MUonE: theory status

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The Evaluation of the Leading Hadronic Contribution To the Muon g-2: Towards the MUonE Experiment

MITP, Mainz, 14-18 November 2022

MUonE: theory status



 ★ G. Abbiendi, C.M. Carloni Calame, U. Marconi, C. Matteuzzi, G. Montagna, O. Nicrosini, M. Passera, F. Piccinini, R. Tenchini, L. Trentadue, G. Venanzoni, *Measuring the leading hadronic contribution to the muon g-2 via μe scattering* Eur. Phys. J. C **77** (2017) no.3, 139 - arXiv:1609.08987 [hep-ph]

* C. M. Carloni Calame, M. Passera, L. Trentadue and G. Venanzoni,

A new approach to evaluate the leading hadronic corrections to the muon g-2

Phys. Lett. B 746 (2015) 325 - arXiv:1504.02228 [hep-ph]

$$a_{\mu}^{\text{HLO}} = \frac{\alpha}{\pi} \int_{0}^{1} dx (1-x) \Delta \alpha_{\text{had}}[t(x)]$$

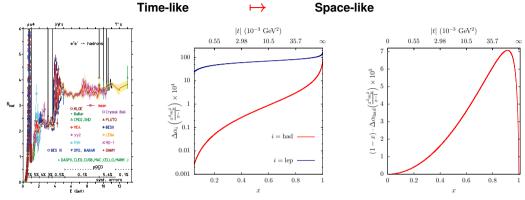
$$t(x) = \frac{x^{2} m_{\mu}^{2}}{x-1} < 0$$

$$e a \text{ Lautrup, Peterman, De Bafael, Phys. Bept. 3 (1972) 193}$$

- The hadronic VP correction to the running of α enters
- * $\Delta \alpha_{had}(t)$ can be directly measured in a (single) experiment involving a space-like scattering process and a_{μ}^{HLO} obtained through numerical integration Carloni Calame. Passera. Trentadue. Venanzoni PLB 746 (2015) 325
- \star A data-driven evaluation of a_{μ}^{HLO} , but with space-like data
- By modifying the kernel function $\frac{\alpha}{\pi}(1-x)$, also a_{μ}^{HNLO} and a_{μ}^{HNNLO} can be provided

Balzani, Laporta, Passera, Phys.Lett.B834 (2022) 137462 talks by E. Balzani and S. Laporta

From time-like to space-like evaluation of $a_{\mu}^{\rm HLO}$



Smooth function

- \mapsto Time-like: combination of many experimental data sets, control of RCs better than O(1%) on hadronic channels required
- → Space-like: in principle, one single experiment, it's a one-loop effect, very high accuracy needed

Abbiendi et al., EPJC 77 (2017) 3, 139

Abbiendi et al., Letter of Intent: the MUonE project, CERN-SPSC-2019-026, SPSC-I-252 (2019)

- Scattering μ's on e's in a low Z target looks like an ideal process (fixed target experiment)
- → It is a pure *t*-channel process at tree level
- \rightarrow The M2 muon beam ($E_{\mu} \simeq 160$ GeV) is available at CERN
- \checkmark $\sqrt{s} \simeq 0.4$ GeV and -0.143 < t < 0 GeV² (and extrapolate to $x \to 1$)
- ightarrow With ~ 3 years of data taking, a statistical accuracy of 0.35% on $a_\mu^{ extsf{HLO}}$ can be achieved talk by G. Abbiendi

$$rac{1}{2}rac{\delta\sigma}{\sigma}\simeqrac{\deltalpha}{lpha}\simeq\delta\Deltalpha_{\mathsf{had}}$$

 $\Delta \alpha_{had}$ is a 0.1% effect in this region \rightarrow to measure it at 1%, σ must be controlled at the 10^{-5} level

experimental

⇒ talks by G. Abbiendi, R.N. Pilato and U. Marconi

theoretical: higher order radiative corrections modify the shapes

$$\sigma = \sigma^{(0)} + \left(\frac{\alpha}{\pi}\right)\sigma^{(1)} + \left(\frac{\alpha}{\pi}\right)^2\sigma^{(2)} + \ldots + \left(\frac{\alpha}{\pi}\right)^n\sigma^{(n)}$$

- order of magnitude estimate, barring infrared logs and setting $c_{i,j} \sim 10$
- $c_{1,1}\left(\frac{\alpha}{\pi}\right)L \sim 0.2$ $c_{1,0}\left(\frac{\alpha}{\pi}\right) \sim 2.5 \cdot 10^{-2}$ • $c_{2,2}\left(\frac{\alpha}{\pi}\right)^2 L^2 \sim 5 \cdot 10^{-3}$ $c_{2,1}\left(\frac{\alpha}{\pi}\right)^2 L \sim 5 \cdot 10^{-4}$ $c_{2,0}\left(\frac{\alpha}{\pi}\right)^2 \sim 5 \cdot 10^{-5}$ • $c_{3,3}\left(\frac{\alpha}{\pi}\right)^3 L^3 \sim 1.5 \cdot 10^{-4}$ $c_{3,1}\left(\frac{\alpha}{\pi}\right)^3 L^2 \sim 1.5 \cdot 10^{-5}$ $c_{3,0}\left(\frac{\alpha}{\pi}\right)^3 L \sim 1.5 \cdot 10^{-6}$
- the most advanced technologies for NNLO calculations and higher order resummation and matching are needed

thanks also to past

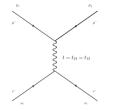
MUonE theory workshops

- Theory Kickoff Workshop, Padova, 4-5 September 2017
- MITP Workshop, Mainz 19-23 February 2018

