A novel mixed species In⁺/Yb⁺ ion optical clock

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The Team "Quantum Clocks and Complex Systems"



Int. Collaborations:

NICT Toyko (J) University of Osaka (J) CMI (Prag, Cz) NPL (London, UK) W. Zurek (Los Alamos NL) R. Nigmatullin (Uni Sydney, Au) Haggai Landa (IBM, IL)



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exchange with Osaka and Tokyo

TRIAC International Joint Laboratory for

International Joint Laboratory for Trapped-Ion Integrated Atomic-Photonic Circuits



"Quantum Clocks and Complex Systems"



建筑 法法律 网络名法法

In⁺ /Yb⁺ crystals – optical clock

Test of Local Lorentz Invariance & Isotope Shifts/search for new bosons

Herausforderung Quantentechnologie





Integrated & Scalable Traps \rightarrow Nanophotonics → Quantum Information

Ion Coulomb Crystals

if I may ...



Hot Ion Plasma/Liquid

 $E_{kin} > E_{pot}$!



Self-organized system! Nonlinear chaotic dynamics

Ion Coulomb Crystals, $E_{pot} > E_{kin}$





Phases in Ion Coulomb Crystals



Lucretius "De rerum naturalium" / Proof according to Epicurus (300 B.C.)

The dilemma of argumentation in his physics lies in the fact that on one hand it wants to be based completely on the perception,

but on the other hand it denies its world explanation directed against the superstition above all with the help of its doctrine of the invisible (more exactly: not perceptible) atoms.

This situation makes it necessary to develop reliable methods with the help of which conclusions can be drawn from the visible to the invisible.

3 Forms of empirical cognition:

- Direct perception
- Based on certain (empirically gained) 'ideas' (induction)
- Logical ways of inferring from phenomena (argumentatively) to the unknown
 - \rightarrow Conclusion by analogy

What happens when a system changes from one equilibrium condition to another?

- Examples for 2nd-order phase transitions:
 - ferro-magnetism \rightarrow para-magnetism
 - metal \rightarrow superconductor
 - early universe



Symmetry breaking phase transitions

What happens when a system changes from one equilibrium condition to another?

- Examples for 2nd-order phase transitions:
 - ferro-magnetism \rightarrow para-magnetism
 - metal \rightarrow superconductor
 - early universe

Spontaneous symmetry _____ breaking of Higgs field



Nature Physics 7, 2 (2011)

Symmetry breaking in ion Coulomb crystals



The Kibble-Zurek Mechanism

- ξ : correlation length
- τ : relaxation time
- defect density: $d = \frac{\xi}{system \ size}$

•
$$\tau \sim \left|\frac{\tau_0}{t}\right|^{\upsilon z}$$

• $\xi \sim \left|\frac{\tau_0}{t}\right|^{\upsilon}$

- KZM: d from ξ at freezeout
- friction (laser cooling) negligible $\implies v = \frac{1}{2}, z = 1$

test of KZM with defined ν , z

del Campo et al., PRL 105, 075701 (2010) Fishman et al., PRB 77, 064111 (2008)

Quenching the control field



Harmonic Ion Traps – Inhomogeneous Case

 Ions in harmonic potential: phase transition spreads out from center!



• Phase front faster than speed of sound!

A. Del Campo et al., Phys. Rev. Lett. 105, 075701 (2010)

The Kibble-Zurek Mechanism

Prediction of KZM

Power law scaling of defect density:

$$d \sim \left|\frac{1}{\tau_Q}\right|^{\nu/(1+\nu z)}$$

test of KZM with defined v, z



Scaling of Defect Creation

Defect Probability as Function of Ramp Velocity:



Nikoghosyan et al., Universality in the dynamics of second-order phase transitions, PRL 116, 080601 (2016)

Puebla et al., Fokker-Planck formalism approach to Kibble-Zurek scaling laws and nonequilibrium dynamics, PRB 95, 134104 (2017)

Topological defects in ion Coulomb crystals

Localized Defect $v_{rad}/v_z \approx 8$



Extended Defect $v_{rad}/v_z \approx 5.5$

Pyka et al., Nat. Commun. 4, 2291 (2013)

