

# Normal Spin Asymmetry in Elastic Electron Scattering

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***Washington, DC, USA***

**MITP Workshop *Precision Tests with Neutral-Current Coherent  
Interactions with Nuclei***

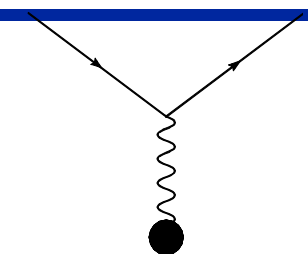
***MITP, Mainz, 24 May 2022***

# Plan of talk

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- **Single-spin asymmetries**
  - Introduction
  - **Optical theorem approach in ep-scattering via two-photon exchange**
  - **Novel features of a single-spin asymmetry**
  - **Comparison with experiment on a nucleon and nuclei**
- **Summary and outlook**

# Elastic Nucleon Form Factors



- Based on one-photon exchange approximation

$$M_{fi} = M_{fi}^{1\gamma}$$

$$M_{fi}^{1\gamma} = e^2 \bar{u}_e \gamma_\mu u_e \bar{u}_p (F_1(t) \gamma_\mu - \frac{\sigma_{\mu\nu} q_\nu}{2m} F_2(t)) u_p$$

- Two techniques to measure

$$\sigma = \sigma_0 (G_M^2 \tau + \varepsilon \cdot G_E^2) \quad : \text{Rosenbluth technique}$$

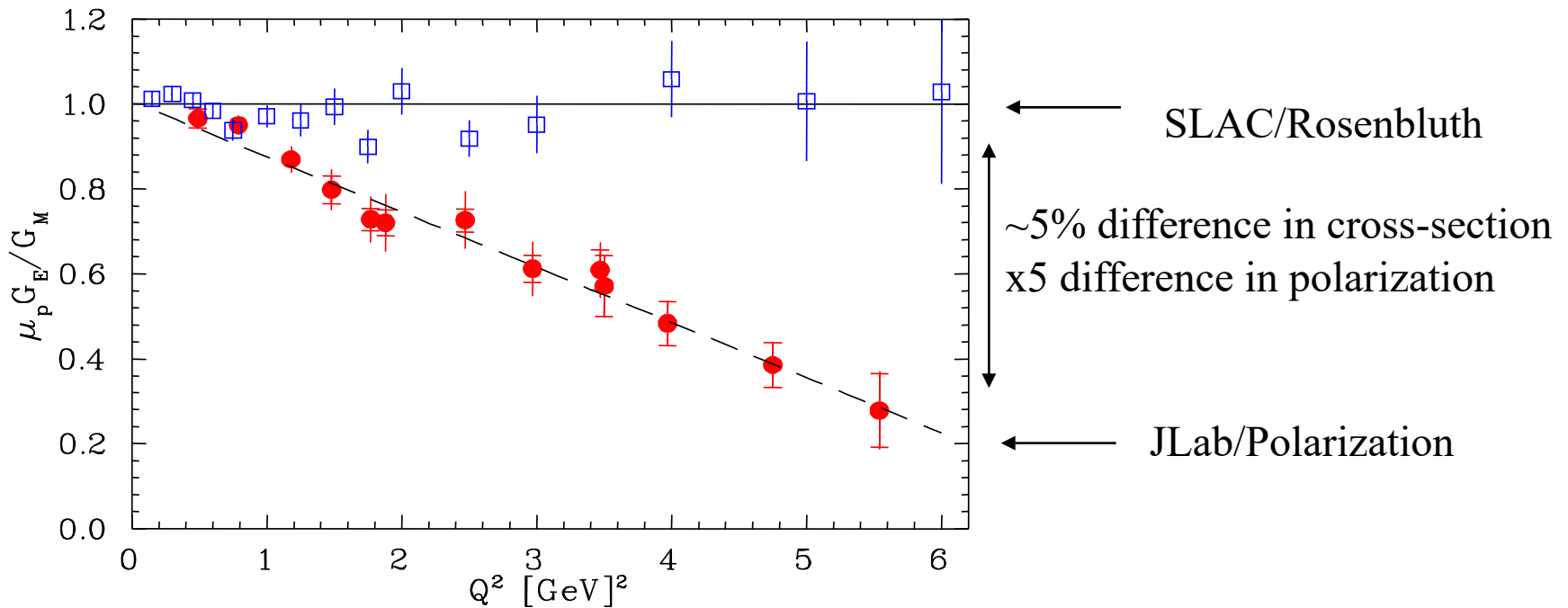
$$\frac{P_x}{P_z} = -\frac{A_x}{A_z} = -\frac{G_E \sqrt{\tau} \sqrt{2\varepsilon(1-\varepsilon)}}{G_M \tau \sqrt{1-\varepsilon^2}} \quad : \text{Polarization technique}$$

$$G_E = F_1 - \tau F_2, \quad G_M = F_1 + F_2$$

$$(P_y = 0)$$

Latter due to: Akhiezer, Rekalov; Arnold, Carlson, Gross

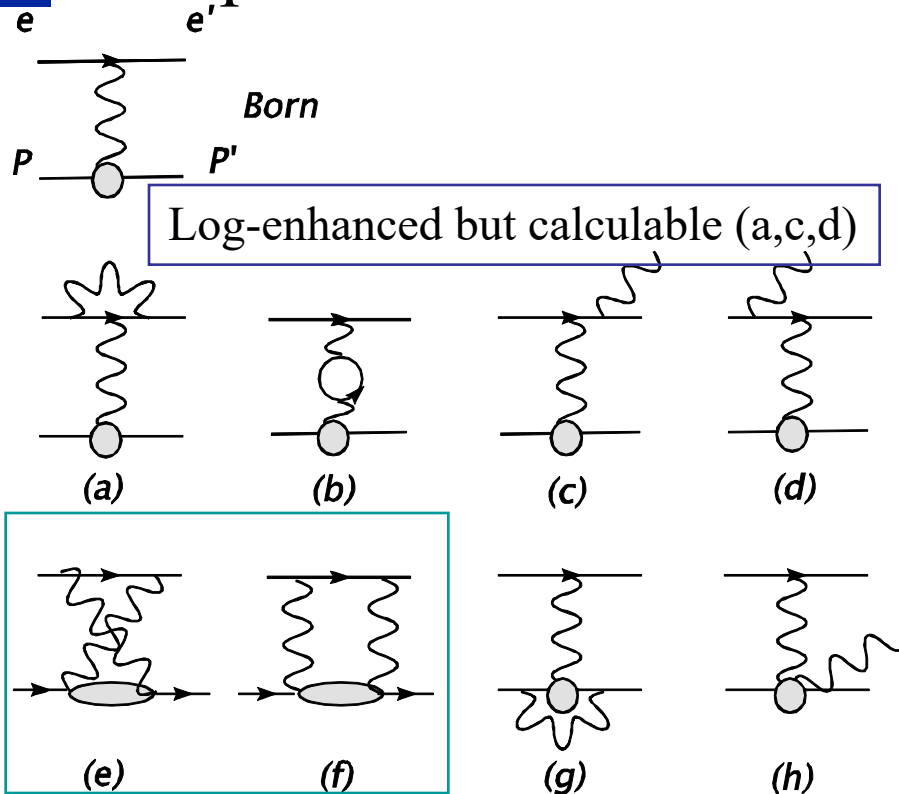
# Ge/Gm Puzzle



- Both early SLAC and JLab experiments on (super)Rosenbluth separations followed  $G_E/G_M \sim \text{const}$ , see I.A. Quattan et al., Phys.Rev.Lett. 94:142301,2005
- JLab measurements using polarization transfer technique give different results (Jones'00, Gayou'02)

