

A novel approach to $\tau \rightarrow \ell + \text{invisible}$

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based on work in collaboration with

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Introduction & Motivation

Motivation

- *“invisible” = an escaping scalar (e.g. ALP)
or vector (e.g. hidden photon)*
- *Such new (light) particles are well motivated*
- *New scalars in the MeV-GeV range w/ larger-than-weak
couplings to 2nd or 3rd generation matter
fully compatible with present data*

see e.g. [Lanfranchi, Pospelov, Schuster, 2021]

- *No compelling reason why these particles should have
flavour-diagonal couplings*

see e.g. [Georgi, Kaplan, Randall, 1986]

Motivation

- *meson or lepton decays at colliders especially suited for such tests*
- *a common hypothesis allowing for minimal model-dependence is that the new particle escapes detection*



$\tau \rightarrow \ell + \text{ALP}$: prototype example of such searches

$\tau \rightarrow \ell + \text{ALP}$: status

- *performed at MARK-III* [Baltrusaitis et al. (MARK-III), 1985]
and ARGUS [Albrecht et al. (ARGUS), 1995]
- *on-going at Belle-II* see e.g. [Tenchini at ICHEP 2020]

Strategy

- *pair-produced τ 's*
- *total E accurately measured*
- *Main bkg's: SM processes w/ undetected particles*



To separate signal from bkg's:

*Estimate the signal-tau momentum
using the visible momenta on the tag side*

"ARGUS method"

**Our approach
to the reconstruction
of the signal- τ momentum**

Rationale

(i) *Signal and bkg's have identical topologies:*

$$\tau (\rightarrow \text{visibles} + \text{invisibles}) + \tau (\rightarrow \text{visibles} + \text{invisibles})$$

Example

$$e^+e^- \rightarrow \boxed{\text{signal}} \tau(\rightarrow \ell\phi) \boxed{\text{tag}} \tau(\rightarrow 3\pi\nu) \quad \text{vs.} \quad e^+e^- \rightarrow \boxed{\text{irreducible bkg.}} \tau(\rightarrow \ell\nu\bar{\nu}) \tau(\rightarrow 3\pi\nu)$$

For such pairwise topologies, there exists an arsenal of kinematic variables, e.g. the “stransverse mass”

(Think of searches of pair-produced SUSY particles)



Let us denote these variables collectively as M_2

M_2 -based strategies

have been widely used in high- p_T searches

We reappraise them

for low-energy pair-prod. leptons or mesons

- M_2 is the 2-decay-chain, Lorentz-invariant generalization of the M_T variable [Smith et al.; Barger et al., 1983]

How to measure the W mass in $W \rightarrow \ell \nu$ (at a hadron collider)

$$m_W^2 \geq m_\ell^2 + m_\nu^2 + 2(E_T^\ell E_T^\nu - \mathbf{p}_T^\ell \mathbf{p}_T^\nu) \equiv M_T^2$$

⇒ the M_T endpoint allows to determine m_W

If one has 2 parents decaying to visibles + invisibles

Would take: $\max\{M_T(\text{branch}_1), M_T(\text{branch}_2)\}$


However, the invisible momenta $\mathbf{k}_{1,2}$, for the 2 branches are not known separately

Only their sum is constrained: $\mathbf{k}_{1T} + \mathbf{k}_{2T} = \mathbf{P}_T^{\text{miss}}$

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 the M_T endpoint allows to determine m_W

If one has 2 parents decaying to visibles + invisibles

$$M_{T2} = \min_{\mathbf{k}_{1T}, \mathbf{k}_{2T}} \left[\max \left\{ M_T(\mathbf{p}_{1T}, \mathbf{k}_{1T}), M_T(\mathbf{p}_{2T}, \mathbf{k}_{2T}) \right\} \right]$$

subject to $\mathbf{k}_{1T} + \mathbf{k}_{2T} = \mathbf{P}_T^{\text{miss}}$

[Lester, Summers, 1999; Barr, Lester, Stephens, 2003]

- M_2 is the 2-decay-chain, Lorentz-invariant generalization of the M_T variable [Smith et al.; Barger et al., 1983]

At Belle II, we don't need the "T"

