# Monte Carlo Program GF-CAIN for Simulations of Photon Emission in Collisions of Partially Stripped Ion Bunches with Laser-Photon Pulses

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#### What is GF-CAIN?

2 Monte Carlo simulations of laser-photon–PSI collisions

#### 3 Numerical results





#### CAIN

- Stand-alone Monte Carlo program for simulations of **beam-beam interactions** involving **high-energy electrons**, **positrons and photons**.
- Written by K. Yokoya et al., KEK, Japan, 1984–2011.
- $\bullet$  Code is a mixture of  ${\rm FORTRAN}$  77 and  ${\rm FORTRAN}$  90/95,  $\sim$  45000 lines in  $\sim$  400 files
  - $\rightarrow$  not well-documented, comments in code scarce.
- Dedicated, elaborate *meta-language* for defining Input/Output (65 pages of description in *User Manual*).
- Output in form of text files with all particle information and TopDrower histograms (no well-defined event record).



#### ABEL → CAIN history

- It started with program called **ABEL** for **beam-beam interactions** (deformation due to Coulomb field and beamstrahlung) in  $e^+e^-$  linear colliders.
- Then, after adding interactions with **laser** beams it was renamed to **CAIN**.
- CAIN 2.0 was written from scratch and allowed for any mixtures of  $e^+$ ,  $e^-$ ,  $\gamma$  and lasers, and multiple-stage interactions (input data format completely refreshed).
- Last version: CAIN 2.42, 27 June 2011, available at: https://ilc.kek.jp/~yokoya/CAIN/Cain242/



#### Physical processes in CAIN 2.42

- O Classical interactions (orbit deform.) due to Coulomb field.
- 2 Luminosity between beams  $(e^+, e^-, \gamma)$ .
- Synchrotron radiation by electrons/positrons (beamstrahlung) and (coherent) pair creation by high-energy photons due to beam field.
- Interactions of high-energy photon or electron/positron beams with laser field, including non-linear effects of field strength.
- Solution Classical and Quantum interactions with const. external field.
- **(**) Incoherent  $e^+e^-$ -pair creation by photons, electrons and positrons.
- Transport of charged particles through magnetic beamline.
- Polarisation effects can be included in most interactions (through polarisation vector for electron/prositron beams, Stokes parameters for photons).



#### Output of CAIN

- Output data (particle properties, luminosities, statistics, etc.) can be written in specified files at any moment of job
   → Can be huge! (for GF up to several GBs)
- Graphical output is written only in TopDrawer format
   → Obsolete!
- ▷ How to use CERN **ROOT** system for data analysis?
  - I For low statistics:

Write particle properties in **CAIN output file** and read them by CERN's **ROOT** data analysis program (in C++).

Por high statistics:

Transfer **CAIN output** to **input** of **ROOT** data analysis program (run concurrently) through UNIX **named (FIFO) pipes**.





Monte Carlo simulations of laser-photon-PSI collisions

# Scattering probability



• Scattering probability for a single particle (PSI) in time step  $\Delta t$ :

$$P(\vec{r},\vec{p},\vec{k},t) = \sigma_{\rm abs}(\vec{p},\vec{k}) \left(1 - \vec{\beta} \cdot \vec{k}/|\vec{k}|\right) n_p(x,y,z,k,t) c \Delta t,$$

where:  $\vec{k}$  – photon wave vector, c – velocity of light,  $\vec{p}$ ,  $\vec{\beta}$  – PSI momentum and relativistic velocity,  $n_p(x, y, z, k, t)$  – local density of laser-photon beam,  $\sigma_{\rm abs}(\vec{p}, \vec{k})$  – cross section for laser-photon absorption by PSI.

