

Sources and Detection Prospects for GHz Gravitational Waves

Jan Schütte-Engel

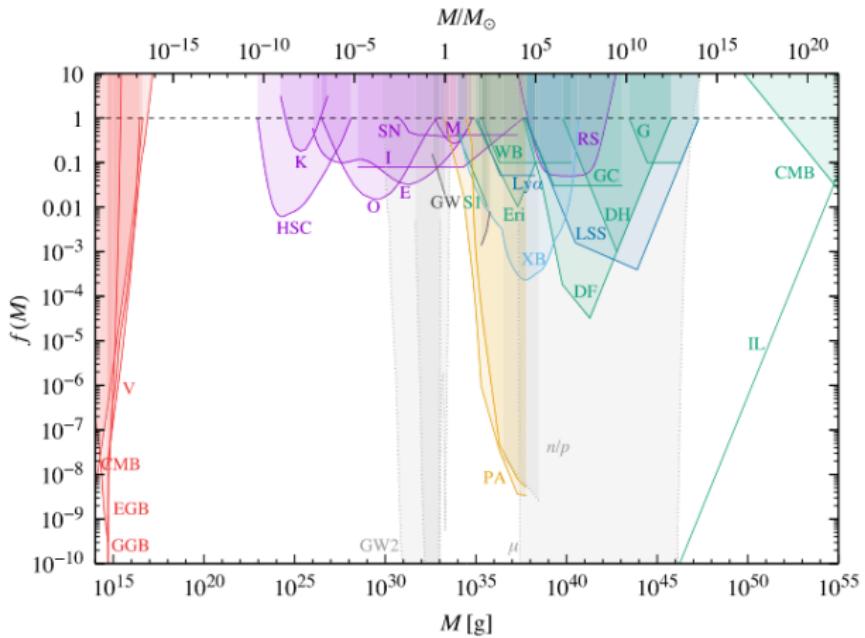
MITP workshop *Probing New Physics with Gravitational Waves*

based on [[Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, JSE 21](#)]

03.08.2022



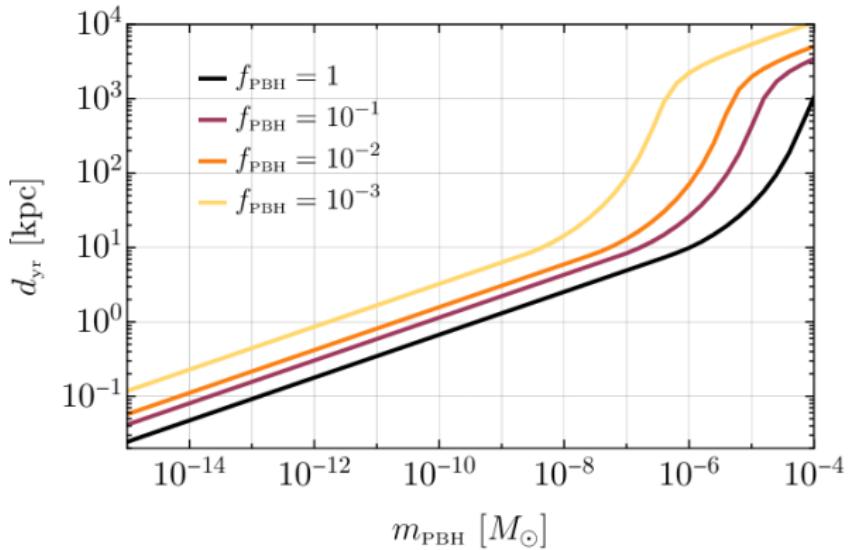
PBHs as dark matter: constraints



[Carr et al. 21]

