

# Radiative corrections to $\mu e$ scattering in MESMER

MITP, Mainz

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# Overview

- Photonic NNLO contributions
  - MESMER
  - NNLO lepton pair contributions
  - $\pi^0$  production
- } Ettore Budassi
- } Clara Lavinia del Pio

Numerical results for  $\mu e$  scattering

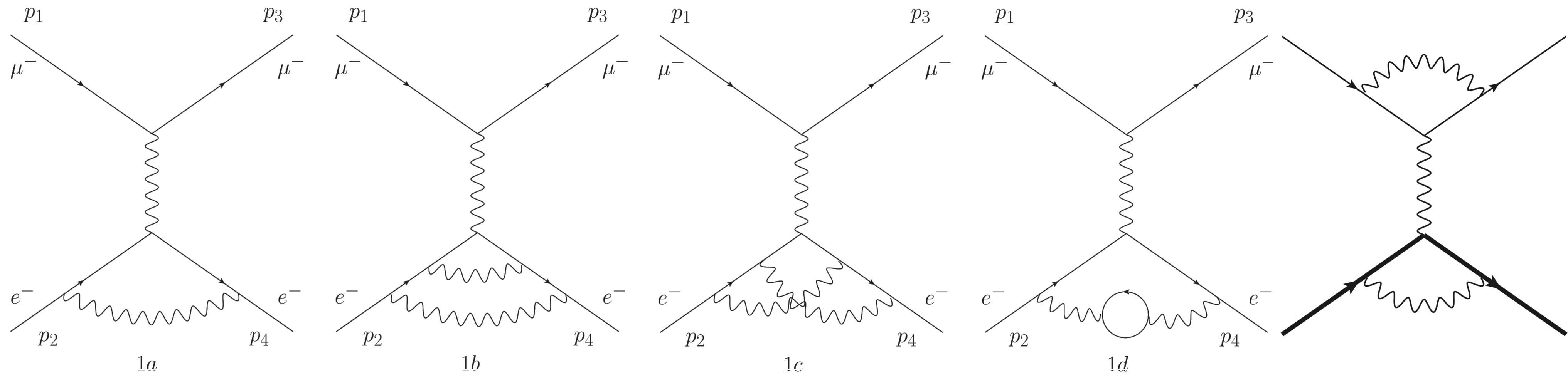
# $\mu e$ scattering at NLO

- Both real and virtual QED corrections to  $\mu e$  scattering are included.
- Weak effects are known, at the level of  $\sim 10^{-5}$  (LO) and  $\lesssim 10^{-6}$  (NLO).
- Full dependence on masses and Radiative Corrections have been studied with **specific kinematical cuts** that mimic the experimental setup of MUonE:
  - **Basic acceptance cuts:**  $\vartheta_e, \vartheta_\mu < 100$  mrad and  $E_e > 1$  GeV;
  - **Acoplanarity cut:**  $\xi = \left| \pi - |\phi_e - \phi_\mu| \right| < \xi_c = 3.5$  mrad.

# NNLO Photonic Corrections

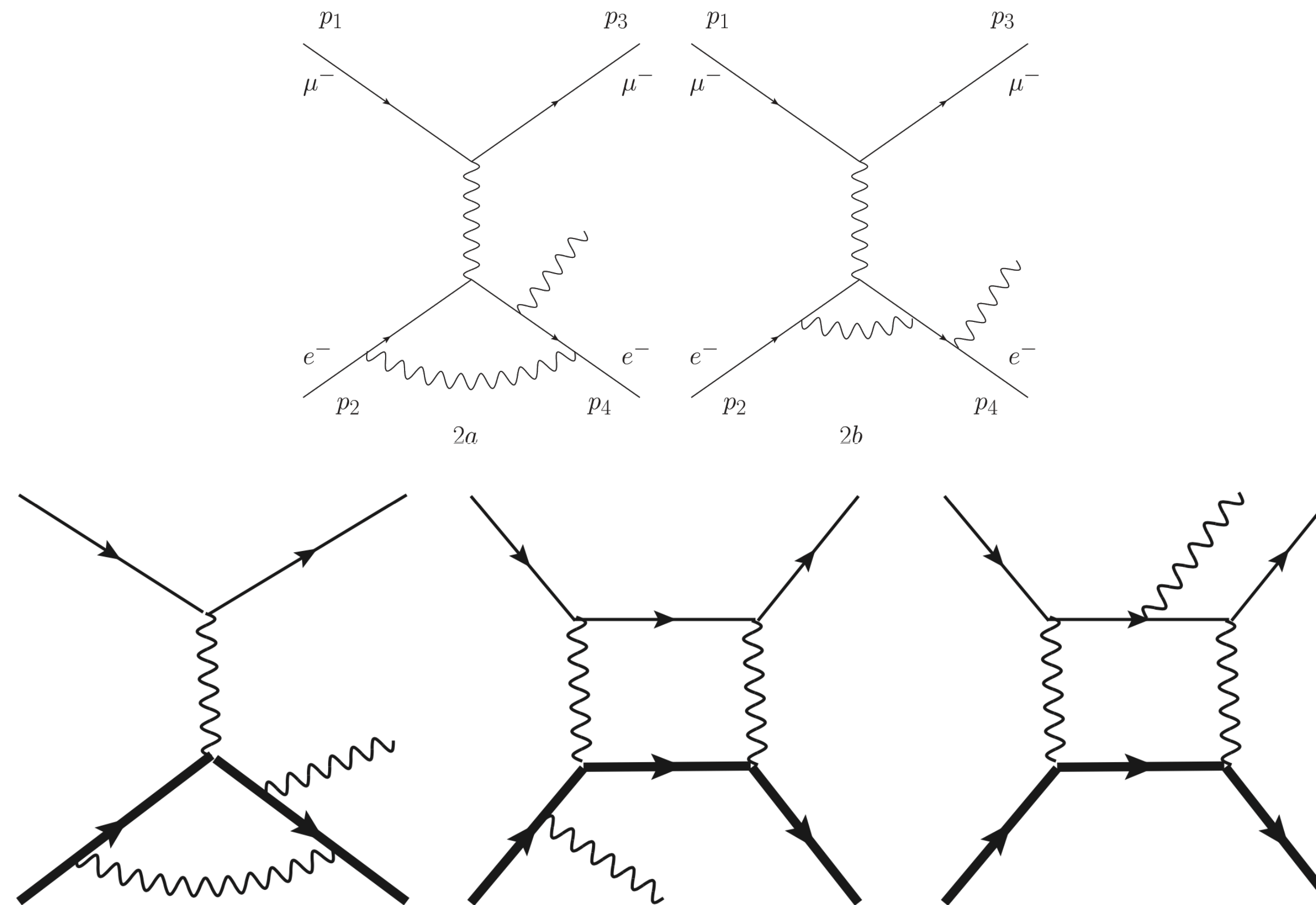
# NNLO Photonic Corrections: exact contributions

- $|\text{NLO virtual diagrams}|^2$
- Virtual NNLO photonic contributions are included exactly for electron or muon leg emission. 2-loop QED vertex from factors taken from Mastrolia and Remiddi.

