

Chiral Effective Field Theory and Nuclear Forces: concepts of chiral EFT

Kai Hebeler

Mainz, July 29, 2022

TALENT school @MITP:
Effective field theories in light nuclei



Basic ideas of effective theories: Multipole expansion

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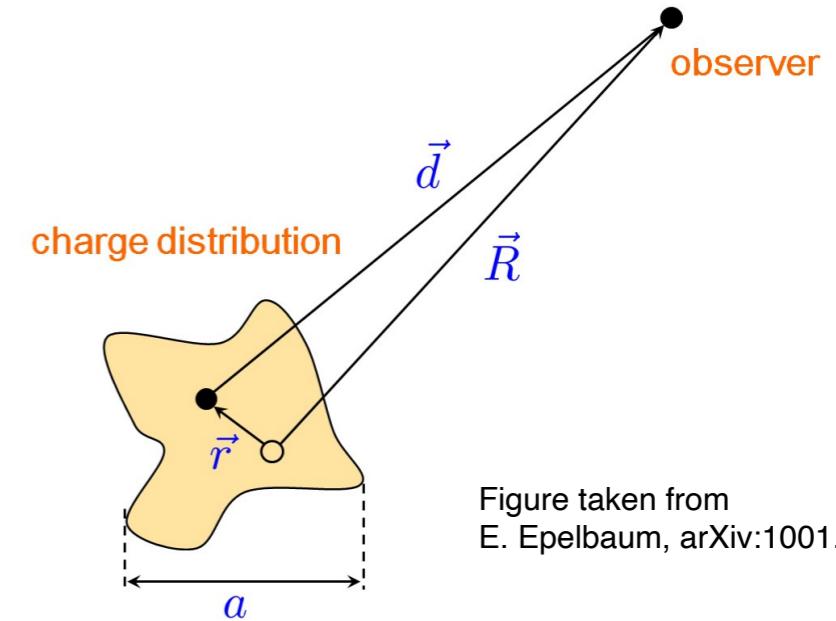


Figure taken from
E. Epelbaum, arXiv:1001.3229

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$$\Phi_{\text{exact}}(\mathbf{R}) = C \int d^3\mathbf{r} \frac{\rho(\mathbf{r})}{|\mathbf{R} - \mathbf{r}|} \quad (\text{set } C = 1)$$

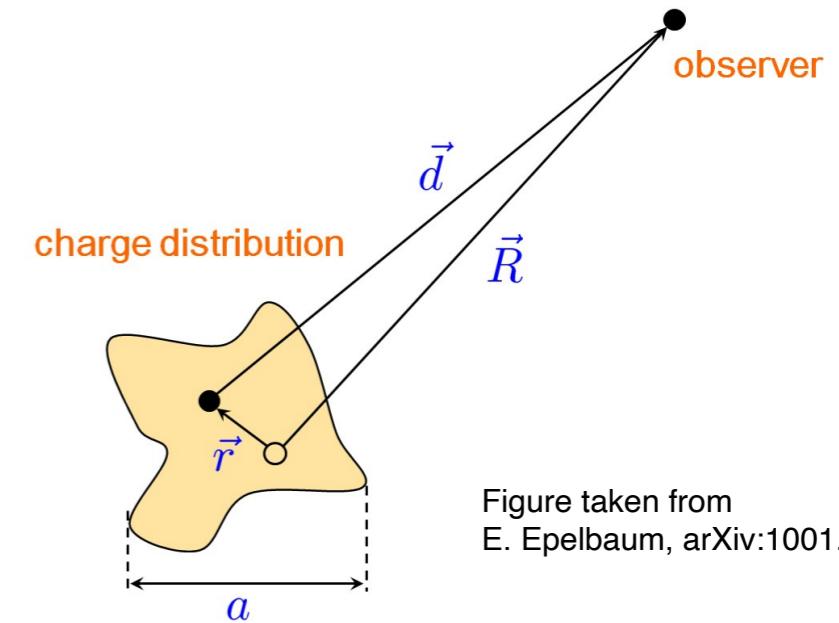
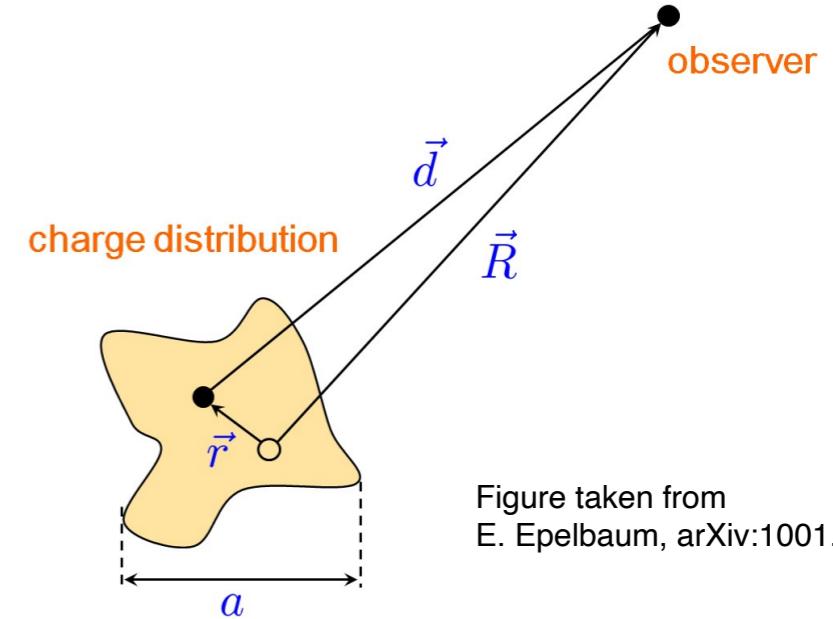


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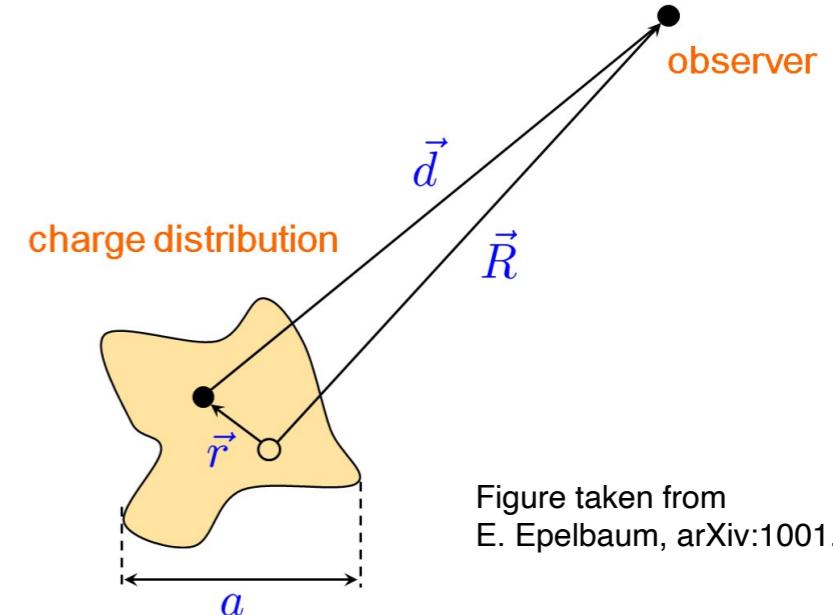
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How can we approximately determine $\Phi(\mathbf{R})$ for $|\mathbf{R}| \gg a$?

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How can we approximately determine $\Phi(\mathbf{R})$ for $|\mathbf{R}| \gg a$?

Expand $\frac{1}{|\mathbf{R} - \mathbf{r}|}$:

$$\Phi(\mathbf{R}) = \frac{q}{R} + \frac{1}{R^3} \sum_i R_i P_i + \frac{1}{6R^5} \sum_{i,j} (3R_i R_j - \delta_{ij} R^2) Q_{ij} + \dots$$

with:

$$q = \int d^3\mathbf{r} \rho(\mathbf{r}),$$

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Monopole

Dipole

Quadrupole

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3. **Fixing constants**: measure $\Phi(\mathbf{R})$ at several locations ($|\mathbf{R}| \gg a$),
determine q, P_i, Q_{ij} and make predictions for other locations

