

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut)

#### Explaining the Hints for Lepton Flavour Universality Violation with Three $S_2$ Leptoquark Generations

Luc Schnell Flavor at the Crossroads April 27, 2022



# 1. Motivation and Setup

1.1 Flavour Anomalies
1.2 Single Leptoquark Solutions
1.3 Leptoquarks with Lepton Flavour
1.4 Three Leptoquark Generations
1.5 Lagrangian



#### **1. Motivation and Setup** 1.1 Flavour Anomalies

Hints for Lepton Flavour Universality Violation



Source: cerncourier.com

#### No Hints for Charged Lepton Flavour Violation



### 1. Motivation and Setup **1.2 Single Leptoquark Solutions**







### 1. Motivation and Setup **1.3 Leptoquarks with Lepton Flavour**

- <u>2107.07518 (Greljo, Soreq, Stangl, Thomsen, Zupan)</u> introduced a theoretical framework for **muoquarks**, LQs that only couple to muons.
- This can e.g. be achieved via appropriate  $U(1)_X$  gauge extensions of the SM.

• 
$$X_H = 0$$
,  $X_{Q_i} = X_{U_j} = X_{D_k} \equiv X_q$  for  $i, j$ ,

• This is an economic framework for **joint explanations of the muon anomalies.** But it gives the muon a special treatment and leaves out the other anomalies.

#### $\rightarrow$ Can this be extended to three generations?



#### k = 1.2.3

•  $X_{\ell_2} \neq X_{\ell_{1,2}}$  for  $\ell = L, E \rightarrow$  if LQs have appropriate charges, they couple exclusively to muons.





#### 1. Motivation and Setup **1.4 Three Leptoquark Generations**

- Three LQ generations that couple to one lepton flavour each (tauquark, muoquark, electroquark)
  - $X_H = 0$ ,  $X_{Q_i} = X_{U_i} = X_{D_k} \equiv X_q$  for i, j, k = 1, 2, 3•  $X_{\ell_1}, X_{\ell_2}, X_{\ell_3}$  pairwise different for  $\ell = L, E$
- This is still satisfied by 234 charge assignments for  $U(1)_X$  where  $-10 \le X_{F_i} \le 10$  for all SM fermions (before it was **273**). A possible solution is  $L_{\mu} - L_{\tau}$ .



Sources: 2103.12504 (Angelescu et al.), 2002.12544 (Bigaran et al.)

- Di-quark couplings not possible  $\rightarrow$  no proton decay
- Radiative generation of charged lepton masses





### 1. Motivation and Setup **1.5 Lagrangian**

$$\mathcal{L}_{LQ} = \left(Y_{ij}^{RL}\bar{u}_{i}\left[\Phi_{2}\cdot L_{j}\right] + Y_{ij}^{LR}\left[\bar{Q}_{i}e_{j}\Phi_{2}\right] + H.c.\right) - \left(M^{2} + Y^{H(1)}\left[H^{\dagger}H\right]\right)\Phi_{2}^{\dagger}\Phi_{2}$$
$$-Y^{H(3)}\left[H\cdot\Phi_{2}\right]^{\dagger}\left[H\cdot\Phi_{2}\right] + \mathcal{L}_{4\Phi},$$
Complete SLQ Lagrangian in 2105.04844 (Crivellin, LS)
Three LQ generations with lepton flavours.
$$\mathcal{L}_{LQ} = \sum_{\ell} \left(Y_{\ell}^{RL}\bar{u}_{i}\left[\Phi_{2,\ell}\cdot L_{\ell}\right] + Y_{\ell}^{LR}\left[\bar{Q}_{i}e_{\ell}\Phi_{2,\ell}\right] + H.c.\right)$$

$$\mathcal{L}_{LQ} = \left( Y_{ij}^{RL} \bar{u}_i \left[ \Phi_2 \cdot L_j \right] + Y_{ij}^{LR} \left[ \bar{Q}_i e_j \Phi_2 \right] + \text{H.c.} \right) - \left( M^2 + Y^{H(1)} \left[ H^{\dagger} H \right] \right) \Phi_2^{\dagger} \Phi_2 - Y^{H(3)} \left[ H \cdot \Phi_2 \right]^{\dagger} \left[ H \cdot \Phi_2 \right] + \mathcal{L}_{4\Phi} ,$$