*Radiative corrections, in particular, a short-range part of 2-photon exchange is a likely origin of the discrepancy*

# Complete radiative correction in $O(\alpha_{em})$



## Radiative Corrections:

- Electron vertex correction (a)
- Vacuum polarization (b)
- Electron bremsstrahlung (c,d)
- Two-photon exchange (e,f)
- Proton vertex and VCS (g,h)
- Corrections (e-h) depend on the nucleon structure
- Meister&Yennie; Mo&Tsai
- Further work by Bardin&Shumeiko; Maximon&Tjon; AA, Akushevich, Merenkov;
- Guichon&Vanderhaeghen'03:  
*Can (e-f) account for the Rosenbluth vs. polarization experimental discrepancy? Look for ~3% ...*

## Main issue: Corrections dependent on nucleon structure

Model calculations:

- Blunden, Melnitchouk, Tjon, Phys.Rev.Lett.**91**:142304,2003
- Chen, AA, Brodsky, Carlson, Vanderhaeghen, Phys.Rev.Lett.**93**:122301,2004
- For review, see AA, Blunden, Hasell, Raue, Prog.Part.Nucl.Phys. 95, 245 (2017)

# Single-Spin Asymmetries in Elastic Scattering

## Parity-conserving

- Observed spin-momentum correlation of the type:

$$\vec{s} \cdot \vec{k}_1 \times \vec{k}_2$$

where  $k_{1,2}$  are initial and final electron momenta,  $s$  is a polarization vector of a target OR beam

- For elastic scattering asymmetries are due to *absorptive part* of 2-photon exchange amplitude

## Parity-Violating

$$\vec{s} \cdot \vec{k}_1$$

# Beam Single-Spin Asymmetry: Early Calculations

- *Spin-orbit interaction of electron scattering off a Coulomb field*

N.F. Mott, Proc. Roy. Soc. London, Set. A 124, 425 (1929); *ibid.* 135, 429 (1932);

- *Interference of one-photon and two-photon exchange Feynman diagrams in electron-muon scattering*: Barut, Fronsdal, Phys.Rev.120, 1871 (1960)

- *Extended to quark-quark scattering SSA in pQCD*: Kane, Pumplin, Repko, Phys.Rev.Lett. 41, 1689 (1978)



Sir Nevill Mott  
Nobel Prize (1977)

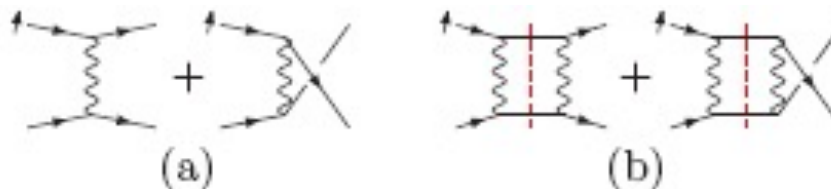
$$\Delta(\vartheta) = \mp 2Z\alpha \frac{v\sqrt{1-v^2}}{1-v^2} \frac{\sin^3(\vartheta/2)}{\sin^2(\vartheta/2)} \frac{1}{\cos(\vartheta/2)} \ln \frac{1}{\sin(\vartheta/2)}.$$

$$A_n \propto \frac{\alpha \cdot m_e \cdot \theta^3}{E}, \text{ for } \theta \ll 1$$

(small – angle scattering)

# Normal Beam Asymmetry in Moller Scattering

- Pure QED process,  $e^-+e^- \rightarrow e^-+e^-$ 
  - Barut, Fronsdal , Phys.Rev.120:1871 (1960): Calculated the asymmetry in first non-vanishing order in QED  $O(\alpha)$
  - Dixon, Schreiber, Phys.Rev.D69:113001,2004, Erratum-ibid.D71:059903,2005: Calculated  $O(\alpha)$  correction to the asymmetry



$$A_n \propto \frac{2M_\gamma \text{Im}(M_{2\gamma})}{M_\gamma^2} \xrightarrow{\sqrt{s} \gg m_e} \alpha \frac{m_e}{\sqrt{s}} f(\theta)$$

SLAC E158 Results (K. Kumar, private communication):

$A_n(\text{exp}) = 7.04 \pm 0.25(\text{stat})$  ppm

$A_n(\text{theory}) = 6.91 \pm 0.04$  ppm



# Single-Spin Target Asymmetry $\vec{s}_T \cdot \vec{k}_1 \times \vec{k}_2$

De Rujula, Kaplan, De Rafael, Nucl.Phys. B53, 545 (1973):

Transverse polarization effect is due to the *absorptive part of the non-forward Compton amplitude for off-shell photons* scattering from nucleons

See also AA, Akushevich, Merenkov, hep-ph/0208260

$$A_{l,p}^{el,in} = \frac{8\alpha}{\pi^2} \frac{Q^2}{D(Q^2)} \int dW^2 \frac{S + M^2 - W^2}{S + M^2} \frac{dQ_1^2}{Q_1^2} \frac{dQ_2^2}{Q_2^2} \frac{1}{\sqrt{K}} B_{l,p}^{el,in}$$

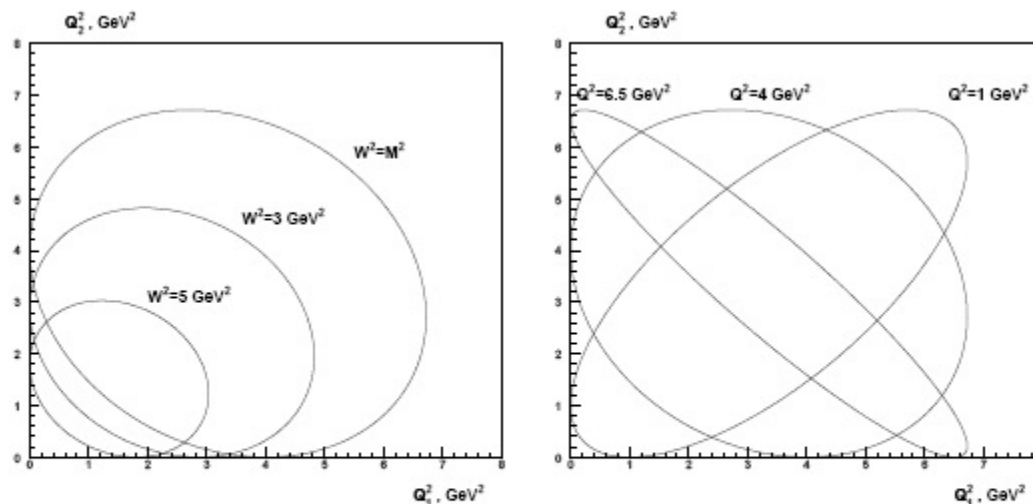
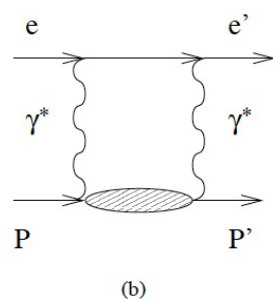
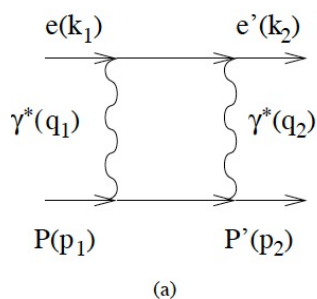