Order Parameter Φ

radial ion separation a:



Topological Soliton:

Langragian of Scalar Field:

$$L = \frac{1}{2} \left(\frac{\partial \Phi}{\partial t} \right)^2 - \frac{1}{2} \left(\frac{\partial \Phi}{\partial z} \right)^2 + \lambda \Phi^2 + A \Phi^4$$

- field configurations, such that their presence can be detected by looking at the values of the field far away from the defect
- cannot be removed by local deformations of the field



The Aubry Phase Transition

Friction at the Atomic Scale



Complex, Self-Organized System with Back-Action

How can we experimentally realize nanofriction?

Vanossi *et al.,* Modeling friction: From nanoscale to mesoscale, Rev. Mod. Phys. 85 (2013) **Proposal: use linear ion chain in optical corrugation potential**

 Theory:
 Benassi et al., Nature Commun 2, 236 (2011)

 Puttivarasin et al., New J. Phys. 13, 075012 (2011)

 Fogarty et al., Phys. Rev. Lett. 115, 233602 (2015) ...



Standing optical wave in cavity

here: Frenkel-Kontorova model \rightarrow gives sine-Gordon equation

 $\mathcal{I} = \frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 + A \sin \phi$

Self-organized crystal with back-action

Locally the kink disturbes the quasi-periodicity of the two atomic layers



 $D = m \omega_{ax}^2$: ion trap $\kappa \sim \frac{e^2}{\pi \epsilon_0} \frac{1}{ma^3}$: ion interaction



Interaction between the ion layers gives corrugation potential!

Order parameter \rightarrow parametrizes symmetry breaking

Order parameter Φ := relative distance to closest ion in other layer



Order parameter \rightarrow parametrizes symmetry breaking



Experimental observation:

Peierls-Nabarro Barriers

= potential energy of kink soliton



First Experimental Observation of the Soft Mode



legend:
Hessian matrix T = 0K
simulations T = 5μK
▲ simulations T = 1mK
▼ exp. data

PN-energy barriers of kink soliton:

 \rightarrow thermal energy allows for switching



Yes, but strong non-linearities due to finite temperature!

How does a soft mode couple to thermal phonon environment?

Finite temperature spectrum at the symmetry-breaking linear-zigzag transition

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 (Dated: December 22, 2020)



Fishman et al., PRB 77, 064111 (2008)

Coupling of soft mode to thermal phonon environment

Equation of motion of phonon *i* under non-linear interaction:

$$V_4 = \frac{1}{2!} \sum_{i=1}^{3N} m\omega_i^2 \Theta_i^2 + \frac{1}{3!} \sum_{ijk=1}^{3N} L_{ijk} \Theta_i \Theta_j \Theta_k + \frac{1}{4!} \sum_{ijkl=1}^{3N} M_{ijkl} \Theta_i \Theta_j \Theta_k \Theta_l$$

$$\ddot{\Theta}_i = -\frac{1}{m} \frac{\partial V_4}{\partial \Theta_i} - \gamma_i \Theta_i + \Xi_i(t)$$

Time average over higher frequency phonons:

$$\ddot{\Theta}_{1} = -\tilde{\omega}_{1}^{2}\Theta_{1} - \frac{1}{2}\nu_{12}^{2}\Theta_{2} + \eta_{1} - \gamma_{1}\Theta_{1} + \Xi_{1}$$

$$\ddot{\Theta}_{2} = -\tilde{\omega}_{2}^{2}\Theta_{2} - \frac{1}{2}\nu_{12}^{2}\Theta_{1} + \eta_{2} - \gamma_{2}\Theta_{2} + \Xi_{2},$$

with
$$\tilde{\omega}_i(T)^2 = \omega_i^2 + \nu_{\text{eff},i}^2 T$$
 $\nu_{\text{eff},i}^2 = \frac{1}{2m} \sum_{k \neq 1,2} M_{iikk} \frac{k_B}{m\omega_k^2}$

Theory for two lowest frequency modes:



dashed lines = new theory

Coupling of soft mode to thermal phonon environment



Theory for two lowest frequency modes:

Conclusion:

also at Linear-to-Zigzag Transition...