MC Program GF-CAIN



## Monte Carlo generation

- $\rightarrow$  Two stages of Monte Carlo simulation:
  - According to probability  $P(\vec{r}, \vec{p}, \vec{k}, t)$  scattering event is sampled using von Neumann rejection method.
  - When scattering event occurs emitted photon is generated, i.e. its energy and angles are generated in PSI rest-frame according to differential cross section, and then event is Lorentz-transformed to LAB frame.
  - ▷ The above is repeated for each **macroparticle**, and then generation moves to the **next** time moment, i.e.  $t + \Delta t$ , ... .
  - One macroparticle represents some number of real particles (PSI) in a bunch (simulations for each real particle may be not feasible if their number is very large!).
  - To each **macroparticle** a Monte Carlo **weight** is assigned which is a ratio of the number of real particles to the number of macroparticles (the smaller weight the better).

#### Cross section

#### • Cross section of photon-absorption by PSI:

[E.G. Bessonov and K.J. Kim, IEEE PAC 1995:2895-2897]

$$\sigma_{\rm abs}(\vec{p},\vec{k}) = \frac{2\pi r_e c f \Gamma}{[\gamma \omega (1-\beta \cos \psi) - \omega_0]^2 + \Gamma^2}$$

- $r_e$  classical electron radius,
- f oscilator strength,
- $\gamma,\beta$  relativistic factor and velocity of PSI,
- $\omega$  incoming photon frequency,
- $\psi$  angle between incoming photon and PSI,

 $\omega_0$  – PSI transition frequency between states 1 and 2,

 $\Gamma = \omega_0^2 r_e fg_1/(cg_2)$  – spontaneous emission half-linewidth,

where  $g_{1,2}$  – degeneracy factors of states 1 and 2, respectively.



#### Emitted photon kinematics

- MC generation of emitted photon in PSI rest-frame
   ⇒ Unpolarised case so far!
  - azimuthal angle  $\phi$ :

 $\phi \in \mathcal{U}(0,1),$ 

where  $\boldsymbol{\mathcal{U}}$  denotes  $\boldsymbol{Uniform}$  distribution,

**2** polar angle  $\theta$ :

 $\cos heta \in \mathcal{U}(-1,1),$ 

(angular frequency  $\omega' (\rightarrow \text{energy } E' = \hbar \omega')$ :

$$\omega' \in \mathcal{L}(\omega'_{\min}, \omega'_{\max}),$$

where  $\mathcal{L}$  – Lorentzian distribution with prob. density funct.:

$$\rho_{\omega_0,\Gamma}(\omega';\omega'_{\min},\omega'_{\max}) = \mathcal{N} \frac{\Gamma}{(\omega'-\omega_0)^2 + \Gamma^2},$$

with 
$$\mathcal{N}^{-1} = \arctan([\omega'_{max} - \omega_0]/\Gamma) - \arctan([\omega'_{min} - \omega_0]/\Gamma)$$



#### Energy spread of laser beam

- CAIN assumes monochromatic laser beam (photon energy spread not important for inverse-Compton scattering).
- For resonant atomic photon absorption laser-beam energy spread can be comparable or even larger than the resonance linewidth, so it has to be taken into account!
- In GF-CAIN it is done in two ways (inside corresponding routines):
  - If  $\sigma_{\bar{\omega}}/\bar{\omega} < \Gamma/\omega_0$ , the laser-photon energy  $E = \hbar\omega$  is generated from the corresponding Gaussian distribution, then the scattering cross section is calculated using the weight corresponding to  $\sigma_{abs}(\vec{p}, \vec{k})$ .
  - **2** Otherwise, the **photon energy** in the PSI-rest frame is generated from the **Lorentzian distribution** of  $\sigma_{abs}(\vec{p}, \vec{k})$ , then the **scattering cross section** is calculated using the **weight** corresponding to the **Gaussian** function of the **laser-energy spread**.
- ▷ In this way, Monte Carlo event generation in GF-CAIN is efficient for an arbitrary resonance linewidth!



