

Conventional and prompt fluxes in neutrino telescopes

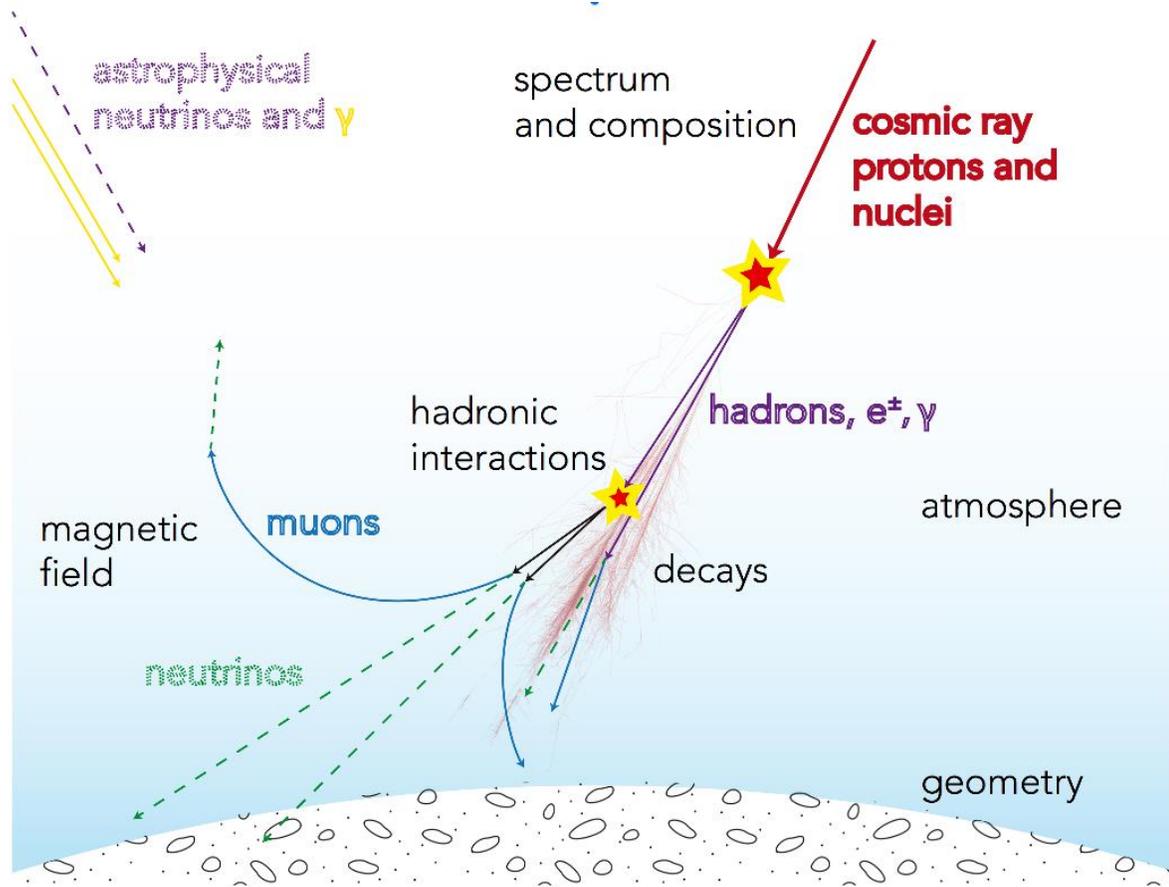
Anatoli Fedynitch
ICRR, University of Tokyo, Japan

