

# Intensity Interferometry for Ultralight Bosonic Dark Matter Detection

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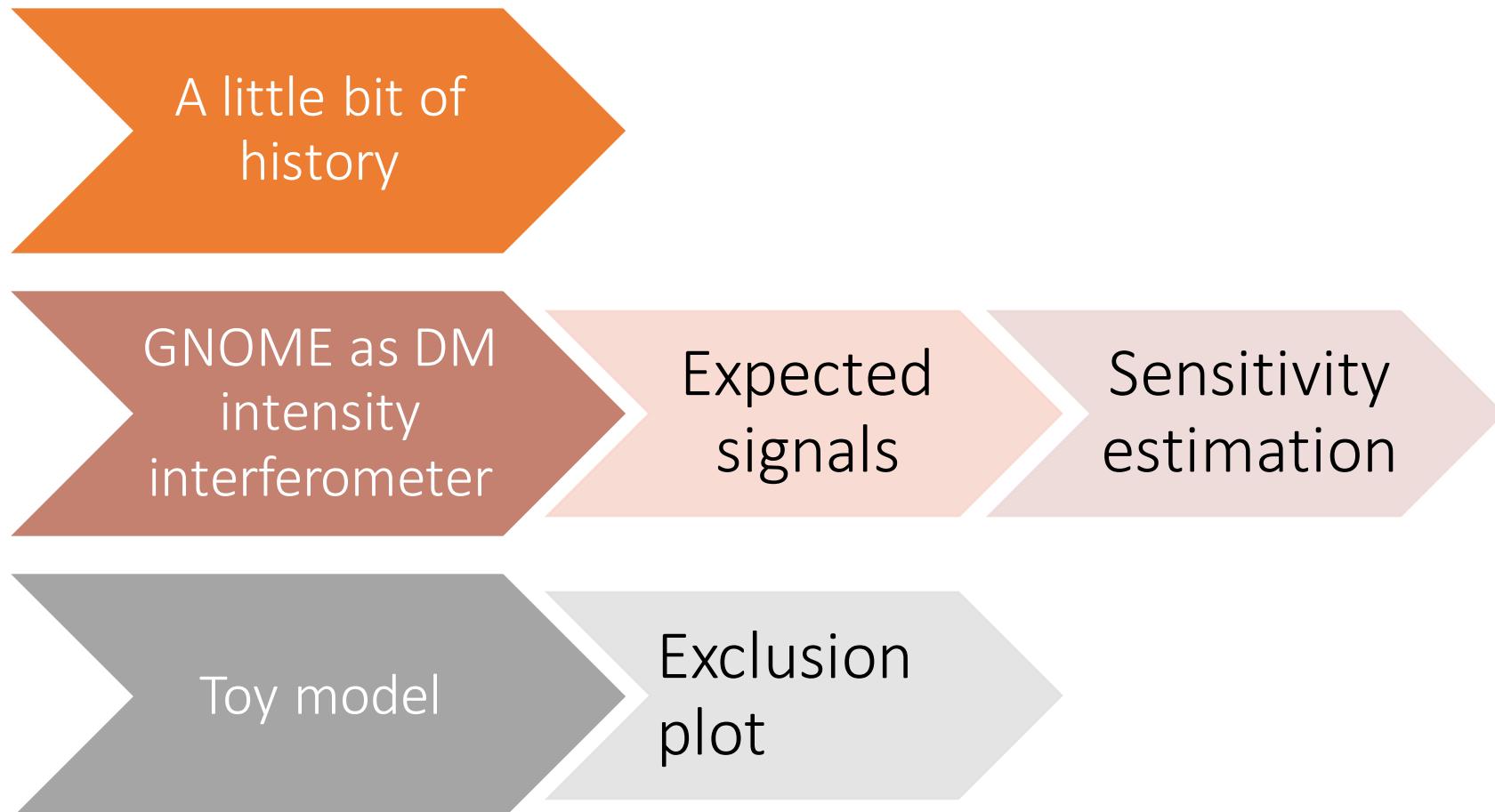
Grzegorz Łukasiewicz (GNOME Team in Kraków)

15.08.2022

MITP Workshop on Searches for Wave-Like Dark Matter with Quantum Networks

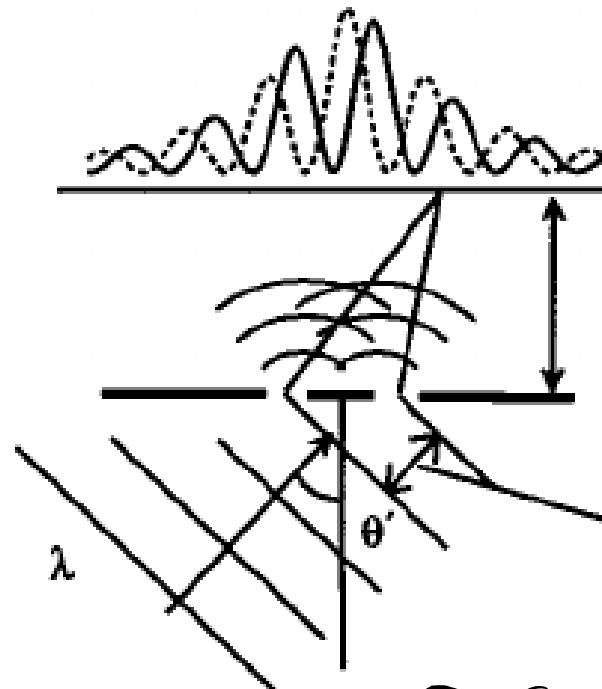


# Outline

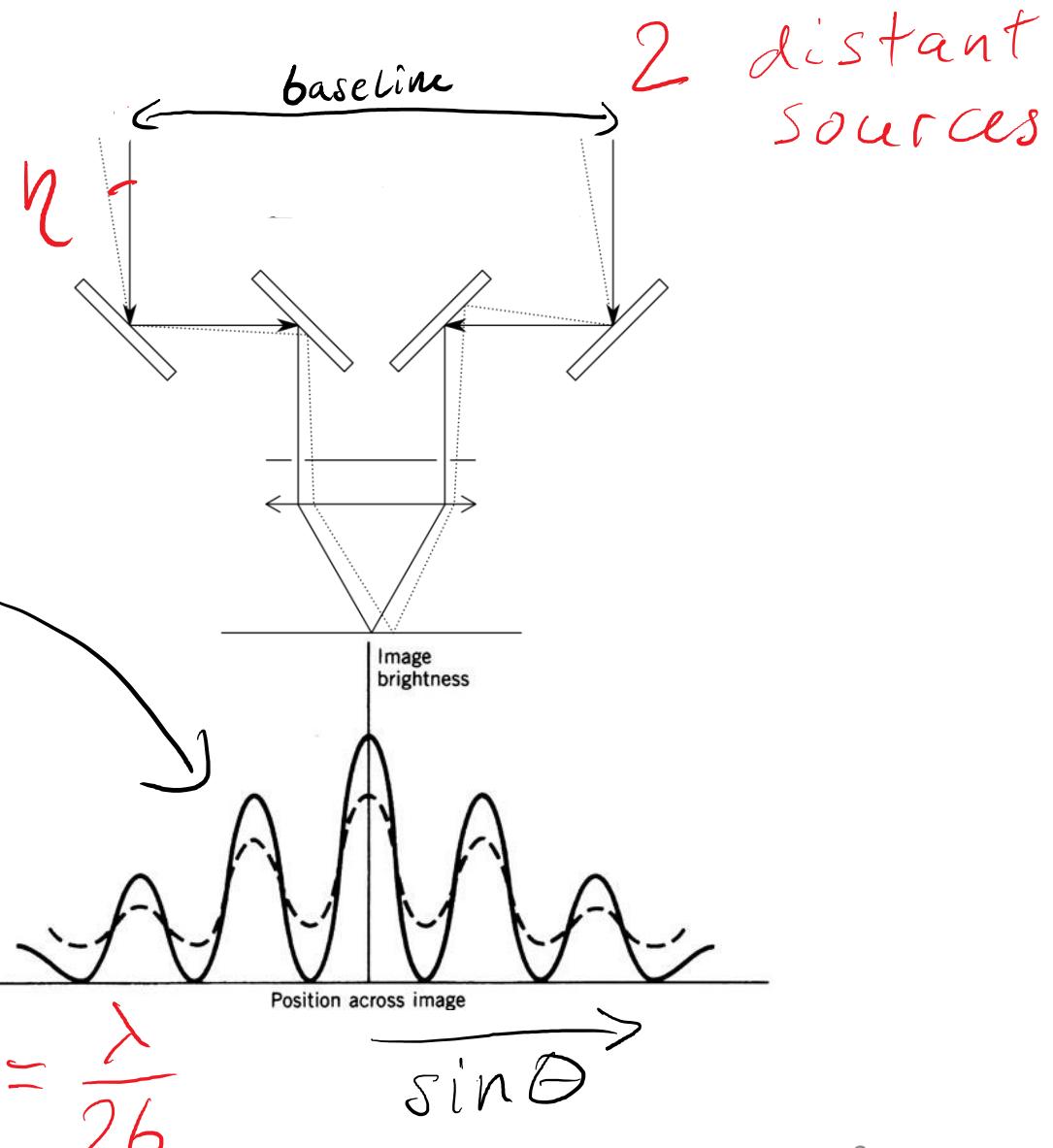


# How to Measure Angular Diameter of a Star?

incident wave  
at an angle



$$\text{resolution; } \Delta n = \frac{\lambda}{26}$$

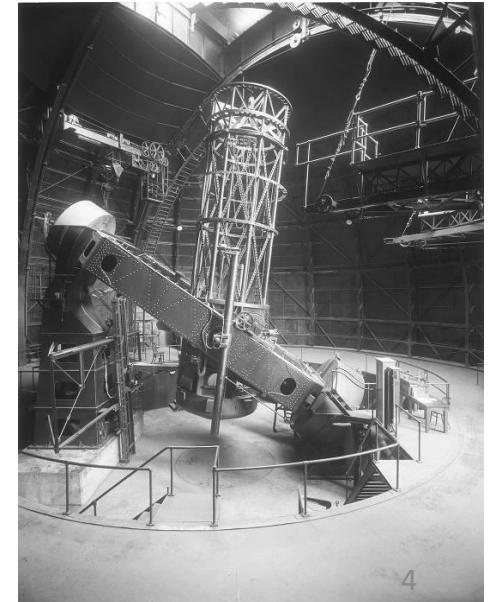
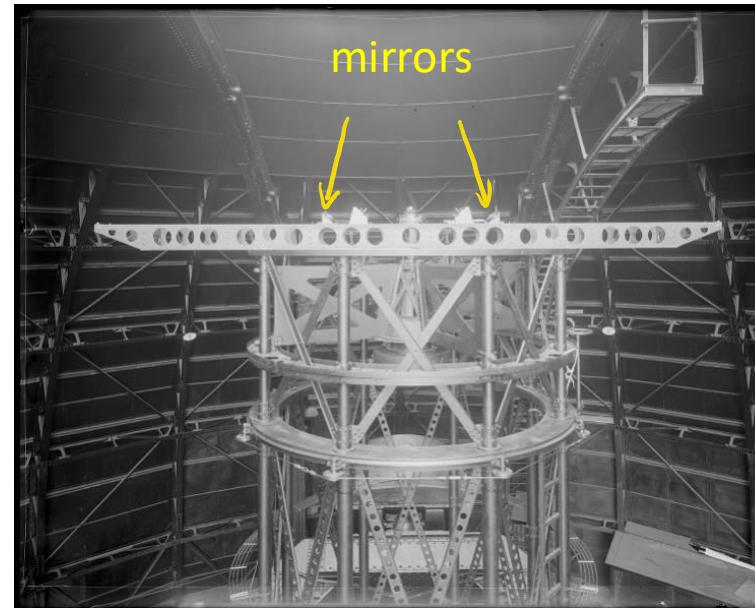
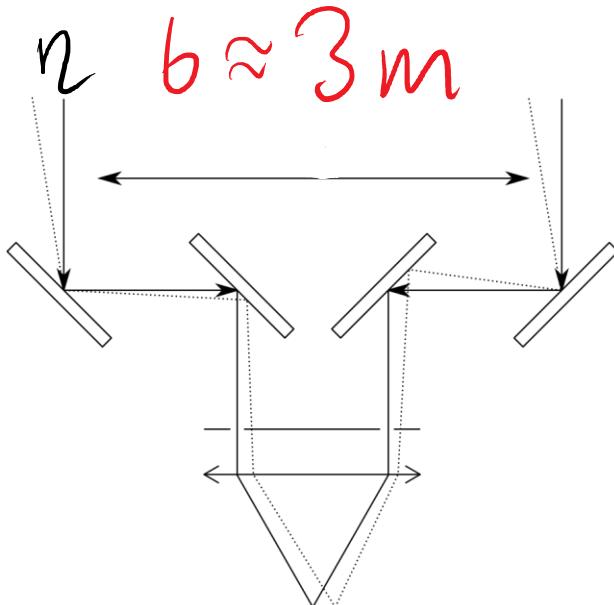