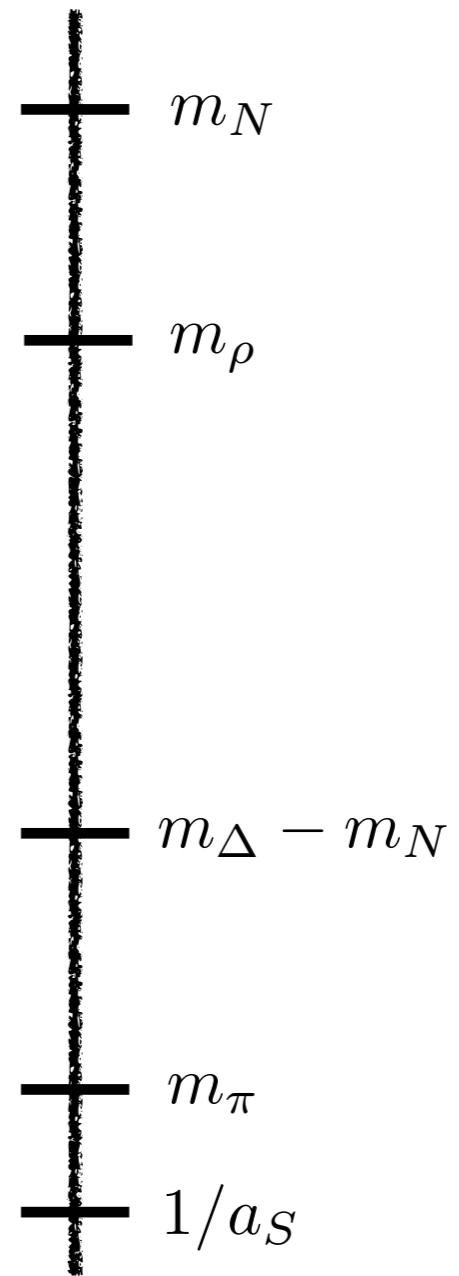
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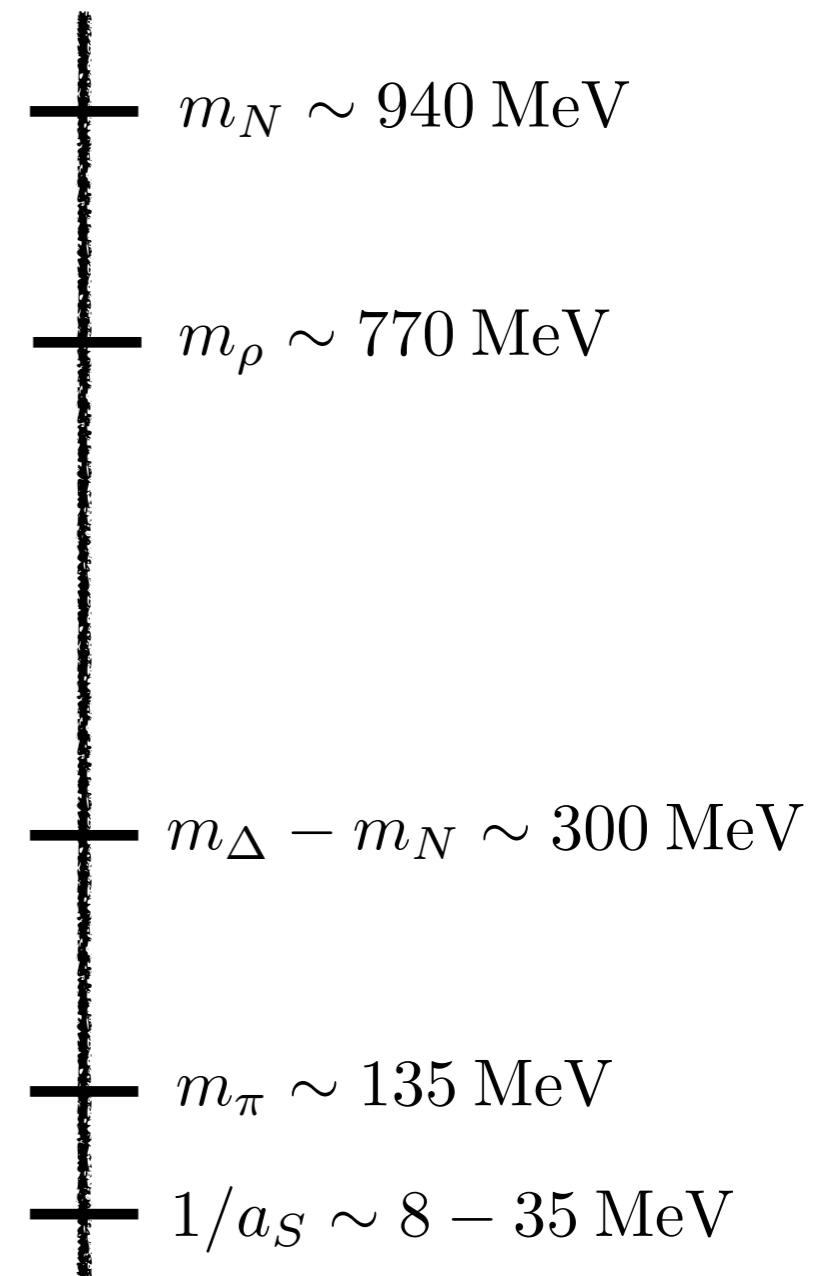
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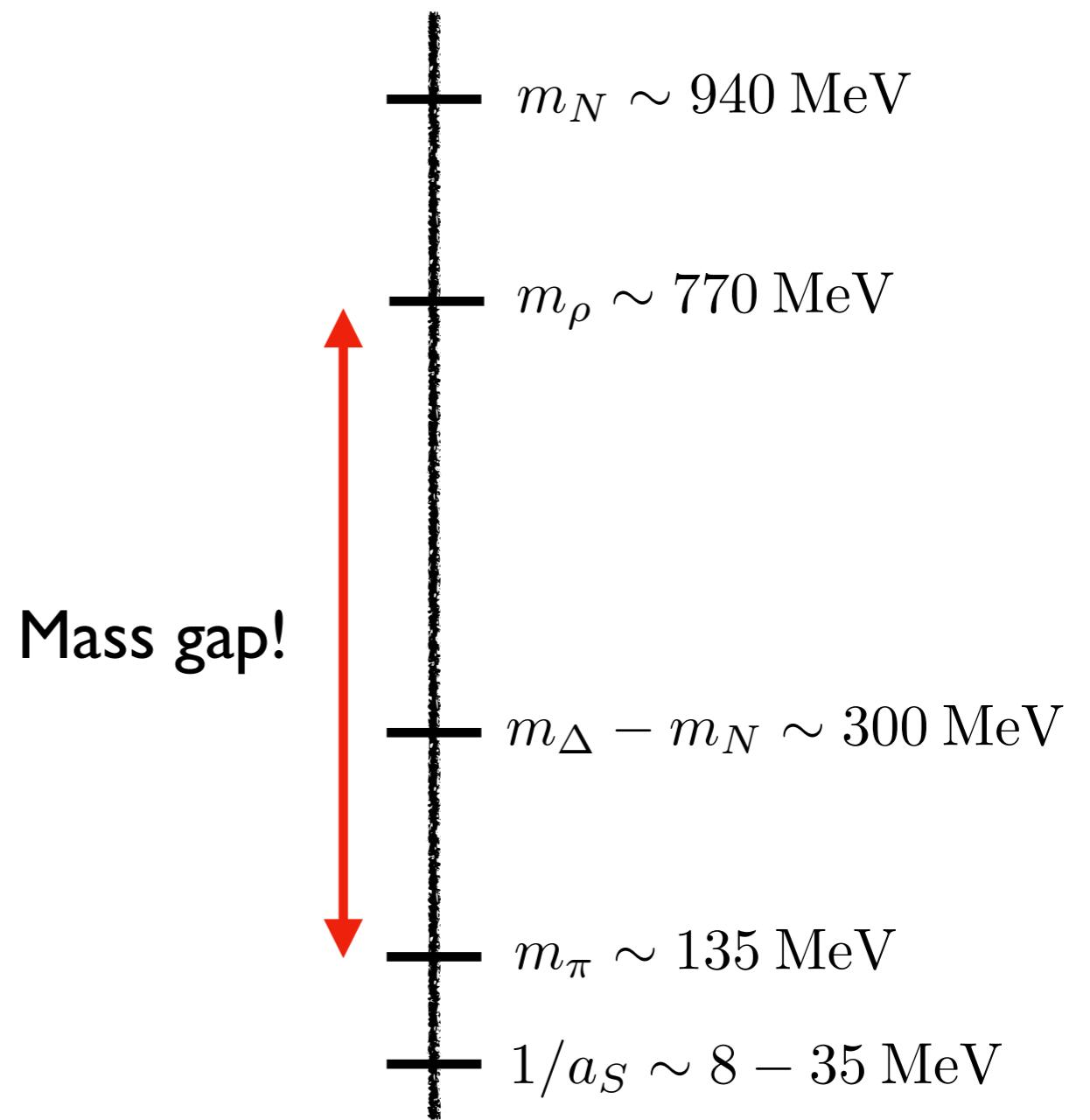
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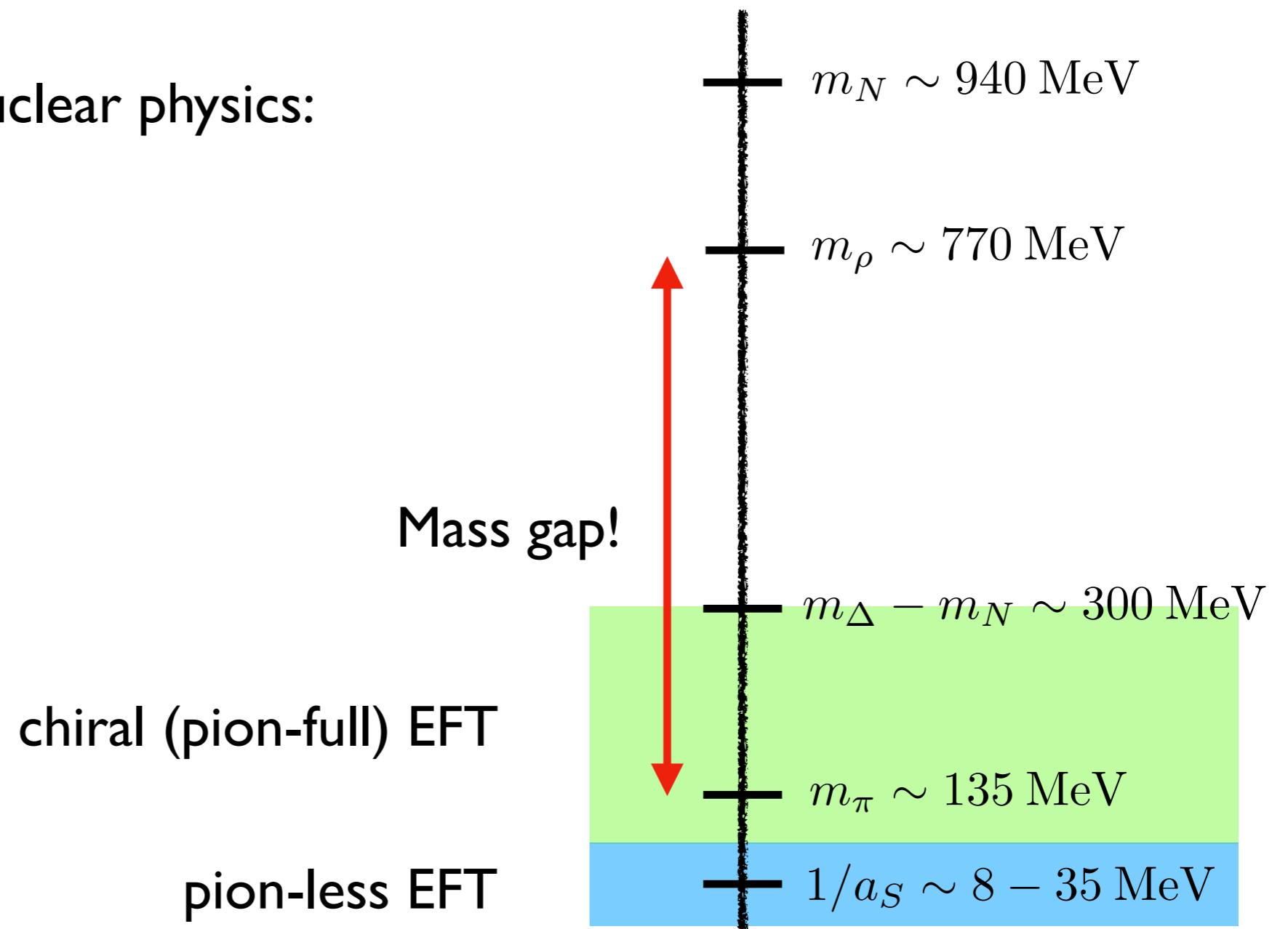
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Basic idea: utilize separation of energy scales!

$$\frac{Q}{\Lambda} \quad \begin{array}{l} \xleftarrow{\hspace{1cm}} \text{typical momenta of nucleons} \\ \xleftarrow{\hspace{1cm}} \text{breakdown scale of EFT } (\sim 500 \text{ MeV}) \end{array}$$

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3. **Fixing of low-energy constants**: perform calculations at a given order in expansion, determine constants by matching to experimental data (e.g. NN scattering observables) and make predictions for other observables, in principle constants can be computed from QCD (cf. multipole expansion)

Symmetries of QCD - chiral symmetry

$$\begin{aligned}\mathcal{L}_{QCD} &= -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + \bar{q}_f(i\gamma^\mu \partial_\mu - m_f)q_f + g\bar{q}\gamma^\mu T_a q A_\mu^a \\ &\equiv -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + \bar{q}_f(i\gamma^\mu D_\mu - m_f)q_f\end{aligned}$$

- f : flavour index, here we consider only light flavors ($f = u, d$)
- ‘covariant derivative’ (independent of flavor):

$$D_\mu = \partial_\mu - ig\bar{q}\gamma^\mu T_a q A_\mu^a$$

- introduce mass matrix (in flavour space):

$$\mathcal{M} = \text{diag}(m_u, m_d)$$

Decompose quark fields (for each flavor) into left- and right-handed **chiral** components via:

$$q_L = P_L q, \quad q_R = P_R q, \quad P_{L/R} = \frac{1}{2} (1 \mp \gamma_5), \quad P_L + P_R = 1$$

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quark number
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quark number broken by
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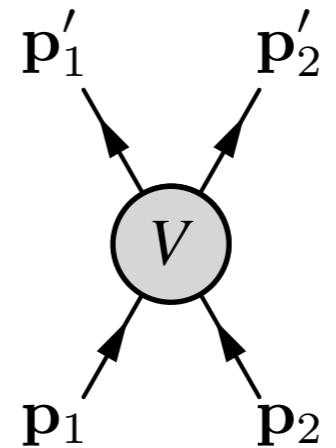
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quark number conservation	broken by anomaly	spontaneously broken down to subgroup $SU(2)_V$ explicitly broken by finite \mathcal{M}
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Regularization schemes for nuclear interactions (here: NN)

Separation of long- and short-range physics



$$\mathbf{p} = (\mathbf{p}_1 - \mathbf{p}_2)/2$$

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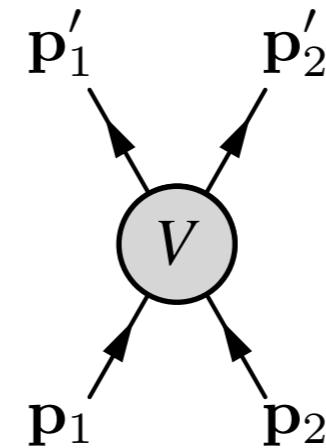
$$\mathbf{q} = (\mathbf{p}_1 - \mathbf{p}'_1)$$

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nonlocal

$$V_{\text{NN}}(\mathbf{p}, \mathbf{p}') \rightarrow \exp \left[- \left((p^2 + p'^2)/\Lambda^2 \right)^n \right] V_{\text{NN}}(\mathbf{p}, \mathbf{p}')$$

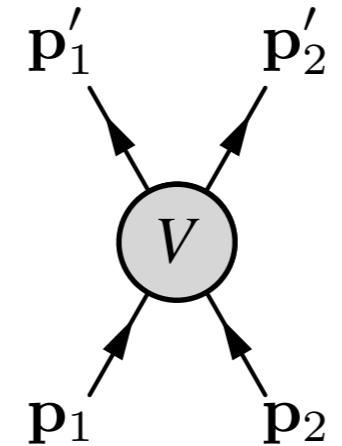


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Epelbaum, Glöckle, Meissner, NPA 747, 362 (2005)
Entem, Machleidt, PRC 68, 041001 (2003)

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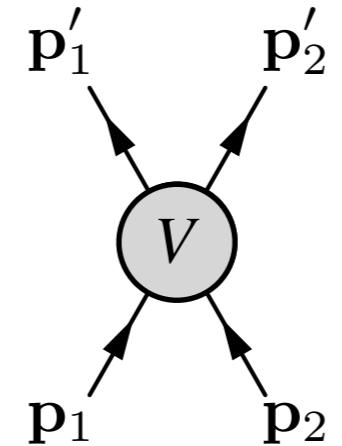
local
(momentum space)

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Regularization schemes for nuclear interactions (here: NN)

Separation of long- and short-range physics



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Epelbaum, Glöckle, Meissner, NPA 747, 362 (2005)
Entem, Machleidt, PRC 68, 041001 (2003)

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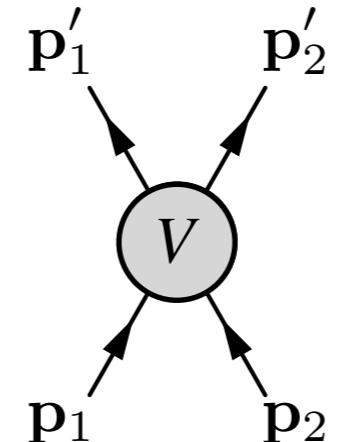
local
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$$\begin{aligned} V_{\text{NN}}^\pi(\mathbf{r}) &\rightarrow \left(1 - \exp \left[- \left(r^2/R^2 \right)^n \right] \right) V_{\text{NN}}^\pi(\mathbf{r}) \\ \delta(\mathbf{r}) &\rightarrow \alpha_n \exp \left[- \left(r^2/R^2 \right)^n \right] \end{aligned}$$

Gezerlis et. al, PRL, 111, 032501 (2013)

