

From Radio Astronomy Interferometry to Table-top Networks for Ultralight Bosons

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Searches for Wave-Like Dark Matter with Quantum Networks



Introduction to Ultralight Bosons

Supermassive Black Holes as Detectors for Ultralight Bosons

Dissecting Ultralight Bosons with Sensor Networks

Prospect

Ultralight Bosons: $\Psi = a, B^\mu$ and $H^{\mu\nu}$

$$-\frac{1}{2}\nabla^\mu a \nabla_\mu a - \frac{1}{4}B^{\mu\nu} B_{\mu\nu} + \mathcal{L}_{\text{EH}}(H) - V(\Psi)$$

- ▶ **Extra dimensions** predict a **wide range of ultralight boson mass**.

Dimensional reduction from higher form fields:

e.g. $g^{MN}(5D) \rightarrow g^{\mu\nu}(4D) + B^\mu(4D)$, $B^M(5D) \rightarrow B^\mu(4D) + a(4D)$.

- ▶ String axiverse/photiverse: **logarithmic mass window**.

In 4D, $m_\Psi \propto e^{-\mathcal{V}_{6D}}$.

- ▶ Ultralight m_Ψ as low as $\sim 10^{-22}$ eV can be naturally predicted.

- ▶ **Coherent waves** dark matter candidates when $m_\Psi < 1$ eV:

$$\Psi(x^\mu) \simeq \Psi_0(\mathbf{x}) \cos \omega t; \quad \Psi_0 \simeq \frac{\sqrt{\rho}}{m_\Psi}; \quad \omega \simeq m_\Psi.$$

Supermassive Black Holes as Detectors for Ultralight Bosons

Superradiance and Gravitational Atom

- ▶ **Gravitational Atom** between BH and axion cloud:

$$\text{BL coordinate : } \Psi^{\text{GA}}(x^\mu) = e^{-i\omega t} e^{im\phi} S_{lm}(\theta) R_{lm}(r), \quad \omega \simeq m\phi + i\Gamma.$$



- ▶ **Rotational and dissipational medium** can amplify the wave around. [Zeldovich 72']
- ▶ **Superradiance** [Penrose, Zeldovich, Starobinsky, Damour et al]: bosons' wave-functions are **exponentially amplified from extracting BH rotation energy** when

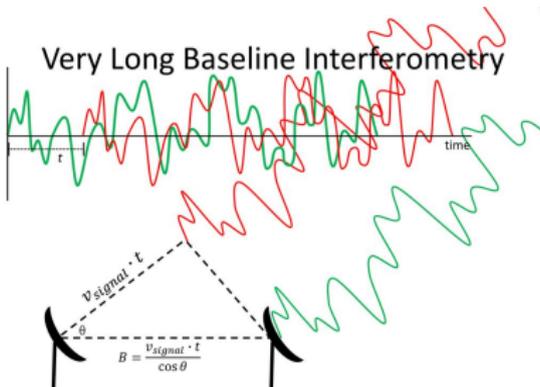
Compton wavelength $\lambda_c \simeq$ gravitational radius r_g .

- ▶ **Supermassive black holes as detectors for ultralight bosons:**

$$M_{\text{BH}} \sim 10^9 M_\odot \leftrightarrow m_\phi \sim 10^{-21} \text{ eV}.$$

Event Horizon Telescope: an Earth-sized Telescope

- ▶ For single telescope with diameter D , the angular resolution for photon of wavelength λ is around $\frac{\lambda}{D}$;
- ▶ VLBI: for multiple radio telescopes, the effective D becomes the **maximum separation between the telescopes**.



- ▶ As good as being able to see



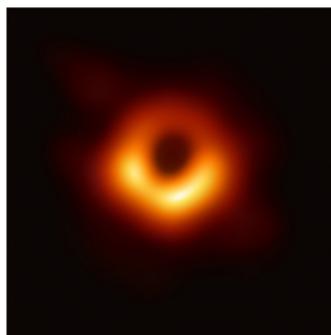
- ▶ on the moon from the Earth.

