

HYPERON-NUCLEAR INTERACTIONS

and

STRANGENESS in NEUTRON STARS



Wolfram Weise
Technische Universität München



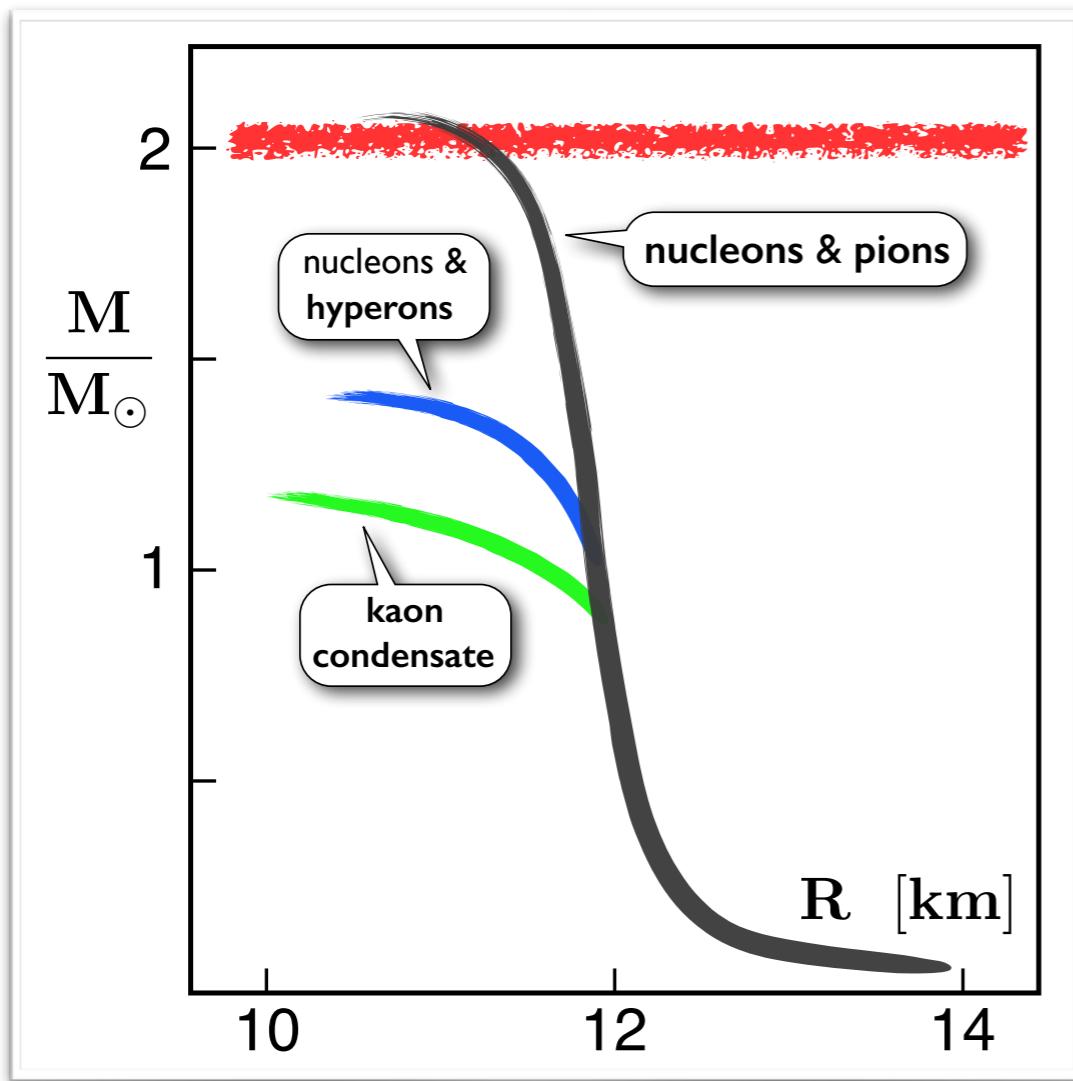
- ★ **Brief overview: neutron star equation-of-state**
 - Constraints from $2 M_{\odot}$ stars and mergers (GW signals)
 - Neutron star matter as a (relativistic) Fermi Liquid

- ★ **Strangeness and baryonic matter**
 - Hyperon-nuclear interactions and three-body forces
 - “Hyperon puzzle” in neutron stars

Key Issue:

Role of **STRANGENESS** (Kaons/Antikaons, Hyperons) in **DENSE BARYONIC MATTER** (Neutron Stars)

- Several heavy neutron stars observed with masses around $M \simeq 2 M_{\odot}$.



- Hyperon puzzle (strangeness puzzle) in neutron star matter:

Appearance of hyperons
(or kaon condensate)



neutron star equation-of-state
far too soft !

- Basic requirements:

- Many-body theory of dense baryonic matter
- Theory of hadronic interactions including strangeness

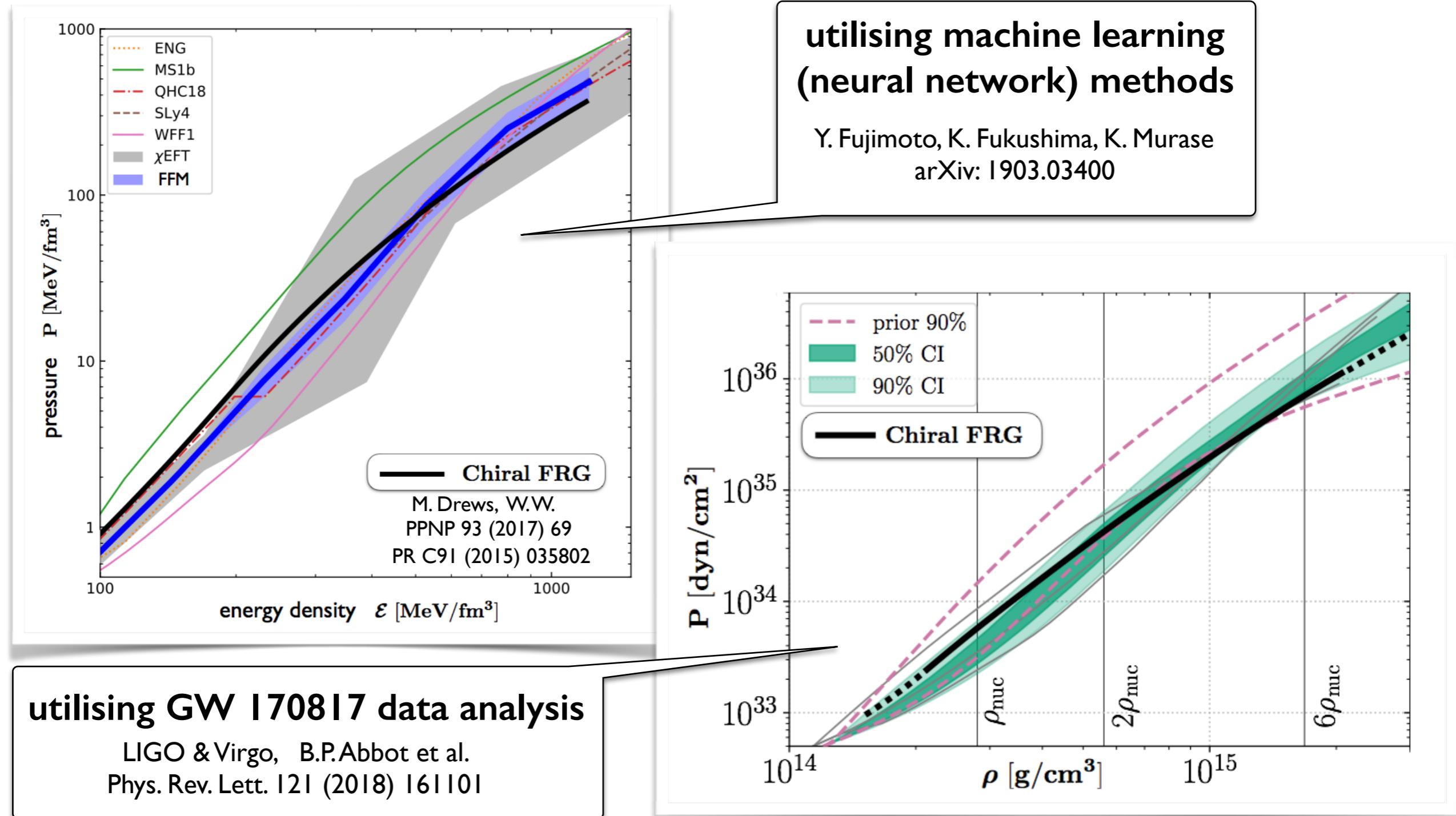
Part I.

Status of Neutron Star Equation-of-State

- **Stiffness constraint:**
EoS must support massive ($\sim 2 M_{\odot}$) neutron stars
- **Compatibility with GW neutron star merger signals**
(tidal deformability)
- **Microscopic EoS and Landau Fermi-Liquid theory**

NEUTRON STAR MATTER Equation of State

- Comparisons: EoS reconstructions from neutron star data and a microscopic equation-of-state



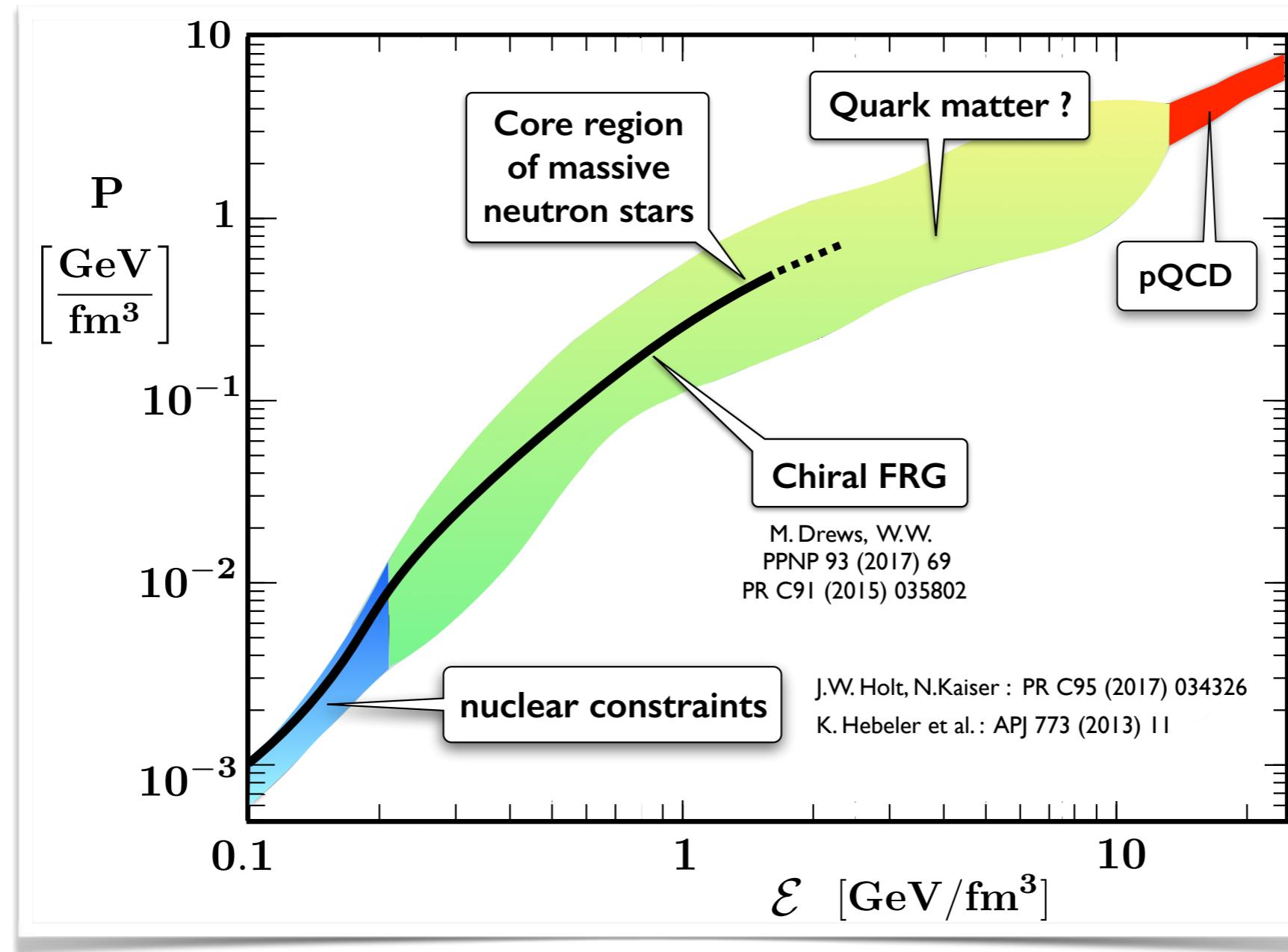
NEUTRON STAR MATTER Equation of State

- incl. $M_{max} \simeq 2 M_\odot$ + GW constraints and extrapolation to pQCD limit

A. Kurkela et al.: *Astroph. J.* 789 (2014) 127

A. Annala et al.: *PRL* 120 (2018) 172703

A. Vuorinen : *Nucl. Phys.* A982 (2019) 36



- Recent news: msec pulsar J0740+6620 with $M = (2.14 \pm 0.10) M_\odot$.

H.T. Cromartie et al.: *Nat. Astron.* 4 (2019) 72 ; arXiv:1904.06759 [astro-ph.HE]

NEUTRON STAR MATTER : Chiral FRG Equation of State

- Chiral Nucleon-Meson Field Theory & Functional Renormalisation Group

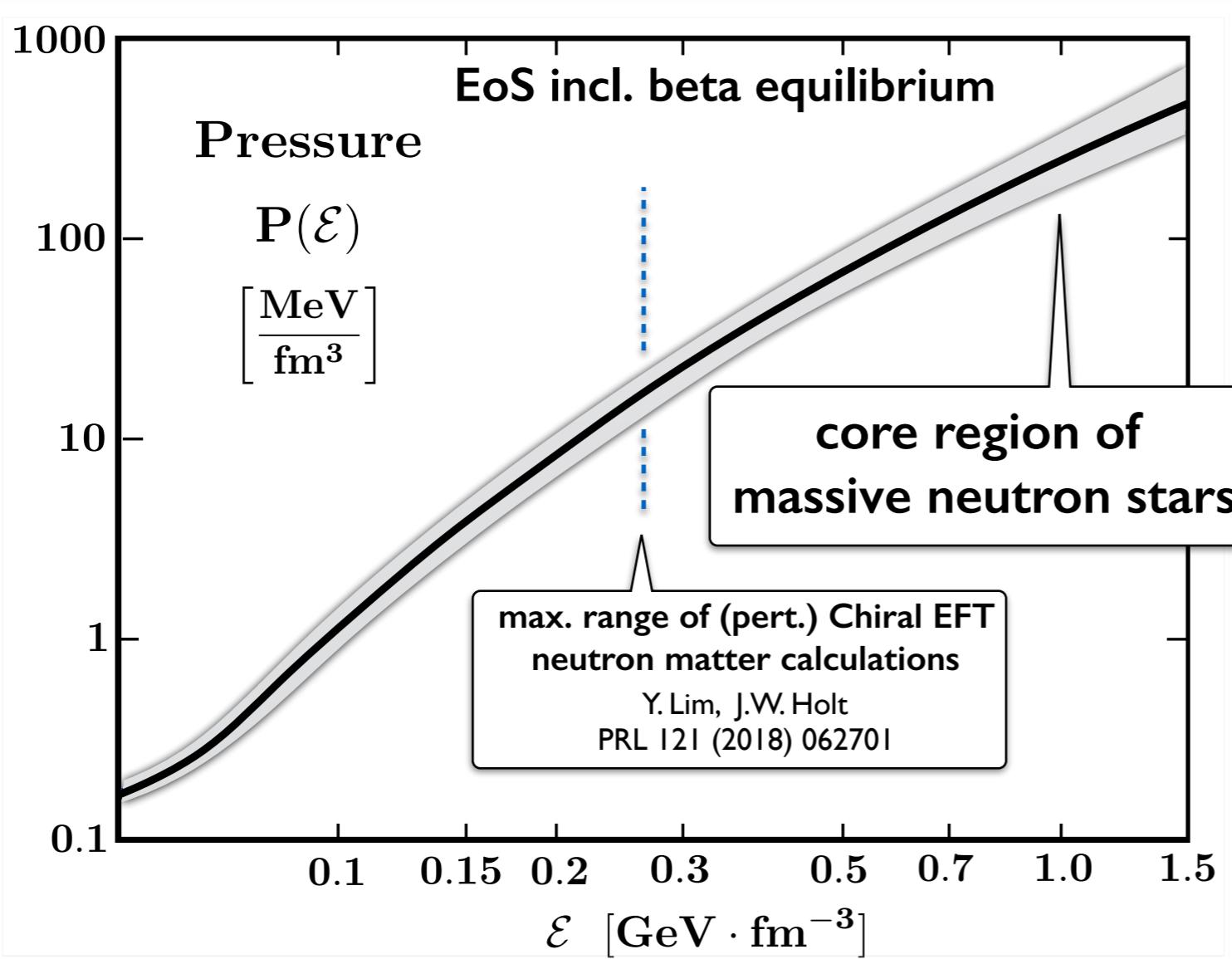
M. Drews, W.W.