# Michelson Stellar Interferometer = Amplitude Interferometer

Betelgeuse angular diameter

measured in 1920  $2,3 \cdot 10^{-7}$  rad

$$\lambda = 575 \text{ nm}$$



# Amplitude Interferometer Limitations

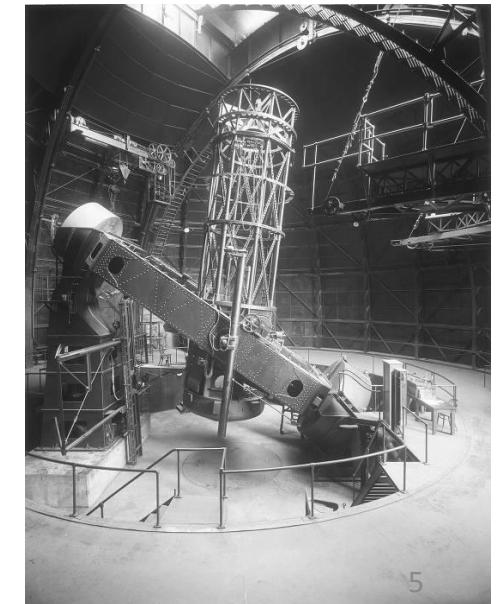
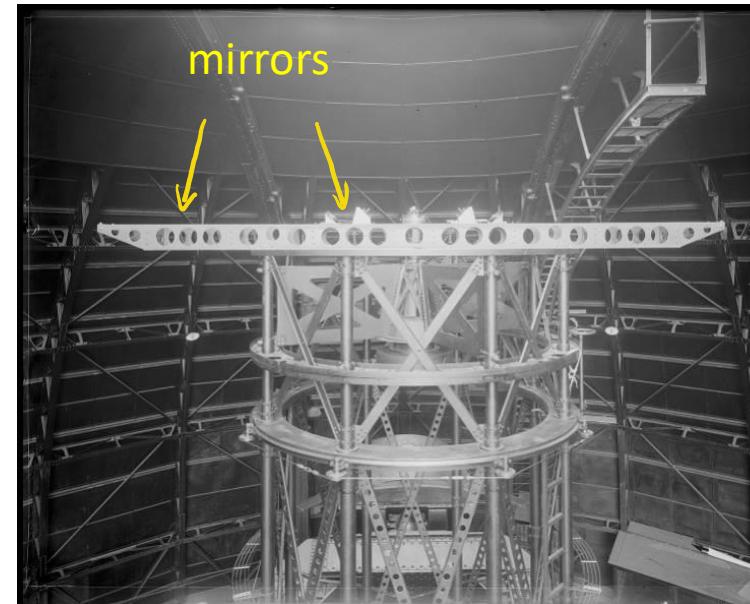
Phase stability in atmosphere and telescope

RF astronomy in 1950's

Long baselines needed

RF signals can travel in cables

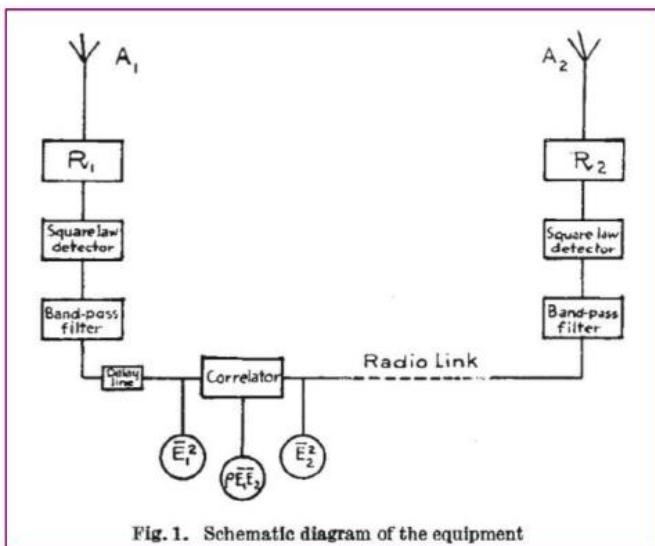
$$\Delta n = \frac{\lambda}{26}$$



# Intensity (Hanbury-Brown-Twiss) Interferometer

Apparent Angular Sizes of Discrete Radio Sources:  
Observations at Jodrell Bank, Manchester R.  
Hanbury Brown, et al, Nature (1952)

R. Hanbury Brown & R.Q. Twiss LXXIV. A new type of interferometer for use in radio astronomy, The London, Edinburgh, and Dublin Philosophical Magazine and Journal of Science, 45:366, 663-682 (1954)



The theory of the instrument is involved, and it will be given in detail elsewhere\*. It can be shown that the value of the cross-correlation coefficient ( $\rho$ ) is given by an expression similar to that for the visibility of the fringes in a Michelson stellar interferometer :

## LXXIV. A New Type of Interferometer for Use in Radio Astronomy

By R. HANBURY BROWN  
Jodrell Bank Experimental Station, Cheshire  
and

R. Q. TWISS  
Services Electronics Research Laboratory, Baldock, Herts.\*

[Received March 20, 1954]

### SUMMARY

A new type of interferometer for measuring the diameter of discrete radio sources is described and its mathematical theory is given. The principle of the instrument is based upon the correlation between the rectified outputs of two independent receivers at each end of a baseline, and it is shown that the cross-correlation coefficient between these outputs is proportional to the square of the amplitude of the Fourier transform of the intensity distribution across the source. The analysis shows that it should be possible to operate the new instrument with extremely long baselines and that it should be almost unaffected by ionospheric irregularities.

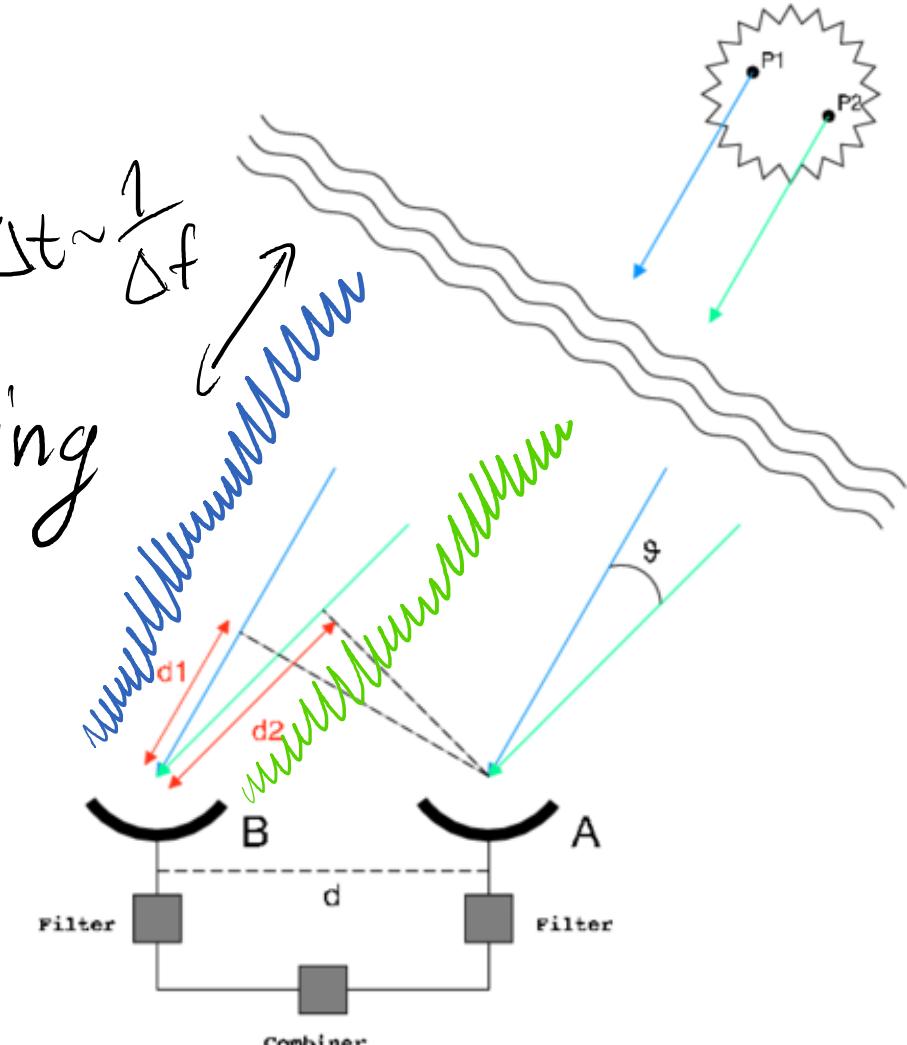
# Intensity Interferometry

Intensity detectors

Phase is lost

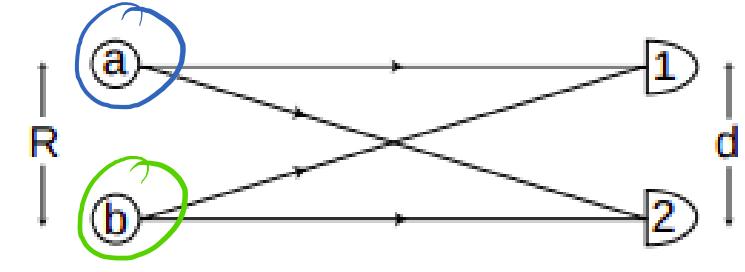
Narrow bandwidth  $\rightarrow$  beating

$$C = \frac{\langle I_A I_B \rangle}{\langle I_A \rangle \langle I_B \rangle}$$



(correlator)

How does it work?



$$A_1 = \frac{1}{L} (\underline{\alpha e^{ikr_{1a} + i\phi_a}} + \underline{\beta e^{ikr_{1b} + i\phi_b}})$$

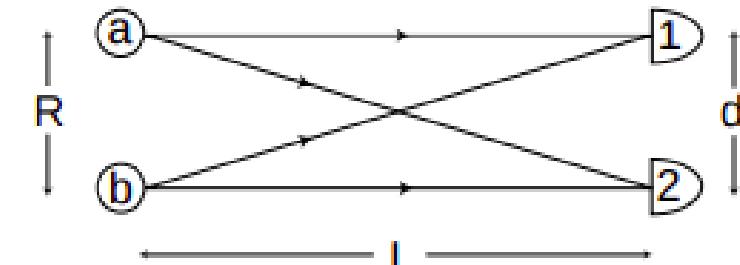
$$I_1 = \frac{1}{L^2} \left( |\alpha|^2 + |\beta|^2 + \cancel{\alpha^* \beta e^{ik(r_{1b} - r_{1a}) + i(\phi_b - \phi_a)}} + \cancel{\alpha \beta^* e^{-ik(r_{1b} - r_{1a}) + i(\phi_b - \phi_a)}} \right)$$

averaged over random phases

$$\langle I_1 \rangle = \langle I_2 \rangle = \frac{1}{L^2} (\langle |\alpha|^2 \rangle + \langle |\beta|^2 \rangle)$$

How does it work?  $I_1 = \frac{1}{L^2} (|\alpha|^2 + |\beta|^2 + \alpha^* \beta e^{ik(r_{1b} - r_{1a}) + i(\phi_b - \phi_a)} + \alpha \beta^* e^{-ik(r_{1b} - r_{1a}) + i(\phi_b - \phi_a)})$

$\langle I_1 \rangle = \langle I_2 \rangle = \frac{1}{L^2} (\langle |\alpha|^2 \rangle + \langle |\beta|^2 \rangle)$

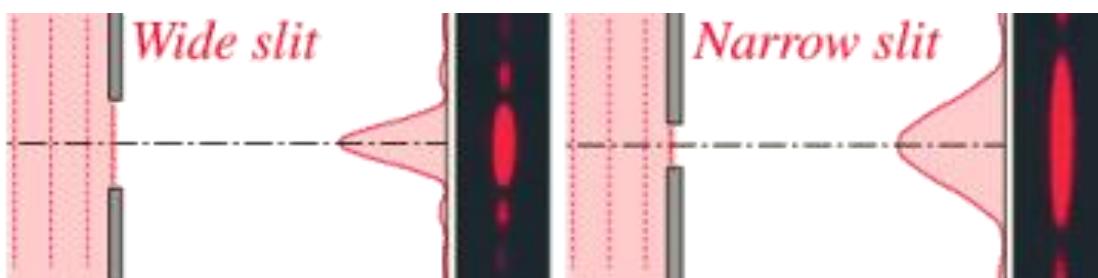


Intensity correlation depends on baseline

$L \gg R$

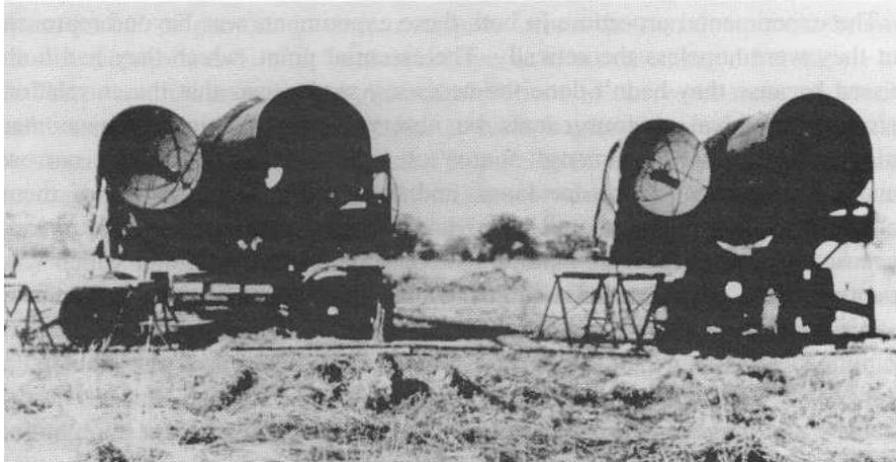
$\langle I_1 I_2 \rangle = \langle I_1 \rangle \langle I_2 \rangle + \frac{2}{L^4} |\alpha|^2 |\beta|^2 \cos(k(r_{1a} - r_{2a} - r_{1b} + r_{2b}))$