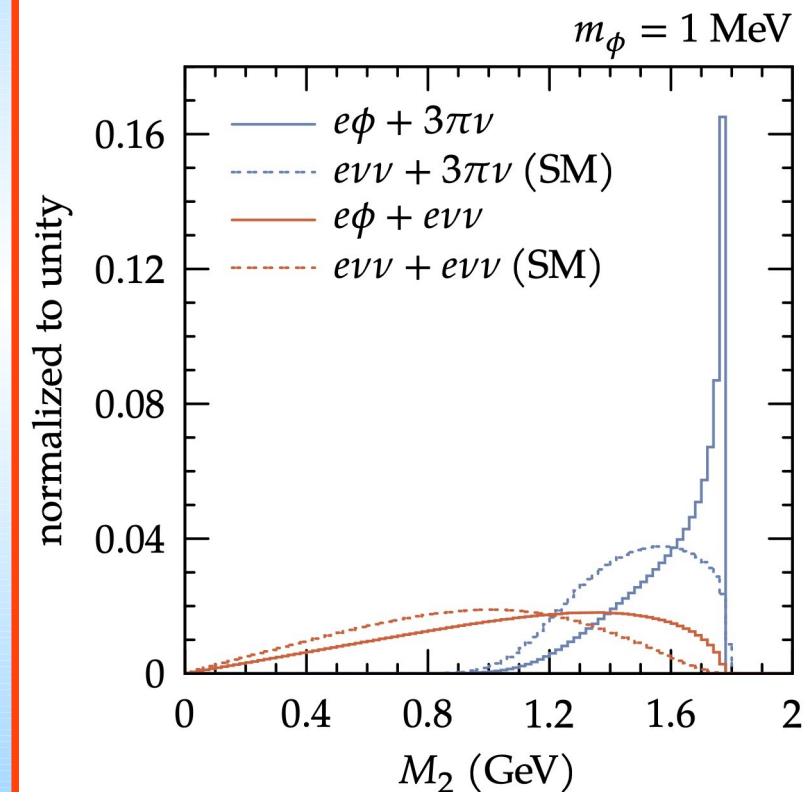
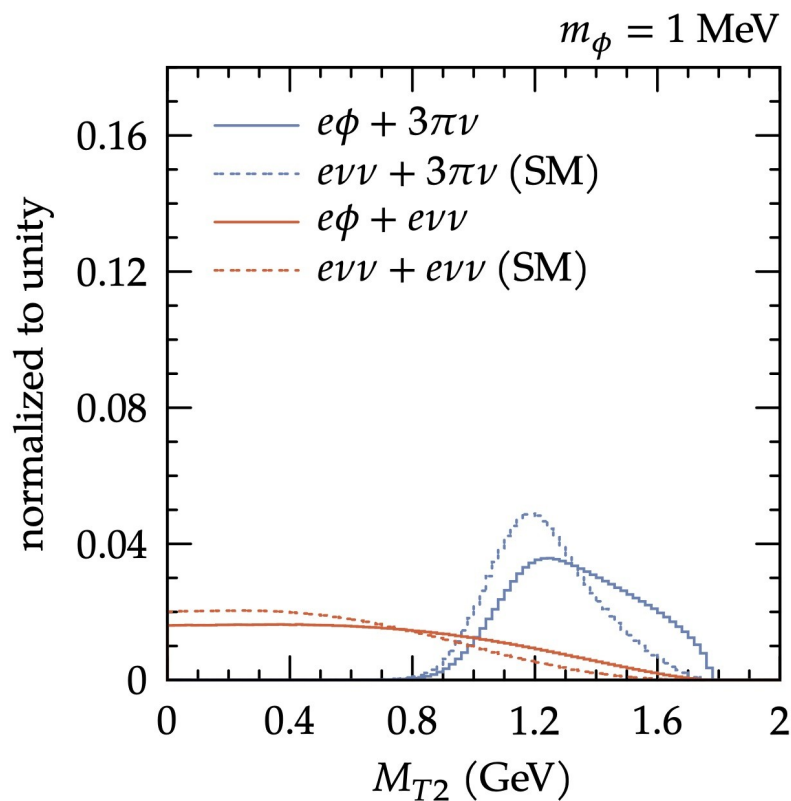
$$M_2 = \min_{\mathbf{k}_1, \mathbf{k}_2} \left[\max \left\{ M(p_1, \mathbf{k}_1), M(p_2, \mathbf{k}_2) \right\} \right]$$
$$\text{subject to } \begin{cases} \mathbf{k}_1 + \mathbf{k}_2 = \mathbf{P}^{\text{miss}}, \\ (p_1 + p_2 + k_1 + k_2)^2 = s \end{cases}$$

[Barr et al., 2011; see also Ross, Serna, 2007; Cho et al., 2014]

M_2 shares most of the features that make M_{T2} very useful



*E.g.: the smaller the number of invisibles,
the more the distrib. is populated towards the upper edge*



This makes these variables “invisible-savvy”

MAOS

$$M_2 = \min_{\mathbf{k}_1, \mathbf{k}_2} \left[\max \left\{ M(p_1, k_1), M(p_2, k_2) \right\} \right]$$

subject to $\begin{cases} \mathbf{k}_1 + \mathbf{k}_2 = \mathbf{P}^{\text{miss}}, \\ (p_1 + p_2 + k_1 + k_2)^2 = s \end{cases}$

*The solution to the minimization can be used as
an estimator of the separate invisible momenta $\mathbf{k}_{1,2}$*



