

Overview of recent measurements of γ production and suppression with the STAR experiment

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1 Introduction

- Upsilon as a probe of quark-gluon plasma
- Production mechanism

2 STAR experiment

3 Υ production in p+p

- p_T and rapidity spectra
- Cross section ratios
- Event activity dependence of Υ production

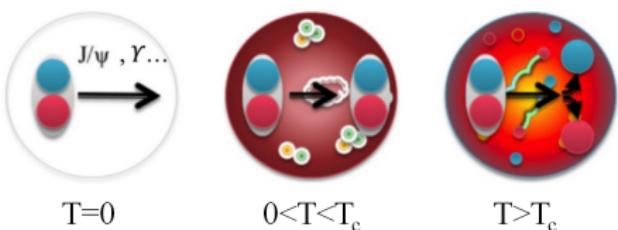
4 Υ production in p+Au

5 Υ suppression in Au+Au

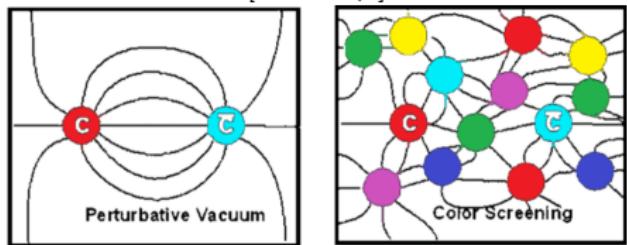
6 Summary

Upsilon in quark-gluon plasma

$$\Upsilon = b \bar{b}$$



[A. Rothkopf]



[O. Matonoha, Hot Quarks 2018]

High mass - produced early

$$m_b = 4.18_{-0.03}^{+0.04} \text{ GeV}/c^2$$

$$m_{\Upsilon(1S)} = 9460.30 \pm 0.26 \text{ MeV}/c^2$$

$$m_{\Upsilon(2S)} = 10023.26 \pm 0.31 \text{ MeV}/c^2$$

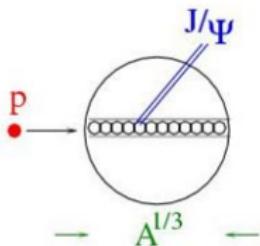
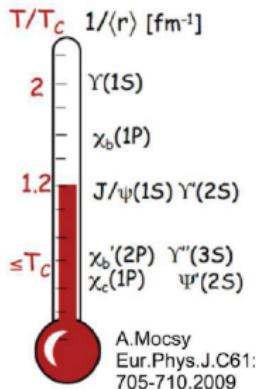
$$m_{\Upsilon(3S)} = 10355.2 \pm 0.5 \text{ MeV}/c^2$$

[Phys. Rev. D 98, 030001 (2018)]

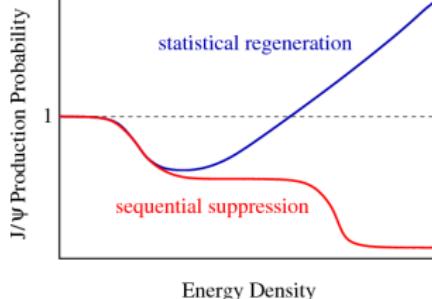
Υ as a probe of quark-gluon plasma

- Υ is sensitive to the QGP properties similarly to J/ψ
[Phys.Lett.B 178(4),416-422(1986)]
- Dissociation due to Debye-like screening when $r_\Upsilon > r_{\text{Debye}} \propto T^{-1}$
- Suppression observed at RHIC and LHC
[Phys.Lett.B 735,127-137(2014)], [Phys.Lett.B. 770,357-359(2017)]
- Sequential suppression, due to lower binding energy for excited Υ states, expected and observed
[Phys.Rev.D 64, 094015], [Phys.Rev.Lett 109, 222301]

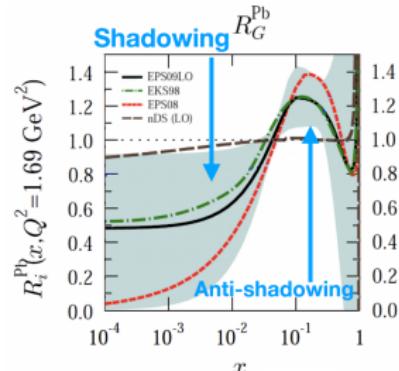
Other effects



[L. Grandchamp, LBNL
2005]



[Nucl.Phys.B (Proc.Suppl.) 214, 3-36(2011)]



[Phys.Rev.C 81 (2010) 064911]

Other modifications to Υ production

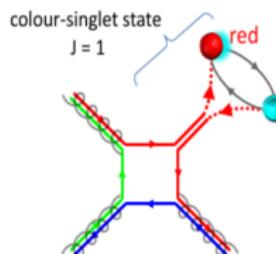
- Regeneration very small for Υ at RHIC
[Phys.Rev.C 96, 054901]
- Feed-down from excited states $\Upsilon(nS) \rightarrow \Upsilon(1S)\pi^+\pi^-$ and $\chi_{bn} \rightarrow \gamma\Upsilon(1S)$
- Cold Nuclear Matter effects - can be studied separately in $p + A$ or $d + A$ collision
 - nuclear absorption
 - comover interactions - very small for $\Upsilon(1S)$
[Phys.Lett.B 503, 104 (2001)]
 - nuclear PDFs: shadowing, anti-shadowing

Production mechanism

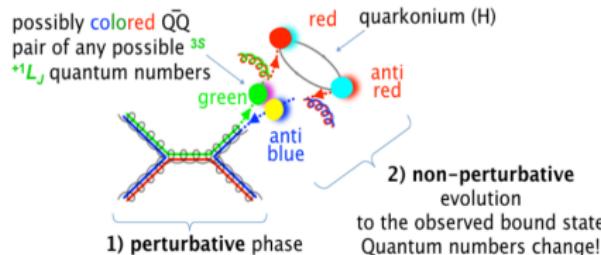
- Still not well understood: hard scattering+non-perturbative hadronization
- Quarkonium measurements provide tests of production models, and thus help to understand QCD

Quarkonium production models

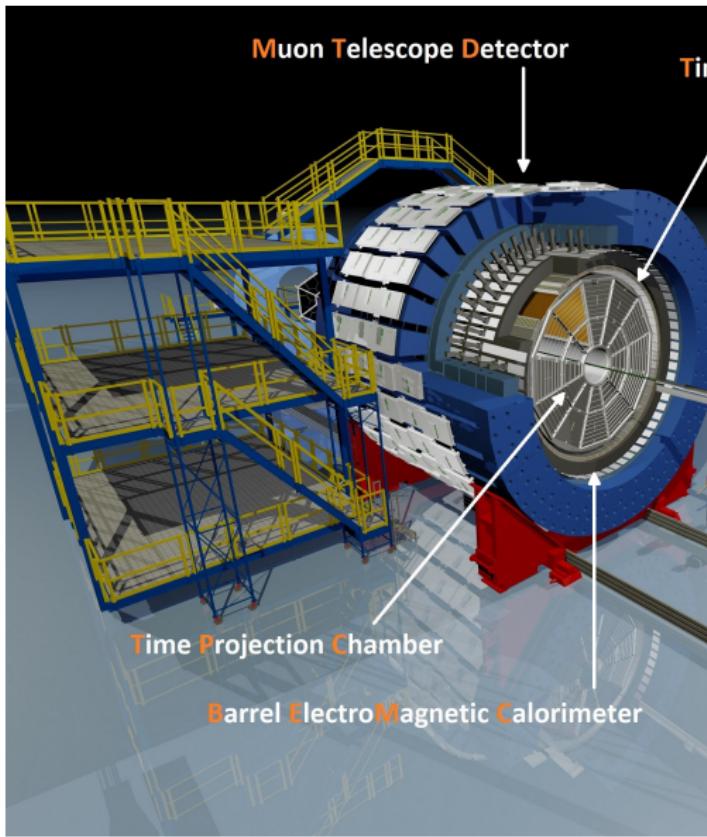
- Color Singlet - $Q\bar{Q}$ produced directly in a color neutral state
- Color Octet - $Q\bar{Q}$ produced in a colored state. Gluon emissions are needed to neutralize color. This is described by long-distance matrix elements (LDMEs) which are assumed universal.
- Color Evaporation Model - color irrelevant. Fixed fractions of $Q\bar{Q}$ pairs evolve into various quarkonium states.



