

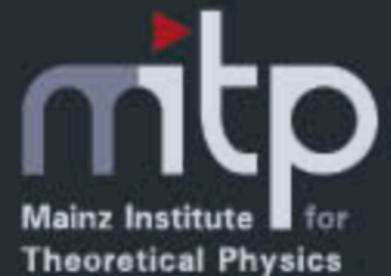
MITP  
SCIENTIFIC  
PROGRAM

Flavor at the Crossroads

19 – 29 April 2022



<https://indico.mitp.uni-mainz.de/e/flavor2022>



# From Light to Heavy $Z'$ to explain $b \rightarrow s\ell^+\ell^-$ data

Claudio Andrea Manzari - 04.28.2022



University of  
Zurich<sup>UZH</sup>



# Heavy $Z'$

<https://inspirehep.net/literature/2014092>

M. Alguero, A. Crivellin, C.A.M., J. Matias

# Heavy Z'

## Setup

- Simple patterns where NP couples solely to muons can explain the discrepancies between the SM and experiment very well.
- Structures with additional LFU contributions can describe data even better...

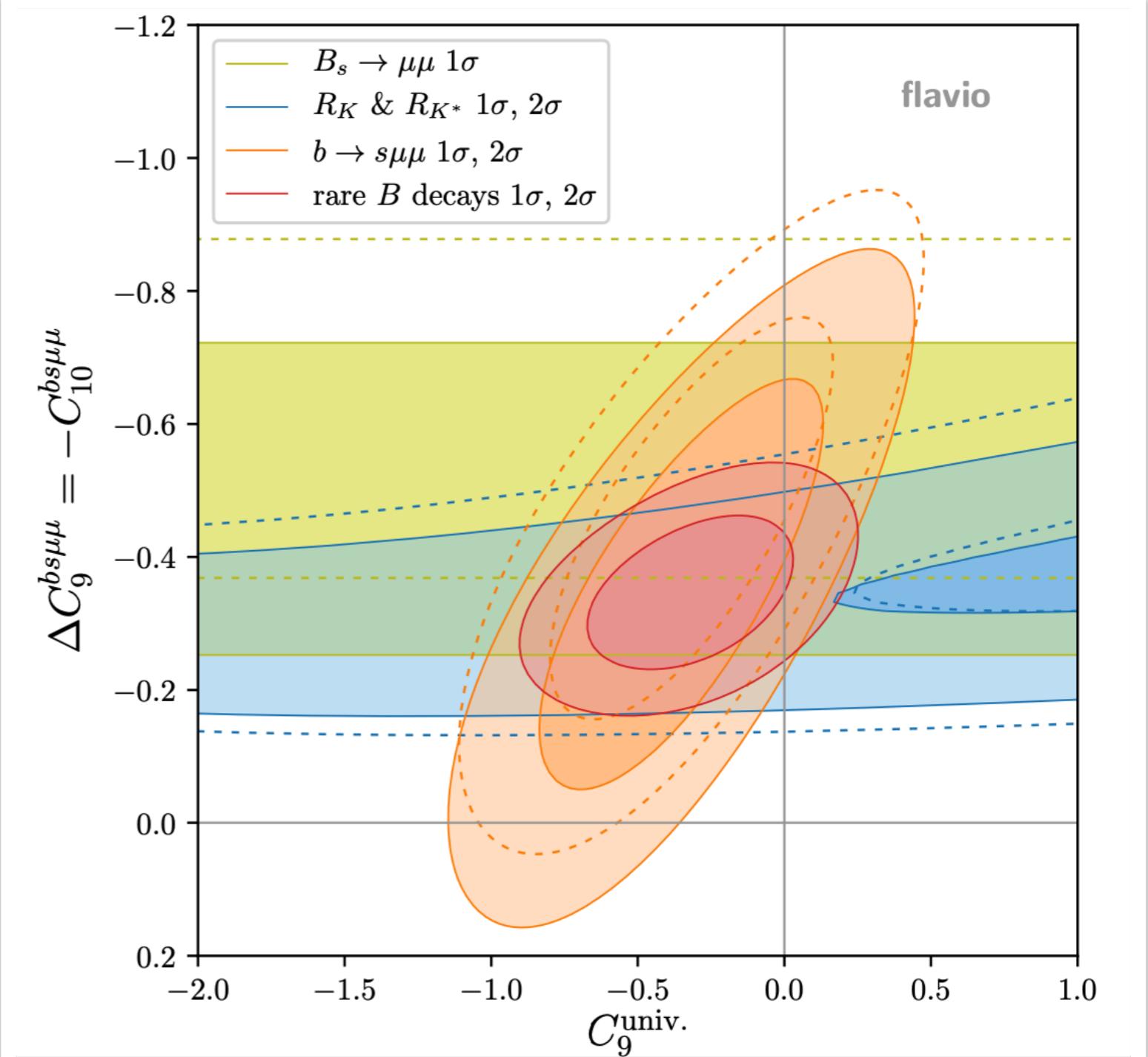
I. 2HDMs

II. Leptoquarks

III.  $SU(2)_L$  Vector Bosons Triplets

IV. Vector-like quarks

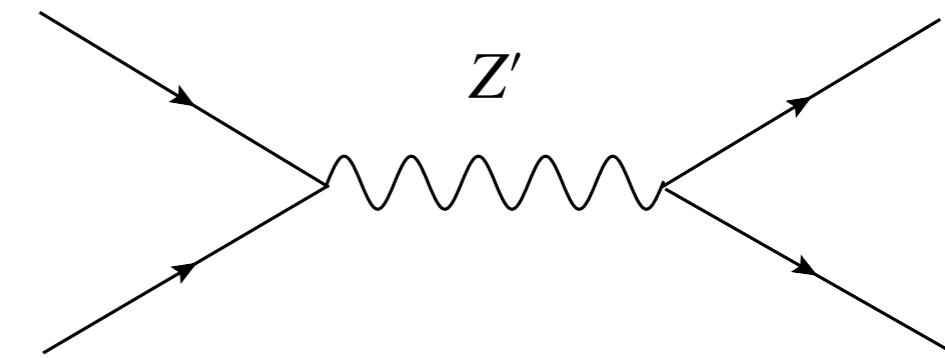
V. Vector Boson Singlet



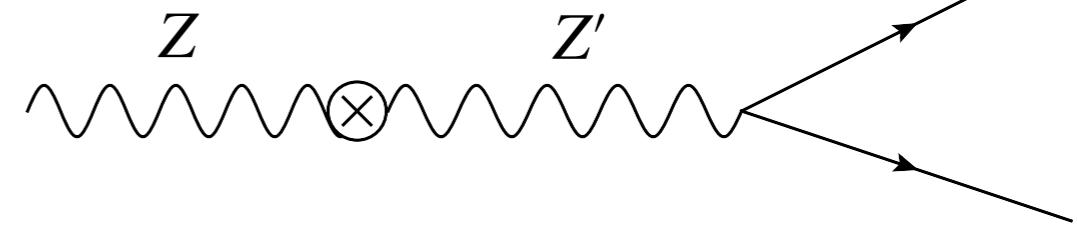
# Heavy $Z'$

## LFUV & LFU

LFUV in  $b \rightarrow s\ell^+\ell^-$



LFU effects via  $Z - Z'$  mixing



## Impact on the EW Fit!

- Constraints from  $Z \rightarrow \ell^+\ell^-$ , etc.
- Prediction of the W mass altered with respect to the SM

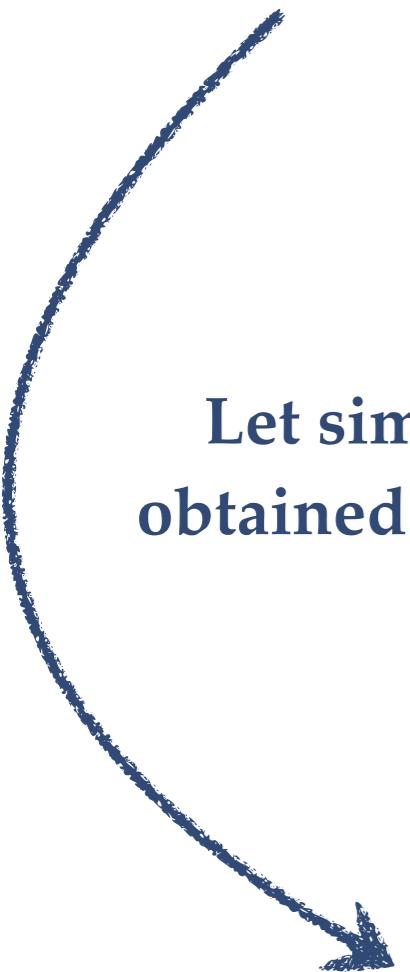
# Heavy $Z'$

## Couplings of the $Z'$

$$\mathcal{L}_{Z'_0}^{\text{int.}} = \bar{u}_j \gamma_\mu (g_{ji}^{uL} P_L + g_{ji}^{uR} P_R) u_i Z_0'^\mu + \bar{d}_j \gamma_\mu (g_{ji}^{dL} P_L + g_{ji}^{dR} P_R) d_i Z_0'^\mu$$

$$+ g_{ji}^{\ell L} (\bar{\nu}_j \gamma_\mu P_L \nu_i) Z_0'^\mu + \bar{\ell}_j \gamma_\mu (g_{ji}^{\ell L} P_L + g_{ji}^{\ell R} P_R) \ell_i Z_0'^\mu$$