$$\textbf{Three LQ generations with} \qquad \textbf{Complete SLQ Lagrangian} \text{ in } 2105.04844 \text{ (Crivellin, LS)}$$

$$\mathcal{L}_{LQ} = \sum_{\ell} \left( Y_{i\ell}^{RL} \bar{u}_i \left[ \Phi_{2,\ell} \cdot L_{\ell} \right] + Y_{i\ell}^{LR} \left[ \bar{Q}_i e_\ell \Phi_{2,\ell} \right] + \text{H.c.} \right) \\ - \left( M_{\ell}^2 + Y_{\ell}^{H(1)} \left[ H^{\dagger} H \right] \right) \Phi_{2,\ell}^{\dagger} \Phi_{2,\ell} - Y_{\ell}^{H(3)} \left[ H \cdot \Phi_{2,\ell} \right]^{\dagger} \left[ H \cdot \Phi_{2,\ell} \right] + \mathcal{L}_{4\Phi} .$$

Source: 2203.10111 (Crivellin, Fuks, LS)



# 2. Phenomenology

2.1 Tauquark

- 2.2 Tauquark and Muoquark
- 2.3 Muoquark
- 2.4 Electroquark
- 2.5 Higgs Couplings



# 2. Phenomenology 2.1 Tauquark

#### $R_{D^{(*)}}$ Anomaly

- >  $3\sigma$  deviation from SM predictions in  $R_{D^{(*)}}$ . •
- In our model we get a **tree-level** contribution to  $C_{S_L} = +4C_T$ , giving an excellent fit to data provided that a complex phase is present.
- Running with **wilson**, fit with **flavio**.



$$R_{D^{(*)}} = \left. \frac{\operatorname{Br}(B \to D^{(*)} \tau \bar{\nu})}{\operatorname{Br}(B \to D^{(*)} \ell \bar{\nu})} \right|_{\ell \in \{e, \mu\}} ,$$



Source: HFLAV Semileptonic Results 2021

$$(\mathcal{O}_{S_L})_{bc\tau\nu_{\tau}} = -\frac{4G_F}{\sqrt{2}} V_{23}^{\text{CKM}} \left(\bar{c}P_L b\right) \left(\bar{\tau}P_L \nu_{\tau}\right) ,$$
  
$$(\mathcal{O}_T)_{bc\tau\nu_{\tau}} = -\frac{4G_F}{\sqrt{2}} V_{23}^{\text{CKM}} \left(\bar{c}\sigma^{\mu\nu}P_L b\right) \left(\bar{\tau}\sigma_{\mu\nu}P_L \nu_{\tau}\right) ,$$





# 2. Phenomenology 2.1 Tauquark

#### $pp \rightarrow \tau \tau$ Tail

- Tree-level *t*-channel contribution ( $\rightarrow$  energy enhancement).
- Simulation with MadGraph\_aMC@NLO+Pythia8, reconstruction with MadAnalysis5 (b-tag inclusive).





Source: CMS-PAS-HIG-21-001













# 2. Phenomenology





## 2. Phenomenology 2.3 Muoquark

#### Lepton Masses, AMMs and EDMs

- The contributions to these observables are related.
- There is a  $m_t$ -enhanced contribution to  $C_7$ .

$$a_{\ell}^{\mathrm{LQ}} = \frac{G_F m_{\ell}^2}{\sqrt{2}\pi^2} \operatorname{Re}\left\{ \left( \mathcal{C}_7 \right)_{\ell\ell} \right\} \quad \text{and} \quad \left| d_{\ell}^{\mathrm{LQ}} \right| = \frac{eG_F m_{\ell}}{2\sqrt{2}\pi^2} \left| \operatorname{Im}\left\{ \left( \mathcal{C}_7 \right)_{\ell\ell} \right\} \right|,$$

The same diagram (without photon) also induces corrections to the lepton masses.

