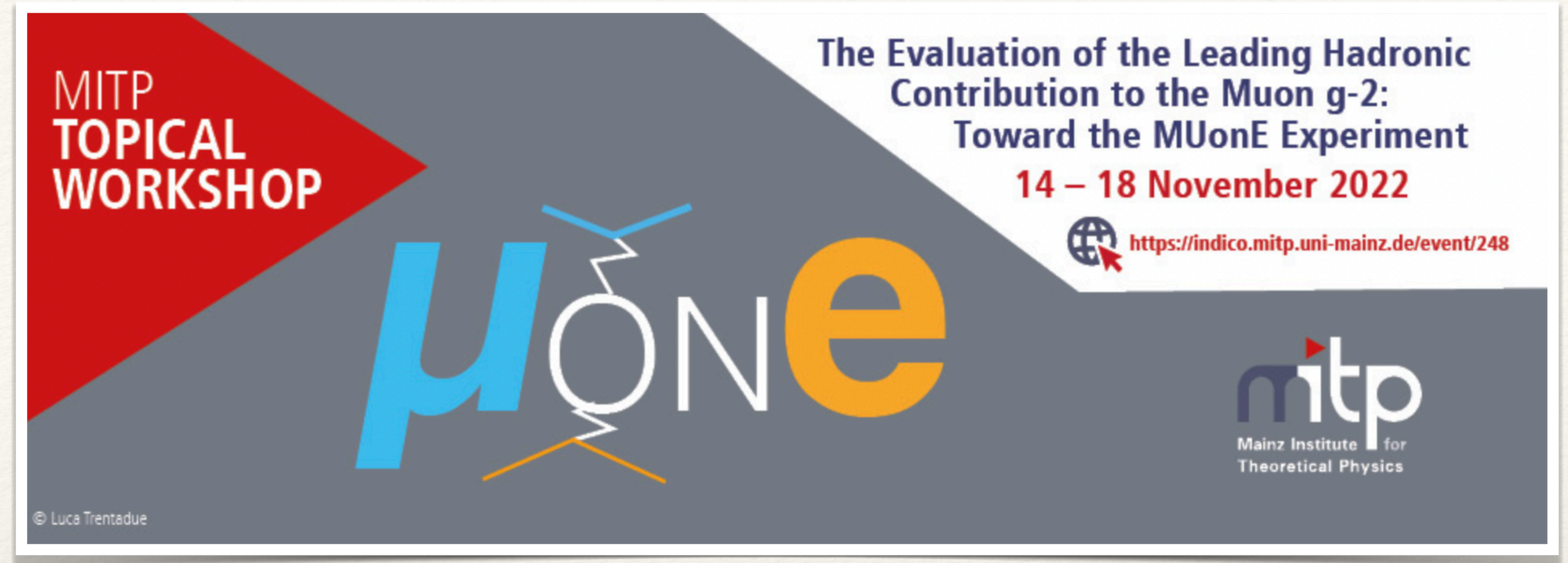


MUONE2022

Mainz Institute for Theoretical Physics

Nov 14-18 2022



MITP
TOPICAL
WORKSHOP

The Evaluation of the Leading Hadronic
Contribution to the Muon $g-2$:
Toward the MUonE Experiment
14 – 18 November 2022

<https://indico.mitp.uni-mainz.de/event/248>

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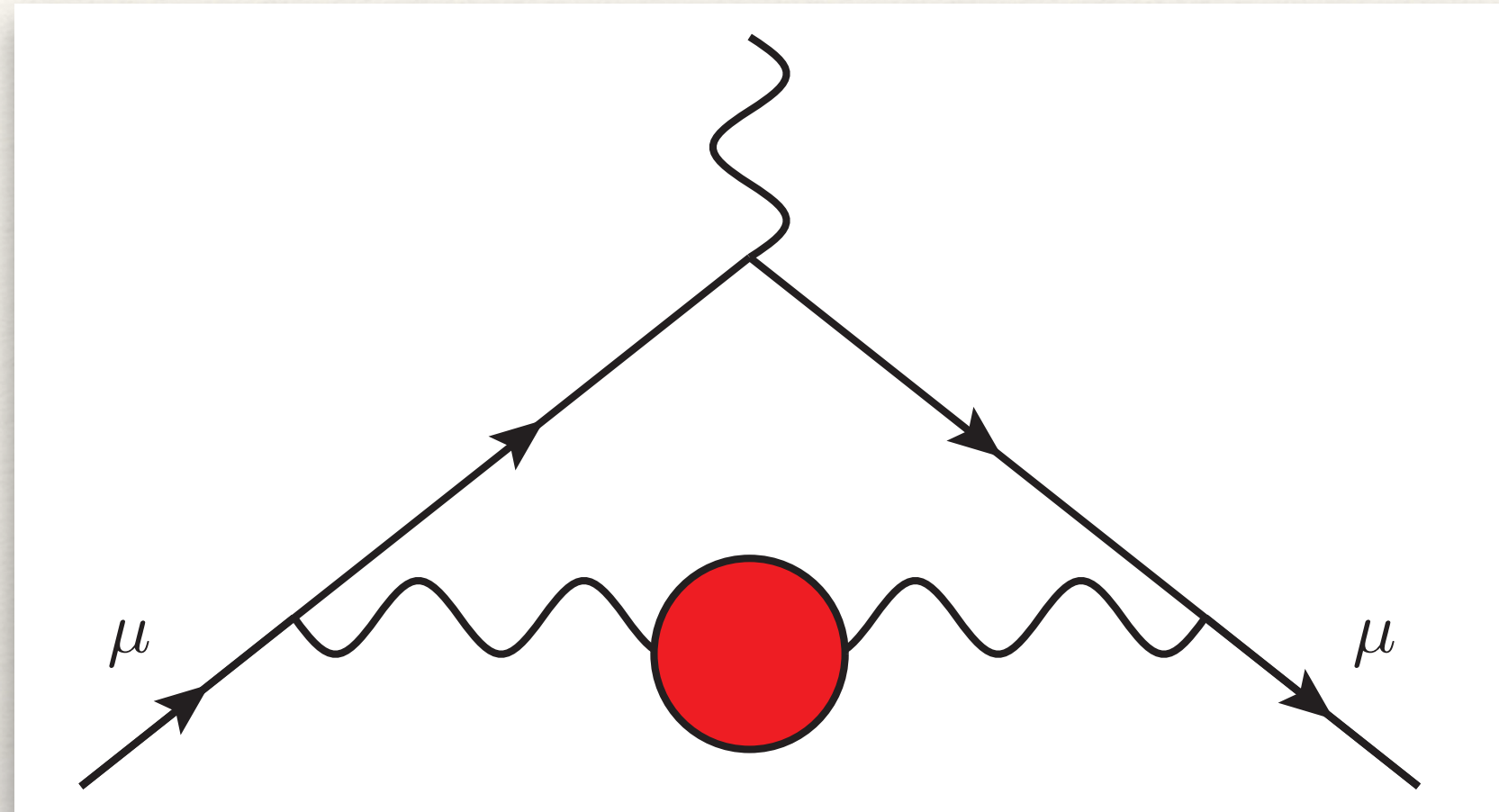
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NLO HVP contributions to the muon $g - 2$ in the space-like region