• 2nd Workstop/ThinkStart, Zürich, 4-7 February 2019

First step towards precision: QED NLO (2018)

• analytical expression for tree level



NLO virtual diagrams

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2}{\lambda(s,m_{\mu}^2,m_e^2)} \left[\frac{(s-m_{\mu}^2-m_e^2)^2}{t^2} + \frac{s}{t} + \frac{1}{2} \right]$$

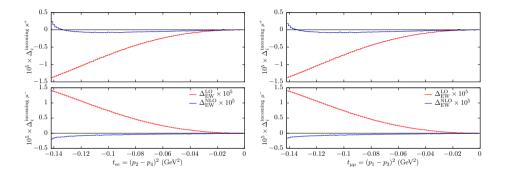
- VP gauge invariant subset of NLO rad. corr.
- factorized over tree-level: $\alpha \rightarrow \alpha(t)$

(Van Nieuwenhuizen 1971, D'Ambrosio 1983, Kukhto et al. 1987, Bardin, Kalinovskaya 1997)

and corresponding real emission diagrams

• NLO matrix elements calculated with finite m_{μ} and m_{e} mass effects and a Monte Carlo program, MESMER, has been taylored to the fixed target kinematics of MUonE

Alacevich, Carloni Calame, Chiesa, Montagna, Nicrosini, Piccinini, JHEP 02 (2019) 155

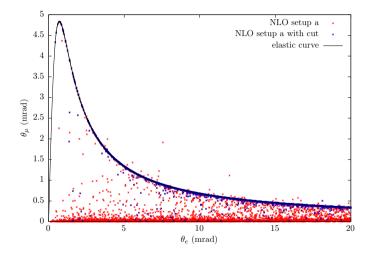


Alacevich, Carloni Calame, Chiesa, Montagna, Nicrosini, Piccinini, JHEP 02 (2019) 155

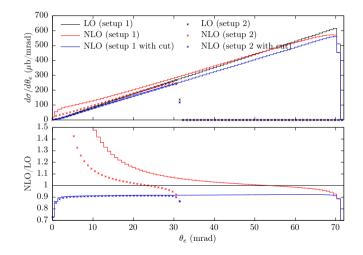
- tree-level Z-exchange important at the 10^{-5} level ($\sim tG_{\mu}/4\pi\alpha\sqrt{2}$ in the Fermi theory)
- purely weak RCs (in QED NLO units) at a few 10^{-6} level

What can we learn from "simple" NLO?

μ -e angle correlation in the lab

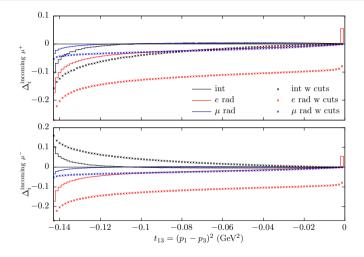


• many points fall out of the $2 \rightarrow 2$ correlation curve $\theta_{\mu} - \theta_{e}$ because of the radiative events



• due to the almost vanishing LO cross cross section for $\theta_e \rightarrow 0$

Radiation from electron leg not always dominant, in particular for "almost" elastic events



interference between electron and muon leg can be of the same order of magnitude as the contribution

from (virtual and real) radiation along the electron leg

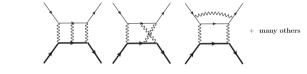
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Second step, towards photonic radiative corrections at NNLO (2020)

- exact calculation of corrections along one lepton line
 - two independent calculations, with different subtraction procedures
 - implemented in Mesmer and McMule

complete calculation (including up-down interference) approximated in Mesmer

• interference of LO $\mu e \rightarrow \mu e$ amplitude with



- NNLO double-virtual amplitudes where at least 2 photons connect the e and μ lines are approximated according to the Yennie-Frautschi-Suura ('61) formalism to catch the infra-red divergent structure
- complete calculation of the amplitude $f^+f^- \rightarrow F^+F^-$ with $m_f = 0, m_F \neq 0$
- R. Bonciani et al., PRL 128 (2022)
- talks by W.T. Bobadilla and J. Ronca

- complete calculation with massification in McMule
 - complete calculation including leading log $\propto \ln(m_e^2/Q^2)$ and m_e^0 terms

talks by T. Engel, M. Rocco, Y. Ulrich

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Carloni Calame et al., JHEP 11 (2020) 028,