+ analogous colour combinations



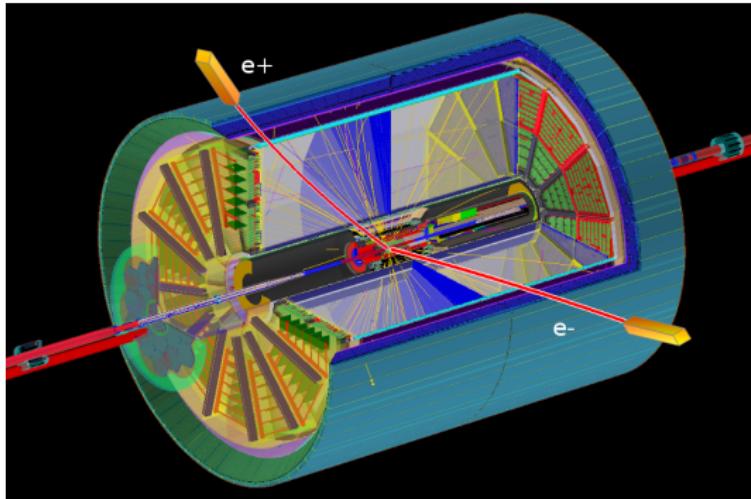
[P. Faccioli, Polarization in LHC physics, Course on Physics at the LHC 2014]



Detectors used for γ studies

- TPC $|\eta| < 1, 0 < \phi < 2\pi$
 - Tracking - momentum measurement
 - Particle identification based on energy loss dE/dx
- BEMC $|\eta| < 1, 0 < \phi < 2\pi$
 - Trigger on high- p_T electrons
 - Electron identification via E/p and EM shower shape
- MTD $|\eta| < 0.5, 45\% \text{ in } 0 < \phi < 2\pi$
 - Dimuon trigger
 - Muon identification with time-of-flight
 - Magnet used as hadron absorber
 - Muons - less bremsstrahlung
- TOF $|\eta| < 1, 0 < \phi < 2\pi$
 - Particle identification based on time-of-flight - not used for γ
 - Fast detector used to remove pile-up for N_{ch} determination

Upsilon measurements in STAR - dielectron channel



Υ decay reconstruction in e^+e^- channel

$$\Upsilon(1S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(1S)} = 2.38 \pm 0.11\%$$

$$\Upsilon(2S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(2S)} = 1.91 \pm 0.16\%$$

$$\Upsilon(3S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(3S)} = 2.18 \pm 0.20\%$$

[*Phys. Rev. D* 98, 030001 (2018)]

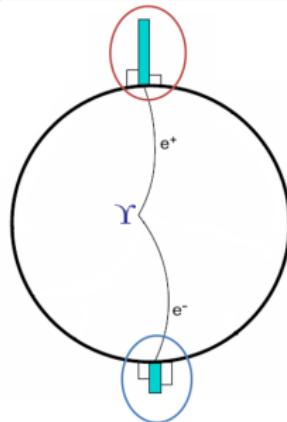
- Project TPC track to the high-energy tower in BEMC, which fired the trigger, and reconstruct a cluster
- Find a partner track and project it to BEMC cluster

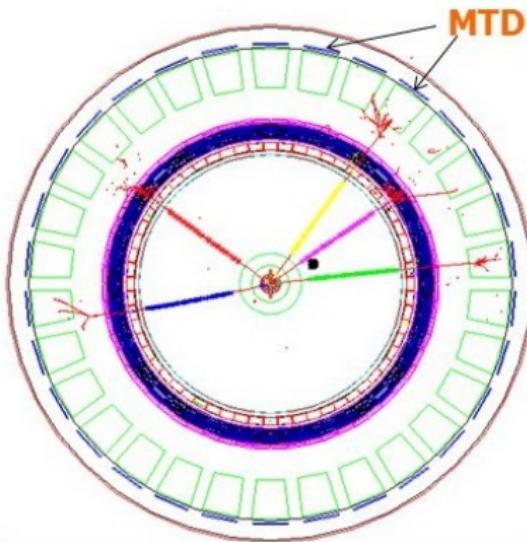
Detectors used

- TPC+BEMC(+TOF for N_{ch})

STAR datasets for $\Upsilon \rightarrow e^+e^-$

- Au+Au $\sqrt{s_{NN}} = 200$ GeV: 2011, 2010
- p+Au $\sqrt{s_{NN}} = 200$ GeV: 2015
- d+Au $\sqrt{s_{NN}} = 200$ GeV: 2008
- p+p $\sqrt{s} = 500$ GeV: 2011
- p+p $\sqrt{s} = 200$ GeV: 2015, 2009





Υ decay reconstruction in $\mu^+\mu^-$ channel

$$\Upsilon(1S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(1S)} = 2.48 \pm 0.05\%$$

$$\Upsilon(2S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(2S)} = 1.93 \pm 0.17\%$$

$$\Upsilon(3S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(3S)} = 2.18 \pm 0.21\%$$

[*Phys. Rev. D* 98, 030001 (2018)]

Detectors used

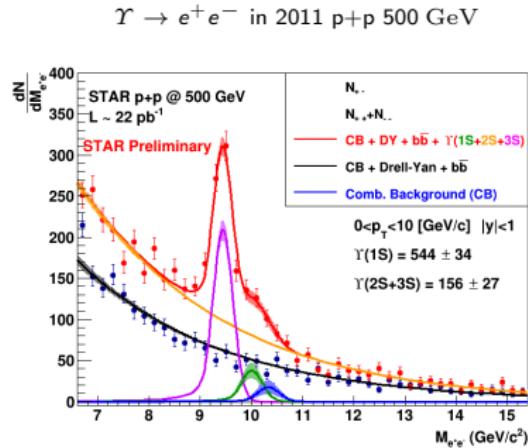
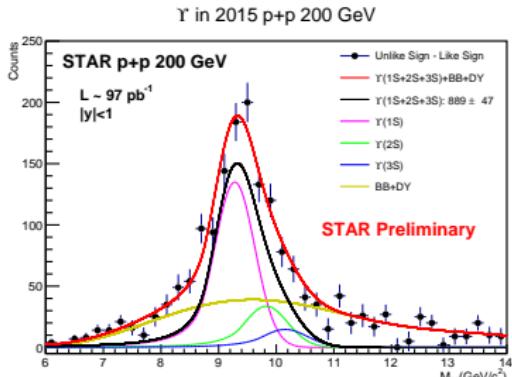
- TPC+MTD