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Università degli Studi di Padova
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Work in Collaboration with S. Laporta and M. Passera, *Phys.Lett.B* 834 (2022) 137462

The hadronic LO contribution: time-like method



$$a_{\mu}^{HVP}(LO) = \frac{\alpha}{\pi^2} \int_{s_0^2}^{\infty} \frac{ds}{s} K^{(2)}(s/m^2) \text{Im}\Pi_h(s)$$

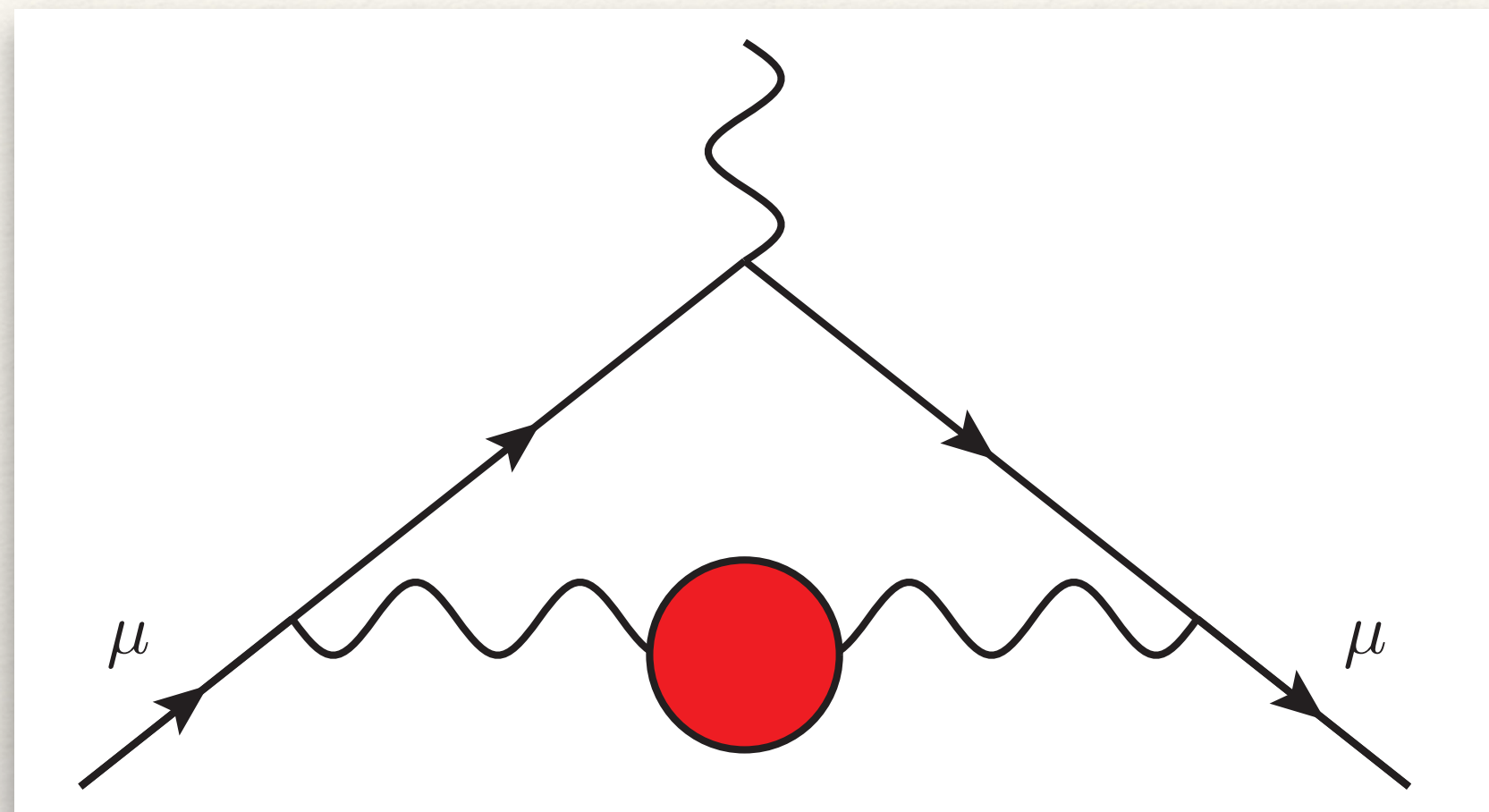
$$\text{Im}\Pi_h(s) = \frac{\alpha}{3} R(s) \quad R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{4\pi\alpha^2/(3s)}$$

$$K^{(2)}(s/m^2) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)}$$

Defining $z = \frac{q^2}{m^2}$ & $y(z) = \frac{z - \sqrt{z(z-4)}}{z + \sqrt{z(z-4)}}$:

$$K^{(2)}(z) = \frac{1}{2} - z + \left(\frac{z^2}{2} - z \right) \ln z + \frac{\ln y(z)}{\sqrt{z(z-4)}} \left(z - 2z^2 + \frac{z^3}{2} \right)$$

