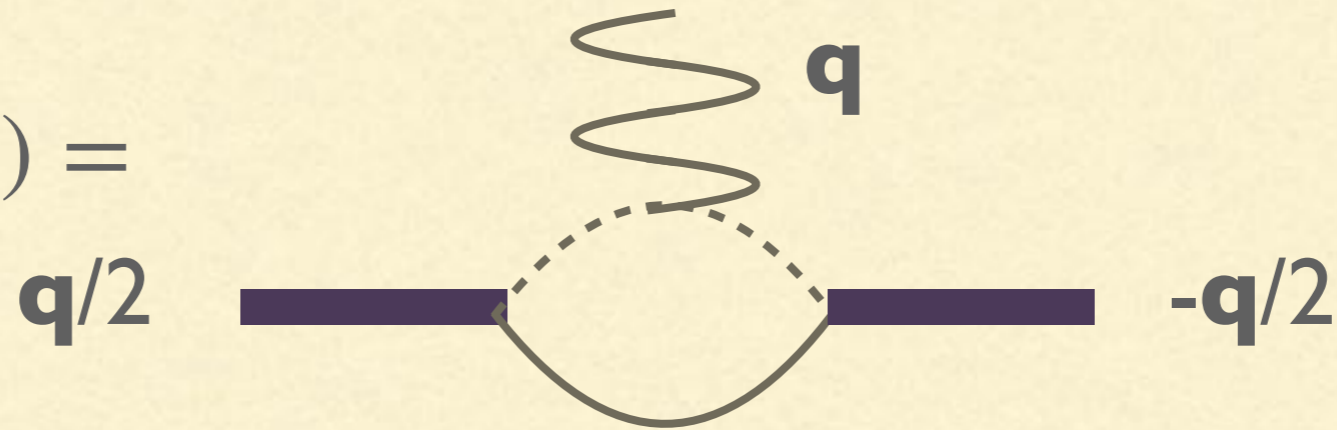
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$$G_c(|\mathbf{q}|) = \frac{A_0^2}{f|\mathbf{q}|} \arctan\left(\frac{f|\mathbf{q}|}{2\gamma}\right)$$

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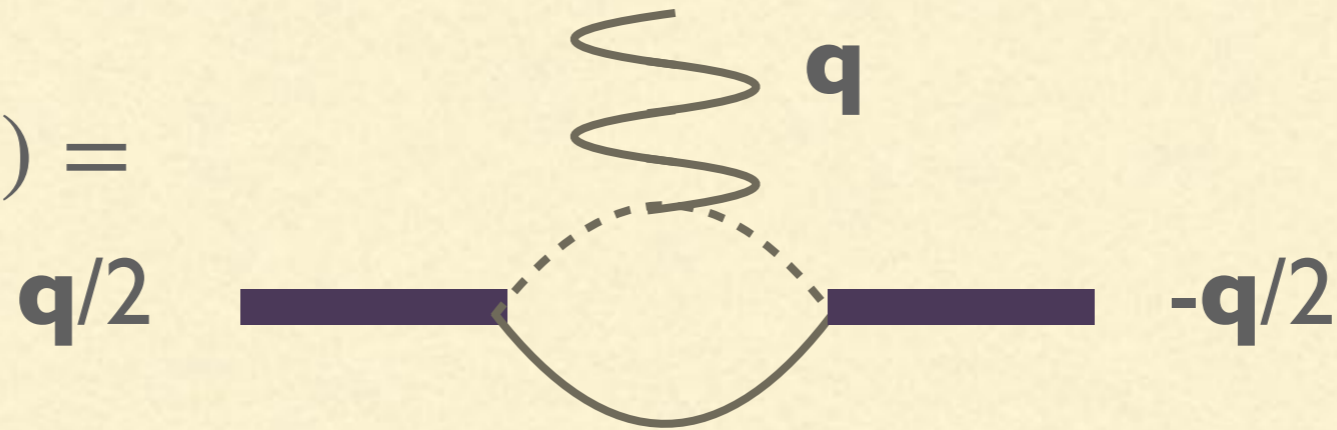