- → soft mode is modified due to non-linear coupling to thermal phonon environment
- → soft-mode sees modified Coulomb environment due to oscillating phonon environment
- \rightarrow Floquet physics



J. Kiethe, L. Timm, H. Landa, D. Kalincev, G. Morigi, T. E. Mehlstäubler <u>Finite-temperature spectrum at the symmetry-breaking linear to zigzag transition</u> Phys. Rev. B, **103**, 104106 (2021)

Is there quantum nanofriction ?

Tunneling over 10s of micrometer? \rightarrow collective excitation! \rightarrow ions move by only a few nm while top. defect moves by 10s µm



Define effective mass of the quasi particle:

$$M_{\text{eff}}(X) = m \sum_{i} \left(\frac{d\vec{r}_{i,C}(X)}{dX}\right)^2$$

Solve Hamiltonian for quasi particle:

$$\hat{H}_s = \hat{P} \frac{1}{2M_{\text{eff}}(\hat{X})} \hat{P} + U(\hat{X})$$

L. Timm, L. A. Rüffert, H. Weimer, L. Santos, T. E. Mehlstäubler Quantum nanofriction in trapped ion chains with a topological defect, Phys. Rev. Research, 3, 043141 (2021)

Quantum nanofriction





Sliding phase:

QUANTUM WORLD

Assumption: Kink = quantum particle in classical PN potential

- Harmonic eigenstates
- Equidistant spectrum

Tunneling regime:

- Barrier splits ground state wavefunction
- Sym. and antisym. pairs

Quasi-classical regime:

- Large barrier
 →tunneling negligible
- Localized states

Consequences for observables in experiments?

Quantum nanofriction

ENERGY SPECTRUM



- Quantum fluctuations: No "softmode"
- Tunneling regime: Eigenstate energy below barrier height but tunneling causes energy gap
- Quasi-classical regime: Degenerate sets of eigenstates, localized left/right

→ Spectroscopic measurement of the energy spectrum after ground-state cooling

TUNNELING DYNAMICS



for T = $0.1 \,\mu\text{K}$

- Initialize a localized state in one potential minimum and monitor evolution
- Oscillation between left and right minimum due to tunneling
- Not observable in quasi-classical regime

\rightarrow At low μ K temperatures tunneling dynamics is observed

Optical Clocks and Trapped Ions

Accuracy of Single Ion Clocks



- 10^4 K deep traps \rightarrow **long trapping times** for single ion (**up to months**)
- ions are trapped at $E = 0 \rightarrow$ no systematic shifts to 1st order
- strong trap potential \rightarrow strong localization ($\sigma \approx nm$)
- high level control of internal (pseudo-spin) & external degrees of freedom (bosonic degree of motion)
- Laser-cooling to mK \rightarrow resolved sideband-cooling to quantum mechanical ground-state of motion!

Current world record in clock accuracy: *Brewer et al., PRL 2019:* accuracy $\Delta v/v = 9.4 \times 10^{-19}$

Accuracy of Single Ion Clocks



3D-Paul trap



single Yb+-ion

Trap Depth ~ 10⁴ K





ponderomotive potential



Nobel Prize 2012:

"for groundbreaking experimental methods, that allow to manipulate and measure single quantum systems."





Test of Fundamental Physics

Clocks as Quantum Sensors for Geodesy



 $\Delta \nu / \nu = 10^{-18} \rightarrow 1 \text{ cm}$ height resolution



gravimeter

→ measures acceleration g= local gradient of potential





clocks

Earth's geoid/ESA

 \rightarrow measure difference in potential

Examples for Geodetic Observations



Example: impact of groundwater body on gravity acceleration δg and gravitational potential δN



S ₀	δg	δΝ
[m]	[nm/s²]	[mm]
10	77.1	0.0001
100	131.5	0.0014
1 000	137.7	0.0141
10 000	138.3	0.1400
100 000	138.3	1.3650
000 000	138.3	10.7170

Resolution of FG5 gravimeter: 10 nm/s² ≈ 10⁻⁹ g

Review article: Mehlstäubler et al., "Atomic Clocks for Geodesy", Rep. Prog. Phys. 81, 6 (2018)

More about Time Dilation...