STAR MTD datasets for $\Upsilon \rightarrow \mu^+ \mu^-$

- 2016 Au+Au $\sqrt{s_{NN}} = 200$ GeV
- 2014 Au+Au $\sqrt{s_{NN}} = 200$ GeV

Track projection to MTD hits

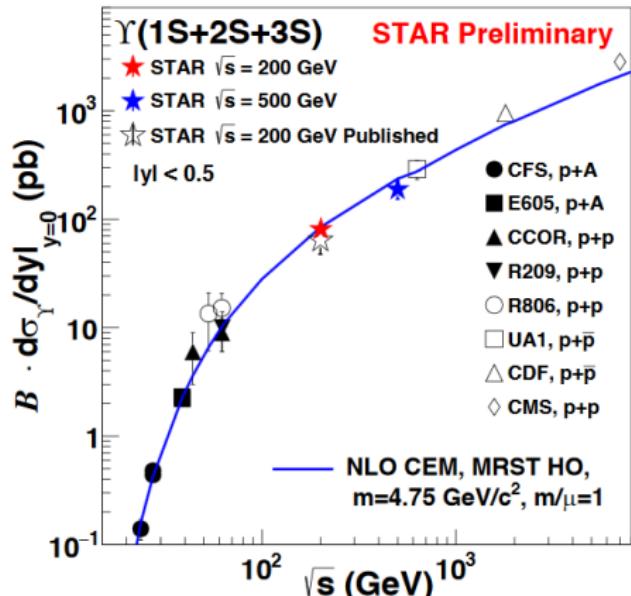
- Tracks are projected from TPC to MTD and matched to hits
- Energy loss in the magnet is included in the track projection procedure



Signal fitting $\Upsilon \rightarrow e^+ e^-$

- Υ signal shapes modeled by 3 Crystal-Ball functions
- Fit to **Unlike-sign (red)** distribution consists of:
 - 3 Crystal-Ball functions (**1S, 2S, 3S** states) - fixed using MC simulation
 - $b\bar{b}$ +Drell-Yan correlated background (**orange**) determined using MC simulation

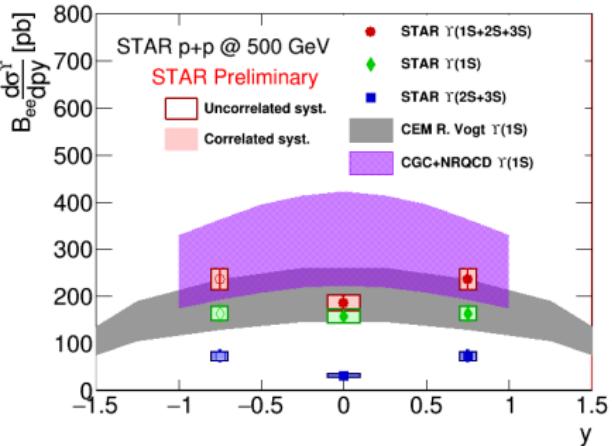
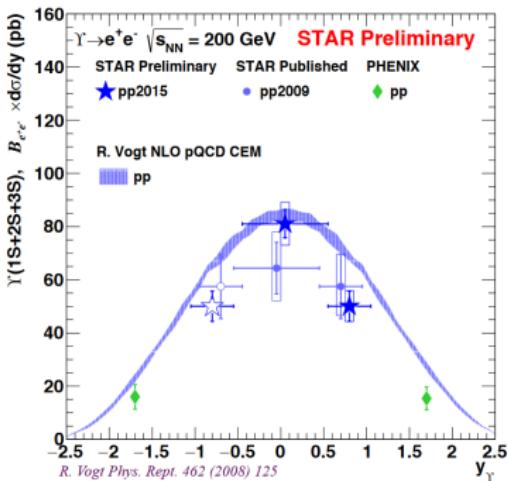
Integrated cross section in p+p



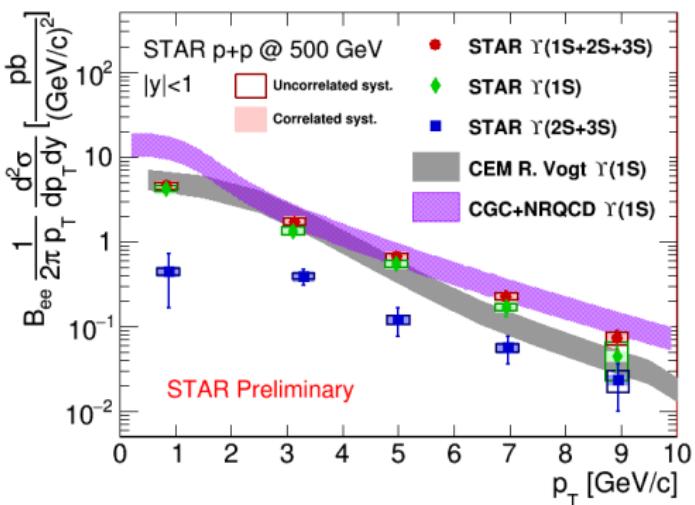
STAR
 [Phys.Lett.B 735,127–137(2014)]
 CDF
 [Phys.Rev.Lett. 88,161802(2002)]
 CMS
 [Phys.Rev.D 83,112004(2010)]
 CF5
 [Phys.Rev.Lett. 39,1240–1242(1977)]
 [Phys.Rev.Lett. 41,684–687(1978)]
 [Phys.Rev.Lett. 42,486–489(1979)]
 [Phys.Rev.Lett. 55,1962–1964(1985)]
 E605
 [Phys.Rev.D 43,2815–2835(1991)]
 [Phys.Rev.D 39,3516(1989)]
 CCOR
 [Phys.Lett.B 87,398–402(1979)]
 L. Camilleri, T.B.W. Kirk, H.D.I. Abarbanel (Eds.)
 E866
 [Phys.Rev.Lett. 100,062301(2008)]
 ISR
 [Phys.Lett.B 91,481–486(1980)]

- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 81 \pm 5(stat) \pm 8(syst)$ pb in p+p collisions at $\sqrt{s} = 200$ GeV
- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 186 \pm 14(stat) \pm 33(syst)$ pb in p+p collisions at $\sqrt{s} = 500$ GeV
- STAR results follow the world data trend
- Consistent with the Color Evaporation Model calculation
 $[Phys.Rep. 462, pp.125–175(2008)]$