$$f = \frac{m_R}{M_c} = \frac{1}{A_c + 1}$$

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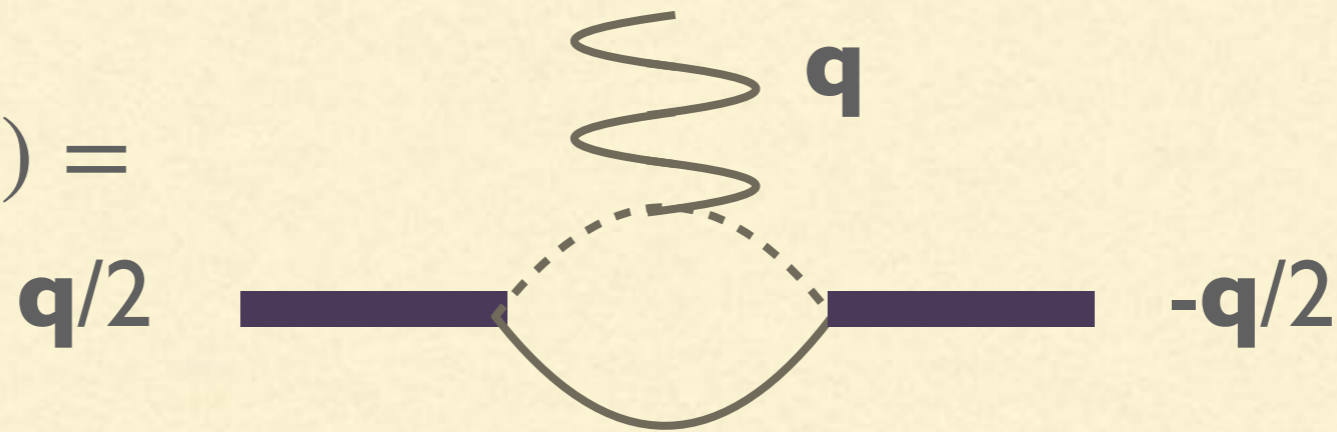
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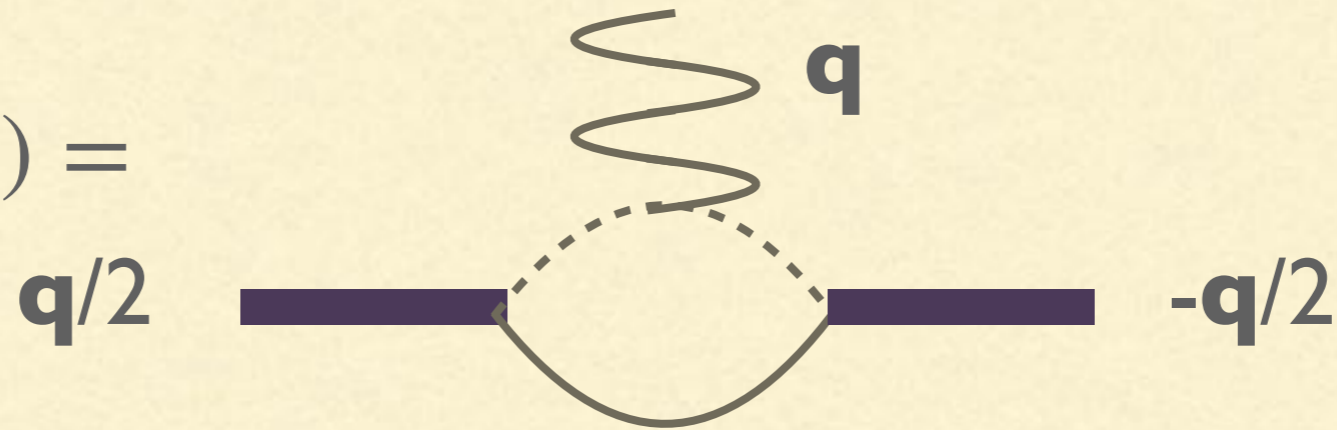
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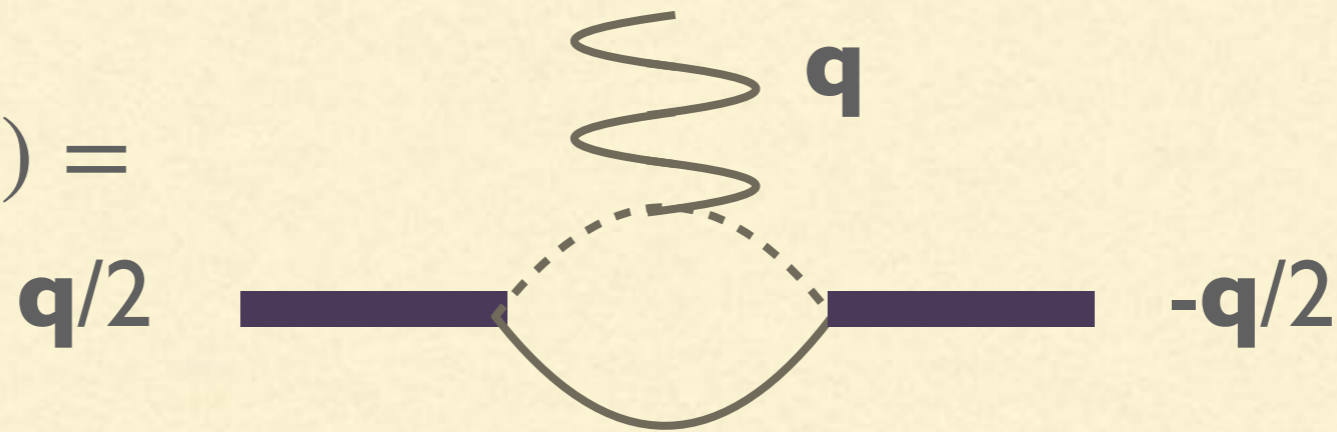
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$$\langle r_E^2 \rangle_{\text{nucleus}} = \frac{f^2 A_0^2}{4\gamma^3} + \langle r_E^2 \rangle_c + \frac{1}{Z_c} \langle r_n^2 \rangle$$

Cf. Ryberg et al.,  
EPJA 56, 7 (2020)







# Electric form factor

$$-i |e| Z_c G_c^{NLO}(|\mathbf{q}|) = \begin{array}{c} \mathbf{q} \\ \text{wavy line} \\ \text{---} \mathbf{q}/2 \quad \text{---} -\mathbf{q}/2 \end{array} + \begin{array}{c} \text{wavy line} \\ \text{---} \end{array}$$

$$G_c(|\mathbf{q}|) = \frac{A_0^2}{f|\mathbf{q}|} \arctan\left(\frac{f|\mathbf{q}|}{2\gamma}\right) + 1 - \frac{A_0^2}{2\gamma}$$

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LO:

$$A_0 = \sqrt{2\gamma}$$

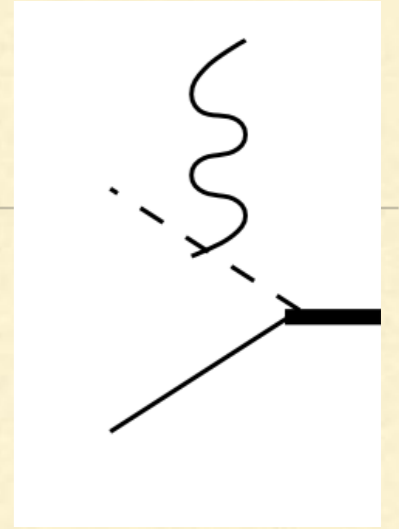
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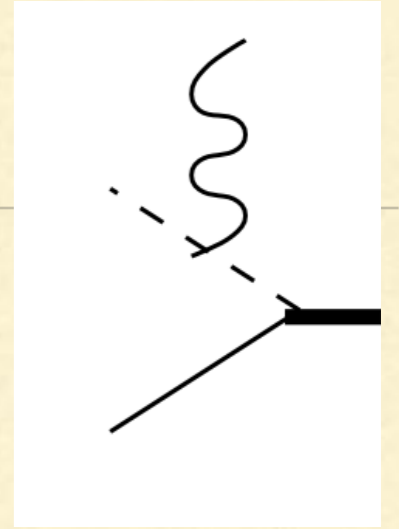
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Chen, Savage (1999)

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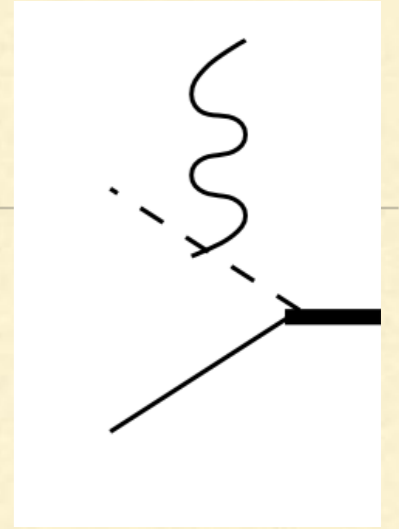
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# E1 photodissociation