Phys. Rev. C91 (2015) 035802

Prog. Part. Nucl. Phys. 93 (2017) 69

- Consistent with (perturbative) ChEFT at low densities

- Reproduces nuclear matter properties (incl. liquid-gas phase transition, ...)



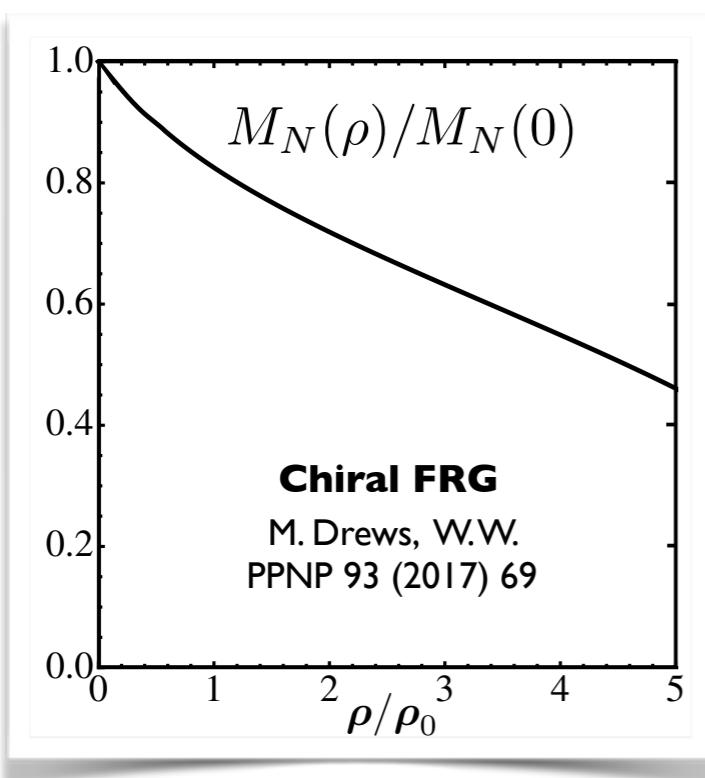
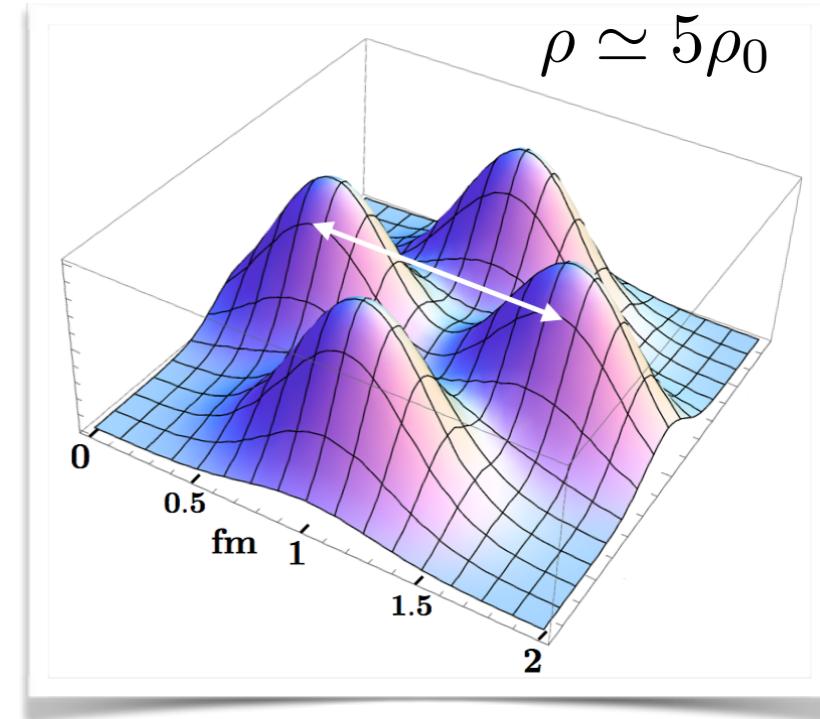
Non-perturbative FRG calculations :

- multi-pion exchange processes
- nucleon-hole excitations
- multi-nucleon correlations
- NO chiral phase transition**

DENSE BARYONIC MATTER in NEUTRON STARS as a RELATIVISTIC FERMI LIQUID

B. Friman, W.W. Rhys. Rev. C100 (2019) 065807

- Neutron Star Matter : Fermi liquid dominantly neutrons + ca. 5 % protons
- **Quasiparticles:**
nucleons “dressed” by their **strong interactions**
... and imbedded in **multi-pion field**
- Neutrons as relativistic quasiparticles
Fermi momentum in neutron star core region :



$$p_F = (3\pi^2 \rho)^{1/3} \sim 0.5 - 0.6 \text{ GeV/c} \gtrsim M_N(\rho)$$

- **Landau effective mass**
- **Chemical potential**

$$\mu = \frac{\partial \mathcal{E}}{\partial \rho} = M_0 + \left(1 + \rho \frac{\partial}{\partial \rho}\right) \frac{E(\rho)}{A} = m^*(\rho) + U(\rho)$$

repulsive many-body correlations

Relativistic Landau Fermi-Liquid Theory (contd.)

- Quasiparticle interaction expanded in Legendre series
Coefficients: Landau Fermi-liquid parameters

- Leading Landau parameters ($T=0$): incompressibility

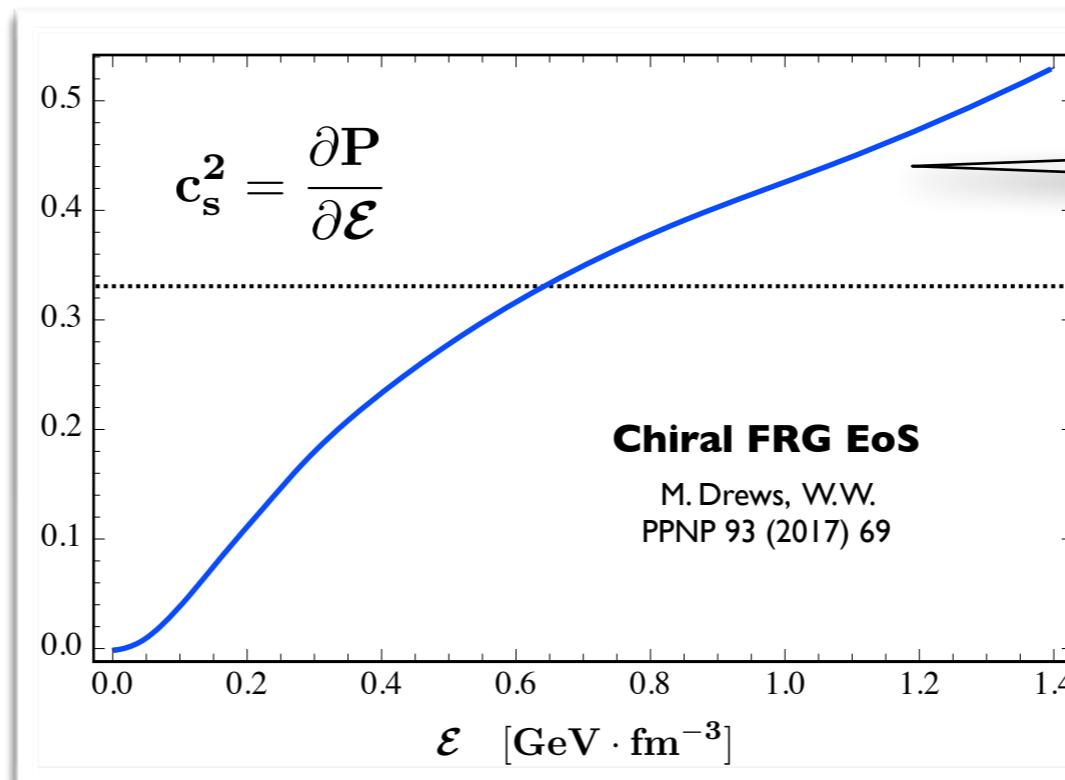
$$1 + F_0(\rho) = \frac{p_F m^*(\rho)}{\pi^2} \frac{\partial^2 \mathcal{E}(\rho)}{\partial \rho^2} = \frac{m^*(\rho)}{3p_F^2} \mathcal{K}(\rho)$$

$$1 + \frac{F_1(\rho)}{3} = \frac{m^*(\rho)}{\mu}$$

Landau effective mass

- Speed of sound

$$c_s^2 = \frac{1}{3} \left(\frac{p_F}{\mu} \right)^2 \frac{1 + F_0}{1 + \frac{1}{3} F_1}$$



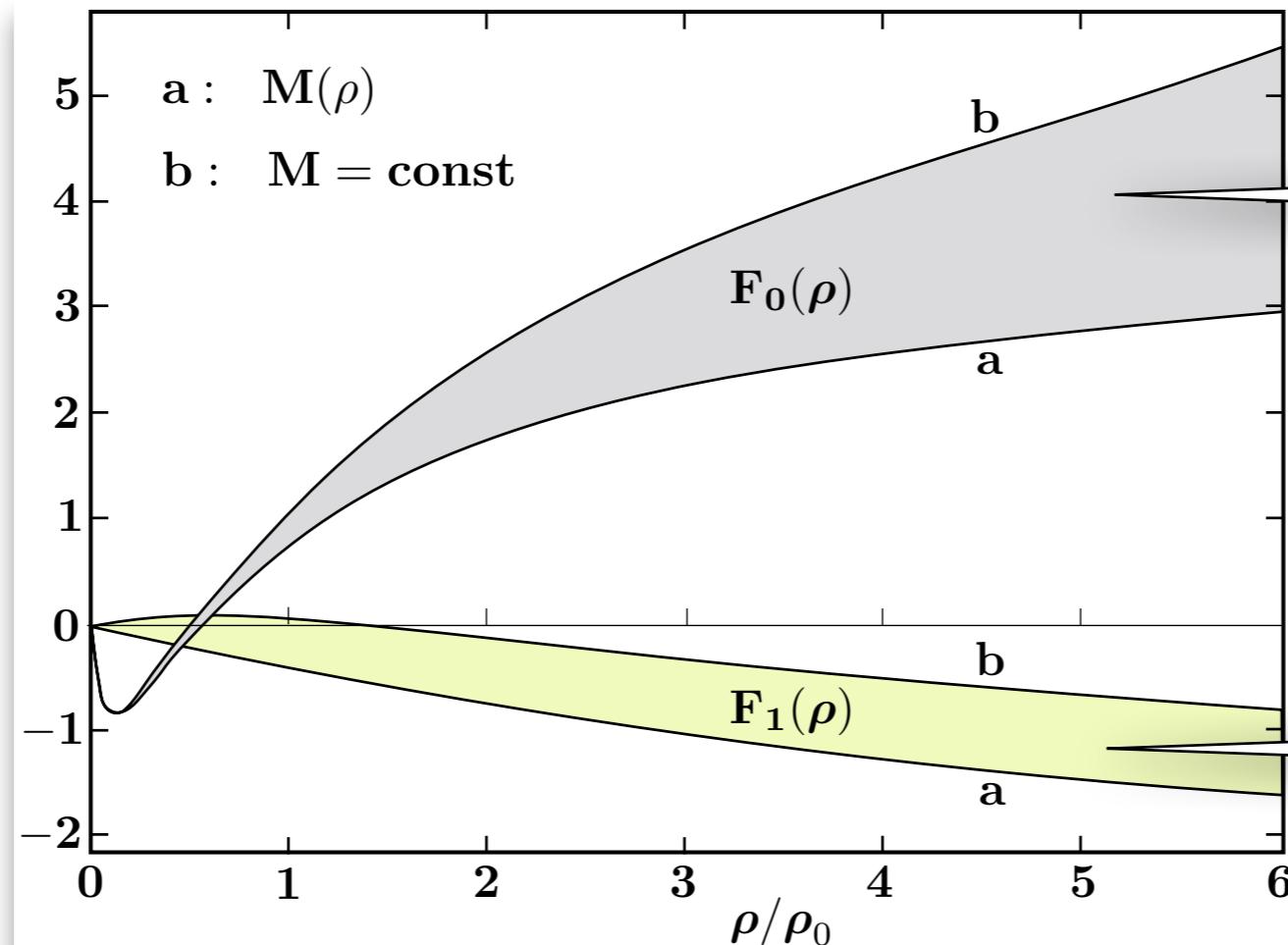
strongly repulsive correlations

$$c_s > \frac{1}{\sqrt{3}}$$

LANDAU FERMI-LIQUID PARAMETERS

- Deduced from neutron star matter EoS calculated using Chiral Nucleon-Meson Field Theory & Functional Renormalisation Group

neutron star matter : strongly correlated Fermi liquid ... but not “extreme” !



strongly repulsive many-body correlations

decreasing effective Fermion mass

B. Friman, W.W. Phys. Rev. C100 (2019) 065807

- Low densities $\rho \lesssim \rho_0$: good agreement with (perturbative) ChEFT results

J.W. Holt, N.Kaiser, W.W.: Phys. Rev. C87 (2013) 014338

LANDAU FERMI-LIQUID PARAMETERS (contd.)

- Comparison with atomic liquid helium-3 in its normal phase at low temperature (3 K)
- Interaction between He-3 atoms:
attractive van der Waals potential plus strongly repulsive short-range core
- Landau parameters of liquid helium-3 at pressures P = (0 - 30) bar:
 $F_0(^3\text{He}) \sim 10 - 70$ $F_1(^3\text{He}) \sim 5 - 13$

... much larger by order of magnitude than
Landau parameters of neutron star matter !!

- Intermediate conclusion :
**Neutron star matter at central densities
is a strongly correlated Fermi system**
... but not as extreme as one might have thought !

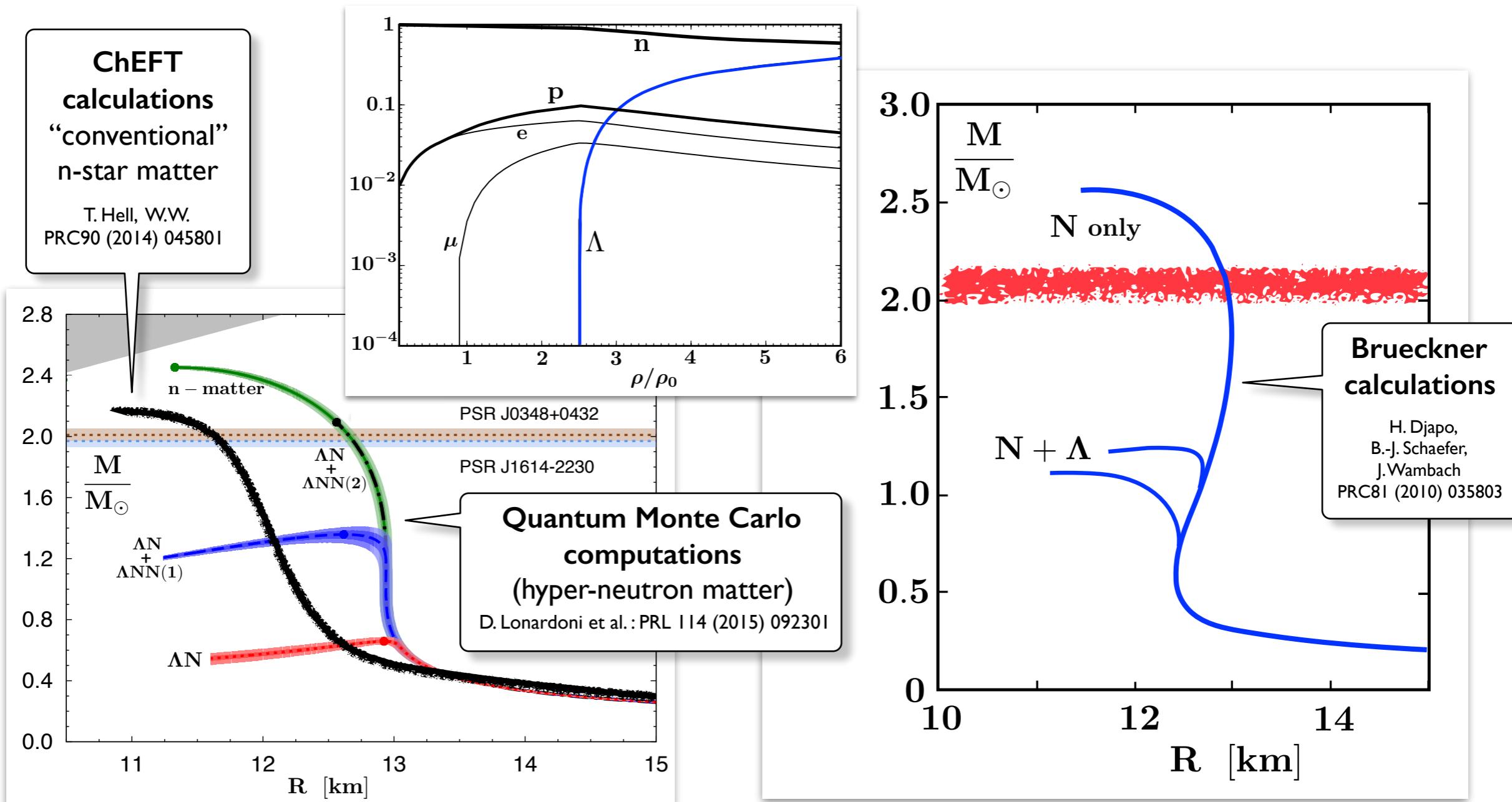
Part II.