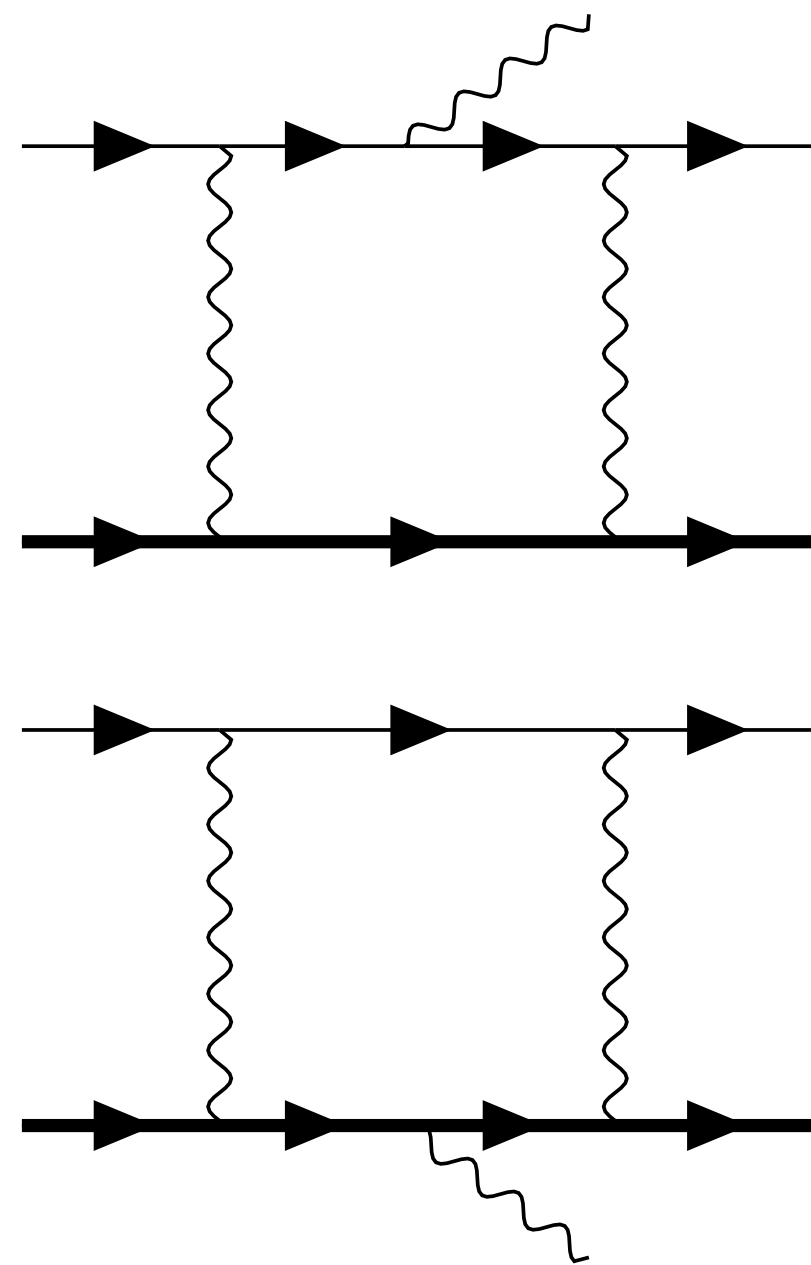
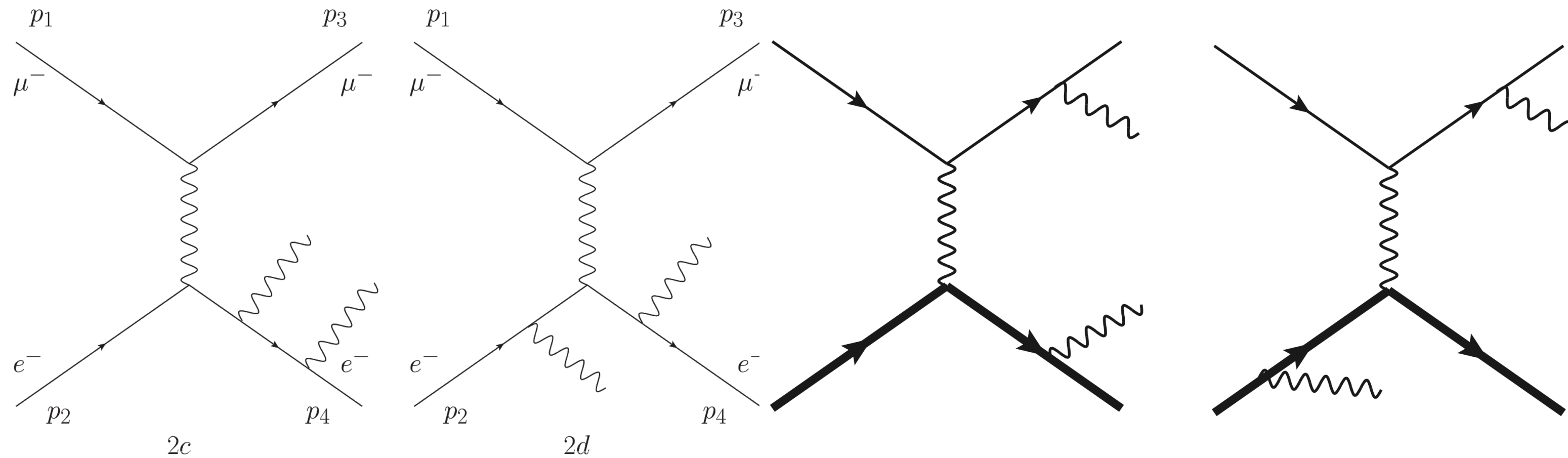