P. Banerjee, T. Engel, A. Signer, Y. Ulrich, SciPost Phys. 9 (2020) 027 talks by E. Budassi, T. Engel, M. Rocco, Y. Ulrich

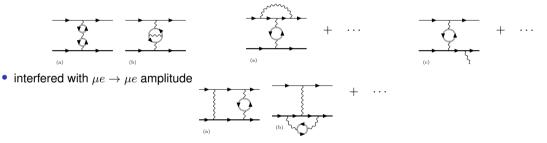
T. Engel, C. Gnendinger, A. Signer and Y. Ulrich, JHEP 02 (2019) 118

talk by E. Budassi

NNLO Virtual leptonic pairs (vacuum polarization insertions) (2021)

talk by C.L. Del Pio

- any lepton (and hadron) in the VP blobs
- interfered with $\mu e \rightarrow \mu e$ or $\mu e \rightarrow \mu e \gamma$ amplitudes

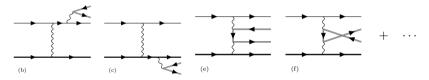


Here the 2-loop integral is evaluated with dispersion relation techniques used e.g. in the past for Bhabha: Actis et al., Phys. Rev. Lett. 100 (2008) 131602; Carloni Calame et al., JHEP 07 (2011) 126

$$\frac{g_{\mu\nu}}{q^2 + i\epsilon} \to g_{\mu\nu}\frac{\alpha}{3\pi}\int_{4m_\ell^2}^\infty \frac{dz}{z}\frac{R_\ell(z)}{q^2 - z + i\epsilon} = g_{\mu\nu}\frac{\alpha}{3\pi}\int_{4m_\ell^2}^\infty \frac{dz}{z}\frac{1}{q^2 - z + i\epsilon}\left(1 + \frac{4m_\ell^2}{2z}\right)\sqrt{1 - \frac{4m_\ell^2}{z}}$$

talk by C.L. Del Pio

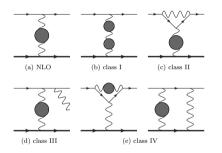
- they also contribute at NNLO accuracy
- squared absolute vaule of

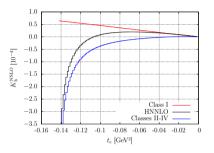


- the emission of an extra electron pair μe → μe e⁺e⁻ is potentially a dramatically large (reducible) background, because of the presence of "peripheral" diagrams
- NNLO virtual and real pairs implemented in Mesmer

NNLO hadronic contributions (2019)

• using the dispersion relation approach





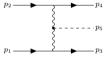
Fael, Passera, Phys. Rev. Lett. 122 (2019) 192001

relevant on the target precision scale

space-like method to calculate hadronic NNLO corrections

M. Fael, JHEP02 (2019) 027

- pion pair production forbidden kinematically with the available \sqrt{s}
- single π^0 production possible



• π^0 production recently calculated and shown to be well below 10^{-5} w.r.t. $\mu e \rightarrow \mu e$

E. Budassi et al., PLB 829 (2022) 137138

talk by C.L. Del Pio

• Possible contamination of New Physics in the $\Delta \alpha(t)$ measurement?

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

- Effects of heavy NP mediators investigated through EFT by considering dim-6 operators
 - excluded (at the 10^{-5} level) by existing data

- Effects of light NP mediators investigated through EFT by considering dim-6 operators
 - spin-0 NP mediators (ALPs)
 - spin-1 NP mediators (Dark Photons, light Z')
 - excluded (at the 10⁻⁵ level) by existing data

• interesting proposals for NP searches at MUonE (new light mediators) in 2 ightarrow 3 processes

• invisibly decaying light Z' in $\mu e \to \mu e Z'$

Asai et al., Phys. Rev. D106 (2022) 5

- a relevant background can be $\mu e
 ightarrow \mu e \pi^0$, in addition to $\mu e
 ightarrow \mu e \gamma$
- long-lived mediators with displaced vertex signatures $\mu e \rightarrow \mu e A' \rightarrow \mu e e^+ e^-$

Galon et al., arXiv:2202.08843

• through scattering off the target nuclei $\mu N \rightarrow \mu N X \rightarrow \mu N e^+ e^-$

Grilli di Cortona and E. Nardi, Phys. Rev. D105 (2022) L111701

- Given its precision requirements, MUonE represents a challenge for present theoretical calculations and simulation tools
- at present we have two Monte Carlo tools, Mesmer and McMule
 - NLO QED corrections cross-checked with perfect agreement
 - NNLO QED corections from single lepton legs in perfect agreement
 - YFS inspired approximation to the full NNLO QED in Mesmer
 - full NNLO QED with electron "massification" in McMule
 - efforts for N³LO started
 - very recent result on three-loop QCD form factors for massive quarks

M. Fael, F. Lange, K. Schönwald and M. Steinhauser, Phys, Rev. D106 (2022) 034029

talk by M. Fael

N³LO kick-off workshop/thinkstart, Durham, 3-5 August 2022 (organized by Y. Ulrich)