### Li-like Pb, H-like Pb and He-like Ca

- PSI's cannot be defined by **CAIN** input they are implemented in **CAIN** routine LNCPGN:
  - Lithium-like  ${}^{208}_{82} {\rm Pb}^{79+}$  in file src/GF/Pb/lncpgn-Pb\_Li-like.f
  - Hydrogen-like  $^{208}_{82}$ Pb<sup>81+</sup> in file src/GF/Pb/lncpgn-Pb\_H-like.f
  - Helium-like  $^{40}_{20}\mathrm{Ca}^{18+}$  in file /src/GF/Ca/lncpgn-Ca\_He-like.f
- They are copied into **CAIN**'s file /src/lncpgn.f in Makefile when the corresponding PSI-run is chosen by a make command, e.g.
  - make run-PbLi
  - make run-PbH
  - make run-CaHe

and then an appropriate input file is read.

- **Spontaneous** emission **delay** and **stimulated** emission have been added important for PoP experiment  $Pb^{79+}$  as well as for  $Ca^{18+}$ 
  - $\rightarrow$  appropriate modifications of CAIN event record as well as 'drift' routines were necessary.
- Other PSI's can be implemented in a similar way not elegant, but easier than modifying complicated **CAIN** input!



MC Program GF-CAIN

#### H-like Pb – input parameters (based on Bessonov et al.)

• PSI beam:  ${}^{208}_{82}\text{Pb}^{81+}$  with transition:  $1s^1 \, {}^2S_{1/2} \rightarrow 2p^1 \, {}^2P_{1/2}$ 

- transition energy:  $\hbar \omega_0 = 68.7 \, {\rm keV}; \ f = 0.416, \ g_1 = 1, \ g_2 = 3$
- ion mass:  $M_i = 193.687 \, {
  m GeV/c^2}$
- ion energy and relative spread:  $E_i = 579 \,\mathrm{TeV}$ ,  $\sigma_E = 2 \cdot 10^{-4}$
- relativistic factor:  $\gamma_i = 2989$
- number of ions per bunch  $N_i = 9.4 \cdot 10^7$
- beta function in IR:  $\beta_x = \beta_y = 0.5 \,\mathrm{m}$
- geometric emittance:  $\epsilon_x = \epsilon_y = 3 \cdot 10^{-9} \,\mathrm{m \, rad}$
- r.m.s transverse beam size:  $\sigma_x = \sigma_y = 38.73 \,\mu\mathrm{m}$
- r.m.s. bunch length  $\sigma_z=15\,{
  m cm}$
- Laser: Gaussian spatial and time profiles
  - photon energy and rel. spread:  $E_{\gamma} = 11.45 \, {\rm eV}$ ,  $\sigma_{\omega} = 2 \cdot 10^{-4}$
  - photon wavelength:  $\lambda_{\gamma} = 108.28\,\mathrm{nm}$
  - pulse energy:  $W_I = 56 \,\mu J$
  - peak power density:  $P_{00} = 1.1 \cdot 10^{13} \, {
    m W/m^2}$
  - r.m.s. transverse beam size at focus:  $\sigma_x = \sigma_y = 25.42 \,\mu\mathrm{m}$
  - Rayleigh length:  $R_{L,x} = R_{L,y} = 7.5 \,\mathrm{cm}$
  - r.m.s. pulse length:  $\sigma_z = 15 \,\mathrm{cm}$



#### Emitted photon energy in LAB

- Number of macroparticles generated in GF-CAIN: 9.4 · 10<sup>7</sup>
- Spontaneous emission delay included (small in this case)



- ightarrow Half of most energetic photons within  $heta < 1/\gamma_{\mathsf{i}}$
- $\rightarrow$  Number of emitted **photons** per **ion**: N<sub> $\gamma$ </sub>/N<sub>i</sub> = 0.11



### Comparisons: GF-CAIN vs. GF-CMCC

 Comparisons with the independent Monte Carlo program GF-CMCC of Camilla Curatolo (INFN-Padova)



 $\rightarrow$  Very good agreement of the two MC programs!