Bloomberg TV "Atomic Clocks are Reinventing Time"

Featuring Don Lincoln (**Fermi Lab**), David Hume (**NIST**) and Tanja Mehlstäubler (**PTB/LUH**).

On Bloomberg TV and YouTube: https://www.youtube.com/watch? v=hzLTgtFaPLY



Ingredients for an optical clock



State-of-the-art with dedicated ¹¹⁵In⁺ set-up

- Self-built lasers at 230 nm and 946 nm (→236 nm) operational for direct In⁺ detection and clock spectroscopy
- In⁺ clock laser (30 cm ULE) operational at 1.1 x 10⁻¹⁶ in 1s *
- drift compensation: low BW transfer-lock to cryogenic Si cavity (PTB)





*Didier et al., Optics Letters 44,1781 (2019)

Instability: statistical error

Mr M. M. M. M. An

How well can we resolve the atomic frequency?



Multi-Ion Approach

New Quantum Clocks ?

<u>Now needed</u>: "experimental methods, that allow to manipulate and measure many-body quantum systems."



Quantum Metrology \leftrightarrow **Quantum Simulation & Information**

Herschbach et al., Appl. Phys. B 107, 891 (2012)

New Quantum Clocks ?



Quantum Metrology \leftrightarrow **Quantum Simulation & Information**

Herschbach et al., Appl. Phys. B 107, 891 (2012)

Scaling up the number of ions for metrology

• Basis for multi-ensemble clocks and new interrogation protocols¹ - in one system

our approach: bottom up, scalable ion traps



1) Rosenband and Leibrandt, arXiv:1303.6357 (2013), Borregaard and Sorensen, PRL 111, 090802 (2013), Kessler et al., PRL 112 (2014), Lebedev et al., PRA 89 (2014)

High-Precision Scalable Chip Ion Traps – AIN

- Laser cut and structured: high-precision, µm-tolerances
- P = 10⁻¹¹ mbar, non-magnetic materials, excellent heat management
- Controlled micromotion, low heating rates, 3D high optical access



Dolezal et al., Metrologia (2015), Keller et al., PRA 99, 1 (2019), ...

Ion traps built for other labs

AIN and Rogers chip traps produced by PTB



Trap for the group of **Prof. Kjeld Eikema (FU Amsterdam)**. The trap will be used for He⁺ 1S-2S spectroscopy to test bound-state quantum electrodynamics and to search for new physics beyond the standard model.



Traps for the group of **Prof. Tanja Mehlstäubler (PTB/LUH)**. 1. A multi-ion clock where In⁺ sympathetically cooled with Yb⁺ ions. 2. A multi-ion system with Yb⁺ ions to study topological defects and many-body physics.



Trap for the group of **Prof. Andreas Schell (LUH)**. The trap will be used for levitation of nanoparticles to investigate if they behave in a classical way or in a quantum mechanical one.



Trap for **Dr. Nils Huntemann (PTB)**. The trap is used for a Yb⁺ ion clock where the ions are sympathetically cooled with Sr^+ ions.

Pysikalisch-Technische Bundesanstat Braunschweig und Berlin

Trap for **Dr. Stephan Hannig (DLR)**. The trap will be used as a portable Al^{*} clock.

Trap for the IDEAL project, **ongoing development** of a chip trap with integrated micro-optics.

Trap for **Prof. Piet Schmidt (PTB/LUH)**. It is a Rogers chip trap based on the PTB design for a Al^*/Ca^* clock. [4]



Trap for the Opticlock project, the trap is in the group of **Prof. Christof Wunderlich** (Universität Siegen). The trap is for a portable multi-ion clock.