# NNLO Photonic Corrections: exact contributions

- Interference of LO  $\mu e \rightarrow \mu e \gamma$  amplitude with



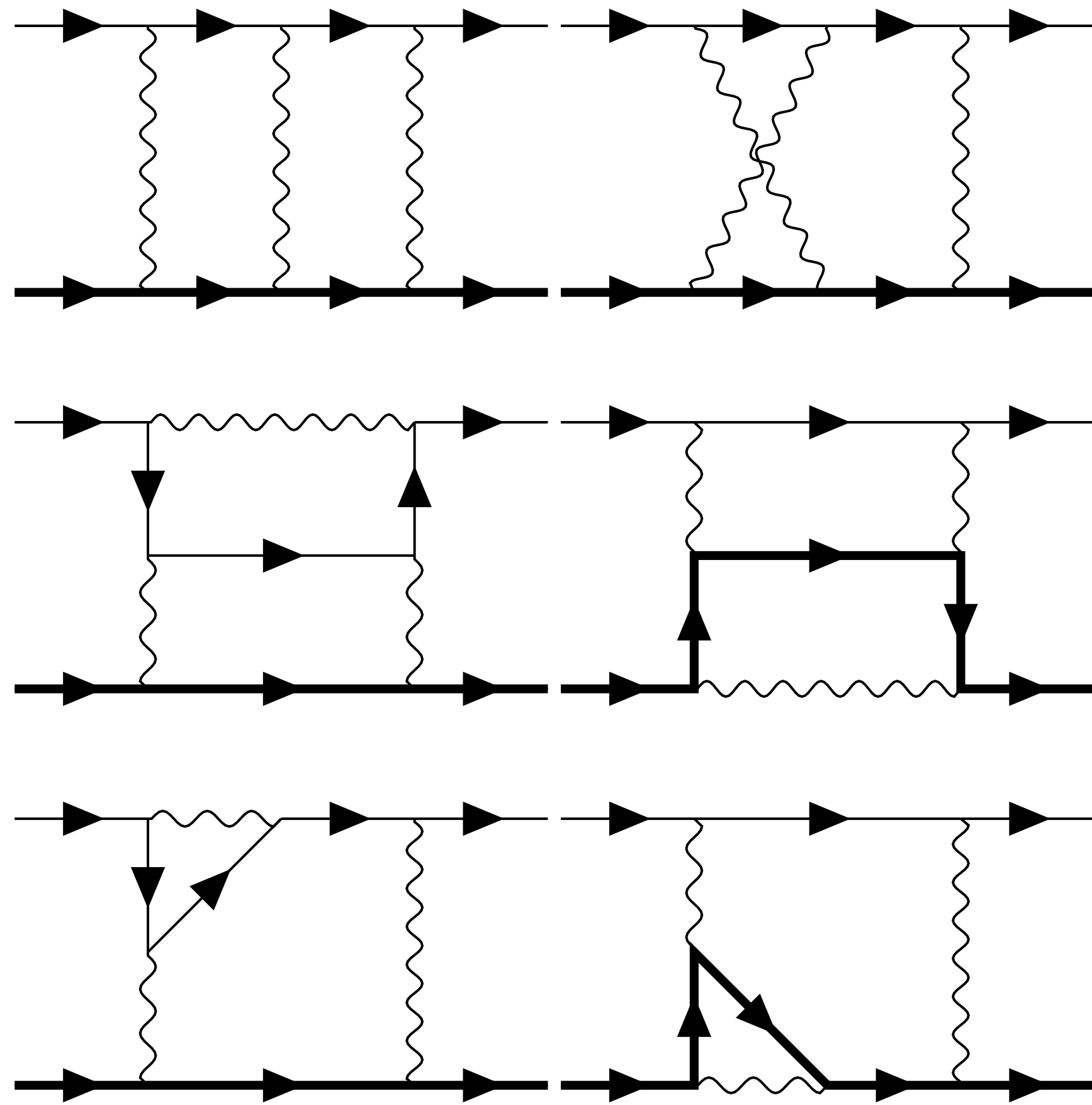
# NNLO Photonic Corrections: exact contributions



- 1-loop corrections to real photon emission exactly included: *e.g.* pentagon diagrams.
- Double real emission included exactly;
- We use a small  $m_\gamma = \lambda$  and the slicing method;
- The gauge-invariant subset of NNLO corrections along the electron line are cross-checked against McMule.



# NNLO Photonic Corrections: approximated contributions



- Of the two-loop virtual diagrams with a virtual photon insertion on top of NLO boxes, only the IR part is included exactly (YFS).
- The non-IR remnants are approximate.
- All photonic NNLO effects weigh at most some % at the Phase Space boundaries.
- Work is in progress for the full  $\mu e \rightarrow \mu e$  at NNLO (Padova&PSI).

# NNLO Photonic Corrections: in formulas...

$$\widetilde{\mathcal{M}}^{\alpha^2} = \underbrace{\mathcal{M}_e^{\alpha^2} + \mathcal{M}_\mu^{\alpha^2} + \mathcal{M}_{e\mu, 1L \times 1L}^{\alpha^2}}_{\text{Exact}} + \underbrace{\frac{1}{2} Y_{e\mu}^2 \mathcal{T} + Y_{e\mu} (Y_e + Y_\mu) \mathcal{T} + (Y_e + Y_\mu) \mathcal{M}_{e\mu}^{\alpha^1, R} + Y_{e\mu} M^{\alpha^1, R}}_{\text{YFS approximated}}$$