Let simplify: the one dimensional scenario with the best fit to data is obtained from a left-handed b-s coupling and a vectorial muon coupling



$$\mathcal{L}_{Z'_0}^{\text{int.}} = \left[ g_{23}^{dL} \bar{s} \gamma_\mu P_L b + g_{23}^{dL*} \bar{b} \gamma_\mu P_L s + g_{ii}^{\ell V} \bar{\ell}_i \gamma_\mu \ell_i \right] Z_0'^\mu$$

# Heavy $Z'$

## Mixing

$$\mathcal{L}_{Z'_0} = -\frac{1}{4} Z'_{0,\mu\nu} Z_0^{\prime\mu\nu} + \frac{\mu_Z'^2}{2} Z'_{0\mu} Z_0^{\prime\mu} + g_{Z'} Z'_{0\mu} Z_0^{\prime\mu} \phi^\dagger \phi - i g_{Z'}^\phi Z_0^{\prime\mu} \phi^\dagger \overleftrightarrow{D}_\mu \phi$$

$$\mathcal{M}^2 = \begin{pmatrix} m_{Z_0}^2 & -\frac{y}{c_W} \\ -\frac{y}{c_W} & m_{Z'_0}^2 \end{pmatrix}, \quad y \equiv \frac{v^2}{2} g_2 g_{Z'}^\phi$$

$$\sin \xi \simeq \frac{y}{c_W m_{Z'_0}^2}$$

$$\begin{pmatrix} Z \\ Z' \end{pmatrix} = \begin{pmatrix} Z'_0 \sin \xi + Z_0 \cos \xi \\ Z'_0 \cos \xi - Z_0 \sin \xi \end{pmatrix} \quad m_Z^2 \simeq m_{Z_0}^2 - \sin \xi^2 m_{Z'_0}^2$$

# Heavy $Z'$

## Electroweak Fit

EW sector completely  
parametrized by 3 parameters  
(+ fermion masses)

$$\alpha, \quad M_{Z_0}, \quad G_F$$



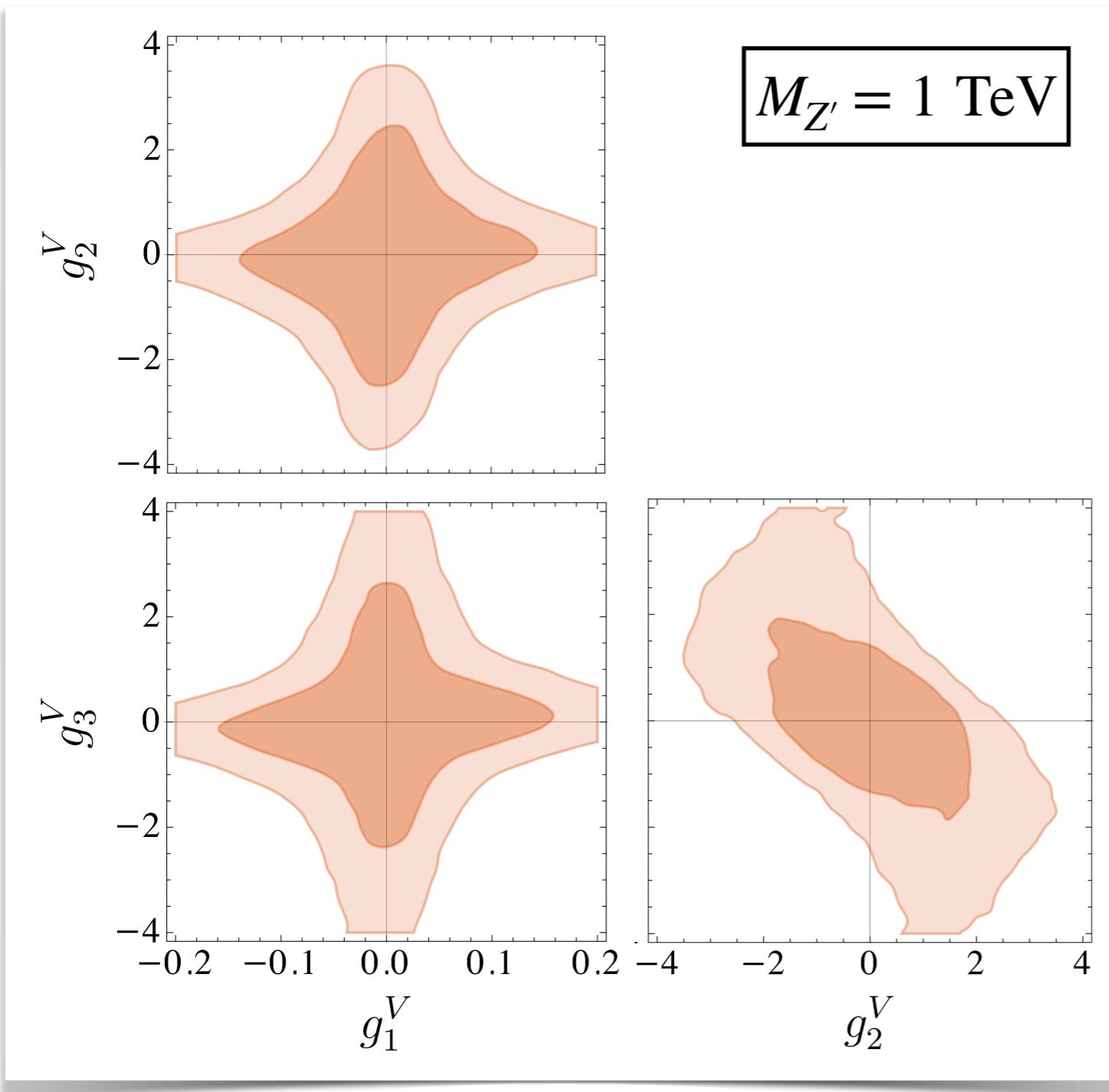
Parameter	Prior
$\Delta\alpha_{\text{had}} \times 10^4$ [16, 17]	276.1(1.1)
$\alpha_s(M_Z)$ [8, 33]	0.1179(10)
$M_H$ [GeV] [8, 38–40]	125.10(14)
$m_t$ [GeV] [8, 41–44]	172.76(30)

*Needed for higher order corrections*

Observable	Measurement	Observable	Measurement
$M_W$ [GeV]	80.379(12)	$\sigma_h^0$ [nb]	41.541(37)
$\Gamma_W$ [GeV]	2.085(42)	$R_e^0$	20.804(50)
$\text{BR}(W \rightarrow \text{had})$	0.6741(27)	$R_\mu^0$	20.785(33)
$\sin^2\theta_{\text{eff}}^e$ (CDF)	0.23248(52)	$R_\tau^0$	20.764(45)
$\sin^2\theta_{\text{eff}}^e$ (D0)	0.23146(47)	$A_{\text{FB}}^{0,e}$	0.0145(25)
$\sin^2\theta_{\text{eff}}^\mu$ (CDF)	0.2315(20)	$A_{\text{FB}}^{0,\mu}$	0.0169(13)
$\sin^2\theta_{\text{eff}}^\mu$ (CMS)	0.2287(32)	$A_{\text{FB}}^{0,\tau}$	0.0188(17)
$\sin^2\theta_{\text{eff}}^\mu$ (LHCb)	0.2314(11)	$R_b^0$	0.21629(66)
$P_\tau^{\text{pol}}$	0.1465(33)	$R_c^0$	0.1721(30)
$A_e$	0.1516(21)	$A_{\text{FB}}^{0,b}$	0.0992(16)
$A_\mu$	0.142(15)	$A_{\text{FB}}^{0,c}$	0.0707(35)
$A_\tau$	0.136(15)	$A_b$	0.923(20)
$\Gamma_Z$ [GeV]	2.4952(23)	$A_c$	0.670(27)