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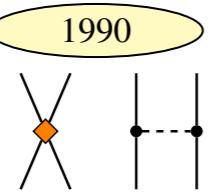
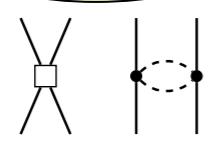
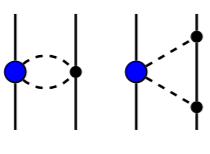
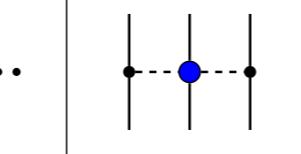
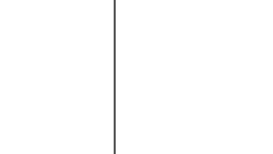
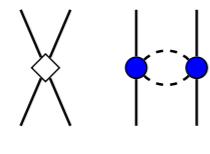
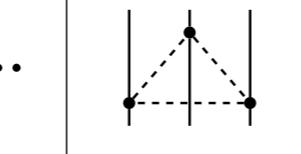
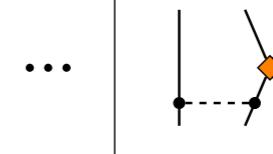
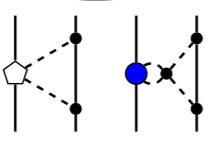
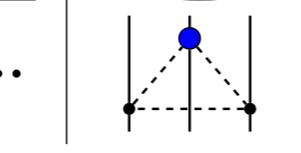
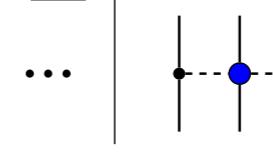
Gezerlis et. al, PRL, 111, 032501 (2013)

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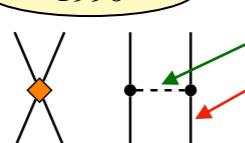
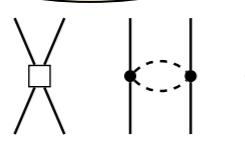
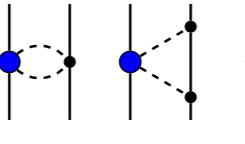
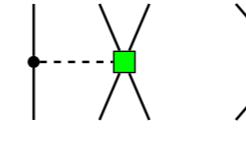
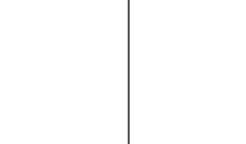
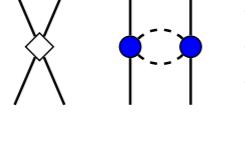
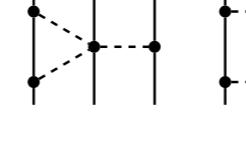
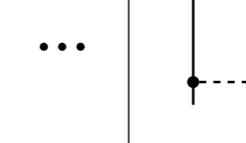
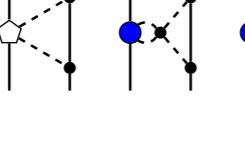
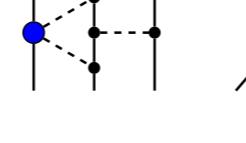
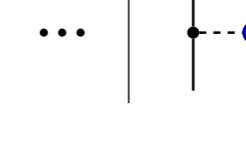
Epelbaum et. al, PRL, 115, 122301 (2015)

Chiral effective field theory for nuclear forces

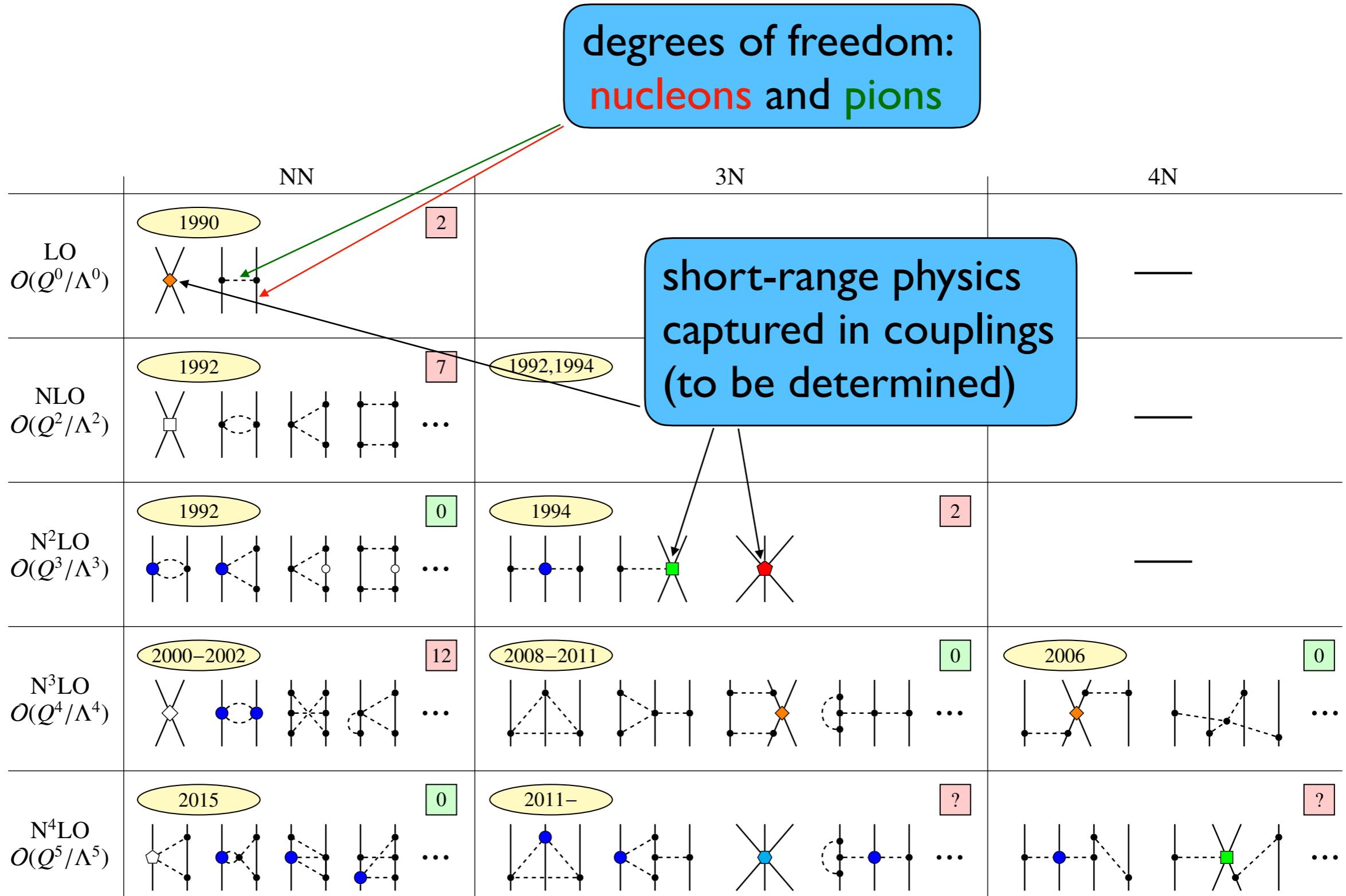
	NN	3N	4N
$O(Q^0/\Lambda^0)$	 1990	—	—
$O(Q^2/\Lambda^2)$	 1992	 1992, 1994	—
$O(Q^3/\Lambda^3)$	 1992	 1994	 2
$O(Q^4/\Lambda^4)$	 2000–2002	 2008–2011	 0
$O(Q^5/\Lambda^5)$	 2015	 2011–	 ?

Chiral effective field theory for nuclear forces

degrees of freedom:
nucleons and pions

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$N^4\text{LO}$ $O(Q^5/\Lambda^5)$	 2015	 2011–	 ?

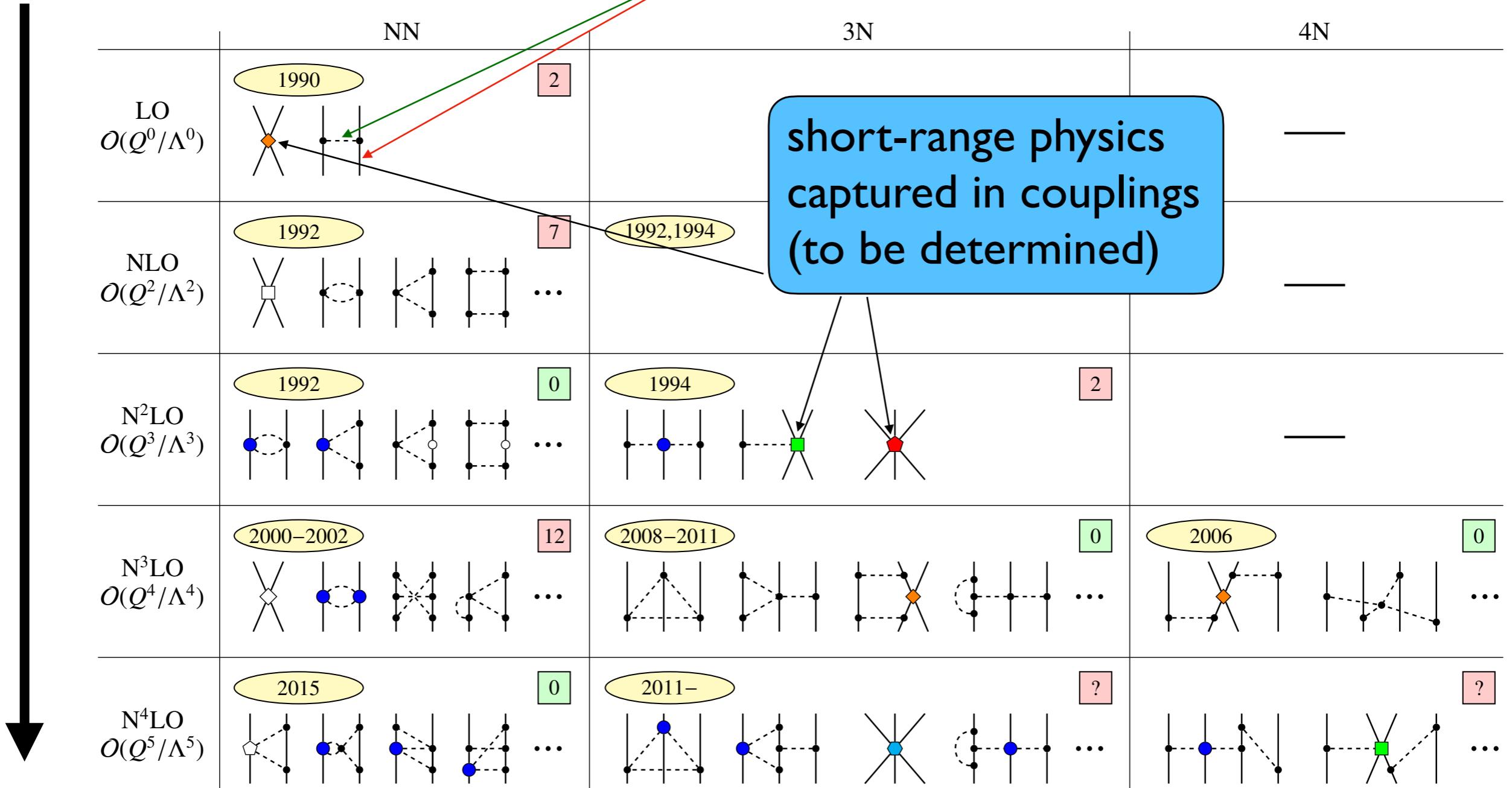
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power-counting:
expand in Q/Λ , error estimates!

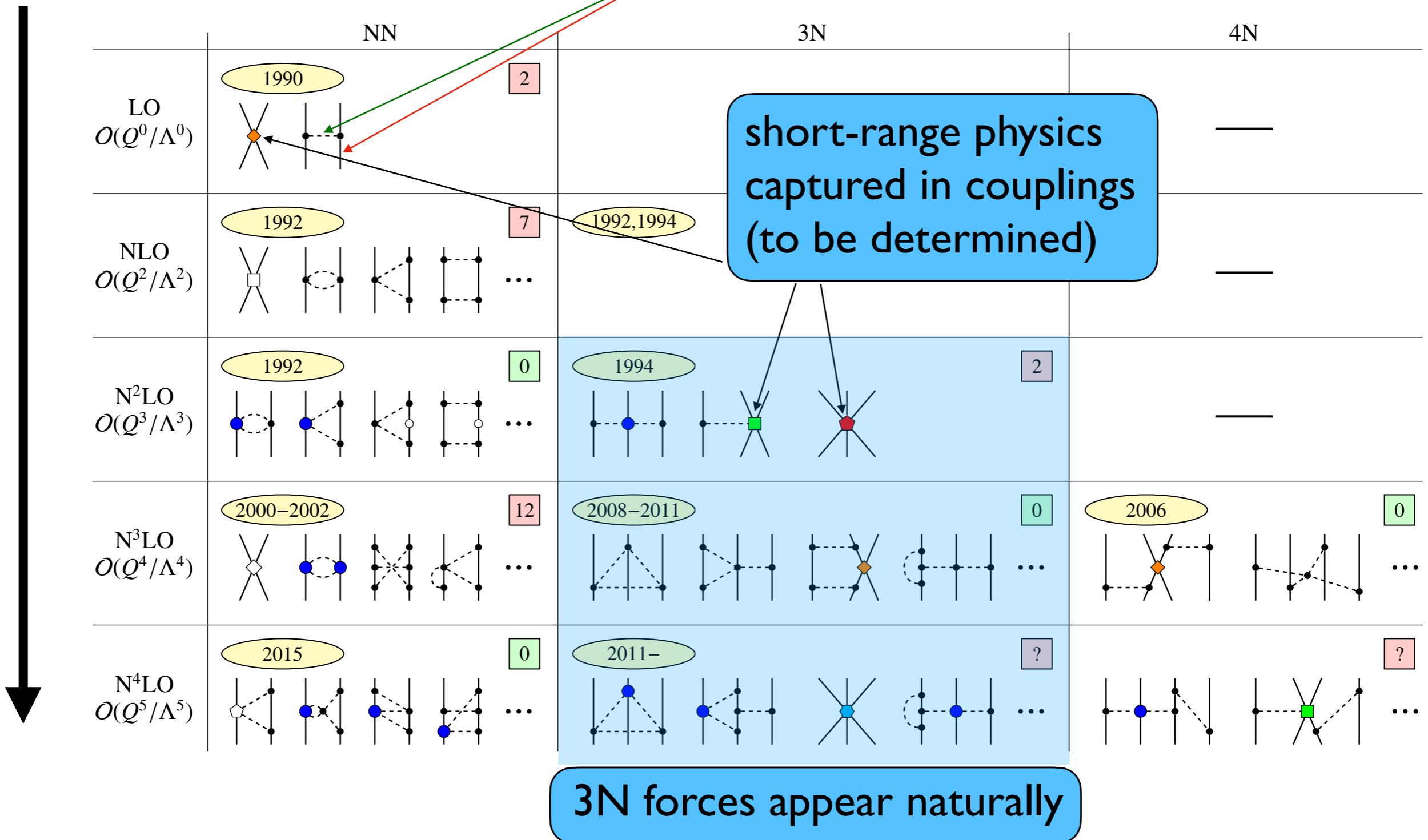
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Chiral EFT uncertainty estimation

