<u>Measurement:</u> $E1_{PNC}$ with another forbidden amplitude such as *M*1, *E*2 or $E1_{Stark}$

$$\begin{array}{ll}
\underline{\text{Theory:}} & |\Psi_{n}(n,J)\rangle = \left|\Psi_{n}^{(0)}(n,J,\pi)\rangle + G_{F} \left|\Psi_{n}^{(1)}(n,J,\pi')\rangle \\
E1_{PNC}^{NSI} = \frac{\langle \Psi_{f} | D | \Psi_{i} \rangle}{\langle \Psi_{f} | \Psi_{i} \rangle} = G_{F} & \frac{\left[\langle \Psi_{f}^{(0)} | D | \Psi_{i}^{(1)} \rangle + \langle \Psi_{f}^{(1)} | D | \Psi_{i}^{(0)} \rangle\right]}{\sqrt{\langle \Psi_{f}^{(0)} | \Psi_{f}^{(0)} \rangle \langle \Psi_{i}^{(0)} | \Psi_{i}^{(0)} \rangle}} & \xrightarrow{\rightarrow \text{Our}} \\
= \frac{G_{F}}{N} \left[\sum_{I \neq i} \frac{\langle \Psi_{f}^{(0)} | D | \Psi_{I}^{(0)} \rangle \langle \Psi_{I}^{(0)} | H_{PNC}^{NSI} | \Psi_{i}^{(0)} \rangle}{E_{i}^{(0)} - E_{I}^{(0)}} + \sum_{I \neq f} \frac{\langle \Psi_{f}^{(0)} | H_{PNC}^{NSI} | \Psi_{I}^{(0)} \rangle \langle \Psi_{I}^{(0)} | D | \Psi_{i}^{(0)} \rangle}{E_{f}^{(0)} - E_{I}^{(0)}} \\
\xrightarrow{\rightarrow \text{Others}} \end{array}$$

 $= Q_W X^{Theory} \rightarrow < 0.5\%$

<u>Combination:</u> Nuclear weak charge " Q_W ".

Calculation for ¹³³Cs in 10⁻¹¹ $(-Q_W/N)iea_0$

$$E1_{PNC}^{NSI}(6S \rightarrow 7S) = \sum_{np_{1/2}} \frac{\langle 7S|D|np_{1/2}\rangle\langle np_{1/2}|H_{PNC}^{NSI}|6S\rangle}{E_{6S}^{(0)} - E_{nP_{1/2}}^{(0)}} + \sum_{np_{1/2}} \frac{\langle 7S|H_{PNC}^{NSI}|np_{1/2}\rangle\langle np_{1/2}|D|6S\rangle}{E_{7S}^{(0)} - E_{np_{1/2}}^{(0)}}$$

= Core (n<6) + Main (n=6-9) + Tail

 -0.002(2) + 0.893(7) + 0.018(5)
 → 1%
 [Blundell et al, Phys. Rev. Lett. 65, 1411 (1990)]
 -0.0020 + 0.8823(17) + 0.0195
 → 0.27 %
 [Porsev et al, Phys. Rev. Lett. 102, 181601 (2009)]
 +0.0018(8) + 0.8823(17) + 0.0238(35)
 → 0.5 %
 [Dzuba et al, Phys. Rev. Lett. 109, 203003 (2012)]

The main objective of the present work is to address the issue of large differences in the core-contribution.

Commonly used many-body methods

Random phase approximation (RPA):

 $|\Psi_n^{(0)}\rangle \to |\Phi_n\rangle$ and $|\Psi_n^{(1)}\rangle \to \Omega_{I,CP}^{(\infty,1)}|\Phi_n\rangle = \Omega_{RPA}^{(1)}|\Phi_n\rangle$

Configuration interaction (CI) method:

$$\left|\Psi_{n}^{(0)}\right\rangle = C_{0}\left|\Phi_{n}\right\rangle + C_{I}\left|\Phi_{I}\right\rangle + C_{II}\left|\Phi_{II}\right\rangle + \cdots$$

Coupled-cluster (CC) method (all-order perturbation): $\begin{aligned} |\Psi_n^{(0)}\rangle &= C_0 |\Phi_n\rangle + C_I |\Phi_I\rangle + C_{II} |\Phi_{II}\rangle + \cdots \\ &= |\Phi_n\rangle + T_I^{(0)} |\Phi_n\rangle + T_{II}^{(0)} |\Phi_n\rangle + \frac{1}{2} T_I^{(0)^2} |\Phi_n\rangle + \cdots \\ &= e^{T_I^{(0)} + T_{II}^{(0)} + \cdots} |\Phi_n\rangle = e^{T^{(0)}} |\Phi_n\rangle \end{aligned}$

