1+1D Entanglement in conformal field theories with boundaries

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Quantum field theory at the boundary, Mainz Virtual Workshop 09/28/2021

Entanglement in many body systems



A and B are entangled if $|\psi_{AUB}\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$

Measures of entanglement

A and B are entangled if $|\psi_{AUB}\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$

Measures of entanglement:

A

von Neumann entropy: $S_A = -\text{Tr}\rho_A \ln \rho_A$, $\rho_A = \text{Tr}_B \rho_{AUB}$ Entanglement Hamiltonian: $\mathcal{H}_A = -\frac{1}{2\pi} \ln \rho_A$

log negativity: $\mathcal{E}_A = \ln \left| \rho_{AUB}^{T_B} \right|$

Contain signatures of quantum critical phenomena, topological order, information scrambling, ...

Eisert *et al* (2010), Haag (1992), Li and Haldane (2008), Vidal and Werner (2004)



Holzhey et al (1994), Cardy and Calabrese (2004), Hastings (2007)

Entanglement entropy in 1+1D conformal-invariant quantum systems A B Entanglement entropy: $S_A = \frac{c}{6} \ln l_A + \cdots$, c = central charge

Ising chain: $H = -\sum_{i} \sigma_{i}^{x} \sigma_{i+1}^{x} - \lambda \sum_{i} \sigma_{i}^{z}$ $H = \sum_{i} n_{i}^{2} - \lambda \sum_{i} \cos(\phi_{i} - \phi_{i+1})$ $[n_{i}, e^{\pm i\phi_{j}}] = \pm e^{\pm i\phi_{j}} \delta_{ij}$ gapped
gapped c = 1 λ $\gamma = 1$ Entanglement entropy in 1+1D conformal-invariant quantum systems with boundaries Α В α $L = l_A + l_B$ $S_A(\alpha, \beta) = \frac{c}{6} \ln \left[\frac{2L}{\pi} \sin \frac{\pi l_A}{L} \right] + S_\alpha + S_\beta + S_0$ **Entanglement entropy:** S_{α} , $S_{\beta} = \ln g_{\alpha}$, $\ln g_{\beta} =$ boundary entropies, $\alpha =$ inherited, $\beta =$ free (from

Easier to extract than conventional thermodynamic quantities like Casimir energy, impurity free energy, ...

entanglement cut)

Cardy and Calabrese (2004, 2009), Cardy and Tonni (2016), Affleck and Ludwig (1991)

Entanglement Hamiltonian in 1+1D conformalinvariant quantum systems with boundaries A B

Entanglement Hamiltonian:

$$\mathcal{H}_{A}(\boldsymbol{\alpha},\boldsymbol{\beta}) = -\frac{1}{2\pi} \ln \rho_{A}(\boldsymbol{\alpha},\boldsymbol{\beta}) \simeq -\frac{1}{2\pi} \ln \frac{e^{-2\pi H_{\boldsymbol{\alpha}\boldsymbol{\beta}}}}{\operatorname{Tr} e^{-2\pi H_{\boldsymbol{\alpha}\boldsymbol{\beta}}}}$$

 $H_{\alpha\beta}$ = boundary CFT Hamiltonian

Entanglement spectrum obtained from standard partition function computation



Cardy and Tonni (2016), AR et al, J. Stat. Mech (2020)

Entanglement entropy in the critical Ising chain with boundaries: Results $DMRG (\simeq matrix)$

Hamiltonian:

$$H = -\sum_{i} \sigma_i^x \sigma_{i+1}^x - \sum_{i} \sigma_i^z + \lambda_0 (\sigma_0^x + \sigma_{L-1}^x)$$

Neumann (free): $\lambda_0 = 0$, Dirichlet (fixed): $\lambda_0 \ll 1$

Central charge =
$$1/2$$

Change in boundary entropy:

$$\Delta S_{N \to D} = \frac{1}{2} \ln 2$$

AR et al, J. Stat. Mech (2020), Affleck and Ludwig (1991)



Entanglement spectrum in the critical Ising chain with boundaries: Results

Boundary states:

 $\left|\frac{\tilde{1}}{16}\right\rangle = \left|0\right\rangle - \left|\epsilon\right\rangle$

Partition function (Neumann boundary

$$\begin{split} |\tilde{0}\rangle &= \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |\epsilon\rangle + \frac{1}{2^{1/4}} |\sigma\rangle, \qquad \text{condition}:\\ |\frac{\tilde{1}}{2}\rangle &= \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |\epsilon\rangle - \frac{1}{2^{\frac{1}{4}}} |\sigma\rangle, \qquad Z_{\text{NN}} = \int_{j=-\infty}^{x} \left| \frac{\alpha}{\sqrt{2}} \right|^{2} \left| \frac{1}{16} \left| j \right| \right|^{2} \chi \left| \tilde{q} \right|, \tilde{q} = e^{-2L_{\text{eff}}} \\ |\frac{\tilde{1}}{16}\rangle &= |0\rangle - |\epsilon\rangle \end{split}$$