EFT expansion of an observable X at a momentum scale p :

$$X(p) = X_{\text{ref}}(p) \sum_{n=0}^{\infty} c_n(p) Q^n$$

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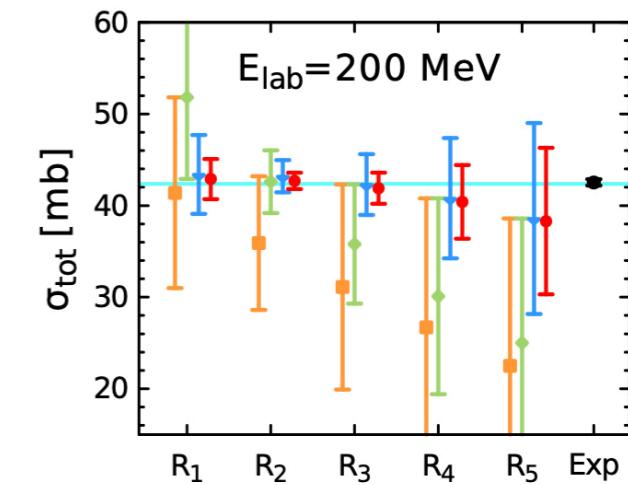
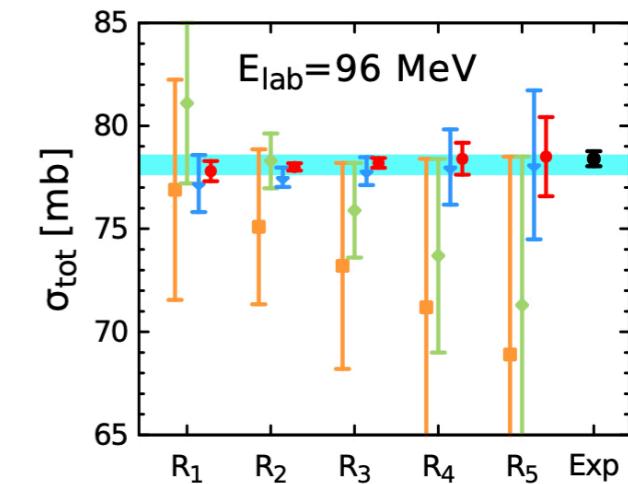
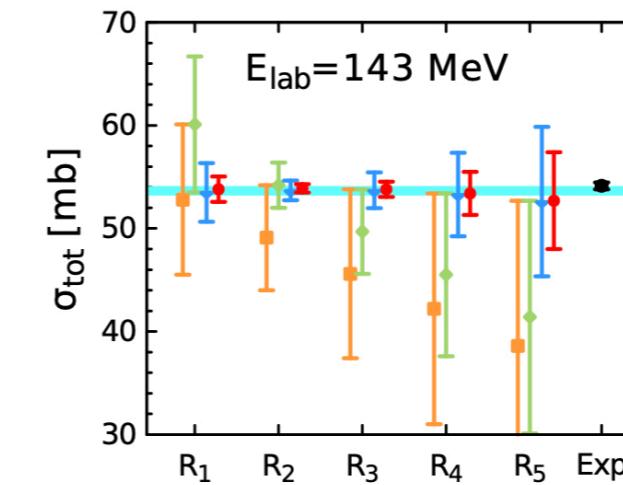
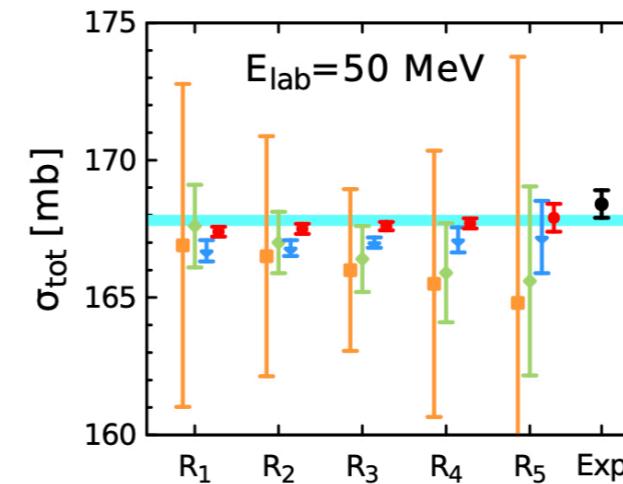
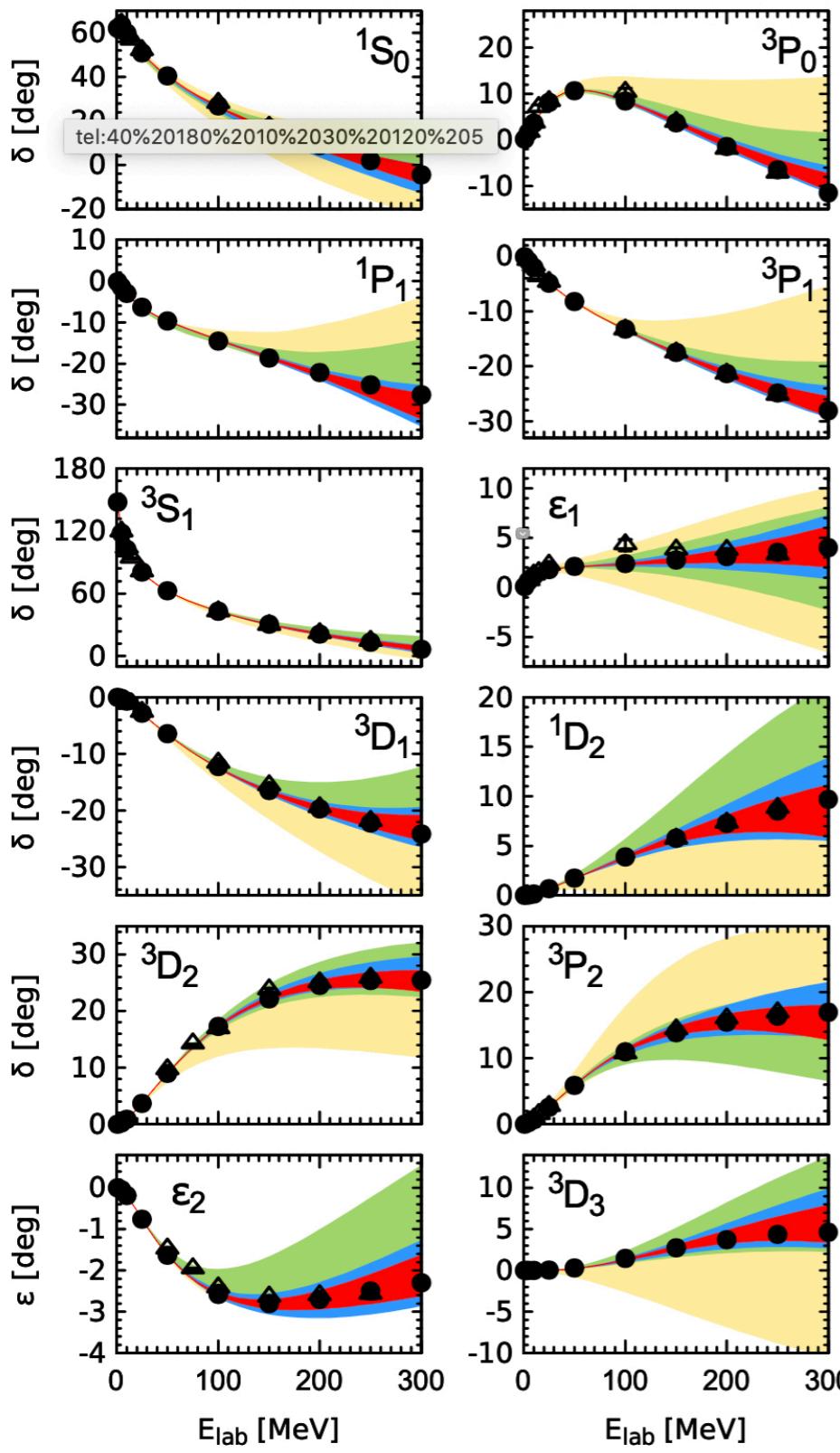
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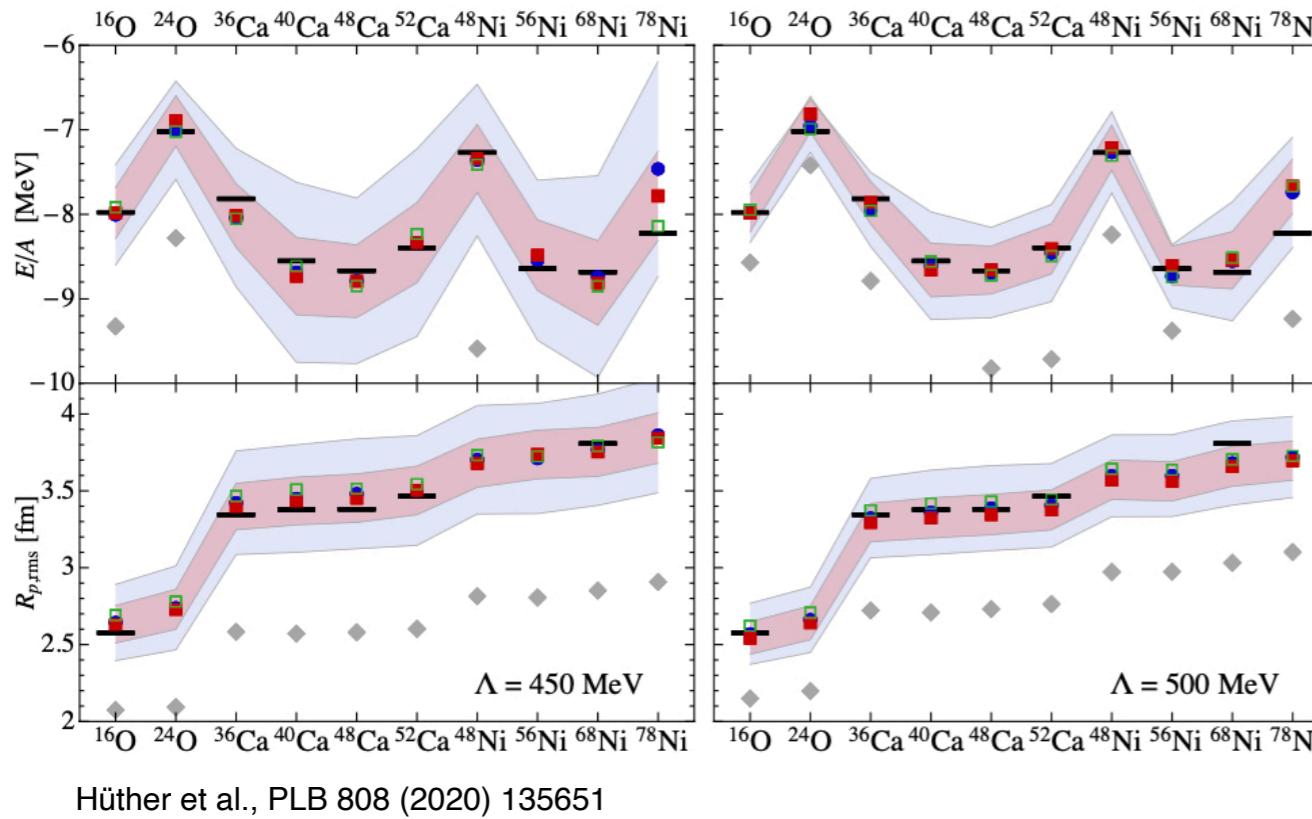
A conservative prescription: $\Delta X^{\text{N}^3\text{LO}}(p) = \max(Q^5 |X^{\text{LO}}(p)|, Q^3 |X^{\text{LO}}(p) - X^{\text{NLO}}(p)|, Q^2 |X^{\text{NLO}}(p) - X^{\text{N}^2\text{LO}}(p)|, Q |X^{\text{N}^2\text{LO}}(p) - X^{\text{N}^3\text{LO}}(p)|)$