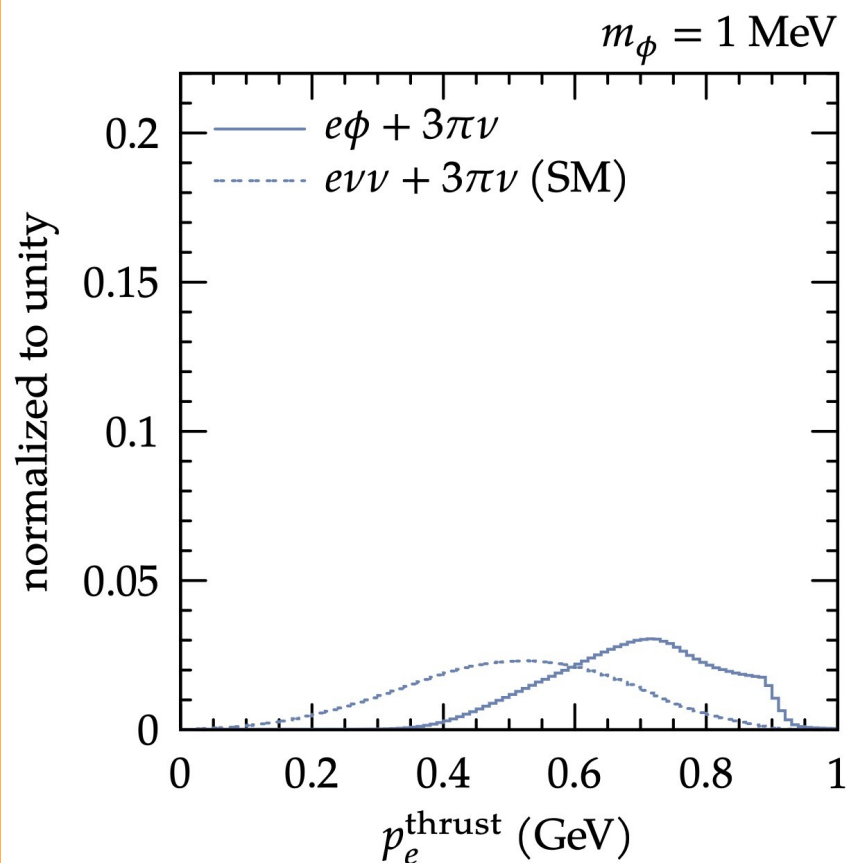
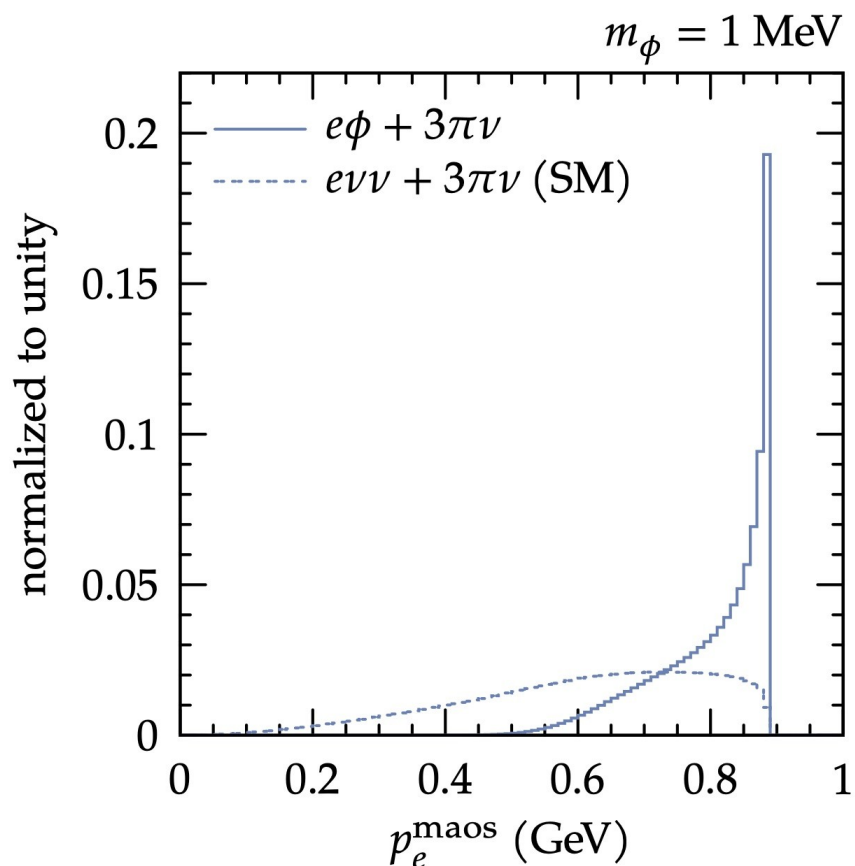
M_2 -Assisted On-Shell (MAOS) invisible momenta