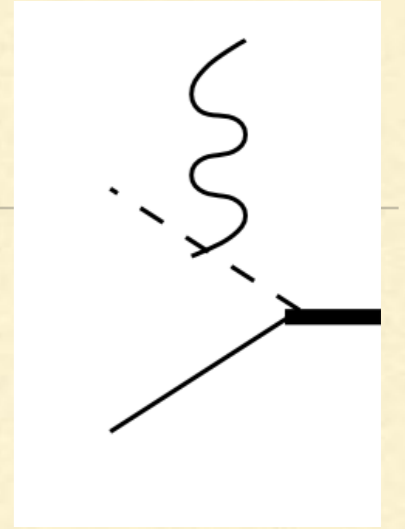


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1. Choose photon momentum to be aligned with z-axis
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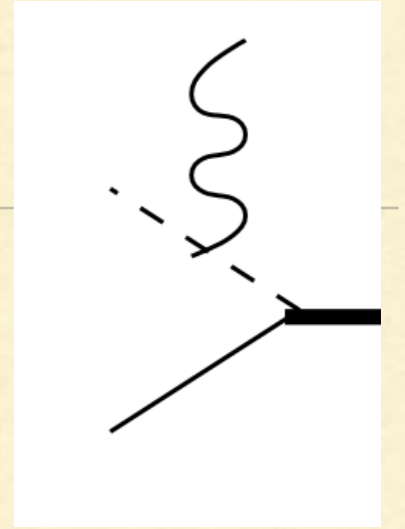
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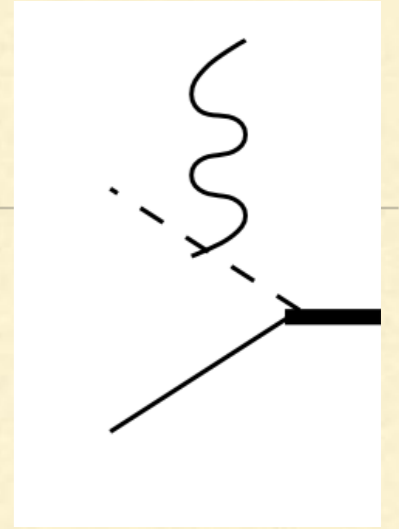
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$$\frac{dB(E1)}{dE} = \sum_{m_l} |\mathcal{M}_{E1}^{(l=1, m_l)}|^2 \frac{d^3k}{(2\pi)^3}$$



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# Universal E1 strength formula

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$$\frac{dB(E1)}{\alpha_{em}dE} = \frac{6m_R}{\pi^2} \frac{Z_c^2}{(A_c + 1)^2} A_0^2 \frac{k^3}{(\gamma^2 + k^2)^4}$$

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$$m_R S_{1n}^2 \frac{dB(E1)}{dE} = \frac{A_0^2}{2\gamma} \frac{3\alpha_{em} Q_{eff}^2}{\pi^2} \frac{x^{3/2}}{(1+x^2)^2}; \quad Q_{eff} = fZ_c = \frac{Z_c}{A+1}; \quad x = E/S_{1n}$$

- Universal E1 strength formula for S-wave halos

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- Short-distance piece of E1 m.e.:  $L_{E1} \left[ d^\dagger \mathbf{E} \cdot (n \overleftrightarrow{\nabla} c) + h.c. \right] \sim \left( \frac{R_{core}}{R_{halo}} \right)^4$

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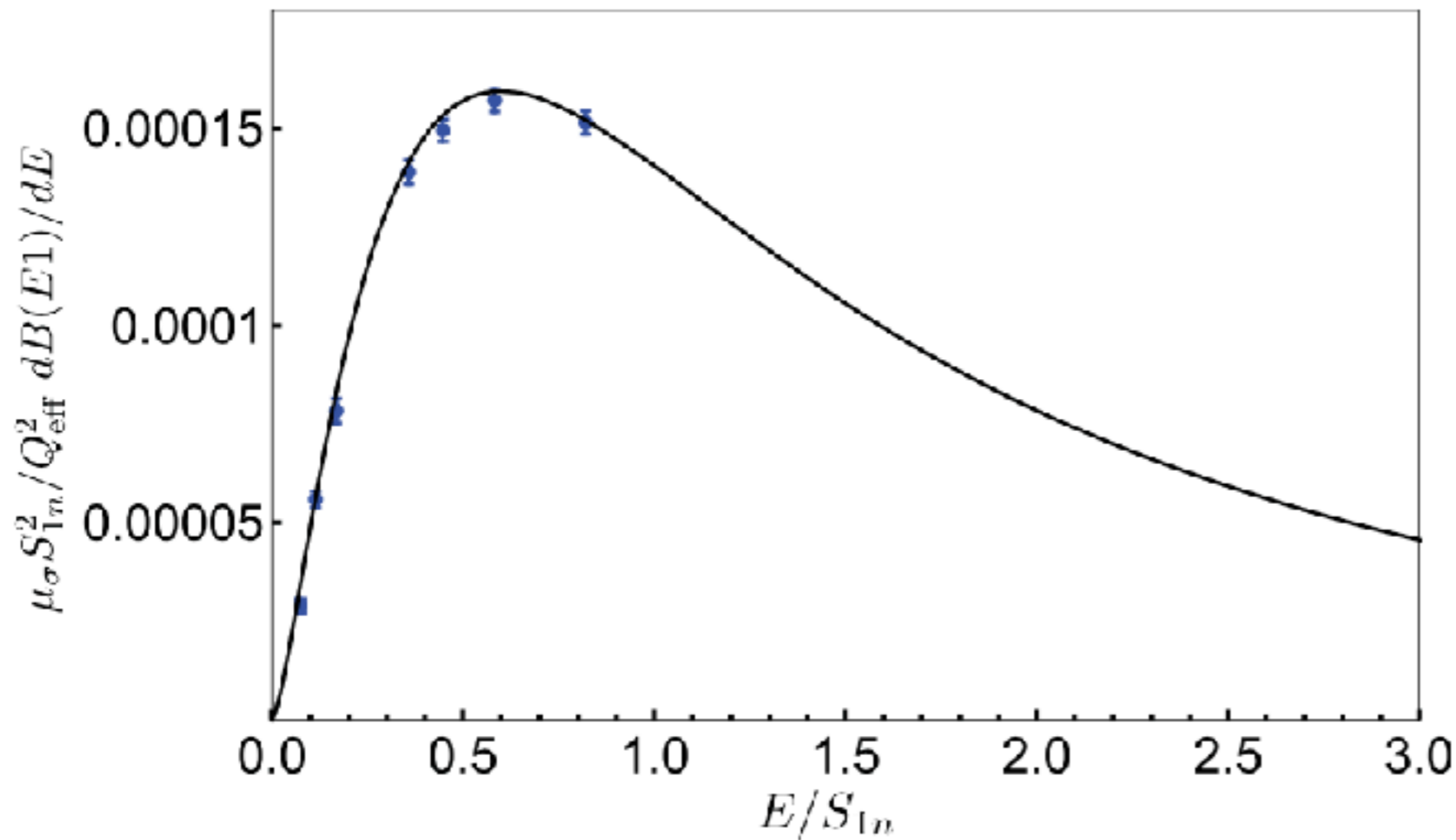
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- Tells you how universal formula “breaks”