Figure 2. Integration region over  $Q_1^2$  and  $Q_2^2$  in Eq.(2) for elastic ( $W^2 = M^2$ ) and inelastic contributions. The latter (left) is given for  $Q^2=4 \text{ GeV}^2$  and two values of  $W^2$ , which is an integration variable in this case. The elastic case is shown on the right as a function of external  $Q^2$ . The electron beam energy is  $E_b = 5 \text{ GeV}$ .

# Sherman function on nuclei

- Theoretical approach by Sherman Phys. Rev. 103, 1601 (1956)
- Mott polarimetry is based on comparing measured analyzing power with theoretical asymmetry: Gay, Dunning, Rev. Sci. Instrum. 63, 1635 (1992); Price, Poelker, Sinclair et al., In: Proc. Protvino 1998, High Energy Spin Physics Symposium, p.554; Tioukine, Aulenbacher and Riehn, Rev. Sci. Instrum. 82, 033303 (2011)

Extended to high energies (100-1000MeV) in

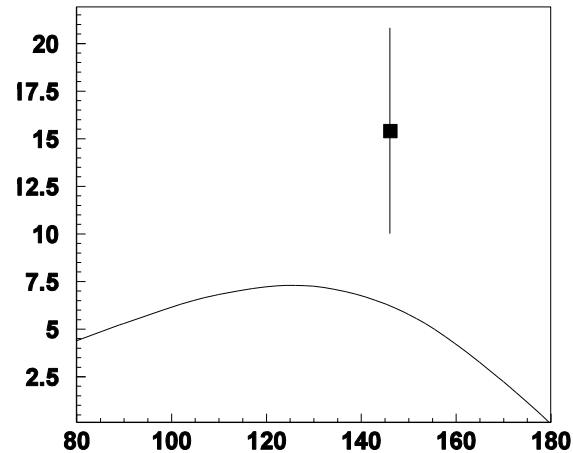
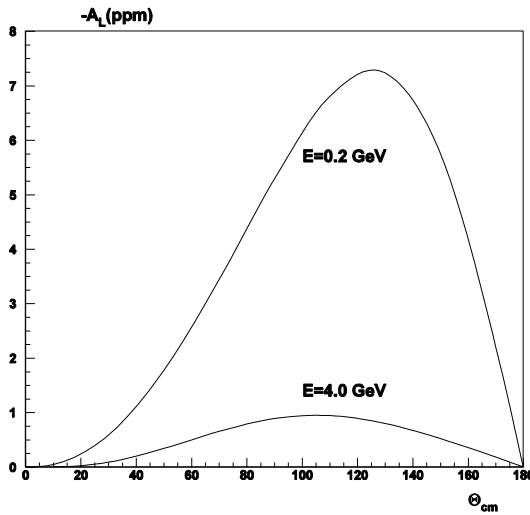
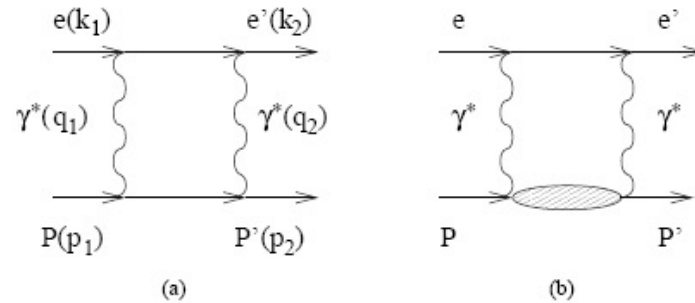
- Cooper and Horowitz, PRC 72, 034602 (2005)
- Jakubassa-Amundsen and Barday, 2012 J. Phys. G: Nucl. Part. Phys. 39 025102

The approaches to calculating Sherman functions on nuclei involve solving Dirac equation in Coulomb field (see Jakubassa-Amundsen's talk)

# Proton Mott Asymmetry at Higher Energies

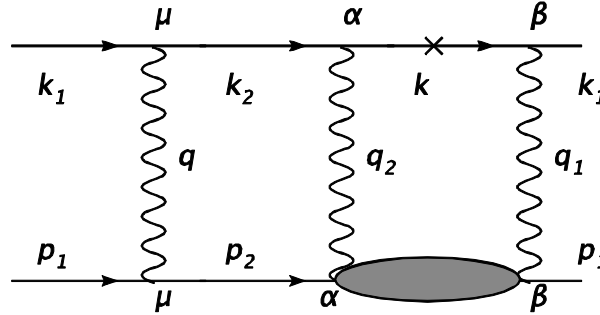
AA, Akushevich, Merenkov,  
 hep-ph/0208260

Transverse beam SSA,  
 units are parts per million



- Asymmetry due to absorptive part of two-photon exchange amplitude; shown is elastic intermediate state contribution
- Nonzero effect first observed by SAMPLE Collaboration (S.Wells et al., PRC63:064001,2001) for 200 MeV electrons

# Beam Normal Asymmetry from Inelastic Intermediate States (hep-ph/0407167)



$$A_n^{e,P} = -\frac{\alpha Q^2}{\pi^2 D(s, Q^2)} \text{Im} \int \frac{d^3 k}{2k_0} \cdot \frac{L_{\mu\alpha\beta} H_{\mu\alpha\beta}}{Q_1^2 Q_2^2}$$

$$L_{\mu\alpha\beta} = \frac{1}{4} \text{Tr}(\hat{k}_2 + m_e) \gamma_\mu (\hat{k}_1 + m_e) (1 - \gamma_5 \hat{\xi}_e) \gamma_\beta (\hat{k} + m_e) \gamma_\alpha$$

$$H_{\mu\alpha\beta} = \frac{1}{4} \text{Tr}(\hat{p}_2 + M) \Gamma_\mu (\hat{p}_1 + M) (1 - \gamma_5 \hat{\xi}_p) T_{\beta\alpha}$$

$$\hat{a} \equiv a_\mu \gamma_\mu$$

$$L_{\mu\alpha\beta} q_\mu = L_{\mu\alpha\beta} q_{2\alpha} = L_{\mu\alpha\beta} q_{1\beta} = H_{\mu\alpha\beta} q_\mu = H_{\mu\alpha\beta} q_{2\alpha} = H_{\mu\alpha\beta} q_{1\beta} = 0$$