# Heavy $Z'$

## Couplings to Leptons



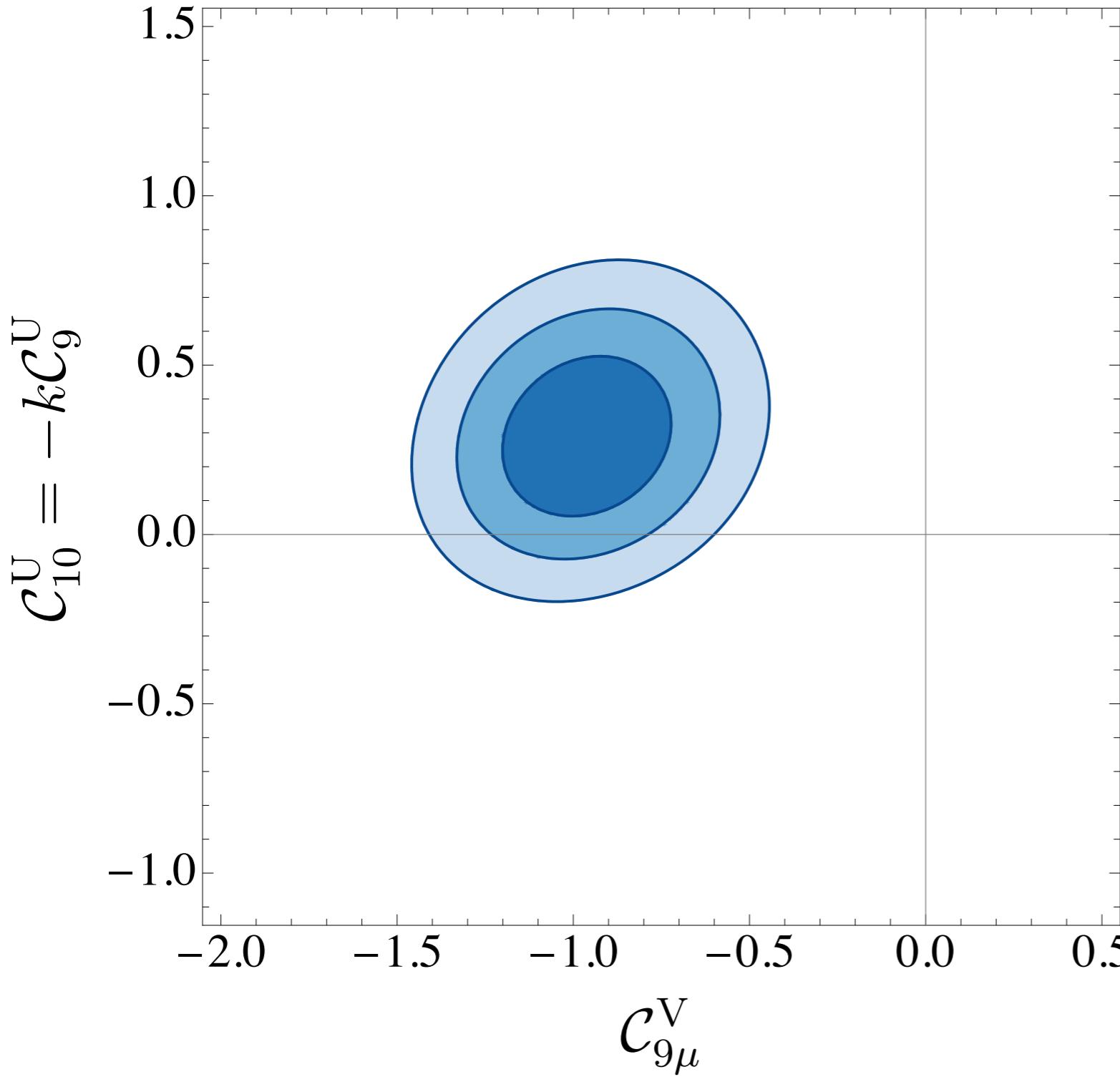
- EW Precision Observables
- Neutrino Trident Production
- LEP-II
- Lepton decays:  
 $\tau \rightarrow \mu \nu_\tau \bar{\nu}_\mu, \tau \rightarrow e \nu_\tau \bar{\nu}_e, \mu \rightarrow e \nu_\mu \bar{\nu}_e$

$$g_1^V \ll g_{2,3}^V \quad \& \quad g_2^V \simeq -g_3^V \equiv g'$$

$L_\mu - L_\tau$

# Heavy $Z'$

$b \rightarrow s\ell^+\ell^-$



$$C_{9\mu}^V = - \frac{\pi^2}{e^2} \frac{4\sqrt{2}g_{23}^{dL}g_{22}^{\ell V}}{G_F m_{Z'}^2 V_{tb} V_{ts}^*}$$

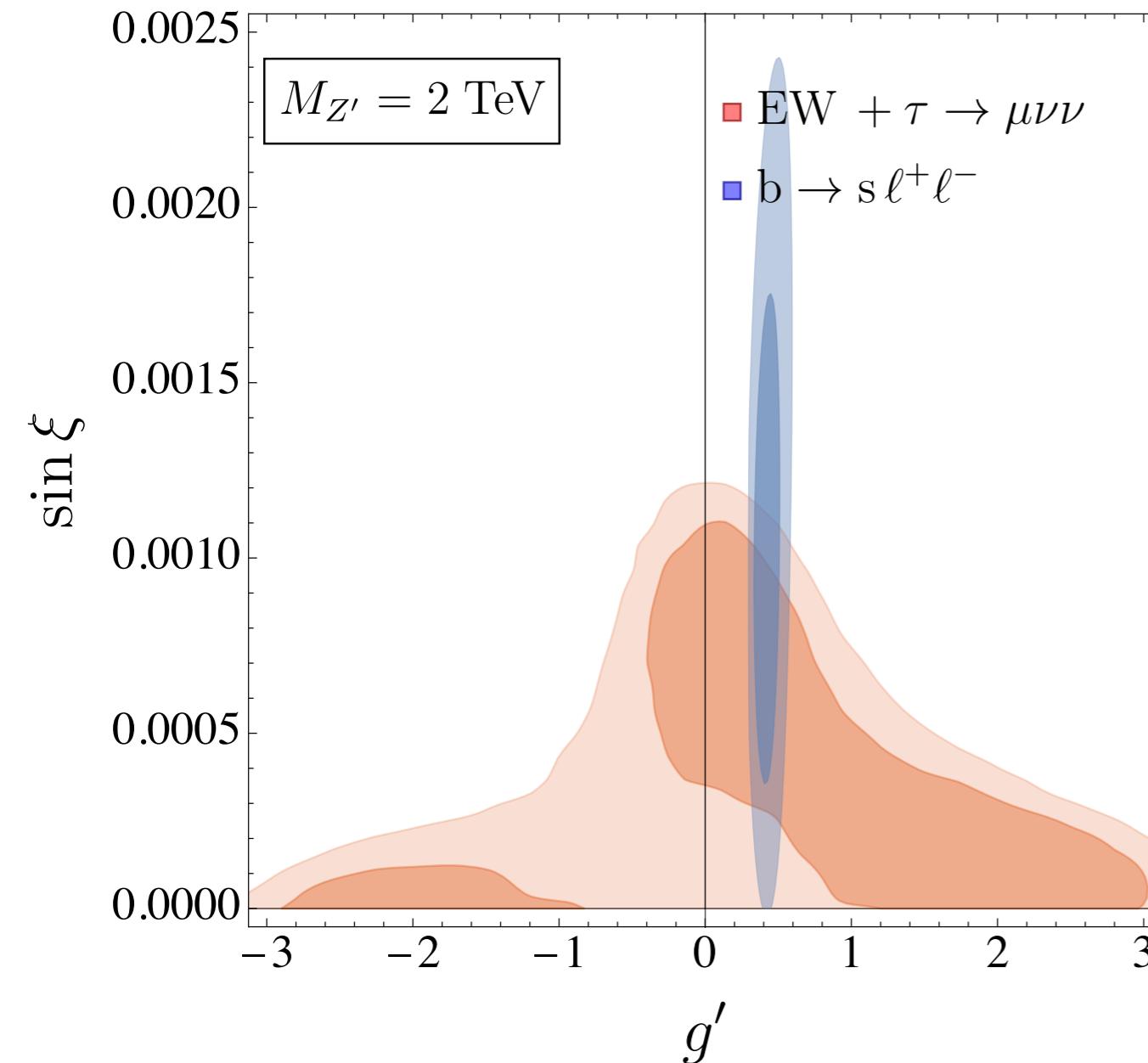
$$\begin{aligned} C_{10}^U &= -k C_9^U \\ &= -\frac{\pi^2}{e^2} \frac{\sqrt{2}g_{23}^{dL} g_2 \sin\xi}{c_W G_F m_{Z'}^2 V_{tb} V_{ts}^*} \end{aligned}$$