Ion Trap Production at PTB Clean Room



Industrial grade fs-laser μm precision machining





Ti and Au sputtering



Precision measurement microscope



Fineplacer[®], precision < 1 μm

• Flip-Chip-Bonding

- Thermo-Compression
- Ribbon Bonding
- Wire Bonding





Future \rightarrow CGPA Size Optically Integrated Ion Traps (incl. detection)

For Quantum Computers and Clocks → for scalability > 50 qubits



Atomic "Clock on a Chip"

• Collaboration: Karan Mehta (ETH, now Cornell); production at LioniX (NL)





Precision Spectroscopy in Trapped Ions

Our atomic candidates:

¹¹⁵In⁺ ...a first *"simple"* approach:

- Large mass, m=115: Time dilation: $\Delta v/v = -E_{kin}/mc^2 = -\overline{v}^2/2c^2$
- Low blackbody shift^[1]: $\Delta v/v = (1.36 \pm 0.1) \times 10^{-17} @ T = 300 K$
- 2nd order Zeeman shift: 4.1 Hz/mT² (²⁷Al⁺: 72 Hz/mT², ¹⁷¹Yb⁺: 52 kHz/mT²)
- Directly detectable transition to ³P₁! direct cooling: T < 100 μK





N. Herschbach et al., Appl. Phys. B 107, 891 (2012)

Our atomic candidates:

- Sympathetic cooling of In⁺ via Yb⁺ ions
- Spatially resolved detection with ROIs on CCD

n

Yb+

+

¹¹⁵In⁺ ...1st order free of quadrupole shift



2019: Measurement based multi-ion uncertainty budget



Keller et al., "Controlling systematic frequency uncertainties at the 10⁻¹⁹ level in linear Coulomb crystals", PRA 99, 1 (2019)

First Step: Development of Scalable Chip Ion Traps

...since dominant uncertainties of optical ion clocks are trap related

- 1) Very low trap warming
- → trap-related BBR shift uncertainty <10⁻¹⁸



Nordmann et al., *Rev. Sci. Instrum.* **99**, 11301 (2020) Dolezal et al., *Metrologia*, **52**, 842-856 (2014);

 2) Low axial excess micromotion across all trapping regions
 → time dilation shifts <10⁻¹⁸



"Probing Time Dilation in Coulomb Crystals in a high-precision Ion Trap", Keller et al., *Phys. Rev. Appl.* 11, 011002 (2019) 3) Low heating rate of <0.5 phonons/s at 1MHz



Kalincev et al., *Quantum Sci. Technol.* 6, 034003 (2021) Brownnutt et al., *Rev. Mod. Phys.* 87, 1419 (2015)

Atomically resolved micromotion of individual ions





Crystal aligns at RF nodal axis; agrees with single ion probe

"Probing Time Dilation in Coulomb Crystals in a high-precision Ion Trap", Keller et al., Phys. Rev. Appl. 11, 011002 (2019)

Atomically resolved micromotion of individual ions



"Probing Time Dilation in Coulomb Crystals in a high-precision Ion Trap", Keller et al., Phys. Rev. Appl. 11, 011002 (2019)

For optimal cooling: Automatized sorting of In⁺/Yb⁺ crystals





6 Yb⁺ and 15 In⁺



6 Yb⁺ and 16 In⁺



Clock scheme of ¹¹⁵In⁺ / ¹⁷²Yb⁺ clock







Clock scheme of ¹¹⁵In⁺ / ¹⁷²Yb⁺ clock

1.0 -

+9/2->+9/2 -9/2->-9/2

Probing on $-9/2 \rightarrow -9/2$ and $+9/2 \rightarrow +9/2$ during one clock cycle to cancel 1st order Zeeman shift



Frequency Ratio Measurement ¹¹⁵In⁺ - Sr - ¹⁷¹Yb⁺



→ EMPIR project "ROCIT"

Three corner hat comparison: Frequency ratio In⁺ -Sr- ¹⁷¹Yb⁺



\rightarrow Sr - ¹⁷¹Yb⁺ mit u = 2.5 x 10⁻¹⁷

Long-term variation f(Sr)/f(Yb⁺) is larger than expected from individual uncertainties

Dörscher et al., Metrologia 58 015005 (2021)

The importance of frequency ratio measurements!



Following results are preliminary! ...not published yet

The Team "Quantum Clocks and Complex Systems"



DFG

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^{exchange} with Osaka and Tokyo