$$m_{\ell}^{\mathrm{LQ}} \approx -\frac{m_{t}N_{c}}{16\pi^{2}} \mathcal{E}_{3}\left(\frac{\mu^{2}}{M_{\ell}^{2}}, \frac{m_{t}^{2}}{M_{\ell}^{2}}\right) \hat{Y}_{3\ell}^{LR} Y_{3\ell}^{RL*} \quad \text{with} \quad \mathcal{E}_{3}\left(x, y\right) = \frac{1}{\epsilon} + 1 + \log\left(x\right) + y\log\left(y\right) \,,$$

$$(\mathcal{O}_{7})_{\ell\ell} = \frac{4G_{F}}{\sqrt{2}} \frac{e}{16\pi^{2}} m_{\ell} (\bar{\ell} \sigma^{\mu\nu} P_{R} \ell) F_{\mu\nu}.$$

 $\rightarrow$  Could there be a common explanation to the lepton masses,  $(g-2)_{\ell}$  and the absence of EDMs?

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## **2. Phenomenology** 2.3 Muoquark











#### **Radiative Mass Generation**

• Model example discussed in <u>2107.07518 (Greljo, Soreq, Stangl,</u> Thomsen, Zupan): radiative mass generation for charged leptons.

$$\mathcal{L} \supset \left( Y_{i\ell}^{RL} \bar{u}_i \left[ \Phi_{2,\ell}^{RL} \cdot L_\ell \right] + Y_{i\ell}^{LR} \left[ \bar{Q}_i e_\ell \Phi_{2,\ell}^{LR} \right] + \right)$$

$$-M_{\ell,LR}^{2} \Phi_{2,\ell}^{LR\dagger} \Phi_{2,\ell}^{LR} - M_{\ell,RL}^{2} \Phi_{2,\ell}^{RL\dagger} \Phi_{2,\ell}^{RL} - \tilde{M}_{\ell}^{2} \left( \Phi_{2,\ell}^{LR\dagger} \Phi_{2,\ell}^{RL} + \right)$$

$$M_{\ell,LR} = M_{\ell,RL} \equiv M_{\ell} \qquad \left( \begin{array}{c} M_{\ell}^{2} & \tilde{M}_{\ell}^{2} \\ \tilde{M}_{\ell}^{2} & M_{\ell}^{2} \end{array} \right)$$

$$real \rightarrow no EDM$$

$$m_{e} \approx \frac{3m_{t}}{16\pi^{2}} Y_{3e}^{LR} Y_{3e}^{RL} \frac{\tilde{M}_{e}^{2}}{M_{e}^{2}} \qquad a_{e} \approx \frac{3}{16\pi^{2}} Y_{3e}^{LR} Y_{3e}^{RL} \frac{m_{\mu}m_{t}\tilde{M}_{e}^{2}}{M_{e}^{4}} \left( -\frac{5}{3} + \frac{4}{3} \log \frac{M_{e}^{2}}{m_{t}^{2}} \right)$$











 $pp \rightarrow ee$  Tail

- CMS found an excess in dielectron events.
- This was also found by ATLAS, but with less significance.
- The data prefers LQ representations interfering **constructively** with the SM contribution.
- $\sim 3\sigma$  improvement can be reached.







Table 4: The dielectron and dimuon event yields for the data, the expected background and the respective significance in the different SRs used in the analysis. The p-value of each observation is defined as the probability, given the background-only hypothesis, of an observation at least as large as that seen in the data. The significance is the Gaussian cumulative density function of the p-value, and negative significances correspond to deficits.

SR	Data	Background	Significance
$e^+e^-$ Const.	19	12.4±1.9	1.28 - 0.72
$e^+e^-$ Dest.	2	3.1±1.1	
$\mu^+\mu^-$ Const.	6	9.6±2.1	- 0.99
$\mu^+\mu^-$ Dest.	1	1.4±0.9	- 0.58

Source: 2006.12946 (ATLAS)



#### **Parity Violation Data**

- LQs induce parity-violating contributions to electron-nucleon interactions.
- This modifies the weak charge  $Q_w$  of nucleons, as measured by  $Q_{\text{weak}}$  / APV.
  - $Q_{\text{weak}}$ : low-energy electron-proton scattering
  - APV : parity-violating transitions in atoms (e.g. 7S - 6S in <sup>133</sup>Cs)



$$\mathcal{L}_{\text{eff}}^{ee} = \frac{G_F}{\sqrt{2}} \sum_{q=u,d,s} \left( C_{1q}^e \left[ \bar{q} \gamma^\mu q \right] \left[ \bar{e} \gamma_\mu \gamma_5 e \right] + C_{2q}^e \left[ \bar{q} \gamma^\mu \gamma_5 q \right] \left[ \bar{e} \gamma_\mu e^{-2q} \right] \right]$$

$$Q_w = -2 \left[ Z \left( 2\mathcal{C}_{1u}^e + \mathcal{C}_{1d}^e \right) + N \left( \mathcal{C}_{1u}^e + 2\mathcal{C}_{1d}^e \right) \right],$$