#### Lithium-like Pb ion for PoP – input parameters (Lol)

- PSI beam:  ${}^{208}_{82}\mathrm{Pb}^{79+}$  with transition:  $1s^22s^1 \ {}^2S_{1/2} \rightarrow 1s^22p^1 \ {}^2P_{1/2}$ 
  - transition energy and lifetime:  $\hbar\omega_0=230.81\,{\rm eV},\;\tau_0=76.6\,{\rm ps}$
  - ion mass:  $M_i = 193.687 \, {
    m GeV/c^2}$
  - ion energy and relative spread:  $E_i = 18.65259 \text{ TeV}, \sigma_E = 2 \cdot 10^{-4}$
  - relativistic factor:  $\gamma_i = 96.3$
  - number of ions per bunch  $N_i = 0.9 \cdot 10^8$
  - Twiss parameters:  $\alpha_x = \alpha_y = 0$ ,  $\beta_x = 70.30 \,\mathrm{m}$ ,  $\beta_y = 44.23 \,\mathrm{m}$
  - geometric emittance:  $\epsilon_x = \epsilon_y = 1.558 \cdot 10^{-8} \,\mathrm{m \, rad}$
  - r.m.s transverse beam size:  $\sigma_x = 1.047 \text{ mm}, \sigma_y = 0.83 \text{ mm}$
  - r.m.s. bunch length  $\sigma_z=6.386\,\mathrm{cm}$
- Laser: Gaussian spatial-time profiles, beam angle: 2.6°
  - photon energy and rel. spread:  $E_{\gamma}=1.2\,\mathrm{eV}$ ,  $\sigma_{\omega}=2\cdot10^{-4}$
  - photon wavelength:  $\lambda_{\gamma} = 1034\,\mathrm{nm}$
  - pulse energy:  $W_l = 5 \,\mathrm{mJ}$
  - peak power density:  $P_{00} = 2.684 \cdot 10^{14} \, {
    m W/m^2}$
  - r.m.s. transverse beam size at focus:  $\sigma_x = \sigma_y = 0.65 \,\mathrm{mm}$
  - Rayleigh length:  $R_{L,x} = R_{L,y} = 5.135 \,\mathrm{m}$
  - r.m.s. pulse length:  $\sigma_z = 0.8394 \,\mathrm{mm}$



### Comparisons: GF-CAIN vs. GF-CMCC

 Comparisons with the independent Monte Carlo program GF-CMCC of Camilla Curatolo (INFN-Padova)



 $\rightarrow$  Very good agreement of the two MC programs!



#### Fraction of excited ions

Predictions of independent Monte Carlo programs GF-CMCC (Camilla Curatolo) GF-Python (Alexey Petrenko) and GF-CAIN





#### Spontaneous emission delay and stimulated emission

- $\bullet\,$  Mean path of PSI in excited state in LAB  $\approx 2.2\,m$ 
  - $\rightarrow$  Two important effects included in GF-CAIN:
    - time delay of spontaneous emission generated from the exponential distribution with the mean-time τ<sub>0</sub>, and the excited ion is propagated until it de-exites: (1) after the generated time τ by spontaneous emission or (2) immediately by stimulated emission, or reaches a given z or t coordinate (e.g. detector) in the exited state.
    - **2** stimulated emission generated according to the probability  $P'(\vec{r}, \vec{p}, \vec{k}, t) = (g_1/g_2)P(\vec{r}, \vec{p}, \vec{k}, t)$ , where  $P(\vec{r}, \vec{p}, \vec{k}, t)$  is the photon-absorption probability and  $g_{1,2}$  are the state-degeneracy factors, and when the event is accepted, the ion returns to the ground state while the two photons are discarded.