Chiral EFT uncertainty estimation: examples

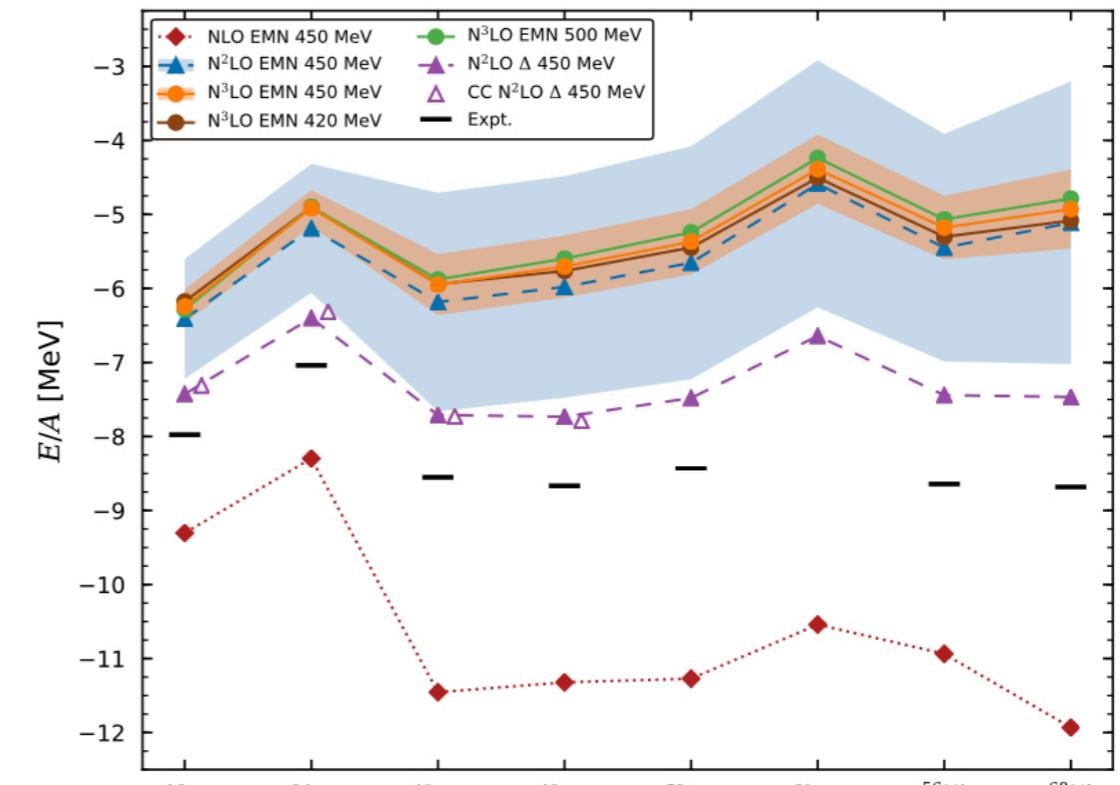


Epelbaum et al., PRL 115 (2015) 122301

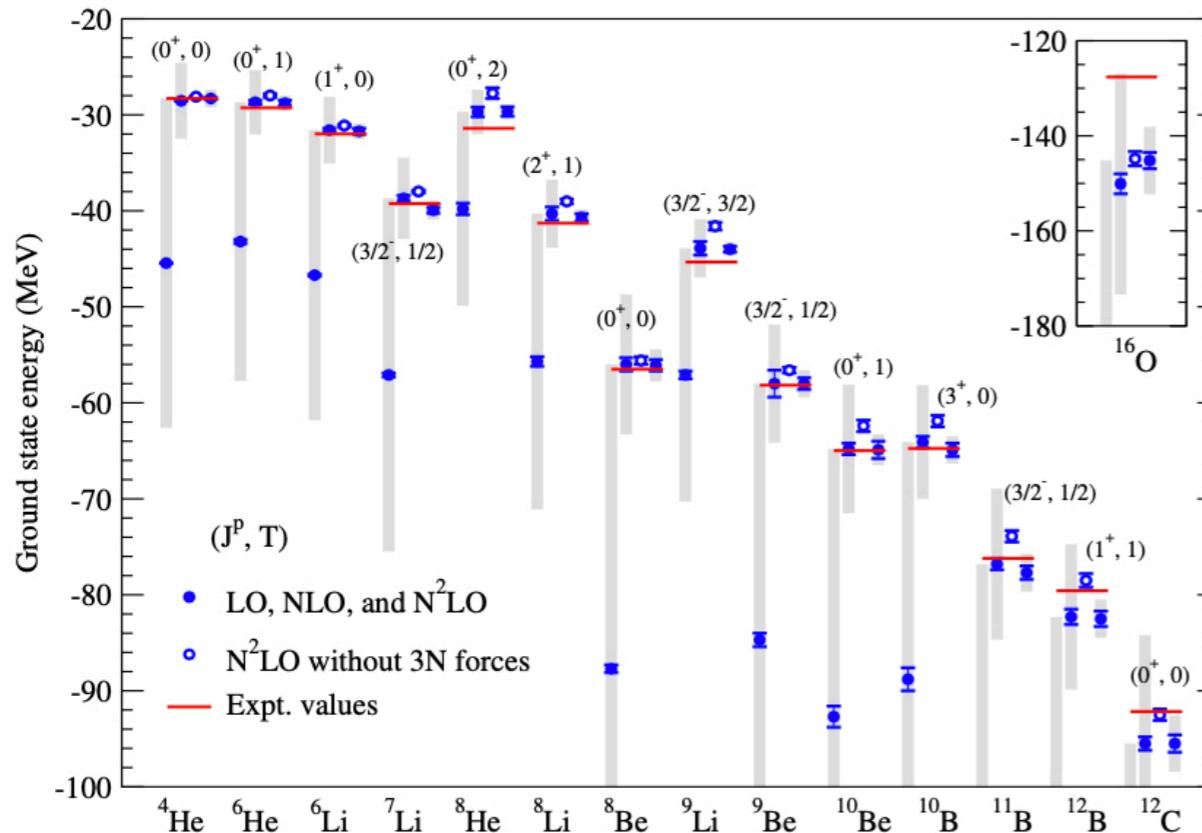
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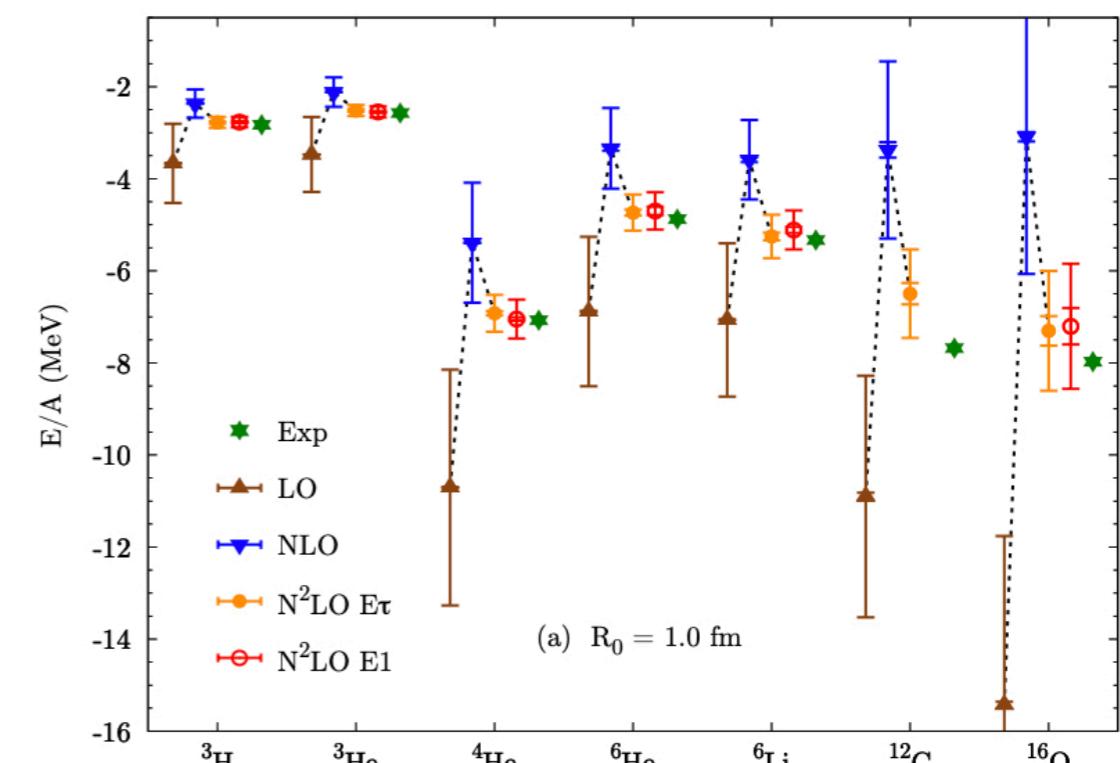
Hüther et al., PLB 808 (2020) 135651