Hyperon-Nuclear Interactions and Strangeness in Dense Matter

- Chiral SU(3) Effective Field Theory of Hyperon-Nucleon Interactions
- Hyperon-NN Three-Body Forces
- “Hyperon Puzzle” in Neutron Stars ?

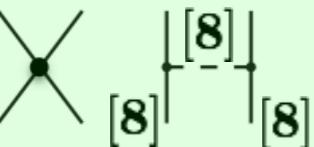
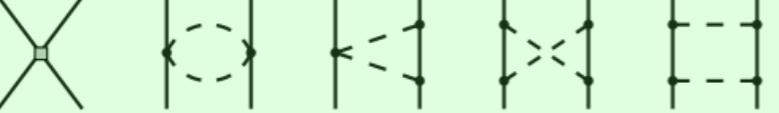
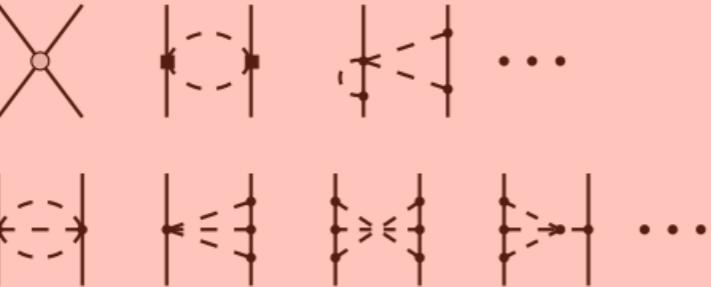


NEUTRON STAR MATTER including HYPERONS

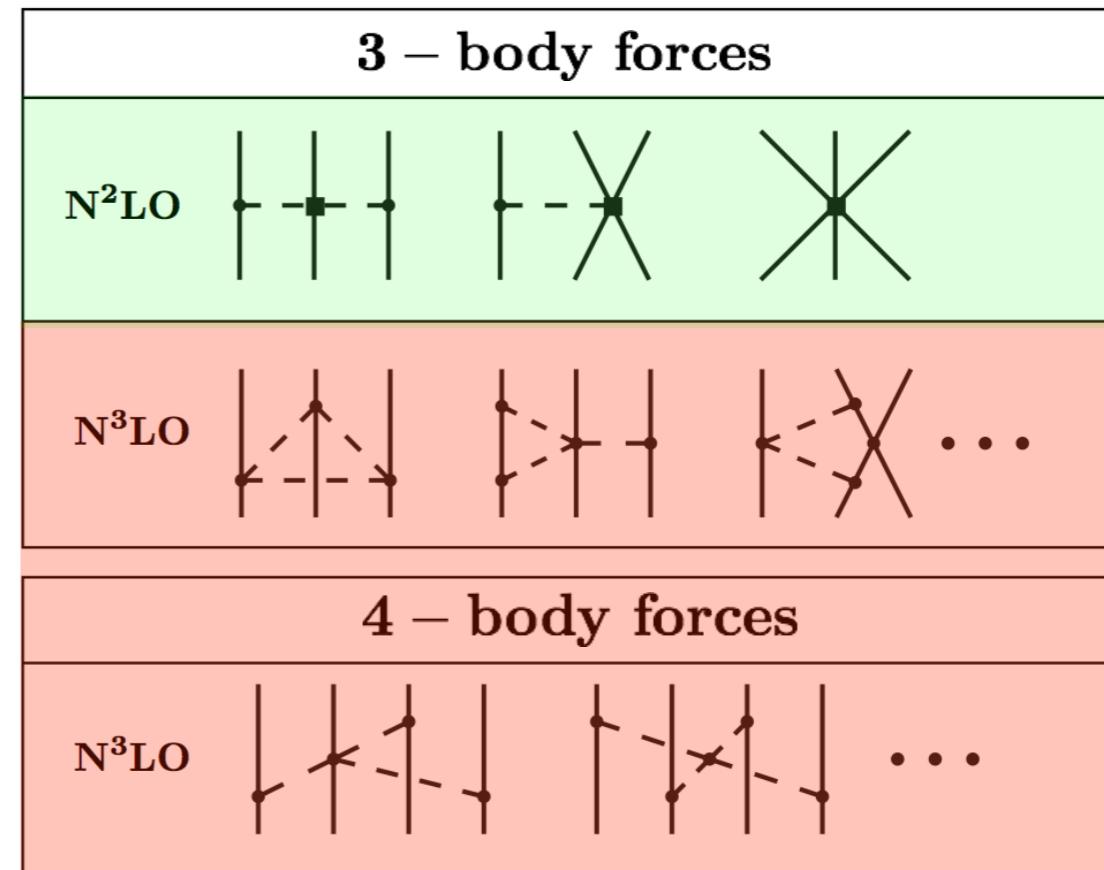


Inclusion of hyperons: EoS too soft to support 2-solar-mass n-stars
unless: strong repulsion in YN and/or $YNN \dots$ interactions

BARYON-BARYON INTERACTIONS from CHIRAL SU(3) EFFECTIVE FIELD THEORY

	BB interactions
LO	
NLO	
N^2LO	
N^3LO	

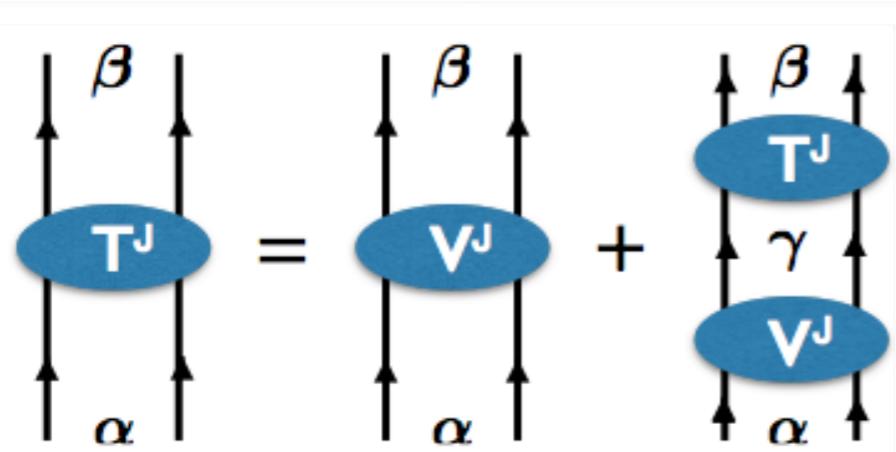
- Systematically organized hierarchy in powers of Q/Λ
(Q : momentum, energy, pion mass)



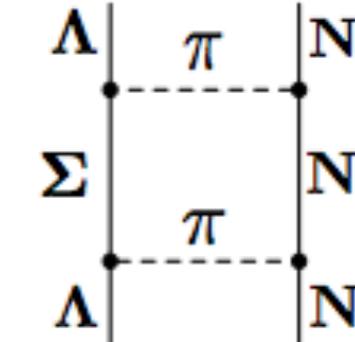
- NN interaction : has reached N^4LO level
- YN interaction : still very limited empirical data base
→ restriction to NLO plus YNN three-body forces



Coupled-Channels Lippmann-Schwinger Equation



example ΛN scattering
inclusion of
 $\Lambda N \leftrightarrow \Sigma N$
coupled channels



- Partial waves $(LS)J$, baryon-baryon channels α, β

$$T^J_{\beta\alpha}(p_f, p_i; \sqrt{s}) = V^J_{\beta\alpha}(p_f, p_i) + \sum_{\gamma} \int_0^{\infty} \frac{dp}{(2\pi)^3} p^2 V^J_{\beta\gamma}(p_f, p) \frac{2\mu_{\gamma}}{p_{\gamma}^2 - p^2 + i\varepsilon} T^J_{\gamma\alpha}(p, p_i; \sqrt{s})$$

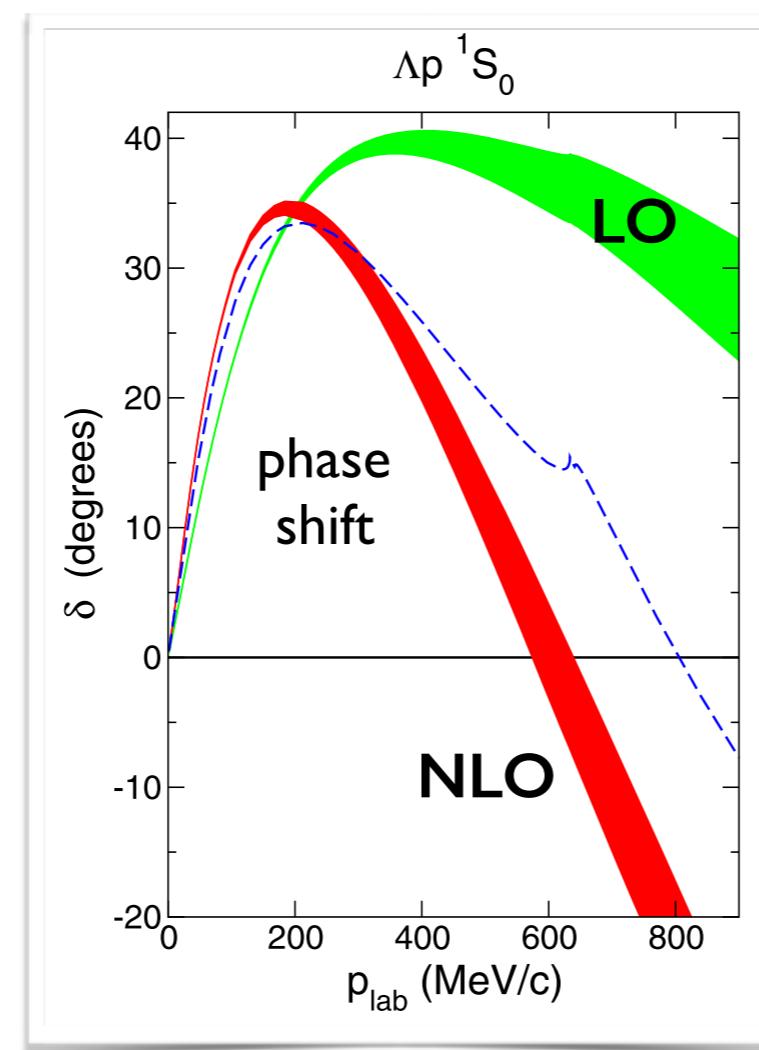
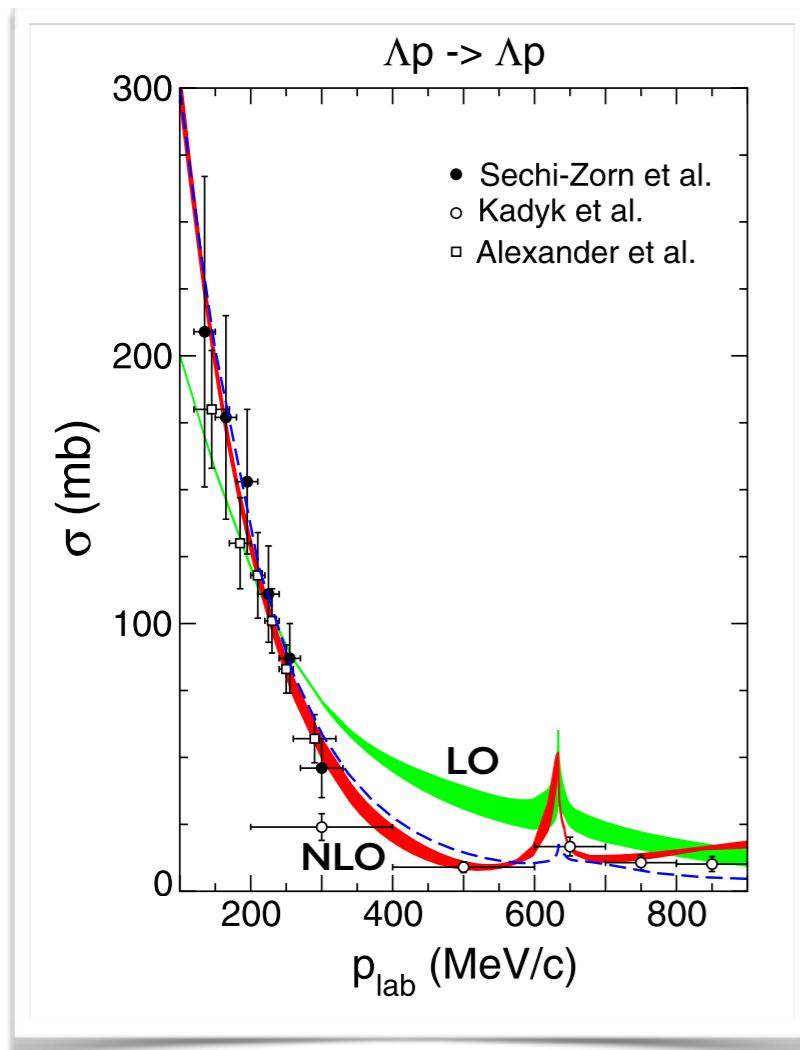
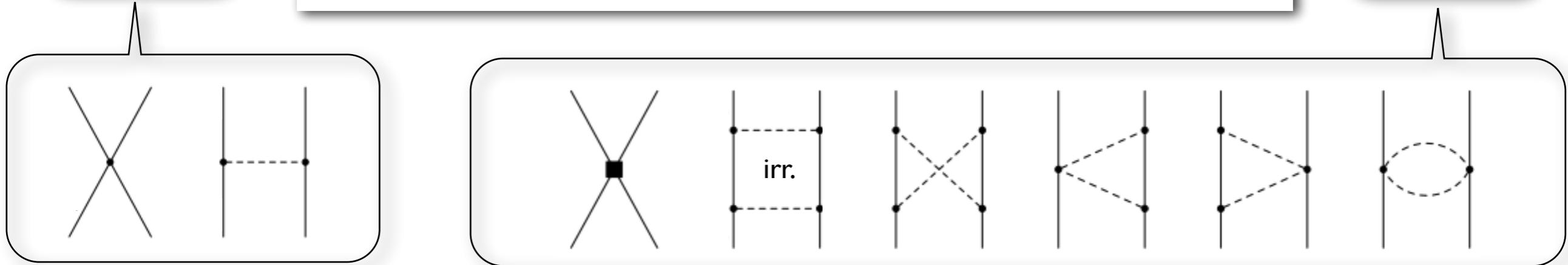
- Input V from chiral SU(3) EFT baryon-baryon interaction at NLO
 - Momentum of intermediate channels γ determined by :
- $$\sqrt{s} = \sqrt{M_{\gamma,1}^2 + p_{\gamma}^2} + \sqrt{M_{\gamma,2}^2 + p_{\gamma}^2}$$
- Momentum space cutoff: 0.5 - 0.6 GeV



