### **Theoretical uncertainties count**

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### **Recent experimental results**



PREx: D. Adhikari et al. (PREX Collaboration) Phys. Rev. Lett. 126, 172502 (2021)

Our final results for  $A_{PV}^{meas}$  and  $F_W$  with the acceptance described by  $\epsilon(\theta)$  and  $\langle Q^2 \rangle = 0.00616 \text{ GeV}^2$  are

 $A_{\rm PV}^{\rm meas} = 550 \pm 16 \,({\rm stat}) \pm 8 \,({\rm syst}) \,\,{\rm ppb}$ 

 $F_W(\langle Q^2 \rangle) = 0.368 \pm 0.013 \,(\text{exp}) \pm 0.001 \,(\text{theo}),$ 

**CREx:** publication expected soon. Data below preliminary.

mean scattering angle:	$\overline{ heta}_{\mathrm{Ca}}$	$4.51\pm0.02$	0.5 -
transferred momentum:	$\langle Q^2 \rangle$	$0.0297\pm 0.0002{\rm GeV^2}$	0.0
	q	$0.873\pm0.006{\rm fm}^{-1}$	0.0
beam energy:	$E_{\rm beam}$	$2182.5\pm1.5~{\rm MeV}$	10
weak charge:	$Q_W$	26.073 (26.074?)	_
parity viol. asymmetry:	$A_{ m PV}^{ m (Ca)}$	$2658.6\pm113.2\mathrm{ppb}$	Final value slightly
weak form factor at $Q^2$ :	$F_W^{(\mathrm{Ca})}$	$0.1297 \pm 4.3\%$	different: see KK's talk
			or check arXiv today



#### **Parity Violating Asymmetry:** definition and simple model (PWBA; $F_W \approx F_n$ ; and $F_{ch} \approx F_p$ )

$$A_{p\nu} = \frac{d\sigma_{+}/d\Omega - d\sigma_{-}/d\Omega}{d\sigma_{+}/d\Omega + d\sigma_{-}/d\Omega}$$

$$A_{pv} = \frac{G_F q^2}{4\pi\alpha\sqrt{2}} \Big[ 4\sin^2\theta_W + \frac{F_n(q) - F_p(q)}{F_p(q)} \Big]$$

... which depends on  $F_n(q) - F_p(q)$ . For  $q \to 0$ , it is approximately,

$$\begin{split} -\frac{q^2}{6} \left( \langle r_n^2 \rangle - \langle r_p^2 \rangle \right) &= -\frac{q^2}{6} \left[ \Delta r_{np} (\langle r_n^2 \rangle^{1/2} + \langle r_p^2 \rangle^{1/2}) \right] \\ &= -\frac{q^2}{6} \left( 2 \langle r_p^2 \rangle^{1/2} \Delta r_{np} + \Delta r_{np}^2 \right) \end{split}$$

variation of  $A_{pv}$  at a fixed q dominated by the variation of  $\Delta r_{np}$ .  $F_p(q)$  well fixed by experiment

### **Dipole Polarizability: definition**

The **linear response** or dynamic polarizability of a **nuclear system excited** from its g.s.,  $|0\rangle$ , to an excited state,  $|v\rangle$ , due to the **action of an external isovector oscillating field** (dipolar in our case) of the form (Fe<sup>iwt</sup> + F<sup>†</sup>e<sup>-iwt</sup>):

$$F_{JM} = \sum_{i}^{A} r^{J} Y_{JM}(\hat{r}) \tau_{z}(i) \ (\Delta L = 1 \rightarrow \text{Dipole})$$

is proportional to the **static polarizability** for small oscillations

$$\alpha = (8\pi/9)e^2 m_{-1} = (8\pi/9)e^2 \sum_{\nu} |\langle \nu | F | 0 \rangle|^2 / E \text{ where } m_{-1} \text{ is}$$
  
the inverse energy weighted moment of the strength function, defined as,  $S(E) = \sum |\langle \nu | F | 0 \rangle|^2 \delta(E - E_{\nu})$ 

## **Dipole Polarizability: simple model**

electric polarizability measures tendency of the nuclear charge distribution to be distorted ( $\alpha \sim \frac{\text{electric dipole moment}}{\text{external electric field}}$ )

 The dielectric theorem establishes that the m<sub>-1</sub> moment can be computed from the expectation value of the Hamiltonian in the constrained ground state H' = H + λD.

