

Flavored dark matter beyond Minimal Flavor Violation

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P. Agrawal, M. Blanke, KG

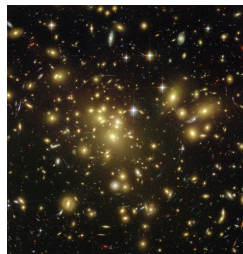
Outline

- 1 **Dark Minimal Flavor Violation**
- 2 **Analysis of flavor observables**
- 3 **Dark Matter Constraints**
- 4 **Combined analysis**

What we know about dark matter

Evidence for dark matter from astrophysics

- galactic rotation curves
- gravitationally lensing
- understanding of clusters
- angular fluctuations in Cosmic Microwave Background



Dark matter properties

- ☐ non-baryonic
- ☐ interacts gravitational through its mass
- ☐ relic density $\Omega_{DM} h^2 = 0.119$
- ☐ stable
- ☐ cold, non-relativistic
- ☐ neutral, no charge and no color

Flavored dark matter

Particle properties of Dark Matter?

- Coupling to SM particles?
- Single particle or entire sector?

→ Considering the role of flavor in the SM

Assumption

→ Dark matter carries flavor and comes in multiple copies



- ✓ non-baryonic
- ✓ neutral, no charge and no color



New coupling to quarks:

$$\lambda^{ij} \bar{d}_{Ri} \chi_j \phi$$

- | | |
|----------|----------------------------------|
| d_{Ri} | - Right-handed down quark |
| χ_j | - Dark matter particle, flavored |
| ϕ | - New scalar, colored |

Flavored dark matter is not new!

Selection only!

- Flavored Dark Matter in Direct Detection Experiments and at LHC
J. KILE, A. SONI (APRIL 2011)
- Dark Matter from Minimal Flavor Violation
B. BATELL, J. PRADLER, M. SPANNOVSKY (MAY 2011)
- Discovering Dark Matter Through Flavor Violation at the LHC
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- Flavored Dark Matter, and Its Implications for Direct Detection and Colliders
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→ BUT common to all these studies

Minimal Flavor Violation

Minimal flavor violation

- **SM:** GIM mechanism controls pattern of FCNC processes formula
- **NP:** large new sources of flavor symmetry breaking are excluded at TeV scale

ISIDORI, NIR, PEREZ (2010)

Example: bounds from B_d -mixing with generic flavor structure

$$M(B_d - \bar{B}_d) \sim \frac{(y_t V_{ti}^* V_{tj})^2}{16\pi^2 M_W^2} + c_{NP} \frac{1}{\Lambda^2} \quad \text{for } c_{NP} \sim \mathcal{O}(1) \quad \Rightarrow \quad \Lambda \geq 10^3 \text{TeV}$$

→ new flavor couplings to quarks cannot be arbitrary but are strongly constrained by flavor observables

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→ new flavor couplings to quarks cannot be arbitrary but are strongly constrained by flavor observables

main features of MFV:

- flavor symmetry is only broken by the SM Yukawa couplings
- CKM matrix is the only source of flavor violation
- FCNCs are naturally suppressed

So, what should we do?

→ Minimal Flavor Violation???

non-MFV



→ DANGEROUS

But interesting if you
know how to handle it!

MFV



→ HARMLESS

But not very exciting.

The new model

- dirac fermionic DM χ carries flavor and couples to quarks via a scalar colored mediator

$$\mathcal{L}_{\text{NP}} = i\bar{\chi}\not{\partial}\chi - m_{\chi}\bar{\chi}\chi + (D_{\mu}\phi)^{\dagger}(D^{\mu}\phi) - m_{\phi}^2\phi^{\dagger}\phi - \lambda^{ij}\bar{d}_{Ri}\chi_j\phi + \text{h.c.} \\ + \lambda_{H\phi}\phi^{\dagger}\phi H^{\dagger}H + \lambda_{\phi\phi}\phi^{\dagger}\phi\phi^{\dagger}\phi$$

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$$\mathcal{L}_{\text{NP}} = i\bar{\chi}\not{D}\chi - m_{\chi}\bar{\chi}\chi + (D_{\mu}\phi)^{\dagger}(D^{\mu}\phi) - m_{\phi}^2\phi^{\dagger}\phi - \lambda^{ij}\bar{d}_{Ri}\chi_j\phi + \text{h.c.} \\ + \lambda_{H\phi}\phi^{\dagger}\phi H^{\dagger}H + \lambda_{\phi\phi}\phi^{\dagger}\phi\phi^{\dagger}\phi$$