Exact

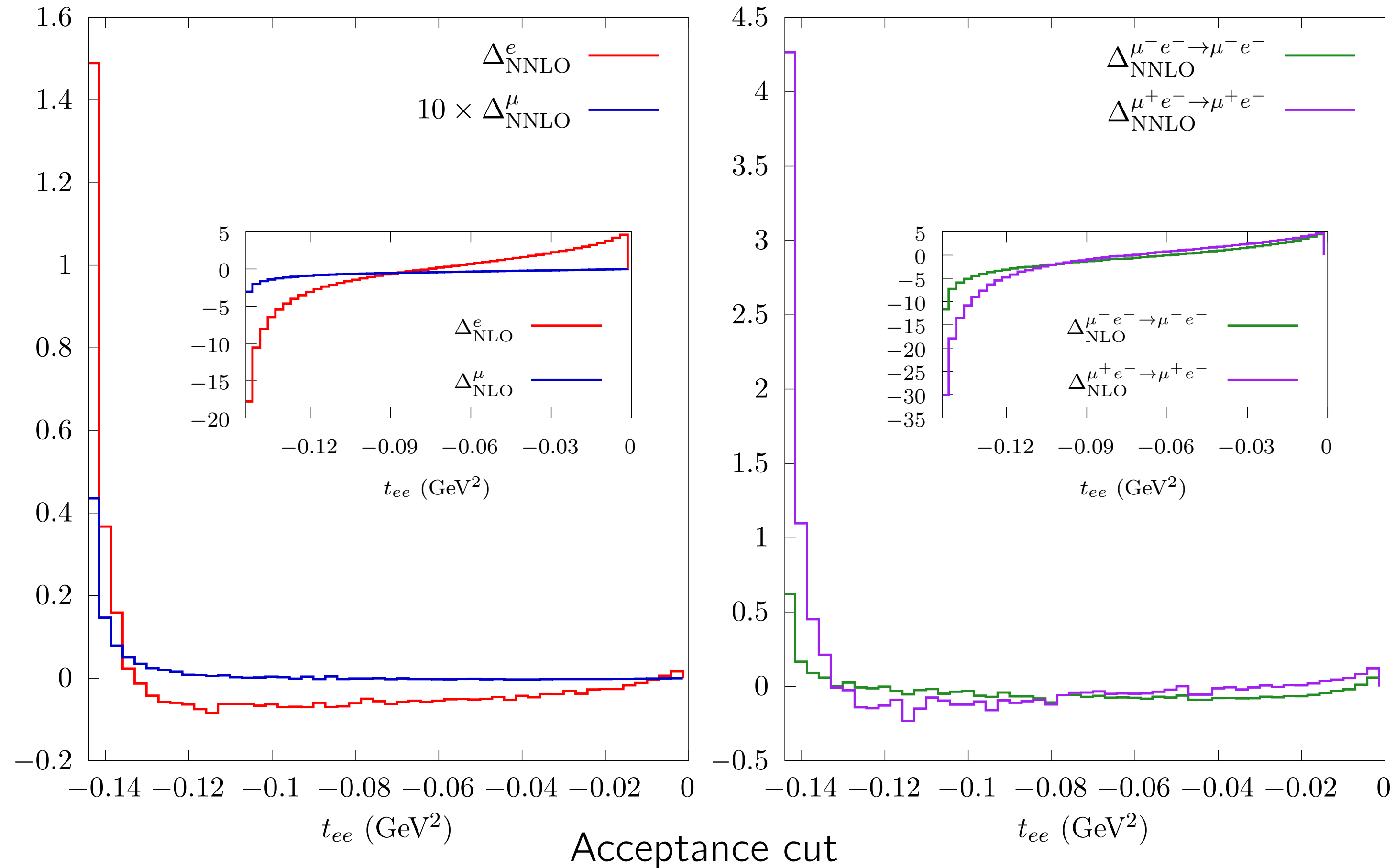
YFS approximated

- $\mathcal{M}_{e(\mu)}^{\alpha^2}$ : two virtual  $\gamma$  are both attached to the electron (muon) leg;
- $\mathcal{M}_{e\mu, 1L \times 1L}^{\alpha^2}$ : 1 virtual  $\gamma$  is attached to each leg.

- This contains all the IR part.
- It misses the non-IR remnant of the two-loop diagrams with two  $\gamma$ s connecting the  $e$  and  $\mu$  legs  $\mathcal{M}^{\alpha^2, R}$ .

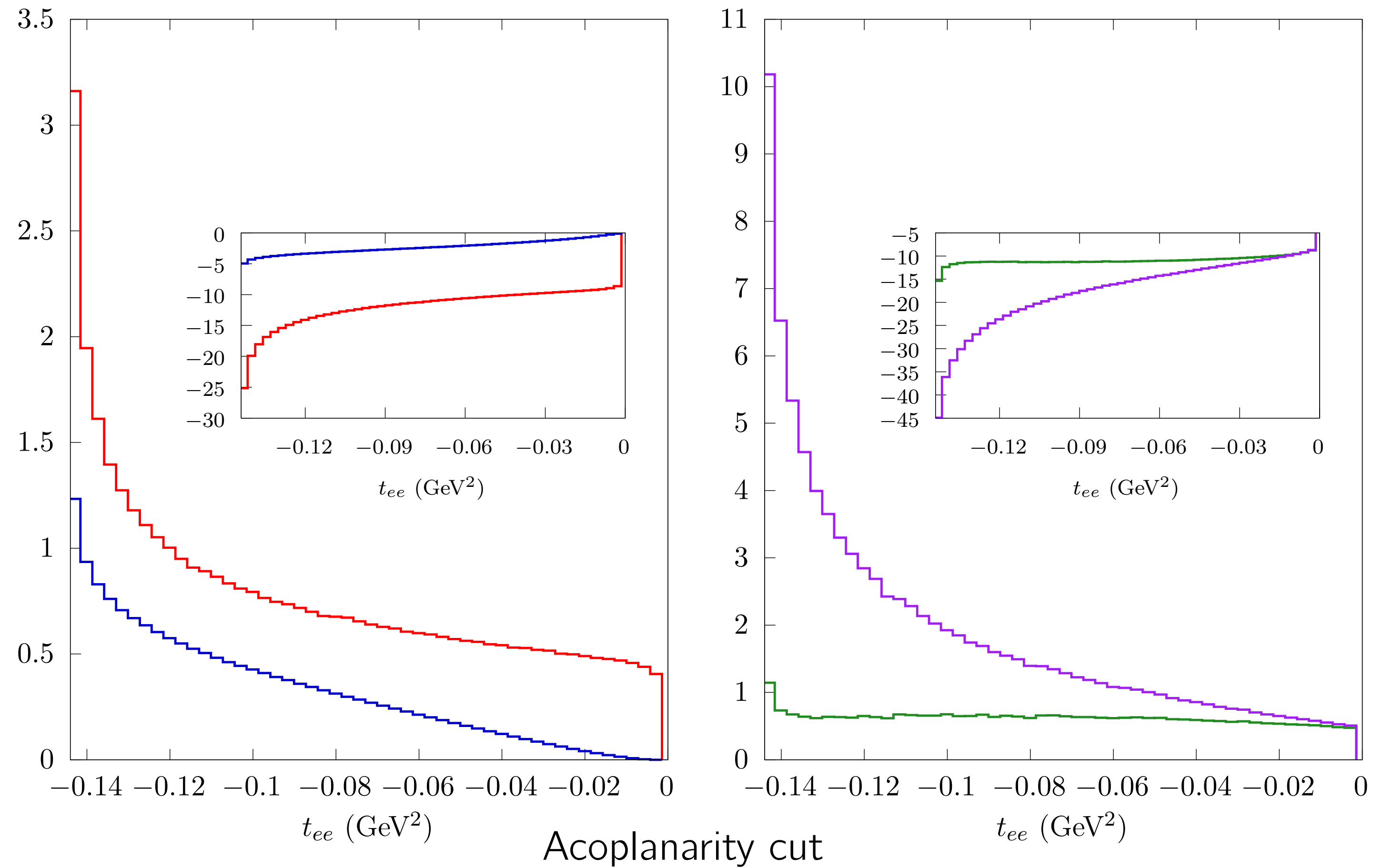
$$Y_{ij} = \begin{cases} \frac{1}{8} \frac{\alpha}{\pi} Q_i^2 [B_0(0, m_i^2, m_i^2) - 4m_i^2 C_0(m_i^2, 0, m_i^2, \lambda^2, m_i^2, m_i^2)] & \text{for } i = j \\ \frac{\alpha}{\pi} Q_i Q_j \vartheta_i \vartheta_j [p_i \cdot p_j C_0(m_i^2, (\vartheta_i p_i + \vartheta_j p_j)^2, m_j^2, \lambda^2, m_i^2, m_j^2) + \frac{1}{4} B_0((\vartheta_i p_i + \vartheta_j p_j)^2, m_i^2, m_j^2)] & \text{for } i \neq j \end{cases}$$