LO

Hyperon - Nucleon Interaction

from CHIRAL SU(3) Effective Field Theory

NLO

- **moderate attraction at low momenta**
→ relevant for hypernuclei

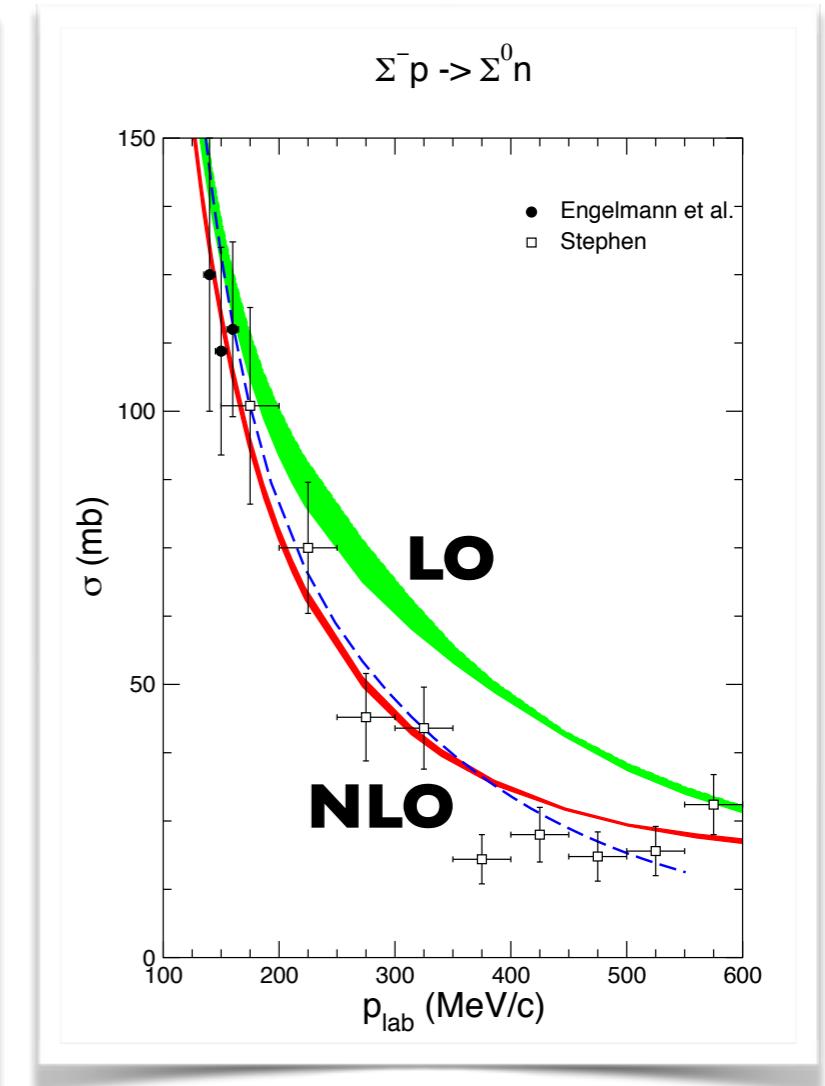
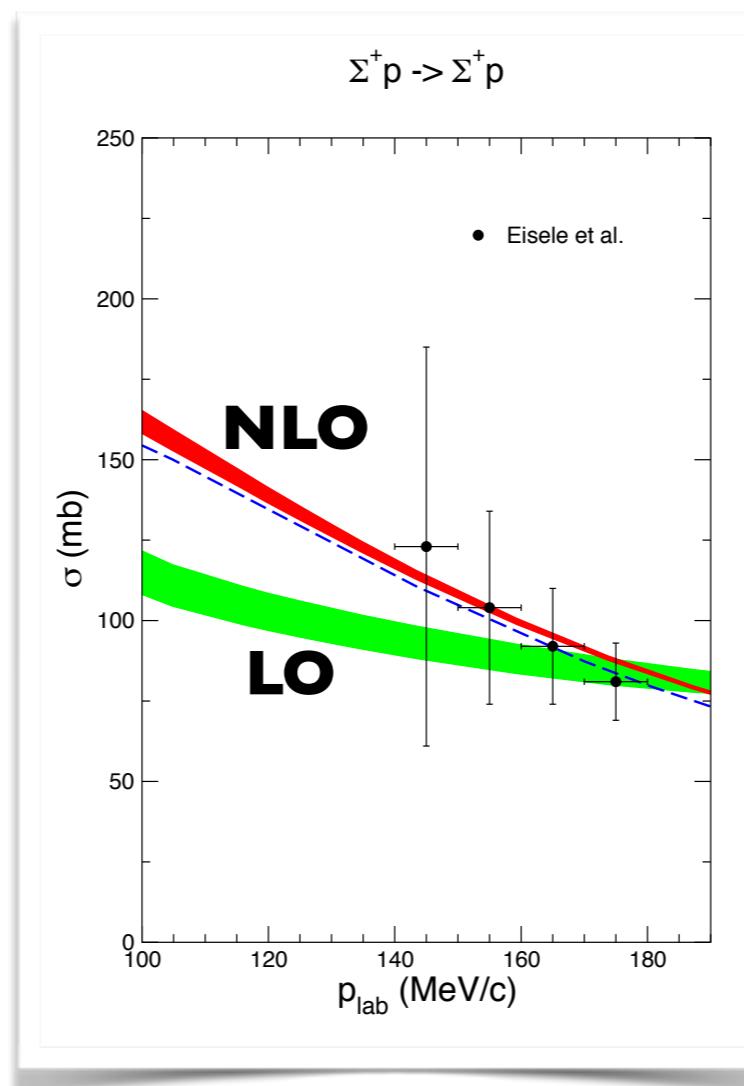
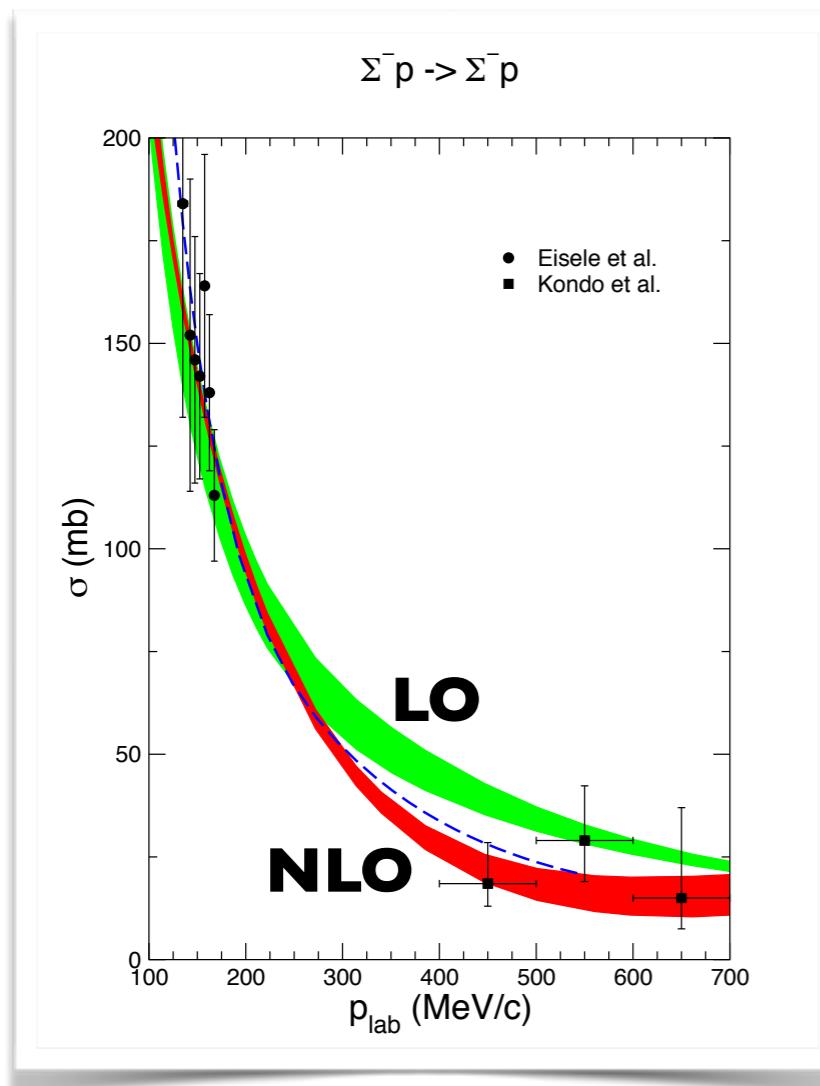
- **increasing repulsion at higher momenta**
→ relevant for dense baryonic matter



Hyperon - Nucleon Interaction (contd.)

J. Haidenbauer, S. Petschauer, N. Kaiser,
U.-G. Meißner, A. Nogga, W.W.
Nucl. Phys. A 915 (2013) 24

- ΣN elastic and charge exchange scattering

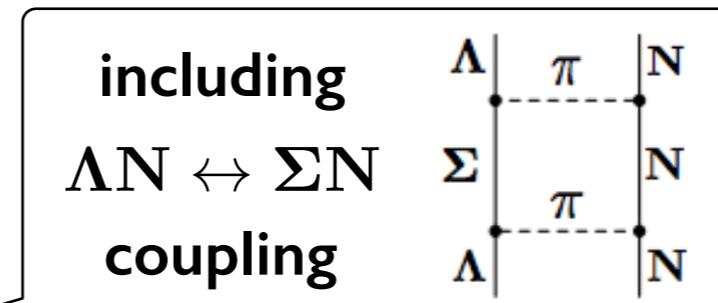
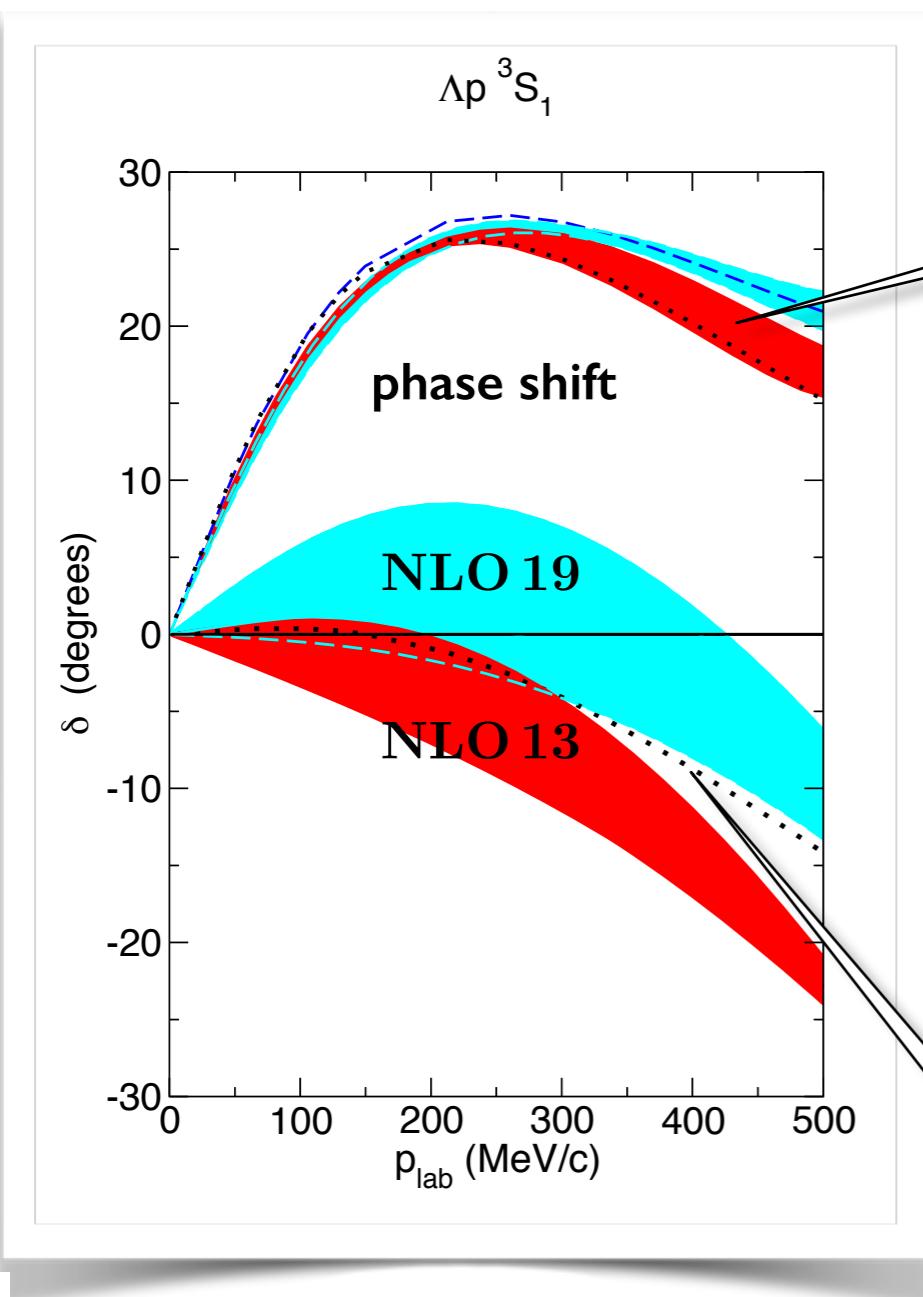


- Quest for much improved hyperon-nucleon scattering data base !

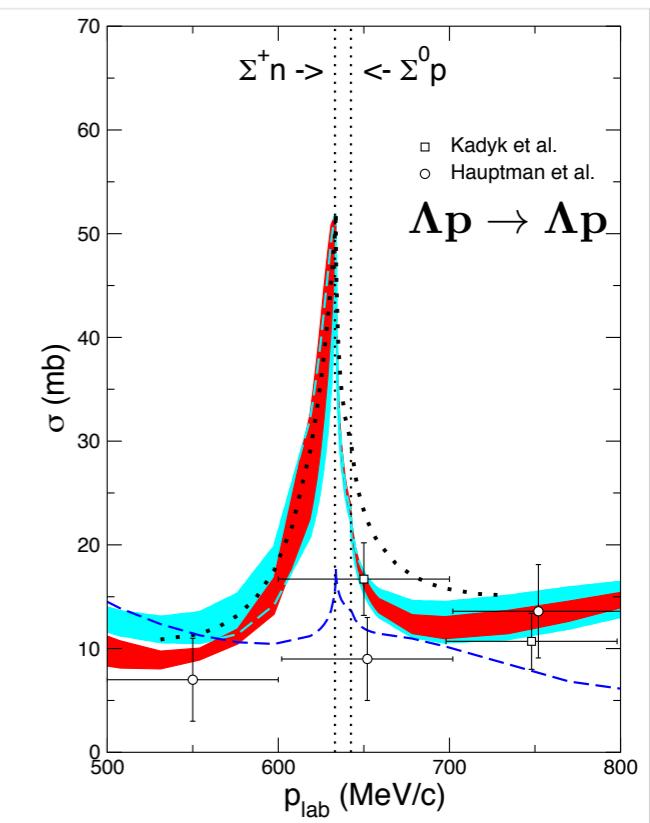
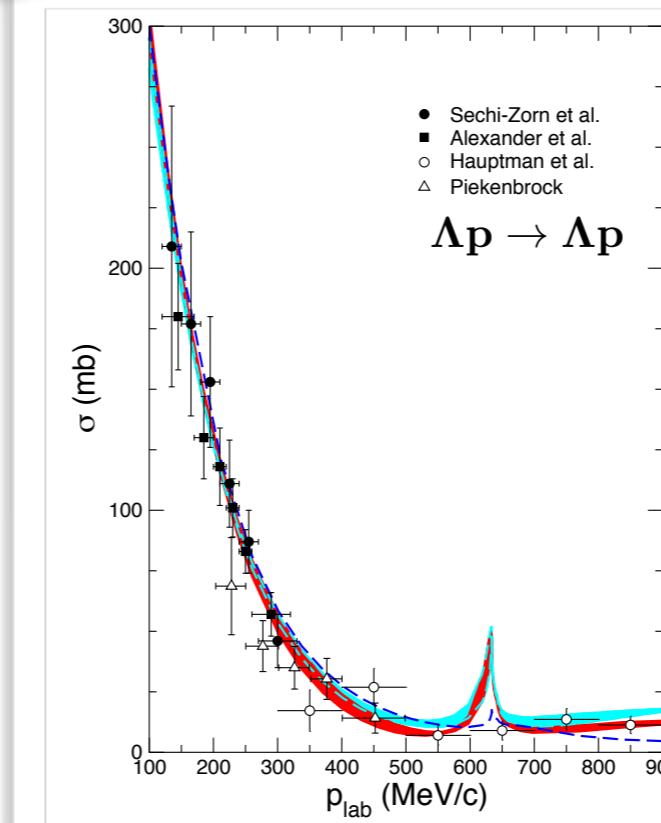
Hyperon - Nucleon Interaction : recent update

J. Haidenbauer, U.-G. Meißner, A. Nogga arXiv:1906.11168 [nucl-th]

- Reduced no. of independent contact terms at NLO by imposing symmetries connecting NN and YN S-waves