Assumption

- Flavor symmetry: $U(3)_q \times U(3)_u \times U(3)_d \times U(3)_{\chi}$

→ only broken by the SM Yukawa couplings and the DM-quark coupling λ

“Dark Minimal Flavor Violation”

(DMFV)

Dark Minimal Flavor Violation - special features

Parametrization of DM-quark coupling

- $U(3)_\chi$ symmetry helps to remove 9 parameters

$$\lambda = U_\lambda D_\lambda$$

U_λ - unitary matrix, 3 mixing angles s_{12}^λ , s_{13}^λ , s_{23}^λ and 3 phases

D_λ - real diagonal matrix, e.g. $D_\lambda = \lambda_0 \cdot \mathbb{1} + \text{diag}(\lambda_1, \lambda_2, -(\lambda_1 + \lambda_2))$

Dark matter mass

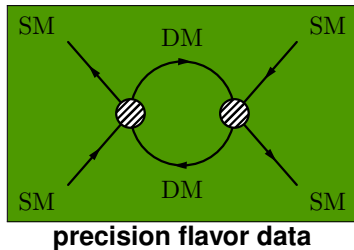
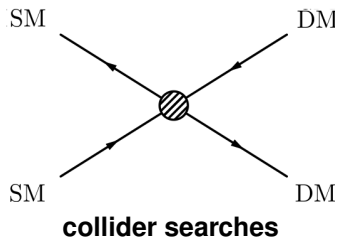
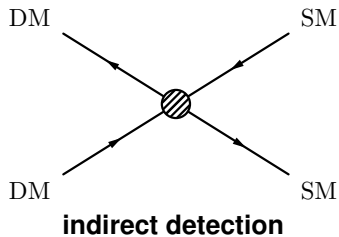
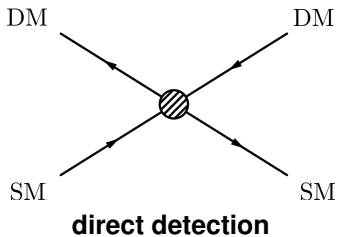
- $U(3)_\chi$ symmetry ensures equal mass for all flavors at tree level
- special form of mass splitting at higher order (loop level)

$$m_{ij} = m_\chi (\delta_{ij} + \eta \lambda_{ik}^\dagger \lambda_{kj}) = (m_\chi + m_\chi \eta D_{\lambda,ii}^2) \delta_{ij}$$

Dark matter stability

- DM stability is guaranteed if Dark Minimal Flavor Violation is exact
→ talk by Monika Blanke

A new set of observables

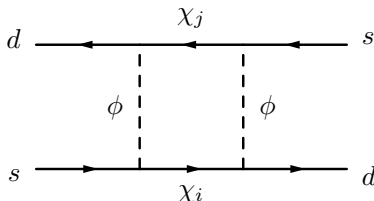


Main goals of our analysis

- 1 Which are the **allowed structures of the coupling matrix λ** beyond MFV?
- 2 What are **viable mass hierarchies** for the flavored dark matter spectrum?
- 3 What happens when **combining flavor and dark matter constraints**?
- 4 How strong are the **mass bounds** on the new particles from LHC searches?

New contributions to meson anti-meson mixing

- new box diagram:



- dominant NP mixing amplitude for the K meson system

$$M_{12}^{K,\text{new}} \sim (\xi_K^*)^2 F(x) \quad \text{where} \quad \xi_K = (\lambda\lambda^\dagger)_{sd} = \sum_{i=1}^3 \lambda_{si} \lambda_{di}^*$$

- ξ_K - involves elements of matrix λ , dependent on the meson system
- $F(x)$ - box loop function

