

Exercises

TALENT School “Effective Field Theories in Light Nuclei”

Introduction to Effective Field Theories

This is a collection of exercises which accompany the lecture. The exercises vary in difficulty and length. You may use mathematica for analytical/numerical calculations where appropriate. Discussions are strongly encouraged.

1 Dimensional analysis I: Planck units

Construct the Planck length l_P , the Planck mass m_P , and the Planck time t_P from the fundamental constants G , \hbar and c . Estimate their size in SI units and compare to typical time scales from particle physics. What is the significance of these units?

2 Dimensional analysis II: hydrogen atom

Use dimensional analysis to estimate the size and energy of the ground state of the hydrogen atom. Which dependencies are determined by dimensional analysis and which are not?

3 Dimensional analysis III: photon-photon scattering

Consider photon-photon scattering in quantum electrodynamics. Estimate the amplitude for photon-photon scattering at energies well below the electron mass, $E_\gamma \ll m_e$. How can you obtain the cross section?

4 Naturalness for a square well potential

Consider the attractive radially symmetric square well potential

$$\begin{aligned} V(r) &= -V_0, & r < r_0, \\ &= 0, & r \geq r_0, \end{aligned}$$

where $V_0 > 0$. The S -wave scattering length for scattering of a particle of mass m off this potential is

$$a = r_0 \left(1 - \frac{\tan(\kappa_0 r_0)}{\kappa_0 r_0} \right), \quad (1)$$

where $\kappa_0 = \sqrt{mV_0}$. Assume a constant probability distribution $P(\kappa_0)$ for κ_0 (r_0 fixed) and derive the resulting probability distribution $P(a)$ for a . Use the relation

$$P(a)da = P(\kappa_0)d\kappa_0.$$

For $\kappa_0 r_0 \gg \pi$ the term $\tan(\kappa_0 r_0)$ in Eq. (1) changes for small variations of $\kappa_0 r_0$ much more strongly than $1/\kappa_0 r_0$. The variation of $1/\kappa_0 r_0$ can thus be neglected. Sketch the result and give an interpretation. What the most probable value of a ?

5 EFT for natural scattering length

Consider the non-relativistic effective field theory with the Lagrangian

$$\mathcal{L} = \psi^\dagger \left(i\partial_t + \frac{\nabla^2}{2m} \right) \psi - \frac{g_2}{4} (\psi^\dagger \psi)^2 - \frac{g_{2,2}}{4} \nabla(\psi^\dagger \psi) \cdot \nabla(\psi^\dagger \psi).$$

- (a) Derive the classical equation of motion for the field ψ by using the Euler-Lagrange equations:

$$\partial_t \frac{\delta \mathcal{L}}{\delta(\partial_t \psi)} + \nabla \frac{\delta \mathcal{L}}{\delta(\nabla \psi)} - \frac{\delta \mathcal{L}}{\delta \psi} = 0.$$

- (b) Consider now 2-body scattering in this EFT. In the lecture, the scattering amplitude in leading order ($\nu = 0$) and next-to-leading order ($\nu = 1$) was discussed. Calculate the contributions at order $\nu = 2$. Determine $g_2(\Lambda)$ and $g_{2,2}(\Lambda)$ at this order.

6 Power Counting

Consider an effective field theory with natural interaction terms. In the lecture it was shown that a Feynman diagram with arbitrary 2-particle exchanges for small momenta k scales like k^ν . The following expression was derived for the exponent ν :

$$\nu = 3L + 2 + \sum_i (2i - 2)V_{2i},$$

where L denotes the number of loops and V_{2i} the number of 2-body interactions with $2i$ derivatives. Generalize this expression to include arbitrary natural N -particle interactions. Give the diagrams for $3 \rightarrow 3$ particle scattering in the lowest three orders in ν .

7 Dimer formalism I

Show that the Lagrangian with an explicit dimer field

$$\mathcal{L}' = \psi^\dagger \left(i\partial_t + \frac{\nabla^2}{2m} \right) \psi + \frac{g_2}{4} d^\dagger d - \frac{g_2}{4} (d^\dagger \psi^2 + (\psi^\dagger)^2 d) - \frac{g_3}{36} d^\dagger d \psi^\dagger \psi + \dots, \quad (2)$$

is equivalent to the usual Lagrangian without dimer field

$$\mathcal{L} = \psi^\dagger \left(i\partial_t + \frac{\nabla^2}{2m} \right) \psi - \frac{g_2}{4} (\psi^\dagger \psi)^2 - \frac{g_3}{36} (\psi^\dagger \psi)^3 + \dots,$$

up to N -body interactions with $N \geq 4$. Use the classical equation of motion for the dimer field. Do you obtain the same observables in 3- and 4-body systems with the two Lagrangians?

8 Dimer formalism II

Consider the two-particle problem in the theory with an explicit dimer field using the Lagrangian in Eq. (2).

1. Which diagrams contribute to the full dimeron propagator? How is the 2-body scattering amplitude obtained from the full dimer propagator?
2. Calculate the full dimeron propagator and the 2-body scattering amplitude.

9 Size of the dimer

In the lecture we discussed the universal properties of the weakly bound dimer for $a \gg l$. The wave function for large relative distances r between the two particles is

$$\psi_D(r) = \frac{\mathcal{N}}{r} e^{-r/a}, \quad r \gg l.$$

How large is the dimer?

10 Asymptotic normalization constants

Consider a system with the scattering amplitude

$$\langle \mathbf{k}' | t(E) | \mathbf{k} \rangle = \frac{2\pi}{\mu} \left[\frac{1}{a} - \frac{r}{2} k^2 + ik \right]^{-1} \quad \text{with } E = k^2/(2\mu), \quad (3)$$

where a and r are the scattering length and effective range of the particles while μ is their reduced mass. Near the bound state pole, the scattering amplitude can be expanded about the pole at $E = -\gamma^2/(2\mu)$:

$$\langle \mathbf{k}' | t(E) | \mathbf{k} \rangle = \frac{Z}{E + \gamma^2/(2\mu)} + \text{regular}. \quad (4)$$

In a bound two-body system, the residue Z is connected to the *asymptotic normalization coefficient* (ANC) A of the bound-state wave function. Near the bound state pole, the full Green's function has the general form

$$\langle \mathbf{k}' | \frac{1}{E - H} | \mathbf{k} \rangle = \frac{\psi(\vec{k}') \psi^*(\vec{k})}{E + \gamma^2/(2\mu)} + \text{regular}, \quad (5)$$

where $\psi(\vec{k})$ is the asymptotic wave function for the S -wave bound state, whose co-ordinate space representation is

$$\psi(\mathbf{r}) = \frac{A}{\sqrt{4\pi}} \frac{\exp(-\gamma r)}{r}, \quad (6)$$

and H is the full Hamiltonian. Relate A to Z in order to determine the ANC.

11 Efimov effect

In a system of three identical bosons with large scattering length, $|a| \gg l$, the Efimov effect occurs. It leads to a geometric spectrum of 3-body bound states in the energy interval $1/ma^2 \ll B_3 \ll 1/ml^2$. For two neighbouring states, the relation

$$\frac{B_3^{(n+1)}}{B_3^{(n)}} \approx \exp(-2\pi/s_0), \quad s_0 \approx 1.00624\dots$$

Give an estimate for the number of Efimov states. What happens in the limit $|a| \rightarrow \infty$?