$\approx \cos\left(\frac{kRd}{L}\right)$

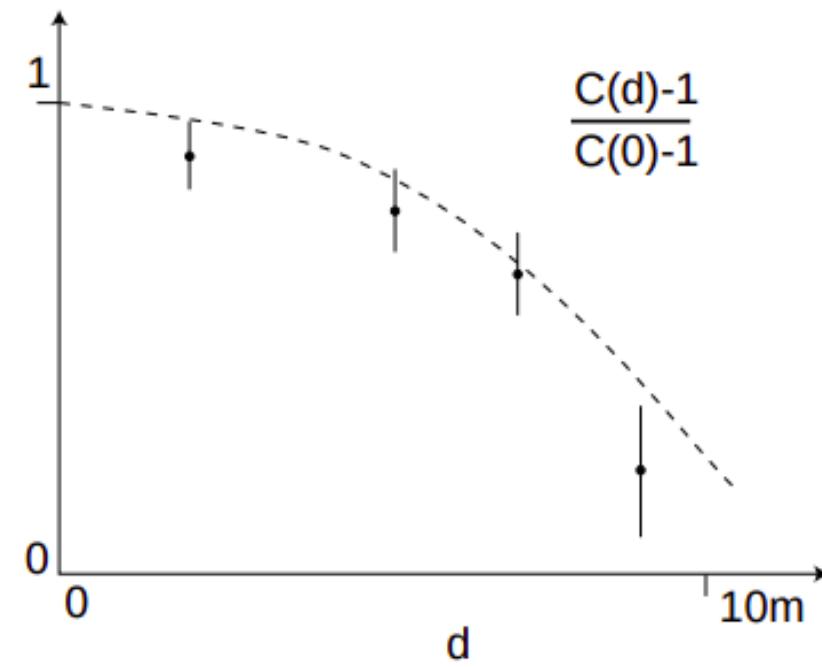


# Intensity Interferometry by HBT

1956 measurement of Sirius angular diameter  $3,7 \times 10^{-8} \text{ rad}$



Military searchlights in Jordell Bank

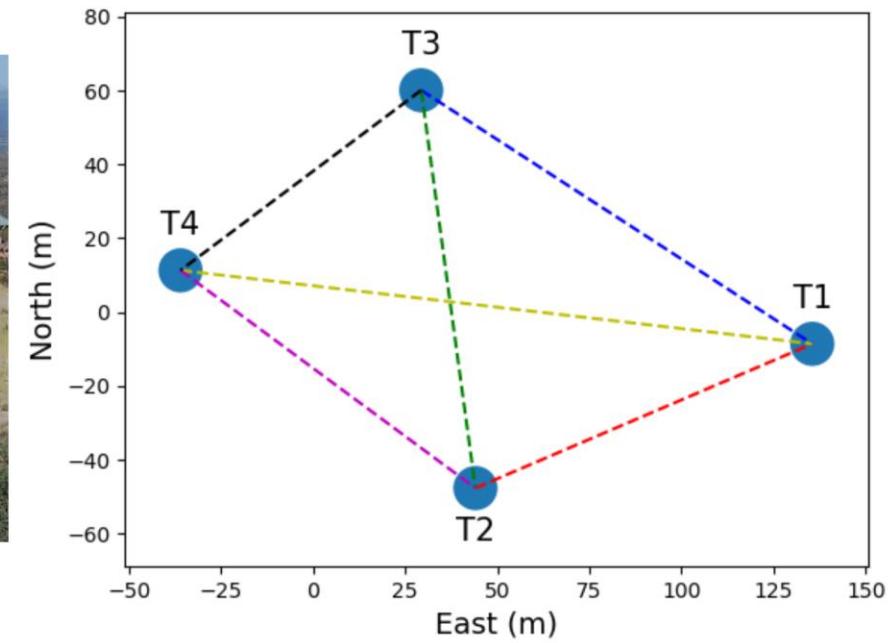


quantum optics?

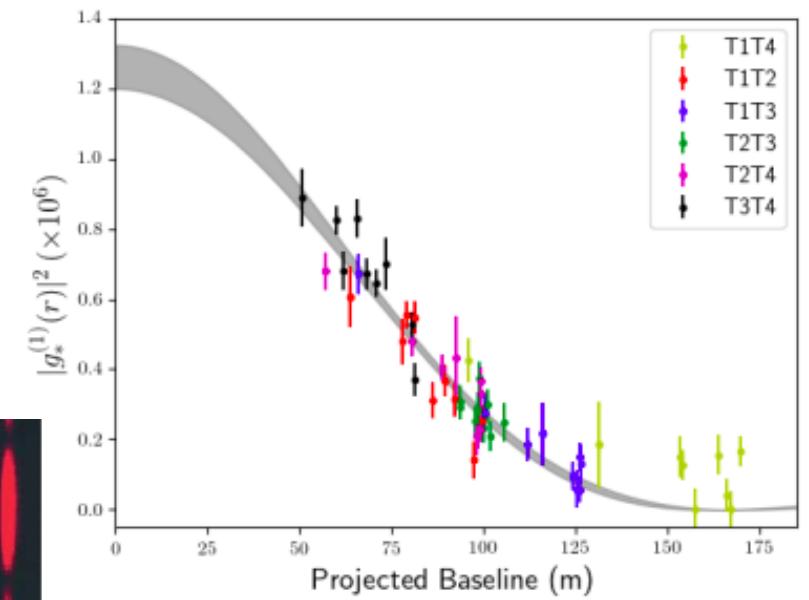
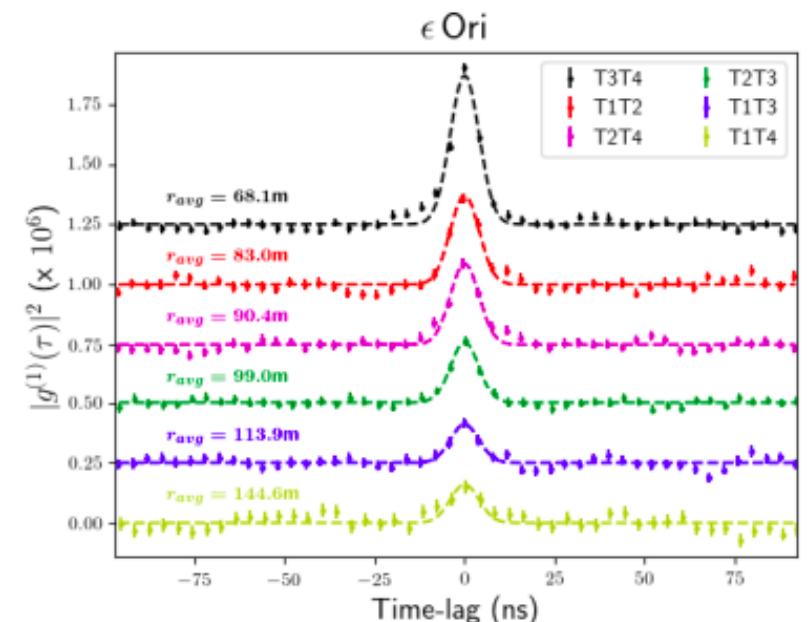
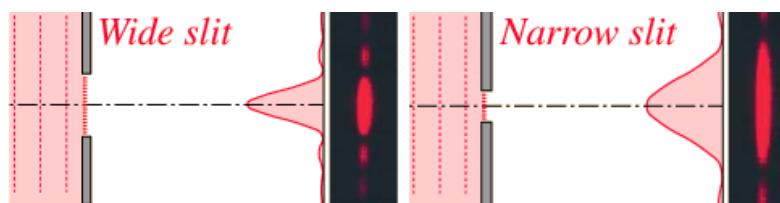
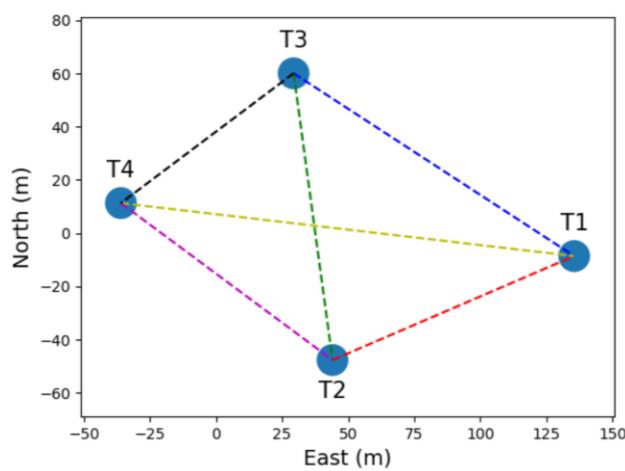


# Back to Present Times

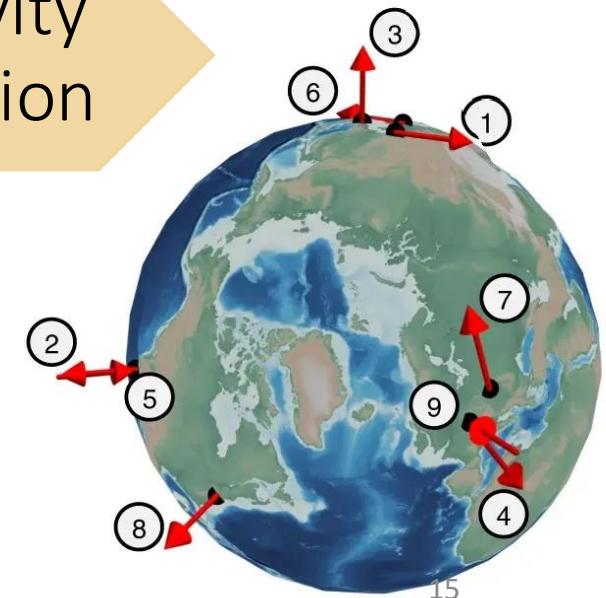
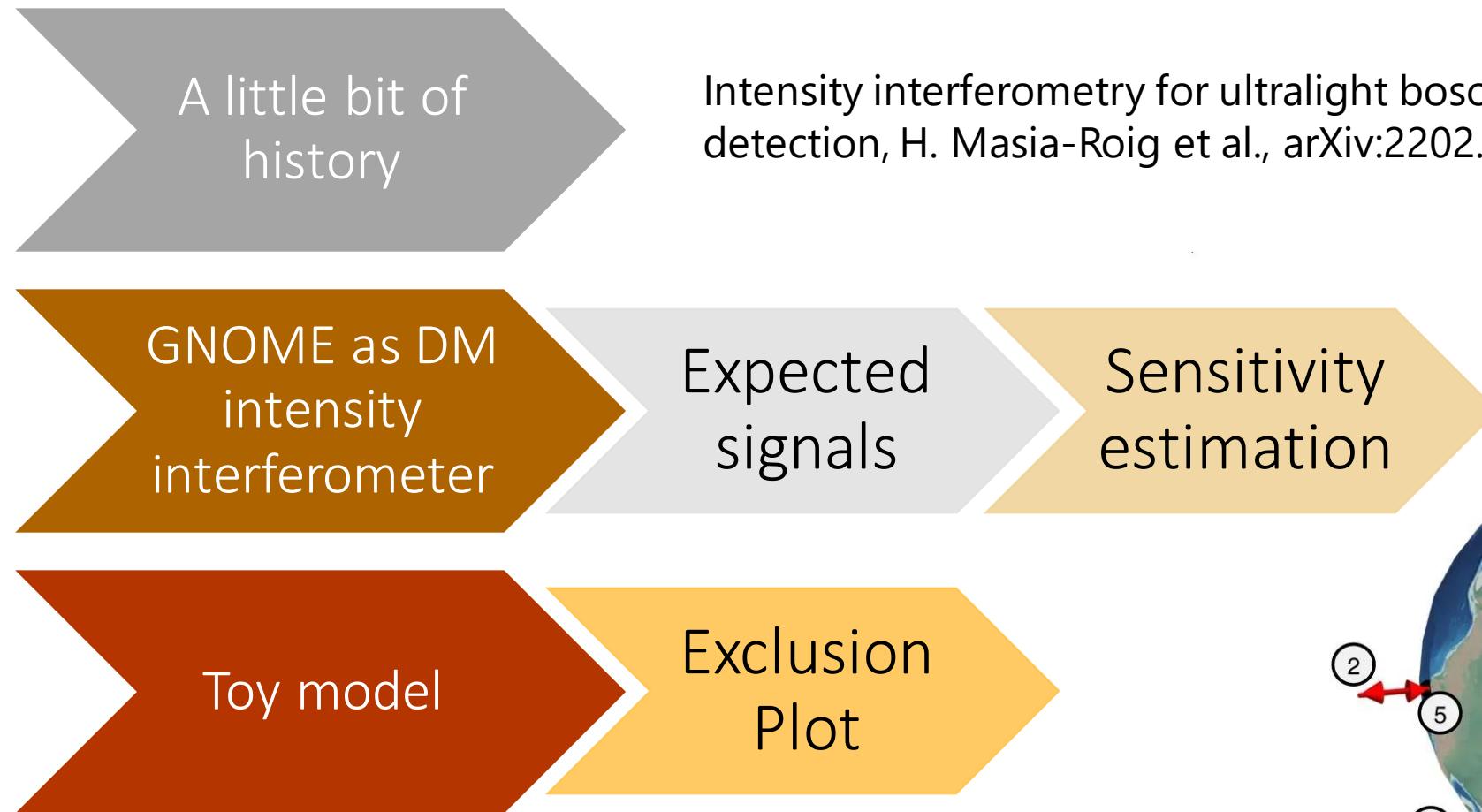
- Very Energetic Radiation Imaging Telescope Array System (VERITAS),  
Arizona, US



# VERITAS



# Dark Matter Intensity Interferometer?



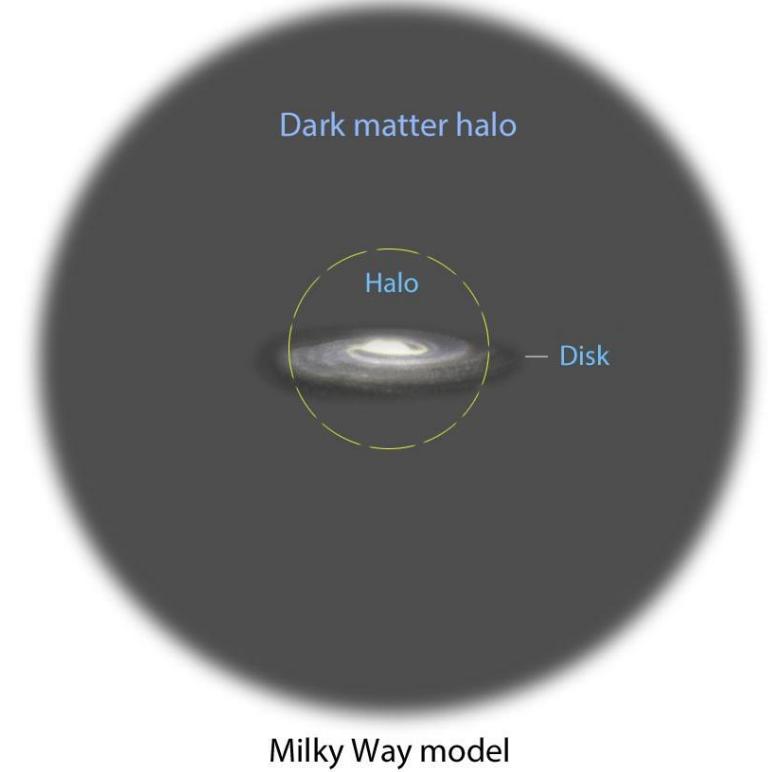
# Ultralight DM

- Single particle – plain wave
- Superposition
- Virialization of DM in the galactic halo