[Cho et al., 2008; Park, 2011]

**With these tools
we then construct
kinematic variables
for S / B discrimination**

Example 1

- With MAOS we can reconstruct $|\mathbf{p}_e|$ in the signal- τ rest frame (RF)
- The same variable can be defined within the “thrust method” (the current state-of-the-art, a generaliz. of the ARGUS method)



Example 2

- Using $\mathbf{k}_{1,2}^{\text{maos}}$ we can construct the ratio

$$\xi_k \equiv \frac{\min\{|\mathbf{k}_1|, |\mathbf{k}_2|\}}{\max\{|\mathbf{k}_1|, |\mathbf{k}_2|\}} \in [0, 1]$$


Underlying rationale

ξ_k distrib. will be populated around 1 for symm. decay chains

This is the case for the $e\nu\bar{\nu} + e\nu\bar{\nu}$ background

Note

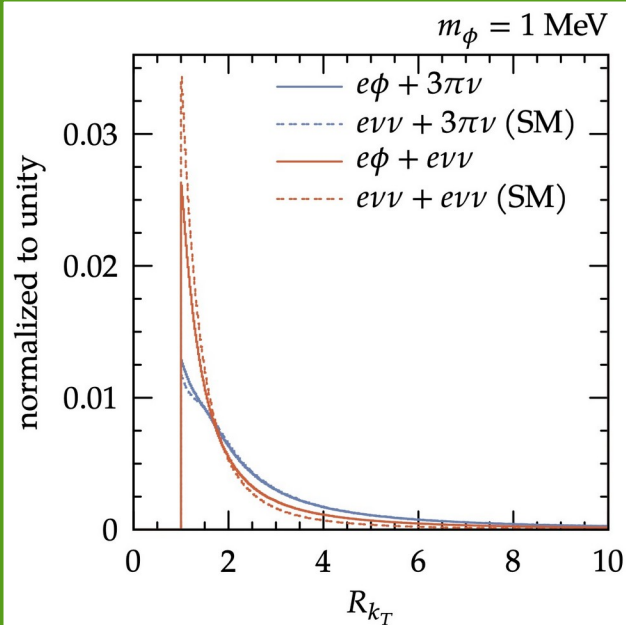
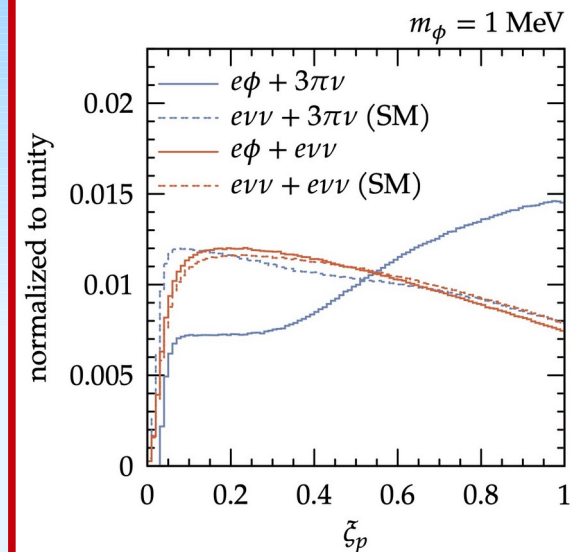
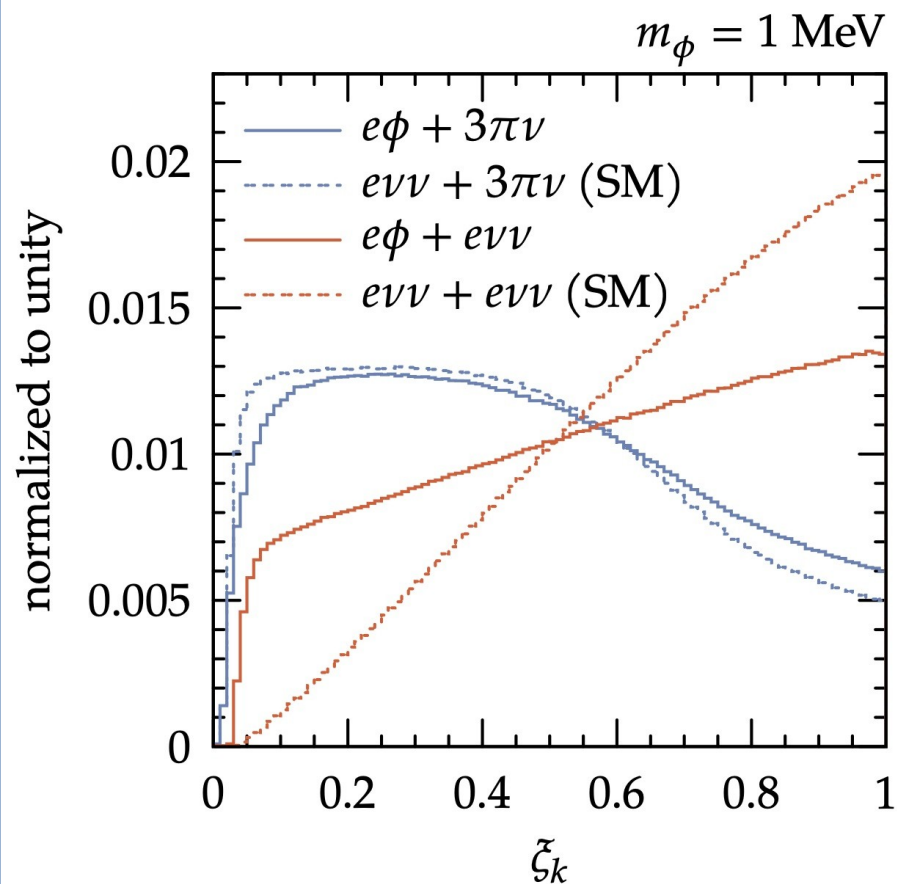
- Earlier literature proposes max / min. [Agashe et al., 2010]
Advantage of min / max : compact domain.
- ξ_k could be defined in the lab or CMS (or yet another) frame.
We use the CMS frame:

 slope differences are magnified

Example 2

- Using $\mathbf{k}_{1,2}^{\text{maos}}$ we can construct the ratio

$$\xi_k \equiv \frac{\min\{|\mathbf{k}_1|, |\mathbf{k}_2|\}}{\max\{|\mathbf{k}_1|, |\mathbf{k}_2|\}} \in [0, 1]$$



Other variables

- We also included variables that do not require MAOS momenta (and show negligible correlation with them)*

M_{recoil} = *invariant mass of the full invisible system*

bkg's have more invisibles than the signal



'typically'

$$M_{\text{recoil}}(\text{bkg}) > M_{\text{recoil}}(\text{sig})$$

$$E_{\text{miss}} = |\mathbf{P}_{\text{miss}}|$$

the more symmetric the 2 decay chains (as in $e\nu\nu e\nu\nu$ bkg)

the more the invisible momenta tend to cancel



the smaller E_{miss}

- *In short, we define a number of kin. variables, whose distrib's are sensitive to the # of invisibles in the decay*
- *These variables can be applied to any search channel with the mentioned decay topology*

***In particular** it can be applied to*

“1 × 3 channel”

signal + tag

$$e\phi + 3\pi\nu$$

irred. bckgr.

$$e\nu\bar{\nu} + 3\pi\nu$$

for which we can compare with the ARGUS method

but also to

“1 × 1 channel”


signal + tag

$$e\phi + e\nu\bar{\nu}$$

irred. bckgr.