Hoppe et al., PRC 100 (2019) 024318

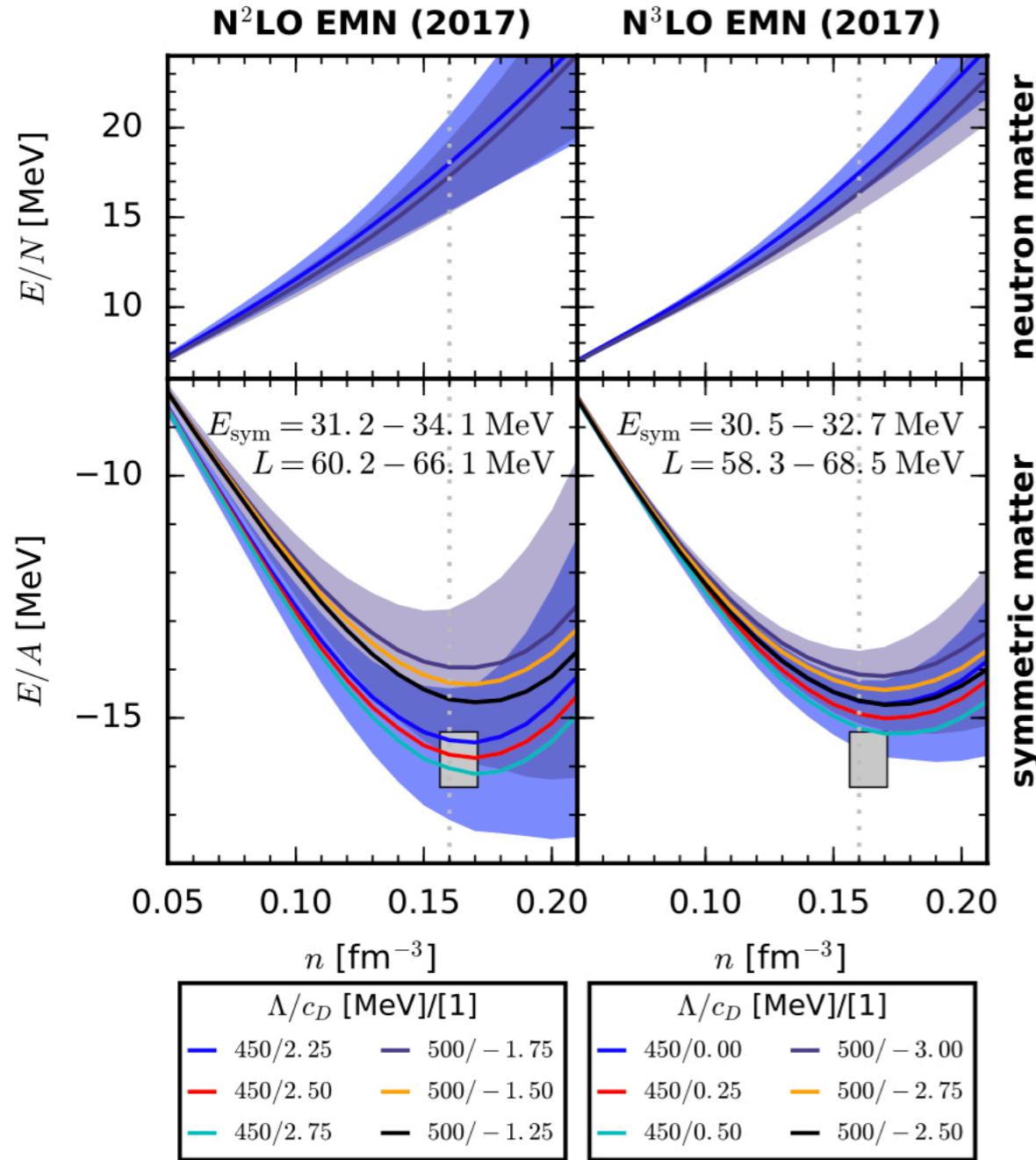


Epelbaum et al., PRC 99 (2019) 024313

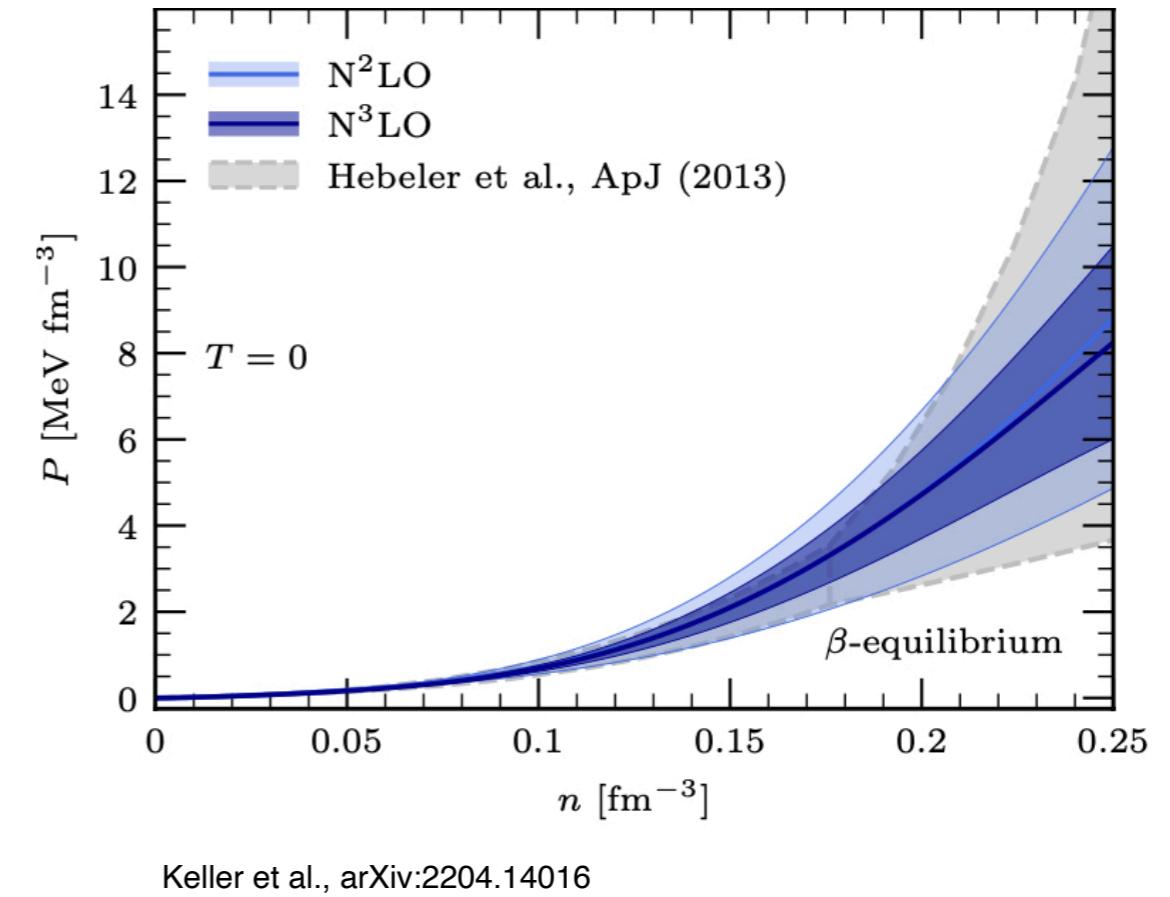


Lonardoni et al., PRC 97 (2018) 044318

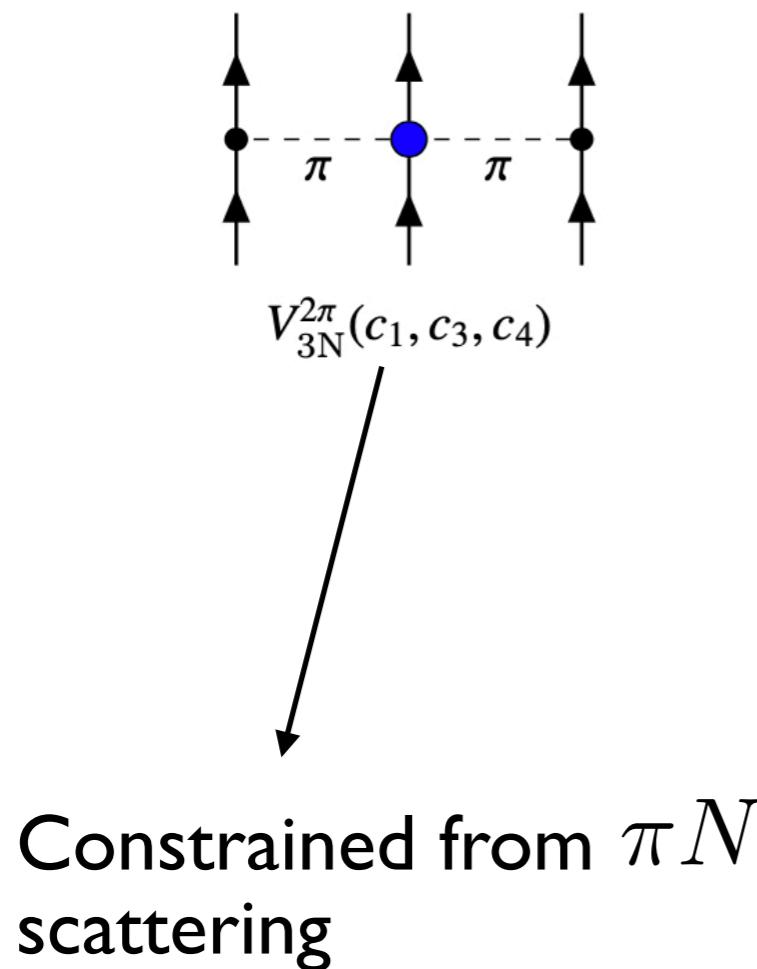
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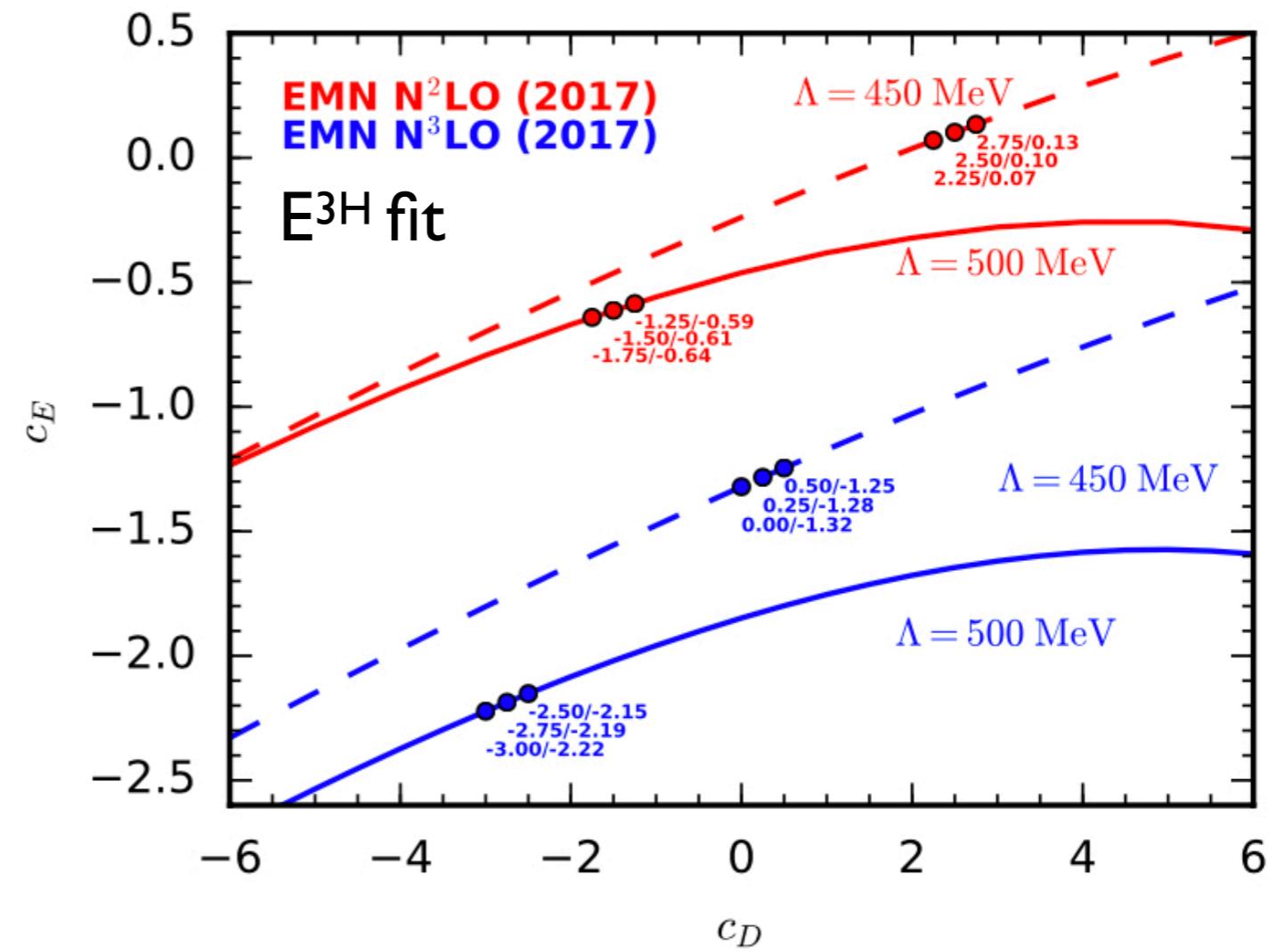
Drischler et al., PRL 122 (2019) 042501



Fits of 3N LECs

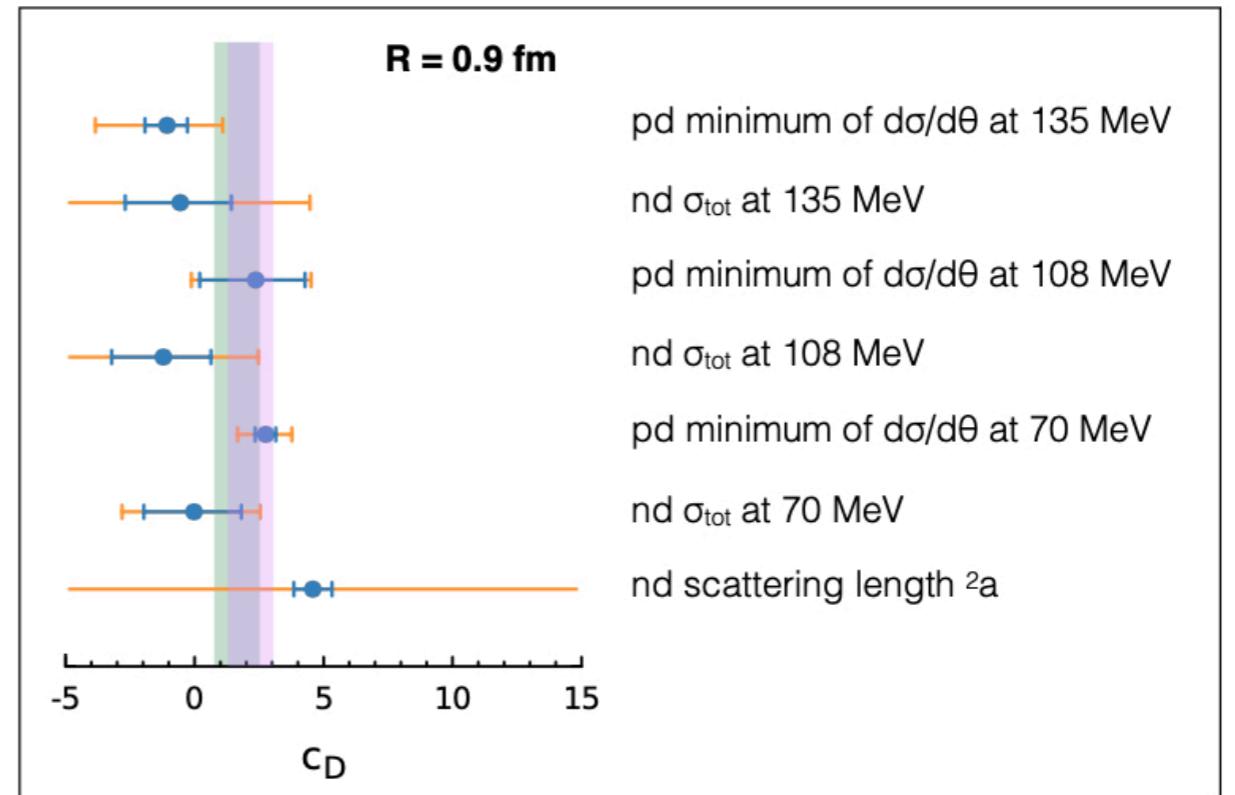
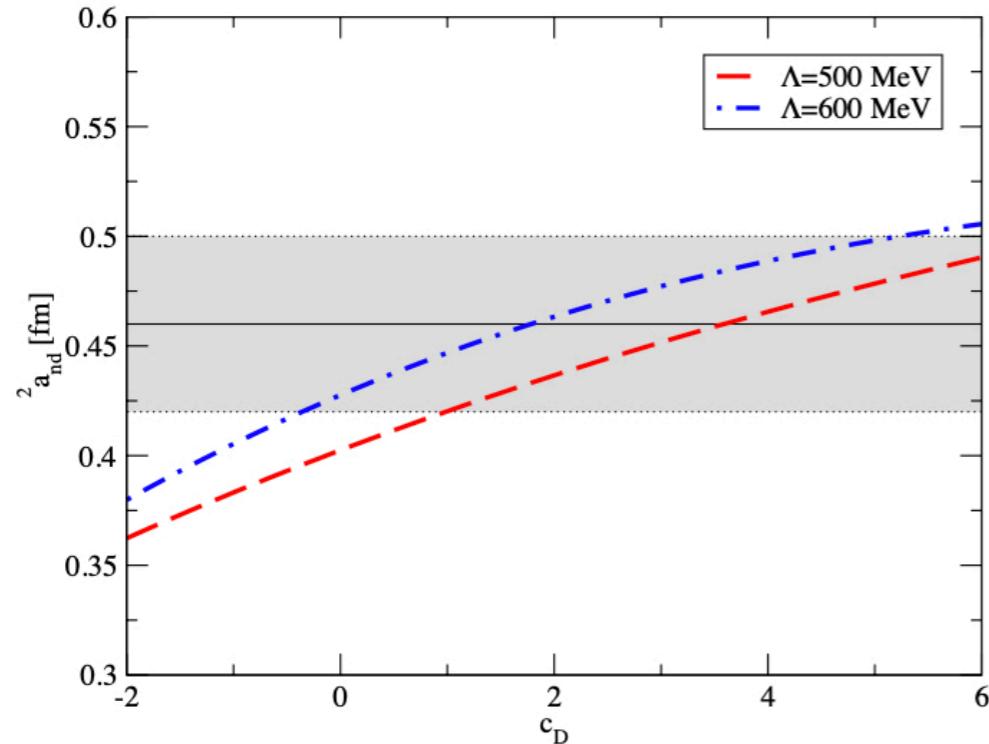