$$\omega_h = \omega_c \left( 1 + \frac{\sigma_h^2}{2c^2} \right)$$

$$\sigma_0 \approx 220 \frac{\text{km}}{\text{s}}$$

Narrow bandwidth



Milky Way model

# Stochastic Fluctuations of Dark Matter Field

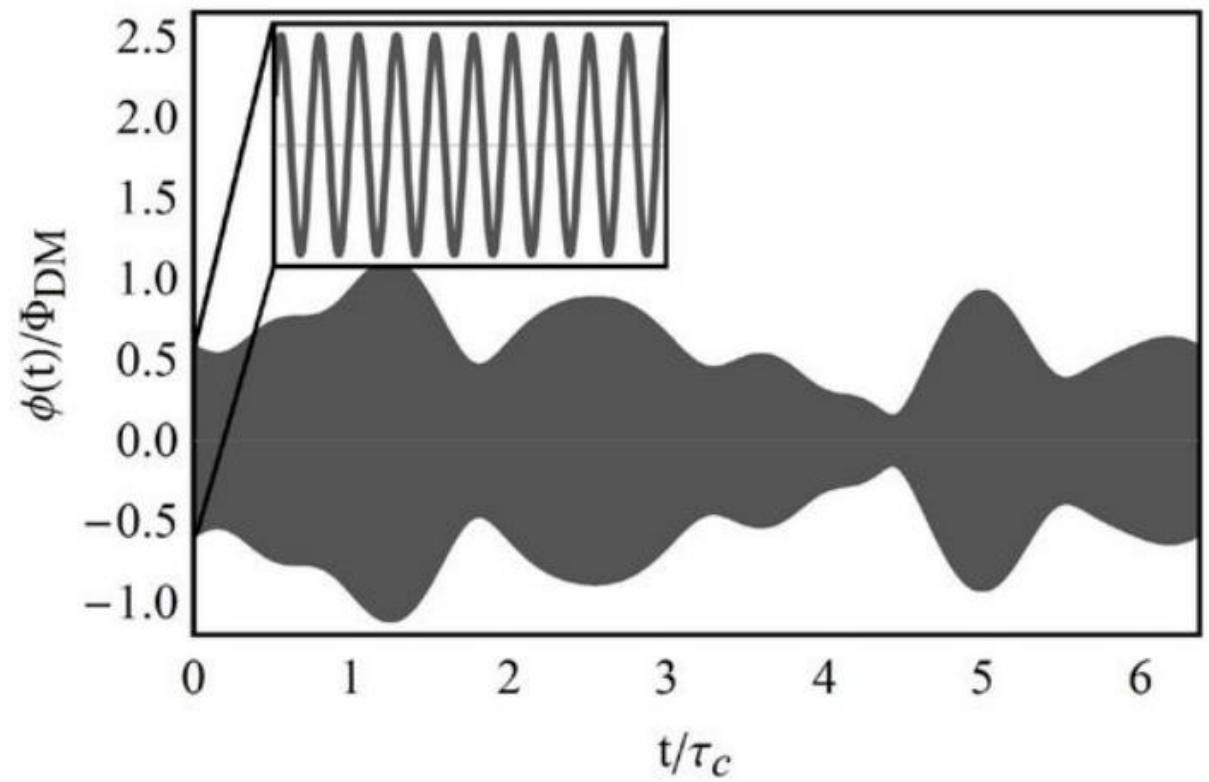
$$\cos(\omega_1 t) + \cos(\omega_2 t) = 2 \cos\left(\frac{\omega_1 - \omega_2}{2} t\right) \cos\left(\frac{\omega_1 + \omega_2}{2} t\right)$$

$\sim \Delta\omega$                              $\sim 2\omega_c$

# Stochastic Fluctuations of Dark Matter Field

$$\phi(\vec{r}, t) \sim \sum_n \cos(\omega_n t - \vec{k}_n \cdot \vec{r} + \theta_n)$$

beating  $\tau_c \sim \frac{1}{\Delta\omega}$

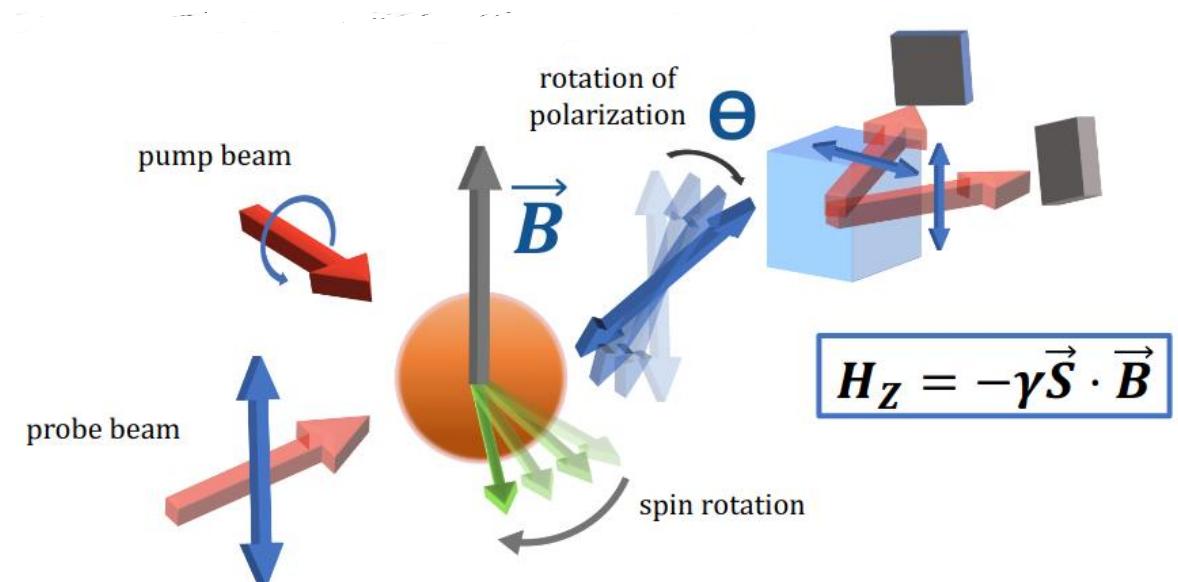


# Intensity Detector - Magnetometer

Quadratic coupling

$$H_\ell \sim \frac{2\hbar c^2}{f_g^2} \vec{S} \cdot \vec{\nabla} \ell^2(\vec{r}, t)$$

Why quadratic coupling?

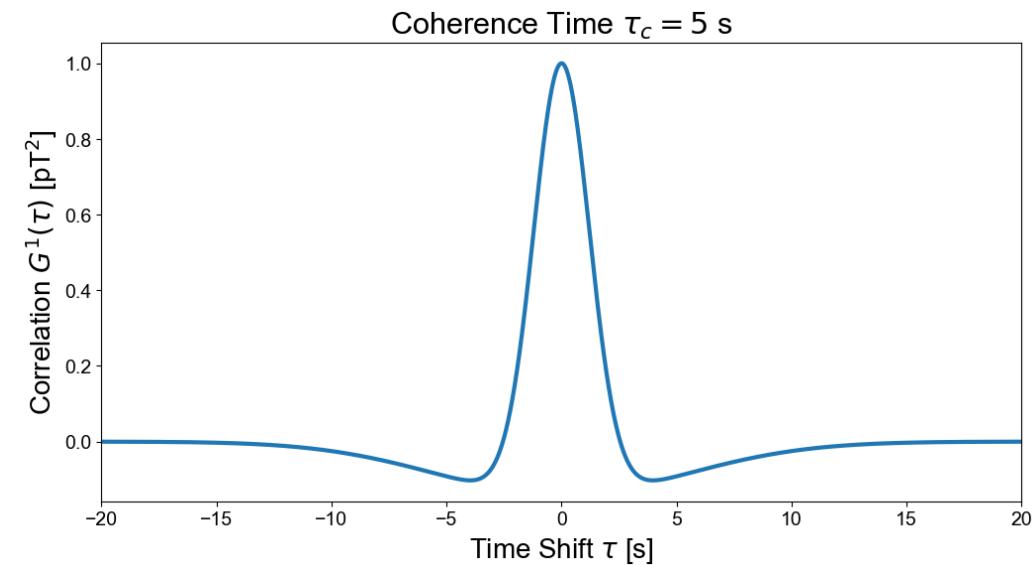
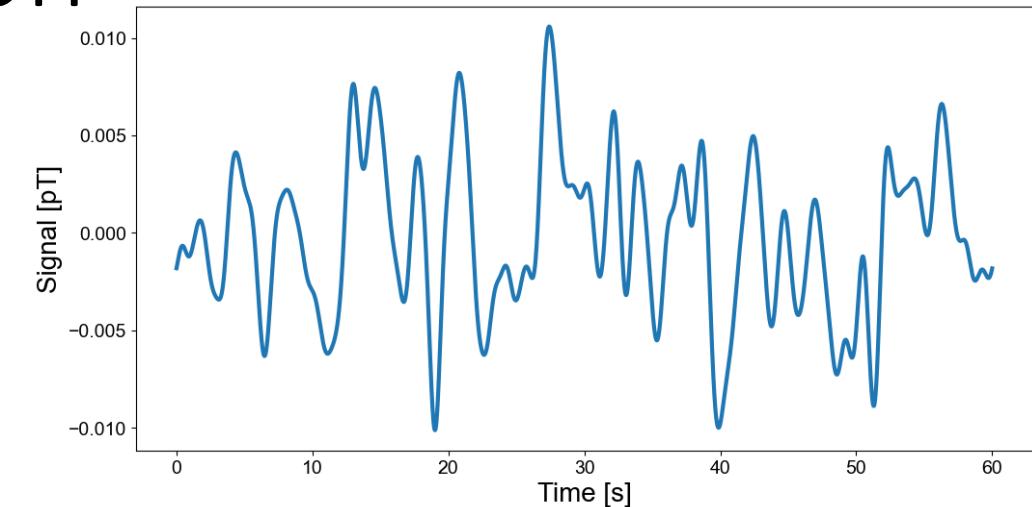
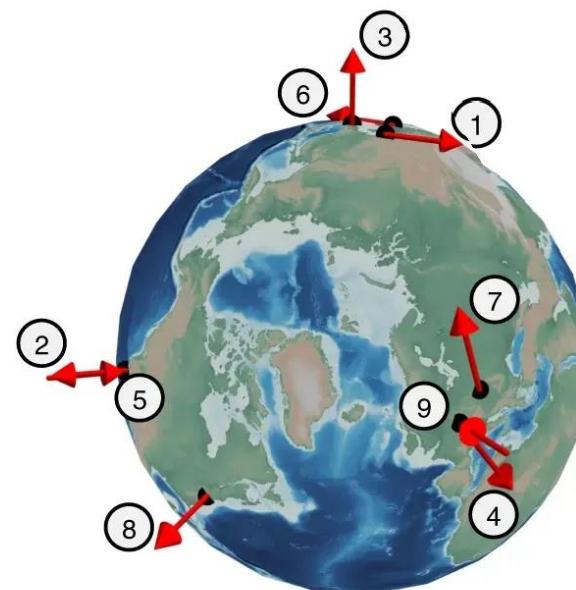
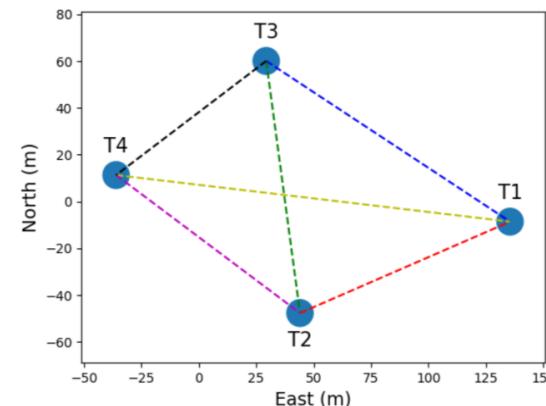
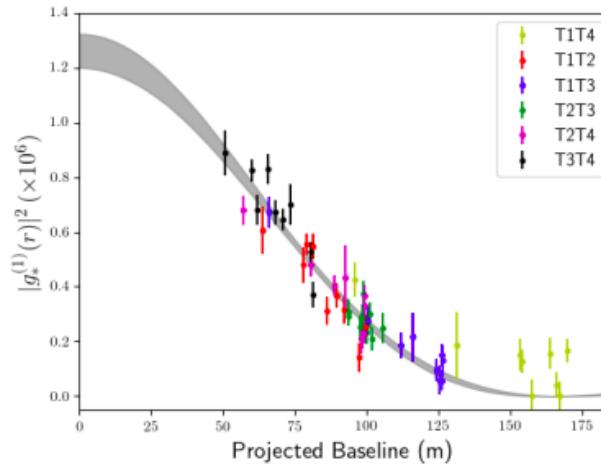
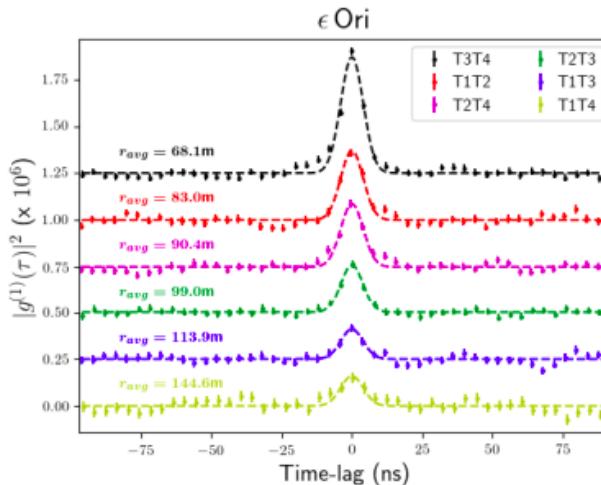