$$k = \frac{1}{1 - 4 s_W^2}$$

# Heavy $Z'$

## Results

PRE - CDF 2022

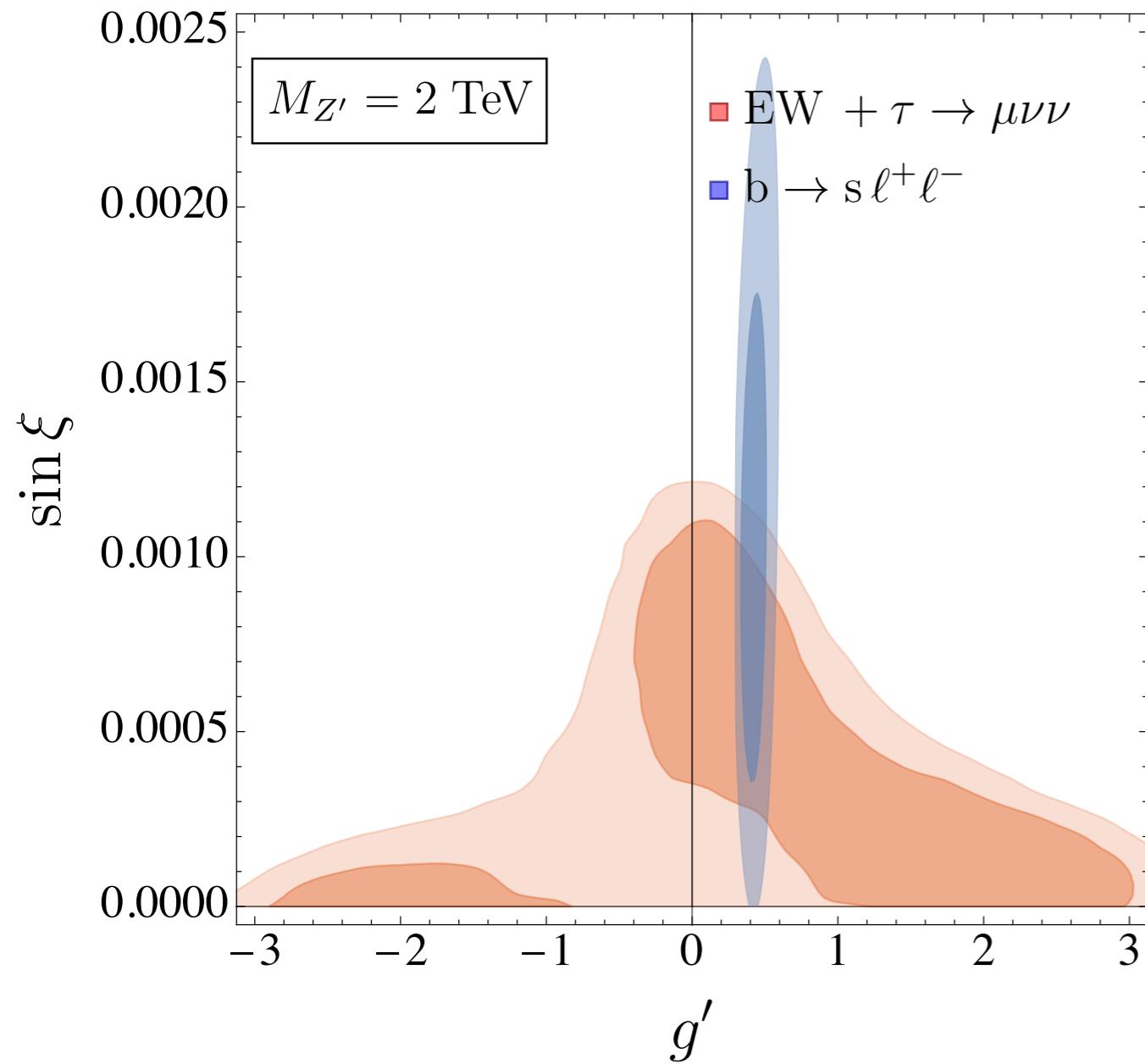


$$m_W^2 \simeq m_Z^2 c_W^2 + \sin \xi^2 m_{Z'_0}^2 c_W^2$$

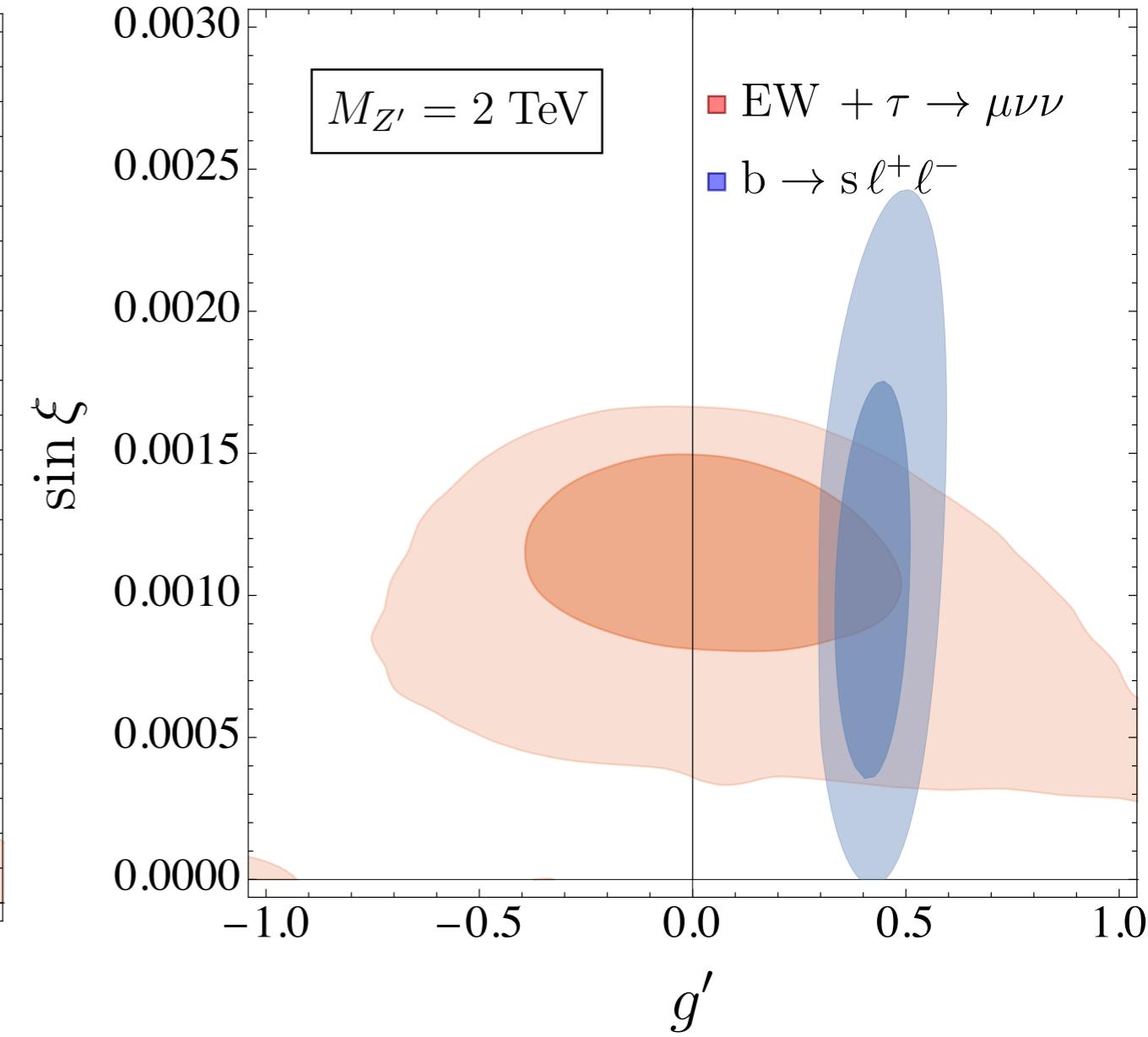
# Heavy $Z'$

## Results

**PRE - CDF 2022**



**POST - CDF 2022**



# Light $Z'$

<https://inspirehep.net/literature/2039365>

W. Altmannshofer, A. Crivellin, C.A.M., G. Inguglia, P. Feichtinger, J.M. Camalich

# Light $Z'$

## Setup

$$m_{Z'} \lesssim 6 \text{ GeV}$$

$$\mathcal{L}_{Z'} \supset \left[ \bar{\mu} \left( g_{\mu\mu}^V \gamma^\mu + g_{\mu\mu}^A \gamma^\mu \gamma_5 \right) \mu + g_{sb}^L \bar{s} \gamma^\mu P_L b + g_\chi \bar{\chi} \gamma^\mu \chi \right] Z'_\mu$$