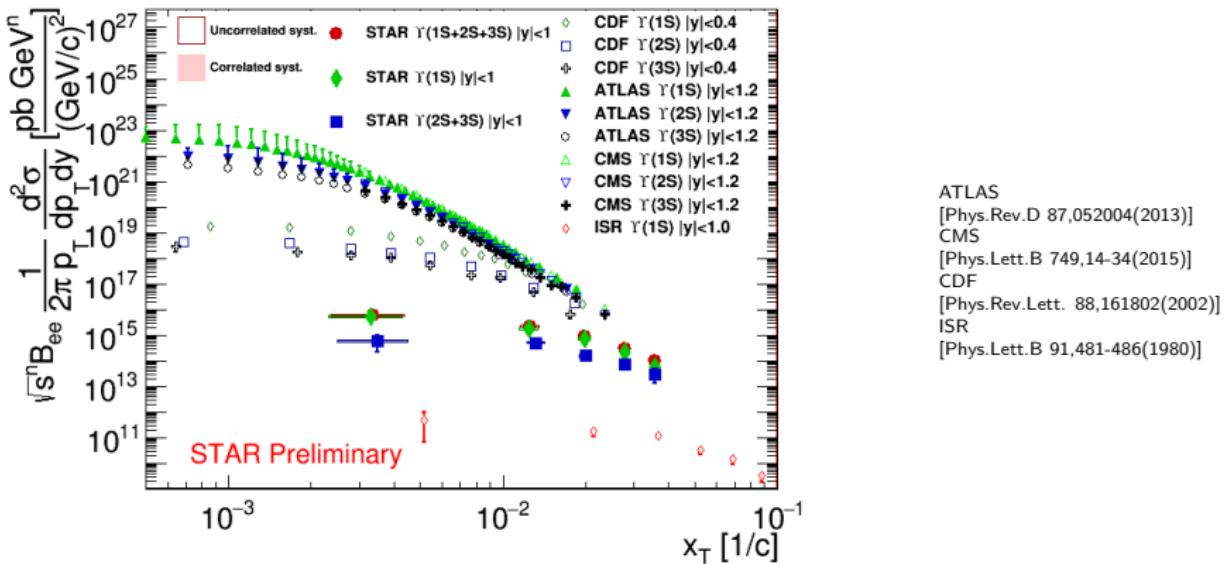
Rapidity dependence in p+p



- STAR data slightly narrower than CEM model at $\sqrt{s} = 200$ GeV
- Flatter rapidity spectrum at $\sqrt{s} = 500$ GeV compared to $\sqrt{s} = 200$ GeV
 - Dip at mid-rapidity for $\gamma(2S + 3S) \approx 2\sigma$ level from flat
 - CEM model (inclusive) consistent with the measurement for $\gamma(1S)$ [*Phys. Rev. C92 (2015) 034909*]
 - CGC+NRQCD (direct) predictions are above the data for $\gamma(1S)$ [*PRD 94, 014028 (2016)*]. [*PRL 113, 192301 (2014)*]

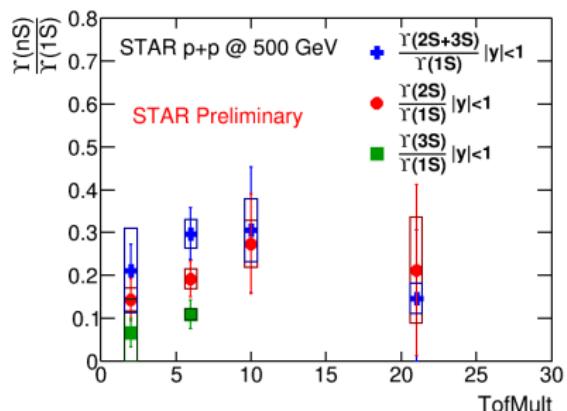
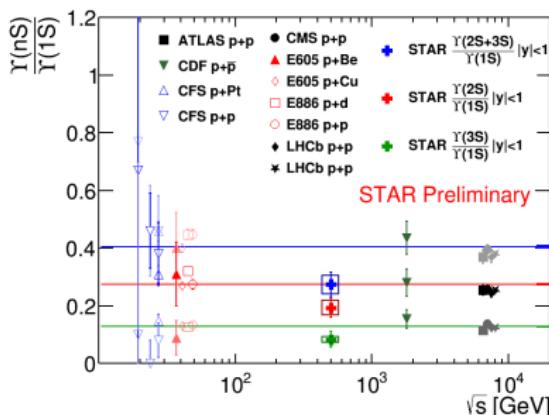


- Color Evaporation Model (CEM) calculation for inclusive $\Upsilon(1S)$
 $[Phys. Rev. C92 (2015) 034909]$
 - Agree with data reasonably well
- Non-relativistic Quantum Chromodynamics coupled with the Color-Glass Condensate formalism (CGC+NRQCD) for direct Υ
 $[PRD 94, 014028 (2016)] [PRL 113, 192301 (2014)]$
 - $\Upsilon(1S)$: model calculation is above the data points. Caveat: additional corrections are needed at low p_T according to authors.



- $x_T = \frac{2p_T}{\sqrt{s}}$, $\sigma^{inv} \equiv E \frac{d^3\sigma}{d^3p} = \frac{F(x_T)}{p_T^n(x_T, \sqrt{s})} = \frac{F'(x_T)}{\sqrt{s}^n(x_T, \sqrt{s})}$
[JHEP06,035(2010)]
- pQCD predicts that spectra of hard processes should follow x_T scaling - check with $n = 5.6$ (number of partons taking active part in the process) obtained for J/ψ
[Phys.Rev.C 80, 041902(2009)]
- No clear scaling observed

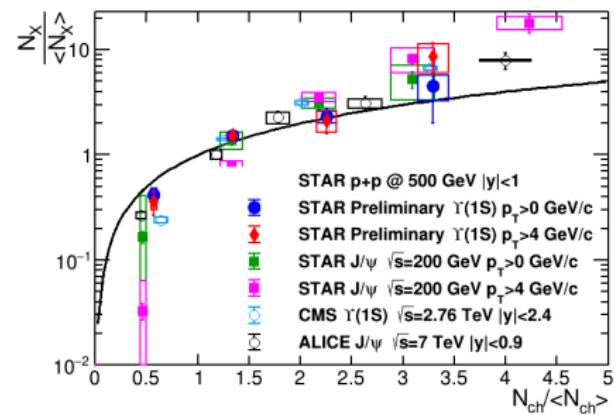
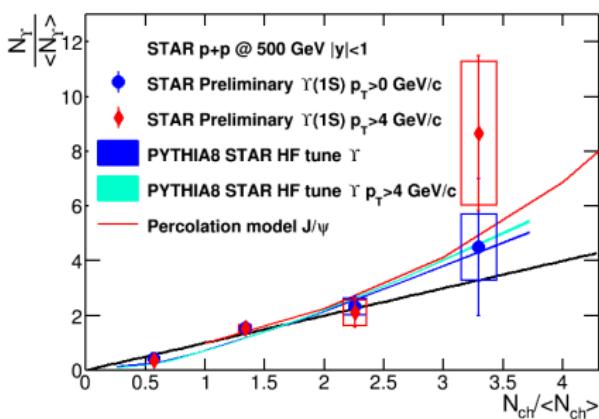
Cross section ratios: $\Upsilon(nS)/\Upsilon(1S)$



[Phys. Rev. C 88, 067901 (2013)]