First-order perturbed RCC method

 $H = H_{at} + G_F H_{PNC}^{NSI}$ and $|\Psi_n\rangle \simeq |\Psi_n^{(0)}\rangle + G_F |\Psi_n^{(1)}\rangle$ First-order eqn.: $(H_{at} - E_n^{(0)}) |\Psi_n^{(1)}\rangle = (E_n^{(1)} - H_{PNC}^{NSI}) |\Psi_n^{(0)}\rangle$ $\Rightarrow \qquad \left| \Psi_n^{(1)} \right\rangle = e^{T^{(0)}} \left(1 + T^{(1)} \right) \left| \Phi_n \right\rangle$ It yields: $E1_{PNC}^{NSI} = \langle \Phi_f | e^{T^{(0)+}} D e^{T^{(0)}} T^{(1)} | \Phi_i \rangle$ $+\langle \Phi_{f}|T^{(1)+}e^{T^{(0)}+}De^{T^{(0)}}|\Phi_{i}\rangle$

Advantages:

- 1. Treats "Core", "Main" and "Tail" contributions on an equal footing unlike the "Sum-over-states" approach; thus enables to estimate errors consistently.
- 2. Incorporates all physical effects including double-core-polarization (DCP) effects implicitly to all orders in correlation effects at a given level of particle-hole excitation.

PHYSICAL REVIEW D 82, 036008 (2010)

Precision determination of weak charge of ¹³³Cs from atomic parity violation

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Revisiting Parity Nonconservation in Cesium

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TABLE I.	Partial	contributions	to	the	$E_{\rm PNC}$	[in
$10^{-11}i(-Q)$	$_W/N)$ a.u.] for Cs in differ	rent a	pproxir	nations.	

Approximation	Core	Main	Tail	Total
RPA ^a	0.0026	0.8705	0.0192	0.8923
$\mathrm{BO}(\hat{\Sigma}^{(2)})^{\mathrm{b}}$	0.0015	0.8641	0.0272	0.8928
$\mathrm{BO}(\lambda\hat{\Sigma}^{(2)})^{\mathrm{c}}$	0.0018	0.8709	0.0244	0.8971
$\mathrm{BO}(\hat{\Sigma}^{(\infty)})^{\mathrm{d}}$	0.0018	0.8711	0.0238	0.8967
$\mathrm{BO}(\lambda\hat{\Sigma}^{(\infty)})^{\mathrm{e}}$	0.0018	0.8678	0.0242	0.8938
Ref. [10] ^f	-0.0020	0.8823	0.0195	0.8998

^aCore polarization but no correlations beyond it. ^bBrueckner orbitals (BO) calculated with the second-order $\hat{\Sigma}$. ^cBO calculated with rescaled second-order $\hat{\Sigma}$. ^dBO calculated with the all-order $\hat{\Sigma}$. ^eBO calculated with rescaled all-order $\hat{\Sigma}$. ^fDHF for the core term; coupled cluster for the main term. TABLE IV. All significant contributions to the E_{PNC} [in $10^{-11}i(-Q_W/N)$ a.u.] for Cs.

Contribution	Value	Source		
Core (<i>n</i> < 6)	0.0018 (8)	This work		
Main $(n = 6-9)$	0.8823 (17)	Ref. [10]		
Tail $(n > 9)$	0.0238 (35)	This work		
Subtotal	0.9079 (40)	This work		
Breit	-0.0055 (1)	Refs. [5,6]		
QED	-0.0029 (3)	Ref. [7]		
Neutron skin	-0.0018 (5)	Ref. [5]		
Total	0.8977 (40)	This work		