Hoferichter et al., Phys. Rept. 625 (2016) 1



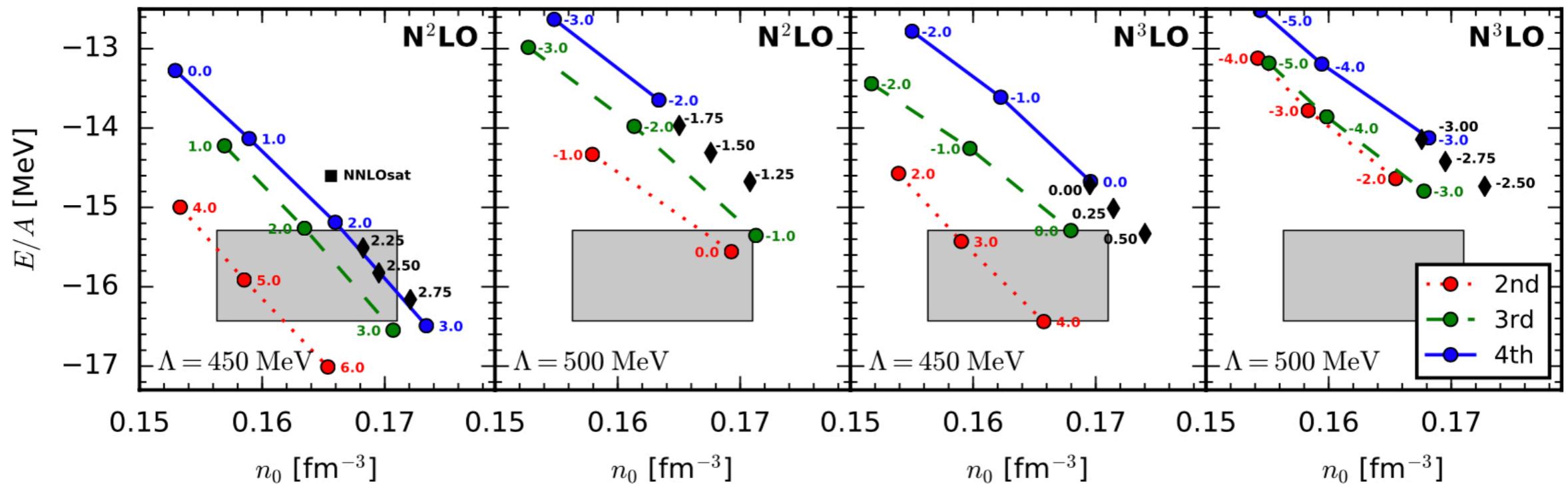
Drischler et al., PRL 122 (2019) 042501

Fits of 3N LECs: three-body scattering cross sections



- a single scattering observable not too constraining (correlated with E^{3H})
- a more global fit using several observables more robust

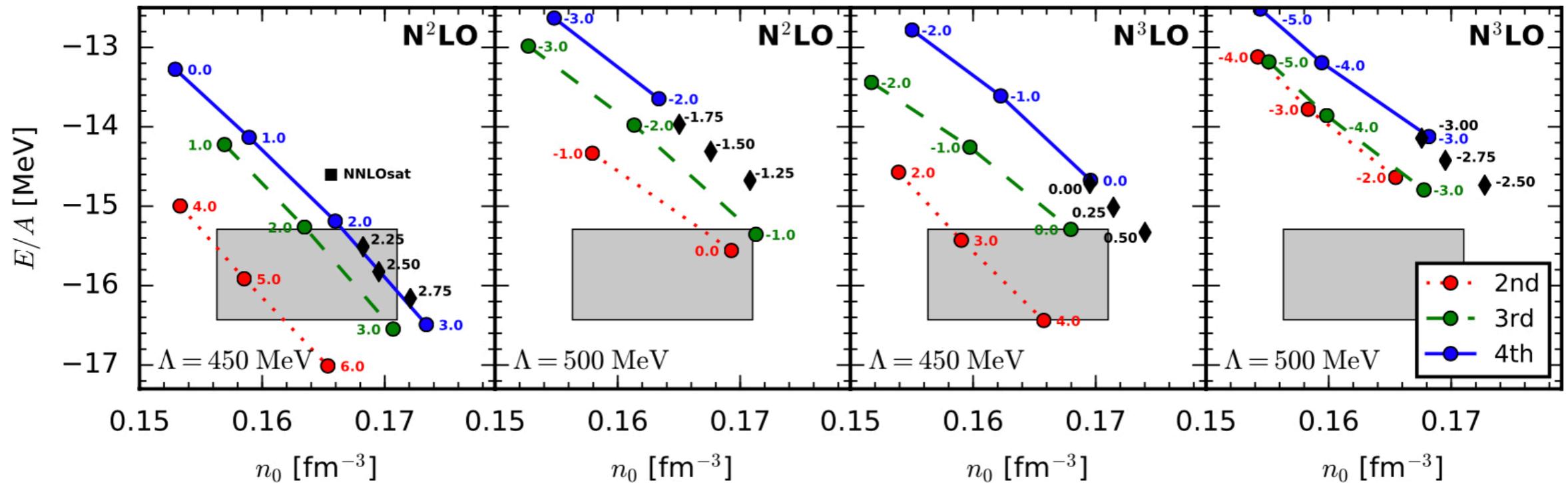
Determination of LECs: From nuclear matter saturation point



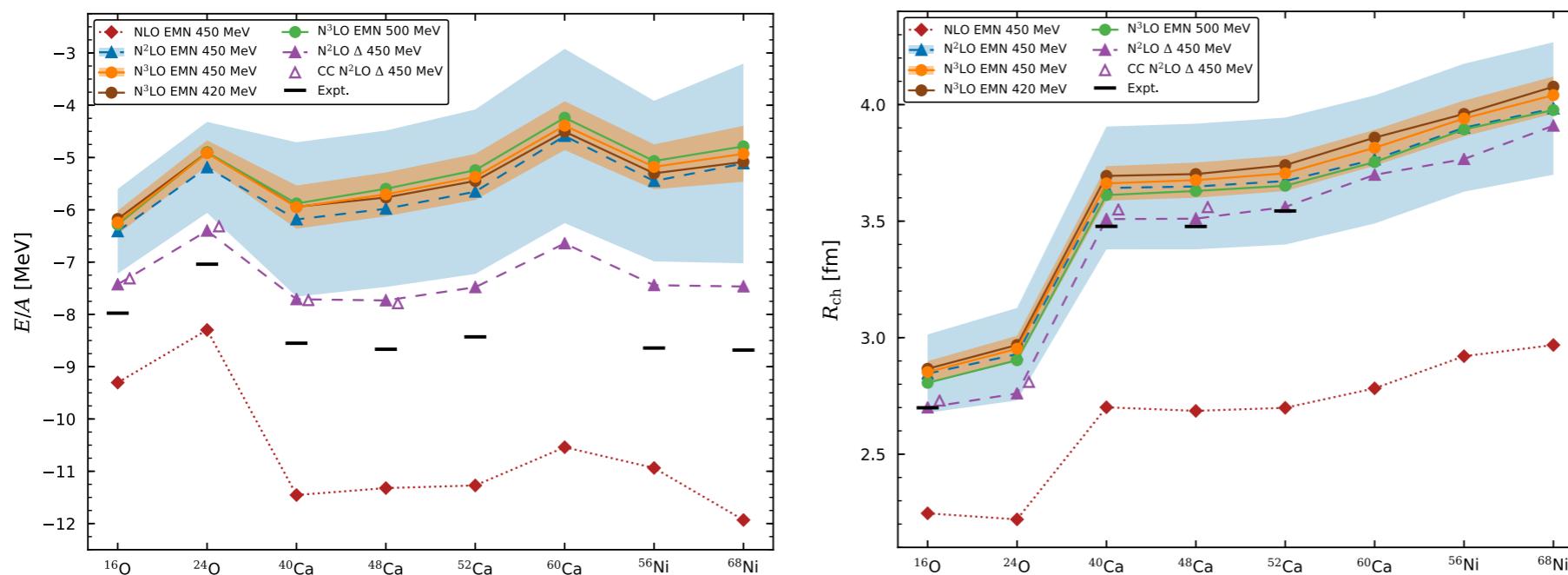
Drischler et al., PRL 122 (2019) 042501

- Use nuclear matter saturation energy and density to constrain LECs
- Reasonable reproduction of both quantities possible
- Results for medium-mass nuclei still not satisfactory

Determination of LECs: From nuclear matter saturation point

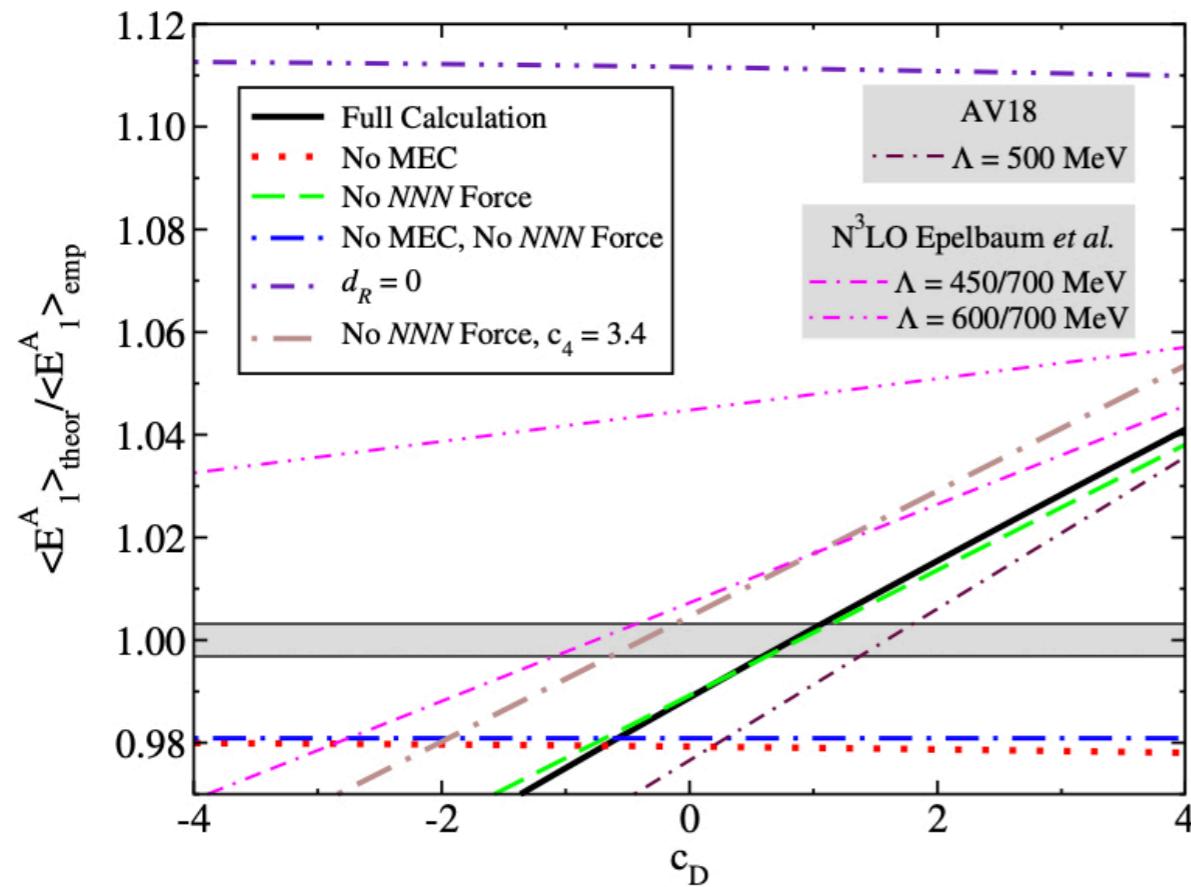


Drischler et al., PRL 122 (2019) 042501

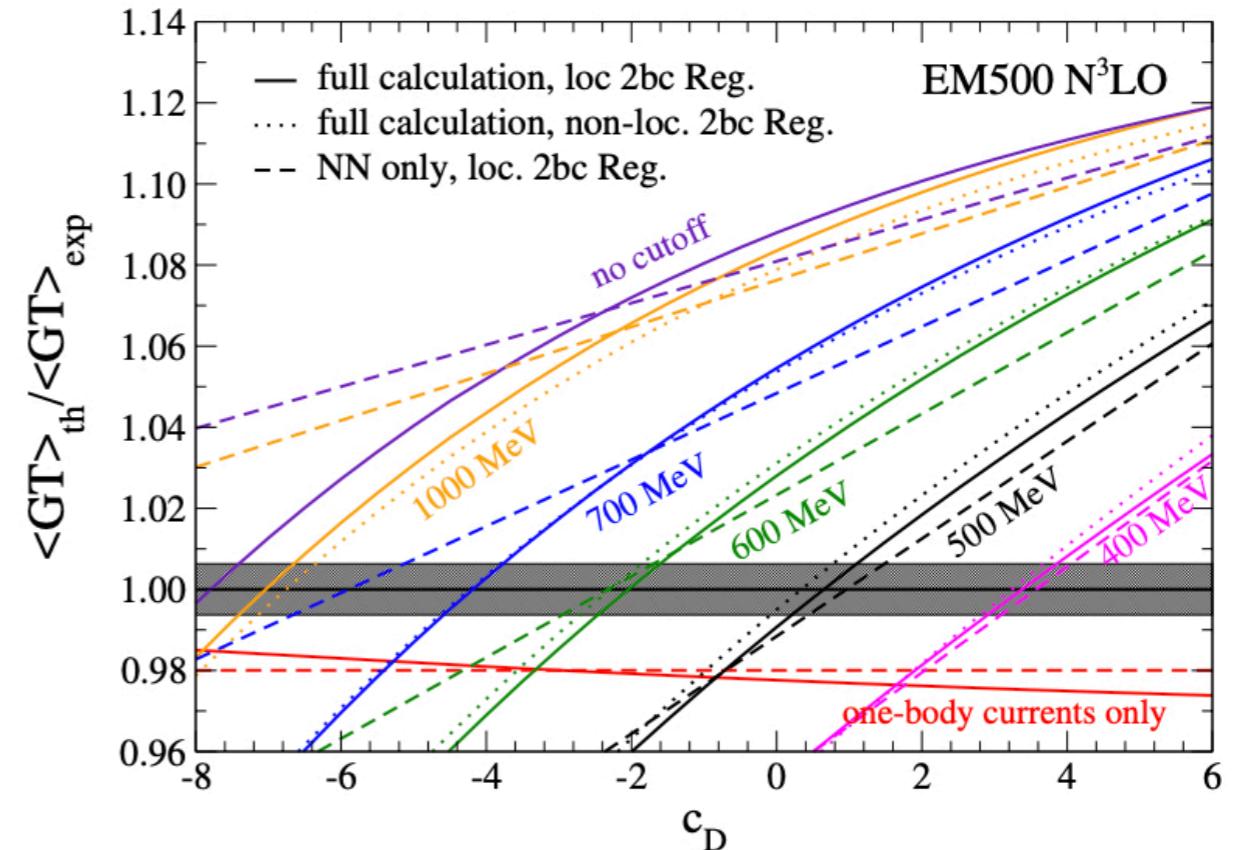


Hoppe et al.
PRC 100, 024318 (2019)