Event Horizon Telescope: an Earth-sized Telescope

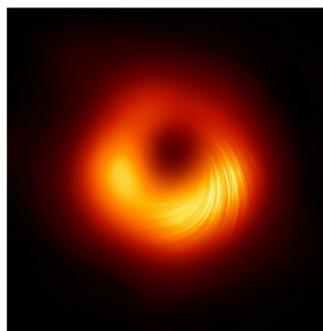
- ▶ For single telescope with diameter D , the angular resolution for photon of wavelength λ is around $\frac{\lambda}{D}$;
- ▶ VLBI: for multiple radio telescopes, the effective D becomes the **maximum separation between the telescopes**.
- ▶ Stokes polarization basis:

$$\begin{pmatrix} I_{IJ} + V_{IJ} & Q_{IJ} + iU_{IJ} \\ Q_{IJ} - iU_{IJ} & I_{IJ} - V_{IJ} \end{pmatrix} \propto \begin{pmatrix} \langle \epsilon_R \epsilon_R^* \rangle_{IJ} & \langle \epsilon_R \epsilon_L^* \rangle_{IJ} \\ \langle \epsilon_L \epsilon_R^* \rangle_{IJ} & \langle \epsilon_L \epsilon_L^* \rangle_{IJ} \end{pmatrix}$$

Total
intensity I



Linear
polarization Q, U



Axion Cloud and Birefringence

- ▶ **Axion cloud** saturates f_a due to self-interactions:



$$r_g m_a \approx \mathcal{O}(1)$$

$$a^{\text{GA}}(x^\mu) \simeq R_{11}(\mathbf{x}) \cos[m_a t - \phi] \sin \theta; \quad a_{\text{max}}^{\text{GA}} \simeq \mathcal{O}(1) f_a; \quad \omega \simeq m_a.$$

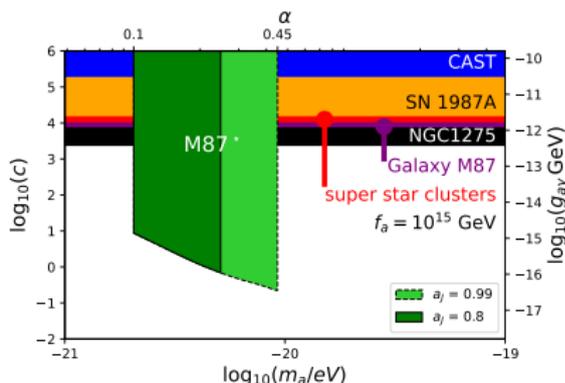
- ▶ **Birefringence** induced from axion-photon couplings:

$$g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} \rightarrow \text{rotate linear polarization orientation } \chi \equiv \arg(Q + i U)/2.$$

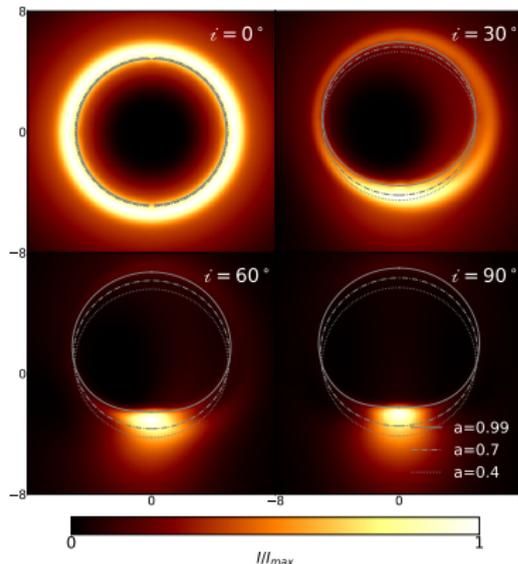
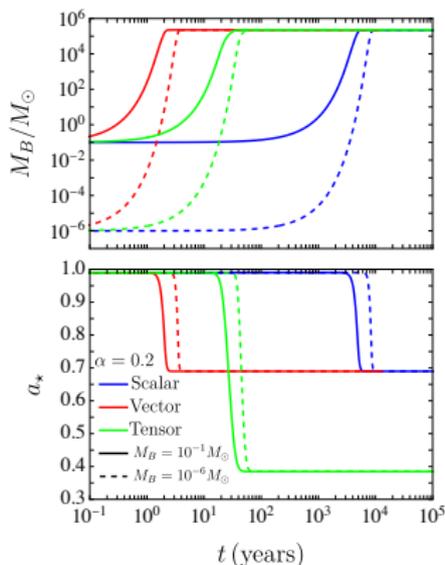
- ▶ **Stringent constraints for $c \equiv 2\pi g_{a\gamma} f_a$** using 21' EHT data: [YC, Liu, Lu, Mizuno, Shu, Xue, Yuan, Zhao, Nature Astronomy 22']