The hadronic LO contribution: time-like method



$$\begin{aligned} a_{\mu}^{HVP}(LO) &= 6895 (33) \times 10^{-11} \\ &= 6939 (40) \times 10^{-11} \\ &= 6928 (24) \times 10^{-11} \\ &= 6931 (40) \times 10^{-11} (0.6\%) \end{aligned}$$

F. Jegerlehner, arXiv:1711.06089

Davier, Hoecker, Malaescu, Zhang, arXiv:1908.00921

Keshavarzi, Nomura, Teubner, arXiv: 1911.00367

WP20 value

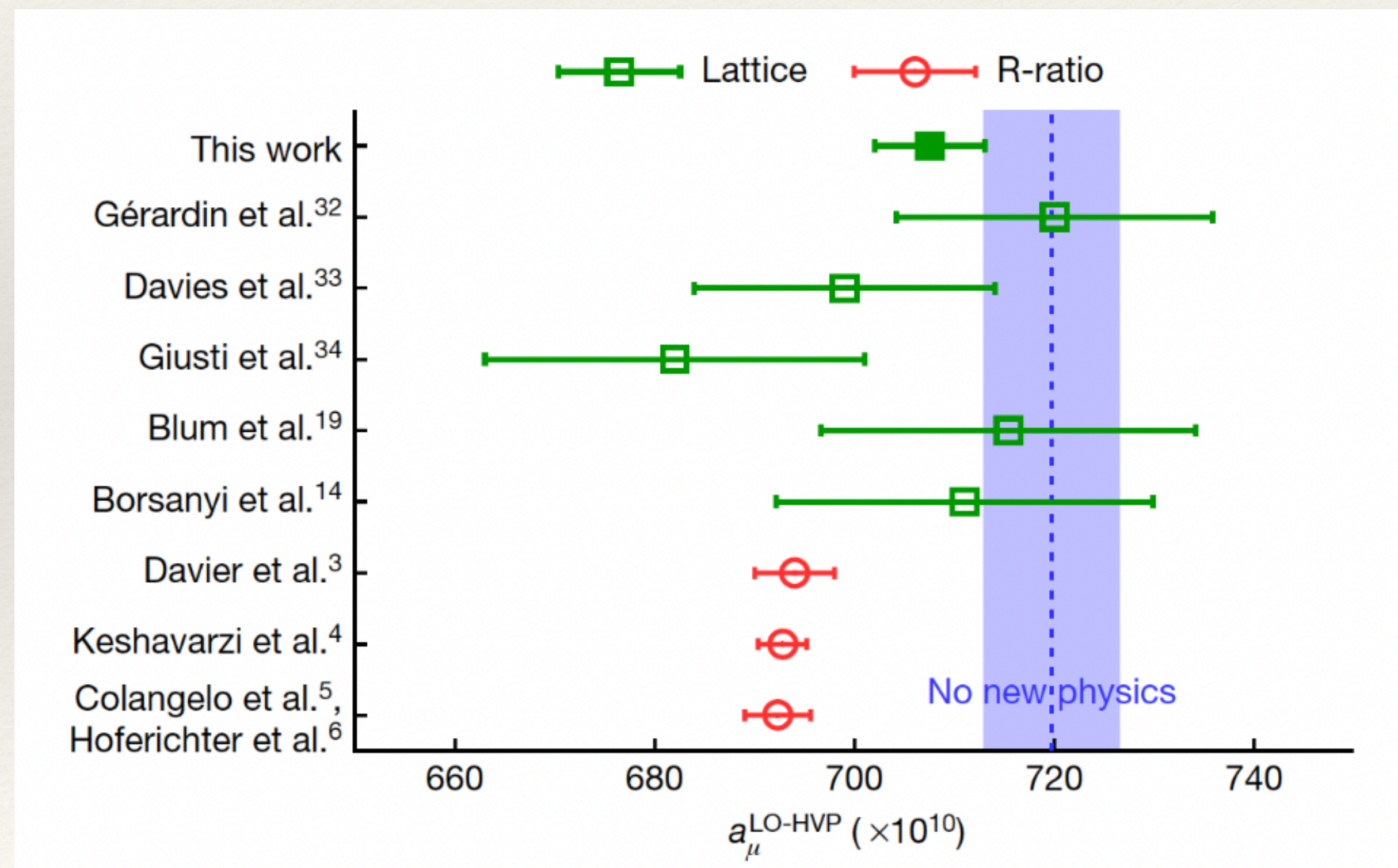
The hadronic LO contribution from lattice QCD

The BMW collaboration presented the first lattice QCD determination of a_μ^{HLO} with a 0.8 % precision

$$a_\mu^{HLO} = 7075(23)_{stat}(50)_{syst}[55]_{tot} \times 10^{-11}$$

2.2 σ tension with the time-like data driven determination.