Gauge invariance essential in cancellation of infra-red singularity for target asymmetry

$$L_{\mu\alpha\beta} H_{\mu\alpha\beta} \rightarrow 0 \quad \text{if} \quad Q_1^2 \text{ and/or } Q_2^2 \rightarrow 0$$

Novel feature of the normal beam asymmetry: After  $m_e$  is factored out, the remaining expression is singular when virtuality of photons reach zero in the loop integral! The expressions are regular for the target SSA, since the photon's virtualities are at hadronic mass scale

$$L_{\mu\alpha\beta} H_{\mu\alpha\beta} \rightarrow m_e \cdot \text{const} \quad \text{if} \quad Q_1^2 \text{ and/or } Q_2^2 \rightarrow 0 \Rightarrow A \sim m_e \log^2 \frac{Q^2}{m_e^2}, m_e \log \frac{Q^2}{m_e^2}$$

Also calculations by Vanderhaeghen, Pasquini (2004); Gorchtein, hep-ph/0505022;

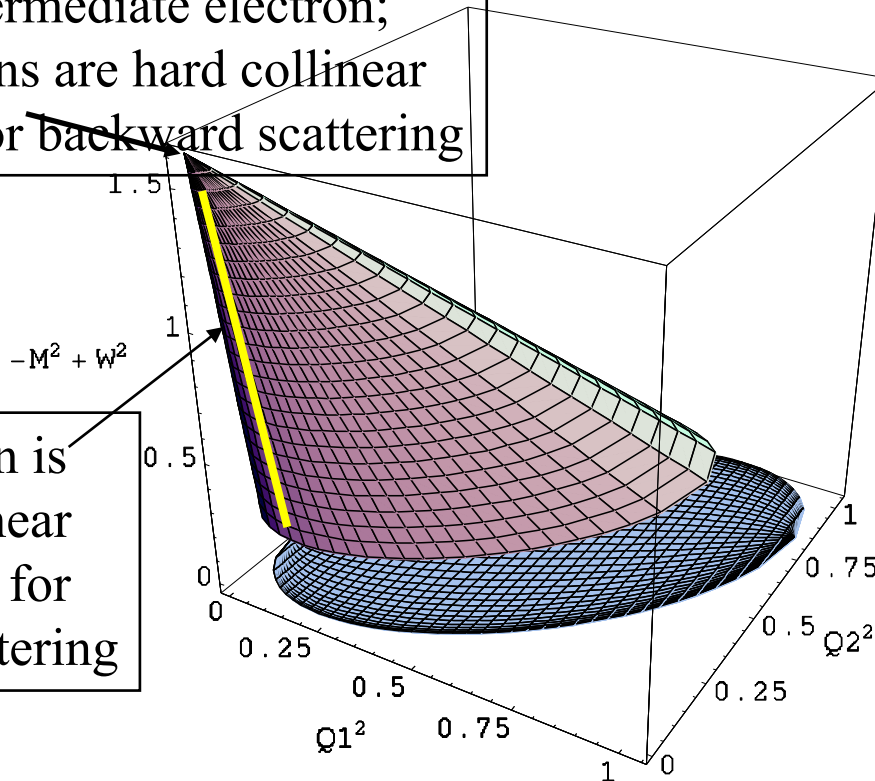
Borisyuk, Kobushkin, Phys. Rev. C 73 (2006) 045210; confirm **quasi-real photon exchange**

# Phase Space Contributing to the absorptive part of $2\gamma$ -exchange amplitude

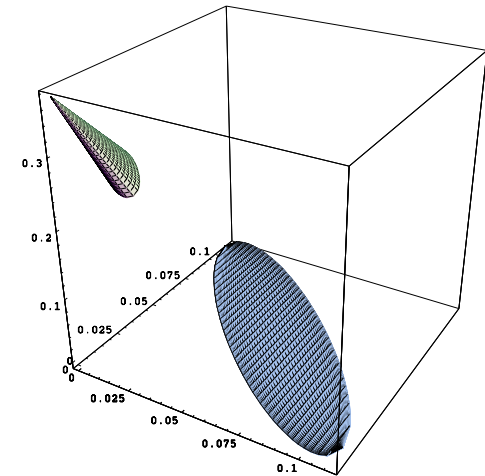
- 2-dimensional integration ( $Q_1^2, Q_2^2$ ) for the elastic intermediate state
- 3-dimensional integration ( $Q_1^2, Q_2^2, W^2$ ) for inelastic excitations

'Soft' intermediate electron;  
Both photons are hard collinear  
Dominates for backward scattering

One photon is hard collinear  
Dominates for forward scattering

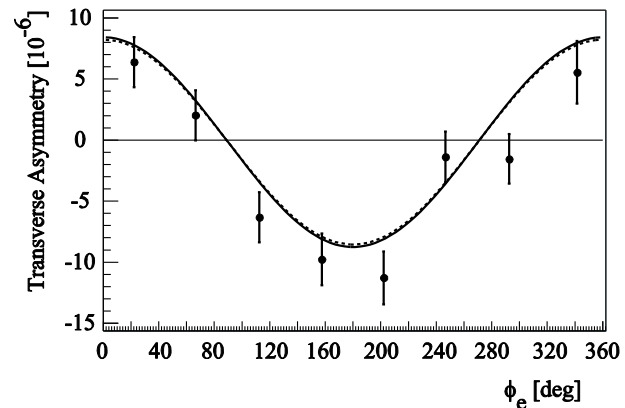


Examples: MAMI A4  
E= 855 MeV  
 $\Theta_{cm}= 57$  deg;  
SAMPLE, E=200 MeV

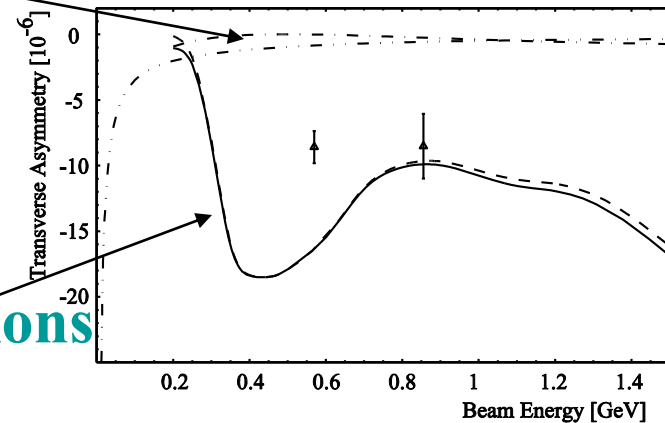


# MAMI data on Mott Asymmetry

- F. Maas et al., [MAMI A4 Collab.]  
Phys.Rev.Lett.94:082001, 2005
- Pasquini, Vanderhaeghen:  
Phys.Rev.C70:045206,2004  
Surprising result: Dominance of  
inelastic intermediate excitations