Shift for
each photon:

$$\Delta\chi \approx g_{a\gamma} \times a^{\text{GA}}(x_{\text{emit}}^\mu)$$



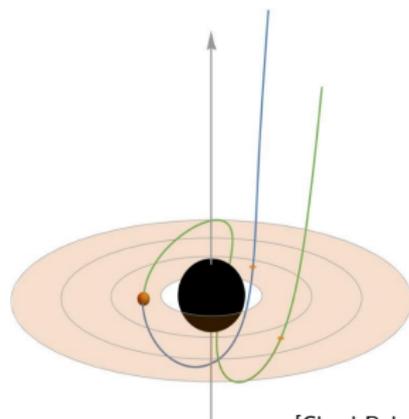
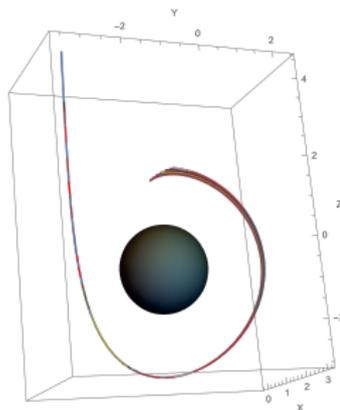
- Superradiant evolution for scalar, vector or tensor:



- Superradiant timescale $\propto M_{BH}$, and is shorter for vector or tensor due to $l=0$ and $j=m=1$ or 2 from intrinsic spin, $\sim \mathcal{O}(10)$ yrs for SgrA*.
- Center of the shadow contour drifts $\sim \mathcal{O}(1)r_g$ once the spin decreases.

Massive Tensor Field (Preliminary)

- ▶ Massive bimetric tensor with minimal coupling $\sim H_{\mu\nu} T_{EM}^{\mu\nu}$.
- ▶ Quadrature mode $H_{\mu\nu} = \text{Re}[R(r)e^{-i\omega t + 2i\phi} \epsilon_{\mu\nu}^R]$ deflects photon geodesics.



- ▶ Astrometry with **photon ring autocorrelation**:
A strain $H \sim 10^{-2}$ leads to **azimuthal lapse oscillation** with $\mathcal{A}_\phi \simeq 5^\circ$.

Dissecting Ultralight Bosons with Sensor Networks

Dark Photon Dark Matter

- ▶ A new **U(1)** vector couples in different portals with SM particles:

$$\epsilon F_{\mu\nu} B^{\mu\nu} + B_\mu \bar{\psi} \gamma^\mu (g_V + g_A \gamma_5) \psi + B_{\mu\nu} \bar{\psi} \sigma^{\mu\nu} (g_M + g_E \gamma_5) \psi.$$

- ▶ Cavity/circuits for kinetic mixing, optomechanics for hidden **U(1)**, spin sensors for dipole couplings...

- ▶ Similar to axion: extra dimensions, misalignment production (or during inflation).

- ▶ Novel aspects: **three polarization degrees of freedom**:

Longitudinal mode: $\vec{\epsilon}_0(\vec{k}) \propto \vec{k}$.

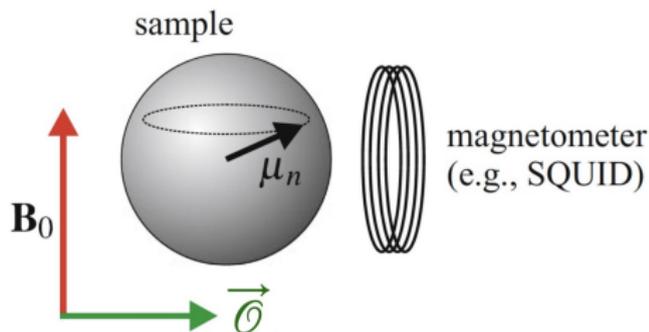
Transverse modes: $\vec{\epsilon}_{R/L} \perp \vec{k}$.

- ▶ Signals projected to the **sensitive direction of a vector sensor**: $\sim \vec{\epsilon} \cdot \hat{l}$.

Spin Precession from Axion Gradient

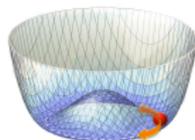
Dipole coupling: $H \propto \vec{\mathcal{O}} \cdot \vec{\sigma}_\psi$.

Effective 'magnetic field' $\vec{\mathcal{O}}$ causes precession of the fermions' spin $\vec{\sigma}_\psi$.
[Graham, Rajendran, Budker et al]



E.g., NMR (Casper), spin-based amplifiers, comagnetometer, magnon ...

▶ **Axion gradient:** $\partial_\mu a \bar{\psi} \gamma^\mu \gamma^5 \psi \rightarrow \vec{\mathcal{O}}_a = \vec{\nabla} a \propto \vec{e}_0$.