- analogous contributions to $B_{d,s} - \bar{B}_{d,s}$ mixing

The $B \rightarrow X_s \gamma$ decay

- effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} \sim (C_7 Q_7 + C_7' Q_7' + \dots)$$

with

$$Q_7 \sim \bar{s}_L \sigma^{\mu\nu} b_R F_{\mu\nu}$$

$$Q_7' \sim \bar{s}_R \sigma^{\mu\nu} b_L F_{\mu\nu}$$

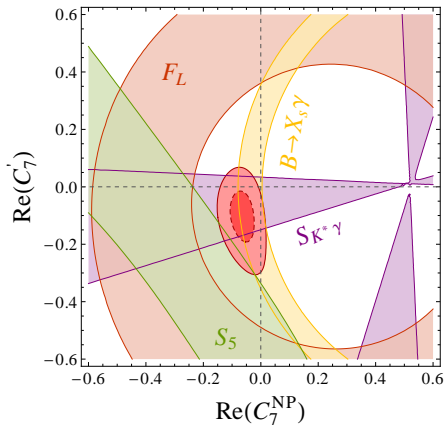
- SM: C_7' is strongly suppressed due to the chiral structure of weak interactions

$$C_{7,\text{SM}}' = \frac{m_s}{m_b} C_{7,\text{SM}}$$

- new contribution

$$|\delta C_7'| \sim 0.04 \left[\frac{500 \text{ GeV}}{m_\phi} \right]^2 \left| \sum_{i=1}^3 \lambda_{si} \lambda_{bi}^* \right|$$

FIGURE FROM: ARXIV:1308.1501
W. ALTMANNSHOFER, D. STRAUB



→ NEGLIGIBLE

Negligible effects in ...

Rare decays $K \rightarrow \pi \nu \bar{\nu}$, $B_{s,d} \rightarrow \mu^+ \mu^-$ **and** $B \rightarrow K^* \mu^+ \mu^-$

- **no box contribution**
since no coupling to leptons in final states
- **Z penguin contribution is zero**
due to chiral structure/new couplings to right-handed quarks only
- **γ penguin is negligible**
estimate from supersymmetric models

Electric dipole elements

- **no relevant contribution** since chirality flips are required

Electroweak precision tests

- γ and Z self-energies (scalar ϕ in the loop) are **highly suppressed**

Strategy for phenomenology

Step 1: **Prealysis of flavor constraints**

- parametrize full mixing amplitude, e.g. B_q mixing

$$M_{12}^{B_q} = C_{B_q} e^{2i\varphi_{B_q}} M_{12}^{B_q, \text{SM}} \quad (q = d, s)$$

SM corresponds to $C_{B_q} = 1$ and $\varphi_{B_q} = 0$

- use results of the model-independent NP fit by the UTfit collaboration
- fix flavor conserving parameters m_ϕ , m_χ and λ_0 to the values

$$m_\phi = 850 \text{ GeV} \quad m_\chi = 200 \text{ GeV} \quad \lambda_0 = 1$$

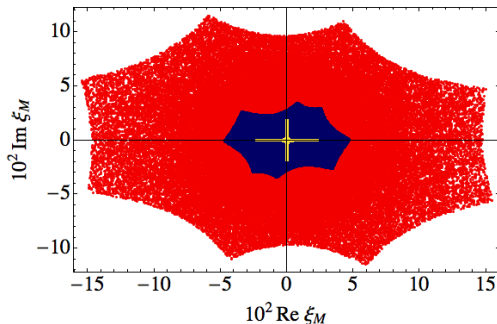
Step 2: **Combined analysis of flavor and DM constraints**

Lessons from flavor physics

$K^0 - \bar{K}^0$ mixing

$B_d^0 - \bar{B}_d^0$ mixing

$B_s^0 - \bar{B}_s^0$ mixing



- strongest constraint comes from $K^0 - \bar{K}^0$ mixing, the CP-violating parameter ε_K
- flavor constraints are under control if

$$\begin{aligned}\xi_M &\sim \lambda \lambda^\dagger \\ &= U_\lambda D_\lambda D_\lambda^\dagger U_\lambda^\dagger = U_\lambda D_\lambda^2 U_\lambda^\dagger \\ &\sim \text{diag}(\star_1, \star_2, \star_3)\end{aligned}$$

→ “Universality”

“Flavor safe” dark matter scenarios

Universality is automatically fulfilled for ...