## Observables

- $b \rightarrow s \ell^+ \ell^-$
- $(g - 2)_\mu$
- $B \rightarrow K^{(*)} + \text{invisible}$
- $pp \rightarrow \mu^+ \mu^- + \text{invisible}$
- $B_s - \bar{B}_s$  mixing
- $e^+ e^- \rightarrow \mu^+ \mu^- + \text{invisible}$

# Light $Z'$

$m_{Z'} \lesssim 6 \text{ GeV}$

$b \rightarrow s\ell^+\ell^-$  &  $(g - 2)_\mu$

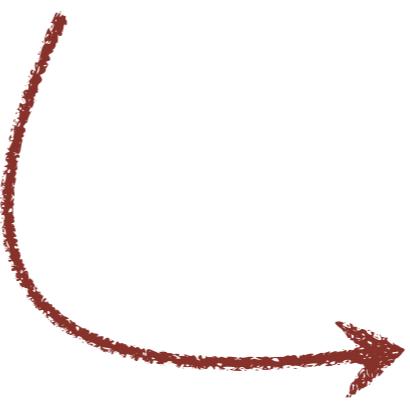
$$\mathcal{L}_{Z'} \supset \left[ \bar{\mu} \left( g_{\mu\mu}^V \gamma^\mu + g_{\mu\mu}^A \gamma^\mu \gamma_5 \right) \mu + g_{sb}^L \bar{s} \gamma^\mu P_{L,R} b + g_\chi \bar{\chi} \gamma^\mu \chi \right] Z'_\mu$$

$b \rightarrow s\ell^+\ell^-$

$(g - 2)_\mu$

$$C_{9(10)} = \frac{g_{sb}^L g_{\mu\mu}^{V(A)}}{q^2 - m_{Z'}^2 + i m_{Z'} \Gamma_{Z'}}$$

$$\Delta a_\mu = \frac{m_\mu^2}{12\pi^2 M_{Z'}^2} \text{Re} \left[ (g_{\mu\mu}^V)^2 - 5(g_{\mu\mu}^A)^2 \right]$$

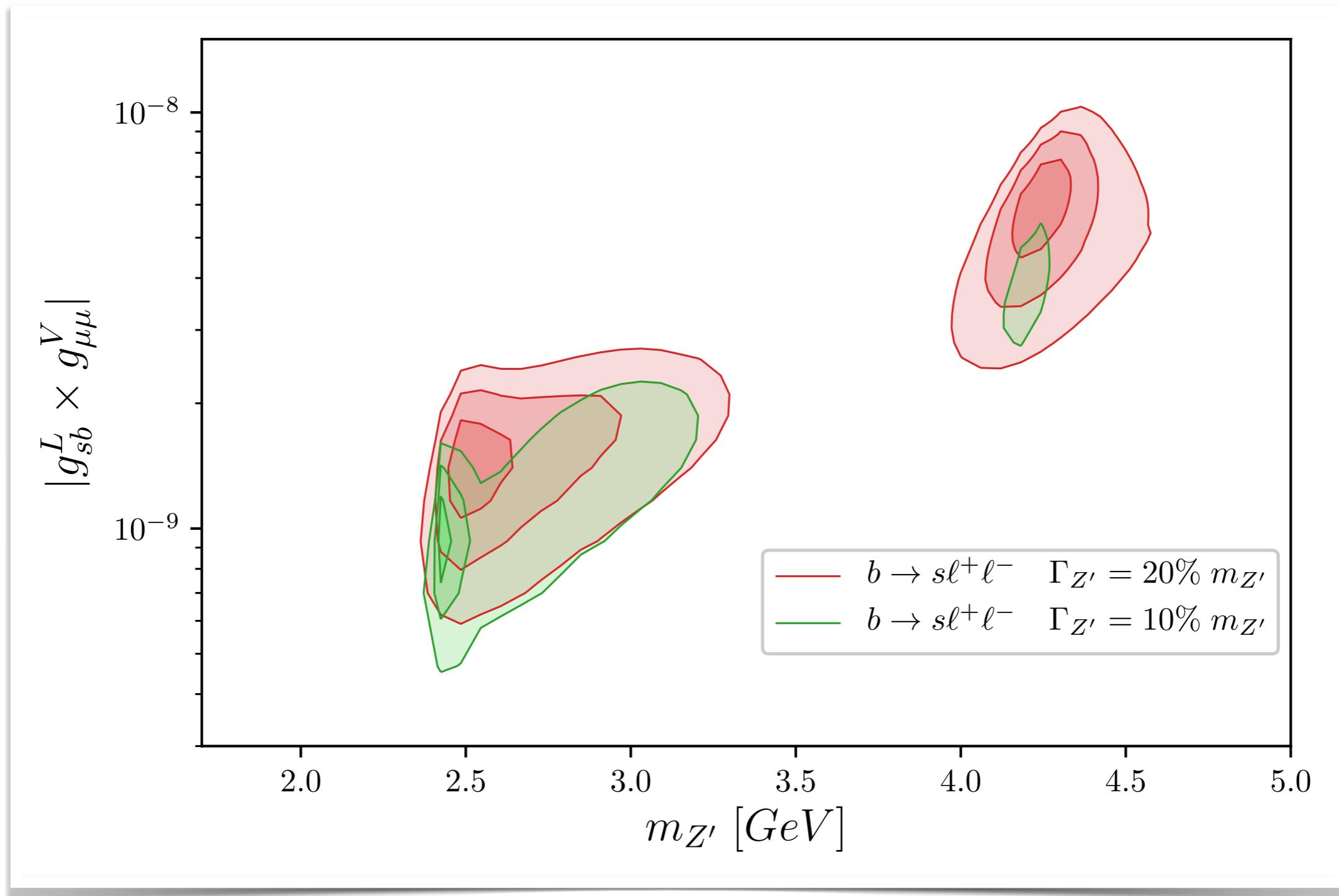


$$g_V = -\sqrt{5}g_A$$

# Light $Z'$

$b \rightarrow s\ell^+\ell^-$

$$C_9 = -\sqrt{5}C_{10}$$



# Light $Z'$

$$B \rightarrow K^{(*)}\nu\bar{\nu}$$

$$\mathcal{B}(B \rightarrow K^{(*)}\chi\bar{\chi}) = \mathcal{B}(Z' \rightarrow \chi\bar{\chi}) \int_{s_{\min}}^{s_{\max}} ds \Gamma_{Z'}(s) \text{BW}(s) \mathcal{B}(B \rightarrow K^{(*)}Z')(s)$$

$$s_{\min} = 4m_\chi^2$$

$$s_{\max} = (m_B - m_{K^{(*)}})^2$$

$$\text{BW}(s) = \frac{\sqrt{s}}{\pi} \frac{1}{(s - m_{Z'}^2)^2 + \Gamma_{Z'}(s)^2 m_{Z'}^2}$$

$$\Gamma_{Z'}(s) = \frac{g_\chi^2}{12\pi\sqrt{s}} \sqrt{1 - 4\frac{m_\chi^2}{s}} (s + 2m_\chi^2)$$