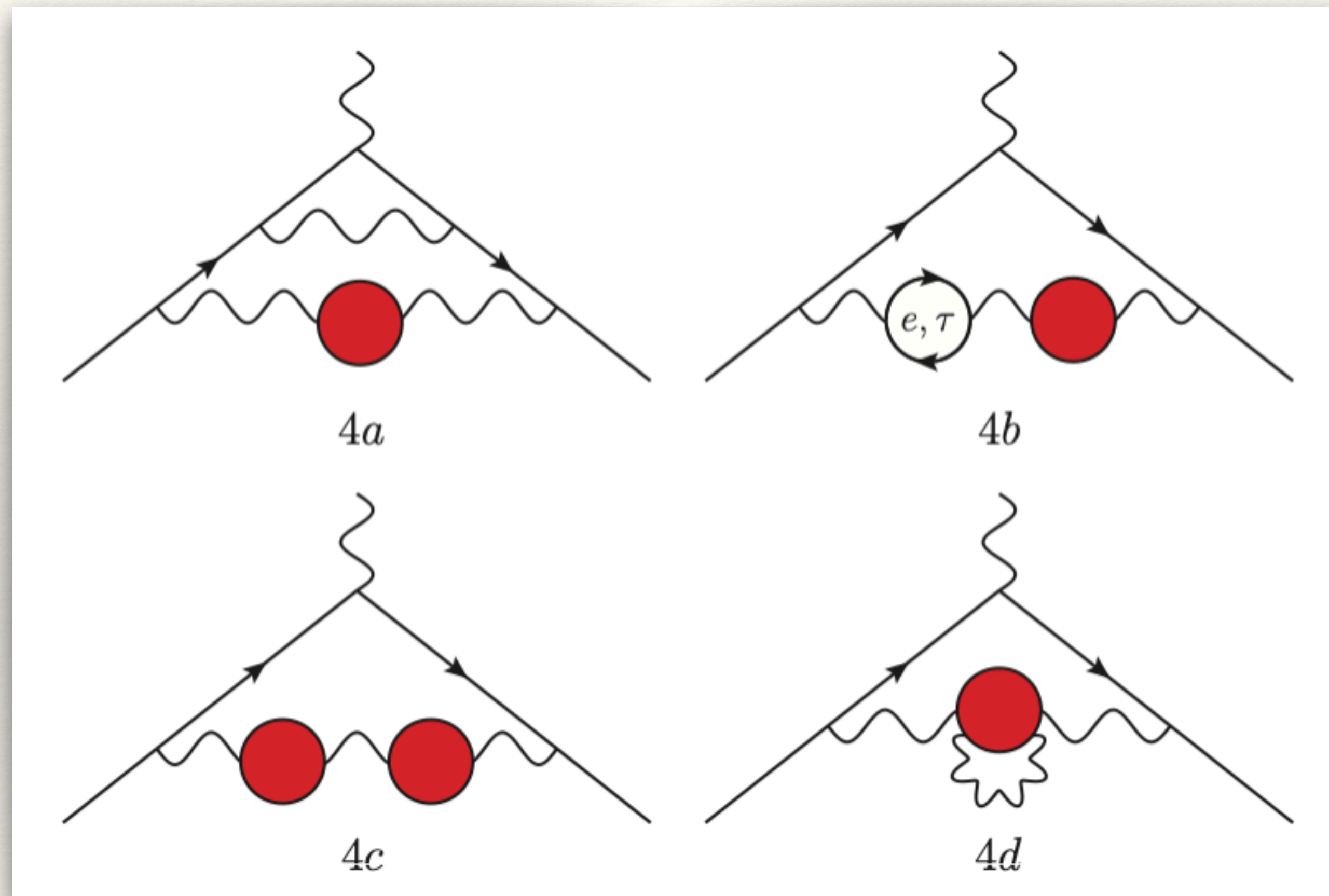
BMW collaboration 2021



Borsanyi et al (BMWc), Nature 2021

The hadronic NLO VP contribution

- $O(\alpha^3)$ contributions containing HVP insertion

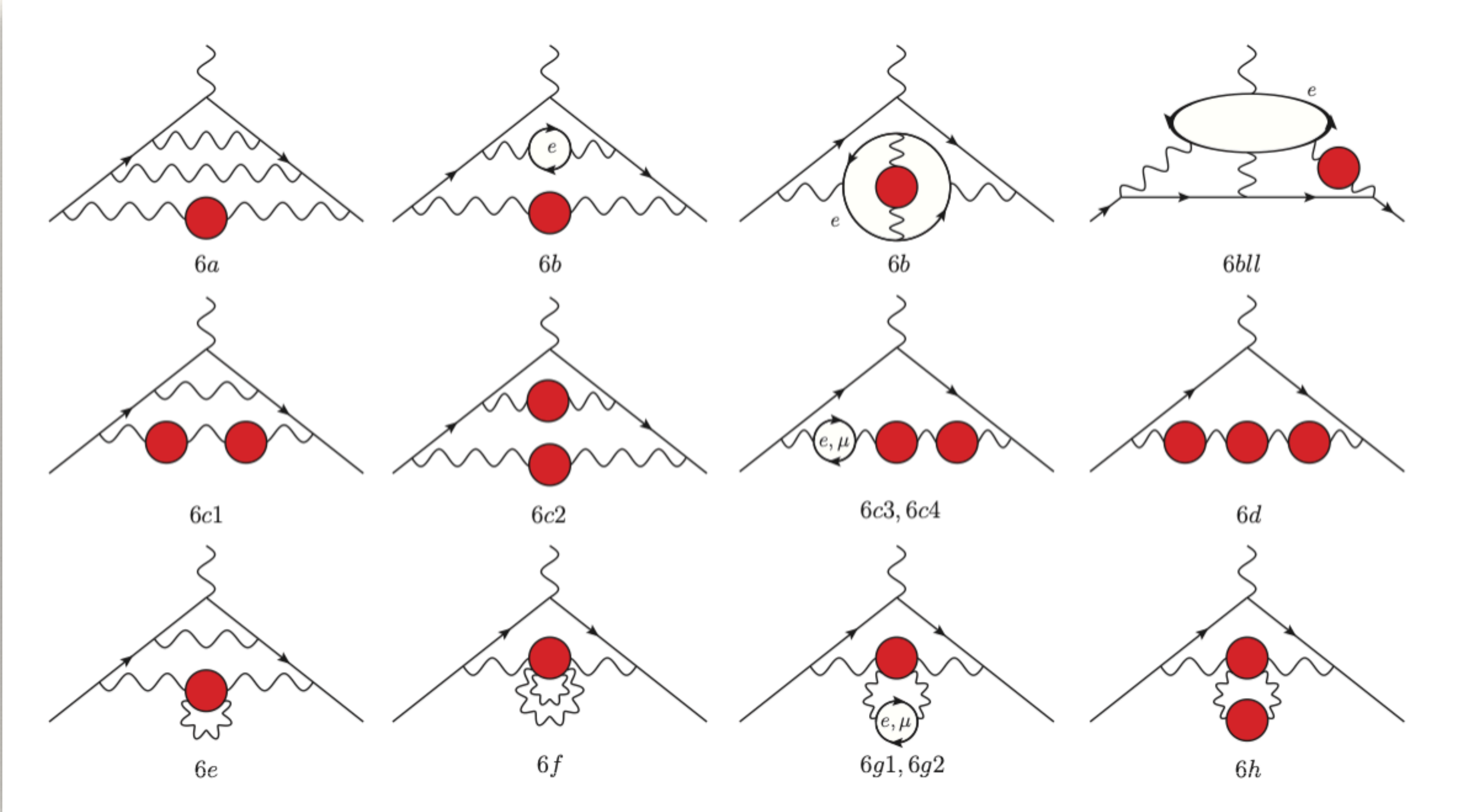


$$a_{\mu}^{HVP}(NLO) = -98.3(7) \times 10^{-11}$$

Krause '96; Keshevarzi, Nomura, Teubner 2019; WP20

The hadronic NNLO VP contribution

- $O(\alpha^4)$ contributions containing HVP insertion

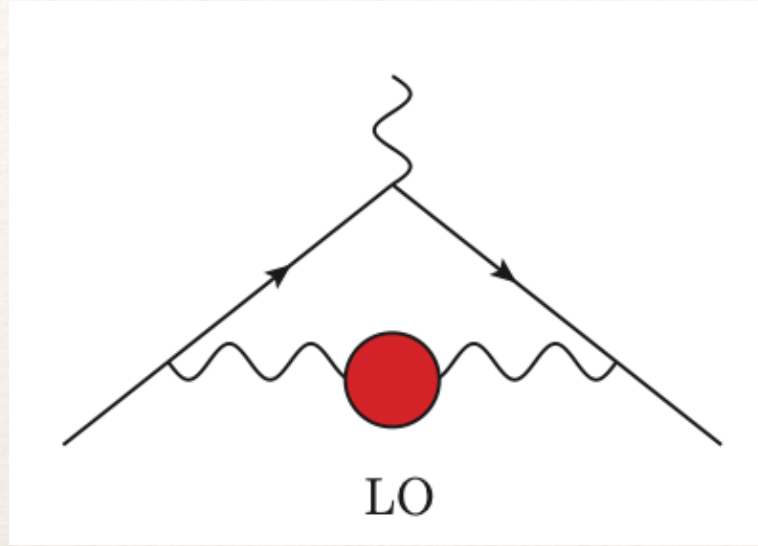


See Laporta's talk

$$a_{\mu}^{HVP}(NNLO) = 12.4(1) \times 10^{-11}$$

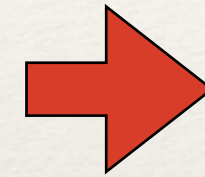
Kurz, Liu, Marquard, Steinhauser 2014

Space-like method: LO hadronic vacuum polarization contribution



$$K^{(2)}(z) = \frac{1}{\pi} \int_{-\infty}^0 dz' \frac{\text{Im}K^{(2)}(z')}{z' - z}, \quad z > 0$$

$$\frac{\Pi(q^2)}{q^2} = \frac{1}{\pi} \int \frac{ds}{s} \frac{\text{Im}\Pi(s)}{s - q^2}, \quad q^2 < 0$$



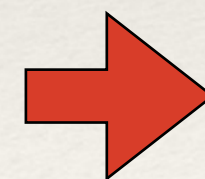
$$a_{\mu}^{HVP}(LO) = -\frac{\alpha}{\pi^2} \int_{-\infty}^0 \frac{dt}{t} \Pi_h(t) \text{Im}K^{(2)}(t/m^2)$$

With the imaginary part:

$$\text{Im}K^{(2)}(z + i\epsilon) = \pi\theta(-z) \left[\frac{z^2}{2} - z + \frac{z - 2z^2 + z^3/2}{\sqrt{z(z-4)}} \right] = \pi\theta(-z) F^{(2)}(1/y(z)), \quad F^{(2)}(u) = \frac{u+1}{u-1} u^2$$