Sensitivity < pT  
100 Hz bandwidth

# GNOME (simulation)

## DM Signature = Correlation

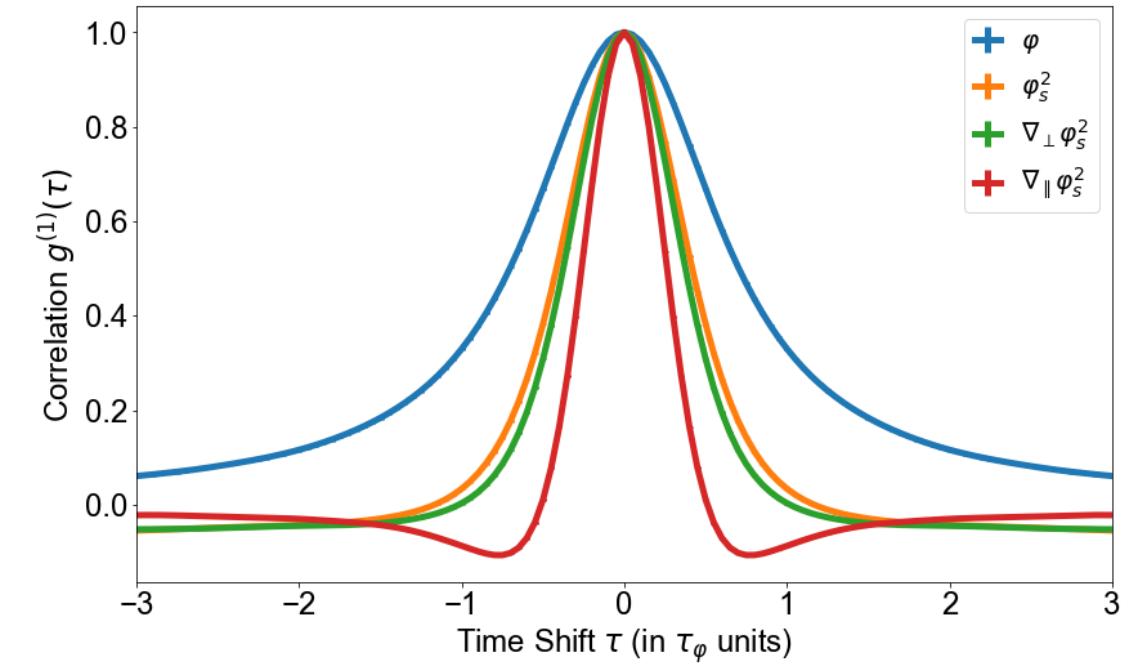
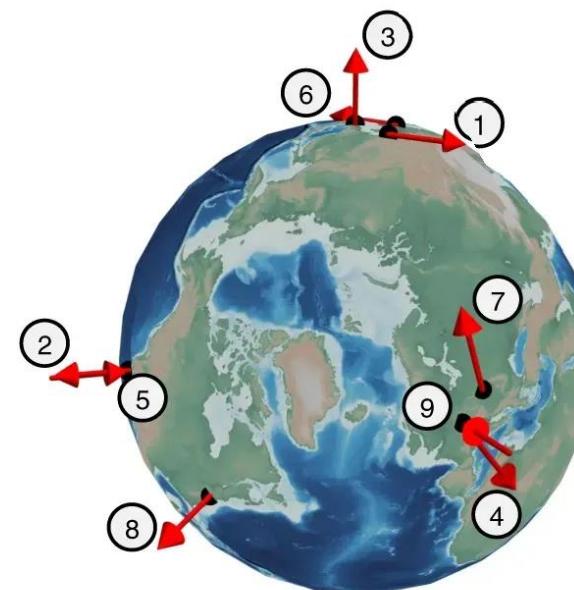
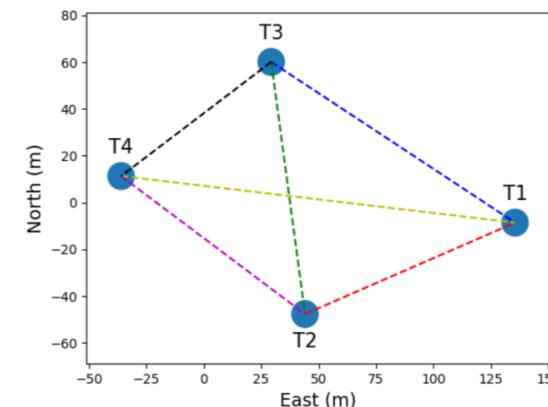
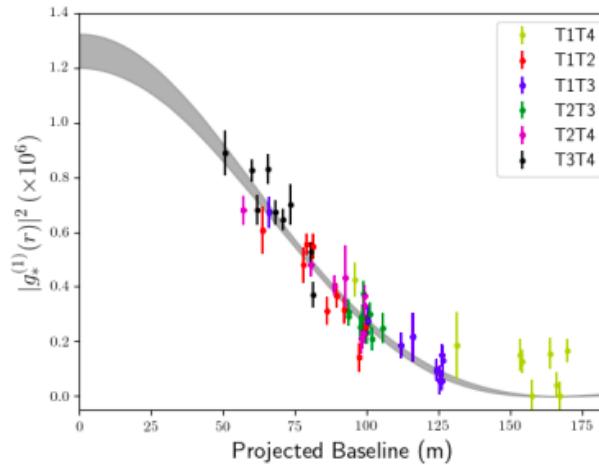
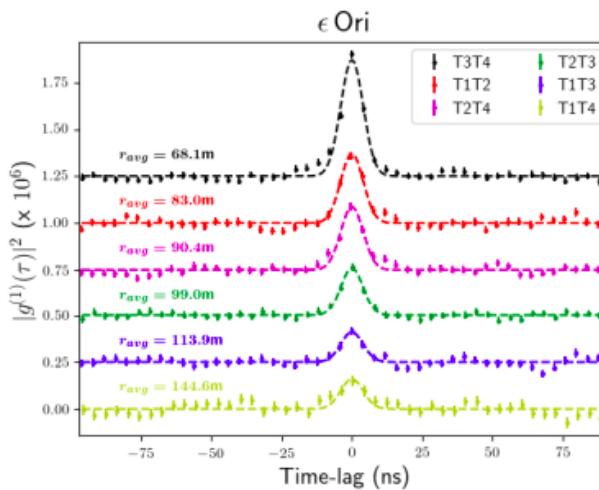
VERITAS