Entanglement spectrum (Neumann):

$$e(j,n) = \frac{L_{\text{eff}}}{48\pi} + \frac{\pi}{L_{\text{eff}}} \left(-\frac{1}{48} + h_j + n \right) + \frac{1}{2\pi} \ln \sum_{k=0,\epsilon} \sum_{m \ge 0} p_k(m) e^{-2L(h_k + m)}$$

degeneracy of $e(j,n) = p_i(n)$

Ishibashi (1988), Cardy (1989), di Francesco et al (1997), AR et al, J. Stat. Mech (2020)

Entanglement spectrum in the critical Ising chain with boundaries: DMRG Results



AR et al, J. Stat. Mech (2020)

Entanglement spectrum in the critical Ising chain with boundaries: DMRG Results



AR et al, J. Stat. Mech (2020)

Entanglement in the 1+1D free, compact boson CFT

Euclidean action:

$$A = \frac{1}{2\pi K} \int \left(\partial_{\mu}\phi\right)^2 dx dt ,$$

 $K = \text{Luttinger parameter, } \left\langle e^{i\phi(0)}e^{-i\phi(r)} \right\rangle \sim \frac{1}{r^{K/2}}$

R =compactification radius

Boundary states:

$$|D(\phi_0)\rangle = \frac{1}{\sqrt{2R\sqrt{\pi}}} \sum_{n} e^{-\frac{in\phi_0}{R\sqrt{\pi}}} \prod_{k>0} e^{-\tilde{a}_{-k}^{\dagger} \tilde{a}_{k}^{\dagger}} |n,0\rangle \qquad \text{Dirichle}^{-\tilde{a}_{-k}} \tilde{a}_{k}^{\dagger} |n,0\rangle$$

$$|N(\tilde{\phi}_0)\rangle = \sqrt{R\sqrt{\pi}} \sum_{n} e^{-\frac{im\tilde{\phi}_0 R}{2\sqrt{\pi}}} \prod_{k>0} e^{+\tilde{a}_{-k}^{\dagger}\tilde{a}_{k}^{\dagger}} |0,m\rangle \qquad \text{Neumann}$$

Callan *et al* (1987), Affleck and Oshikawa (1997)

Entanglement in the free, compact boson CFT: Results

Euclidean action:

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Change in boundary entropy from Neumann to Dirichlet:

$$\Delta S_{N \to D} = \frac{1}{2} \ln \frac{2}{K}$$

degeneracy p(l) = # of integer partitioning of l

Entanglement spectrum (Neumann):

$$\varepsilon_N(k,l) = \varepsilon_N(0,0) + \frac{\pi}{L_{\text{eff}}} \left(\frac{K}{2} \frac{k^2}{2} + \frac{l}{2}\right)$$

exact form available dim. of primary fields

ds descendant level

AR et al, J Stat. Mech. (2020), see also: Saleur (1998)

Entanglement in the free, compact boson CFT: DMRG results for the quantum rotor chain

Hamiltonian with nearest-neighbor interaction:

$$\left[n_i, e^{\pm i\phi_j}\right] = \pm e^{\pm i\phi_j} \,\delta_{ij}$$



Entanglement in the free, compact boson CFT: DMRG results for the quantum rotor chain

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8

Δε_N(0,

N(k, l)

Hamiltonian with nearest-neighbor interaction:

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$$H = \sum_{i} n_i^2 + \epsilon \sum_{i} n_i n_{i+1} - \lambda \sum_{i} \cos(\phi_i - \phi_{i+1})$$

Free, compact boson for $\lambda > 1, \epsilon \ll 1$

Free, compact boson for $\lambda > 1, \epsilon \ll 1$

Entanglement spectrum (Neumann):

AR *et al,* J Stat. Mech. (2020)

Chose λ , ϵ such that K = 0.192

Non-generic behavior, characteristic of an integrable model

p(4) = 5

p(3) = 3

Entanglement in the free, compact boson CFT: DMRG results for the quantum rotor chain

Hamiltonian with nearest-neighbor interaction:

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 $H = \sum_{i} n_i^2 + \epsilon \sum_{i} n_i n_{i+1} - \lambda \sum_{i} \cos(\phi_i - \phi_{i+1})$ Free, compact boson for $\lambda > 1, \epsilon \ll 1$ 6 $\begin{array}{c|c}
 \hline c \\
 \hline c \\$ Entanglement spectrum (Dirichlet): $p_{\sigma}(2) = 1$ $0 \quad p_{\sigma}(0) = 1$ $0 \quad 5 \quad 10 \quad 2$ rescaled $\varepsilon_D(n) = \varepsilon_D(0) + \frac{\pi}{2I_{Loff}}n$ entanglement 15 energies n