October 07
MITP HQHP workshop



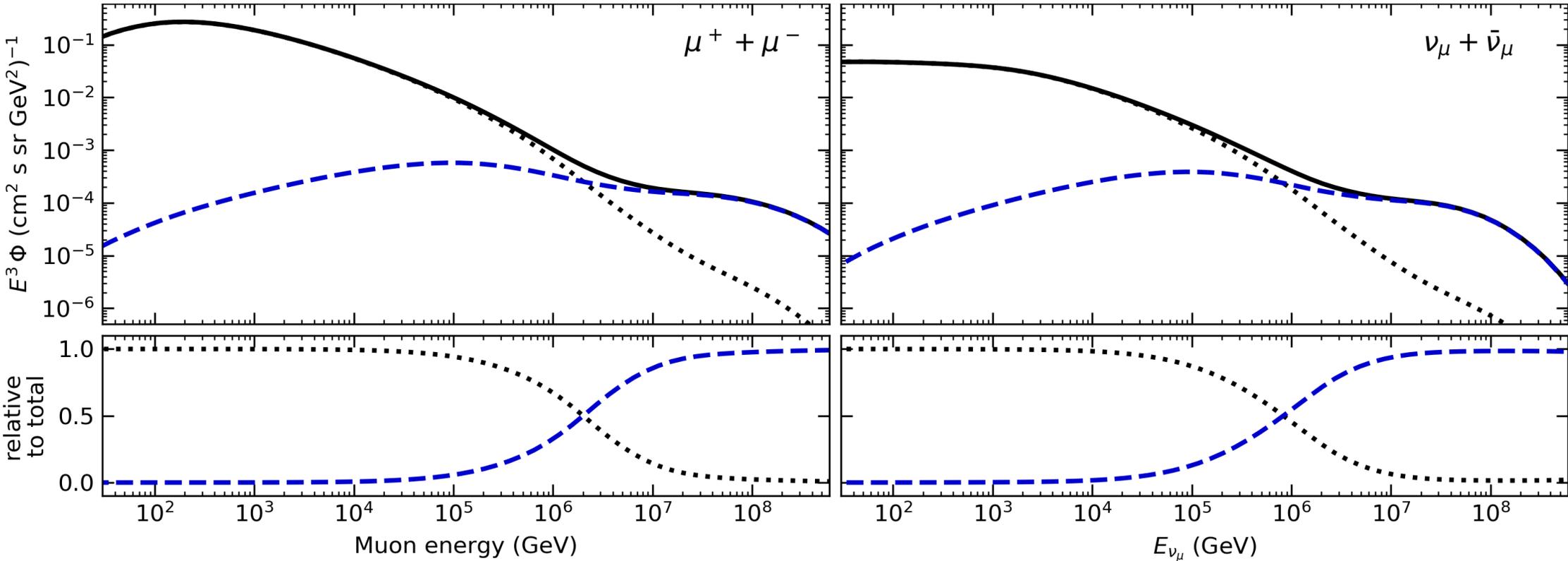
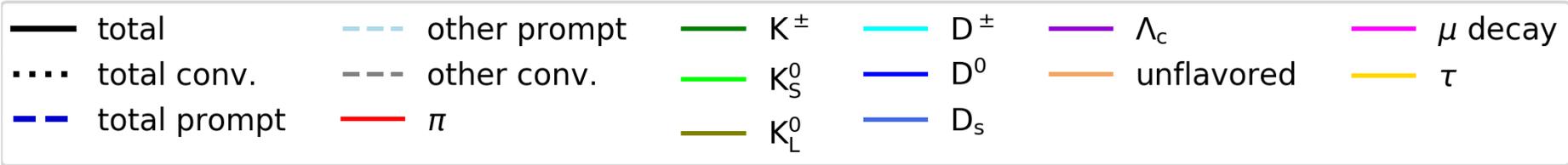
Atmospheric neutrinos