Elastic intermediate  
state



Inelastic excitations  
dominate

# Special property of Mott asymmetry

- Mott asymmetry above the nucleon resonance region
  - (a) does not decrease with beam energy
  - (b) is enhanced by large logs
- (AA, Merenkov, PL B599 (2004)48; hep-ph/0407167v2 (erratum) )
- Reason for the unexpected behavior: exchange of hard collinear quasi-real photons and diffractive mechanism of nucleon Compton scattering
  - For  $s \gg -t$  and above the resonance region, the asymmetry is given by:

$$A_n^e(\text{diffractive}) = \sigma_p \frac{(-m_e)\sqrt{Q^2}}{8\pi^2} \cdot \frac{F_1 - \tau F_2}{F_1^2 + \tau F_2^2} (\log(\frac{Q^2}{m_e^2}) - 2) \cdot \text{Exp}(-bQ^2)$$

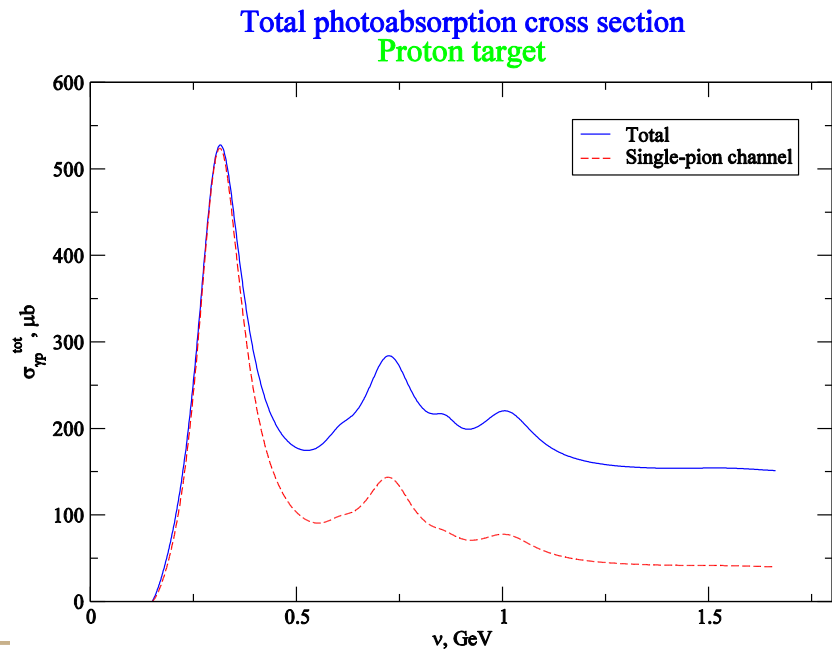
Compare with asymmetry caused by Coulomb distortion at small  $\theta \Rightarrow$   
 may differ by orders of magnitude depending on scattering kinematics

$$A_n^e(\text{Coulomb}) \propto \alpha \frac{m_e}{\sqrt{s}} \theta^3 \rightarrow A_n^e(\text{Diffractive}) \propto \alpha m_e (\sqrt{s}) \theta \cdot R_{\text{int}}^2$$

# Input parameters

For small-angle ( $-t/s \ll 1$ ) scattering of electrons with energies  $E_e$ , normal beam asymmetry is given by the energy-weighted integral

$$A_n \propto \frac{1}{E_e^2} \int_{\nu_{th}}^{E_e} d\nu \cdot \nu \sigma_{\gamma p}^{tot}(\nu; q_{1,2}^2 \approx 0)$$



$\sigma_{\gamma p}$  from N. Bianchi at al., Phys.Rev.C54 (1996)1688 (resonance region) and Block&Halzen, Phys.Rev. D70 (2004) 091901



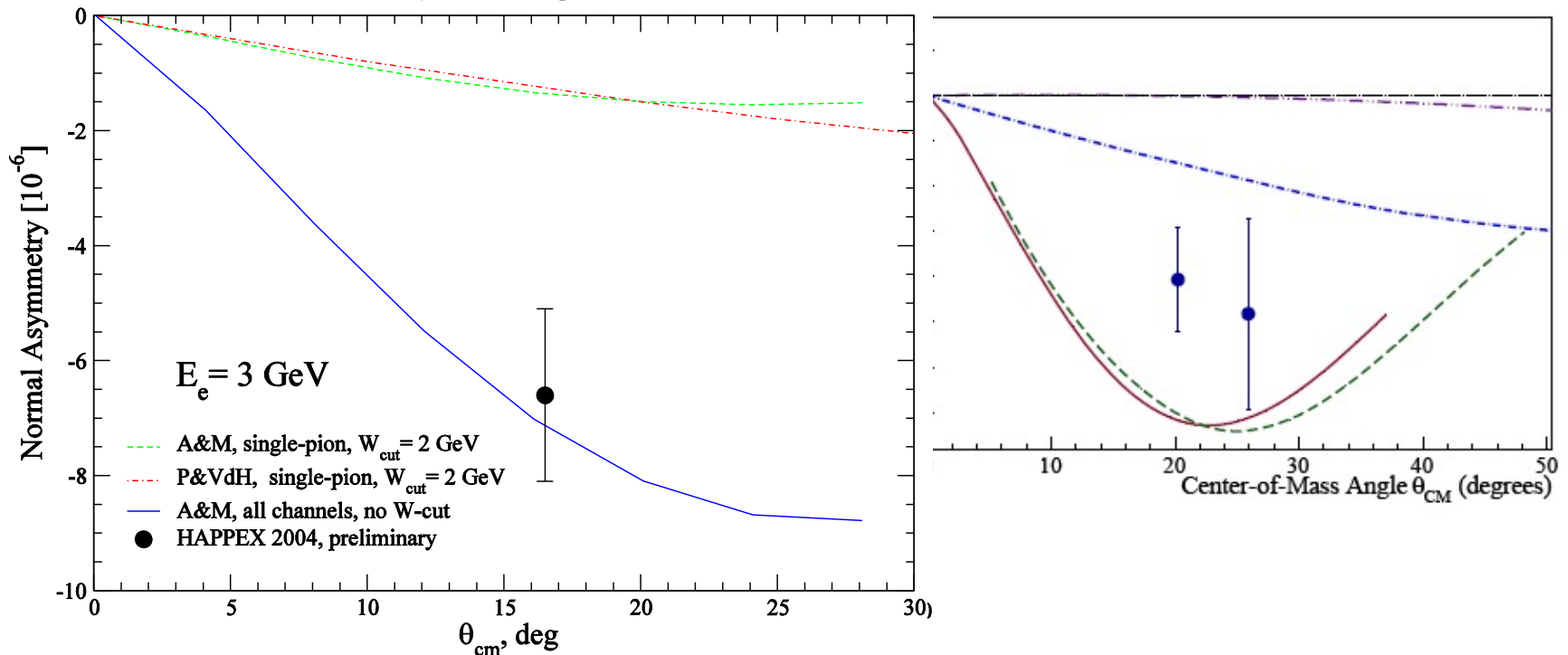
# Predictions vs experiment for Mott asymmetry