$$e\nu\bar{\nu} + e\nu\bar{\nu}$$

for which the ARGUS method is not available



Analysis and Results

Analysis Structure

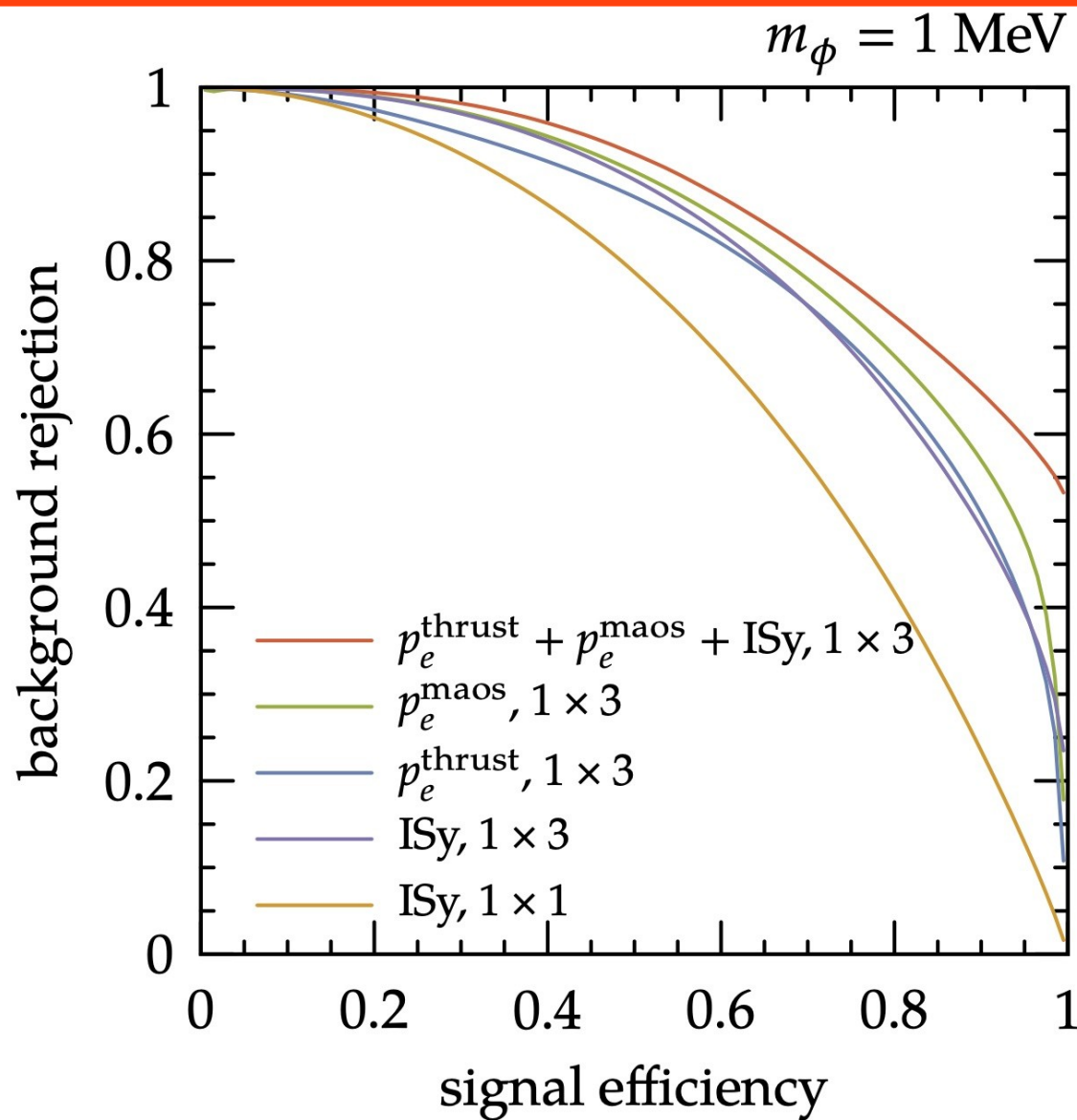
- We collectively denote our kin. variables (M_2 , $\xi_{k,p}$, E_{miss} , M_{recoil}) as 'invisible-savvy'  'ISy' classifier

Note: p_e^{maos} is considered separately (i.e. it's not part of ISy) to allow for a direct comparison with p_e^{thrust}