# Results



Deuteron

Data: Tornow et al. (2003)  
Analysis: Chen, Savage (1999);  
Hammer, Ji, DP (2017)

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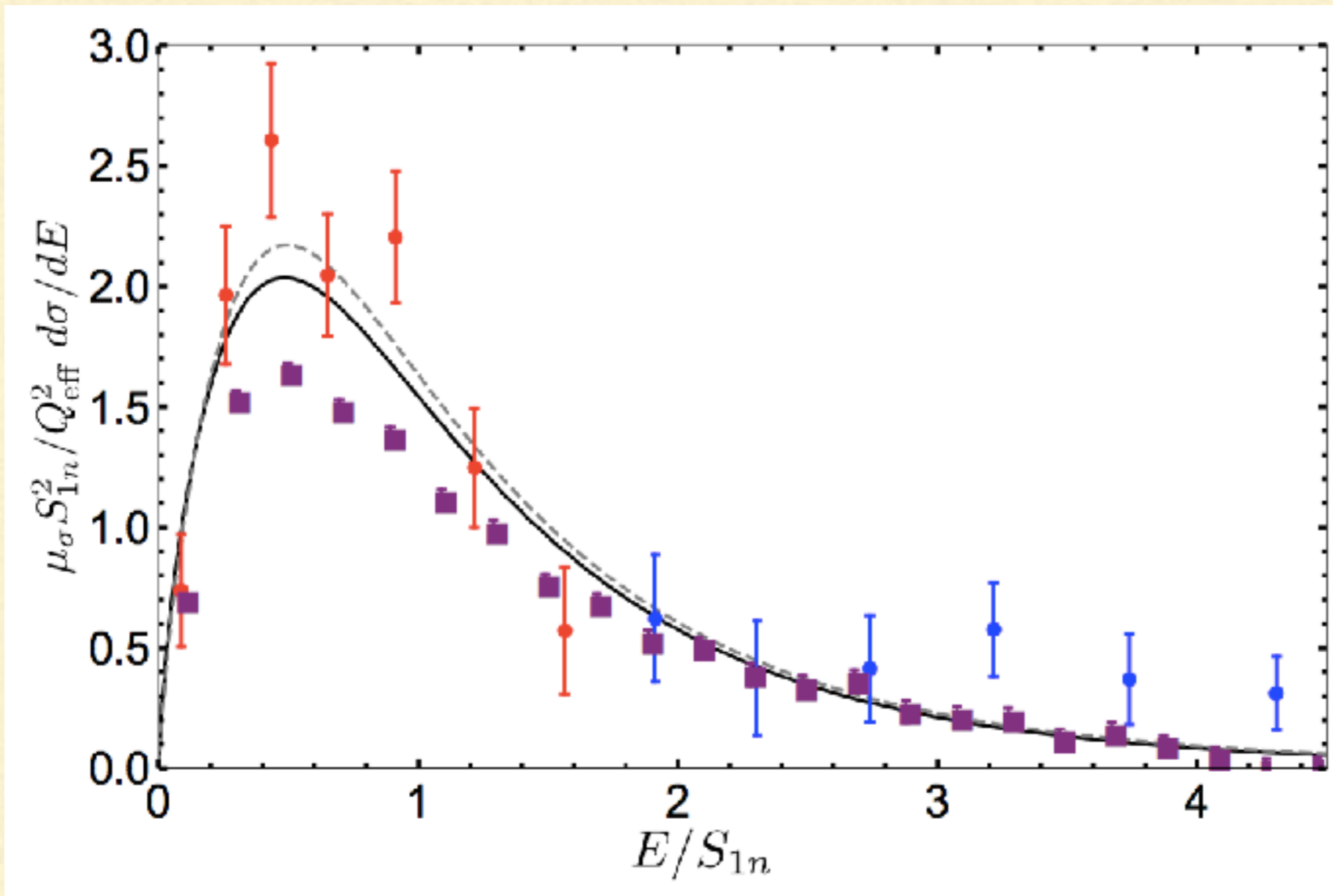
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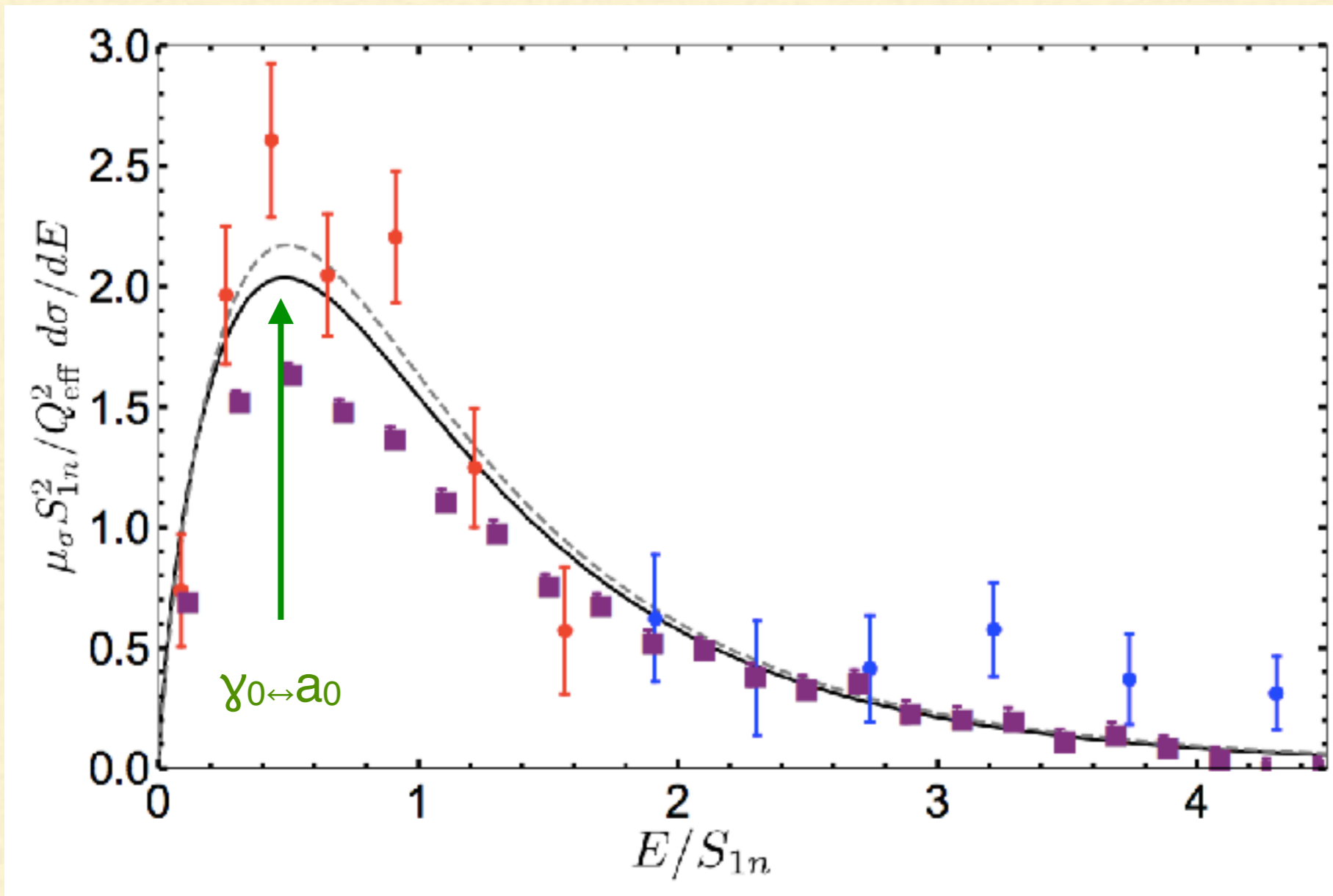