Determination of LECs: triton beta decay half life



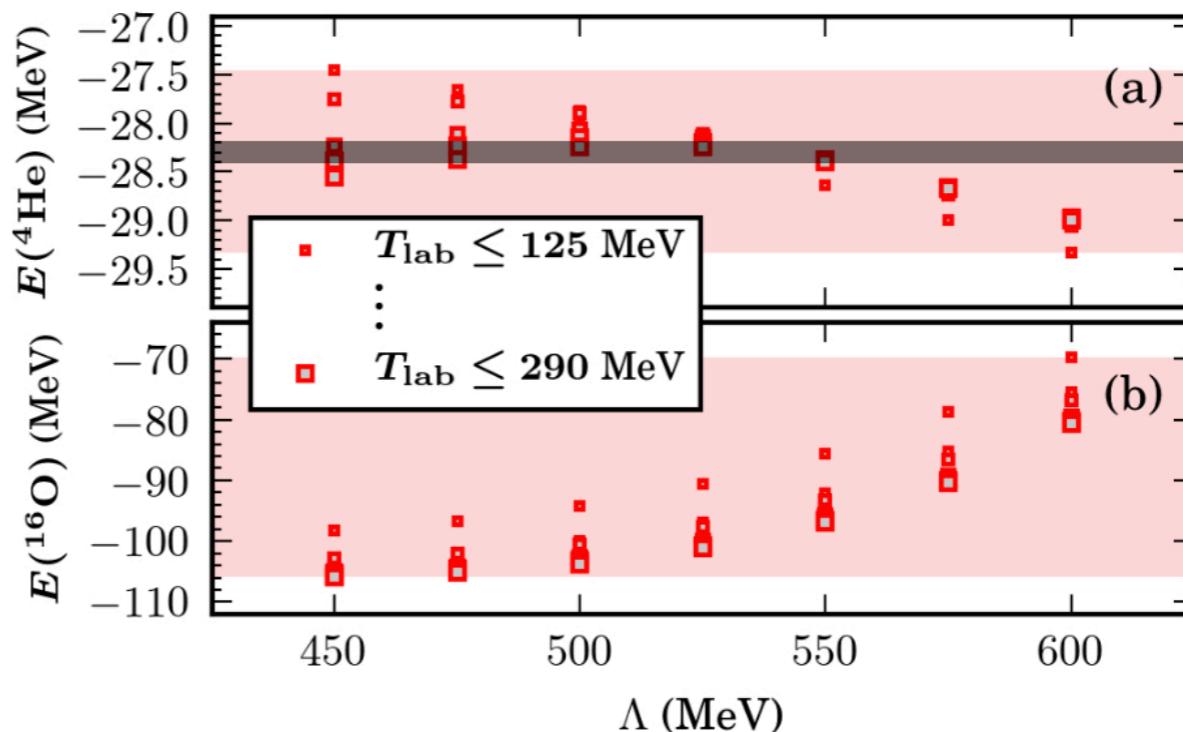
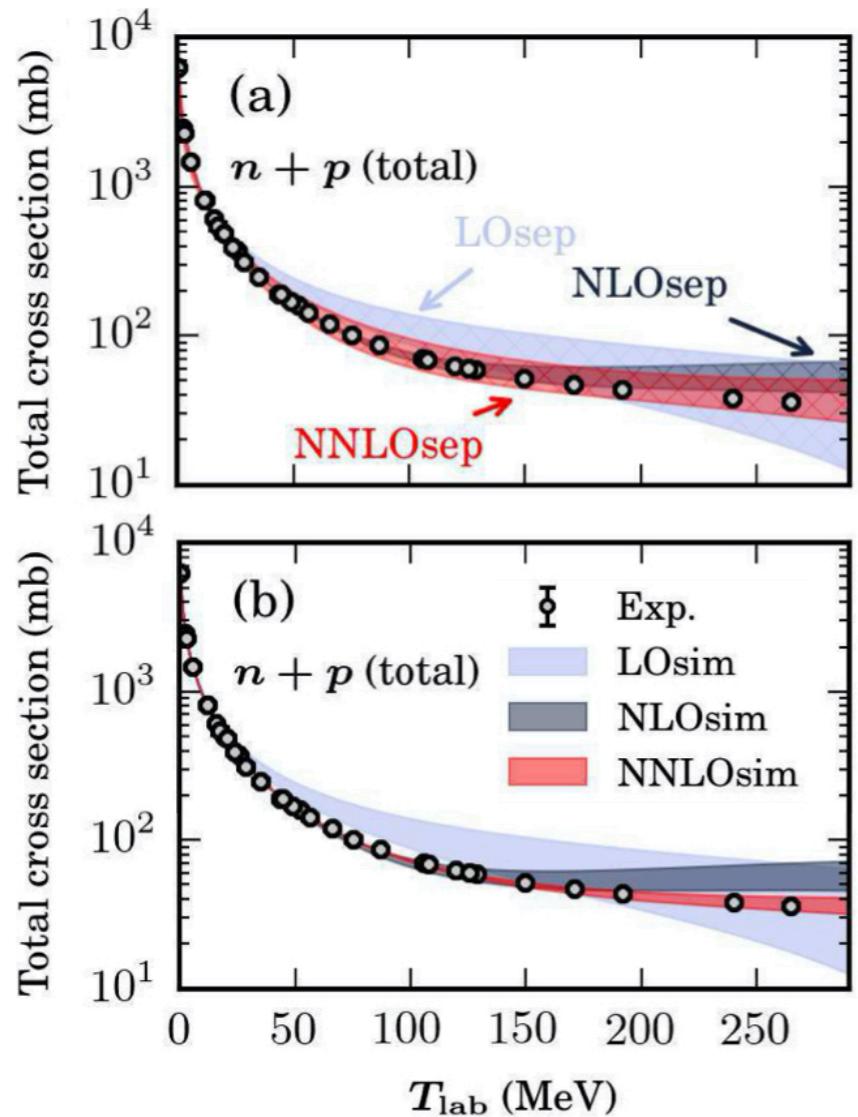
Gazit et al., PRL103 (2009) 102502



Klos et al., EPJA 53 (2017) 168

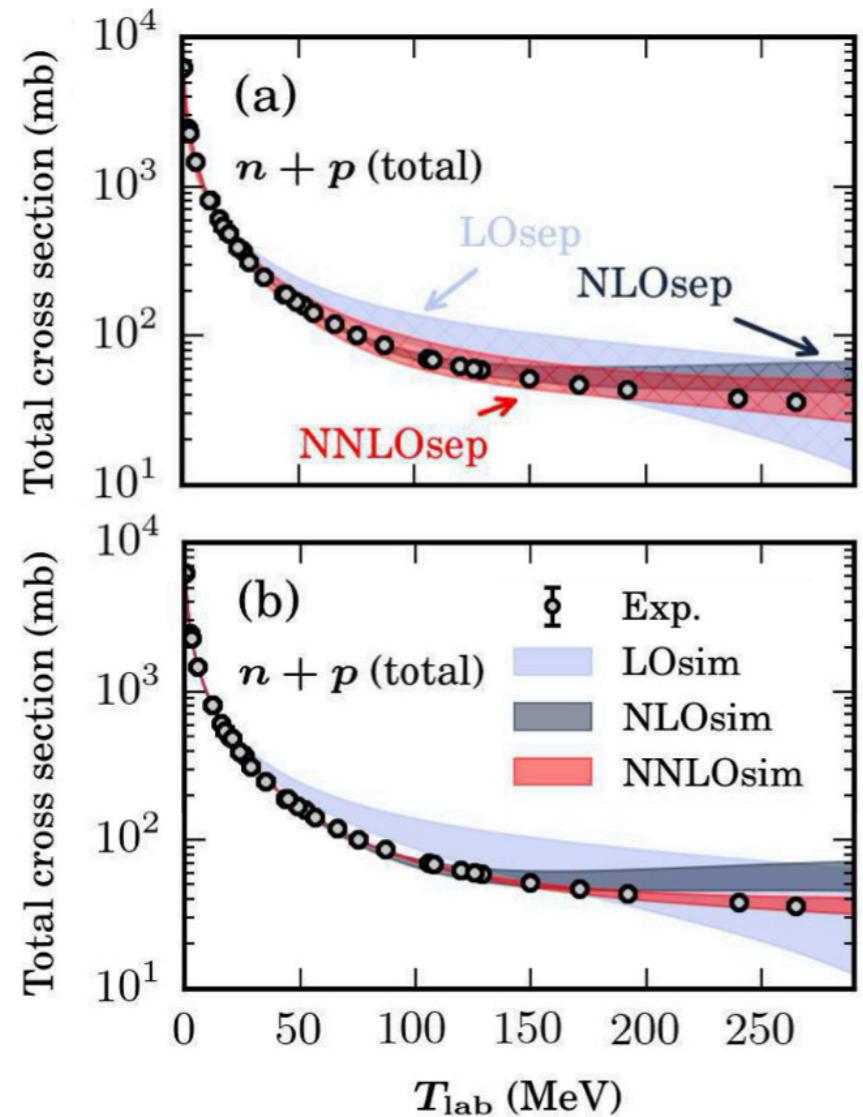
- utilize that electroweak 2b current contributions are proportional to c_D
- Triton beta decay half life much less correlated with E^{3H}
- how to choose the cutoffs consistently in currents and interaction
(continuity equation?)

Determination of LECs: Simultaneous fit of NN and 3N interactions

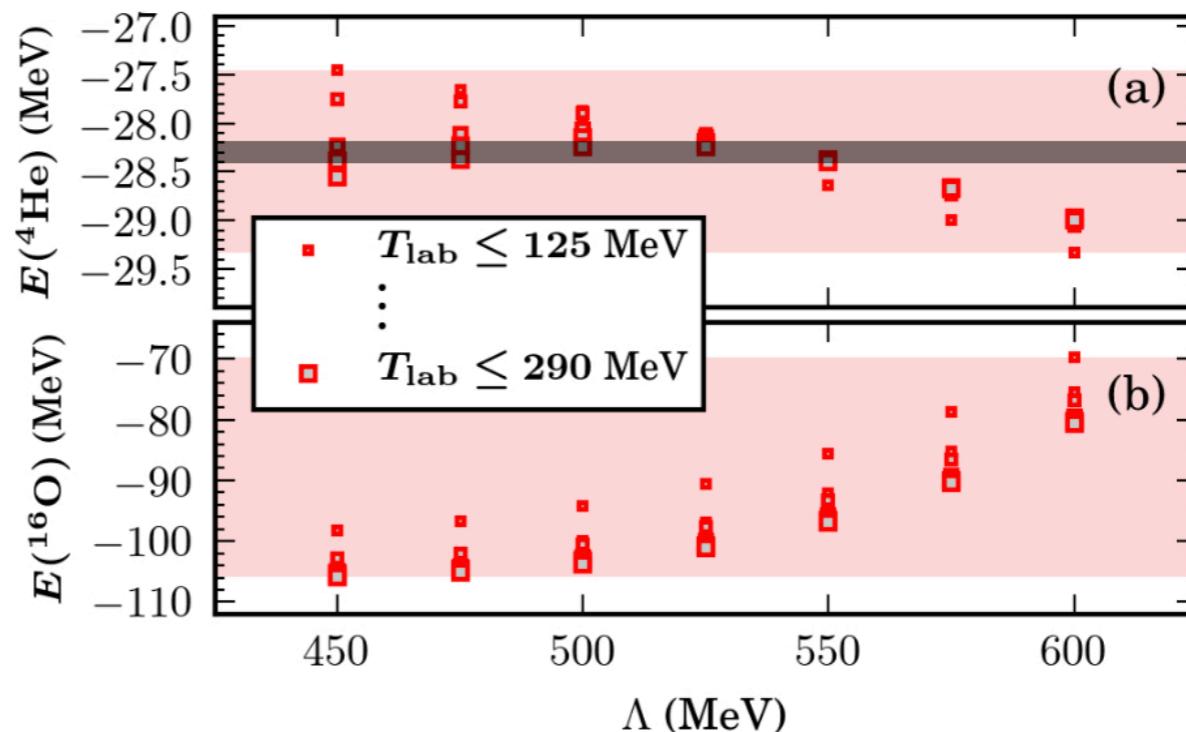


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- Computational challenging due to high dimension of parameter space
- Indications that simultaneous fits lead to more systematic EFT convergence
- Results for heavier systems not consistent with experimental results

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Carlsson et al.,
PRX 6, 011019 (2016)



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