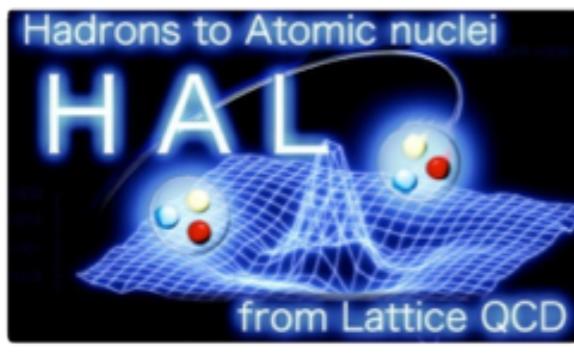
blue : NLO 19
red : NLO 13



without $\Lambda N \leftrightarrow \Sigma N$ coupling



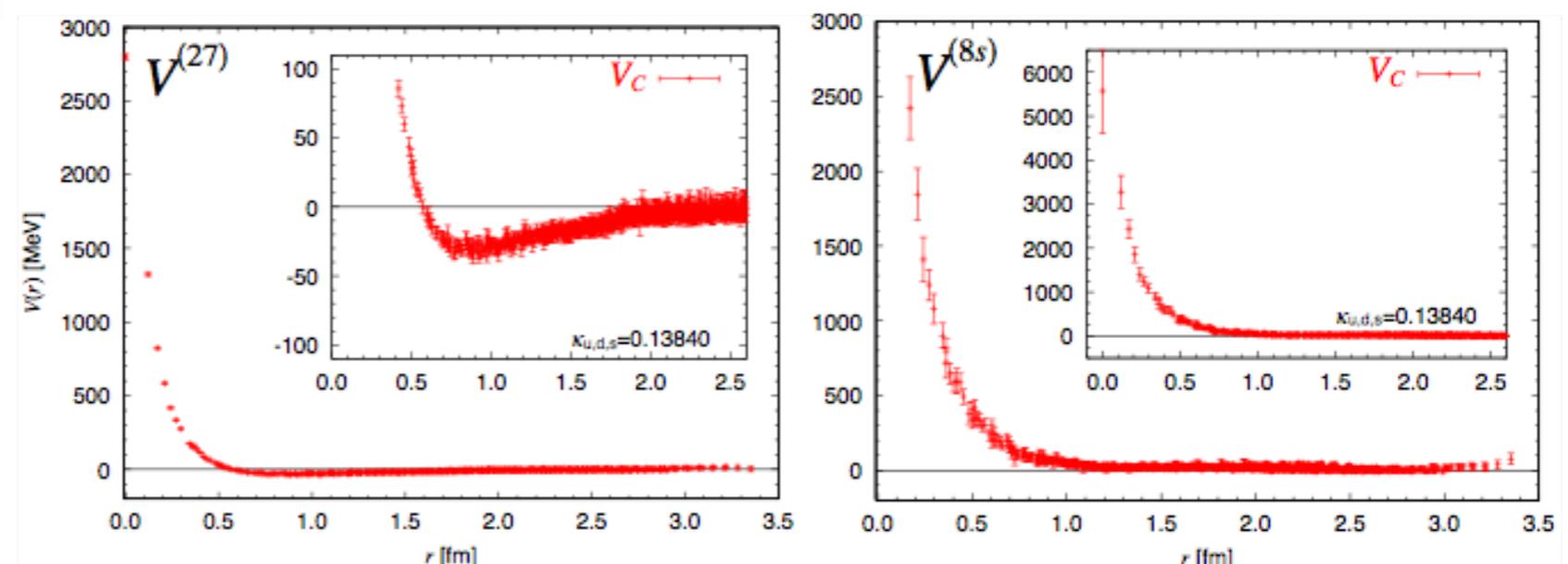
Hyperon - Nucleon Interactions from Lattice QCD



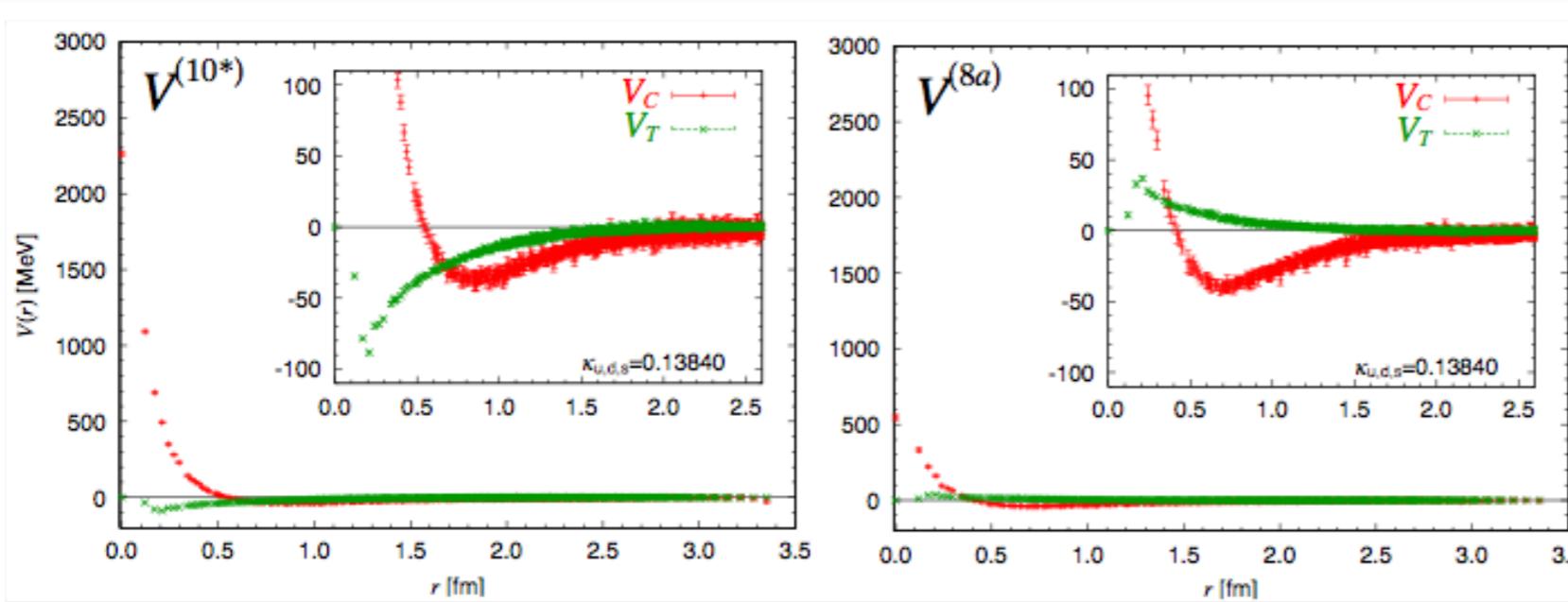
$$m_{ps} = 0.47 \text{ GeV}$$

T. Inoue et al.
 PTP 124 (2010) 591
 Nucl. Phys. A881 (2012) 28

$$\Lambda N(^1S_0) = \frac{9}{10}[27] + \frac{1}{10}[8_s]$$



$$\Lambda N(^3S_1) = \frac{1}{2}[10^*] + \frac{1}{2}[8_a]$$



**Short-range
repulsive core
in all channels**

... towards
physical
quark masses

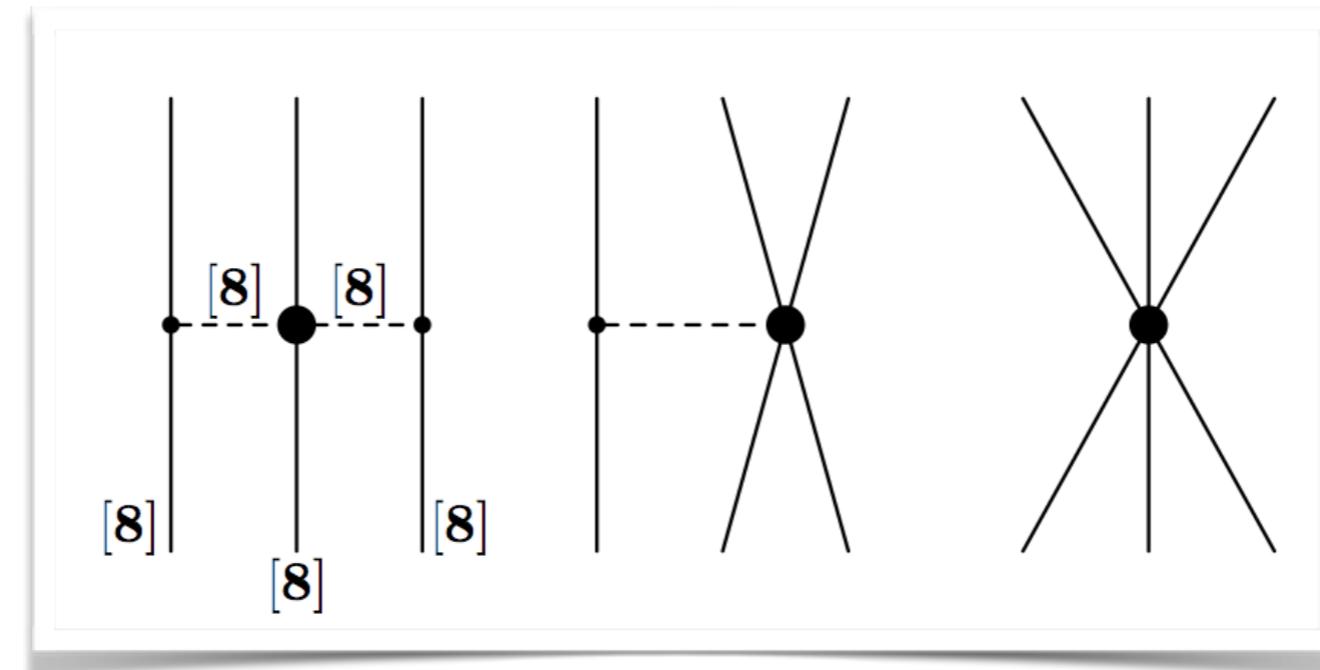
HYPERON - NUCLEON - NUCLEON THREE-BODY FORCES from CHIRAL SU(3) EFT

S. Petschauer, N. Kaiser, J. Haidenbauer, U.-G. Meißner, W.W.

Phys. Rev. C93 (2016) 014001

- Chiral $SU(3)_L \times SU(3)_R$ Effective Field Theory:
interacting pseudoscalar meson & baryon octets + contact terms

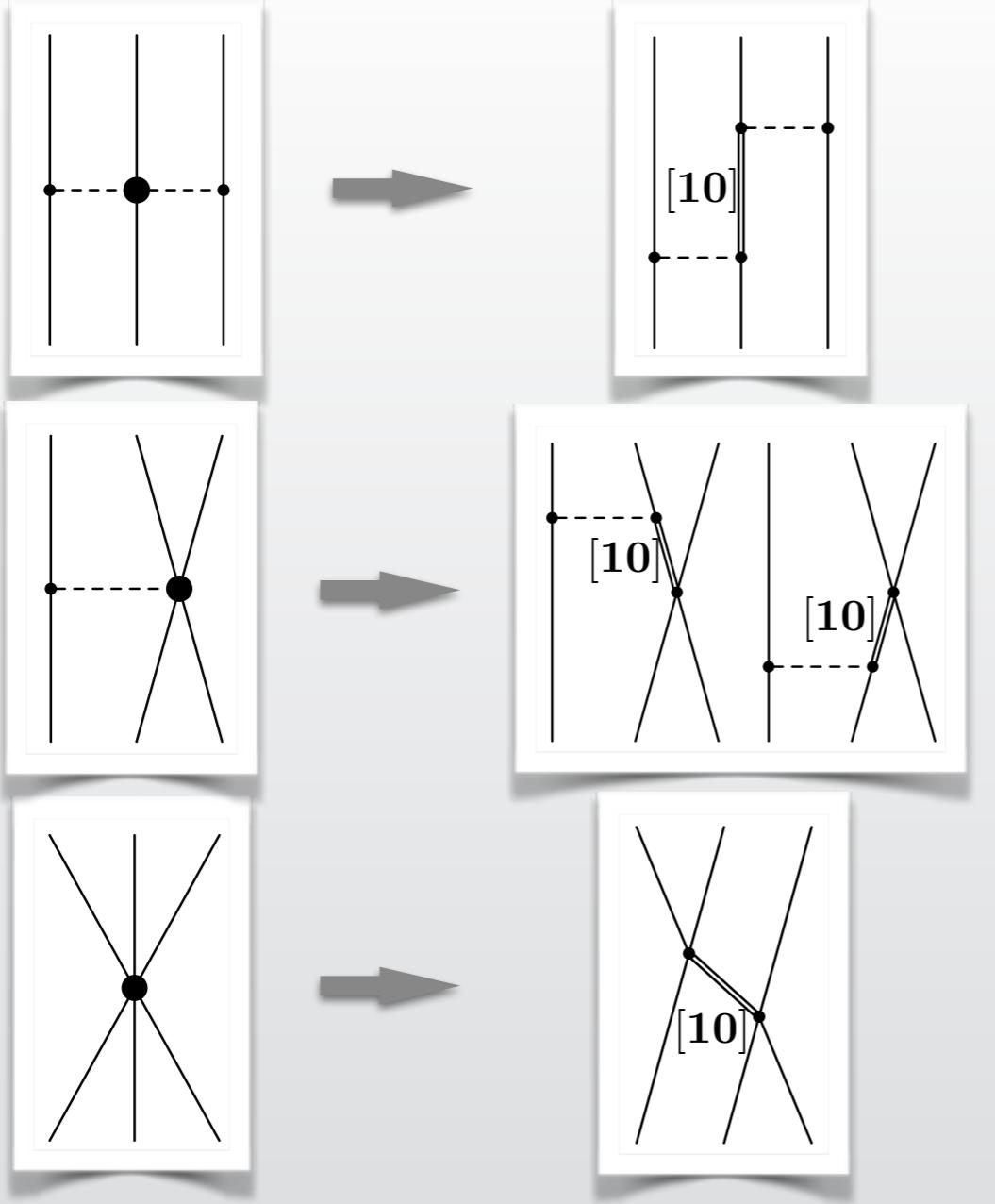
3-baryon
sector at
NNLO



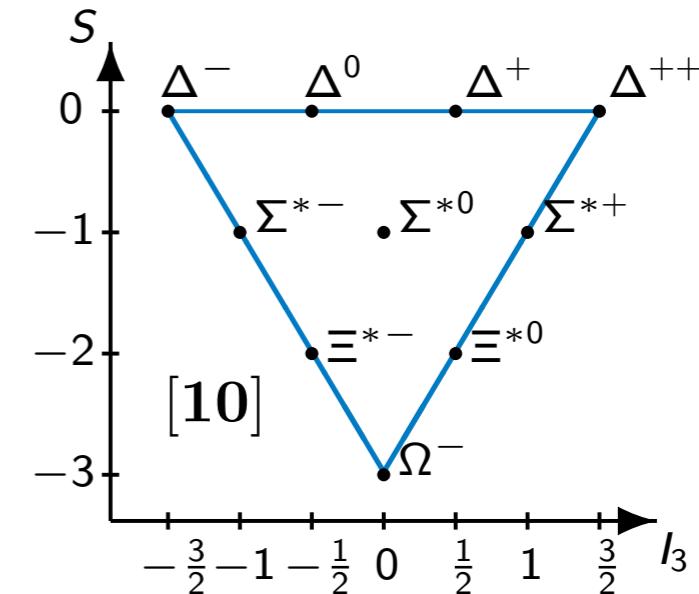
- $SU(3)$ symmetry used to restrict no. of independent parameters
- Further reduction through decuplet dominance

Decuplet Dominance in YNN three-body forces

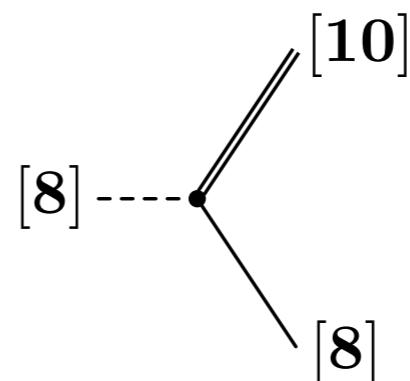
- Estimates of YNN 3-body interactions assuming dominant decuplet (Σ^* , Δ) intermediate states



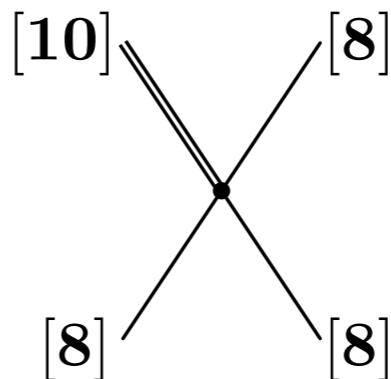
promotion from **NNLO** to **NLO**



- ... much reduced set of parameters -
Basic vertices :