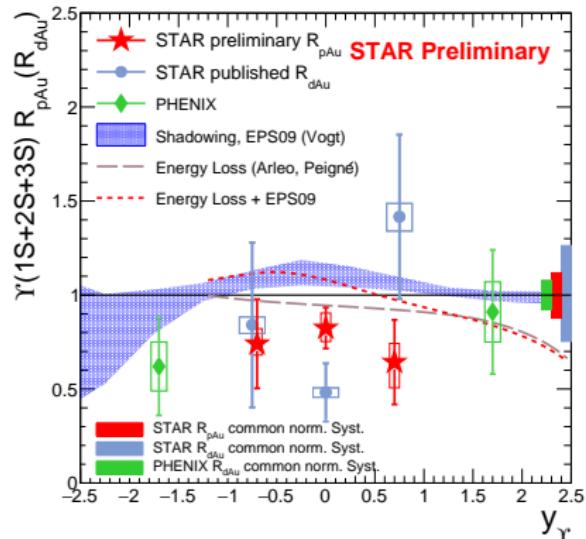
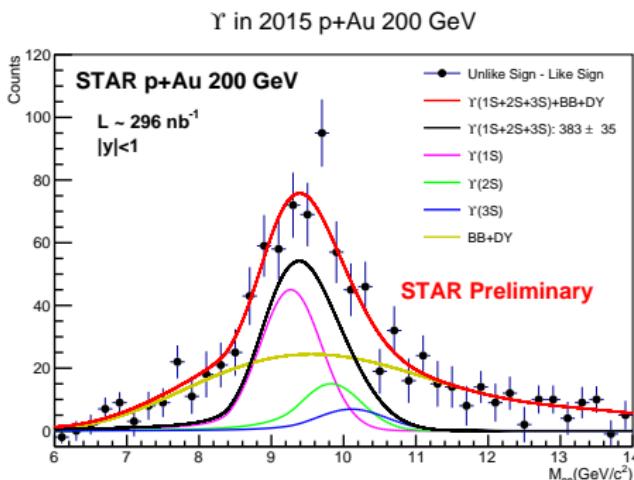
- TofMult: number of tracks matched to TOF within $|\eta| < 1$, $p_T > 0.2 \text{ GeV}/c$
- Boxes correspond to uncorrelated systematic uncertainties (correlated uncertainties largely cancel out)
- Cross section ratios measured in 500 GeV p+p collisions are slightly below (within 2σ) world data average, shown as solid lines in the left plot.
- Right plot: No strong multiplicity dependence observed.

Υ production vs. event activity



- Normalized $\Upsilon(1S)$ yield vs. normalized multiplicity (a measure of event activity)
- Data consistent with a linear rise (black line), with a hint for stronger-than-linear rise for $\Upsilon(1S)$ above $p_T > 4$ GeV/c
- Similar trend at RHIC and LHC for Υ and J/ψ
 $[JHEP04, 103(2014)], [Nucl. and Part. Phys. Proc., 276-278, pp. 261-264(2016)], [Phys. Lett. B 712, 165-175(2012)]$
- Hints of interaction between strings of color field in high multiplicity collisions or Υ production in MPI
 $[Phys. Rev. C, 86, 034903 (2012)]$

Υ production in p+Au

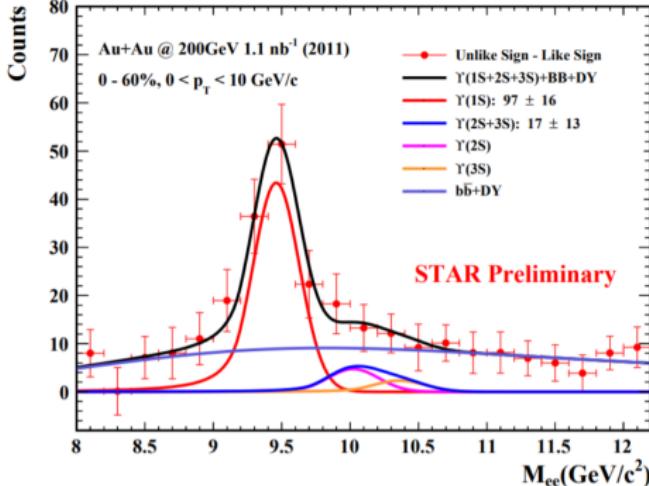


[Phys.Lett.B 735(2014)127], [Phys. Rev. C 87, 044909], [JHEP 03(2013)122]

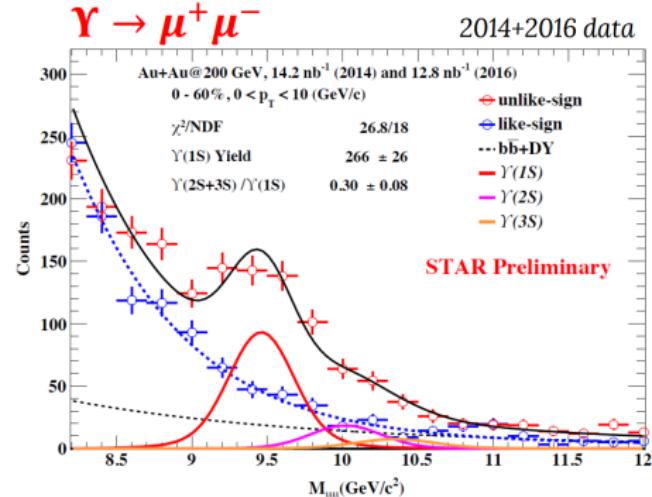
$\Upsilon(1S + 2S + 3S)$

- Significantly improved precision over published R_{dAu} results
 - $R_{pAu}|_{|y| < 0.5} = 0.82 \pm 0.10(\text{stat.})^{+0.08}_{-0.07}(\text{syst.}) \pm 0.10(\text{glob.})$
- Indication of $\Upsilon(1S + 2S + 3S)$ suppression in p+Au collisions

$\Upsilon \rightarrow e^+e^-$



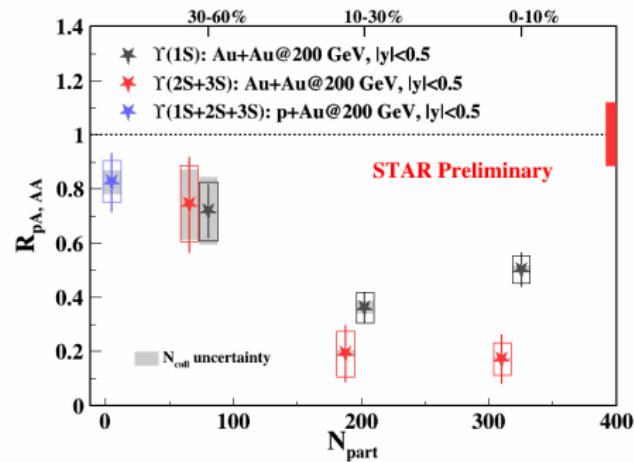
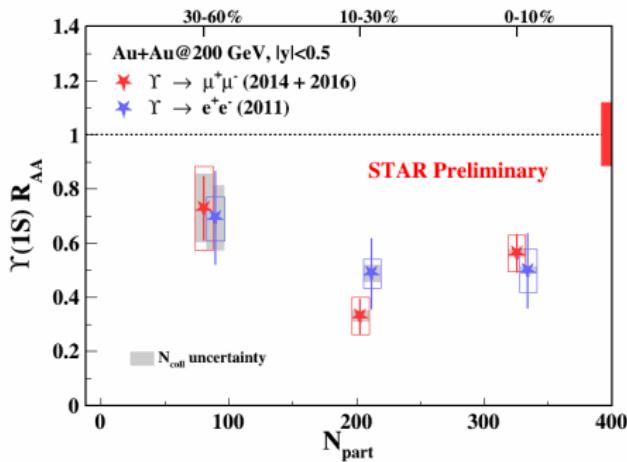
$\Upsilon \rightarrow \mu^+\mu^-$



[Pengfei Wang, QM2018]