▶ **Dark photon with dipole couplings:**

$$B_{\mu\nu} \bar{\psi} \sigma^{\mu\nu} \psi \rightarrow \vec{\mathcal{O}}_{\text{MDM}} = \vec{\nabla} \times \vec{B} \propto \vec{e}_{R/L};$$

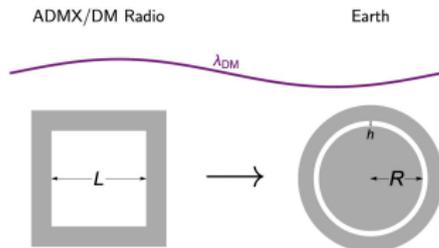
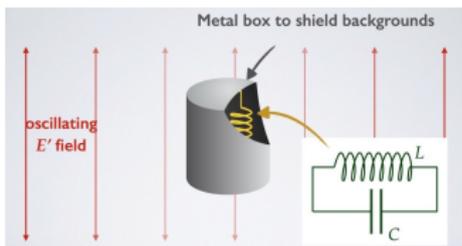
$$B_{\mu\nu} \bar{\psi} \sigma^{\mu\nu} i \gamma^5 \psi \rightarrow \vec{\mathcal{O}}_{\text{EDM}} = \partial_0 \vec{B} - \vec{\nabla} B^0 \propto \begin{cases} \vec{e}, & m \gg |p|, \\ \vec{e}_{R/L}, & m \ll |p|. \end{cases}$$

Identification of the couplings?

Kinetic Mixing and Hidden U(1) Dark Photon

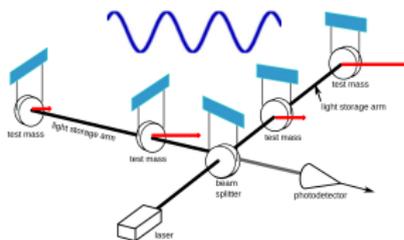
- ▶ **Effective currents:** $\hat{\epsilon} \rightarrow \vec{J}_{\text{eff}}$.

Kinetic mixing U(1) $\sim F_{\mu\nu} B^{\mu\nu}$ shows up in **circuit/cavity**. [Chaudhuri et al 15'] or geomagnetic fields [Fedderke et al 21'];



- ▶ **Force:** $\hat{\epsilon} \rightarrow \vec{F}$.

U(1) B-L & B shows up in **optomechanics** [Graham et al 15', Pierce et al 18'] or **astrometry** [Graham et al 15', Xue et al 19' 21'].



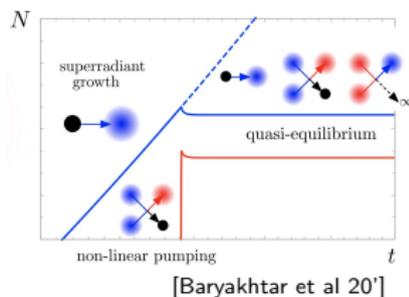
General Axion & Dark Photon Stochastic Background

► **Cosmological isotropic background** [CaB, Dror et al 21']:

Thermal freeze out,
Topological defect decay,
Parametric resonance/tachyonic instability of inflaton,
...

► **Sources from a specific direction:**

Cold stream of dark matter,
Emissions from superradiant clouds.
Dipole radiations from $U(1)'$ charged binaries
...



Microscopic nature: spin, interaction.

Macroscopic property: spectrum, anisotropy and macroscopic polarization.

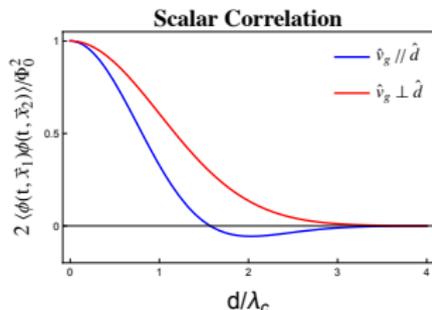
Scalar Field Interferometry

Two point correlation function of the scalar field [Derevianko 18']:

$$\begin{aligned}\langle a(\vec{0})a(\vec{d}) \rangle &= \frac{\rho_a}{\bar{\omega}} \int d^3\vec{v} \frac{f_{\text{DM}}(\vec{v})}{\omega} \cos \left[m_a \vec{v} \cdot \vec{d} \right] \\ &\propto \exp \left[-\frac{d^2}{2\lambda_c^2} \right] \cos \left[m_a \vec{v}_g \cdot \vec{d} \right].\end{aligned}$$

where $f_{\text{DM}}(\vec{v}) \propto \exp\left[-\frac{(\vec{v}-\vec{v}_g)^2}{2v_{\text{vir}}^2}\right]$ and \vec{v}_g is the Earth velocity in the halo.