❶ **universal scenario**

(black):

$$\lambda_1 = \lambda_2 = 0$$

❷ **12-degeneracy**

(blue): $\lambda_1 = \lambda_2$

❸ **13-degeneracy**

(red): $\lambda_2 = -2\lambda_1$

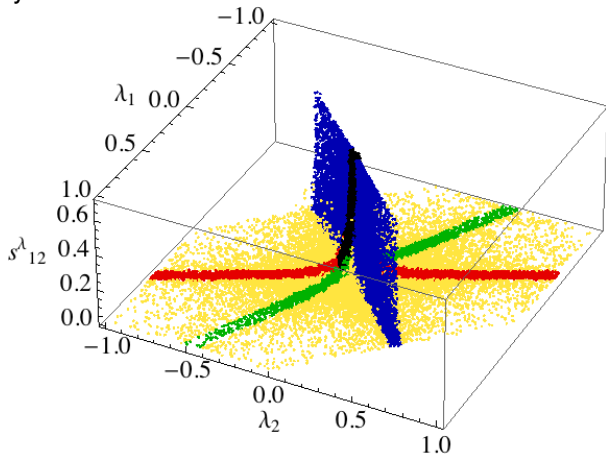
❹ **23-degeneracy**

(green):

$$\lambda_2 = -1/2\lambda_1$$

❺ **small mixing**

(yellow): arbitrary D_λ



RECALL: $D_\lambda = \lambda_0 \cdot \mathbb{1} + \text{diag}(\lambda_1, \lambda_2, -(\lambda_1 + \lambda_2))$

$$m_{ij} = (m_\chi + m_\chi \eta D_{\lambda,ii}^2) \delta_{ij}$$

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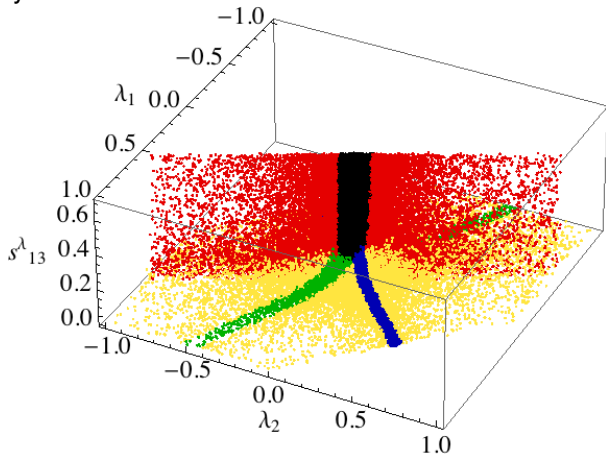
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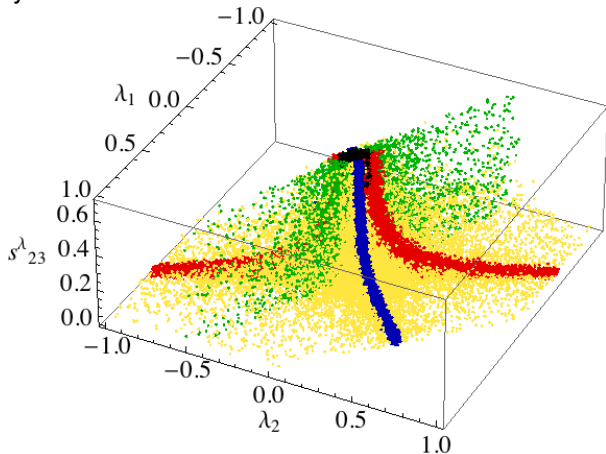
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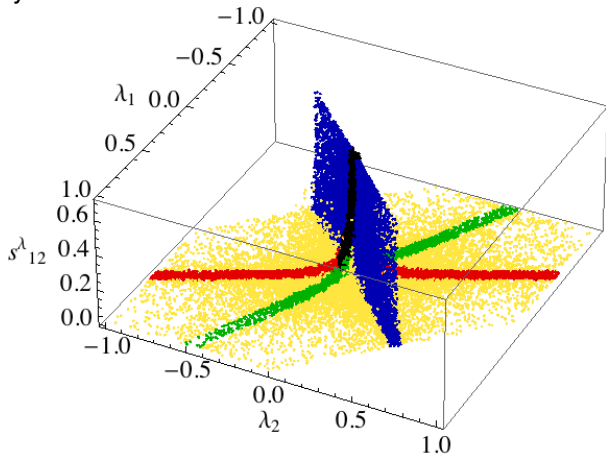
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Recovering the Minimal Flavor Violation limit