# NNLO Photonic Corrections: results



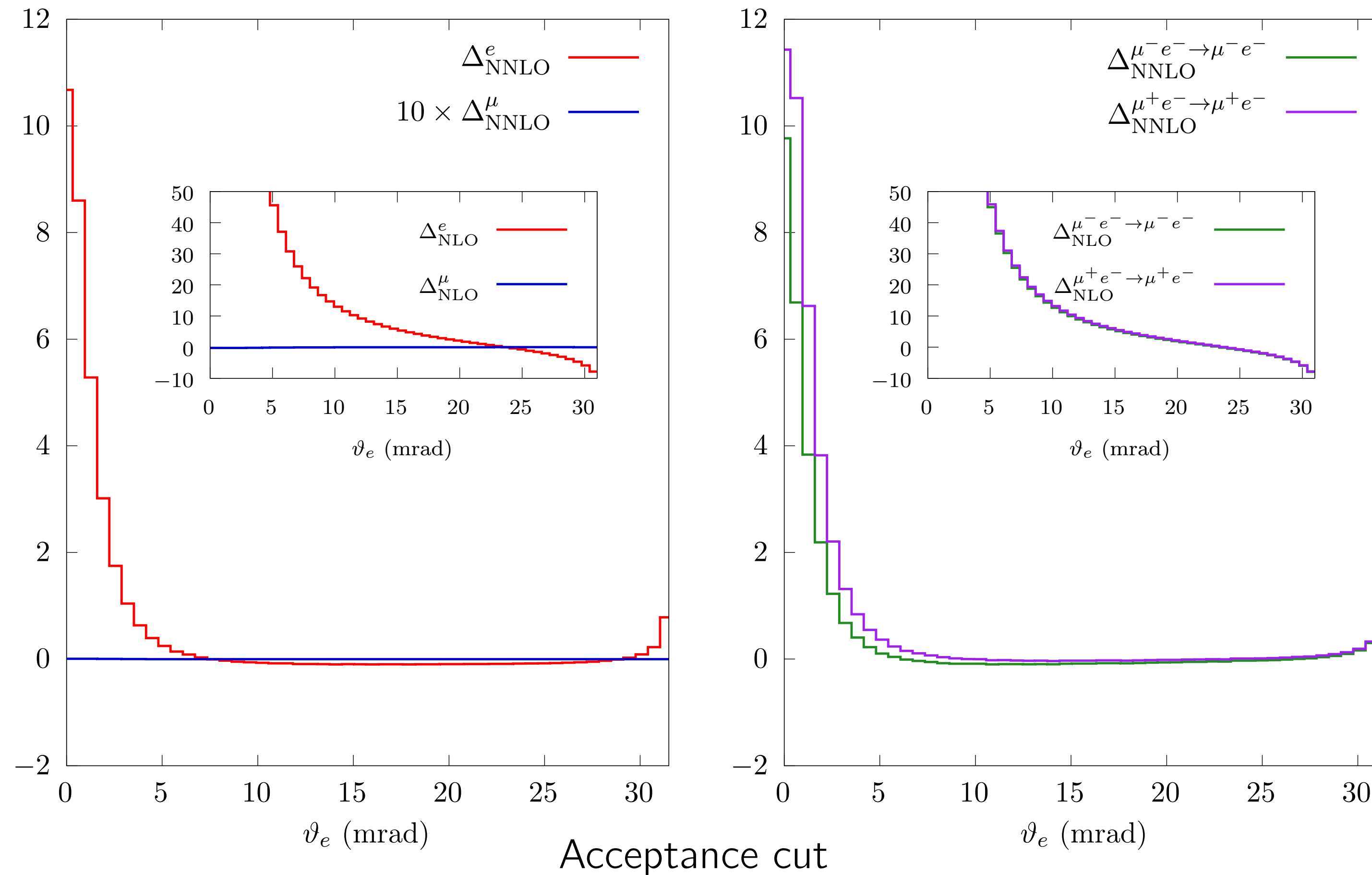
$$\Delta_i^{\text{NNLO}} = \frac{d\sigma_i^{\text{NNLO}} - d\sigma_i^{\text{NLO}}}{d\sigma_i^{\text{LO}}} \times 100$$