Adopting the Droplet Model:

$$m_{-1} \approx \frac{A \langle r^2 \rangle^{1/2}}{48J} \left( 1 + \frac{15}{4} \frac{J}{Q} A^{-1/3} \right)$$

within the same model, connection with the neutron skin thickness:

$$\alpha_{\rm D} \approx \frac{A \langle r^2 \rangle}{12J} \left[ 1 + \frac{5}{2} \frac{\Delta r_{\rm np} + \sqrt{\frac{3}{5}} \frac{e^2 Z}{70J} - \Delta r_{\rm np}^{\rm surface}}{\langle r^2 \rangle^{1/2} (I - I_{\rm C})} \right]$$

# Analysis based on EDFs correlations between A<sub>pv</sub>, ΔR<sub>ch</sub> and L



**EDFs** callibrated to **reproduce** the **binding energy** and **charge radii** of some **selected nuclei**. In some cases pseudo-observables related to the EoS are also used



D. Adhikari et al. Phys. Rev. Lett. 126, 172502 (2021)

# Analysis based on EDFs correlations between α<sub>D</sub>, ΔR<sub>ch</sub> and L



X. Roca-Maza, et al. Phys. Rev. C 88, 024316 (2013)

X. Roca-Maza, et al. Phys. Rev. C 92, 064304 (2015)

# New analysis taking into account model error quantification

#### How to improve current analysis?

- Include theoretic statistical errors and correlations within a given EDF parametrization
- Fitting procedure including experimental data not only on **B** and  $R_{ch}$  but also on  $A_{Pv}$  and/or  $\alpha_{D}$  ("informed" EDFs)
- Extension of available EDFs to account for missing systematic uncertainties and more flexibility
- Other issues related to theory?

. . .

## PREx: theo. uncertainty budget for Apv

# Uncertainties in the determination of the Form Factors is smaller than typical EDFs statistical uncertainties



Thin blue bars: statistical model uncertainties (related to neutron and proton densities)

Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz Phys. Rev. Lett. 127, 232501 (2021)

# EDF predictions with error ellipsoids for $A_{PV}$ and $\alpha_{D}$ in <sup>208</sup>Pb

**Correlation ellipsoids** within each **EDF** show **similar correlations** than the **systematic** study with many **EDFs** ↔ **have we learnt something**?



# How these models perform for *B* and *R*<sub>ch</sub> in <sup>208</sup>Pb?



→ The **residuals** of the charge **radius** (a) and binding **energy** (b) of  $^{208}$ Pb for the theoretical models.

→ The grey bands around the perfect match indicate the typical performance of EDFs (i.e. typical r.m.s. deviation taken over all nuclei where correlation effects are small)

Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz Phys. Rev. Lett. 127, 232501 (2021)

## A<sub>Pv</sub> versus α<sub>D</sub> in well callibarted EDFs



### Fitting A<sub>PV</sub> and α<sub>D</sub>



Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz (to be submitted 2022)

# Are EDFs incompatible with experimental data?



*Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz (to be submitted 2022)* 

#### **Does a richer EDF improve?**



Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz (to be submitted 2022)



# **Other example: minimal model** assumptions and statistical analysis



Reed Essick, et al. Phys. Rev. Lett. 127, 192701 (2021)

Astro+PREx-II

Non-parametric equation of state representation derived from observations of neutron stars with minimal modeling assumptions.

The resulting **astrophysical** constraints from heavy pulsar masses, LIGO/Virgo, and NICER clearly favor "small" values of the **neutron skin** and L.

**Combining astrophysical** data with **PREX-II** and chiral effective field theory constraints yields

> J=33.0 ± 2.0 MeV I = 53 + 15 MeV $R_{skin} = 0.17 \pm 0.04 \text{ fm}$

# Conclusions

- Current EDFs show strong systematic and statistic correlations between  $A_{PV}$  and  $\Delta R_{np}$  or  $\alpha_{D}J$  and  $\Delta R_{np}$
- Fitting masses and radii do not give enough information on  $A_{PV}$ ,  $\alpha_{PJ}$ and  $\Delta R_{np}$  (additional reason for the strong correlation)
- Extending the fitting protocol to include A<sub>PV</sub> and α<sub>D</sub> do not change the correlations above since experimental errors on these observables are still large → model predictions remain biased by masses and radii
- More accurate measurements on A<sub>PV</sub> could point to model deficiencies, while EDFs seem to accommodate better α<sub>P</sub>
- **Current EDFs** are able to overlap within  $\sim 1\sigma$  all experimental **data** (except A<sub>PV</sub> in <sup>48</sup>Ca where SV-ext gives the best result being away  $\sim 1.5\sigma$ from experiment)
- Ideally: measure A<sub>PV</sub> at different kinematics or A<sub>PV</sub> on different nuclei with same (or better) accuracy → better constraints to models

### Collaborators

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