Changing the variable in the dispersive integral $t \rightarrow y \rightarrow x$: $t(x) = \frac{m^2 x^2}{1-x^2}$, $x = 1 + y$

$$\Delta\alpha_h(t) = -\Pi_h(t)$$

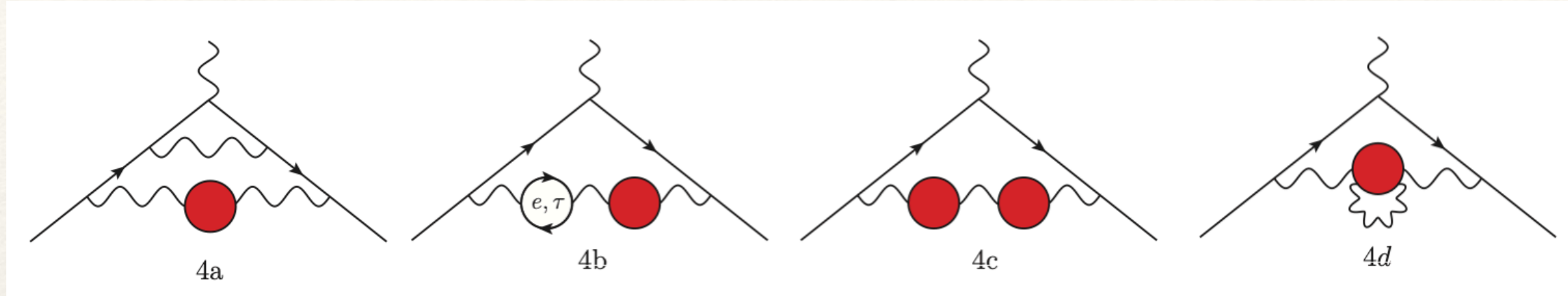


$$a_{\mu}^{HVP}(LO) = \frac{\alpha}{\pi} \int_0^1 dx \kappa^{(2)}(x) \Delta\alpha_h(t(x))$$

$$\kappa^{(2)}(x) = 1 - x$$

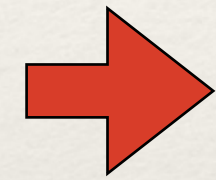
Lautrup, Peterman, de Rafael 1972

Space-like method: NLO HVP contribution



$$a_{\mu}^{(4a)} = \frac{\alpha^2}{\pi^3} \int_{s_0}^{\infty} \frac{ds}{s} 2K^{(4)}(s/m^2) \text{Im}\Pi_h(s)$$

$K^{(4)}$ from Barbieri Remiddi 1975



$$a_{\mu}^{(4a)} = -\frac{\alpha^2}{\pi^3} \int_{-\infty}^0 \frac{dt}{t} \Pi_h(t) \text{Im}K^{(4)}(t/m^2)$$

$$\begin{aligned} F^{(4)}(u) = & R_1(u) + R_2(u) \ln(-u) \\ & + R_3(u) \ln(1+u) + R_4(u) \ln(1-u) \\ & + R_5(u) [4 \text{Li}_2(u) + 2 \text{Li}_2(-u) \\ & + \ln(-u) \ln((1-u)^2(1+u))] \end{aligned}$$

$$\text{Im}K^{(4)}(z + i\epsilon) = \pi \theta(-z) F^{(4)}(1/y(z))$$

$$R_1 = \frac{23u^6 - 37u^5 + 124u^4 - 86u^3 - 57u^2 + 99u + 78}{72(u-1)^2 u(u+1)},$$

$$R_2 = \frac{12u^8 - 11u^7 - 78u^6 + 21u^5 + 4u^4 - 15u^3 + 13u + 6}{12(u-1)^3 u(u+1)^2},$$

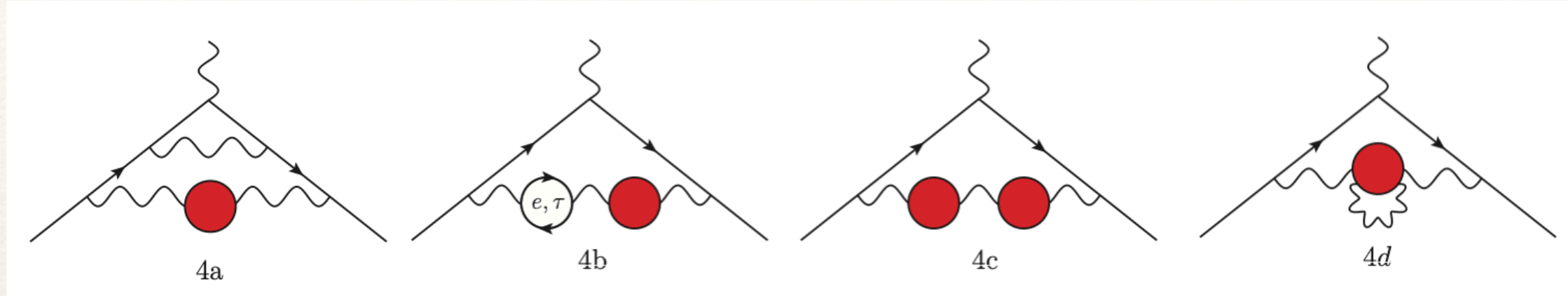
$$R_3 = \frac{(u+1)(-u^3 + 7u^2 + 8u + 6)}{12u^2},$$

$$R_4 = \frac{-7u^4 - 8u^3 + 8u + 7}{12u^2},$$

$$R_5 = -\frac{3u^4 + 5u^3 + 7u^2 + 5u + 3}{6u^2}.$$

Obtained independently by Nesterenko, arXiv: 2112.05009

Space-like method: NLO hadronic vacuum polarization contribution



$$\rightarrow a_{\mu}^{(4a)} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^1 dx \kappa^{(4)}(x) \Delta\alpha_h(t(x)),$$

$$\kappa^{(4)} = \frac{2(2-x)}{x(x-1)} F^{(4)}(x-1)$$

$$\rightarrow a_{\mu}^{(4b)} = \left(\frac{\alpha}{\pi}\right) \int_0^1 dx \kappa^{(2)}(x) \Delta\alpha_h(t(x)) \times 2 \left[\Delta\alpha_e^{(2)}(t(x)) + \Delta\alpha_{\tau}^{(2)}(t(x)) \right]$$

$$\rightarrow a_{\mu}^{(4c)} = \left(\frac{\alpha}{\pi}\right) \int_0^1 dx \kappa^{(2)}(x) [\Delta\alpha_h(t(x))]^2$$

Where $\Pi_{\ell}^{(2)}(t) = -\Delta\alpha_{\ell}^{(2)}(t)$ and $\Pi_{\ell}^{(2)}(t) = \frac{\alpha}{\pi} \left(\frac{8}{9} - \frac{\beta_{\ell}^2}{3} + \beta_{\ell} \left(\frac{1}{2} - \frac{\beta_{\ell}^2}{6} \right) \ln \frac{\beta_{\ell} - 1}{\beta_{\ell} + 1} \right)$, with $\beta_{\ell} = \sqrt{1 - 4m_{\ell}^2/t}$.