# GNOME (simulation)

## DM Signature = Correlation

VERITAS

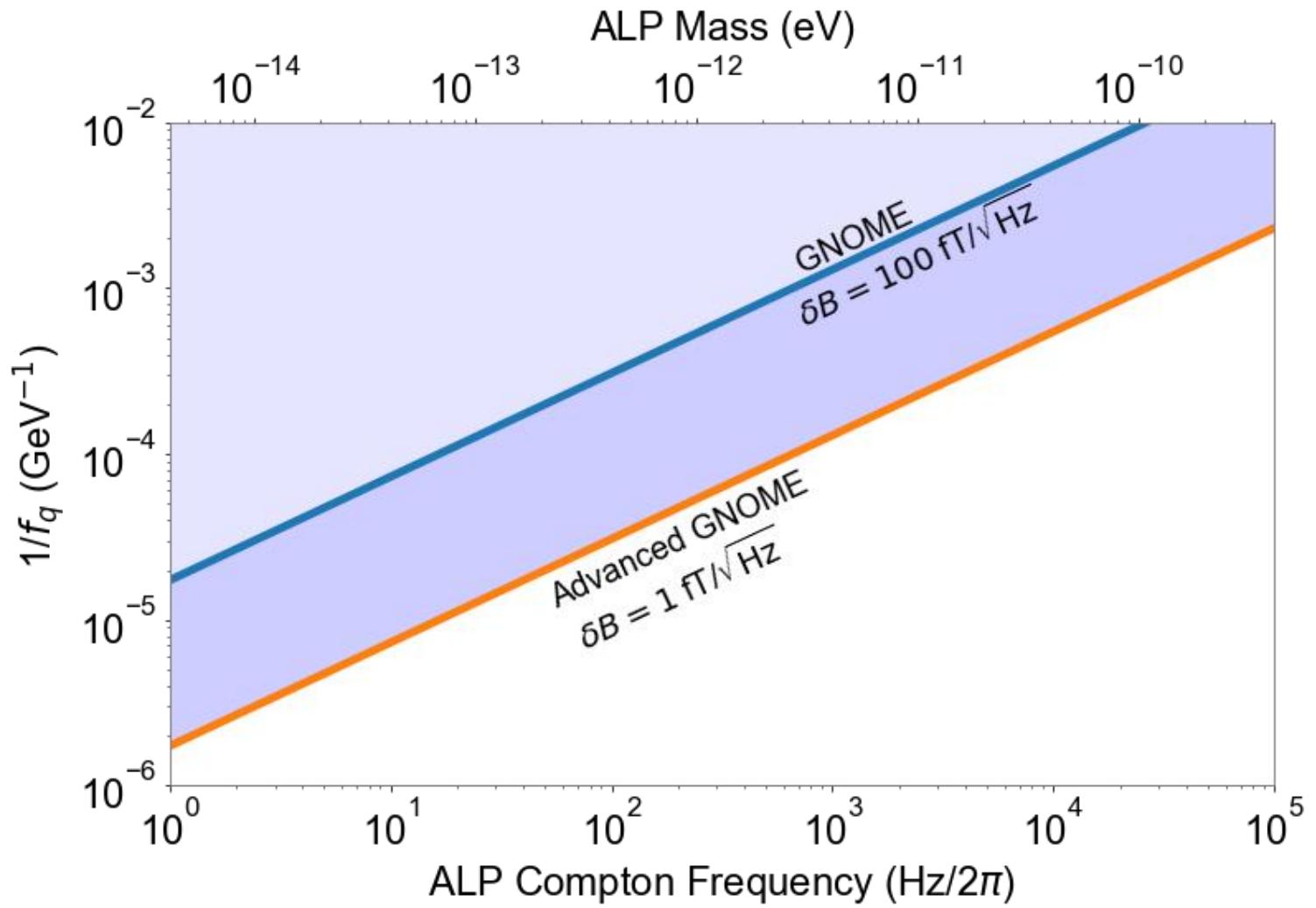


# Sensitivity Estimation

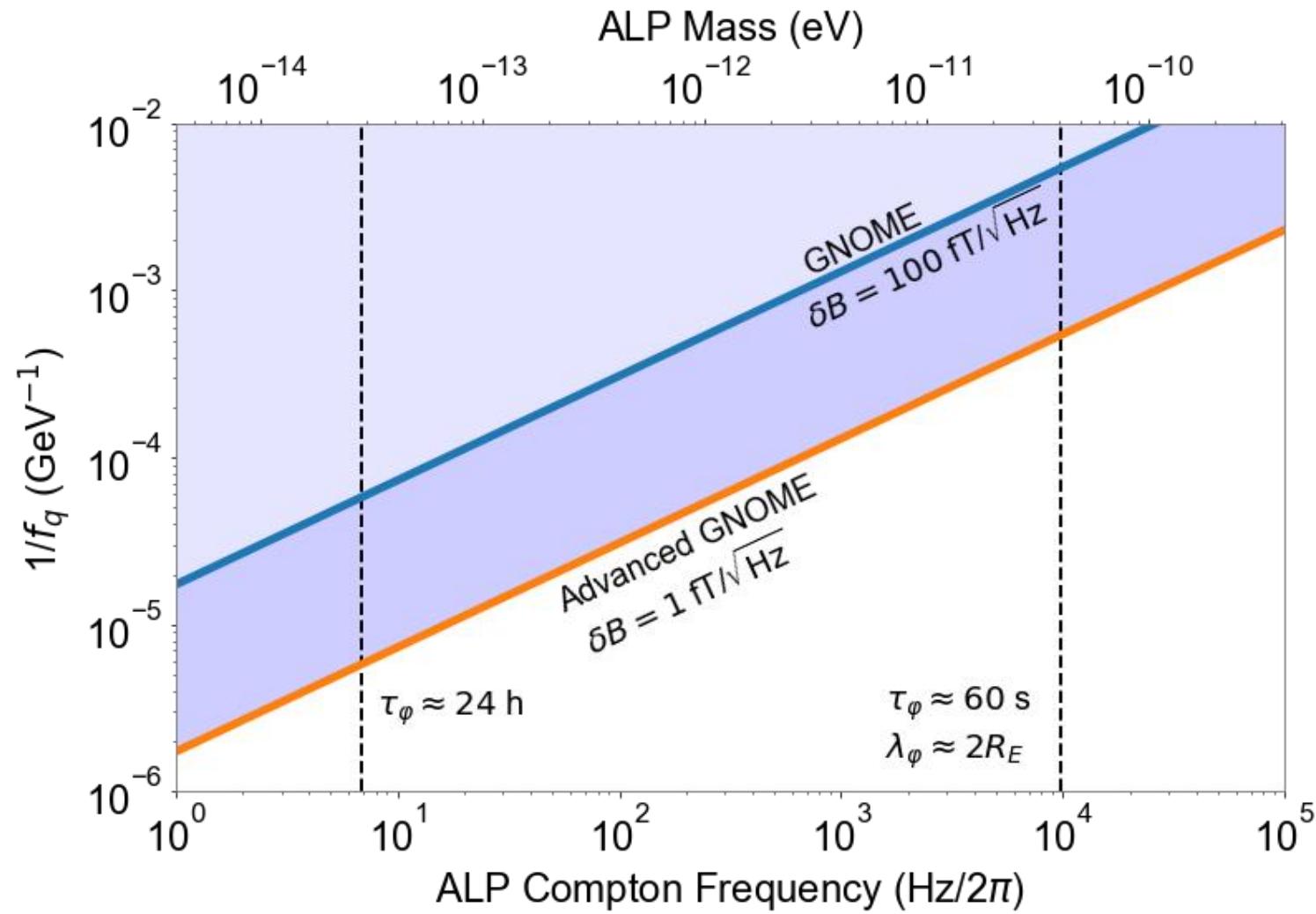
10 sensors

100 days

$\delta B(f) = \text{const.}$

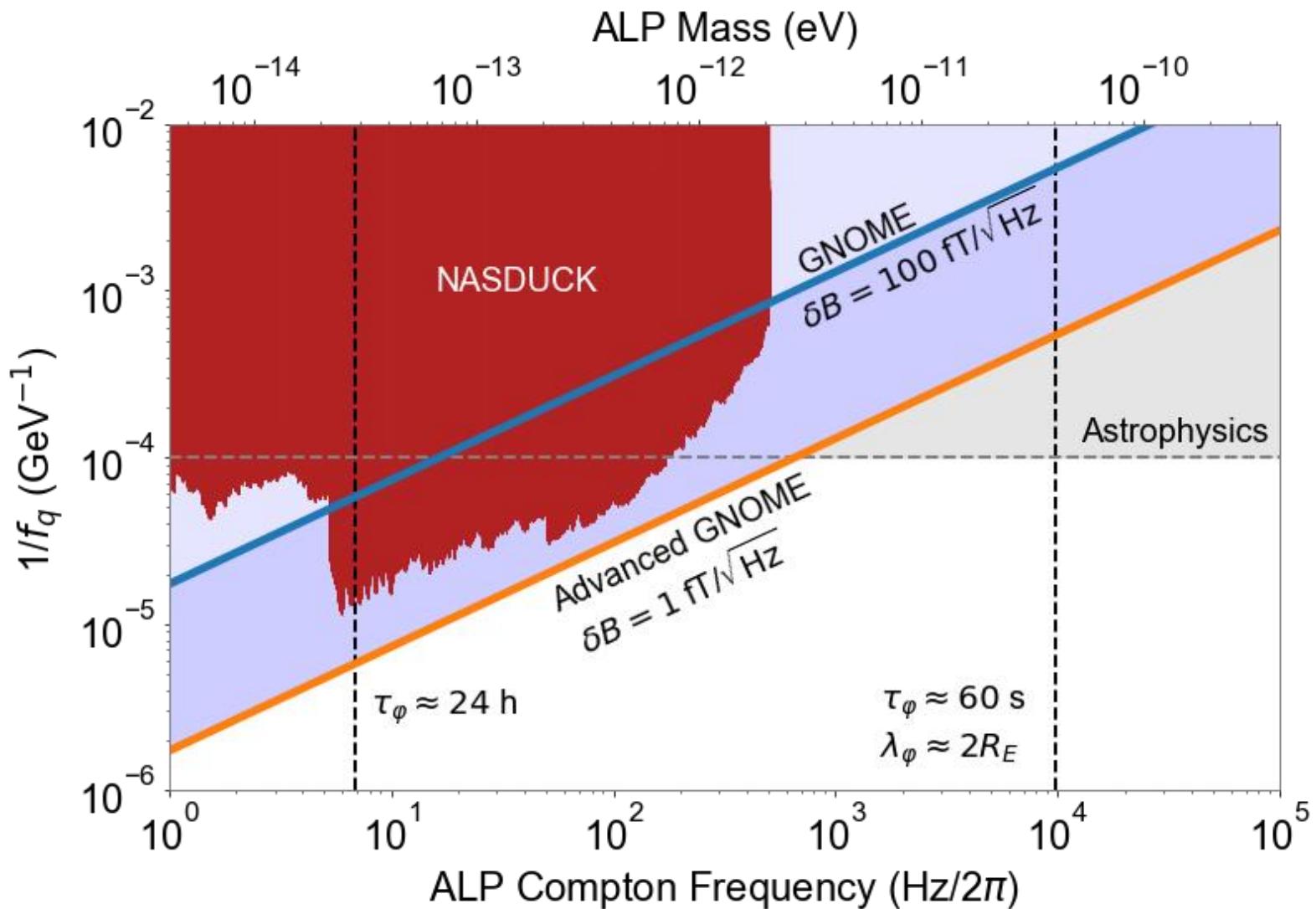


# ALP Mass Range Limitations

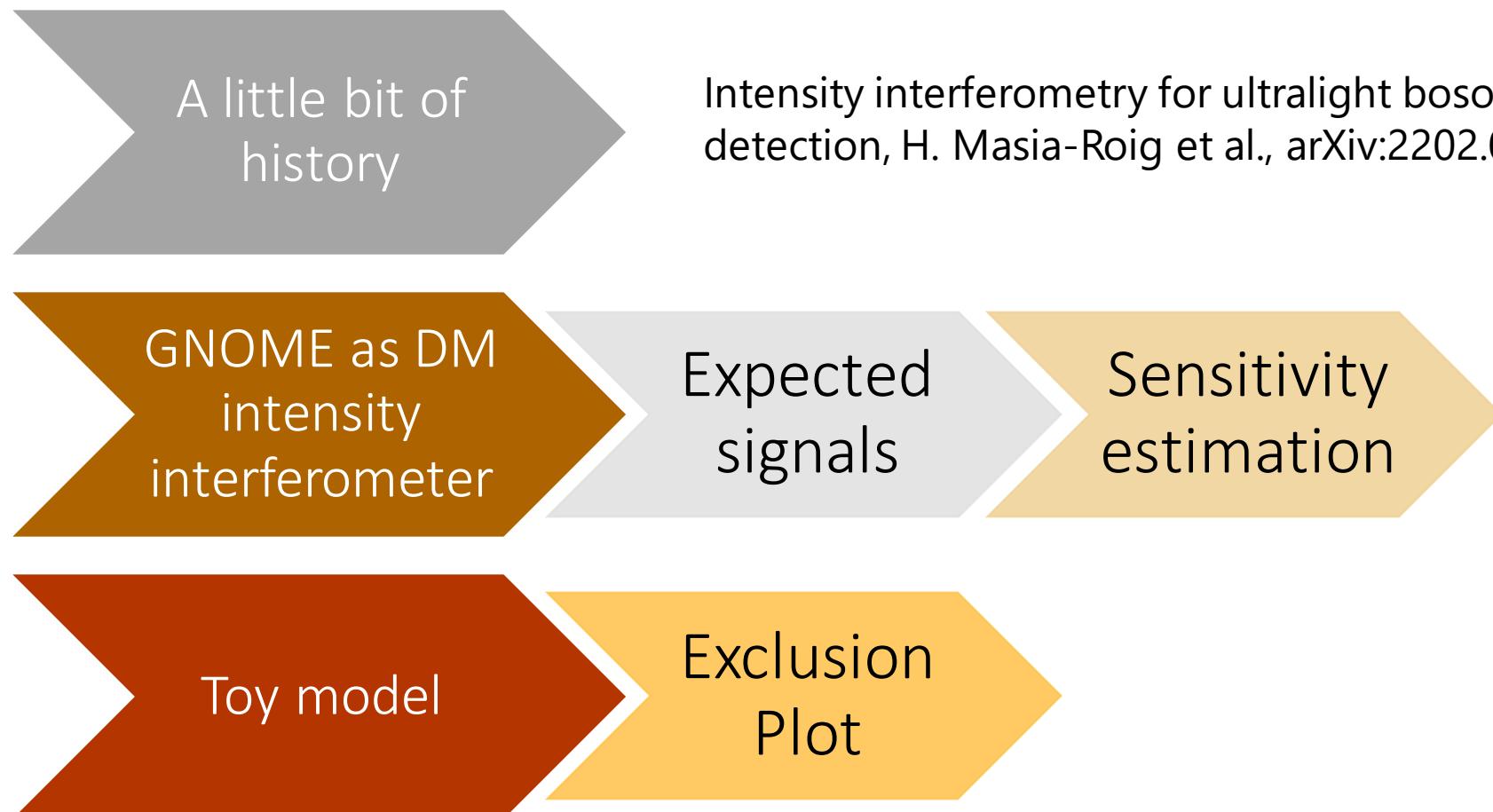


# Current Bounds

Wide frequency bandwidth thanks to  
intensity interferometry!

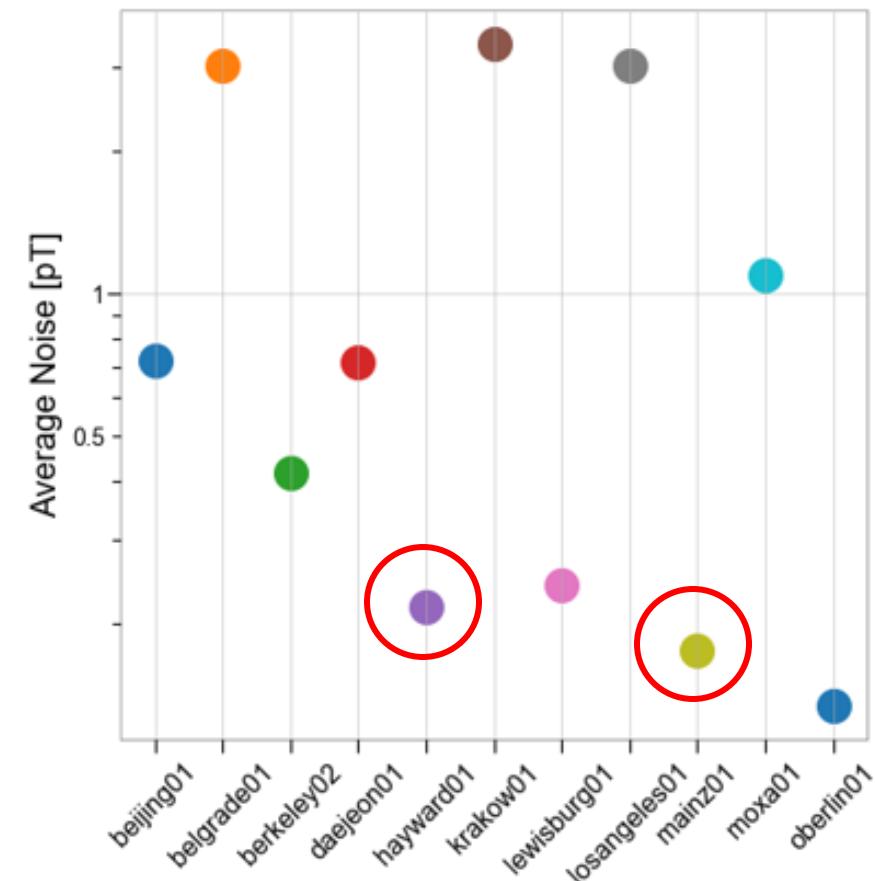


# Dark Matter Intensity Interferometer?

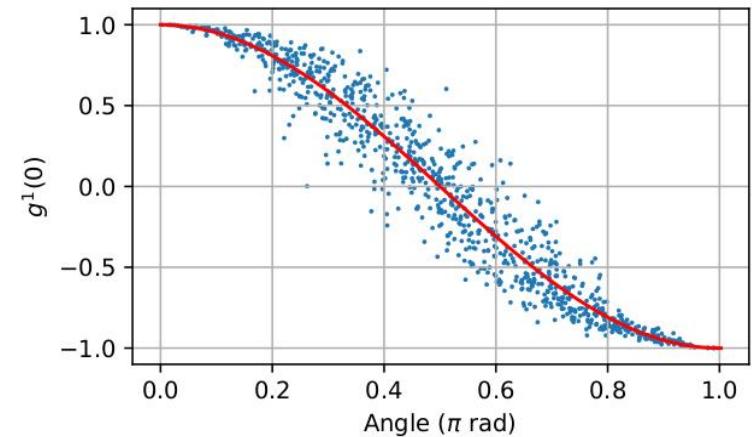
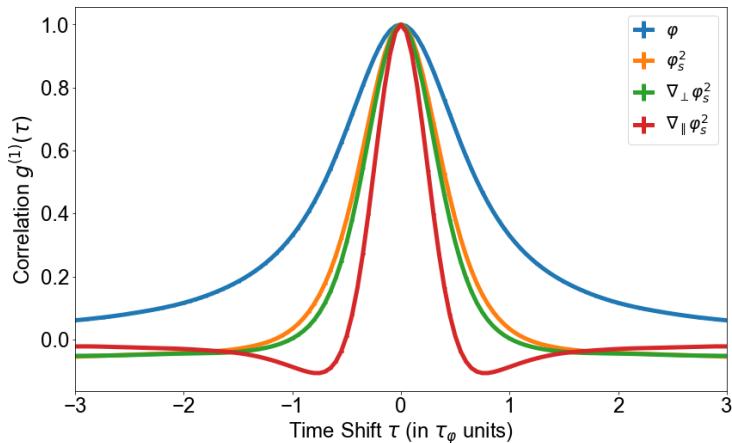


# Toy Model

$$G^1(0)$$



$$|G^1(0)| = \max G^1(\tau \gg \tau_c) \approx 0$$



dot products of sensitive axes

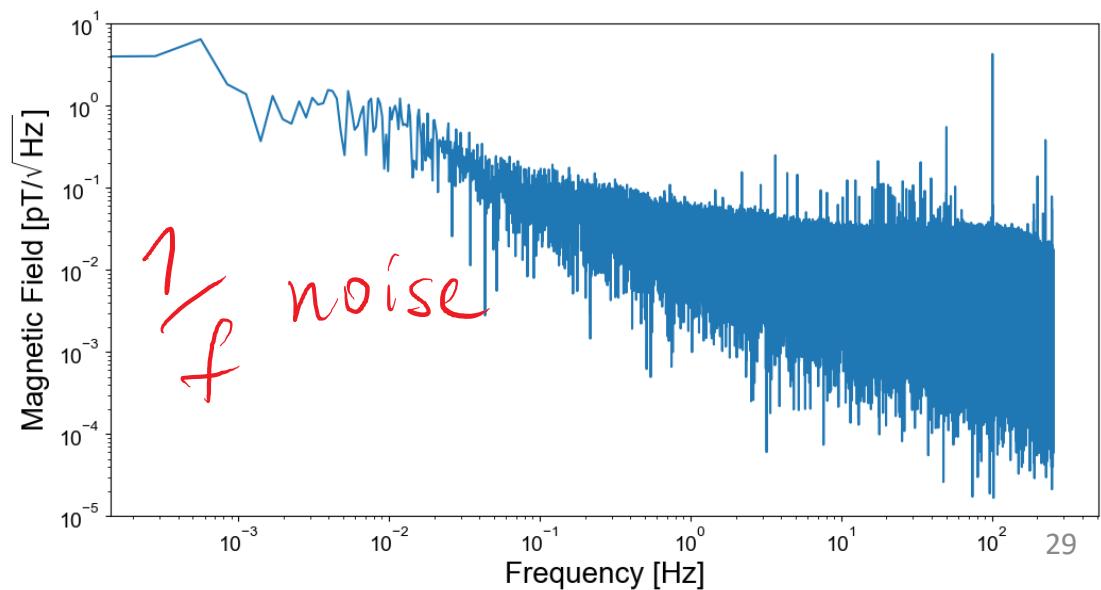
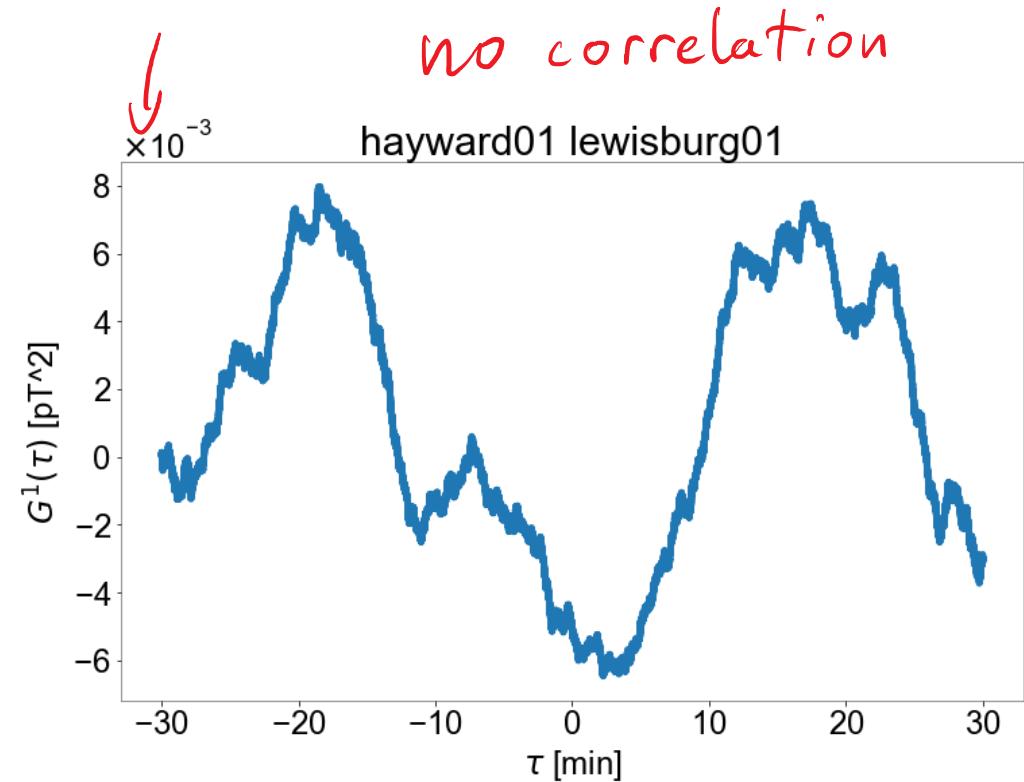
	beijing01	belgrade01	berkeley02	daejeon01	hayward01	krakow01	lewisburg01	losangeles01	mainz01	moxa01	oberlin01
beijing01											
belgrade01											
berkeley02											
daejeon01	0.986011										
hayward01					0.999989						
krakow01				0.950574							
lewisburg01					0.821179			0.821956			
losangeles01									-0.991463		
mainz01						-0.982689				0.754957	
moxa01										0.893642	
oberlin01							0.798336				28