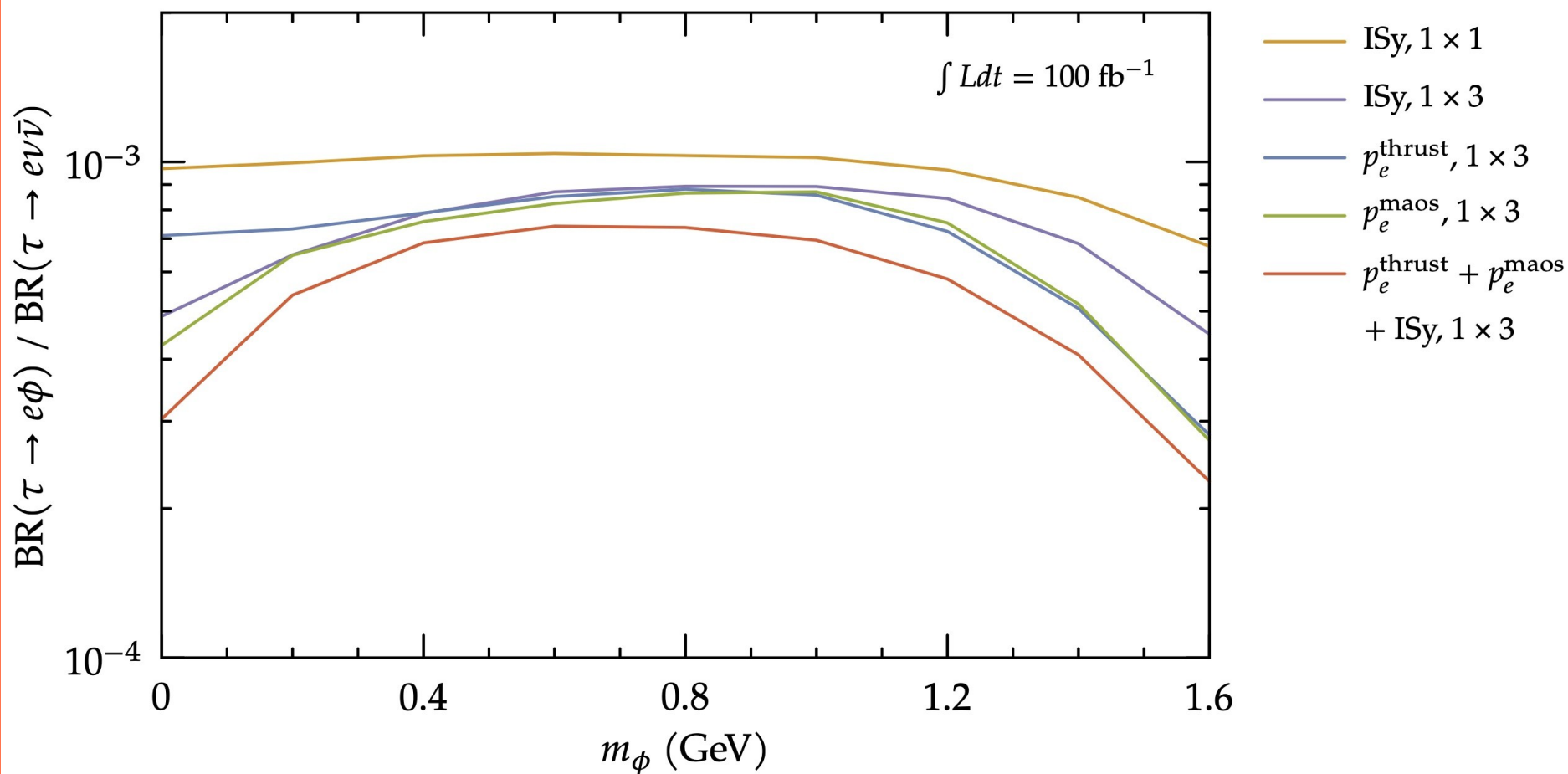
- We then consider the following cases

var's, channel	aim
(a) $p_e^{\text{thrust}}, 1 \times 3$	reproduce the on-going Belle-II analysis, and validate our setup
(b) $p_e^{\text{maos}}, 1 \times 3$	compare MAOS with thrust 'directly' i.e. on a single, common variable
(c) ISy, 1×3	compare ISy with cases (a) and (b)
(d) ISy, 1×1	apply ISy to the 1×1 channel and check its performance compared to 1×3
(abc) $p_e^{\text{thrust}} + p_e^{\text{maos}} + \text{ISy}, 1 \times 3$	