Calculation for ¹³³Cs in 10⁻¹¹ $(-Q_W/N)iea_0$

RCC term	RLCCSD	RCCSD	RCCSDT	Matha	l Coro	Main	 	Fytro
	Core cont	tributions		Method	i Core	Main	Tall	Extra
$\overline{D}T_1^{(1)}$	-0.0534	-0.0410	-0.0410	This w	ork (ab initio	values)		
$T_1^{(1)\dagger}\overline{D}$	0.0519	0.0392	0.0392	DHF	-0.0017	0.7264	0.0137	
Others	-0.0001	-0.0001	~ 0.0	RCCSI	0 -0.0019	0.8623	0.0357	
Total	-0.0016	-0.0019	-0.0018	RCCSI	DT - 0.0018	0.8594	0.0391	0.0026
$\overline{D}S_{1i}^{(1)}$	Valence (Main+T -0.1663	`ail) contribution -0.1913	s -0.1874	Ref. [1] Ref. [1] Ref. [1]	$6]^{\dagger} -0.002(2)$ $3]^{\dagger} -0.0020$ $4]^{\dagger} 0.0018(8)$	0.893(7) 0.8823(17) 0.8678 $0.8823(17)^{\ddagger}$	0.018(5) 0.0195 0.0238(35)	5)
$S_{1f}^{(1)\dagger}\overline{D}$	2.0603	1.8064	1.7925	† Cont	sing additions	l contribution from	the $0\pi^2 P$	
$S_{1f}^{(0)\dagger} \overline{D} S_{1i}^{(1)}$	-0.3045	-0.2336	-0.2288	TADLE	IN D	i contribution from	the $9p P_1$	1/2 state.
$S_{1f}^{(1)\dagger}\overline{D}S_{1i}^{(0)}$	-0.5529	-0.4218	-0.4147	TABLE	IX. Progress	es in the atomic c	alculations	over the
$\overline{DS}_{2i}^{(1)}$	-0.0357	-0.0263	-0.0257	years.				
$S_{2t}^{(1)\dagger}\overline{D}$	0.0006	0.0009	0.0004	Year	Result	Approach	F	Reference
$T_{2}^{(b)\dagger}DS_{2i}^{(1)}$			-0.0019	1989	0.908(9)	$Ab \ initio$]	Ref. [15]
$S_{24}^{(1)\dagger}DT_{2}^{(0)}$			-0.0007	1990	0.909(4)	Sum-over-stat	les	Ref. [16]
$T_{(1)\dagger}^{(1)\dagger}\overline{D}S_{(0)}^{(0)}$			-0.0004	2001	0.901	Scaled optimal en	nergy l	Ref. [19]
$S^{(0)\dagger}\overline{D}T^{(1)}$			-0.0004	2002	0.904(5)	$Ab \ initio$]	Ref. [41]
$C_{3f}^{(0)\dagger} \overline{D} C_{2}^{(1)}$			-0.0000	2005	0.904	Ref. $[41]$ +QED	corr.	Ref. [20]
$S_{2f} DS_{3i}$			-0.0006	2009	0.8906(24)	Sum-over-sta	te l	Ref. [13]
$S_{3f}^{(c)}DS_{2i}^{(c)}$	0.0000	0.0000	0.0007	2012	0.8977(40)	Ref. $[13]$ +core of	corr. I	Ref. [14]
Others	-0.0608	-0.0363	-0.0343	2020	0.8914(27)	$Ab \ initio$	Г	This work
Total	0.9407	0.8980	0.8985					
OFD			-0.0055(5)	Our	orror octi	mation is may		ictont
Extra	-0.0026(3) OUP error estimation is more consistent						Istent	
e = e PNC	e = e PNC correction [40] 0.0003 as we treat all three contributions on						on	
Final	correction [10]		0.8914(27)	an e	nual-footi	ng in the RCC	` metho	h
			. /		γμαι-ιυυιί	ing in the NCC	, memo	/ u.

New value for nuclear weak charge

Combining our results of $E1_{PNC}$ and β with the precisely measured $Im(E1_{PNC}/\beta) = 1.5935(56)$ mV/cm [11], where Im means imaginary part, for the $6s \ ^2S_{1/2} - 7s \ ^2S_{1/2}$ transition in 133 Cs, we get $Q_W^{\text{at}} =$ $-73.43(25)_{ex}(23)_{th}$. After taking into account nuclear skin effect [43], we get

$$Q_W = Q_W^{\text{at}} + \Delta Q_W^{N-P}$$

= -73.43(25)_{ex}(23)_{th} + 0.064
= -73.37(25)_{ex}(24)_{th} (45)

This results in the difference between the value of Q_W obtained from our calculation and the SM value $Q_W^{\text{SM}} = -73.23(1)$ [4] as $\Delta Q_W \equiv Q_W - Q_W^{\text{SM}} = -0.14(35)$.

Summary & Outlook

- Our RCC method treats the ``Core", ``Main" and ``Tail" contributions to E1_{PNC} on an equal footing.
- ✤ It also accounts for DCP contributions implicitly.
- Our calculation demonstrates ``Core" contribution is agreeing with Porsev et al (2009 & 2010).
- It estimates uncertainties to ``Core", ``Main" and ``Tail" in a consistent manner and unambiguously.
- We are developing RCC methods to remove nonterminating series in the calculations.
- \Leftrightarrow S \rightarrow D transition or experiments in other systems.
- A novel technique using optical lattices is proposed.

Acknowledgement



B. P. Das

A. Kastberg

T. Aoki

Y. Sakemi



Thank You!