- ▶ **Velocity fluctuation** $\sim v_{\text{vir}}$ leads to **decoherence at dB length scale**.
- ▶ **Negative correlation** appears when $\vec{d} // \vec{v}_g$.
- ▶ **Localization** with $\sigma_\theta \propto \lambda/d$ and **Daily modulation** due to the self-rotation of the Earth. [Foster, Kahn et al 20']



Vector Sensor Interferometry For Isotropic Backgrounds

A pair of vector sensors separated by a baseline \vec{d} : [YC, Jiang, Shu, Xue, Zeng, PRR 22']

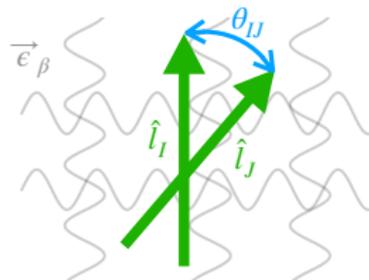
$$\mathcal{F}(\vec{d}, \vec{l}_I, \vec{l}_J) \propto \langle (\vec{O}(t, \vec{x}_I) \cdot \hat{l}_I)(\vec{O}(t, \vec{x}_J) \cdot \hat{l}_J) \rangle, \quad \vec{d} \equiv \vec{x}_I - \vec{x}_J.$$

For isotropic sources $f_{\text{iso}}(\rho, \hat{\Omega}) = \frac{f_{\text{iso}}(\rho)}{4\pi\rho^2}$:

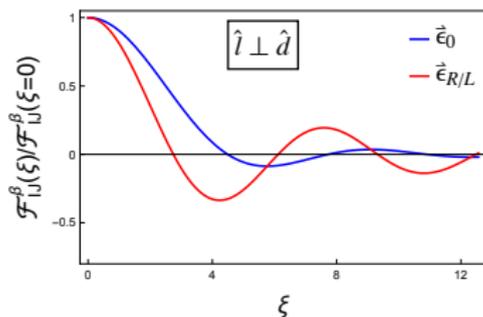
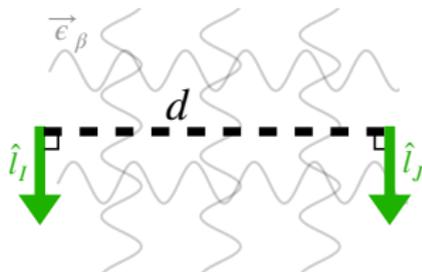
- ▶ **Dipole correlation** for each mode of \vec{e} at $d = 0$.

$$\mathcal{F} \propto \hat{l}_I \cdot \hat{l}_J = \cos\theta_{IJ}$$

Any deviation is a sign of **anisotropy**.



- ▶ **Finite baseline distinguishes \vec{e}_0 from $\vec{e}_{R/L}$** at $\xi \equiv p_0 d \approx 4$.



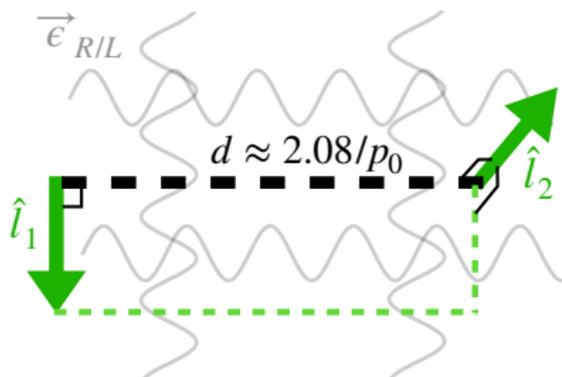
Vector Sensor Interferometry For Isotropic Backgrounds

A pair of vector sensors separated by a baseline \vec{d} : [YC, Jiang, Shu, Xue, Zeng, PRR 22']

$$\mathcal{F}(\vec{d}, \vec{l}_i, \vec{l}_j) \propto \langle (\vec{\mathcal{O}}(t, \vec{x}_i) \cdot \hat{l}_i)(\vec{\mathcal{O}}(t, \vec{x}_j) \cdot \hat{l}_j) \rangle, \quad \vec{d} \equiv \vec{x}_i - \vec{x}_j.$$

For isotropic sources $f_{\text{iso}}(p, \hat{\Omega}) = \frac{f_{\text{iso}}(p)}{4\pi p^2}$:

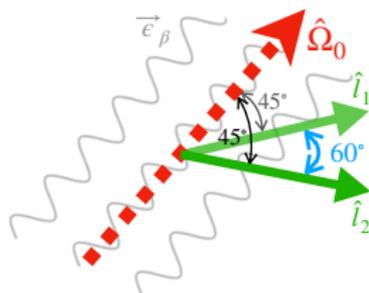
- ▶ A **twisted setup** can **identify the macroscopic circular polarization**.