S efficiency vs. *B* rejection

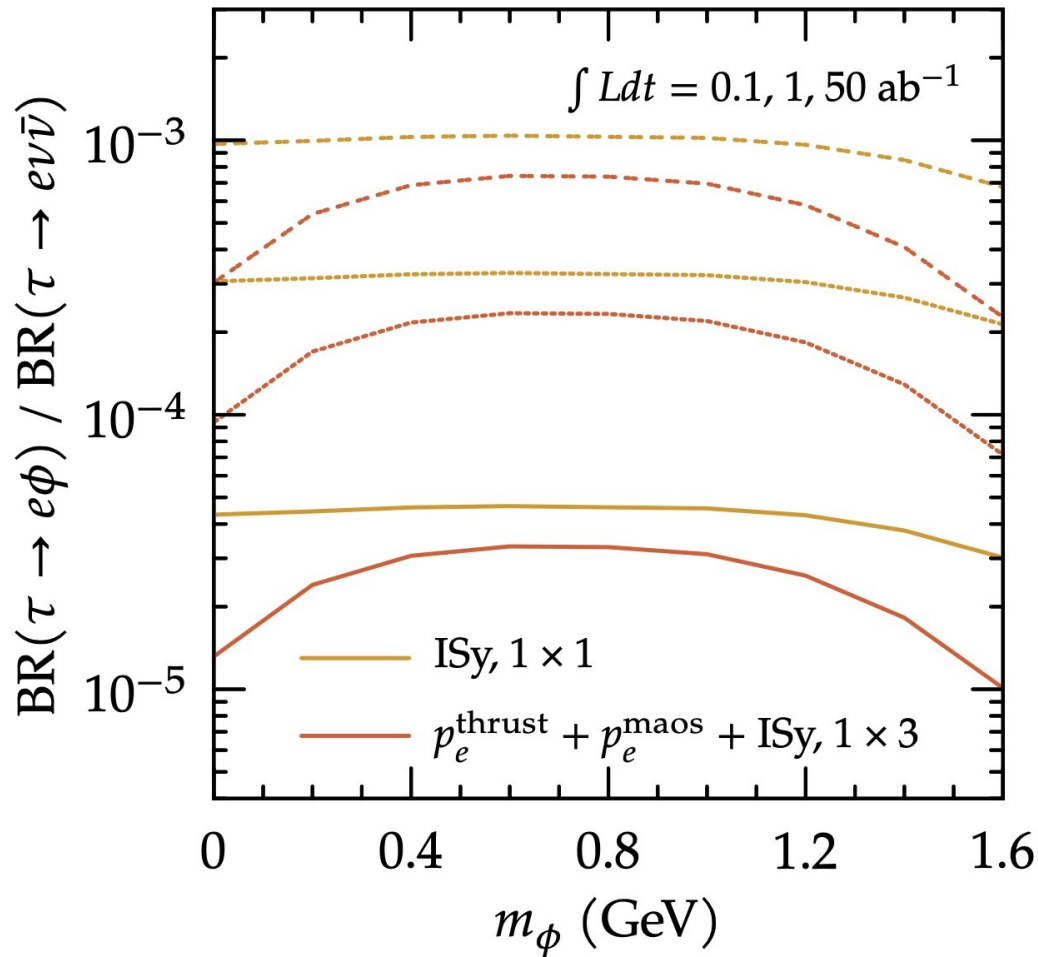


Upper Limit on $BR(\tau \rightarrow e\phi)$



Our method has a clear edge at the two mass endpoints

Upper Limit on $BR(\tau \rightarrow e\phi)$



With 3 benchmark
Belle-II luminosities we get

$$BR(\tau \rightarrow e\phi) \leq$$

$$5.4 \cdot 10^{-5} \quad (\int L dt = 0.1/\text{ab})$$

$$1.7 \cdot 10^{-5} \quad (\int L dt = 1.0/\text{ab})$$

$$2.4 \cdot 10^{-6} \quad (\int L dt = 50/\text{ab})$$

for $m_\phi = 1 \text{ MeV}$


Our method improves by
a factor 3x the limits obtained
with the current strategy

WIP

$B \rightarrow K \tau \mu$

at Belle / Belle II

Motivation

- *B anomalies suggest NP dominantly coupled to the 3rd gen of down-type fermions* [Glashow et al., 2015]
(and w/ hierarchically smaller couplings to the lighter gen's)
 - *Flavour mixing*
 *Dominant (flavoured) effects in $b \rightarrow s$ & final states w/ τ , including LFV ones*
 - *Above obs. made $SU(2)_L$ -compliant in* [Bhattacharya et al., 1412.7164]
thus paving the way for joint explanations of $b \rightarrow s$ and $b \rightarrow c$
see also intro of [Greljo et al., 1506.01705]
 - *Limitations imposed by data on the above picture discussed in*
[Buttazzo et al., 1706.07808]
- A viable avenue: minimally broken $SU(3)^5$* [Barbieri et al., 2011-'12]

$B \rightarrow K \tau \mu$ at Belle

(only correlated charges shown)

signal: $B_{\text{sig}}^+ \rightarrow K_{\text{sig}}^+ \tau \ell_{\text{sig}}$

tag:
(example) $B_{\text{tag}}^- \rightarrow D^0 (\rightarrow K_{\text{tag}}^- \pi^+) \pi^-$ fully reconstructed

- Search usefully reduced to a “bump hunt” in the τ_{sig} decay products

$$\begin{aligned} M_{\text{recoil}}^2 &\equiv (p_{e^+e^-}^* - p_{B_{\text{tag}}}^* - p_{K_{\text{sig}}\ell_{\text{sig}}}^*)^2 \\ &= m_{B_{\text{tag}}}^2 + m_{K_{\text{sig}}\ell_{\text{sig}}}^2 - 2(E_{B_{\text{tag}}}^* E_{K_{\text{sig}}\ell_{\text{sig}}}^* + |\mathbf{p}_{B_{\text{tag}}}^*| |\mathbf{p}_{K_{\text{sig}}\ell_{\text{sig}}}^*| \cos \theta) \end{aligned}$$

- Search could also use semi-lep. tag decays. But then B_{tag} momentum not known, and likewise $\cos \theta$



Enter M_2 ...

Conclusions

- *The M_{T2} variable and its generalizations were conceived for high- p_T events such as pair production of SUSY particles.*

We port these ideas to low-energy processes

- *We devise a novel search strategy, that we apply to pair production of tau leptons, with one decaying as*

$$\tau \rightarrow \ell + \phi \quad \text{with } \phi \text{ a new, elusive particle}$$

- *Our strategy improves the decay-rate sensitivity by 3x compared to the state-of-the-art method.*

(This holds for small m_ϕ , which is the most interesting case.)

- *Our strategy has a vast domain of applicability. To B decays as well!*

How we get from S / B to the BR limit

- Let N_s (N_b) be the # of generated signal (bkg.) events

➡ the bkg. weight is $w_b = B / N_b$

- By def., B is also given by

$$B = \sigma_{\tau\tau} \times \text{BR}(\tau\tau \rightarrow \text{bkg}) \times \int \mathcal{L} dt$$

- The 95% CL stat. significance $\sigma(w_s)$ is given by

$$\sigma(w_s) \simeq S / \sqrt{S + B} = w_s N_s / \sqrt{w_s N_s + w_b N_b} = 1.96$$

➡ we can invert in terms of w_s

- Using $S = w_s N_s$ we finally have

$$\text{BR}(\tau \rightarrow e \phi) / \text{BR}(\tau \rightarrow \text{bkg}) = S / B$$

Note:

The Belle-II analysis determines the BR limit through a template fit instead.

The two procedures are thus independent