Space-like method: NLO hadronic vacuum polarization contribution

Chakraborty et al. (arXiv: 1806.08190) provided an approximated expression for the space-like kernel function. They started from :

$$\begin{aligned}
 K^{(4)}(s) = & 2 \frac{m^2}{s} \left\{ \left[\frac{223}{54} - 2\zeta(2) - \frac{23}{36} \ln \frac{s}{m^2} \right] \right. \\
 & + \frac{m^2}{s} \left[\frac{8785}{1152} - \frac{37}{8} \zeta(2) - \frac{367}{216} \ln \frac{s}{m^2} + \frac{19}{144} \ln^2 \frac{s}{m^2} \right] \\
 & + \frac{m^4}{s^2} \left[\frac{13072841}{432000} - \frac{883}{40} \zeta(2) - \frac{10079}{3600} \ln \frac{s}{m^2} + \frac{141}{80} \ln^2 \frac{s}{m^2} \right] \\
 & \left. + \frac{m^6}{s^3} \left[\frac{2034703}{16000} - \frac{3903}{40} \zeta(2) - \frac{6517}{1800} \ln \frac{s}{m^2} + \frac{961}{80} \ln^2 \frac{s}{m^2} \right] \right\}
 \end{aligned}$$

With $r = \frac{m^2}{s}$

Krause '97

They exploited generating integral representation to fit the r^n and $r^n \ln r$, but not the $r^n \ln^2 r$ ones.

$$m^2 \int_0^1 \frac{dx P(x)}{m^2 x + s} = \frac{m^2}{s} \sum_n a_n \left(\frac{m^2}{s} \right)^n$$

for the power terms

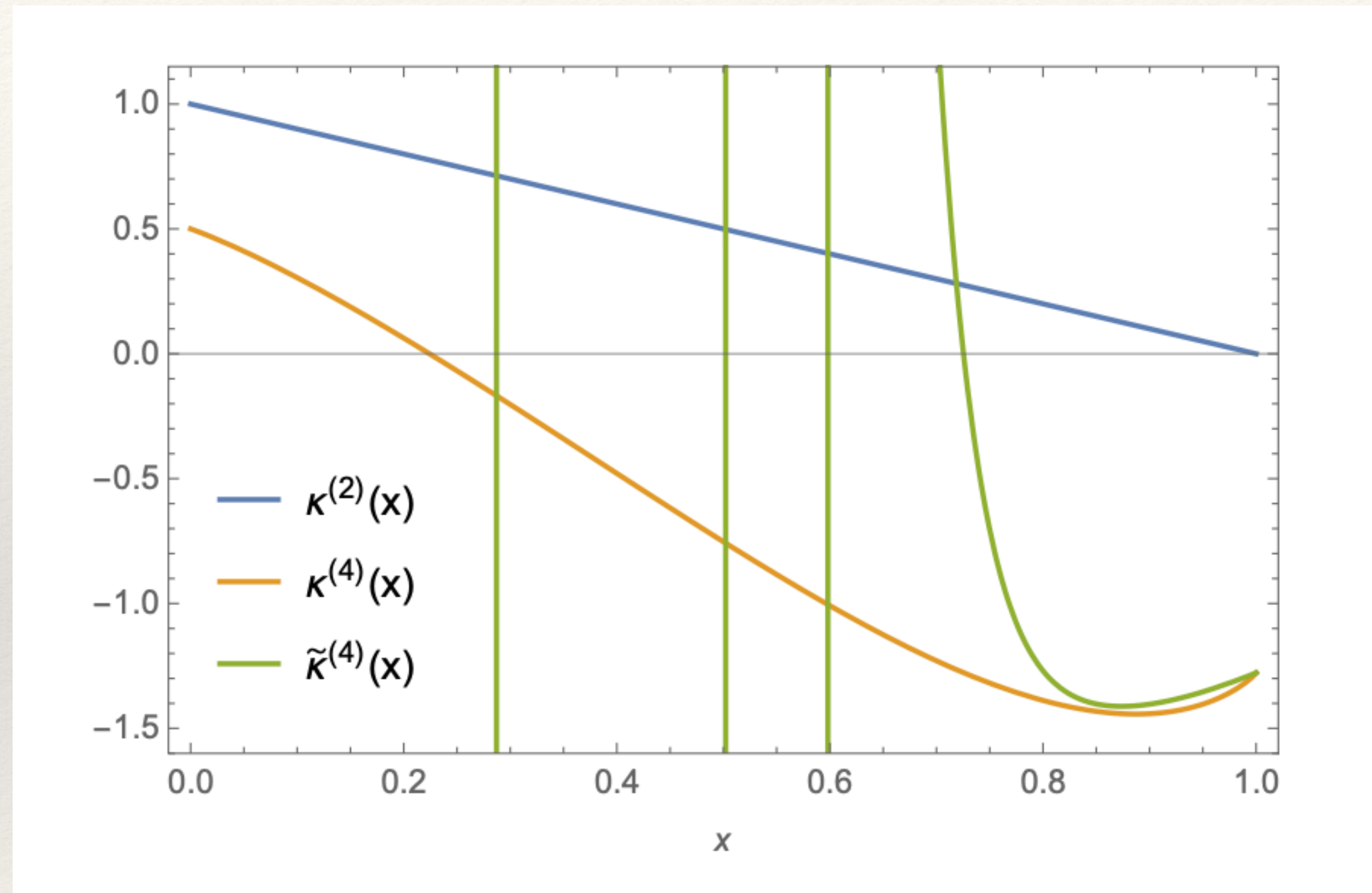
$$m^2 \int_0^1 \frac{dx G(x)}{sx + m^2} = G_1(m^2/s) + G_2(m^2/s) \ln \left(\frac{s}{m^2} \right)$$

for the log terms

Groote et al. 2002

Space-like method: NLO hadronic vacuum polarization contribution

After simple changes of variables their approximation can be compared to our exact function $\kappa^{(4)}(x)$.

