The search for exotic nuclear structure far from stability

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Introduction

Stable nuclei are qualitatively described by "simple" models

- (semi-empirical) liquid-drop model
- (basic) shell model

New techniques enable *ab initio* methods (*A*-body models)

What happens far from stability?

- Experimentally, Radioactive-Ion Beams (RIB) available since 80s
- \Rightarrow study of structure far from stability
- \Rightarrow discovery of exotic structures
 - halo nuclei
 - shell inversions
 - . . .

Nuclear Landscape



• ~ 300 stable nuclei @ $Z \simeq N$ up to ⁴⁰Ca @ N > Z for A > 40

stable nuclei

compact

- magic numbers
- RIB allow to study radioactive nuclei
- Terra incognita between driplines n-dripline unknown beyond O

Basic features in nuclear structure

- Liquid-drop model
- Shell model
- 2 Ab initio nuclear models
- 3 Radioactive-Ion Beams
- Oddities far from stability
 - Halo nuclei
 - Change in shell structure

5 Summary

Charge distributions in (stable) nuclei



- constant density ρ_0 out to the surface (saturation)
- same skin thickness t

(Stable) nuclei look like liquid drops of radius $R \propto A^{1/3}$

Liquid-drop model

Binding energy per nucleon B(Z, N)/A has smooth behaviour Bethe-Weizsäcker semi-empirical mass formula

$$B(Z, N) = a_{\nu}A - a_{s}A^{2/3} - a_{C}\frac{Z(Z-1)}{A^{1/3}} - a_{Sym}\frac{(A-2Z)^{2}}{A}$$

Variation from the semi-empirical mass formula



More bound systems at Z or N = 2, 8, 20, 28, 50, 82, 126magic numbers

 \Rightarrow shell structure in nuclei as in atoms?

Shell model

Developed in 1949 by M. Goeppert Mayer and H. Jensen

As electrons in atoms, nucleons in nuclei feel a mean field and arrange into shells

Spin-orbit coupling is crucial to get right ordering of shells



Example

Shell model explains the higher stability at some Z and N

It predicts the spin and parity of ground state of most nuclei and some of their excited levels, e.g. $^{17}{\rm O}$ and $^{17}{\rm F}$



Nowadays

- Can we go beyond these models?
- Can we build ab initio models?
- i.e. based on first principles
 - nucleons as building blocks
 - realistic N-N interaction

A-body Hamiltonian

Nuclear-structure calculations : A nucleons (Z protons+N neutrons) Relative motion described by the A-body Hamiltonian

$$H = \sum_{i=1}^{A} T_i + \sum_{j>i=1}^{A} V_{i_j}$$

 \Rightarrow solve the A-body Schrödinger equation

$$H|\Psi_n\rangle = E_n|\Psi_n\rangle$$

 $\{E_n\}$ is the nucleus spectrum

Realistic N-N interactions

 V_{ij} not (yet) deduced from QCD \Rightarrow phenomenological potentials fitted on *N*-*N* observables : d binding energy, *N*-*N* phaseshifts Ex. : Argonne V18, CD-Bonn,...



Light nuclei calculations



[R. Wiringa, Argonne]

Three-body force

Need three-body forces to get it right...

$$H = \sum_{i=1}^{A} T_i + \sum_{j>i=1}^{A} V_{ij} + \sum_{k>j>i=1}^{A} V_{ijk} + \cdots$$

But there is no such thing as three-body force...

They simulate the non-elementary character of nucleons \Rightarrow include virtual Δ resonances, \bar{N} ...

$$\begin{array}{c|c} \overline{\pi,\rho,\omega} & \underline{\Delta,N^*} \\ \overline{\pi,\rho,\omega} & \overline{N} \\ \overline{\pi,\rho,\omega} & \overline{\pi,\rho,\omega} \end{array}$$

Phenomenological 3-body interaction fitted on A > 2 levels : IL2 Alternatively, derived from EFT

Effective Field Theory

EFT is an effective quantum field theory based on QCD symmetries with resolution scale Λ that selects appropriate degrees of freedom : nuclear physics is not built on quarks and gluons, but on nucleons and mesons

EFT provides the nuclear force with a systematic expansion in Q/Λ

- gives an estimate of theoretical uncertainty
- naturally includes many-body forces



Expansion of the EFT force





Solving the Schrödinger equation $H |\Psi_n\rangle = E_n |\Psi_n\rangle$

 Ψ usually developed on a basis { $|\Phi_{[\nu]}\rangle$ } :

$$\Psi_n \rangle = \sum_{[\nu]} \langle \Phi_{[\nu]} | \Psi_n \rangle | \Phi_{[\nu]} \rangle$$