# Correlation

1 hour long segments

- subtract linear background
- filter 60 Hz

averaged over  
1 month



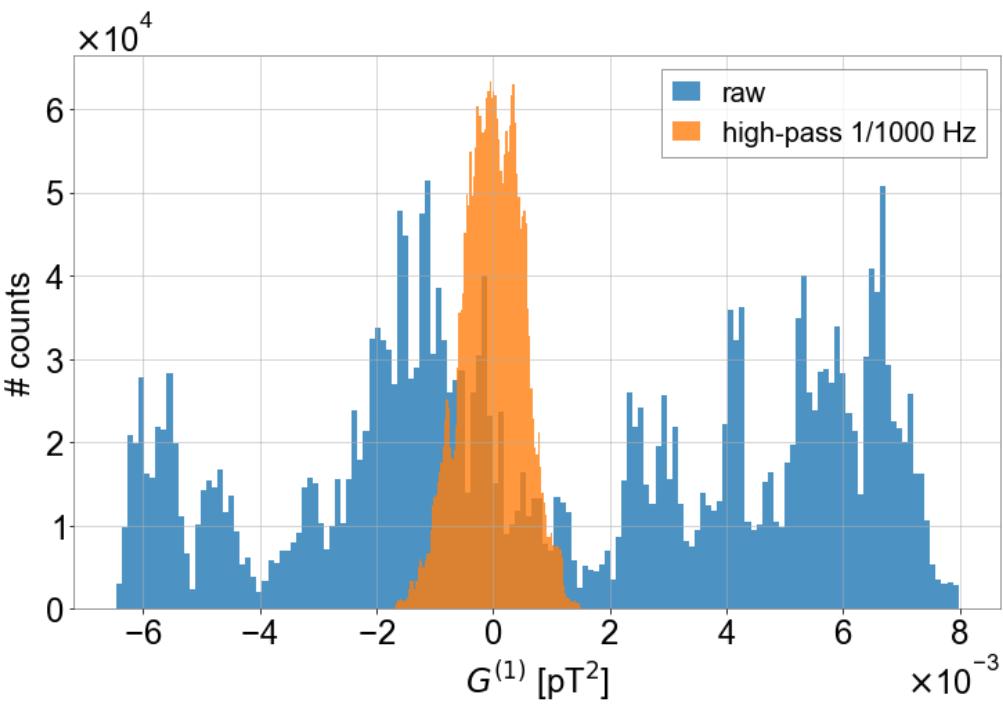
no correlation  $\rightarrow$  how to put constraints?

$$f_q^2 \lesssim \frac{\hbar^3 S_{DM} U_0}{g_F M_B M_Q} \frac{1}{B_q(\min)}$$

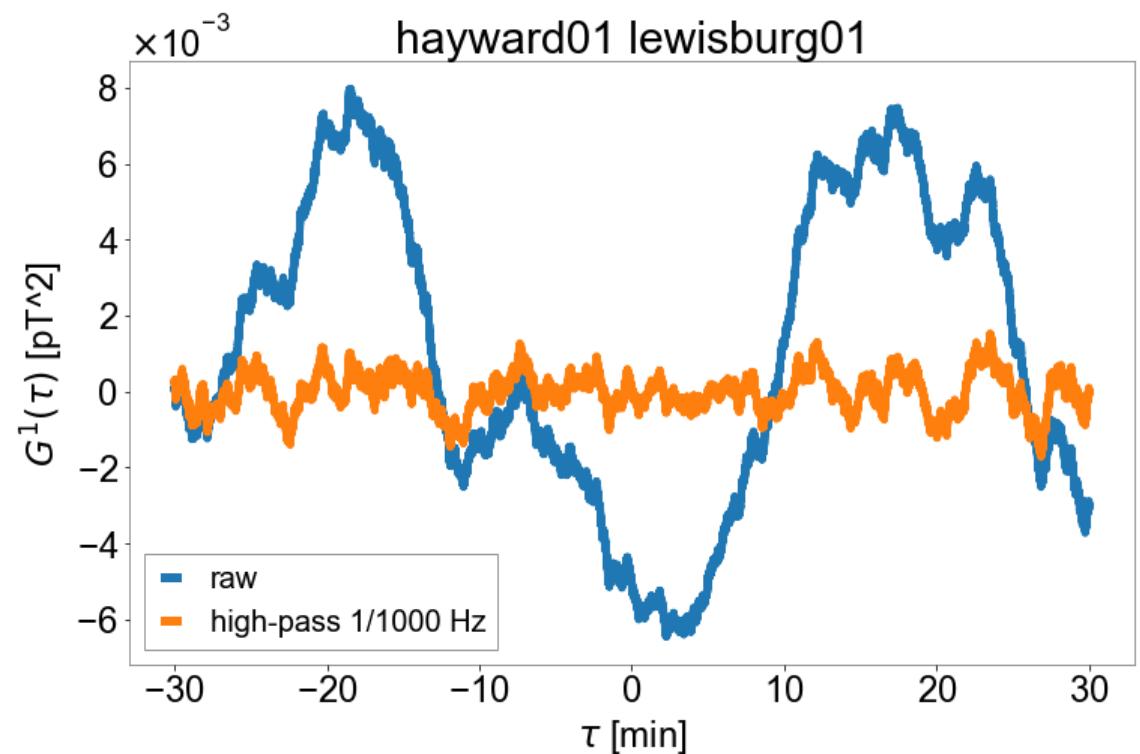
↑                      ↑  
parameters          minimum detectable  
                          pseudo-magnetic field

(defined by the noise  
of the correlation)

# $B_2$ (min) estimation

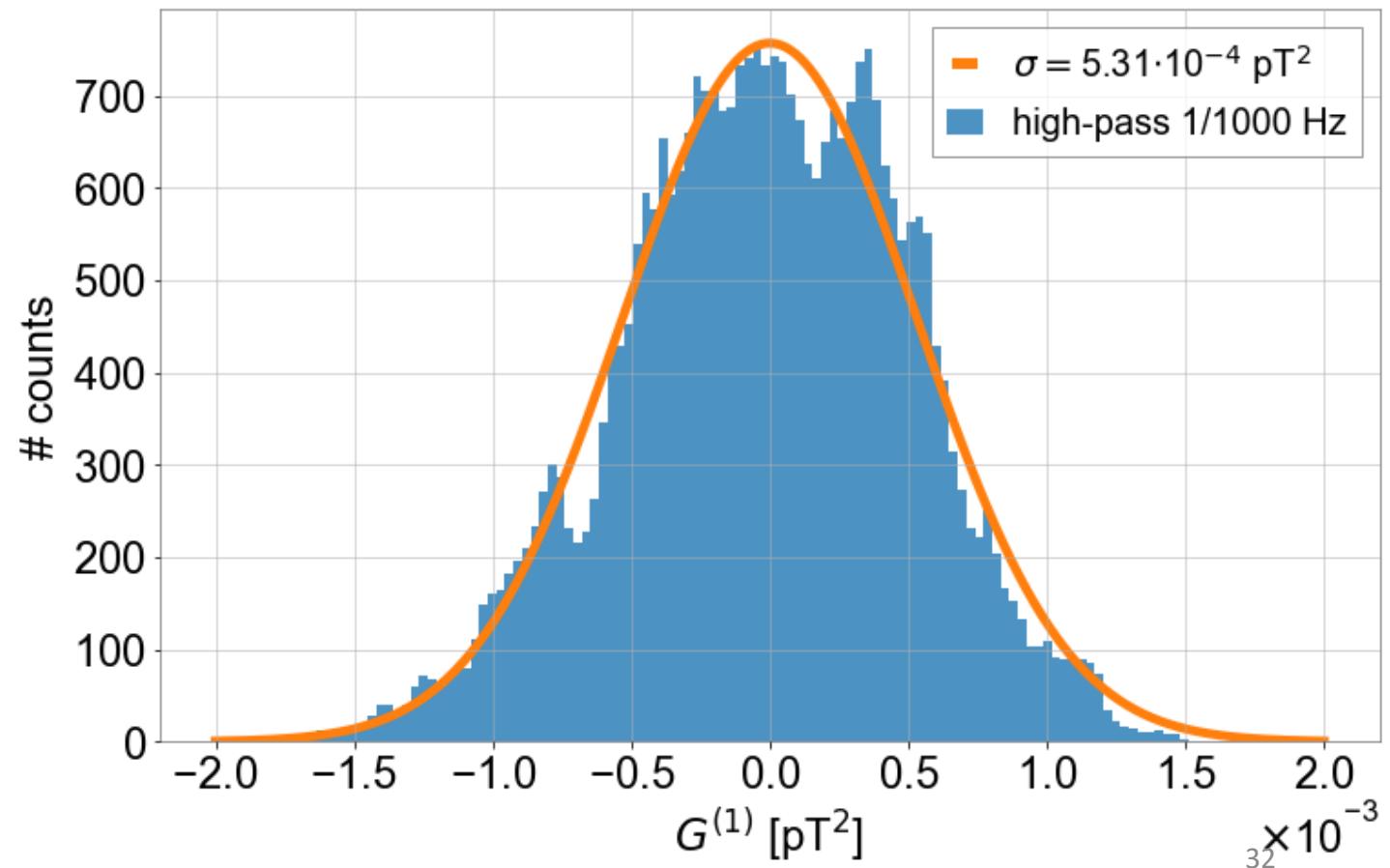


- high-pass filter
- confidence: 3,3355 $\sigma$ ?