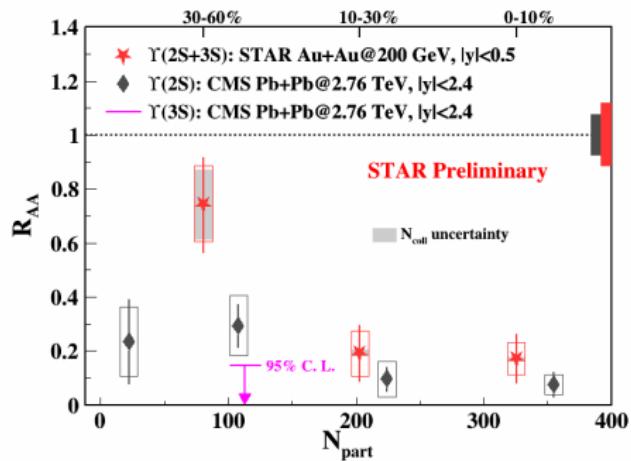
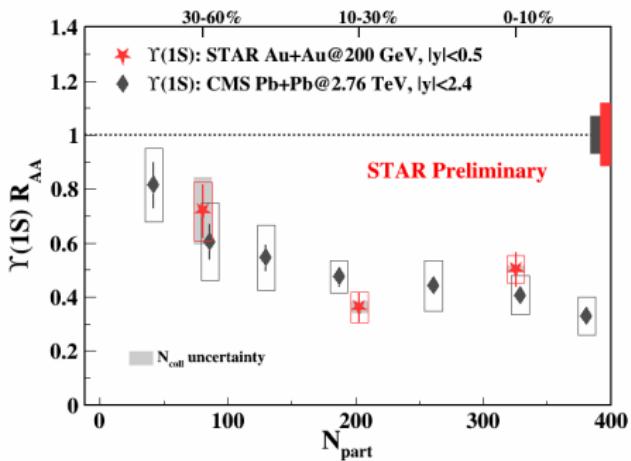
Signal fits

- 3 Crystal Ball fits for $\Upsilon \rightarrow e^+e^-$
- 3 Gaussian fits for $\Upsilon \rightarrow \mu^+\mu^-$, because of less bremsstrahlung



R_{AuAu} measured by STAR

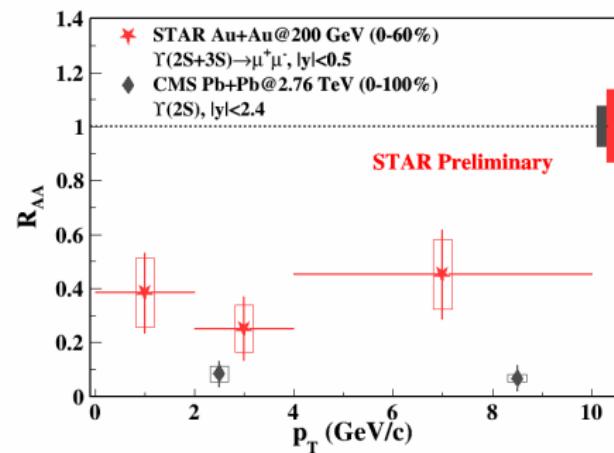
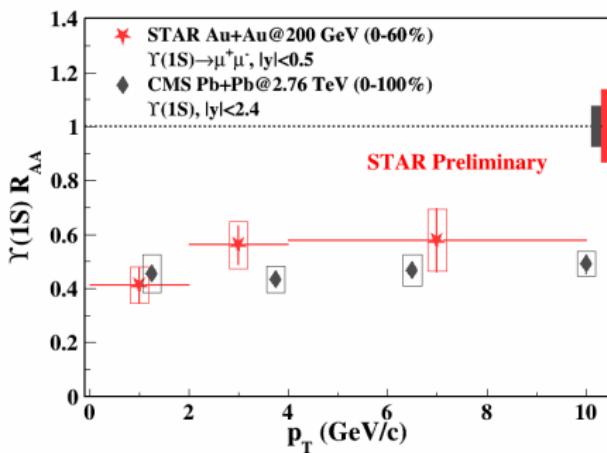
- Dielectron and dimuon results consistent with each other
- Both results combined for better precision
- Stronger suppression of $\Upsilon(2S + 3S)$ than $\Upsilon(1S)$ in central collisions



[Phys.Lett.B 770(2017)357-379]

STAR vs. CMS

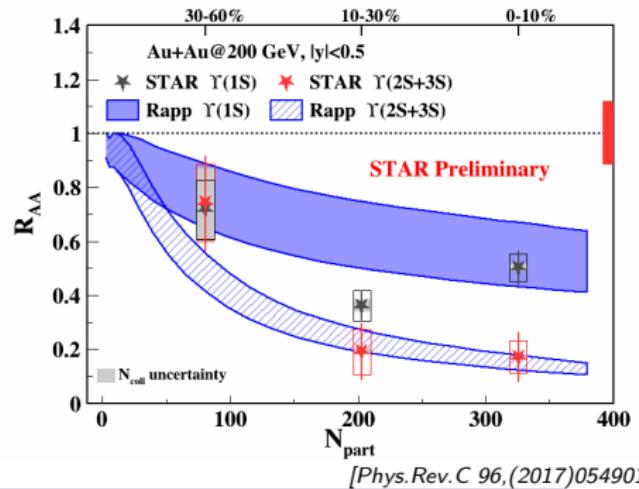
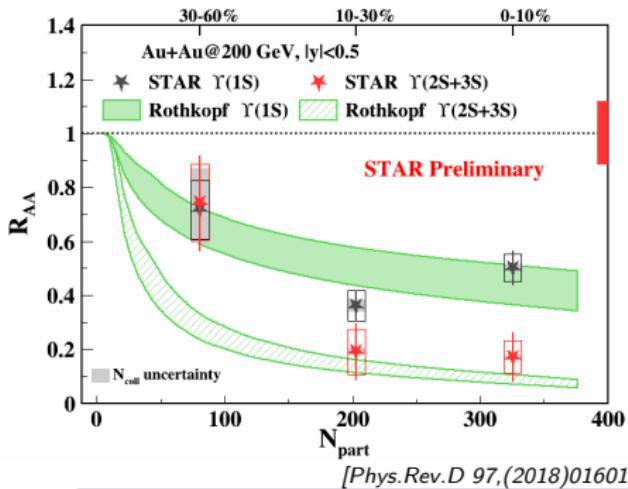
- Similar suppression for $\Upsilon(1S)$, despite higher medium temperature at the LHC
 - Regeneration? Larger at LHC than at RHIC
 - CNM effects
- Indication of smaller suppression for $\Upsilon(2S + 3S)$ at RHIC than at LHC



[Phys.Lett.B 770(2017)357-379]

Transverse momentum dependence

- Similar suppression for $\Upsilon(1S)$ at RHIC and LHC
- Indications of stronger suppression of high- p_T $\Upsilon(2S + 3S)$ at LHC than at RHIC
- Both consistent with flat dependence vs. p_T



Models

- Kroupaa, Rothkopf, Strickland
 - Lattice QCD-vetted potential for heavy quarks in hydrodynamic-modeled medium
 - No regeneration, no CNM effects
- De, He, Rapp
 - Quarkonium in-medium binding energy described by thermodynamic T-matrix calculations with internal energy potential (strongly bound scenario)
 - Includes both regeneration and CNM effects
- Both models agree with STAR $\Upsilon(1S)$ data
- Rothkopf's model underestimates the STAR $\Upsilon(2S + 3S)$ results for 30 – 60% centrality