AR *et al,* J Stat. Mech. (2020)

Chose λ , ϵ such that K = 0.192

Entanglement in 1+1D CFTs perturbed by a primary field α β α α

Hamiltonian:

$$H = H_{\rm CFT} + \lambda' \int dx \, \Phi$$

e.g., Ising chain with $\lambda \neq 1$:

$$H = -\sum_{i} \sigma_{i}^{x} \sigma_{i+1}^{x} - \lambda \sum_{i} \sigma_{i}^{z}$$

Entanglement entropy:

 $S_A = \frac{c}{6} \ln \xi + \cdots$ c = central charge, ξ = correlation length Entanglement Hamiltonian:

$$\mathcal{H}_A(\alpha, \beta) \sim -\frac{1}{2\pi} \ln \frac{e^{-2\pi H_{\Phi}\beta}}{\mathrm{Tr}e^{-2\pi H_{\Phi}\beta}}$$

Entanglement gap: $\varepsilon \sim 1/\ln \xi$

does not require lattice integrability!

Peschel (1999, 2001), Cho et al (2017)

Entanglement spectrum of the quantum sine-Gordon model : DMRG results

Hamiltonian with nearest-neighbor interaction:

 $H = H_0 + H'$

$$H_0 = \sum_i n_i^2 + \epsilon \sum_i n_i n_{i+1} - \lambda \sum_i \cos(\phi_i - \phi_{i+1})$$

Free, compact boson for $\lambda > 1, \epsilon \ll 1$

 $H' = -\mu \sum_{i} \cos \phi_i$

Predicted degeneracies: 1,1,1,2,2,3,4,...

Lattice integrable regularization: XYZ spin-chain



AR et al, Nucl. Phys. B (2021)

The XYZ regularization of the quantum sine-Gordon¹ model

Baxter's XYZ spin-chain

$$H_{XYZ} = -\frac{1}{2} \sum_{i=1}^{L} \left[J_x \sigma_i^x \sigma_{i+1}^x + J_y \sigma_i^y \sigma_{i+1}^y + J_z \sigma_i^z \sigma_{i+1}^z \right], J_x > J_y \ge |J_z|$$

$$i\beta\phi$$

XYZ to sine-Gordon operator mapping: $\sigma^+ \sim e^{\frac{\tau + \tau}{2}}$

Relation to eight-vertex model parameters (in the principal regime):

$$\Gamma = \frac{J_y}{J_x}, \Delta = \frac{J_z}{J_x},$$
$$\frac{2\sqrt{k}}{1+k} = \sqrt{\frac{1-\Gamma^2}{\Delta^2 - \Gamma^2}}, -i \operatorname{sn}(i\lambda, k) = \frac{1}{\sqrt{k}} \sqrt{\frac{1-\Gamma}{1+\Gamma}}$$

Baxter (1982), Luther (1975), Lukyanov (1997, 2003)

The XYZ regularization of the quantum sine-Gordon² model

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Entanglement spectrum obtained from corner transfer matrix spectrum, level spacing: $\varepsilon_{XYZ} = \frac{\pi\lambda}{I(k)}$

AR et al, Nucl. Phys. B (2021)

Baxter (1982)

Entanglement spectrum of the quantum sine-Gordon model : DMRG vs corner transfer matrix results



Entanglement spectrum for integrable vs nonintegrable regularizations





22

Outlook: experimental realization with superconducting quantum electronic circuits





 $\varphi_0 = \hbar/2e$

The free, compact boson CFT with ²⁴ superconducting quantum electronic circuits



$$H = \epsilon \sum_{i} n_{i} n_{i+1} - \lambda \sum_{i} \cos(\phi_{i} - \phi_{i+1})$$

Free, compact boson for $\lambda > 1, \epsilon \ll 1$

AR et al, J. Stat. Mech. (2020)

Where are the experiments?



Devoret group (Yale, 2009): N = 43



Roch group (Grenoble , 2019): $N \sim 1500$

Probing strongly-interacting quantum field theories





Gershenson group (Rutgers, 2012): N = 6, 24

Esteve group (Saclay, 2013): $N \sim 100$



Manucharyan group (Maryland, 2018, 2019): $N \sim 33000$

Early experiments: Delft (1990-s)

Quantum circuits as analog free boson QFT simulators



Maryland group (2018, 2019): *N* ~ 33000

Quantum circuit





Grenoble group (2019): *N* ~ 1500

Quantum circuit



More interesting field theories with quantum circuits: AR and H. Saleur, Phys. Rev. B (2019), AR *et al*, Nucl. Phys. B (2021)

Thank You! Hubert Saleur (CEA Saclay) Johannes Hauschild (UC Berkeley)