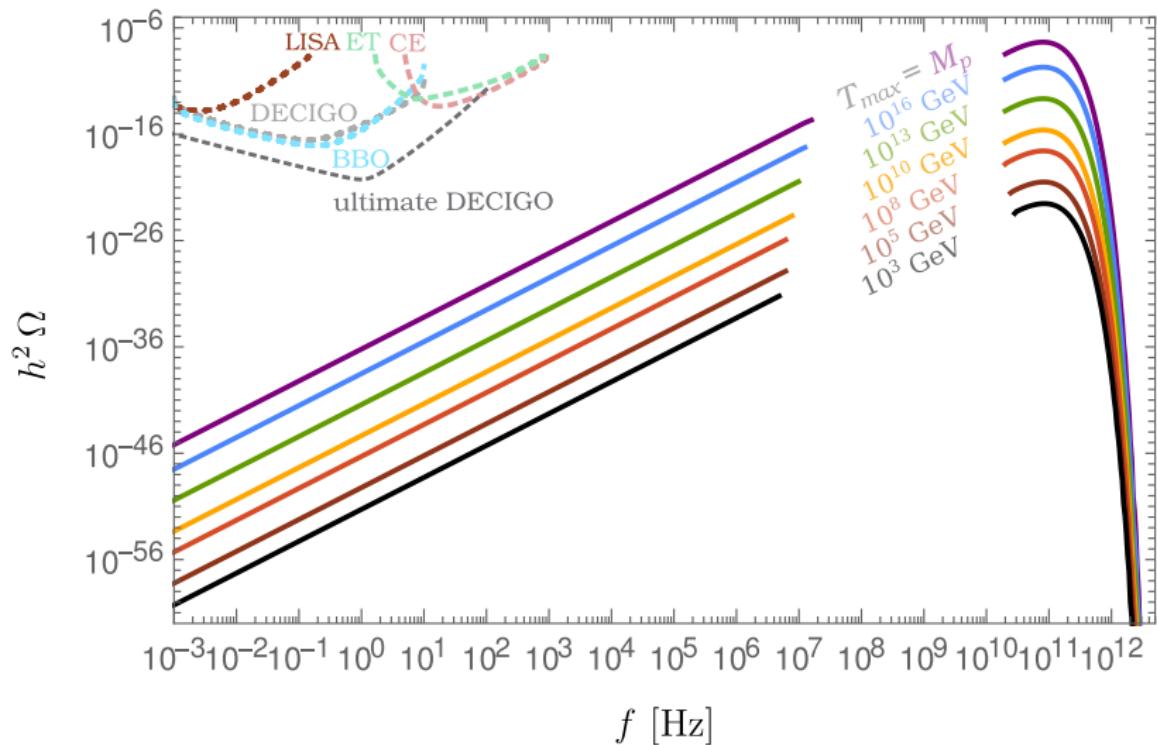
Evaporation (red), lensing (magenta), dynamical effects (green), gravitational waves (black), accretion (light blue), CMB distortions (orange), large-scale structure (dark blue) and background effects (grey).

PBHs merger rate



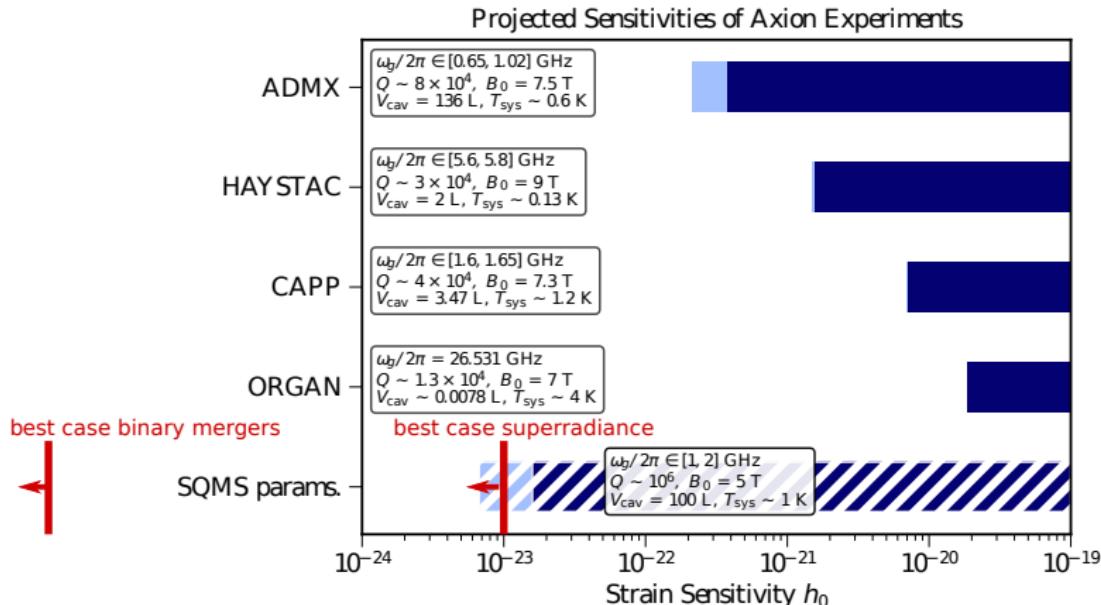
[Franciolini et al. 2205.02153]