Summary

p+p collisions at $\sqrt{s} = 200 \text{ GeV}$ and $\sqrt{s} = 500 \text{ GeV}$

- The $\Upsilon(1S)$ spectra can be reasonably described by CEM calculations.
- Flatter rapidity distribution for Υ at $\sqrt{s} = 500 \text{ GeV}$ than at $\sqrt{s} = 200 \text{ GeV}$.
- Measured $\frac{\Upsilon(nS)}{\Upsilon(1S)}$ vs. multiplicity at 500 GeV - no strong dependence.
 - Ratios slightly lower than world data.
- Dependence of Υ production on event activity.
 - Similar trends observed for J/ψ and Υ at RHIC and LHC.
 - Predictions from PYTHIA8 and Percolation model can qualitatively describe the trend observed in data.

p+Au collisions at $\sqrt{s} = 200 \text{ GeV}$

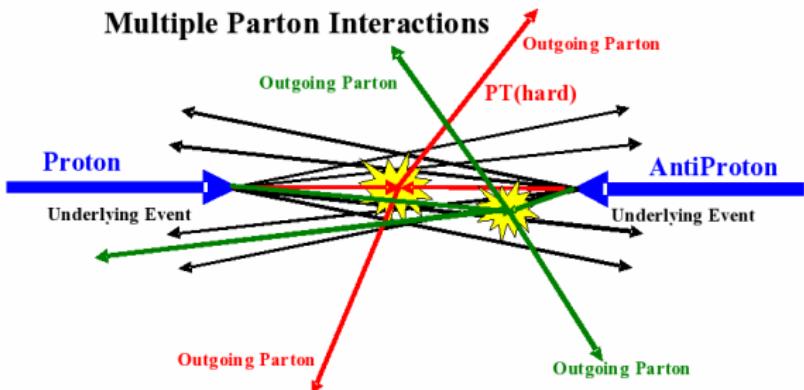
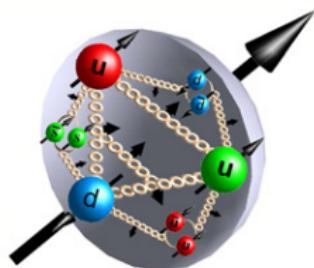
- Indication of Υ suppression

A+A collisions at $\sqrt{s} = 200 \text{ GeV}$

- Consistent results from dielectron and dimuon channels - combined for better precision
- Similar suppression of $\Upsilon(1S)$ at RHIC and LHC
- Stronger suppression of $\Upsilon(2S + 3S)$ than $\Upsilon(1S)$ in central collisions
 - Sequential suppression
 - Hint of smaller suppression at RHIC than at LHC
- Data consistent with model calculations

Thank you for attention!

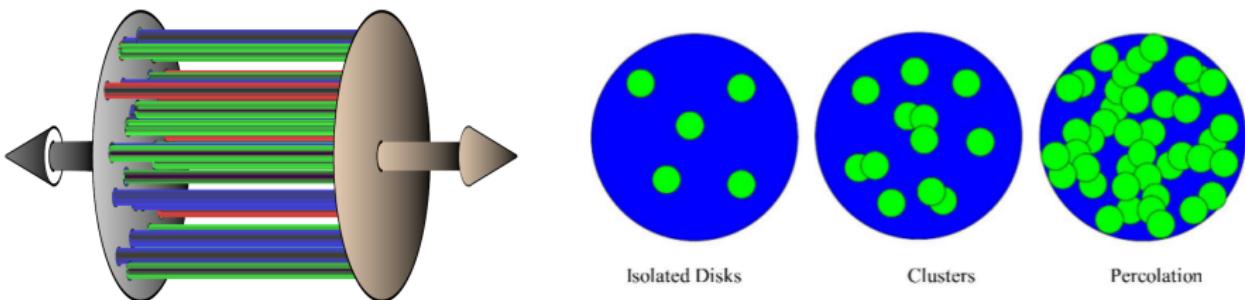
Multiple parton interactions (MPI)



<https://www.bnl.gov/rhic/images/proton-with-gluons-300px.jpg>

<http://www.desy.de/~jung/multiple-interactions/may06/mi-rick.gif>

- Protons are complex objects consisting of constituent quarks, sea quarks and gluons.
- Multiple parton interactions (MPI) may happen in $p + p$ collision - implemented in PYTHIA.
 - Besides the main hard process, there may be additional hard and soft processes in MPI.
- As implemented in PYTHIA8, heavy quarks can also be produced during MPI.
- MPI together with initial- (ISR), final-state radiation (FSR) and beam remnants define the event activity, which can be characterized experimentally using the charged particle multiplicity.



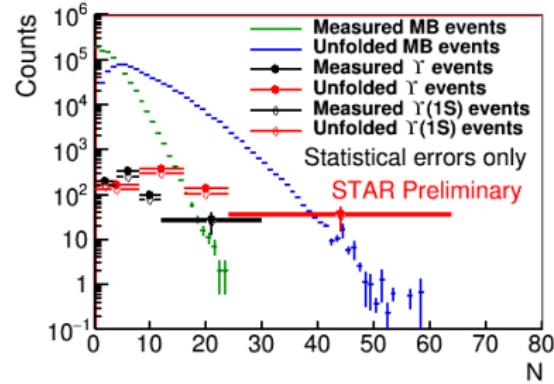
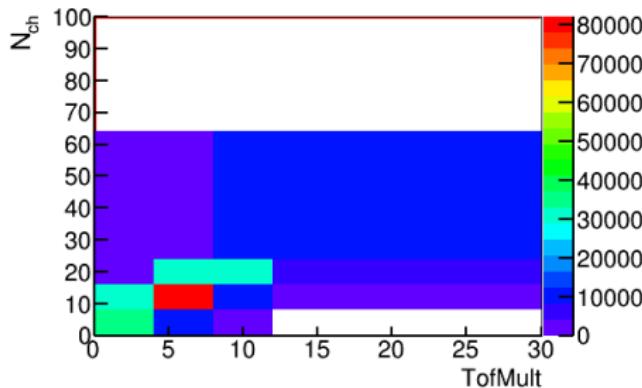
[Ann.Rev.Nucl.Part.Sci.60, 463-489(2010)] [Proceedings of SPIE, 100313U(2016)]

- Models particle production originating from strings of color field formed in $p + p$ collisions.
- Soft particle production damped by interaction of overlapping strings.
- Predicts quadratic dependence of normalized yield for particles from hard processes vs. normalized charged particle multiplicity in high multiplicity events.

$$\frac{N_{hard}}{\langle N_{hard} \rangle} = \langle \rho \rangle \left(\frac{\frac{dN_{ch}}{d\eta}}{\left\langle \frac{dN_{ch}}{d\eta} \right\rangle} \right)^2 \quad [\text{Phys.Rev. C, 86, 034903 (2012)}]$$