- Minimal Flavor Violation:
quark-dark matter coupling has a very specific structure

$$\lambda = \alpha \mathbb{1} + \beta Y_d^\dagger Y_d$$

where

$$Y_d = \frac{\sqrt{2}}{v} \text{diag}(m_d, m_s, m_b)$$

- mass pattern is fixed through

$$m_{ij} = m_\chi (\delta_{ij} + \eta \lambda_{ik}^\dagger \lambda_{kj})$$

- only a small subset of specific points fulfills condition
→ near the 12-degeneracy line, where ALL mixing angles are small

→ DMFV is clearly BEYOND MFV
only the concept is similar

Mass hierarchies in the dark sector

Flavor observables do not fix the mass spectrum m_{χ_i} !

☹ *d*-flavored dark matter ☹

→ severely constrained by direct detection experiments and LHC searches

s- and *b*-flavored dark matter

→ similar for flavor physics and direct detection

☺ *b*-flavored dark matter ☺

→ *b*-jet signatures at colliders

→ possible explanation of γ ray signal from galactic center

P. AGRAWAL, B. BATELL, D. HOOPER, T. LIN (2014)

WE ASSUME ALWAYS:

→ *b*-flavored dark matter

$$\begin{aligned} m_{\chi_b} &< m_{\chi_d}, m_{\chi_s} \\ D_{\lambda,33} &> D_{\lambda,11}, D_{\lambda,22} \end{aligned}$$

Recall:

$$m_{\chi_i} = m_{\chi}(1 - |\eta| D_{\lambda,ii}^2)$$

Dark matter phenomenology

Step 1: Preanalysis of flavor constraints

- so far: flavor conserving parameters have been fixed

Step 2: Combined analysis of flavor and DM constraints

- parameters are varied as follows

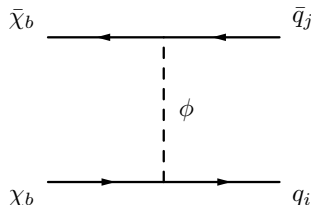
$$m_\phi = 850 \text{ GeV} \quad m_\chi : \text{ free} \quad \lambda : \text{ free}$$

- 1 “single-flavor” freeze-out:
for large mass splittings $\gtrsim 10\%$ between DM flavors only
lightest flavor remains in the thermal bath
- 2 “two-flavor” freeze-out:
if small mass splittings $\lesssim 1\%$ between DM flavors multiple
flavors can be present at freeze-out
- 3 (“three-flavor” freeze-out)

Relic abundance for single-flavor freeze-out

- relic abundance of the dark matter is set by annihilation

$$\langle \sigma v \rangle_{bb} = \frac{3 D_{\lambda,33}^4 m_{\chi_b}^2}{32\pi(m_{\chi_b}^2 + m_\phi^2)^2}$$



- and is determined by solving the Boltzmann equation for the dark matter number density n at late times

$$\frac{dn}{dt} + 3Hn = - \underbrace{\langle \sigma v \rangle_{eff}}_{2.2 \times 10^{-26} \text{cm}^3/\text{s}} (n^2 - n_{eq}^2)$$

- | | |
|----------------------------------|---|
| n | - dark matter number density |
| H | - Hubble constant |
| n_{eq} | - equilibrium number density of χ |
| $\langle \sigma v \rangle_{eff}$ | - $\langle \sigma v \rangle_{eff} = \frac{1}{2} \langle \sigma v \rangle$ |



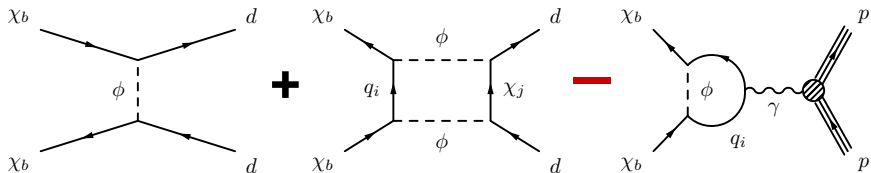
relic density

Constraints from direct detection

- spin-independent contribution to the WIMP-nucleus scattering

$$\sigma_n^{SI} = \frac{\mu_n^2}{\pi} (Zf_p + (A - Z)f_n)^2$$

relevant processes:



→ cancelation between the direct detection box diagram and the one-loop photon contribution