Use fit to experimental data on  $\sigma_{yp}$  (dotted lines include only one-pion+nucleon intermediate states)

G0 data Phys.Rev.Lett.99:092301,2007

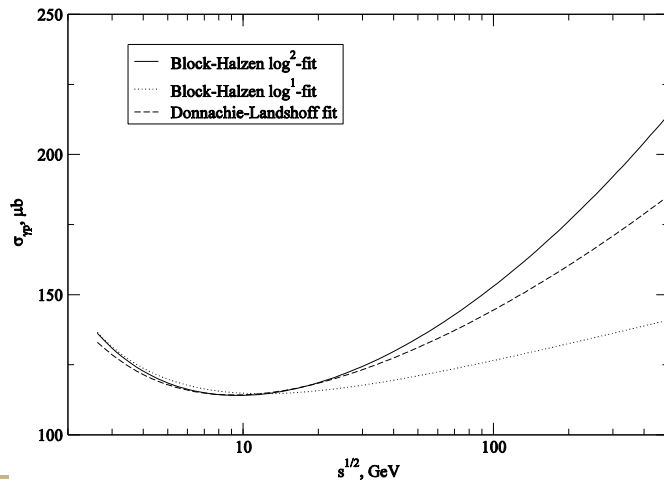
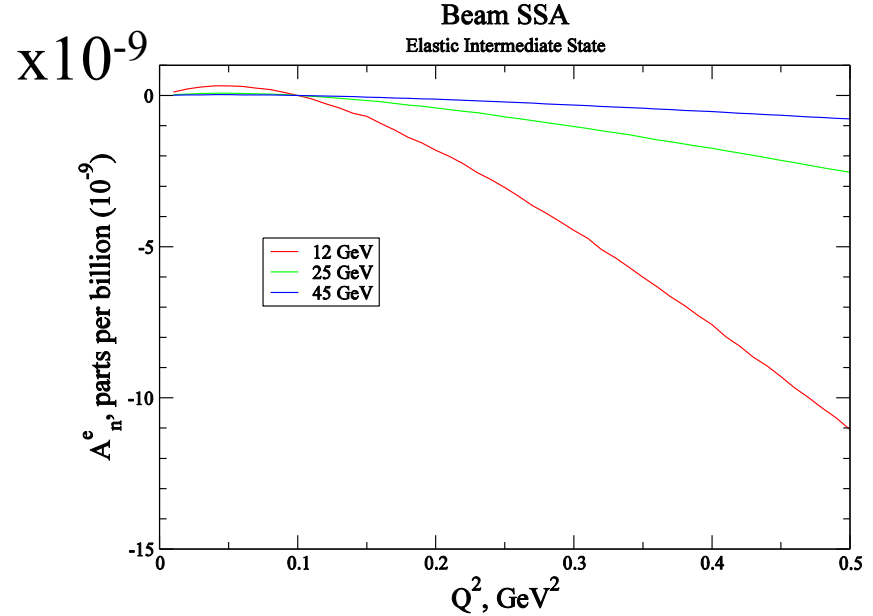
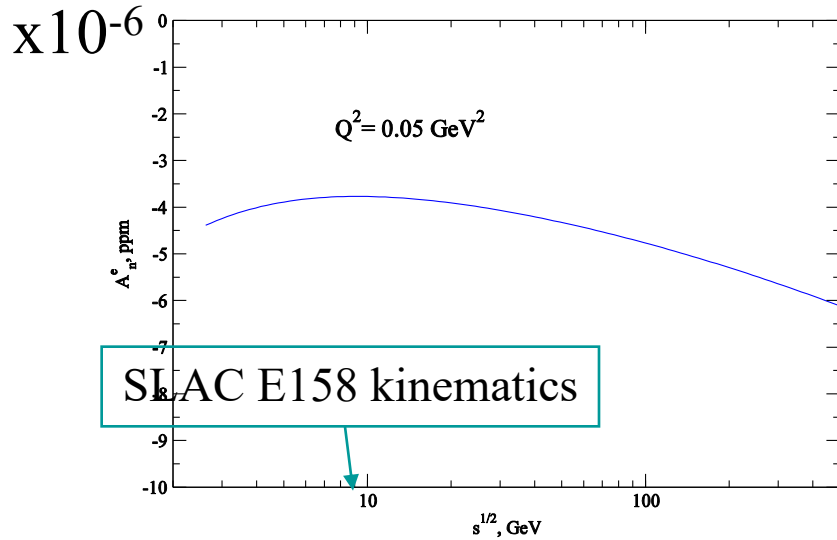
Normal beam asymmetry for elastic ep-scattering

Unitarity-based model predictions



HAPPEX

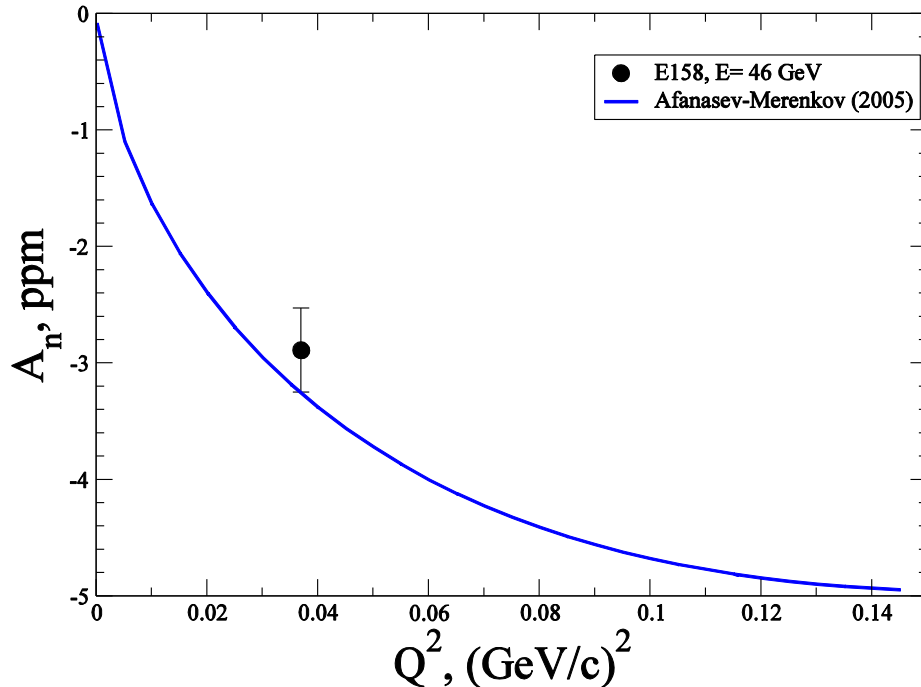
# Predict no suppression for Mott asymmetry with energy at fixed $Q^2$