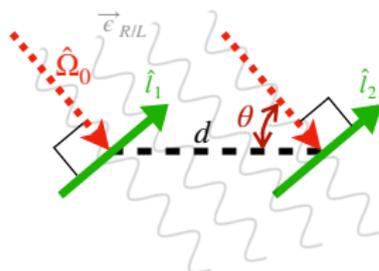
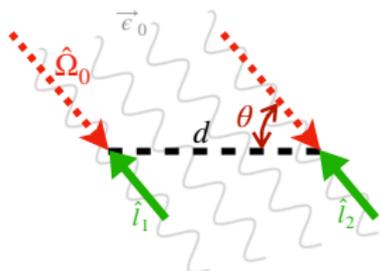
Right and left handed DP respond differently to such setup.

Sources from a specific direction $f_{\text{str}}(\rho, \hat{\Omega}) = \frac{f_{\text{str}}(\rho)}{\rho^2} \delta^2(\hat{\Omega} - \hat{\Omega}_0)$:

- ▶ **Short baseline** limit with $d = 0$:
The optimal arrangements of the sensors are **the same for \vec{e}_0 and $\vec{e}_{R/L}$** , reaching $\sigma_{\Omega} \approx 1/\text{SNR}$.



- ▶ **Long baseline** limit:
The sensitive directions should **overlap with the signals as much as possible** with $\sigma_{\theta} \approx 1/(\text{SNR} \rho d)$.



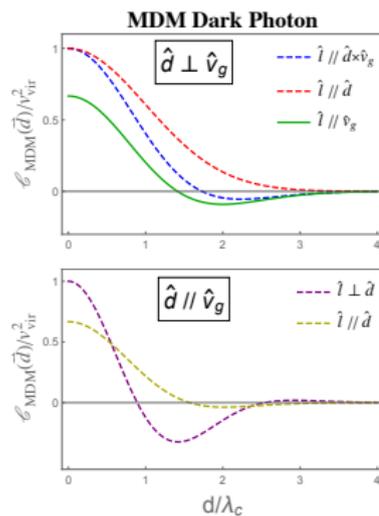
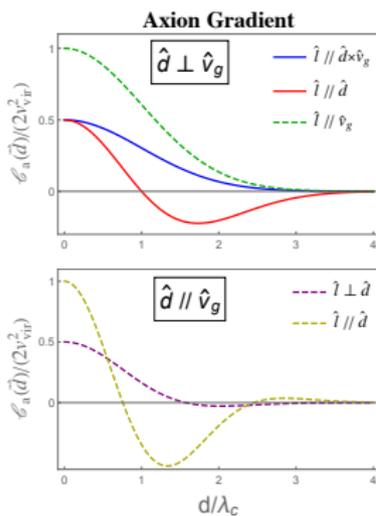
Multi-messenger astronomy with **GNOME** [Dailey et al 21']!

Axion Gradient and MDM DP Dark Matter

3 × 3 matrix of vector correlation: $\mathcal{C}(\vec{d})_{IJ} \propto \langle (\vec{O}(t, \vec{x}_I) \cdot \hat{l}_I) (\vec{O}(t, \vec{x}_J) \cdot \hat{l}_J) \rangle$ with $f_{\text{DM}}(\vec{v}) \propto \exp[-(\vec{v} - \vec{v}_g)^2 / (2v_{\text{vir}}^2)]$.

- ▶ 5 possibilities when two \hat{l}_i align: [YC, Jiang, Shu, Xue, Zeng, PRR 22']

$$\vec{O}^a \propto \vec{E}_0$$



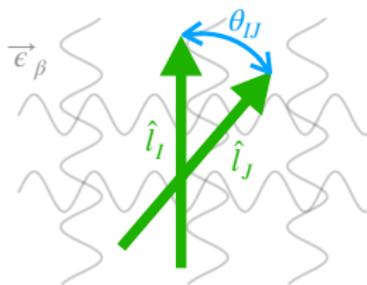
$$\vec{O}^{\text{MDM}} \propto \vec{E}_{R/L}$$

- ▶ Axion and MDM DP **have totally different spatial correlations.**

Dipole Angular Correlation [YC, Jiang, Shu, Xue, Zeng, PRR 22']

For $f_{\text{DM}}(\vec{\nu}) \propto \exp[-(\vec{\nu} - \vec{\nu}_g)^2 / (2v_{\text{vir}}^2)]$,

Tune \vec{l}_1 and \vec{l}_2 with **certain directions at the same location**:



$$\begin{aligned}\Gamma(\vec{l}_1, \vec{l}_2) &= (\vec{l}_1)^T \cdot \mathcal{C}(0) \cdot \vec{l}_2 \\ &= \begin{cases} \frac{v_{\text{vir}}^2}{2} \vec{l}_1 \cdot \vec{l}_2 + \frac{1}{2} (\vec{l}_1 \cdot \vec{v}_g) (\vec{l}_2 \cdot \vec{v}_g) & \text{Axion Gradient;} \\ \frac{v_{\text{vir}}^2}{2} \vec{l}_1 \cdot \vec{l}_2 - \frac{1}{6} (\vec{l}_1 \cdot \vec{v}_g) (\vec{l}_2 \cdot \vec{v}_g) & \text{MDM DP;} \\ \frac{1}{6} \vec{l}_1 \cdot \vec{l}_2 & \text{EDM DP.} \end{cases}\end{aligned}$$

- ▶ **Universal dipole angular correlation:** $\vec{l}_1 \cdot \vec{l}_2 = \cos \theta$, in contrast with monopole or quadrupole (H.D. curve) for stochastic GW searches.
- ▶ \vec{v}_g brings in anisotropy, with different signs for axion gradient and MDM DP.