$^{19}\text{C}$

Data: Nakamura et al. (1999, 2003);  
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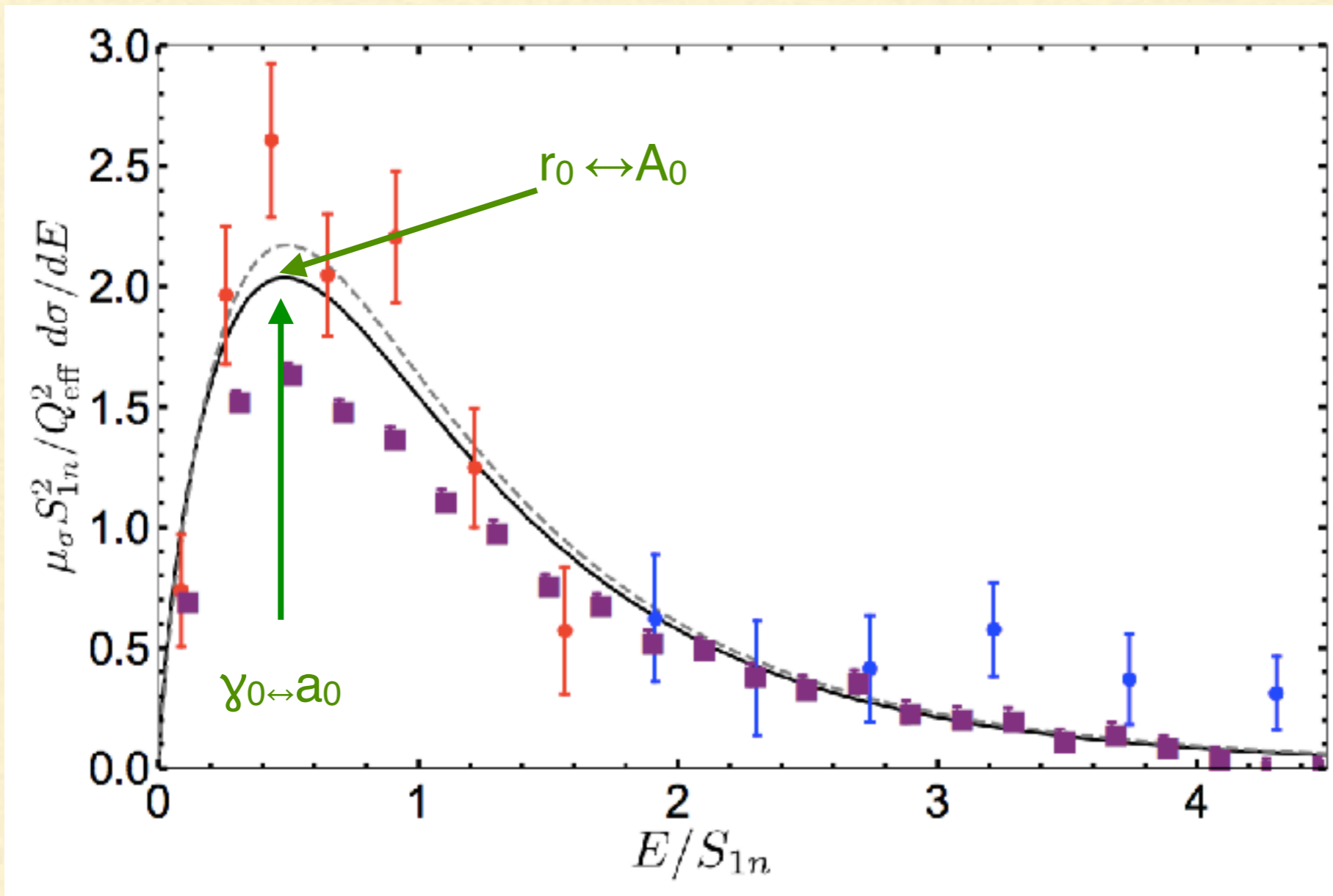
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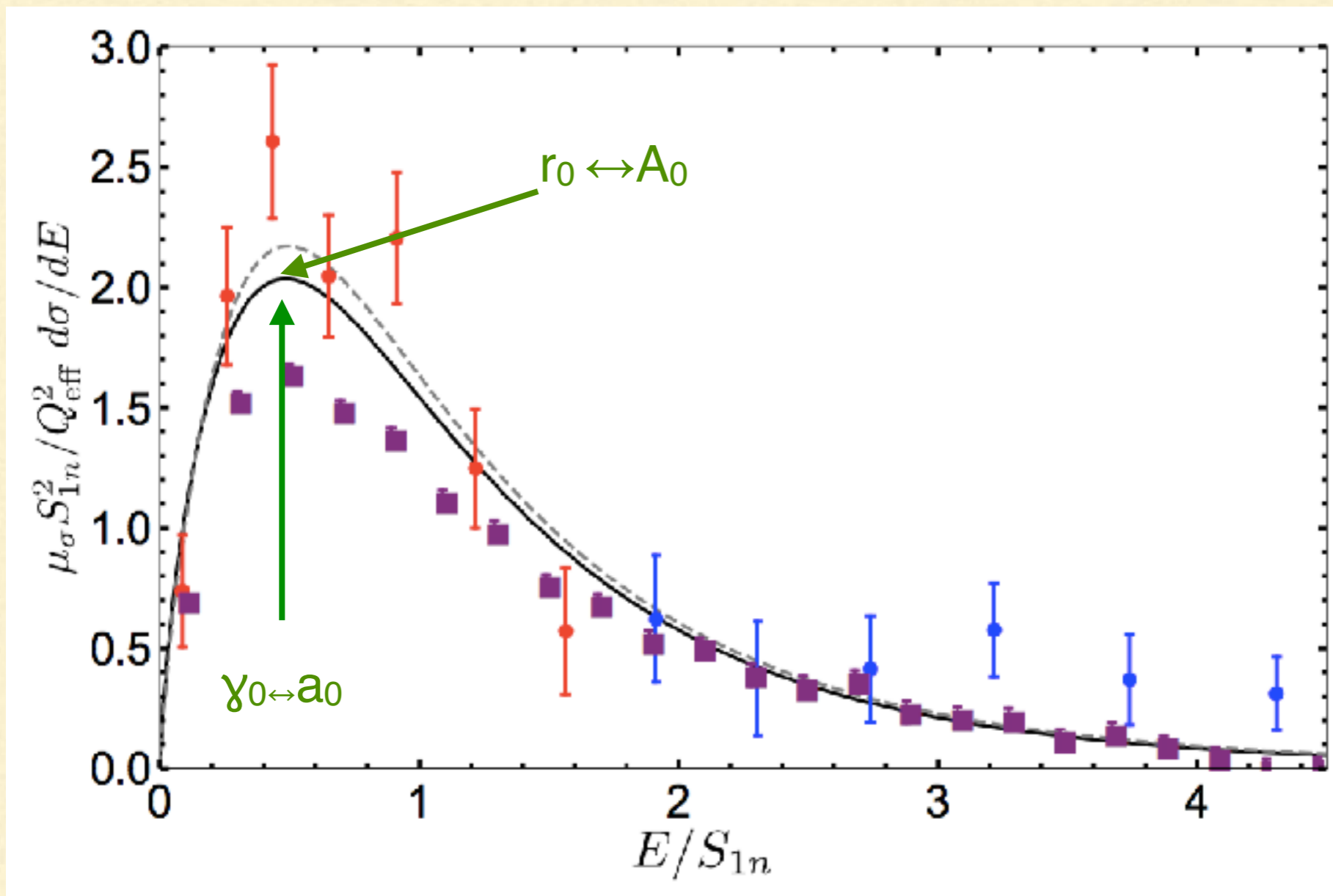