- At 45 GeV predict beam asymmetry parts-per-million (diffraction) vs. parts-per billion (Coulomb distortion)

# Comparison with E158 data

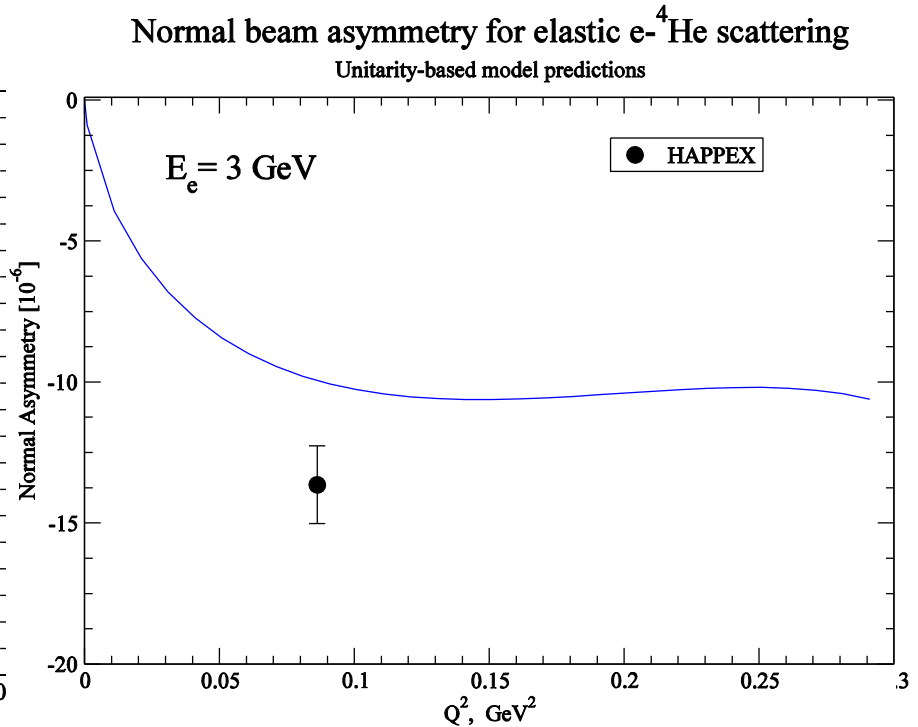
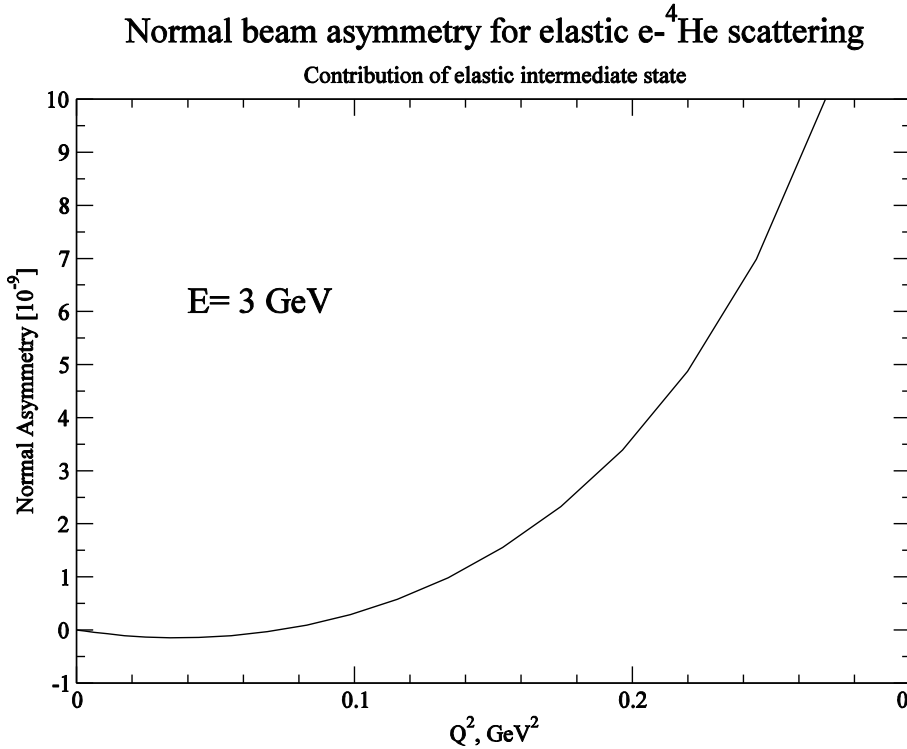
Elastic e+p scattering  
Normal beam asymmetry



- **SLAC E158:**  
 **$A_n = -2.89 \pm 0.36(\text{stat}) \pm 0.17(\text{syst})$  ppm**  
**(K. Kumar, private communication)**
- **Theory (AA, Merenkov):**  
 **$A_n = -3.2$  ppm**
- **Good agreement justifies application of this approach to the real part of two-boson exchange (Gorchtein et al and  $\gamma Z$  box calculations for small-angle scattering)**

# Mott Asymmetry on Nuclei

- Important systematic correction for parity-violation experiments ( $\sim -10$ ppm for HAPPEX on  $^4\text{He}$ ,  $\sim -5$ ppm for PREX on Pb), see AA *arXiv:0711.3065 [hep-ph]*; also Gorchtein, Horowitz, Phys.Rev.C77:044606,2008
- Coulomb distortion: only  $10^{-10}$  effect (Cooper&Horowitz, Phys.Rev.C72:034602,2005)



**Five orders of magnitude** enhancement in HAPPEX kinematics due to excitation of inelastic intermediate states in  $2\gamma$ -exchange (AA, Merenkov; use Compton data from Erevan)

# JLAB Experiments: HAPPEX, PREX

- Abrahamyan et al. New Measurements of the Transverse Beam Asymmetry for Elastic Electron Scattering from Selected Nuclei, PRL 109, 192501 (2012)

TABLE I. Kinematic values for the various targets.

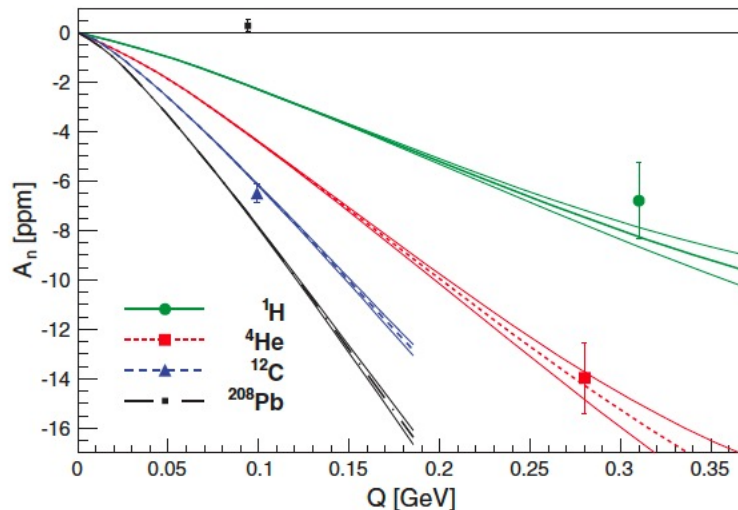
Target	H	<sup>4</sup> He	<sup>12</sup> C	<sup>208</sup> Pb
$\theta$	6°	6°	5°	5°
$Q^2(\text{GeV}^2)$	0.0989	0.0773	0.009 84	0.008 81
$E_b(\text{GeV})$	3.026	2.750	1.063	1.063
$\langle \cos\phi \rangle$	0.968	0.967	0.963	0.967

TABLE III. The measured  $A_n$  and derived  $\hat{A}_n$  values [Eq. (2)] for the four nuclei along with the corresponding total uncertainties  $A/Z$  and  $Q$ .