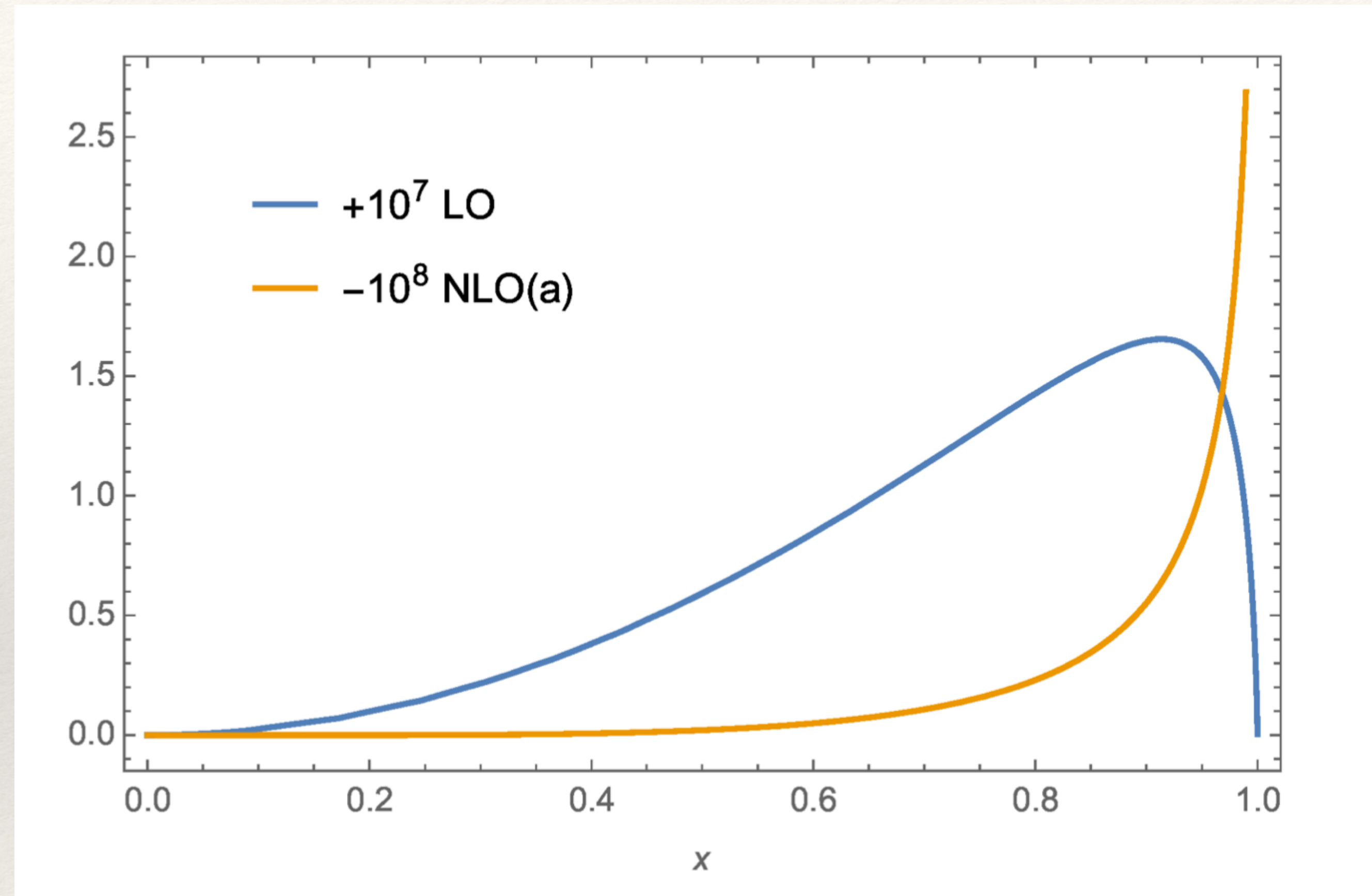


The authors added a $O(10\%)$ uncertainties to their final result.



This uncertainty can be eliminated using our exact formula $\kappa^{(4)}(x)$.

Space-like method: NLO hadronic vacuum polarization contribution



- The LO integrand has a peak at $x \sim 0.914$
- The NLO integrand has an integrable logarithmic singularity at $x \rightarrow 1$

Conclusions

- ❖ We provide simple analytic expressions to calculate the HVP contributions to the muon $g - 2$ in the space-like region up to NLO.
- ❖ These results can be employed at MUonE, as well as in lattice QCD computations.
- ❖ An existing approximation for the NLO kernel induced large uncertainties. These can be eliminated using the exact expression $\kappa^{(4)}(x)$.
- ❖ Calculation of higher-order HVP corrections require a precise treatment of the QED radiative corrections to the HVP. MUonE will naturally include all of these corrections in the space-like approach.
- ❖ These results allow to compare time-like and space-like calculations of $a_\mu^{(HVP)}$ at NLO accuracy.

Thank you for your attention!