**Belle II:**

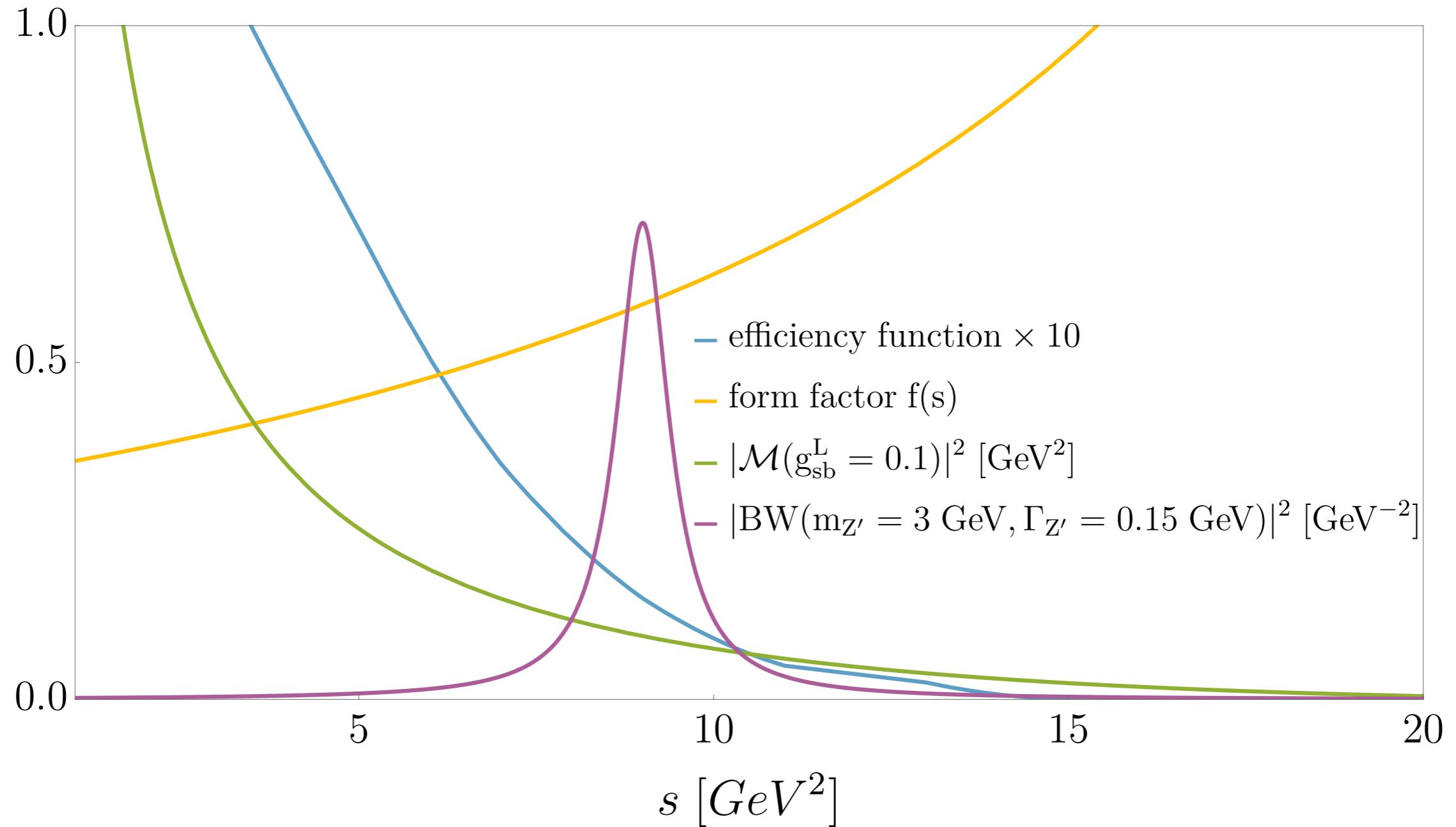
$$\mathcal{B}(B^+ \rightarrow K^+\nu\bar{\nu}) < 4.1 \times 10^{-5}$$