- ▶ **Correlations of vector sensors** can identify the macroscopic property and the microscopic nature of the bosonic background:

Coupling type, macroscopic polarization and localization/anisotropy ...

→ **Multi-messenger astronomy/cosmology!**

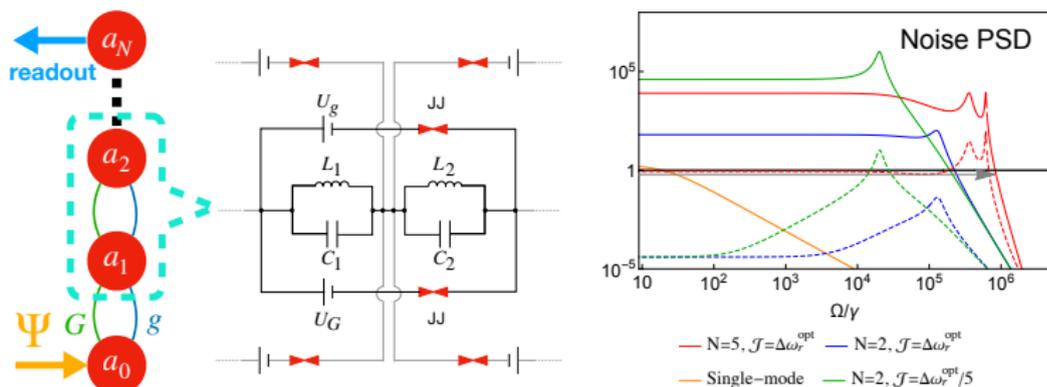
- ▶ How to improve sensitivity based on those information?
- ▶ Prospect: sensitivity of multi-mode systems, tensor-like correlations.

Simultaneous Resonant and Broadband Detection

- ▶ Standard quantum limit for power law detection:

$$\text{SNR}^2 \propto \text{range where } S_{\text{int}} \gg S_r. \quad [\text{Chaudhuri, Irwin, Graham, Mardon, 19'}]$$

- ▶ Scan bandwidth can be significantly increased in a multi-mode system:

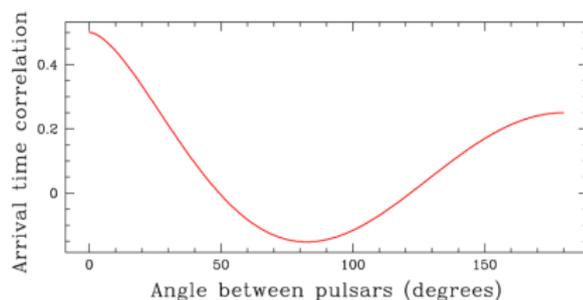


- ▶ New quantum limit for multi-mode resonators. [YC, Liu, Shu, Song, Yang, Zeng, 22']
[YC, Jiang, Ma, Shu, Yang, PRR 22']

Tensor-like Angular Correlations

- ▶ Pulsar timing array for stochastic GW background:

Angular correlation shows Hellings-Downs curves.



- ▶ Microscopic tensor nature shows up in macroscopic correlations.
- ▶ On-going: Hellings-Downs in table-top detectors? e.g., high frequency GW or massive tensor dark matter.

Thank you!

Appendix

Property of Ultralight Dark Matter

Galaxy formation: virialization $\rightarrow \sim 10^{-3}c$ velocity fluctuation, thus kinetic energy $\sim 10^{-6}m_\Psi c^2$.

Effectively coherent waves:

$$\Psi(\vec{x}, t) = \frac{\sqrt{2\rho_\Psi}}{m_\Psi} \cos\left(\omega_\Psi t - \vec{k}_\Psi \cdot \vec{x} + \delta_0\right).$$

▶ Bandwidth: $\delta\omega_\Psi \simeq m_\Psi \langle v_{\text{DM}}^2 \rangle \simeq 10^{-6}m_\Psi$, $Q_\Psi \simeq 10^6$.

▶ Correlation time: $\tau_\Psi \simeq \text{ms} \frac{10^{-6}\text{eV}}{m_\Psi}$.