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## 2. Phenomenology **2.5 Higgs Couplings**

**Oblique Corrections** 

- We looked into this **before** the CDF II measurement.
- LQ-Higgs terms that break  $SU(2)_L$  spontaneously lead to W-mass corrections.

$$S^{\text{LQ}} \approx -\frac{7N_c v^2}{36\pi} \sum_{\ell} \frac{Y_{\ell}^{H(3)}}{M_{\ell}^2} \qquad T^{\text{LQ}} \approx +\frac{N_c v^2}{96\pi^2 \alpha} \sum_{\ell} \left(\frac{Y_{\ell}^{H(3)}}{M_{\ell}}\right)^2$$



$$\begin{split} S &= -\frac{4s_w^2 c_w^2}{\alpha m_Z^2} \left( \Pi_{ZZ}(0) - \Pi_{ZZ}(m_Z^2) + \Pi_{\gamma\gamma}(m_Z^2) + \frac{c_w^2 - s_w^2}{c_w s_w} \Pi_{Z\gamma}(m_Z^2) \right) \\ T &= \frac{\Pi_{WW}(0)}{\alpha m_W^2} - \frac{\Pi_{ZZ}(0)}{\alpha m_Z^2} \\ U &= -\frac{4s_w^2 c_w^2}{\alpha} \left( \frac{\Pi_{WW}(0) - \Pi_{WW}(m_W^2)}{c_w^2 m_W^2} - \frac{\Pi_{ZZ}(0) - \Pi_{ZZ}(m_Z^2)}{m_Z^2} \right) \\ &+ \frac{s_w^2}{c_w^2} \frac{\Pi_{\gamma\gamma}(m_Z^2)}{m_Z^2} + 2\frac{s_w}{c_w} \frac{\Pi_{Z\gamma}(m_Z^2)}{m_Z^2} \right), \end{split}$$





Source: 2006.10758 (Crivellin et al.)









## 2. Phenomenology **2.5 Higgs Couplings**

• Without CDF II measurement:



Source: 2203.10111 (Crivellin, Fuks, LS)



## 2. Phenomenology **2.5 Higgs Couplings**

• With CDF II measurement:





Source: 2006.10758 (Crivellin et al.)

Contributions from  $\Phi_{2,\ell}$  ( $Y^{H(3)} = \pm 1.0, \pm 2.0, \pm 3.0$ ) Global fit de Blas *et al.* (standard average,  $1\sigma$ ,  $2\sigma$ ) Global fit de Blas *et al.* (conservative average,  $1\sigma, 2\sigma$ )

Global fit Ellis *et al.*  $(1\sigma, 2\sigma)$ 

- Best fit point
- SM point

 $-2\Delta \log \mathcal{L} = -39.0$ 

- $S_2$  still yields a good fit, but very large couplings are needed.
- $S_1 S_3$  can accomplish large contributions to  $\Delta T$  $(\rightarrow$  lepton flavor violation?).

Source: 2204.03996 (Athron et al.), 2204.09031 (Bhaskar et al.)





## 3. Conclusions



# **3.** Conclusions

- The hints for LFUV in multiple lepton generations and the absence of LFV motivate lepton flavoured LQs.
- We examined a corresponding model with three  $S_2$  generations.
- It can provide explanations to deviations in
  - $\ \ \, R_{D^{(*)}} \,, \qquad \ \ \, \triangleright \, pp \to ee \,, \qquad \ \ \, \triangleright \, (a_e) \,.$  $(b \to s \ell^+ \ell^-), \quad (pp \to \tau \tau),$ •  $a_{\mu}$ , • (S, T, U),
- There is an interesting relation between the LQ contributions to the charged lepton masses, AMMs and EDMs.



# Thank you for your attention!