Multiplicity distribution via unfolding

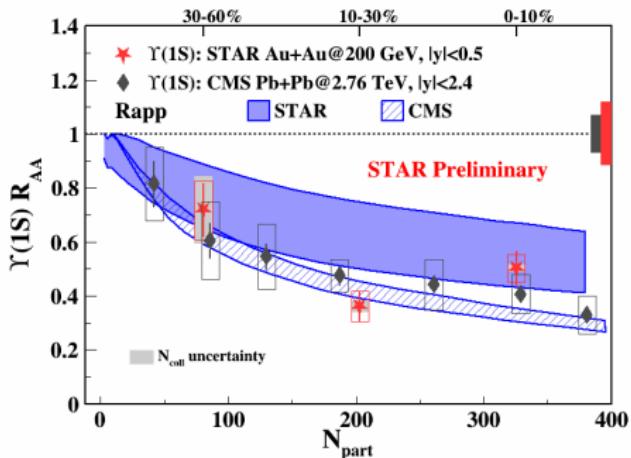
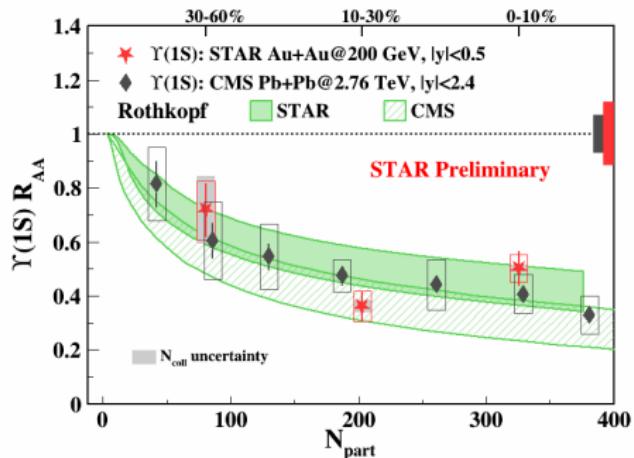
Response matrix for γ events



Unfolding method used for multiplicity dependent studies

- ① A response matrix is obtained using the PYTHIA8 event generator for both min-bias and γ events taking into account reconstruction efficiency
- ② The measured distributions are unfolded with their respective response matrices
- ③ This procedure yields the unfolded (true) distribution

STAR and CMS $\Upsilon(1S)$ vs. models

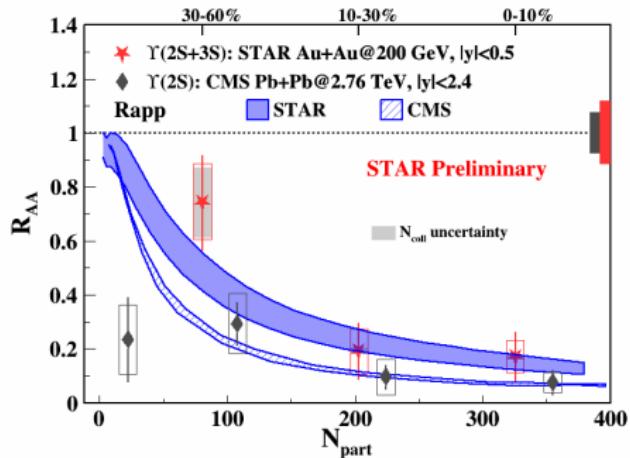
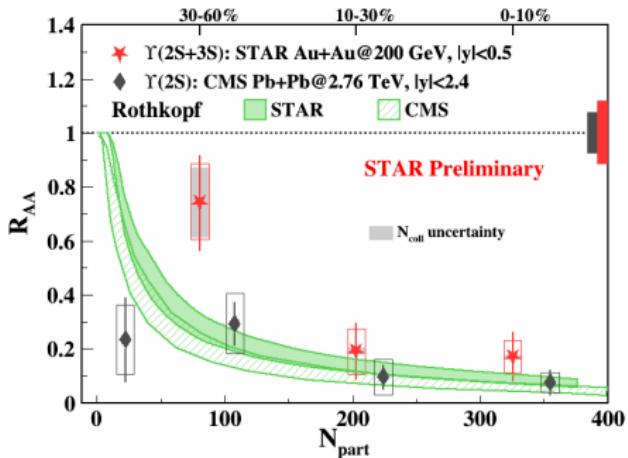


[Phys. Rev.D 97,(2018)016017], [Phys. Rev.C 96,(2017)054901]

$\Upsilon(1S)$ vs. models

- Both models consistent with the data

STAR and CMS $\Upsilon(2S + 3S)$ vs. models

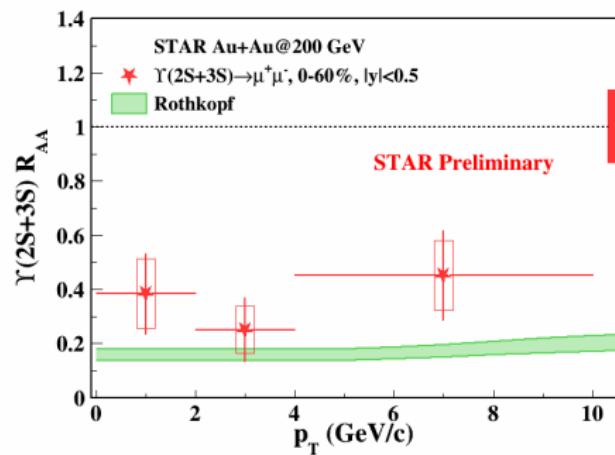
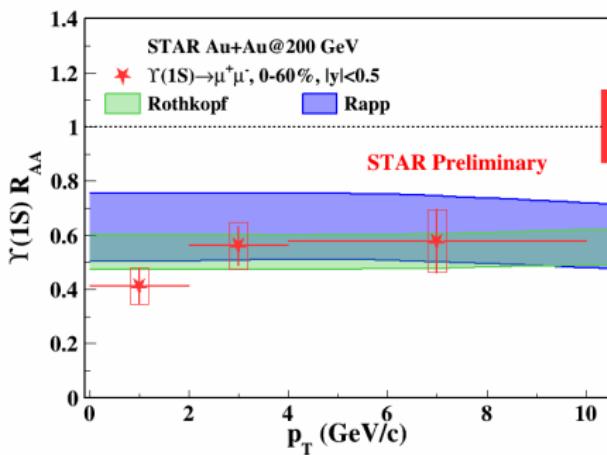


[Phys. Rev.D 97,(2018)016017], [Phys. Rev.C 96,(2017)054901]

$\Upsilon(2S + 3S)$ vs. models

- Both models consistent with the data in central and semi-central collisions

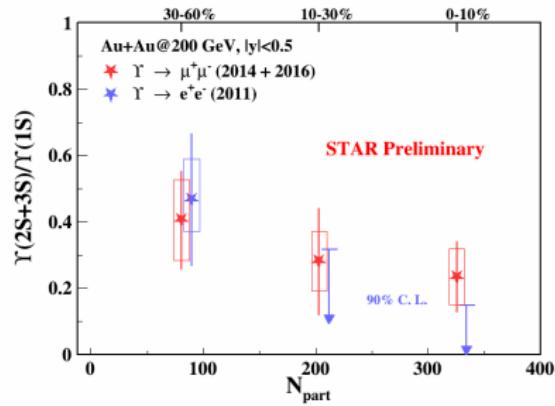
R_{AA} vs. p_T vs. models



[Phys.Rev.D 97,(2018)016017], [Phys.Rev.C 96,(2017)054901]

R_{AA} vs. p_T vs. models

- Both models consistent with the data
- Rothkopf's model slightly lower than $\Upsilon(2S + 3S)$
- Flat vs. p_T



$\frac{\gamma(2S+3S)}{\gamma(1S)}$ vs. N_{part}

- Both models consistent with the data