CGMB spectrum



[Ringwald, **JSE**, Tamarit 20]

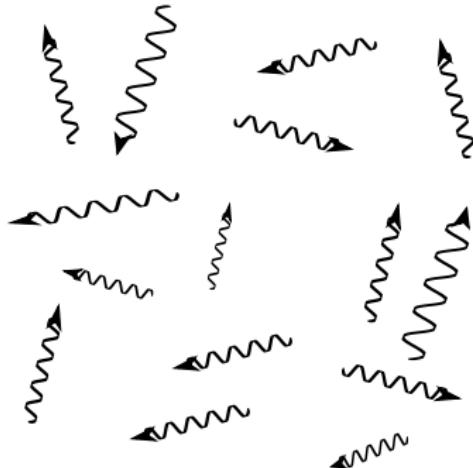
Sensitivity of existing axion experiments



[Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, **JSE 21**]

Existing axion experiments only need to reanalyze their data!

Sensitivity to stochastic GW background

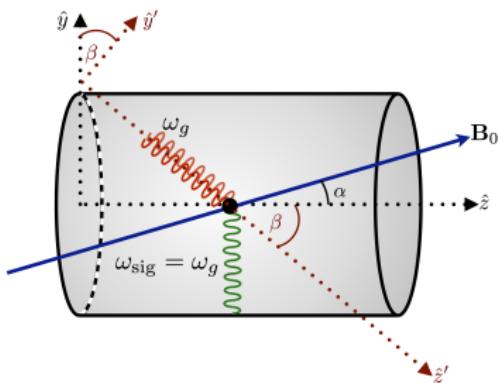


- $\text{SNR} = \frac{S_{\text{sig}}}{S_{\text{noise}}} = \frac{\omega_n Q}{4\pi T} |\eta|^2 B_0^2 V_{\text{cav}} S_h(\omega_g)$
- $\Omega_{\text{GW}}(\omega) = \frac{2}{3} \frac{\omega^3}{H_0^2} S_h(\omega)$

- $\Omega_{\text{GW}} = 8 \times 10^{10} \times \left(\frac{(0.2)^2}{|\eta|^2} \right) \left(\frac{10 \text{ T}}{B_0} \right)^2 \left(\frac{\omega_n}{1 \text{ GHz}} \right)^2 \left(\frac{1 \text{ m}^3}{V_{\text{cav}}} \right) \left(\frac{10^{12}}{Q} \right) \left(\frac{T_{\text{sys}}}{10 \text{ mK}} \right)$
- Cosmologically produced GW backgrounds $\Omega_{\text{GW}} < 10^{-6}$
- Without tricks the detection prospects are not great for stochastic GW backgrounds.