Power law detection is used to make integration time longer than τ_Ψ .

▶ Correlation length: $\lambda_d \simeq 200 \text{ m} \frac{10^{-6}\text{eV}}{m_\Psi} \gg \lambda_c = 1/m_\Psi$.

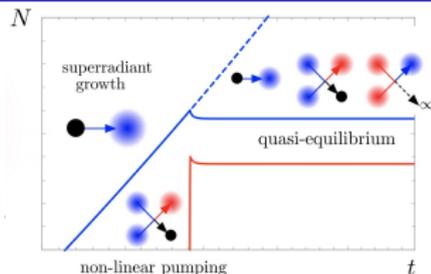
Sensor array can be used within λ_d .

Axion Wave from Saturating Axion Cloud

- ▶ **Self interaction saturating phase**

where $a_{\max} \simeq f_a$.

[Yoshino, Kodama 12', Baryakhtar et al 20']



- ▶ Two level state with 2, 1, 1 and 3, 2, 2. Annihilations between 3, 2, 2 lead to **'ionized' axion wave** with velocity $v \sim \alpha/6$:

$$B_a \simeq 3 \times 10^{-24} \text{ T} \times C_N \left(\frac{\alpha}{0.1} \right)^4 \left(\frac{1\text{kpc}}{r} \right), \quad [\text{Baryakhtar et al 20'}]$$

- ▶ For BH $\sim 10M_\odot$, superradiance happens for $m_a \sim 100$ Hz axion. **Axion gradient/DP signals are expected!**

Multi-messenger astronomy with GNOME, ngEHT and PTA!

- ▶ **Localization of the source ?**

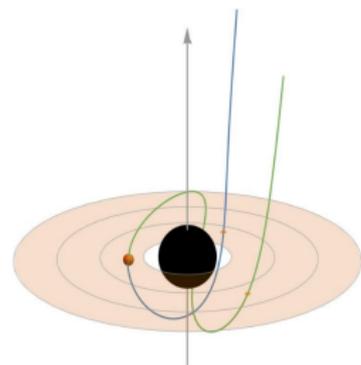
Azimuthal Lapse

- ▶ At low inclination angles,

photon ring autocorrelation:

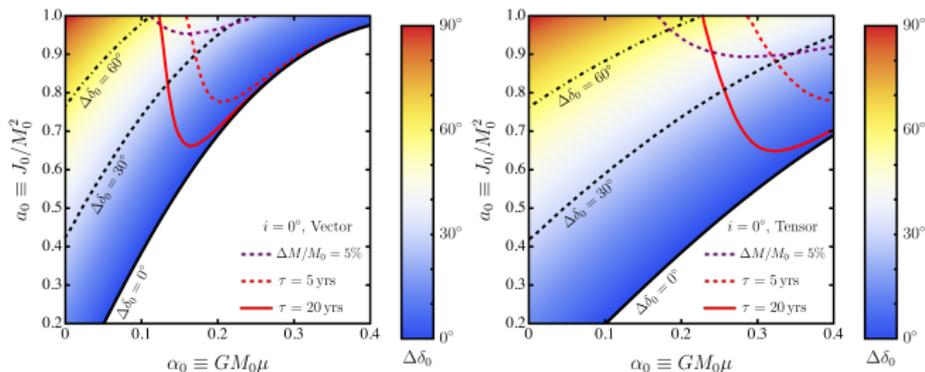
$$\mathcal{C}(T, \varphi) \equiv \iint dr dr' r r' \langle \Delta I(t, r, \phi) \Delta I(t+T, r', \phi+\varphi) \rangle$$

peaks at $T = \tau_0$ and $\varphi = \delta_0$,
where δ_0 is the azimuthal lapse.



[Chael Palumbo]

- ▶ δ_0 is sensitive to spin evolution due to frame dragging.



- ▶ Axion-DP coupling:

$$\frac{1}{2}\partial_\mu a \partial^\mu a - m_a^2 f_a^2 \left[1 - \cos\left(\frac{a}{f_a}\right)\right] - \frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{\alpha}{4f_a}aB_{\mu\nu}\tilde{B}^{\mu\nu}.$$

- ▶ Rolling a leads to **different dispersions between R/L -handed dark photon**:

$$\omega_{L/R}^2 = p^2 \mp p \frac{\alpha}{f_a} a'.$$

- ▶ Tachyonic instability: **exponential increase** of mode with **negative ω^2** .
- ▶ Potential **chiral spectrum**. **How to identify the macroscopic circular polarization?**

Global Gravitational Wave Detector Network

- ▶ **Localization** due to **long baseline** $\sigma_\theta \propto \lambda_h/R_E$.
- ▶ **Macroscopic polarization** from **correlation of detectors**.