Others: Babar, Belle

# Light $Z'$

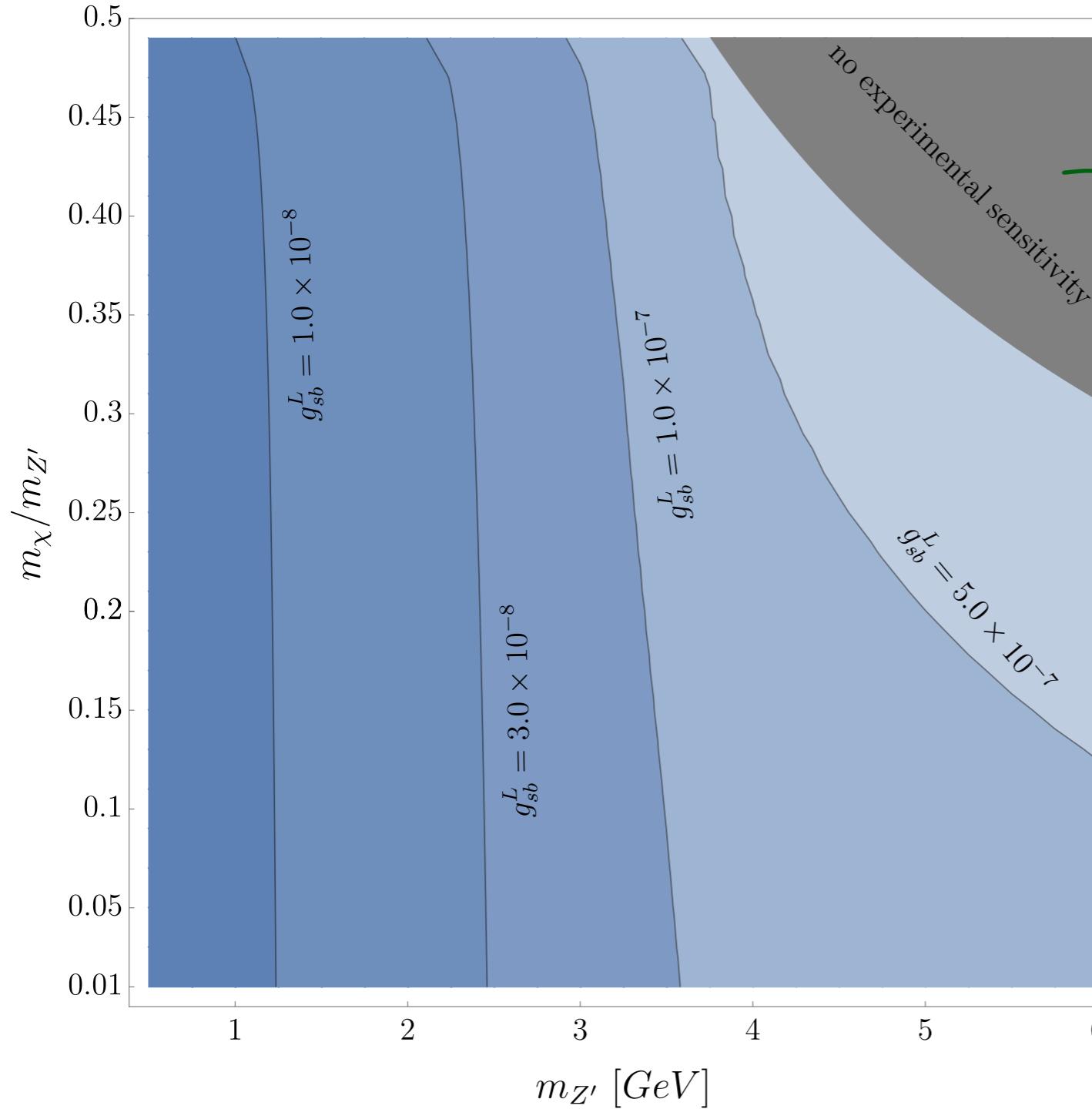
$$B \rightarrow K^{(*)}\nu\nu$$

$$B \rightarrow K\nu\nu$$

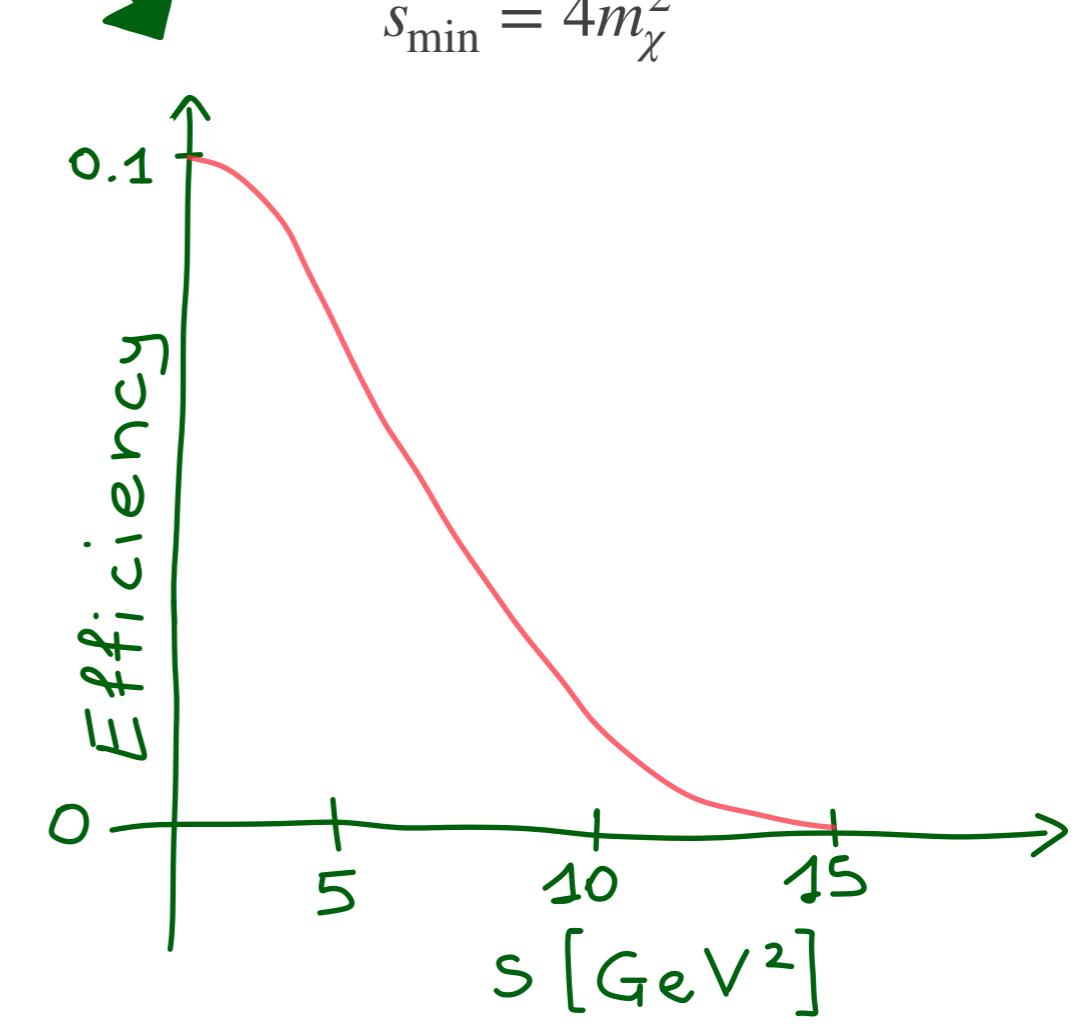


# Light $Z'$

## Recasted Belle II analysis



- A large  $m_\chi$  weakens the bounds on  $g_{sb}^L$



# Light $Z'$

## Other bounds

### $B_s - \bar{B}_s$ mixing

- Light  $Z'$  masses: OPE in  $m_{Z'}/m_B \rightarrow$  weaker than  $B \rightarrow K^{(*)}\nu\nu$
- For  $m_{Z'} \sim m_B$ : no reliable treatment yet

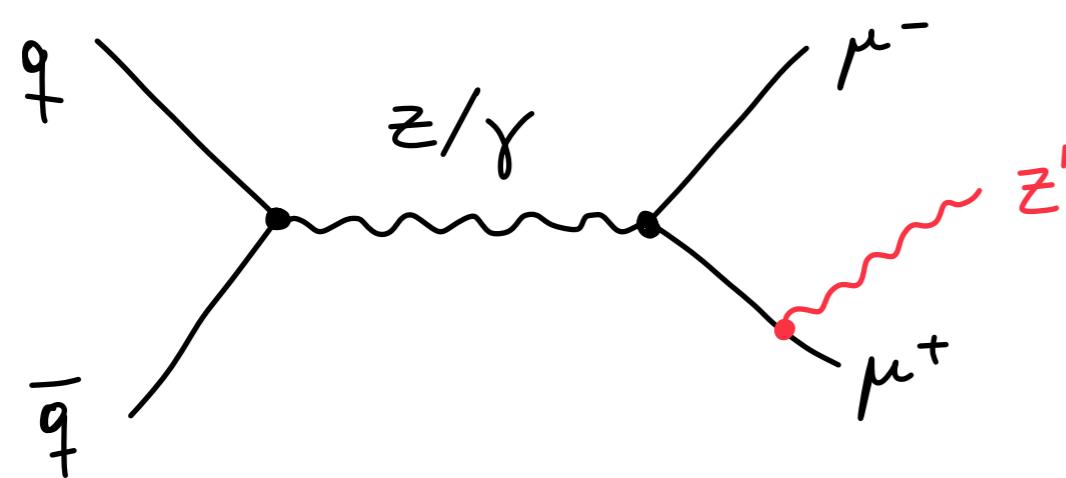
$e^+e^- \rightarrow \mu^+\mu^- + \text{invisible}$

- Recasted Belle II analysis for a sizeable  $Z'$  width:
  - $276 \text{ pb}^{-1}$  [hep/1912.11276](#)
  - $50 \text{ fb}^{-1}$
  - $5 \text{ ab}^{-1}$

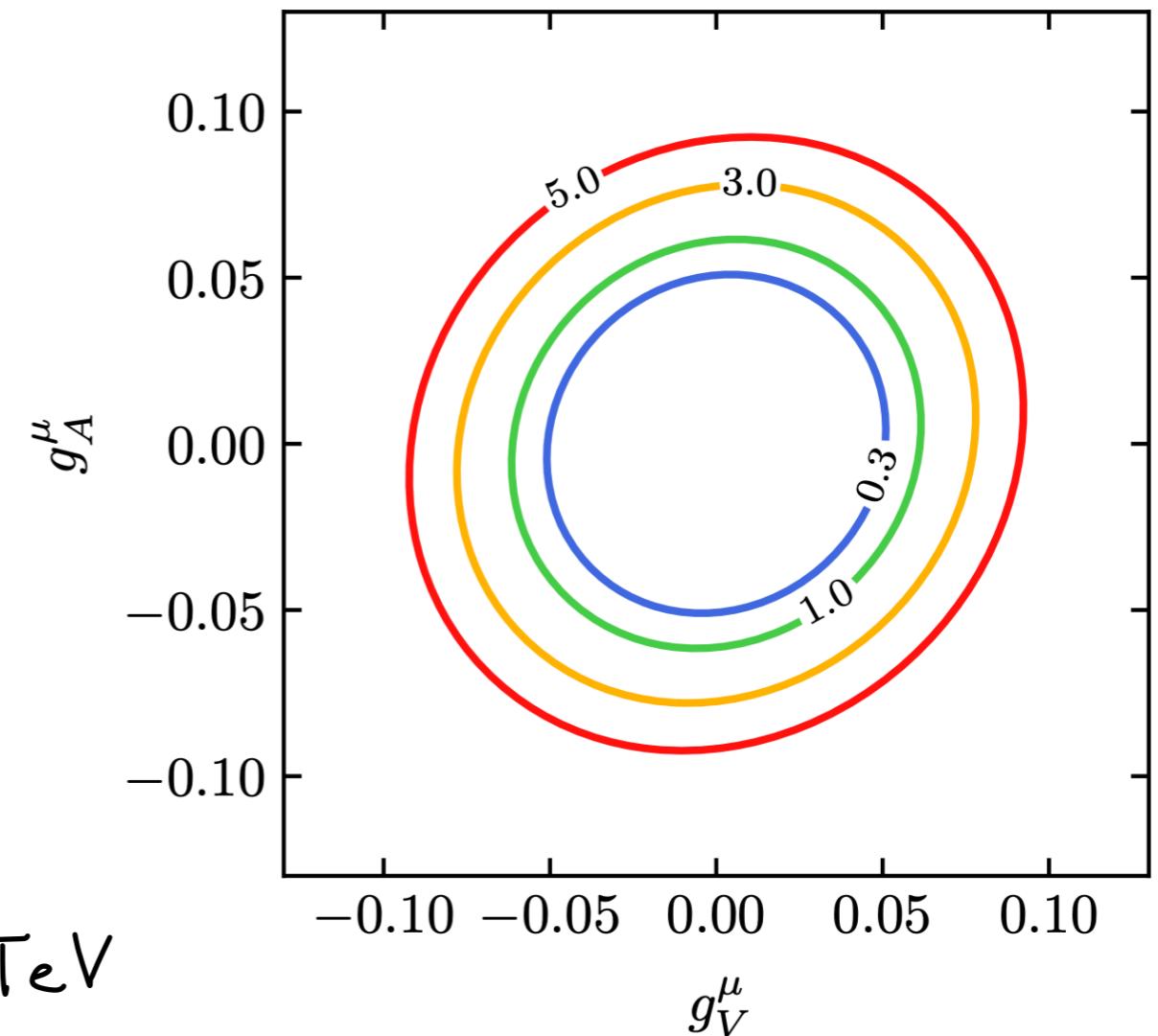
**Limits on  $g_{\mu\mu}^V$  of the order  $10^{-1} - 10^{-2}$**