Solving the Schrödinger equation reduces to matrix diagonalisation

$$\langle \Phi_{[\mu]} | H | \Psi_n \rangle = \sum_{[\nu]} \langle \Phi_{[\mu]} | H | \Phi_{[\nu]} \rangle \langle \Phi_{[\nu]} | \Psi_n \rangle$$

= $E_n \langle \Phi_{[\mu]} | \Psi_n \rangle$

 \Rightarrow need to build an efficient set of basis states $\{|\Phi_{\nu}\rangle\}$

Clear short review paper : [Bacca EPJ Plus 131, 107 (2016)]

Solving the Schrödinger equation on the lattice Alternatively, solve Schrödinger equation numerically on the lattice (like in lattice QCD)



Review paper : [D. Lee Prog. Part. Nucl. Phys. **63** 117 (2009)] TALENT lecture in 2016

[see D. Lee's talk on Monday]





[Hebeler et al. Annu. Rev. Nucl. Part. Sci. 65, 457 (2015)]

Different *ab initio* models predict similar result All require 3N forces to reproduce the dripline at ²⁴O

What happens far from stability?

Liquid-drop and shell models are fair models of stable nuclei What happens away from stability ?

In 80s Radioactive-Ion Beams were developed Enable study of nuclear structure

[see T. Nakamura on Thursday and I. Tanihata on Friday]

- are radioactive nuclei compact?
- are shells conserved far from stability?

Reactions involving radioactive nuclei useful in astrophysics

[see 2nd part, C. Gustavino on Tuesday and A. Tumino on Friday]

How?

Idea : break a heavy nuclei into pieces to produce exotic isotopes

• ISOL : Fire a proton at a heavy nucleus



• In-flight : Smash a heavy nucleus on a target



Where?



In-flight projectile fragmentation



high-energy primary beam of heavy ions (e.g. ¹⁸O, ⁴⁸Ca, U...) on thin target of light element (Be or C) \Rightarrow fragmentation/fission produces many exotic fragments at $\approx v_{beam}$ Sorted in fragment separator

Used for high-energy reactions (KO, breakup...)

[see T. Nakamura talk on Thursday]

Examples : RIKEN, NSCL (MSU), GSI, GANIL

RIBF @ RIKEN



Superconducting Ring Cyclotron



Largest superconducting cyclotron in the world Delivers a U beam at 350AMeV

Halo nuclei

Exotic structure discovered by I. Tanihata [PLB 160, 380 (1985)] Very large matter radius ($R \gg A^{1/3}$)

Seen as core + one or two neutrons at large distance

- Light, neutron-rich nuclei
- small S_n or S_{2n}

One-neutron halo ${}^{11}\text{Be} \equiv {}^{10}\text{Be} + n$ ${}^{15}\text{C} \equiv {}^{14}\text{C} + n$

Two-neutron halo ${}^{6}\text{He} \equiv {}^{4}\text{He} + n + n$ ${}^{11}\text{Li} \equiv {}^{9}\text{Li} + n + n$



 ^{11}Li

0

[T. Nakamura's talk on Thursday and

²⁰⁸Pb

posters of C. Hebborn, L. Moschini and N. Sokołowska on Monday]

Change in magic numbers

Far from stability usual magic numbers disappear :

• 10 He (Z = 2, N = 8) is unbound

•
28
O (Z = 8, N = 20) is unbound

and new magic numbers appear : N = 6, 14, 16...

One possible explanation is the blocking effect of tensor force

[see I. Tanihata's talk on Friday]

Tensor force is a component of NN interaction

$$V^{T} = v_{T}(r) S_{12}$$

where $S_{12} = 4 \left[\frac{3}{r^{2}} (s_{1} \cdot r) (s_{2} \cdot r) - s_{1} \cdot s_{2} \right]$
$$= \frac{6}{r^{2}} (S \cdot r)^{2} - 2 S^{2}$$

This interaction can lead to $\Delta L = 2$ excitation in np pairs which is responsible for $Q_d \neq 0$ and $\mu_d \neq 0$

Tensor blocking

Tensor force responsible for a significant part of $\boldsymbol{\alpha}$ binding

[Myo et al. PTP **117**, 257 (2007)]

In asymmetric nuclei $(N \gg Z)$ these couplings are Pauli blocked



[Myo et al. PRC 76, 024305 (2007)]

Explains why magic $N = 8 \rightarrow 6$ in He and $N = 20 \rightarrow 14, 16$ in O [see also T. Nakamura's talk on Thursday]

Summary

Liquid-drop and shell model describe qualitatively stable nuclei Nowadays *ab initio* nuclear-structure models from first principles

RIBs enable us to study nuclear structure far from stability New exotic structure discovered :

- halo nuclei diffuse halo around a compact core
- shell inversions or shell collapse
- nuclei beyond the dripline (resonant ground state)

RIB can be used to study reactions of astrophysical interest...