One constant
($C = \frac{3}{4}g_A \approx 1$ from $\Delta \rightarrow N\pi$)



Two constants
(H_1, H_2)
(Typical magnitude $|H_i| \sim f_\pi^{-2}$)

Pauli-forbidden
in NN sector

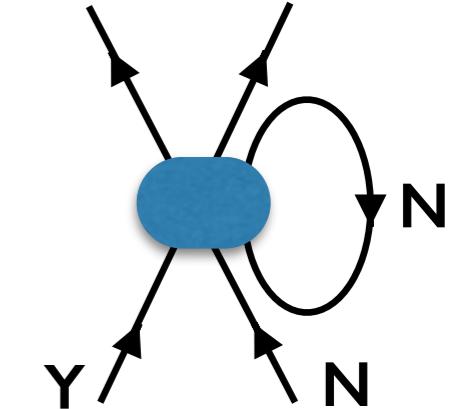
Density-dependent EFFECTIVE HYPERON - NUCLEON INTERACTION from CHIRAL THREE-BARYON FORCES

S. Petschauer, J. Haidenbauer, N. Kaiser, U.-G. Meißner, W.W.

Nucl. Phys. A957 (2017) 347

$$V_{12}^{\text{eff}} = \sum_B \text{tr}_{\sigma_3} \int_{|\vec{k}| \leq k_f^B} \frac{d^3 k}{(2\pi)^3} V_{123}$$

- Example: **Λ -neutron density-dependent effective interaction in a nuclear medium (protons + neutrons)**



$$V_{\Lambda n}^{\text{eff}, \pi\pi} = \frac{C^2 g_A^2}{2f^4 \Delta} [\rho_n + 2\rho_p] + \mathcal{F}(k_F^p, k_F^n; p, q) \quad \text{repulsive}$$

$$V_{\Lambda n}^{\text{eff}, \pi} = \frac{CH g_A}{9f^2 \Delta} [\rho_n + 2\rho_p] + \mathcal{G}(k_F^p, k_F^n; p, q) \quad +/-$$

$$V_{\Lambda n}^{\text{eff}, ct} = \frac{H^2}{18\Delta} [\rho_n + 2\rho_p] \quad (H = H_1 + 3H_2) \quad \text{repulsive}$$

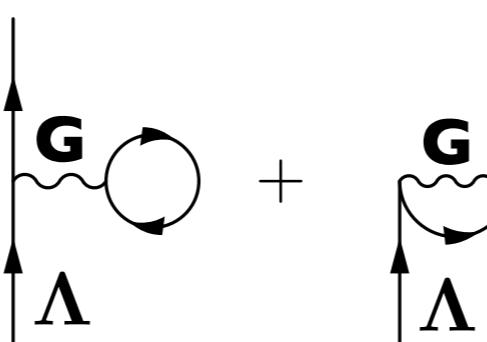
- Decuplet-octet mass difference $\Delta = M_{[10]} - M_{[8]} = 270 \text{ MeV}$

- Coupling parameters : $C = \frac{3}{4}g_A \simeq 1 \quad -\frac{1}{f^2} \lesssim H \lesssim +\frac{1}{f^2}$ (dim. arguments natural size)

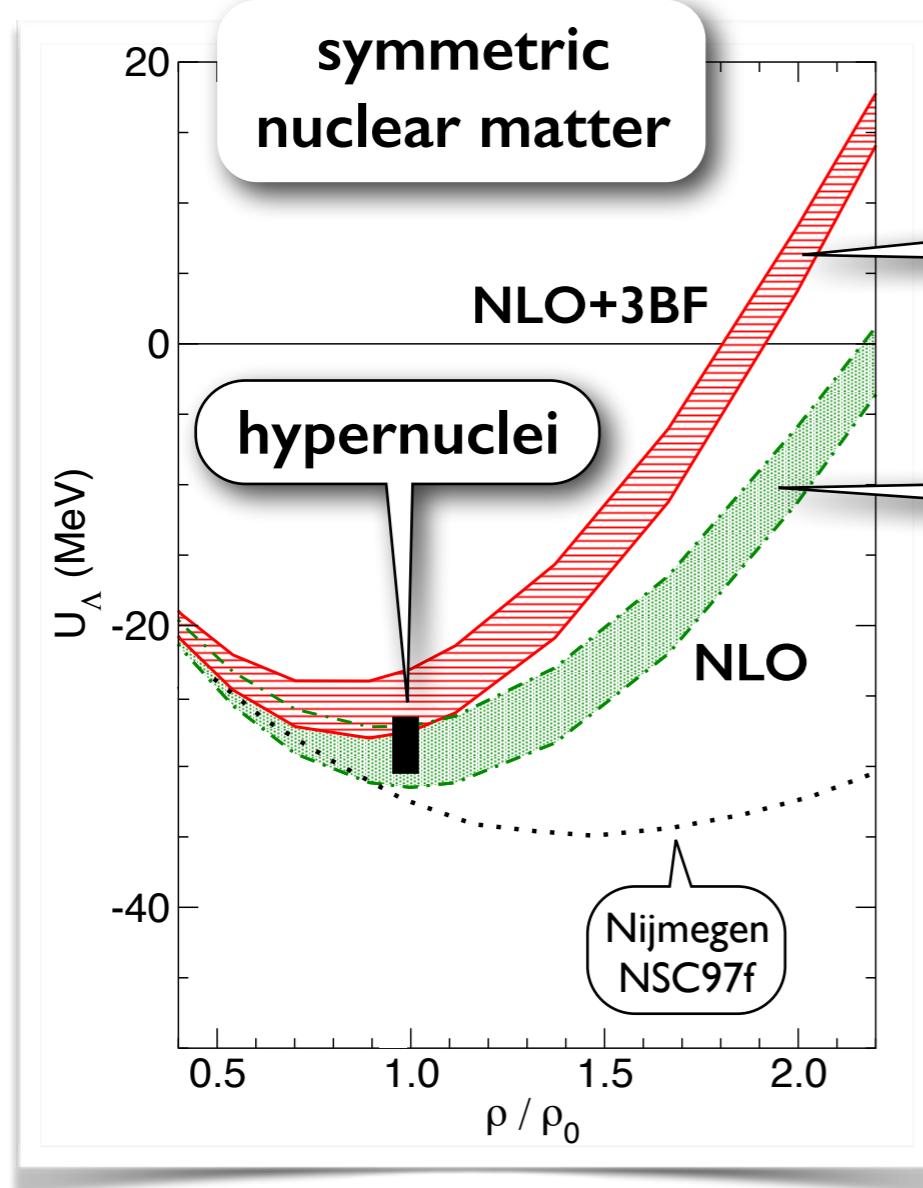


Density dependence of Λ single-particle potential

- Brueckner calculations using chiral SU(3) interactions



$$G(\omega) = V + V \frac{Q}{e(\omega) + i\epsilon} G(\omega)$$



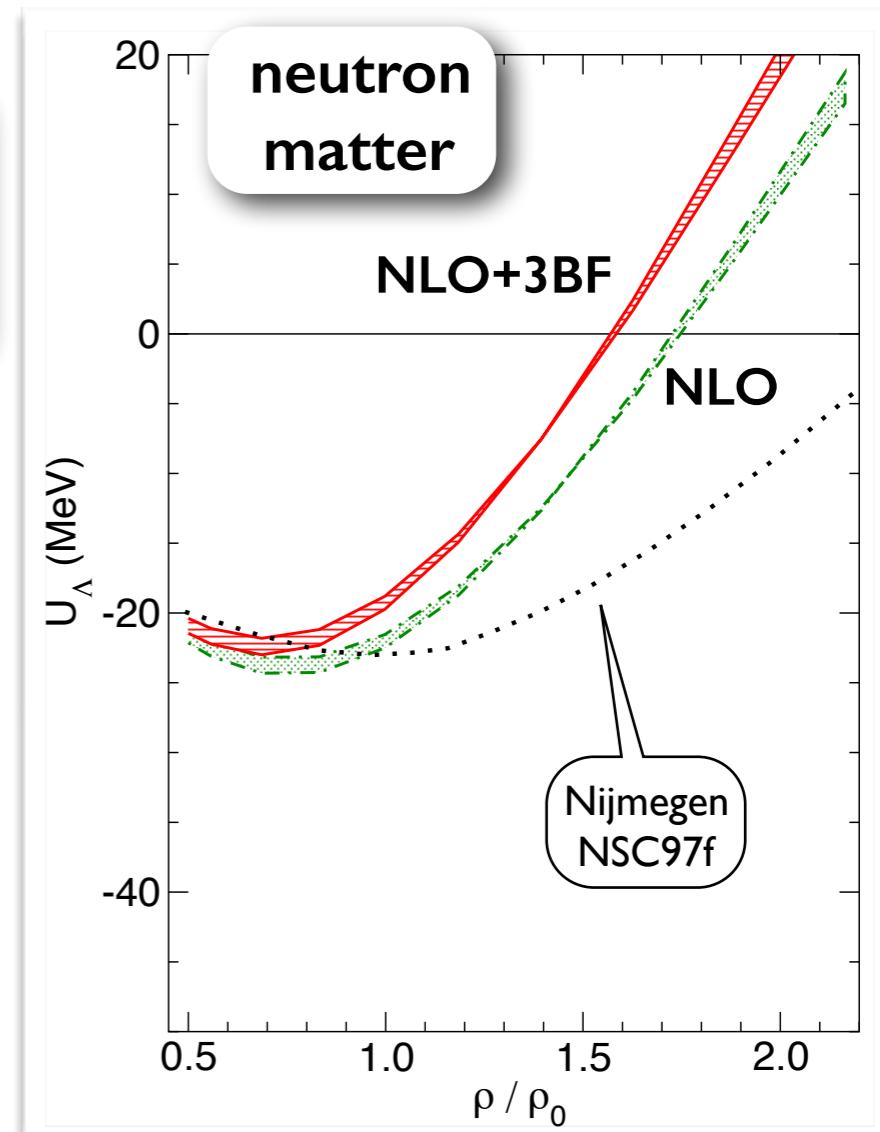
**Chiral SU(3)
2- and 3-body
forces**

2-body only

using NLO13:

J. Haidenbauer,
U.-G. Meißner,
N. Kaiser,
W.W.

Eur. Phys. J.
A53 (2017) 121



- ... towards a possible solution of the “hyperon puzzle” ?



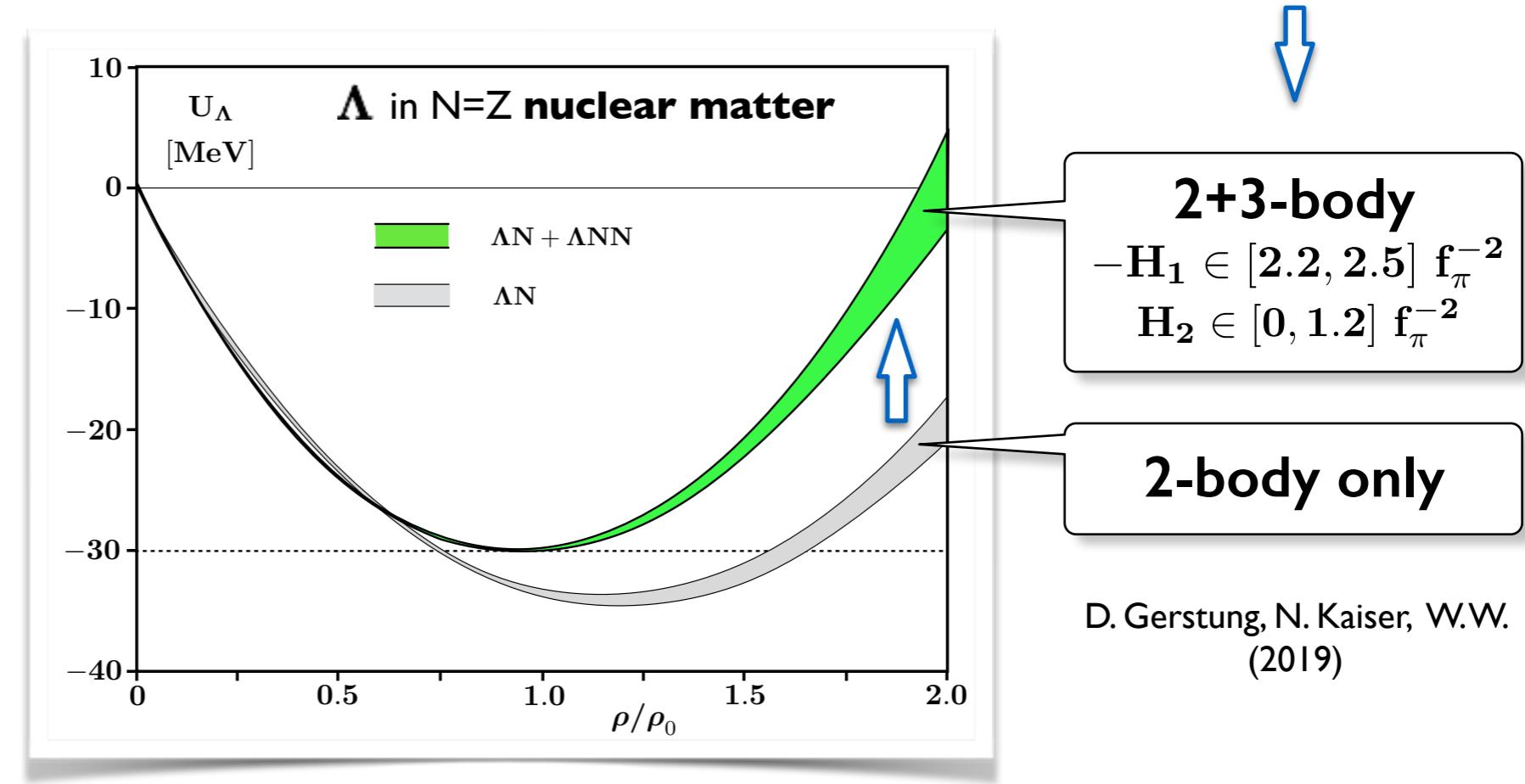
Density dependence of Λ single particle potential (contd.)