→ we apply the bounds from the LUX experiment

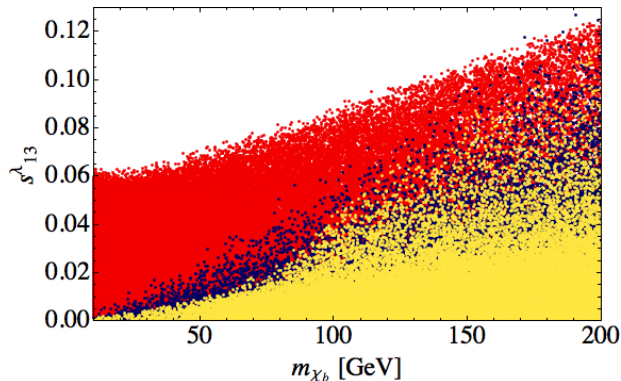
Results of combined analysis

constraints imposed:

Relic abundance
fixes $D_{\lambda,33}$

- LUX only
- flavor only
- LUX & flavor

single-flavor freeze-out



→ non-trivial interplay of dark matter and flavor constraints

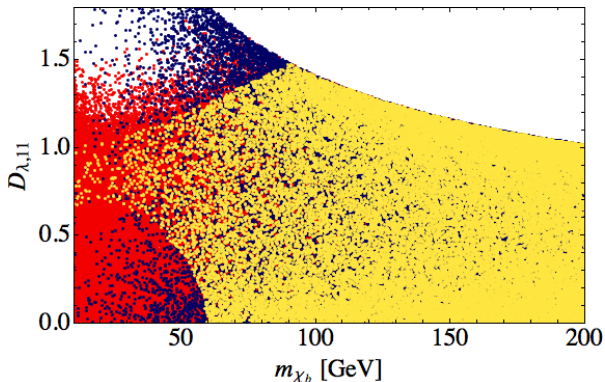
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→ upper and a lower bound on the size of $D_{\lambda,11}$

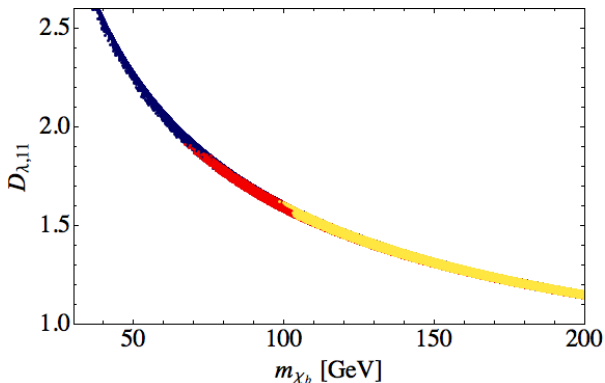
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 $D_{\lambda,11}$ & $D_{\lambda,33}$

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- LUX & flavor

two-flavor freeze-out e.g. 13-degeneracy



→ 13-degeneracy scenario ruled out below $m_{\chi_b} \simeq 100 \text{ GeV}$

Conclusions

- the mechanism to generate the flavor structure of the SM is unknown, assuming a similar mechanism in the dark sector suggests

“Dark Minimal Flavor Violation”

flavor symmetry, enlarged by an additional $U(3)_\chi$, is only broken by the new coupling λ and SM Yukawas

- the lightest dark matter particle is stable if DMFV is exact

→ talk next week by Monika Blanke

- “flavor-safe” scenarios - beyond MFV - can be identified

these can be directly used for further study

- non-trivial interplay of DM and flavor phenomenology

A brief look at collider phenomenology

new particles within the reach of LHC

Dark Minimal Flavor Violation

- χ_i are nearly degenerate
- new particles have to be pair-produced
- particle spectrum is similar to simplified models of squarks and neutralinos in the MSSM

1. dark matter fermion χ_b and the heavier flavors $\chi_{d,s}$

- $\chi_{d,s}$ decay to χ_b produces soft particles (jets, photons) + missing E_T