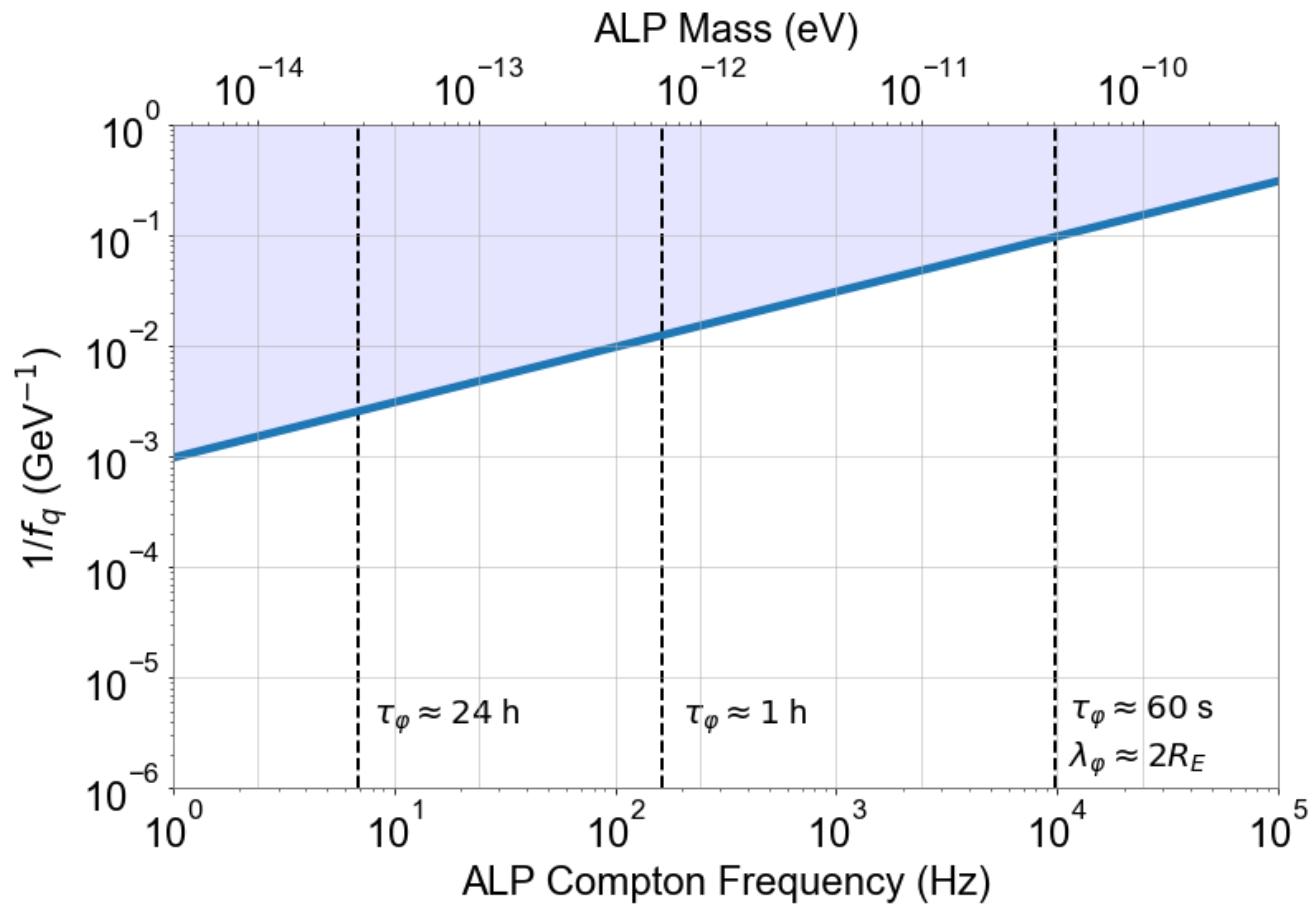


$$C = 90\% \ (1,5\sigma) \rightarrow \beta_g(\min) \approx 30 \text{ fT}$$

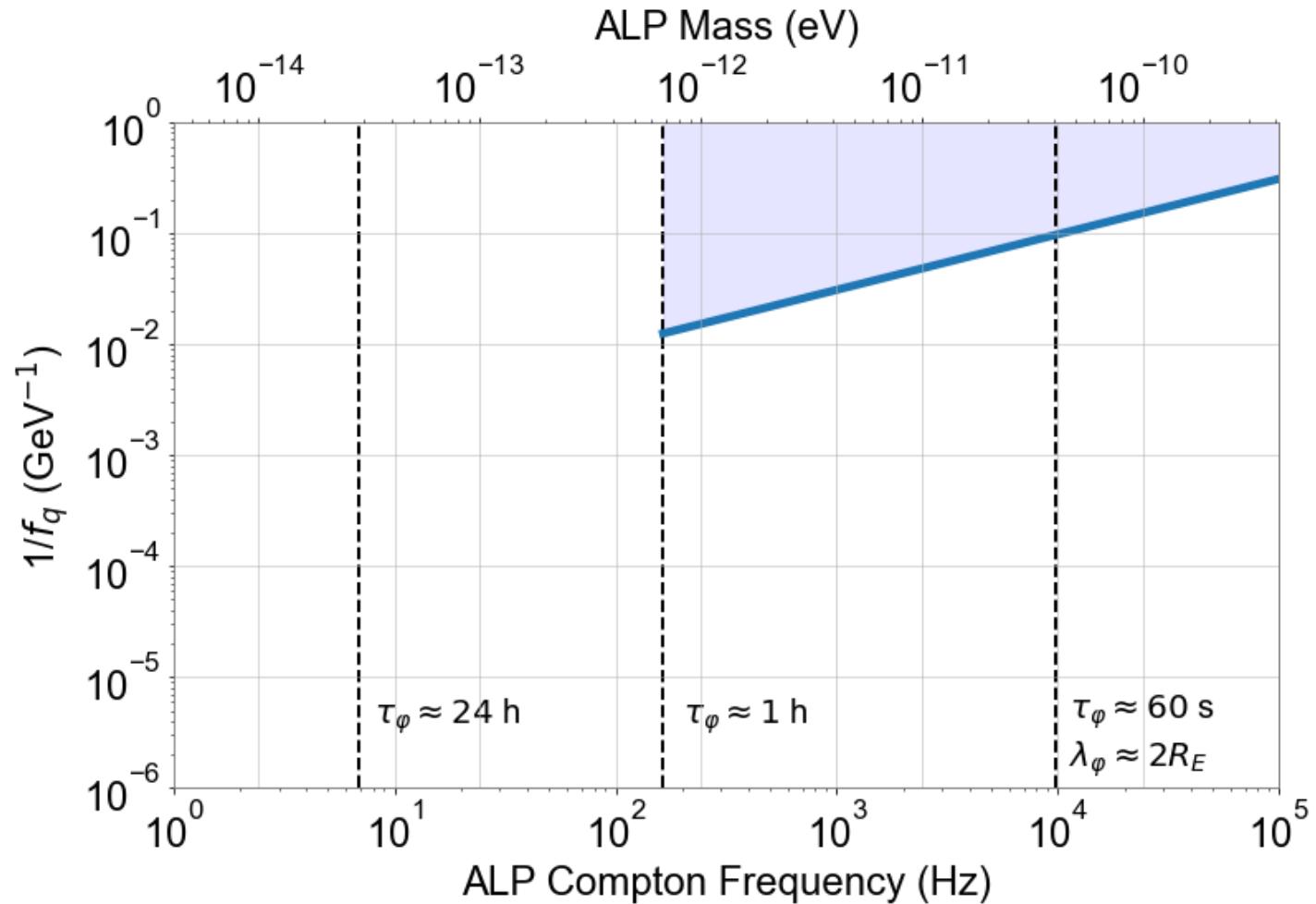
$$C = 99,9\% \ (3,3355\sigma) \rightarrow \beta_g(\min) \approx 42 \text{ fT}$$



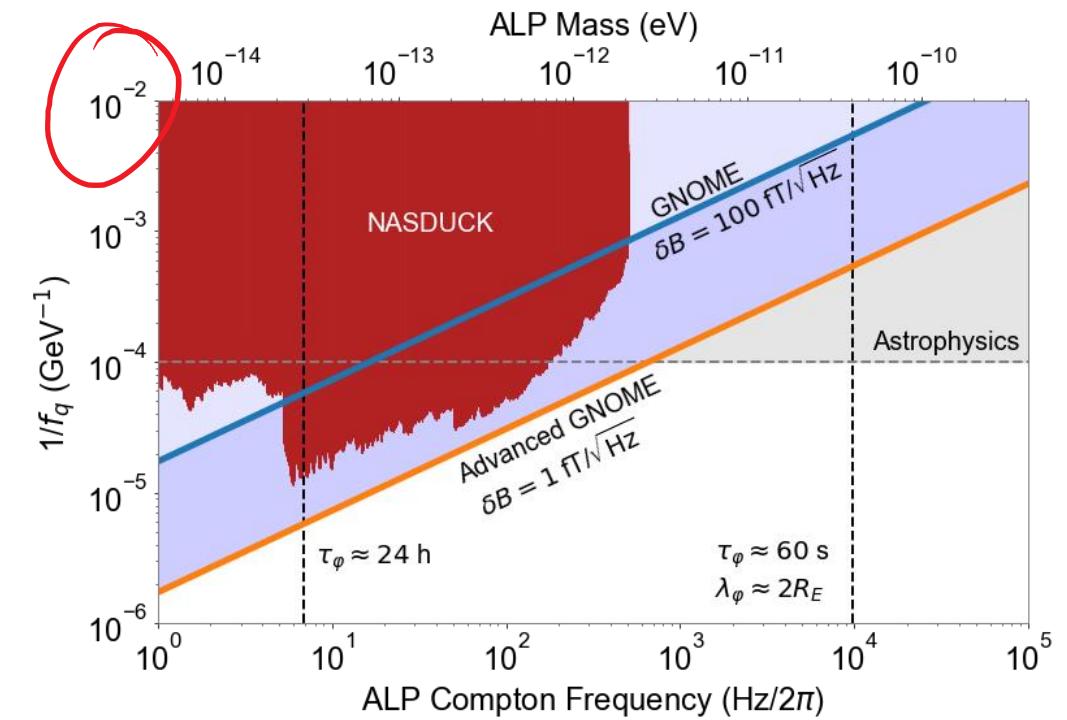
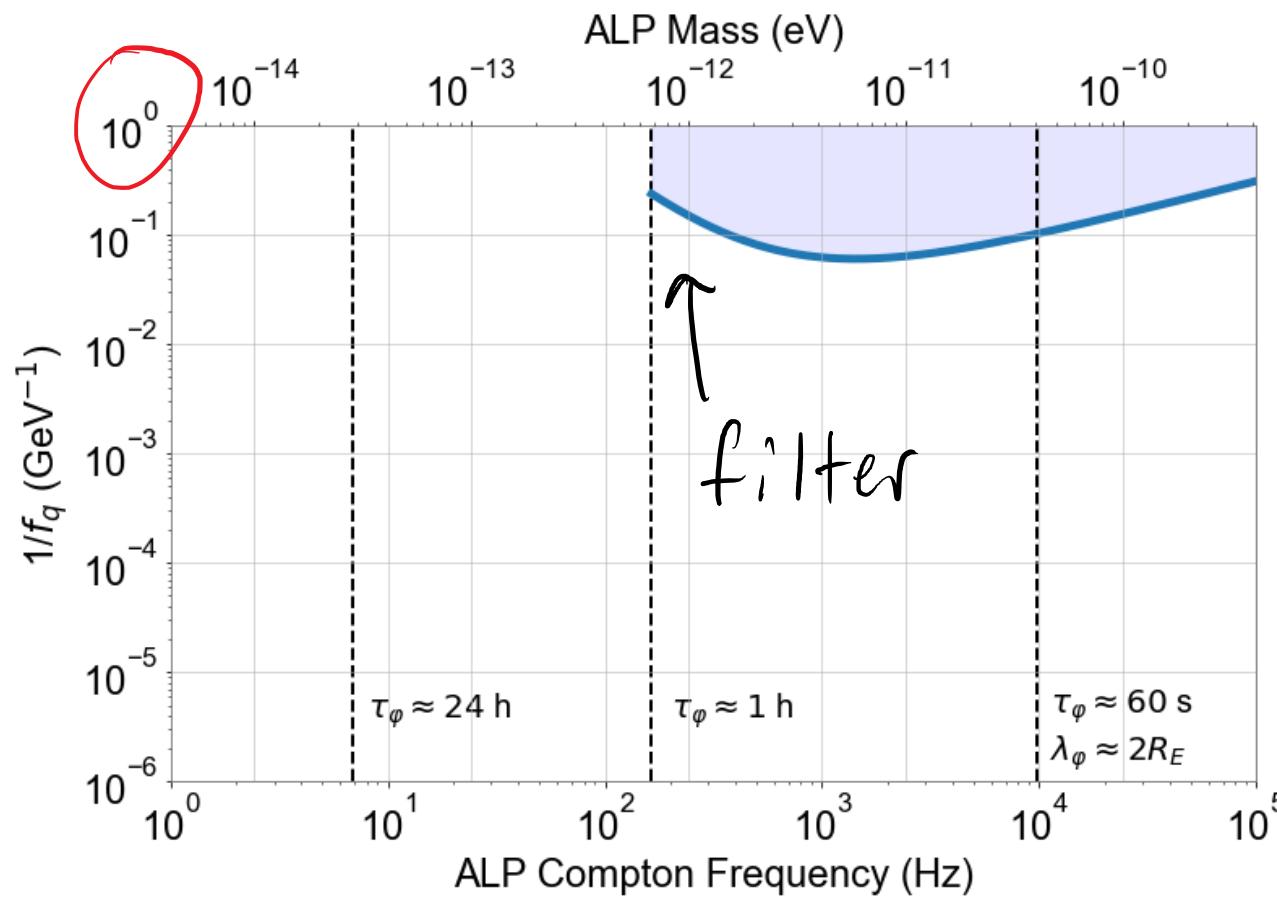
# Toy Model Exclusion Plot



# Toy Model Exclusion Plot



# Toy Model Exclusion Plot



# Outlook

- DM intensity interferometry is possible!
- Fitting lineshape instead of looking at  $G1(0)$  alone
- Need for longer coherence times correlation (Hayward team)
- Low frequency magnetic noise can be suppressed in co-magnetometers (Advanced GNOME)

# Bibliography and Sources

- p. 1: CfA/Rick Peterson / <https://www.cfa.harvard.edu/facilities-technology/telescopes-instruments/very-energetic-radiation-imaging-telescope>
- p. 3 top-right, p. 4 bottom-center: [https://en.wikipedia.org/wiki/Michelson\\_stellar\\_interferometer](https://en.wikipedia.org/wiki/Michelson_stellar_interferometer)
- p. 4 top-right: <https://pl.wikipedia.org/wiki/Betelgeza>
- p. 4 bottom-right: <https://hdl.huntington.org/digital/collection/p15150coll2/id/899>
- p. 3 bottom-right: [https://www.researchgate.net/figure/a-Schematic-diagram-of-the-Michelson-Pease-stellar-interferometer-The-incoming-rays\\_fig3\\_314247837](https://www.researchgate.net/figure/a-Schematic-diagram-of-the-Michelson-Pease-stellar-interferometer-The-incoming-rays_fig3_314247837)
- p.3 left: [https://www.researchgate.net/figure/Effect-of-the-incident-angle-to-the-slit\\_fig2\\_234961268](https://www.researchgate.net/figure/Effect-of-the-incident-angle-to-the-slit_fig2_234961268)
- p. 7 left: <https://www.aanda.org/articles/aa/pdf/2009/45/aa11739-09.pdf>
- p. 9, p. 12 center: <https://www.actaphys.uj.edu.pl/fulltext?series=Reg&vol=29&page=1839>
- p.12 right: <https://en.wikipedia.org/wiki/Sirius>
- p. 13: <https://veritas.sao.arizona.edu/the-science-of-veritas/veritas-results/516-stellarinterferometry>
- p. 14: <https://arxiv.org/pdf/2007.10295.pdf>
- p. 17: <https://medium.com/roaming-physicist/have-we-really-found-dark-matter-46601f62671f>
- p. 21 center: Afach, S., Buchler, B.C., Budker, D. et al. Search for topological defect dark matter with a global network of optical magnetometers. *Nat. Phys.* **17**, 1396–1401 (2021). <https://doi.org/10.1038/s41567-021-01393-y>
- History of HBT interferometry and picture p. 12 left based on lecture by [https://www.ictp-saifr.org/wp-content/uploads/2022/06/SaoPauloJune2022\\_compressed.pdf](https://www.ictp-saifr.org/wp-content/uploads/2022/06/SaoPauloJune2022_compressed.pdf)
- Nasduck exclusion Itay M. Bloch et. al. *New constraints on axion-like dark matter using a Floquet quantum detector*, Science Advances Vol. 8, Issue 5, (2022) <https://doi.org/10.1126/sciadv.abl8919>