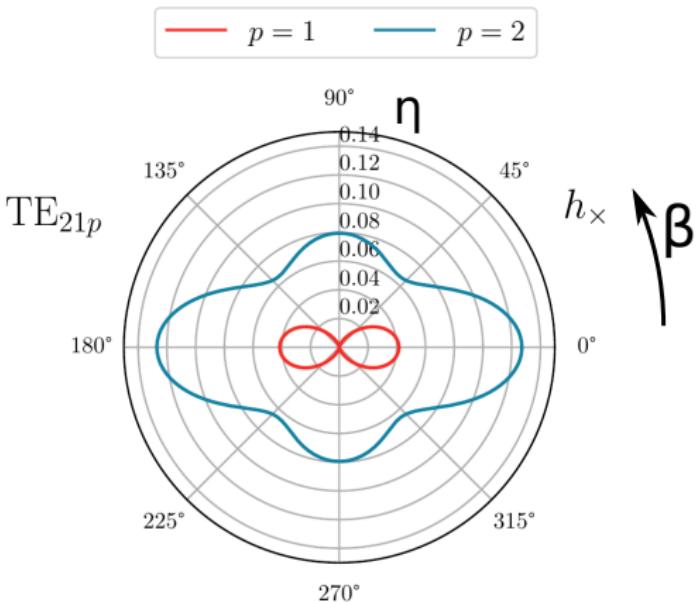
Why the frame matters

Proper detector frame result ($\alpha = 0$)



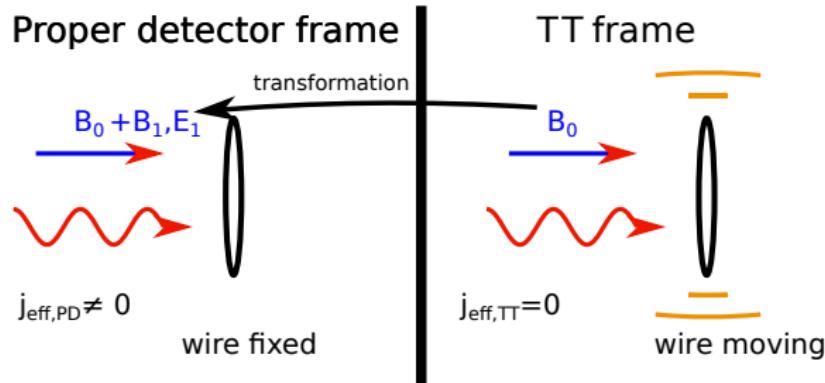
If TT-frame metric is used:

$\alpha = 0, \beta = 0 \Rightarrow \mathbf{j}_{\text{eff}} = 0 \Rightarrow$ No signal.
But this is wrong!



[Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, JSE 21]

Toy example



$$\begin{aligned} t_{\text{TT}} &\simeq t - \frac{i}{4} \omega_g (x^2 - y^2) h_+ e^{i\omega_g t}, \quad x_{\text{TT}} \simeq x - \frac{1}{2} x (1 - i\omega_g z) h_+ e^{i\omega_g t}. \\ y_{\text{TT}} &\simeq y + \frac{1}{2} y (1 - i\omega_g z) h_+ e^{i\omega_g t}, \quad z_{\text{TT}} \simeq z - \frac{i}{4} \omega_g (x^2 - y^2) h_+ e^{i\omega_g t} \end{aligned}$$

- Wire $U_\mu = (1, 0, 0, 0)$
- Signal: E^1 induces current in wire

$$\boldsymbol{E} \simeq \frac{i}{2} B_0 \omega_g h_+ e^{i\omega_g t} (y, x, 0), J_{\text{sig,PD}} = \sigma \boldsymbol{E}$$

- Wire moves $U_{\text{TT},\mu} \neq (1, 0, 0, 0)$
- Signal: moving wire in static B -field induces current in wire

$$J_{\text{sig,TT}}^i \simeq \frac{i}{2} \sigma B_0 \omega_g h_+ e^{i\omega_g t} (y, x, 0)$$

Conclusions

- High frequency GWs very well motivated
- Cavities cannot probe PBH inspirals but test best case PBH superradiance scenarios
- Detection of stochastic GW background with cavities is out of reach
- Signal calculation in cavity: Use proper detector frame metric resummed to all orders
- Existing axion experiments only need to reanalyze data to set limits.

Thank you for your attention

Backup