2. coloured scalar mediator ϕ

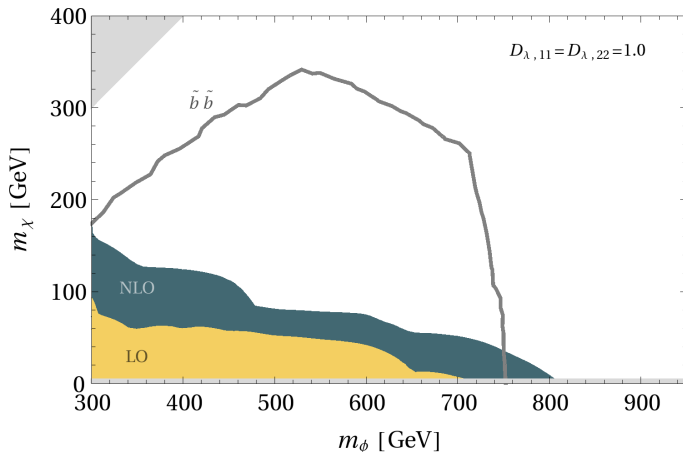
- pair-produced through QCD and through t -channel χ_d exchange
- decay $\phi \rightarrow q_i \chi_i$ with branching ratios given by $D_{\lambda,ii}^2$
 $\Rightarrow bb + \cancel{E}_T, bj + \cancel{E}_T, jj + \cancel{E}_T$ signatures

Mass bounds from dijet constraints

- CMS (& ATLAS) put strong bounds on sbottom and squark pair-production

CMS-PAS-SUS-13-018

- bound on cross-section can be applied to DMFV

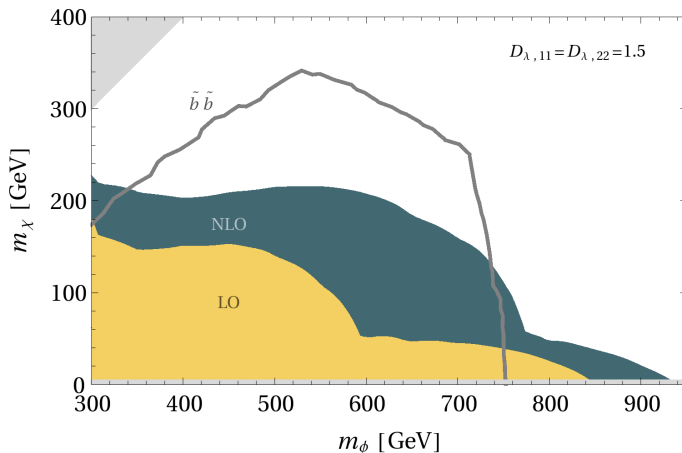


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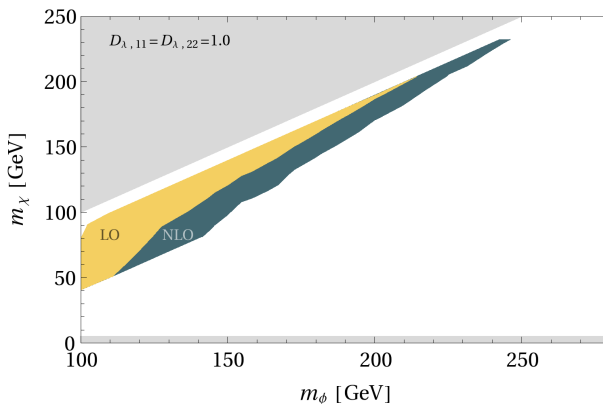
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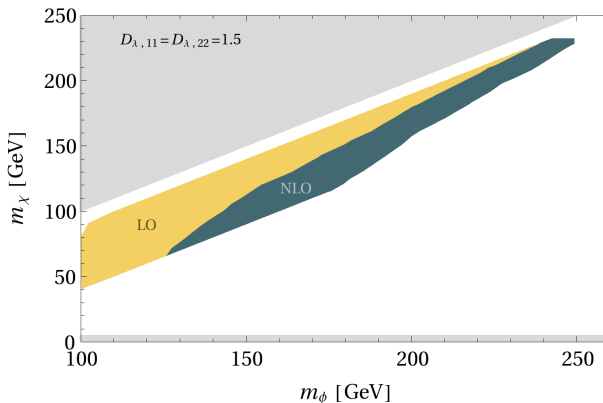
Constraints from Monojet Searches

- monojet searches sensitive to ϕ pair-production if decay products are soft
- constraint on the compressed region $m_\chi \lesssim m_\phi$