<b>GF-CAIN</b> simulation results at $z = 6 \mathrm{m}$ :	$N_{\gamma}/N_{i}$
No spontaneous emission delay:	20.1%
With spontaneous emission delay:	15.7%
With spont. emission delay and stimulated emission	13.3%



#### Photon x-coordinate and radius distributions at z = 6m





#### Photon radius vs. energy distributions at z = 6m





#### Example for Doppler cooling of PSI beam

- ullet Laser energy lowered by  $2\sigma_\omega$  w.r.t. resonance energy
  - excited ions
- ground-state ions



### Helium-like Ca ion – input parameters (transmutations)

- PSI beam:  $^{40}_{20}\mathrm{Ca}^{18+}$  with transition:  $1s^2$   $^1S_0 \rightarrow 1s^12p^1$   $^1P_1$ 
  - transition energy and lifetime:  $\hbar\omega_0=3.9023775\,{\rm keV},\;\tau_0=8.8\,{\rm fs}$
  - ion mass:  $M_i = 37.332 \, {
    m GeV/c^2}$
  - ion energy and relative spread:  $E_i = 39.72 \,\mathrm{TeV}$ ,  $\sigma_E = 2 \cdot 10^{-4}$
  - relativistic factor:  $\gamma_i = 1064$
  - number of ions per bunch  $N_i = 3 \cdot 10^9$
  - Twiss parameters:  $\alpha_x = \alpha_y = 0$ ,  $\beta_x = \beta_y = 50 \,\mathrm{m}$
  - geometric emittance:  $\epsilon_x = \epsilon_y = 3 \cdot 10^{-10} \,\mathrm{m \, rad}$
  - r.m.s transverse beam size:  $\sigma_x = \sigma_y = 0.1225 \,\mathrm{mm}$
  - ullet r.m.s. bunch length  $\sigma_z=15\,{\rm cm}$
- Laser: Gaussian spatial-time profiles, beam angle: 0°
  - photon energy and rel. spread:  $E_{\gamma} = 1.833824 \, {\rm eV}$ ,  $\sigma_{\omega} = 2 \cdot 10^{-4}$
  - photon wavelength:  $\lambda_\gamma = 676.1\,\mathrm{nm}$
  - pulse energy:  $W_l = 0.5 \,\mathrm{mJ}$
  - peak power density:  $P_{00} = 2.822 \cdot 10^{13} \, {
    m W/m^2}$
  - r.m.s. transverse beam size at focus:  $\sigma_x = \sigma_y = 0.15 \,\mathrm{mm}$
  - Rayleigh length:  $R_{L,x} = R_{L,y} = 41.81996$ , cm
  - r.m.s. pulse length:  $\sigma_z = 1.49896\,\mathrm{cm}~(\sigma_t = 50\,\mathrm{ps})$



#### Photon emission angle and energy

- ▷ Repetition rate: 20 MHz
- $\triangleright$  Emission rate: N<sub> $\gamma$ </sub>/N<sub>PSI</sub>  $\approx$  5
- $\rightarrow$  For  $\sigma_t = 500 \text{ ps: } N_{\gamma}/N_{PSI} \approx 30$



in/dE<sub>7</sub> [MeV<sup>-1</sup> s<sup>-1</sup>]

0

7

E, [MeV]

#### Summary

- CAIN code has been customised to compile with gfortran (GNU Fortran) and run on Linux and macOS systems
   → with the use of customised Makefile
- CAIN Monte Carlo program has been debugged and adapted to laser-photon pulse collisions with PSI beams of  ${}^{208}_{82}\text{Pb}^{81+}$ ,  ${}^{208}_{82}\text{Pb}^{79+}$  and  ${}^{40}_{20}\text{Ca}^{18+}$  (Gamma Factory)  $\Rightarrow$  GF-CAIN.
- **Spontaneous** emission **delay** and **stimulated** emission have been implemented important for PoP experiment.
- **GF-CAIN** output has been interfaced with **ROOT** data analysis program via UNIX **named (FIFO) pipes**.
- Good agreement with independent Monte Carlo event generators GF-CMCC of Camilla Curatolo and GF-Python of Alexey Petrenko.
- $\bullet\,$  Statistics of  $\sim 10^8$  macroparticles can be generated on medium PC