# NNLO Photonic Corrections: results



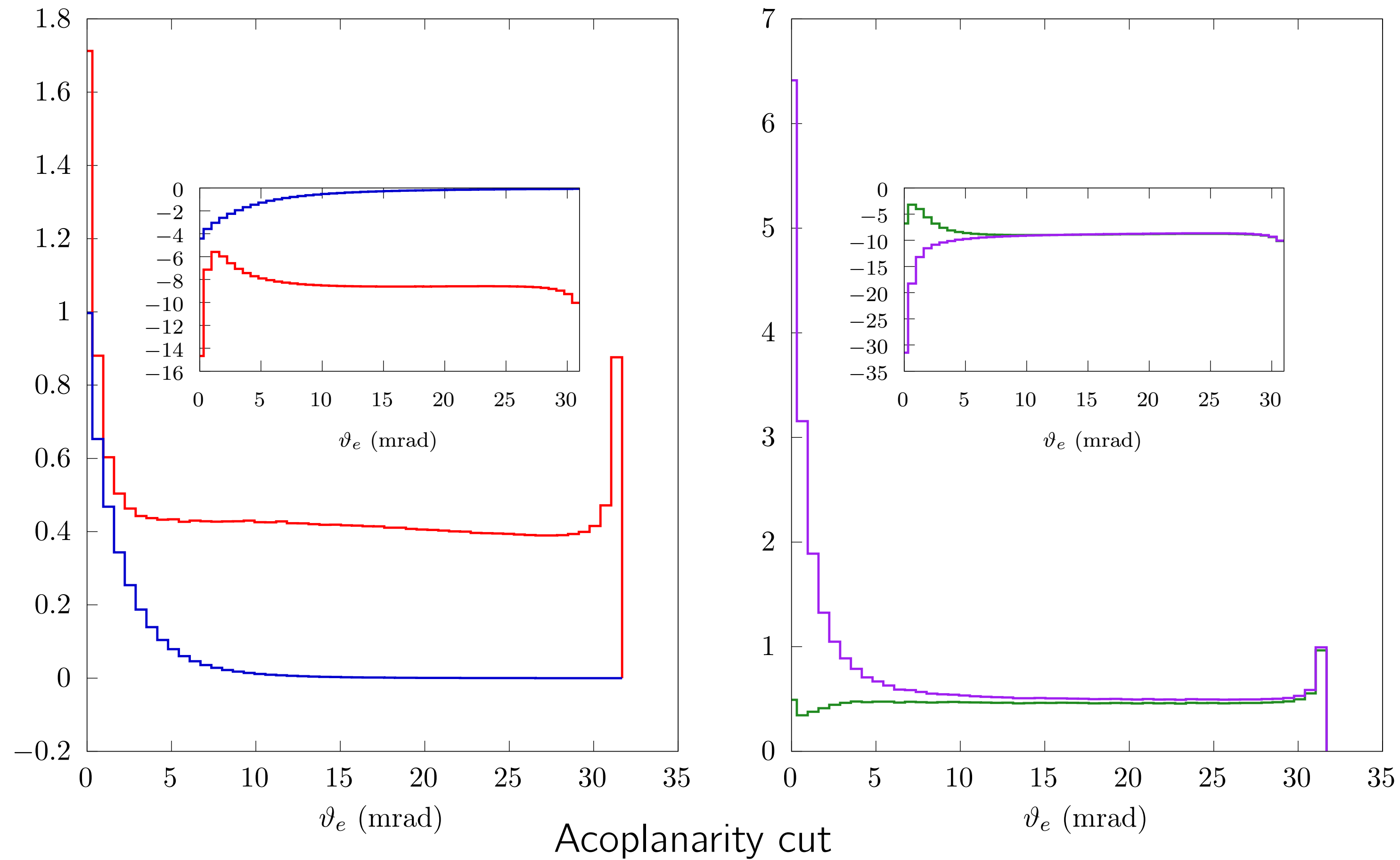
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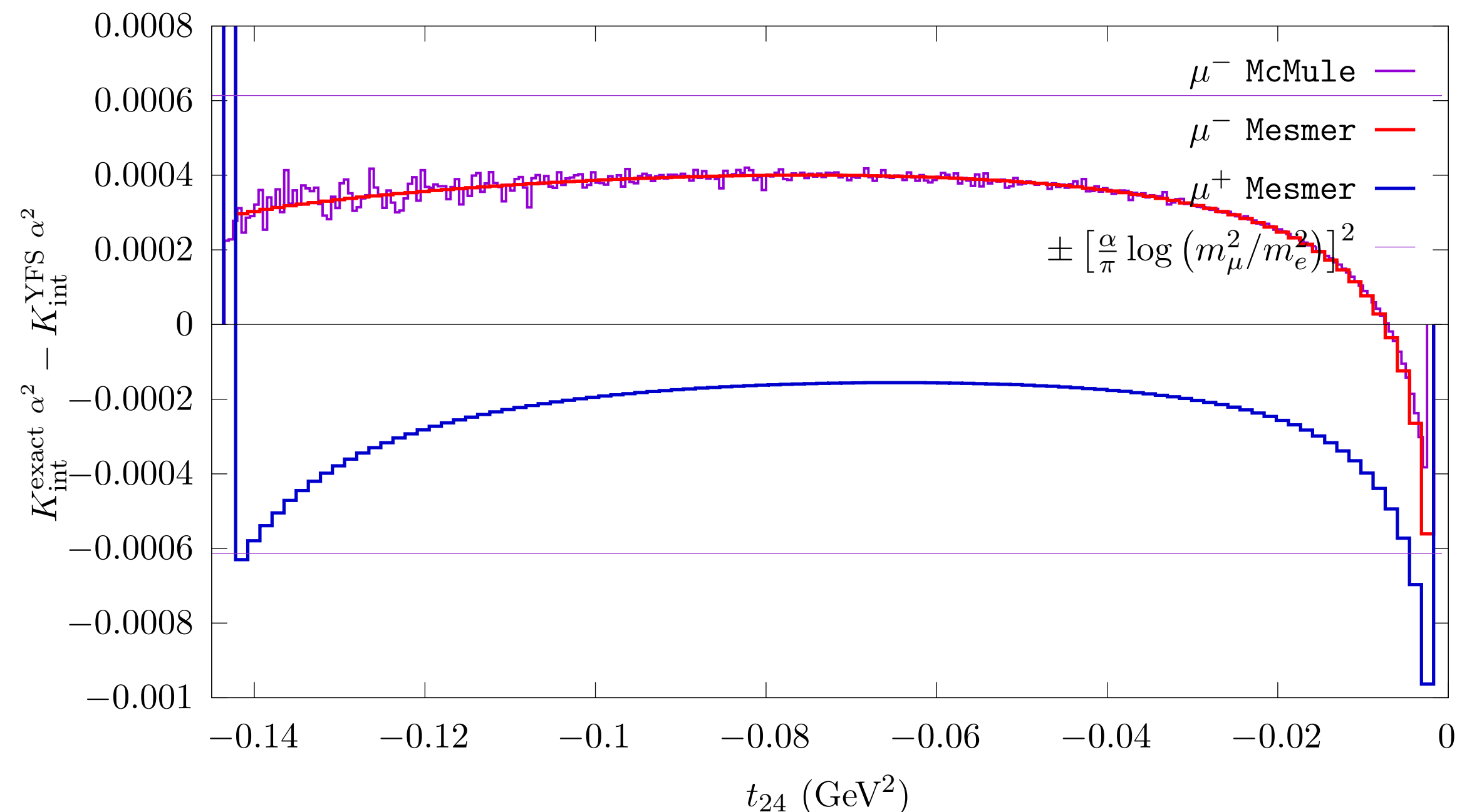
# NNLO Photonic Corrections: results