$pp \rightarrow \mu^+\mu^- + \text{invisible}$

F. Bishara, U. Haisch, P.F. Monni [hep/1705.03465](#)



ATLAS  $3.2 \text{ fb}^{-1} \sqrt{s} = 13 \text{ TeV}$



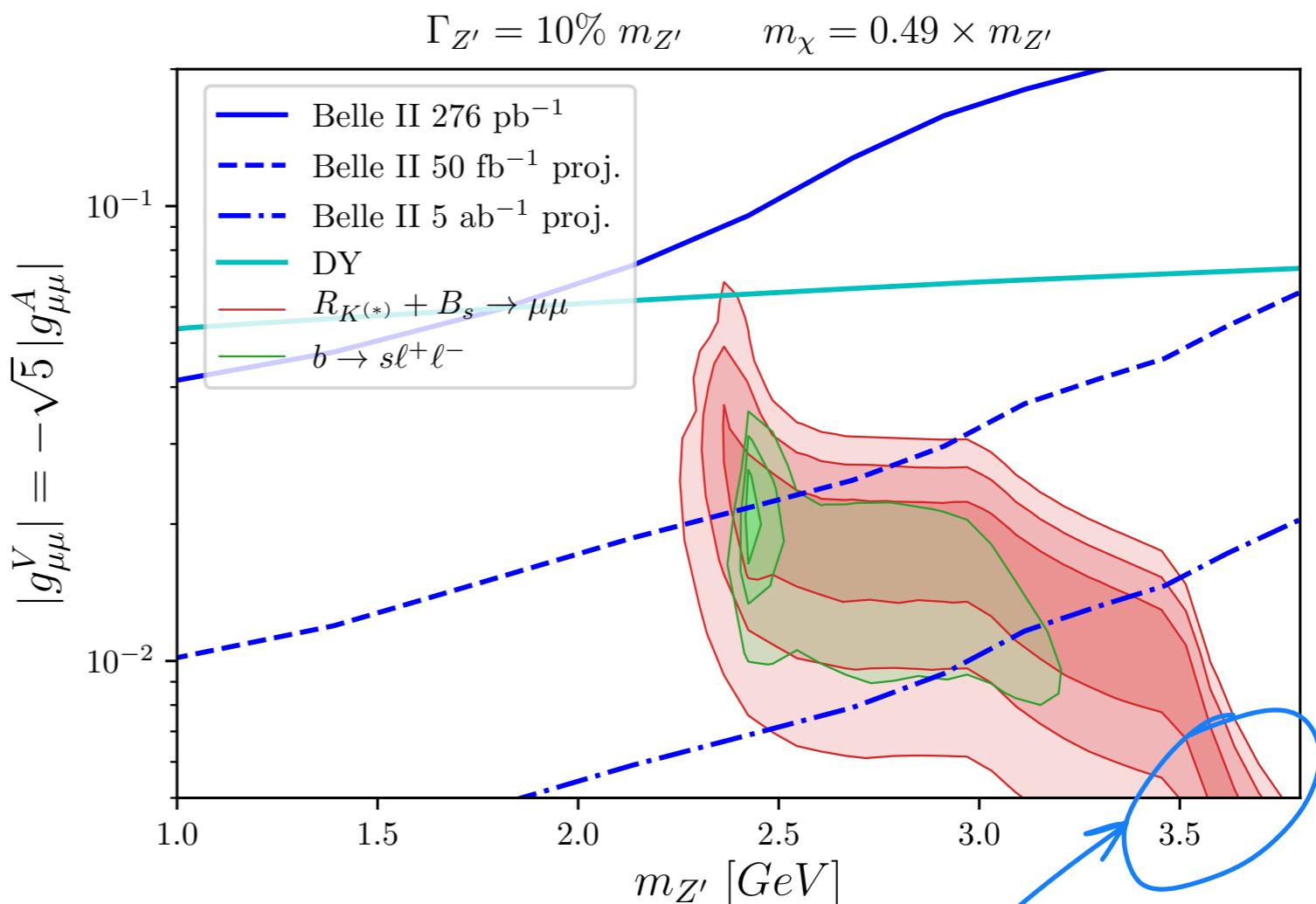
# Light $Z'$

## Bounds on $g_{\mu\mu}^V$

Maximum  $g_{bs}^L$  allowed by  $B \rightarrow K^{(*)}\nu\nu$

Cannot exclude a light  $Z'$  yet

$$\Gamma_{Z'} = 15\% m_{Z'} \quad m_\chi = 0.10 \times m_{Z'}$$



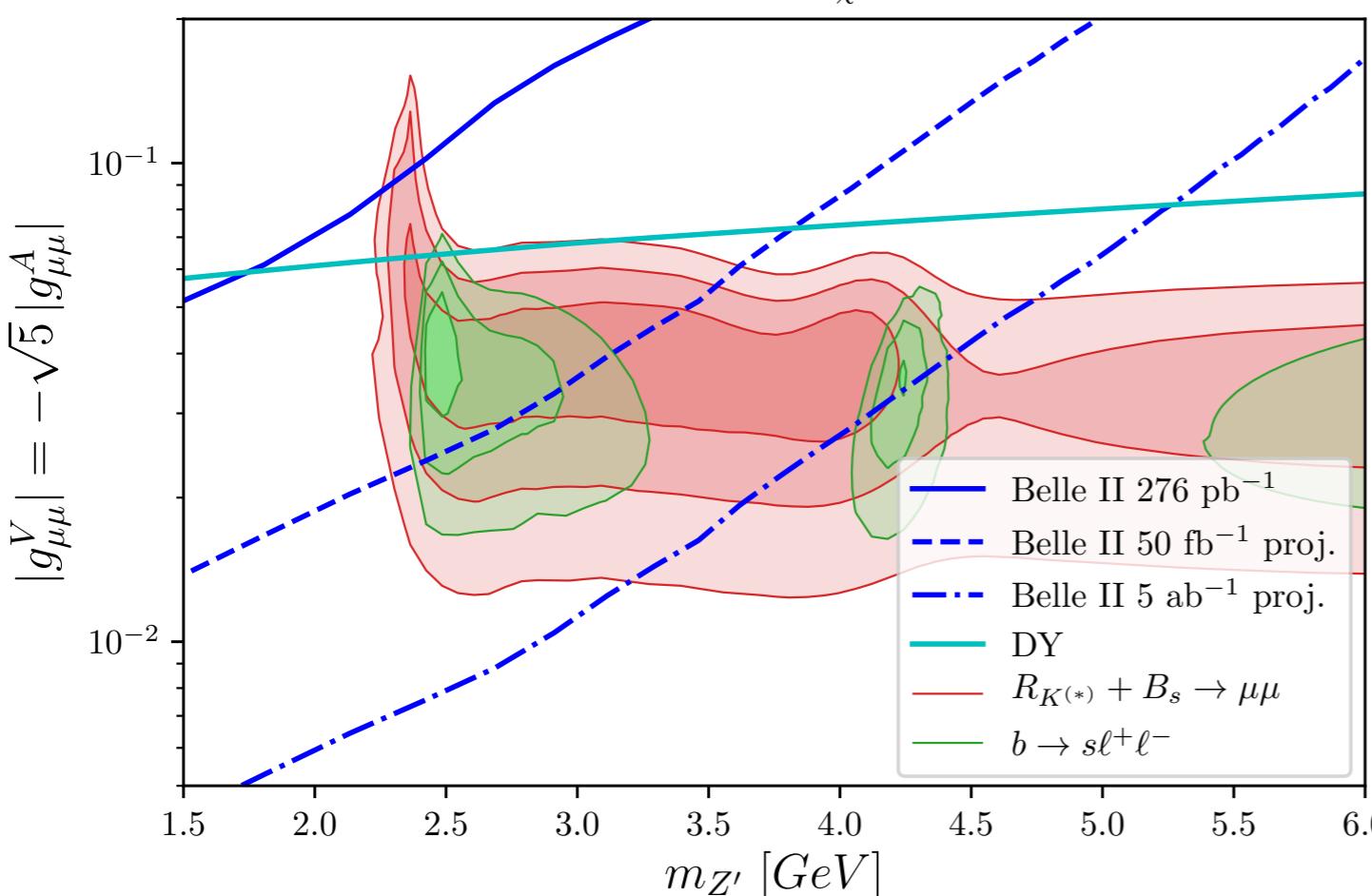
Belle efficiency drops  
⇒  $g_{bs}^L$  can be very large

- Future Belle II analysis can exclude

$$m_{Z'} \sim 2.5 \text{ GeV}$$

- $m_{Z'} \sim 4 \text{ GeV}$  would not lead to

$$R(K^{(*)}) > 1 \text{ in high-}q^2 \text{ bins}$$



# Light $Z'$

## Bounds on $g_{sb}^L$

Alternatively we can maximize  $g_{\mu\mu}^V$  from  $e^+e^- \rightarrow \mu^+\mu^- + \text{invisible}$