- Further extended Brueckner - Bethe - Goldstone calculations

$$G_{\alpha\beta}(\omega; \rho) = V_{\alpha\beta}(\rho) + V_{\alpha\gamma}(\rho) \frac{Q}{e(\omega) + i\epsilon} G_{\gamma\beta}(\omega; \rho)$$

- Coupled-channels G-matrix including explicit Λ NN \leftrightarrow Σ NN three-body interactions
- Chiral NN (N3LO) + YN (NLO) interactions + NNN & YNN 3-body forces

Strong additional repulsion from YNN three-body forces

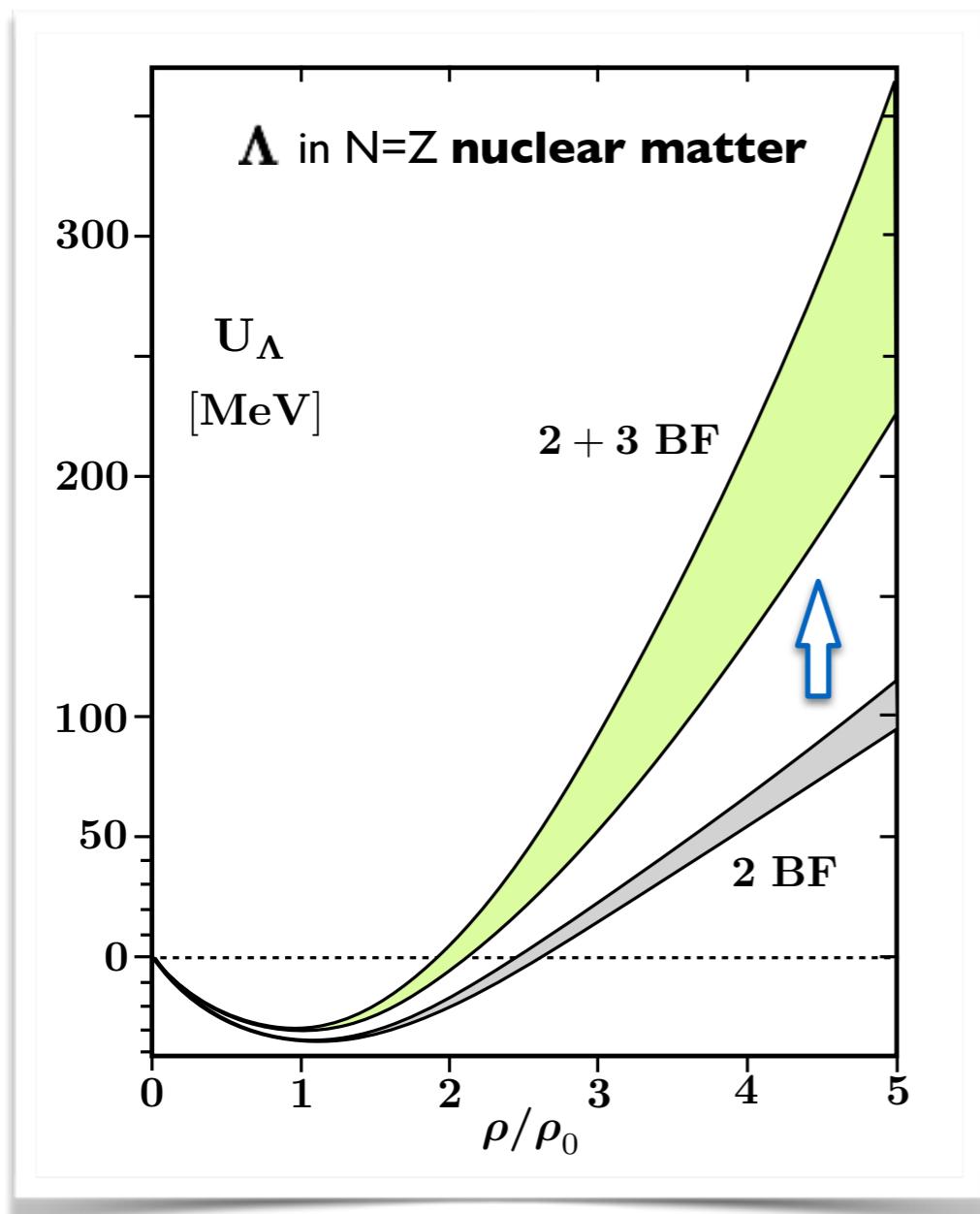


- Constrained by hypernuclear physics : $U_\Lambda(\rho = \rho_0) \simeq -30$ MeV

A. Gal, E. Hungerford, D. Milner
Rev. Mod. Phys. 88 (2016) 035004

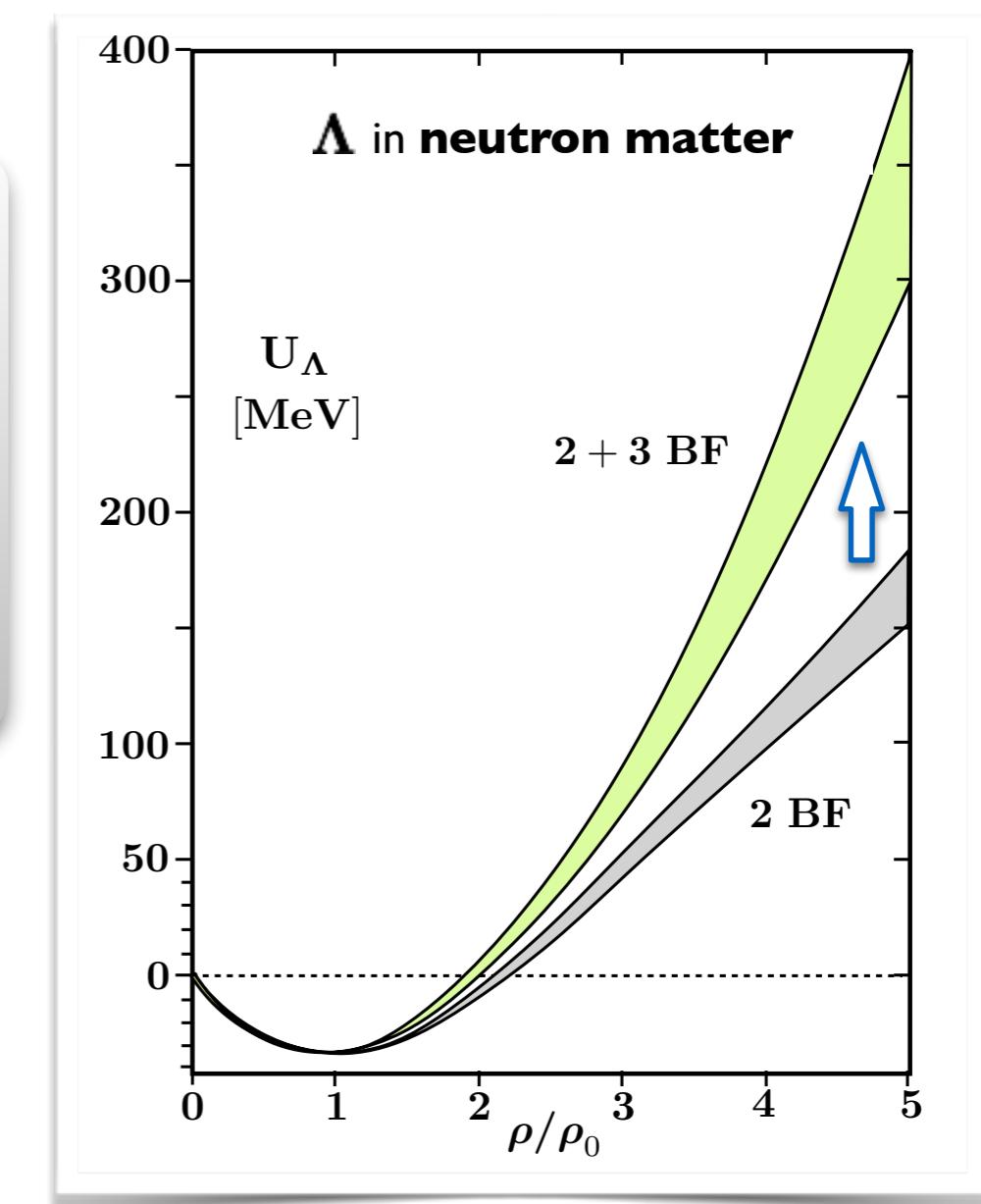
Density dependence of Λ single-particle potential (contd.)

- Extrapolations to high baryon densities
- Chiral NN (N3LO) + YN (NLO) interactions + NNN & YNN 3-body forces
- Coupled-channels G-matrix including explicit Λ NN $\leftrightarrow \Sigma$ NN three-body interactions



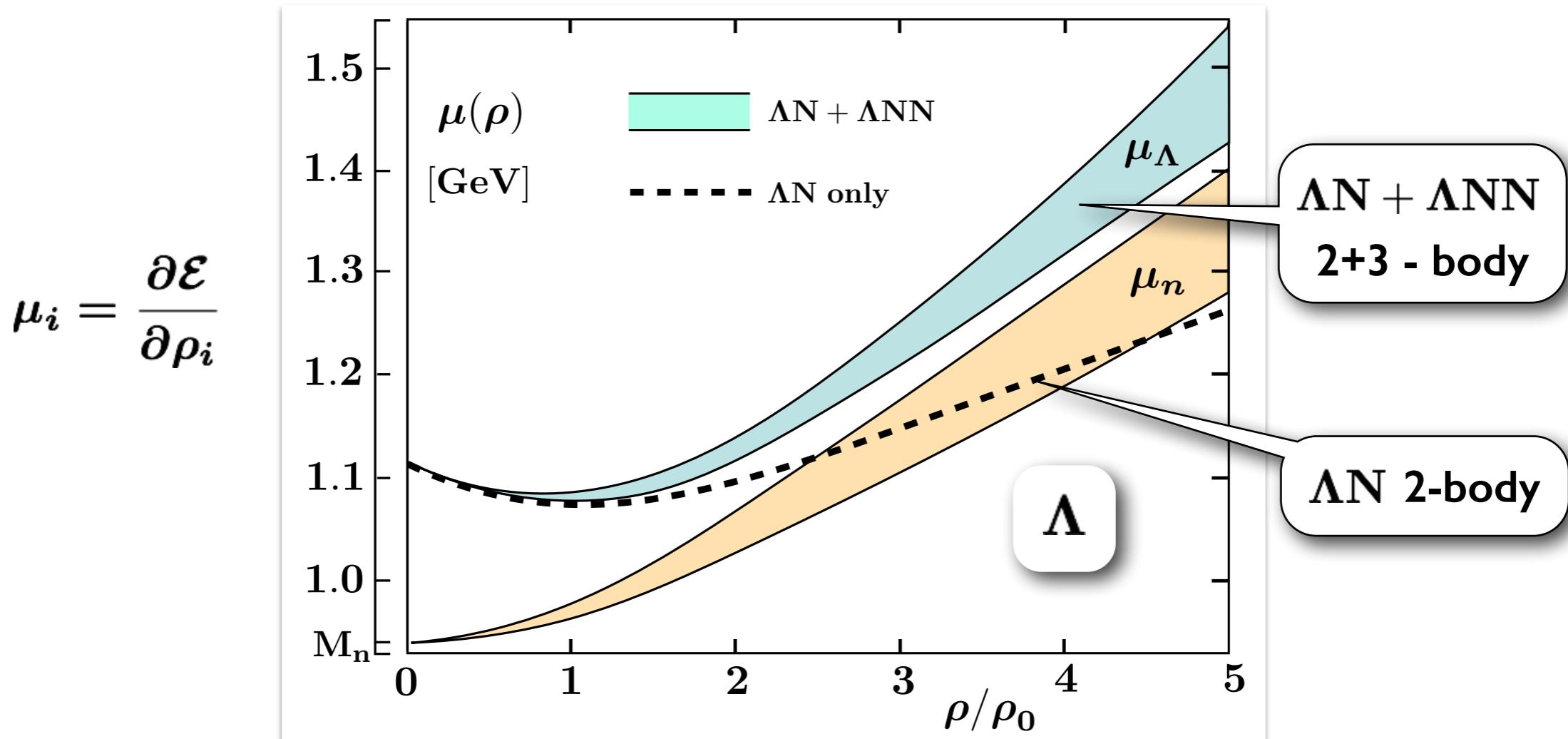
**Strong
additional
repulsion
from YNN
three-body
forces**

D. Gerstung,
N. Kaiser,
W.W.
(2019 / 2020)



Λ Hyperons in Neutron Stars ?

- Onset condition for appearance of Λ hyperons in neutron stars :
Equality of chemical potentials $\mu_\Lambda = \mu_n$



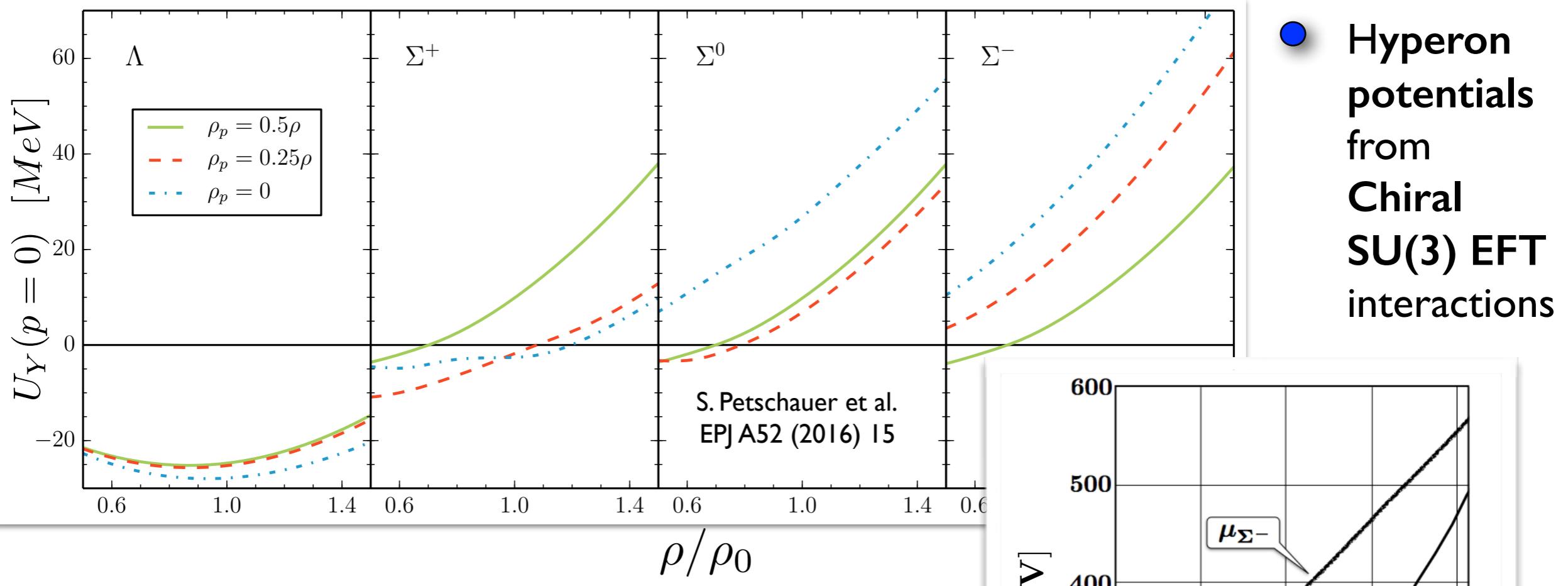
- Hyperon chemical potential in neutron matter from Chiral SU(3) EFT interactions

D. Gerstung, N. Kaiser, W.W. (2019 / 2020)

- Neutron chemical potential in neutron star matter from Chiral EFT + FRG equation-of-state

M. Drews, W.W. PPNP 93 (2017) 69

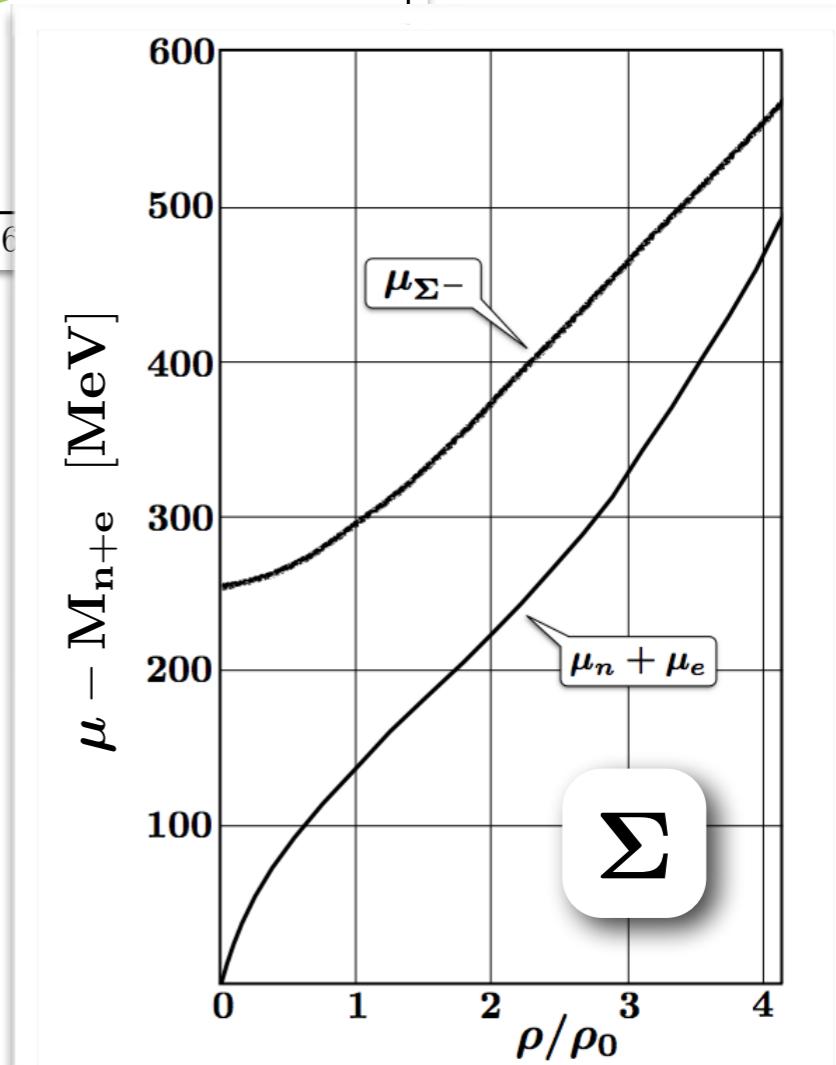
Σ Hyperons in Neutron Stars ?