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# Exact NNLO photonic corrections

- The complete two-loop corrections to  $f\bar{f} \rightarrow F\bar{F}$  have been calculated by Bonciani *et. al* for  $m_f = 0$ . Using **crossing symmetry** it can be used for  $\mu e \rightarrow \mu e$ ;
- The amplitudes with  $m_e = 0$  can undergo the **massification** procedure, to get the collinear divergences in terms of  $\ln(Q^2/m_e)$ ;
- NNLO double boxes are CPU expensive ( $>1$  s/event on a single core).



Difference between YFS-approximated and exact NNLO photonic K factor.

R. Bonciani *et. al*, PRL 128 (2022) 2.  
T. Engel *et al*. JHEP 02 (2019) 118

**PRELIMINARY!**

MESMER



# Relevant references

- Alacevich *et al.*, Muon-electron scattering at NLO, [JHEP 02 \(2019\) 155](#)
- Carloni Calame *et al.*, Towards muon-electron scattering at NNLO, [JHEP 11 \(2020\) 028](#)
- E. Budassi *et al.*, NNLO virtual and real leptonic corrections to muon-electron scattering, [JHEP 11 \(2021\) 098](#)
- E. Budassi *et al.*, Single  $\pi^0$  production in  $\mu e$  scattering at MUonE, [PLB 829 \(2022\) 137138](#)

# What can MESMER calculate?

- NLO QED corrections: single real or virtual  $\gamma$ ;
- NNLO photonic corrections (approximate);
- Exact NNLO lepton pair contributions and  $\mu e \rightarrow \mu e l^+ l^-$ , with  $l = e, \mu$  (not public yet);
- $\mu e \rightarrow \mu e \pi^0$ , with  $\pi^0 \rightarrow \gamma\gamma$  (not public yet).

# MESMER

MESMER (**M**uon **E**lectron **S**cattering with **M**ultiple  
**E**lectromagnetic **R**adiation)

Fully differential Monte Carlo event generator for high-precision simulation of  $\mu e$  scattering at low energies, developed for the MUonE experiment.