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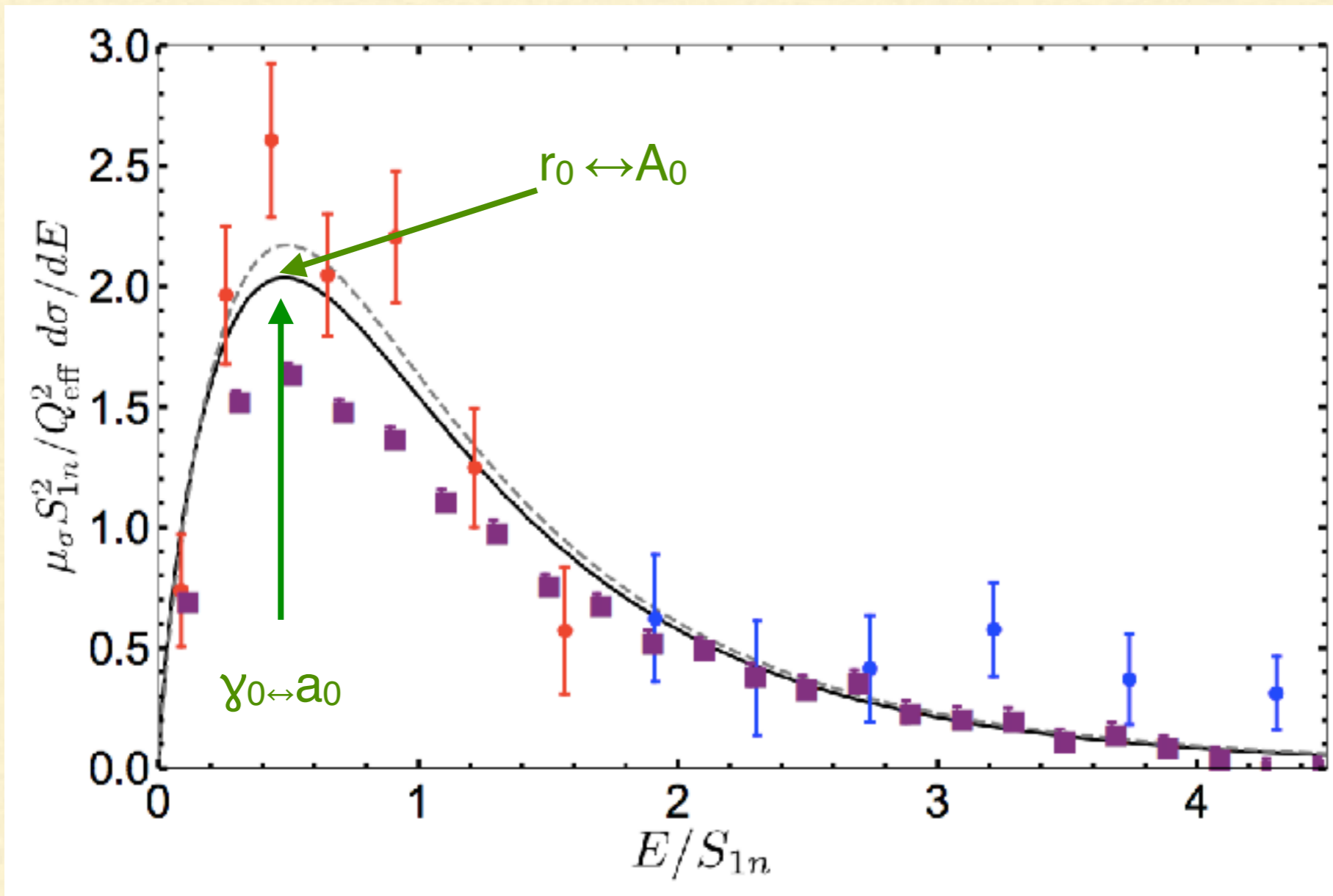