- Σ -nuclear potentials are **repulsive**
- Condition for appearance of Σ^- in neutron star matter :

$$\mu_{\Sigma^-} = \mu_n + \mu_e = 2\mu_n - \mu_p$$

$$\mu_i = \frac{\partial \epsilon}{\partial \rho_i}$$



SUMMARY & OUTLOOK

- ★ **Neutron star matter:
EoS constrained by increasing amounts of observational data**
 - ▶ $M_{max} \simeq 2 M_\odot$, $R \sim 11 - 13 \text{ km}$ ($R_{1.4 M_\odot} = 11.0^{+0.9}_{-0.6} \text{ km}$) C.D. Capano et al., arXiv:1908.10352
 - ▶ **Tidal deformability from GW signals of binary n-star merger**
- ★ **ChEFT, (non-perturbative) FRG and Fermi-Liquid Theory**
 - ▶ Conventional EoS consistent with “stiffness” constraints
 - ▶ No chiral phase transition in n-matter up to at least $\rho > 5 \rho_0$
 - ▶ **New keywords: continuity of quark and hadron matter ? Superfluidity...**

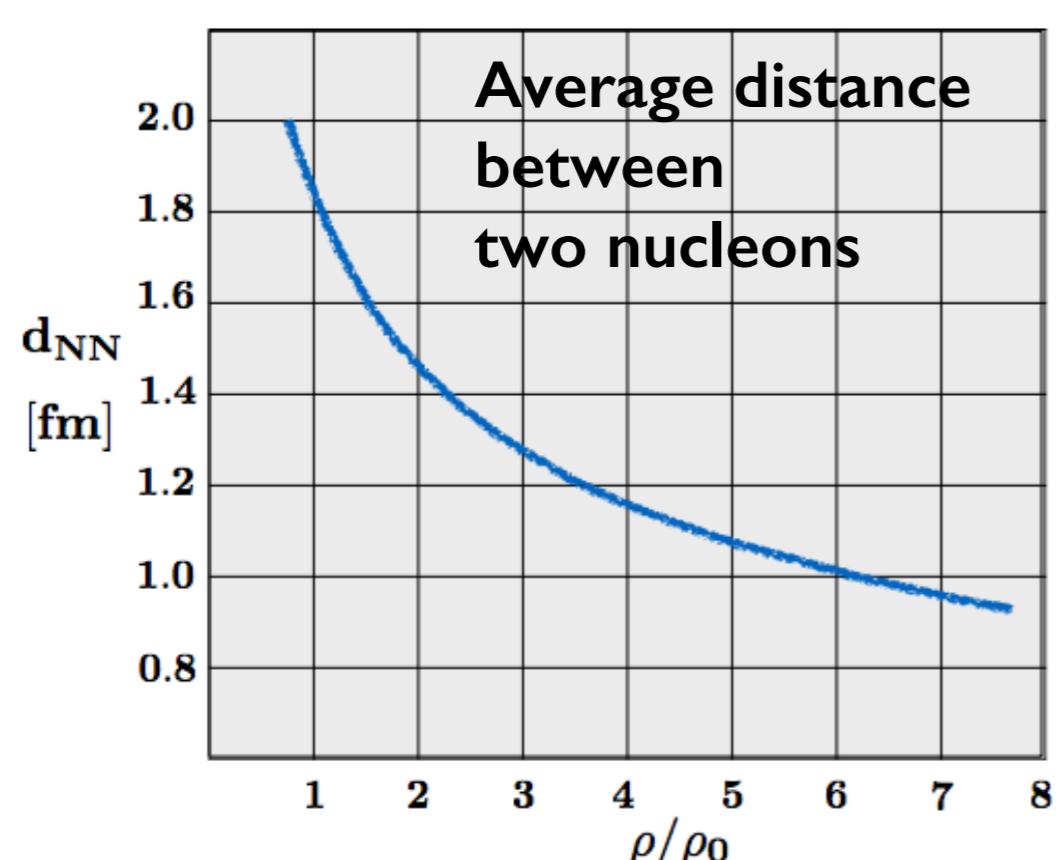
Rept. Prog. Phys. 81 (2018) 056902 arXiv:1903.08963 arXiv: 1908.09360
- ★ **Strangeness in the core of neutron stars ?**
 - ▶ Hyperon-nuclear interactions from Chiral SU(3) EFT
 - ▶ Repulsive YNN three-body forces are capable of resolving hyperon puzzle
 - ▶ **But: quest for much improved hyperon-nucleon/nuclear data base !!**

Appendix

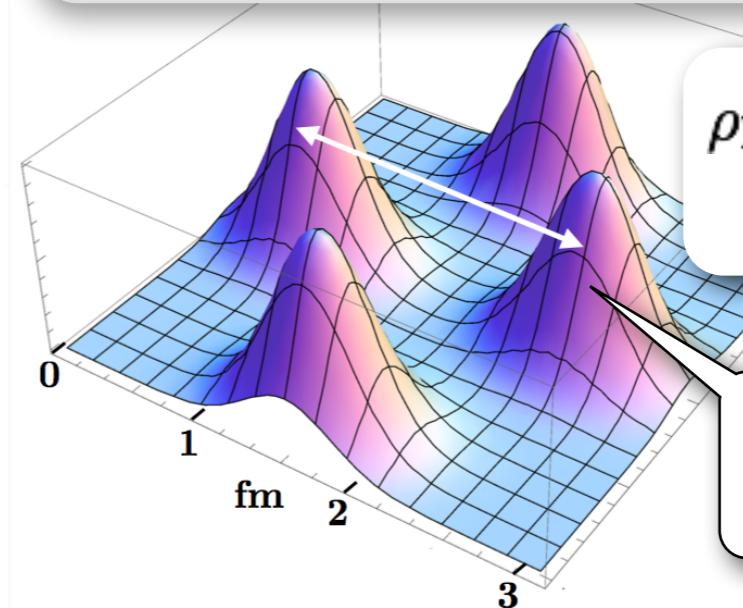
Supplementary Materials

Chiral Nucleon-Meson Field Theory and Functional Renormalisation Group

Densities and Distance Scales in Baryonic Matter



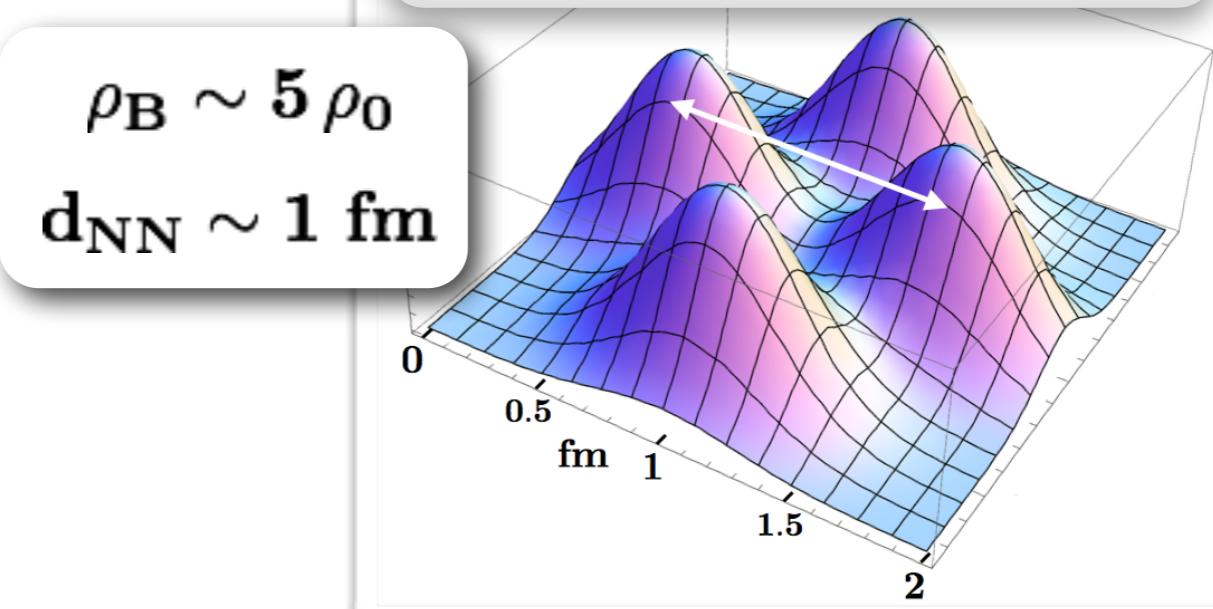
normal nuclear matter



$$\rho_B = \rho_0 = 0.16 \text{ fm}^{-3}$$
$$d_{NN} \simeq 1.8 \text{ fm}$$

baryonic core of the nucleon

neutron star core matter



$$\rho_B \sim 5 \rho_0$$

$$d_{NN} \sim 1 \text{ fm}$$

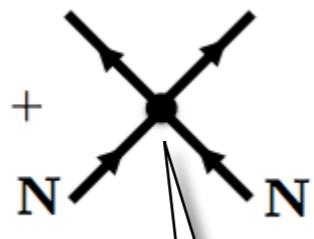
- (Multi-)pionic field in the space between baryonic core sources
- Quark cores of nucleons overlap (percolate) at baryon densities

$$\rho_B > 5 \rho_0$$

Mesons, Nucleons, Dense Baryonic Matter and Functional Renormalisation Group

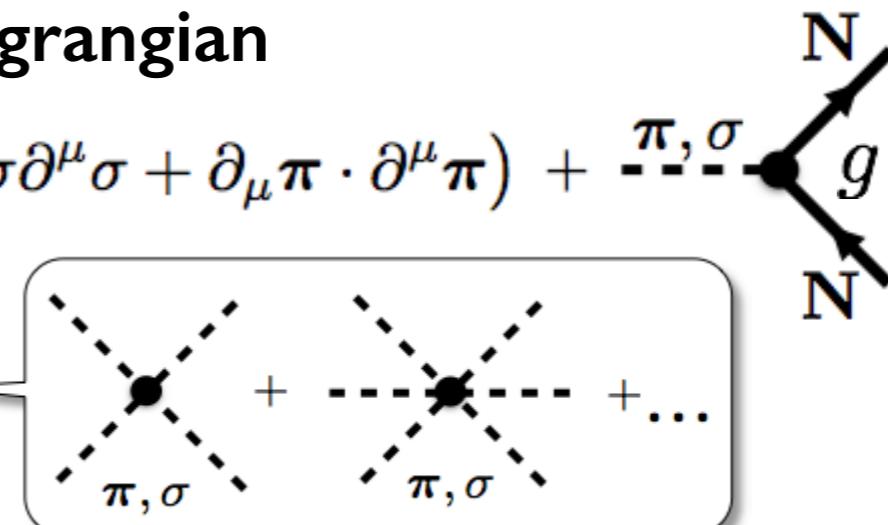
- Chiral nucleon - meson Lagrangian

$$\mathcal{L} = \bar{\mathbf{N}} i\gamma_\mu \partial^\mu \mathbf{N} + \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma + \partial_\mu \pi \cdot \partial^\mu \pi) + \dots$$



- $\mathcal{U}(\pi, \sigma)$

isoscalar & isovector
current-current interactions
mediated by heavy vector fields



- Nambu-Goldstone boson π and “heavy” σ
- Potential $\mathcal{U}(\sigma, \pi)$: polynomial in $\chi = \pi^2 + \sigma^2$ constructed to reproduce vacuum physics and equilibrium nuclear matter
- Symmetry-breaking mass term + $\chi^2 \log \chi$ term

Review: M. Drews, W.W. : Prog. Part. Nucl. Phys. 91 (2017) 347

Renormalisation Group strategies

k-dependent action

$$k \frac{\partial \Gamma_k[\Phi]}{\partial k} = \frac{1}{2} \text{Tr} \left[k \frac{\partial R_k}{\partial k} \cdot \left(\Gamma_k^{(2)}[\Phi] + R_k \right)^{-1} \right]$$

$$\Gamma_{k=\Lambda}[\Phi] = S$$

UV

$\Gamma_k[\Phi]$

C.Wetterich:
Phys. Lett. B 301 (1993) 90

scale regulator R_k

$\bullet \quad \Gamma_{k=0}[\Phi] = \Gamma[\Phi]$
IR

- Thermodynamics:

$$k \partial_k \bar{\Gamma}_k(T, \mu) = \left(\left(\text{nucleons} \right) + \left(\text{pions} \right) \right) \Big|_{T, \mu} - \left(\left(\text{nucleons} \right) + \left(\text{pions} \right) \right) \Big|_{T=0, \mu_0}$$

- Non-perturbative treatment of :

- multi-pion exchange processes
- nucleon-hole excitations
- multi-nucleon correlations

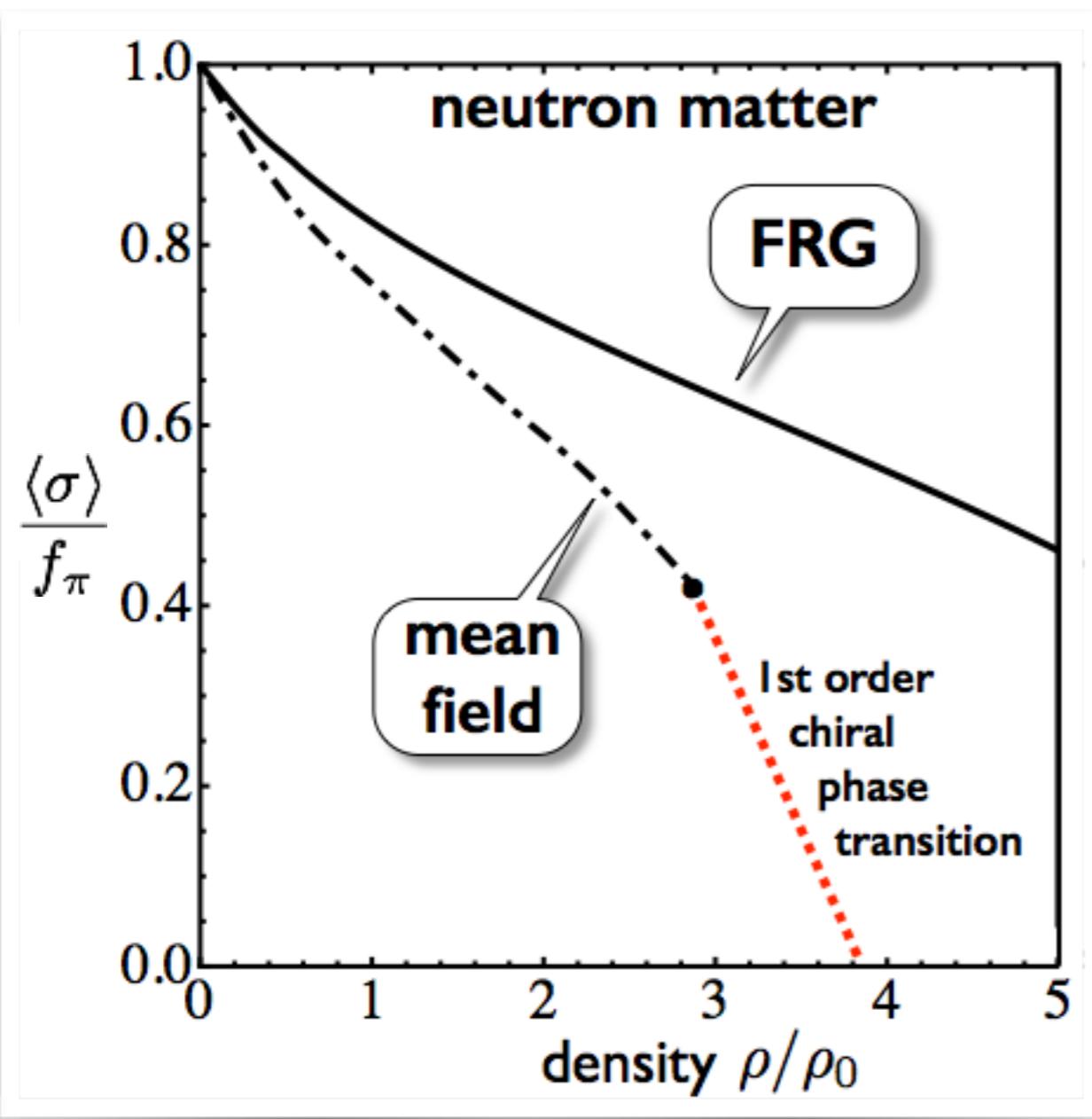
CHIRAL ORDER PARAMETER in NEUTRON MATTER

- Chiral Nucleon-Meson field theory and Functional Renormalization Group

M. Drews, W.W.

Phys. Rev. C91 (2015) 035802

Prog. Part. Nucl. Phys. 93 (2017) 69



- Chiral order parameter :
sigma field
in-medium pion decay constant
$$\langle\sigma\rangle_\rho = f_\pi^*(\rho)$$

Important role of **fluctuations**
(pionic, nucleon-hole,
many-body correlations)
beyond mean-field approximation:

**NO first-order
chiral phase transition
in the n-star density range**