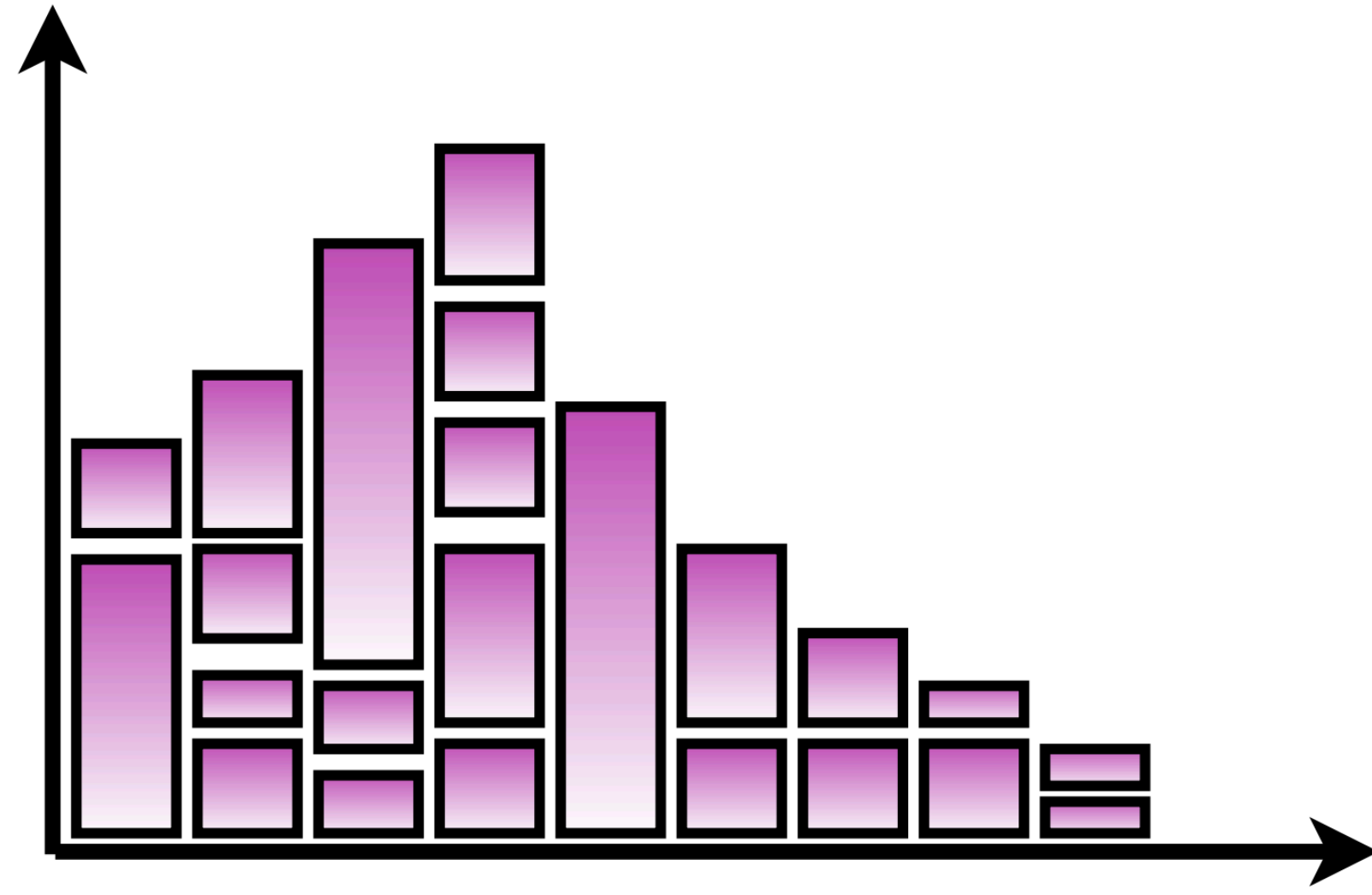
# The code

- The code can be found in [this github](#) repository.
- MESMER is mostly written in Fortran 77 language.
- Some external libraries are used for one-loop integrals and pseudo-random number generation;
- HVP is included using Jegerlehner's, Keshavarzi-Nomura-Teubner's and Ignatov's routines.

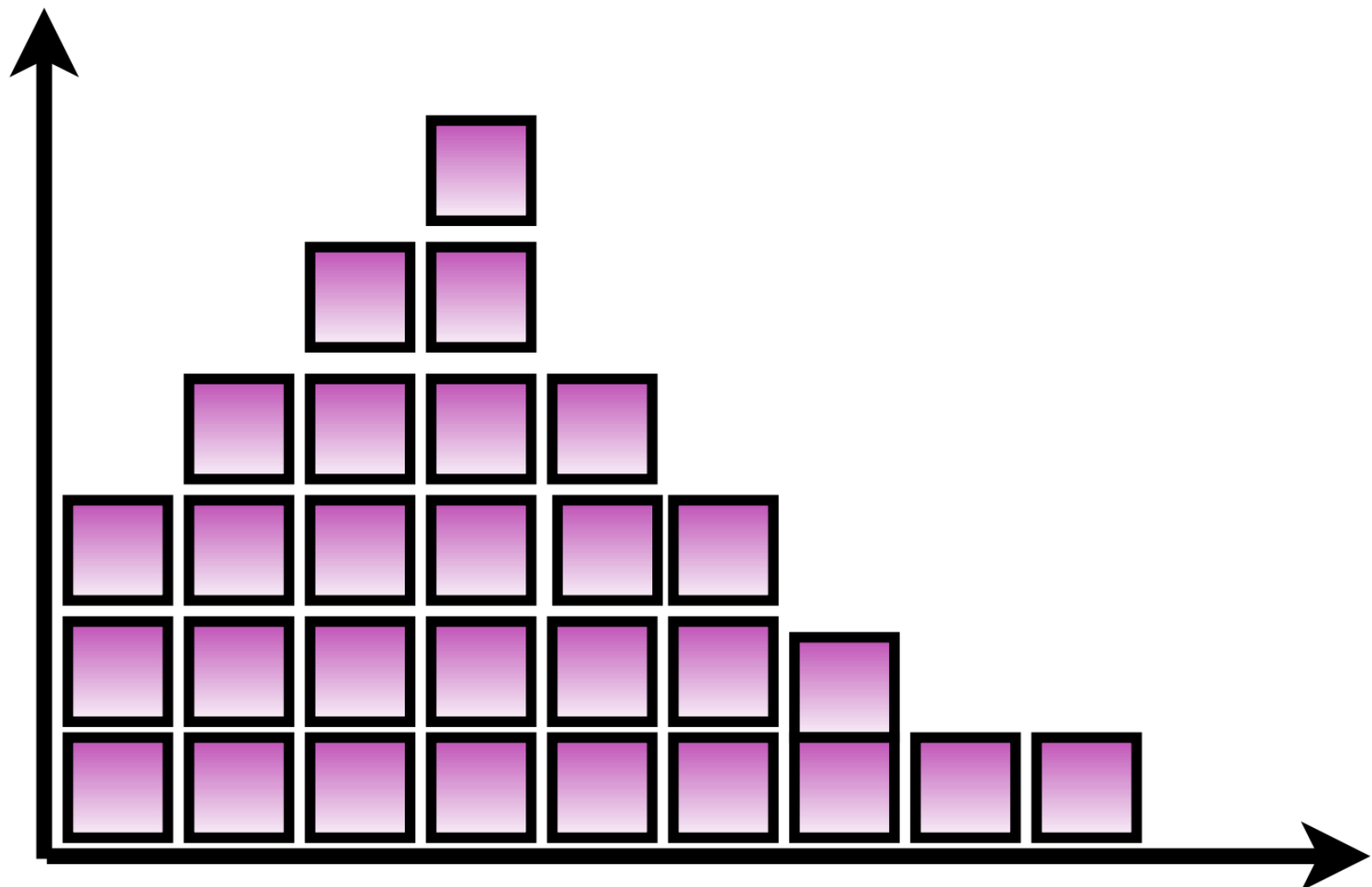
# Parameters & output

- Some kinematical constraints are imposed at generation level. `cuts.F` is the routine that selects kinematical cuts and can be modified according to the user's needs;
- A generic incoming muon momentum can be fed in input event by event, for realistic simulations (beam profile);
- Events are stored in a format that has been thought out with the experimental colleagues (weights, momenta *etc.*);
- Differential distributions are written in specific files, ready to be plotted (*e.g.* with gnuplot).
- A C/C++ interface is provided to use the code as a library in a driver simulation program.

# Weighting vs. Un-weighting



- Events have their weight;
- It needs to be carried throughout the whole detector simulation;
- Fast generation.



- Events are not weighted: they are distributed according to the cross section;
- Slow due to the un-weighting procedure.
- It is tricky to guess a correct maximiser for weights before the generation.

# Reweighting

$$\sigma = \sum_{i=0}^n \frac{w_i}{n} \quad \sigma_{reweighted} = \sum_{i=0}^n \frac{w_i}{n} \times r_i; \quad r_i = \frac{w_i^{rew}}{w_i}$$

- In the same run we can calculate different re-weighting coefficients  $r_i$  to study the inclusion/exclusion of different contributions (e.g. VP, HVP parametrisations *etc.*);
- No need to re-generate another MC sample to account for different effects!
- Since the MC sample is the same, one can exploit the statistical correlation on the weights to reduce errors.

Now it's Clara's turn.