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Determine S-wave  $^{18}\text{C}$ -n scattering parameters  $\Leftrightarrow$   $^{19}\text{C}$  ANC from dissociation data.

# Results



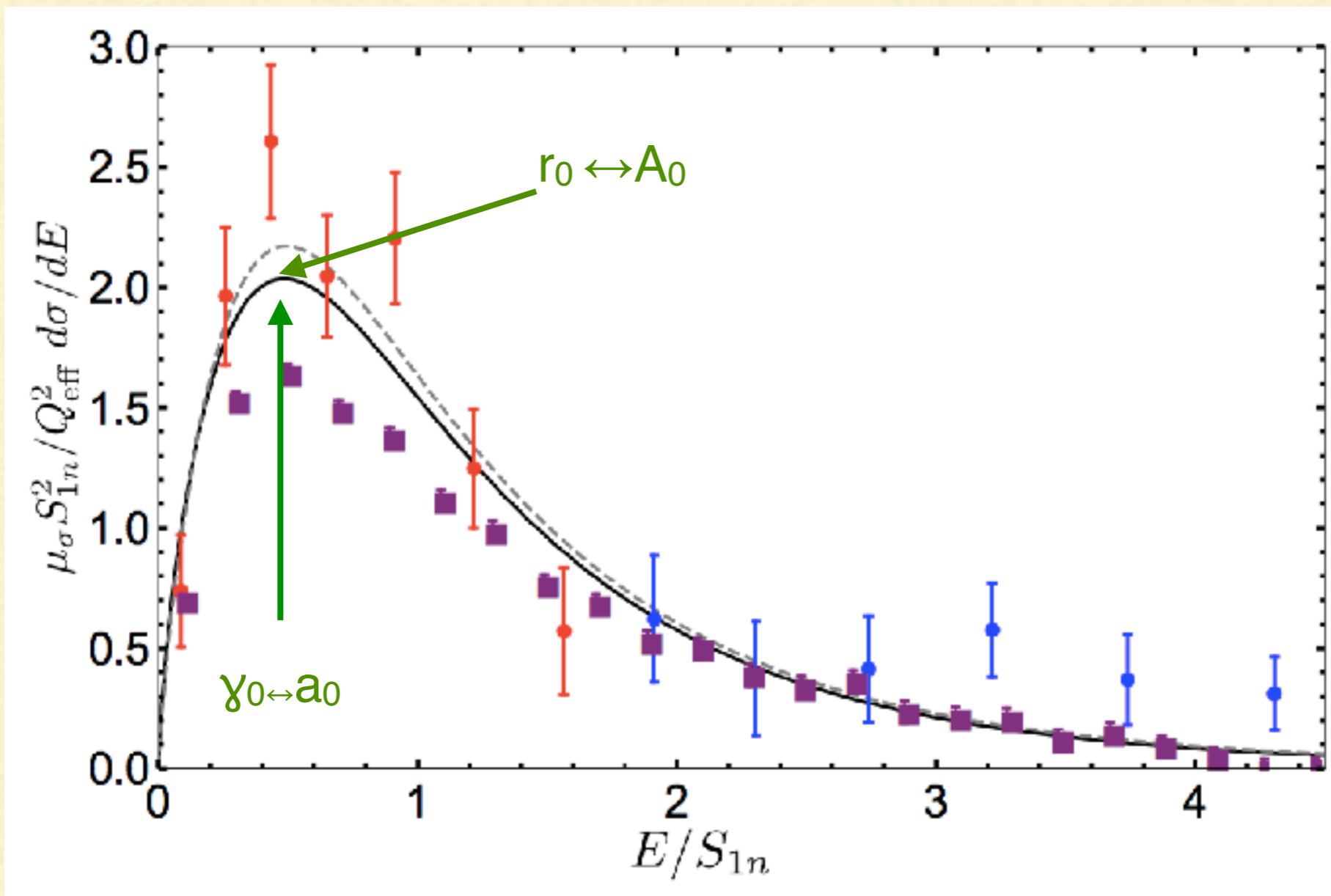
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- For  $^{19}\text{C}$ :
  - $a = (7.75 \pm 0.35(\text{stat.}) \pm 0.3(\text{EFT})) \text{ fm};$
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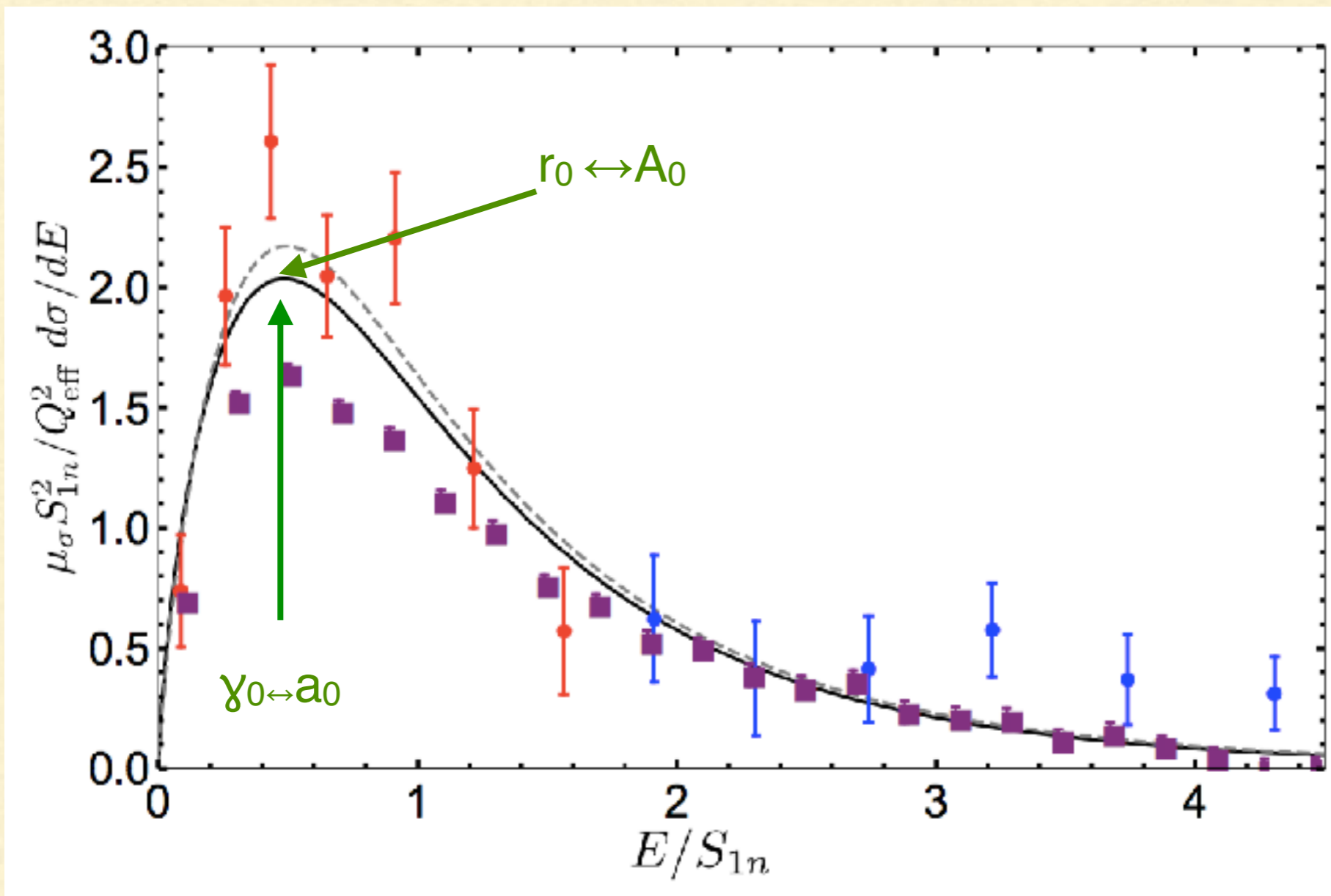
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Combine Halo EFT with reaction  
theory: Capel, Hammer, Hebborn,  
Moschini, Yang, DP, et al.