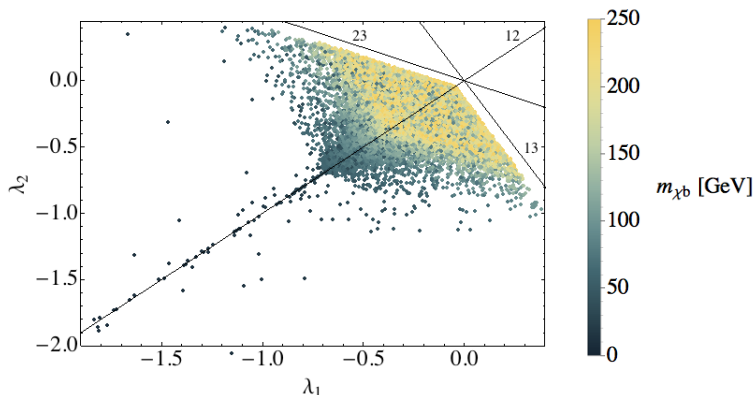
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Recovering flavor scenarios

Example: single-flavor freeze-out

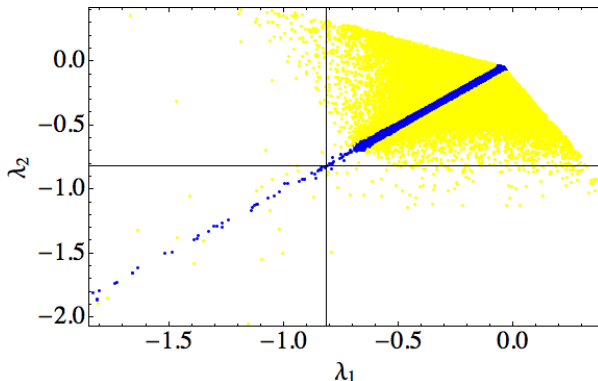


Recall: $D_\lambda = \lambda_0 \cdot \mathbb{1} + \text{diag}(\lambda_1, \lambda_2, -\lambda_1 - \lambda_2)$

→ small $m_{\chi b}$ implies sizeable non-universality $\lambda_{1,2} \neq 0$

Recovering flavor scenarios

Example: single-flavor freeze-out



Recall: $D_\lambda = \lambda_0 \cdot \mathbb{1} + \text{diag}(\lambda_1, \lambda_2, -\lambda_1 - \lambda_2)$

→ only 12-degeneracy and small mixing scenario survive

Dark matter stability

Consider the operator

SIMILAR PROOF IN MFV: ARXIV:1105.1781

B. BATELL, J. PRADLER, M. SPANNOWSKY

$$\mathcal{O} \sim \chi \dots \bar{\chi} \dots \phi \dots \phi^\dagger \dots q_L \dots \bar{q}_L \dots u_R \dots \bar{u}_R \dots d_R \dots \bar{d}_R \dots G \dots S$$

invariant under ...

- **QCD** if the number of $SU(3)_c$ triplet minus the number of $SU(3)_c$ antitriplets is a multiple of three
- **flavor symmetry** if $Y_u \dots Y_u^\dagger \dots Y_d \dots Y_d^\dagger \dots \lambda \dots \lambda^\dagger \dots$

	Invariance	Condition
I	$SU(3)$	$(N_\phi - N_{\phi^\dagger} + N_q + N_u + N_d - N_{\bar{q}} - N_{\bar{u}} - N_{\bar{d}}) \bmod 3 = 0$
II	$U(3)_q$	$(N_q - N_{\bar{q}} + N_{Y_u} - N_{Y_u^\dagger} + N_{Y_d} - N_{Y_d^\dagger}) \bmod 3 = 0$
III	$U(3)_u$	$(N_u - N_{\bar{u}} - N_{Y_u} + N_{Y_u^\dagger}) \bmod 3 = 0$
IV	$U(3)_d$	$(N_d - N_{\bar{d}} - N_{Y_d} + N_{Y_d^\dagger} + N_\lambda - N_{\lambda^\dagger}) \bmod 3 = 0$
V	$U(3)_\chi$	$(N_\chi - N_{\bar{\chi}} - N_\lambda + N_{\lambda^\dagger}) \bmod 3 = 0$

$$\sum \text{II+III+IV+V-I} \quad (N_\chi - N_{\bar{\chi}} - N_\phi + N_{\phi^\dagger}) \bmod 3 = 0$$

$\rightarrow \chi$ and ϕ decays into SM fields forbidden