Target	H	<sup>4</sup> He	<sup>12</sup> C	<sup>208</sup> Pb
$A_n(\text{ppm})$	-6.80	-13.97	-6.49	0.28
$\sigma(A_n)(\text{ppm})$	$\pm 1.54$	$\pm 1.45$	$\pm 0.38$	$\pm 0.25$
$\sqrt{Q^2}(\text{GeV})$	0.31	0.28	0.099	0.094
$A/Z$	1.0	2.0	2.0	2.53
$\hat{A}_n(\text{ppm/GeV})$	-21.9	-24.9	-32.8	+1.2
$\sigma(\hat{A}_n)(\text{ppm/GeV})$	$\pm 5.0$	$\pm 2.6$	$\pm 1.9$	$\pm 1.1$

$$A_n = \hat{A}_n \frac{QA}{Z},$$

The formula captures dependence of the asymmetry on  $A$  and  $Z$ , and power dependence on  $Q$



Agreement for all lighter nuclei except <sup>208</sup>Pb

# JLAB Experiments: CREX data

- Adhikari et al (PREX and CREX Collab), New Measurements of the Beam-Normal Single Spin Asymmetry in Elastic Electron Scattering over a Range of Spin-0 Nuclei, PRL **128**, 142501 (2022)

TABLE I.  $A_n$  measurement kinematics.

$E_{\text{beam}}$ (GeV)	Target	$\langle\theta_{\text{lab}}\rangle$ (deg)	$\langle Q^2\rangle$ (GeV <sup>2</sup> )	$\langle\cos\phi\rangle$
0.95	<sup>12</sup> C	4.87	0.0066	0.967
0.95	<sup>40</sup> Ca	4.81	0.0065	0.964
0.95	<sup>208</sup> Pb	4.69	0.0062	0.966
2.18	<sup>12</sup> C	4.77	0.033	0.969
2.18	<sup>40</sup> Ca	4.55	0.030	0.970
2.18	<sup>48</sup> Ca	4.53	0.030	0.970
2.18	<sup>208</sup> Pb	4.60	0.031	0.969

Recently Coulomb distortion and inelastic excitations were considered in a unified approach: Koshchii, Gorchtein, Roca-Maza, Spiesberger C 103, 064316 (2021) but

**<sup>208</sup>Pb Asymmetry remains an unsolved mystery or “PREX Puzzle”**

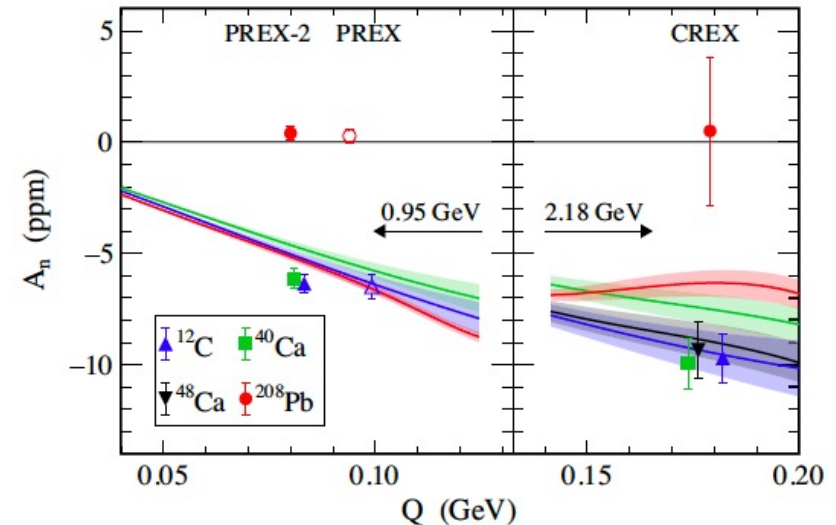


FIG. 2.  $A_n$  measurements from PREX-2, PREX (open circle and triangle, previously published [20]), and CREX at beam energies of 0.95 GeV, 1.06 GeV, and 2.18 GeV, respectively. The solid lines show theoretical calculations from [26] at 0.95 GeV and 2.18 GeV together with their respective one sigma uncertainty bands. The color of each band represents the calculation for the same color data point. Overlapping points are offset slightly in  $Q$  to make them visible.

# Inclusive Electroproduction of Pions (Singles)

- Reaction  $p(e_{\text{pol}}, \pi)X$ 
  - Parity-conserving spin-momentum correlation  $\vec{s}_e \cdot \vec{k}_e \times \vec{k}_\pi$
  - Introduced in Donnelly, Raskin, Annals Phys. 169, 247 (1986)
    - Can be shown to be a) due to  $R_{TL}$  response function (=fifth structure function) and b) not to integrate to zero after integration over momenta of the scattered electron
  - This is **NOT** a two-photon exchange effect (but suppressed by an electron mass)
    - Order-of magnitude estimate:  $A_n(ep \rightarrow \pi X) \sim A_{LT}(ep \rightarrow e' \pi N) * m_e / E' / \sin(\theta_e)$ 
      - Use MAMI data  $A_{LT}(ep \rightarrow e' \pi N) \sim 7\%$ , from Bartsch et al Phys.Rev.Lett.88:142001,2002  $\Rightarrow A_n(ep \rightarrow \pi X) \sim 250\text{ppm}$
    - See C.E. Carlson et al.

# Summary: SSA in Elastic ep-Scattering

- Collinear photon exchange present in (light particle) beam SSA
- Models violating EM gauge invariance **encounter collinear divergence** for target SSA
- VCS amplitude in *beam asymmetry* is enhanced in different kinematic regions compared to *target asymmetry*
- *Beam asymmetry is unsuppressed with energy in forward angles, follows the magnitude of total photoproduction cross section*
- Strong-interaction dynamics for Mott asymmetry in small-angle ep-scattering above the resonance region is *soft diffraction*
  - *For the diffractive mechanism  $A_n$  is a) not suppressed with beam energy and b) does not grow with  $Z$  ( $\sim A/Z$ )*
- *Good agreement for scattering on nuclei **except  $^{208}\text{Pb}$*** 
  - ***$^{208}\text{Pb}$  remains unresolved and requires detailed studies***