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# The three-body scattering problem

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- Consider the Schrödinger equation for the quantum-mechanical three-body scattering problem. Let's start with just pairwise potentials:

$$\left[ -\frac{\nabla_1^2}{2m_1} - \frac{\nabla_2^2}{2m_2} - \frac{\nabla_3^2}{2m_3} + V_{12}(\mathbf{r}_{12}) + V_{23}(\mathbf{r}_{23}) + V_{31}(\mathbf{r}_{31}) \right] \Psi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = E\Psi$$

- Two (related) problems:
    - Disconnected diagrams are present in solution for total wave function  $\Psi$
    - What boundary condition should we impose?
-



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# Solving problem 1

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- Let's define 2B t-matrices that are embedded in the 3B Hilbert space
- Or, equivalently, think about states  $|\psi_3^{(+)}\rangle \otimes |\mathbf{q}_3\rangle$ , where  $|\psi_3^{(+)}\rangle$  is the solution of the Schrödinger equation in which particle 3 is a “spectator”, while the (12) wave function has spherical outgoing wave boundary conditions

$$\left[ -\frac{\nabla_1^2}{2m_1} - \frac{\nabla_2^2}{2m_2} + V_{12}(\mathbf{r}_{12}) \right] \psi_3(\mathbf{r}_{12}) = E_{12}\psi_3(\mathbf{r}_{12})$$

- Questions:
    1. If  $|\psi_3^{(+)}\rangle \otimes |\mathbf{q}_3\rangle$  is a three-body eigenstate of energy E what is the energy  $E_{12}$  that appears in this Schrödinger equation
    2. Is the (12) system in its center-of-mass?
    3. If not, write the equivalent one-body Schrödinger equation, i.e., the equation for  $\psi_3(\mathbf{r}_{12})$  in the co-ordinate  $\mathbf{r}_{12}$
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**“The physics is all in the boundary conditions”**

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- And suppose I am interested in elastic neutron- $^{19}\text{C}$  scattering, what would the outgoing wave be?

$$f(q_3, \theta_3)\psi(\mathbf{r}_{12})\frac{e^{iq_3 r_3}}{r_3}$$



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# “The physics is all in the boundary conditions”

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- Suppose I want to do neutron- $^{19}\text{C}$  scattering? What would the incoming wave be?

$$\psi(\mathbf{r}_{12})e^{i\mathbf{q}_3 \cdot \mathbf{r}_3}$$

- And suppose I am interested in elastic neutron- $^{19}\text{C}$  scattering, what would the outgoing wave be?

$$f(q_3, \theta_3)\psi(\mathbf{r}_{12})\frac{e^{iq_3 r_3}}{r_3}$$

- But if I were interested in outgoing states with  $^{18}\text{C}$ -(nn) states, what would the outgoing wave be?
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- So we're going to need to specify different pieces of  $\Psi$  in order to specify what kind of reaction we're interested in
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# The Faddeev equations

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- Decompose  $|\Psi\rangle = |\Phi\rangle + |\Psi_1\rangle + |\Psi_2\rangle + |\Psi_3\rangle$
  - Where  $|\Psi_i\rangle = G_0 V_i |\Psi\rangle$  and  $|\Phi\rangle$  is a three-body plane wave; separate wave function according to the *last* interaction before particles go to the detector: that defines three separate outgoing boundary conditions
  - Show that the resulting  $|\Psi\rangle$  solves the 3B Schrödinger equation  $(E - H_0) |\Psi\rangle = (V_1 + V_2 + V_3) |\Psi\rangle$
  - Manipulate  $|\Psi_1\rangle = G_0 V_1 (|\Phi\rangle + |\Psi_1\rangle + |\Psi_2\rangle + |\Psi_3\rangle)$  into a formal equation for  $|\Psi_1\rangle$
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$$|\Psi_i\rangle = G_0 V_i |\Phi\rangle \equiv G_0 T_i |\Phi\rangle \Rightarrow T_i = t_i + t_i \sum_{j \neq i} G_0 T_j$$

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