Ingredients for high-precision atmospheric neutrino flux calculation

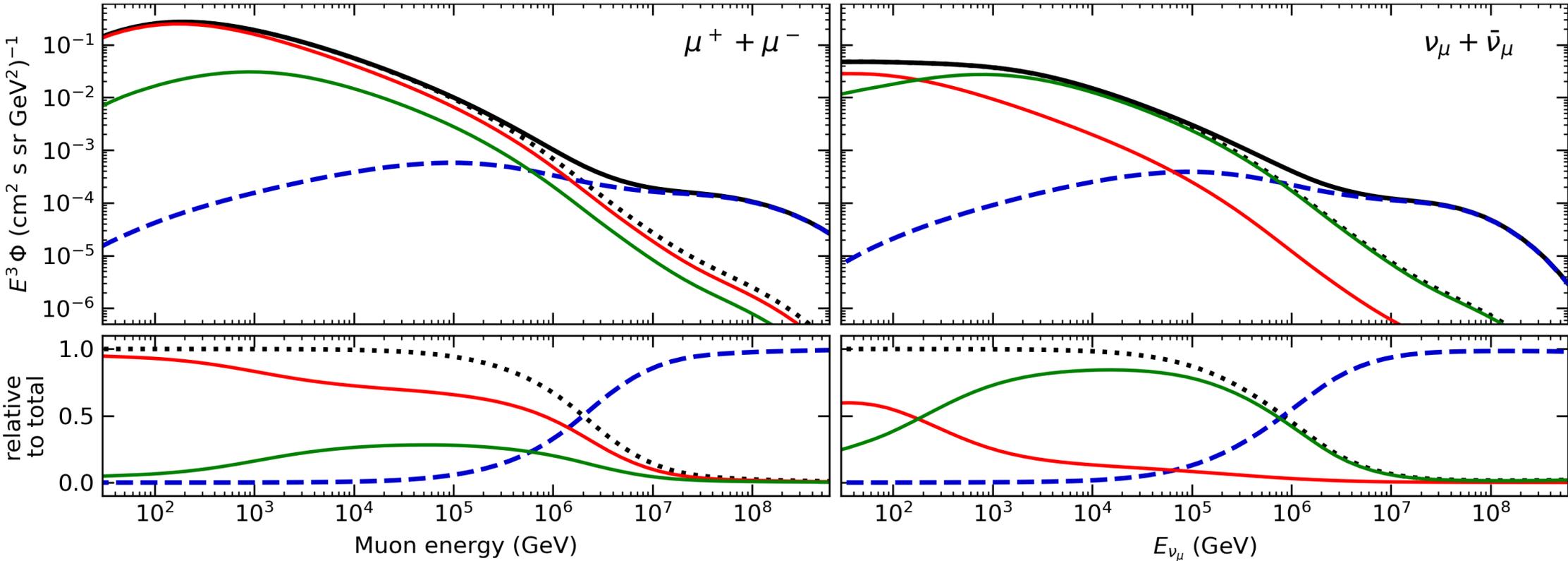
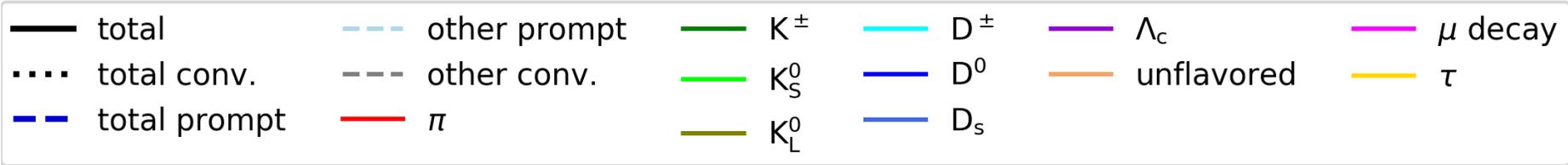


- For high precision calculations all phenomena need accurate modeling
- Uncertain “ingredients”:
 - Cosmic ray spectrum and composition
 - Hadronic interactions
 - Atmosphere (dynamic, depends on use case)
 - (Rare) decays
 - Geometry, magnetic fields, solar modulation
- No clear prescription how to handle uncertainties.
- Energy range MeV – EeV!

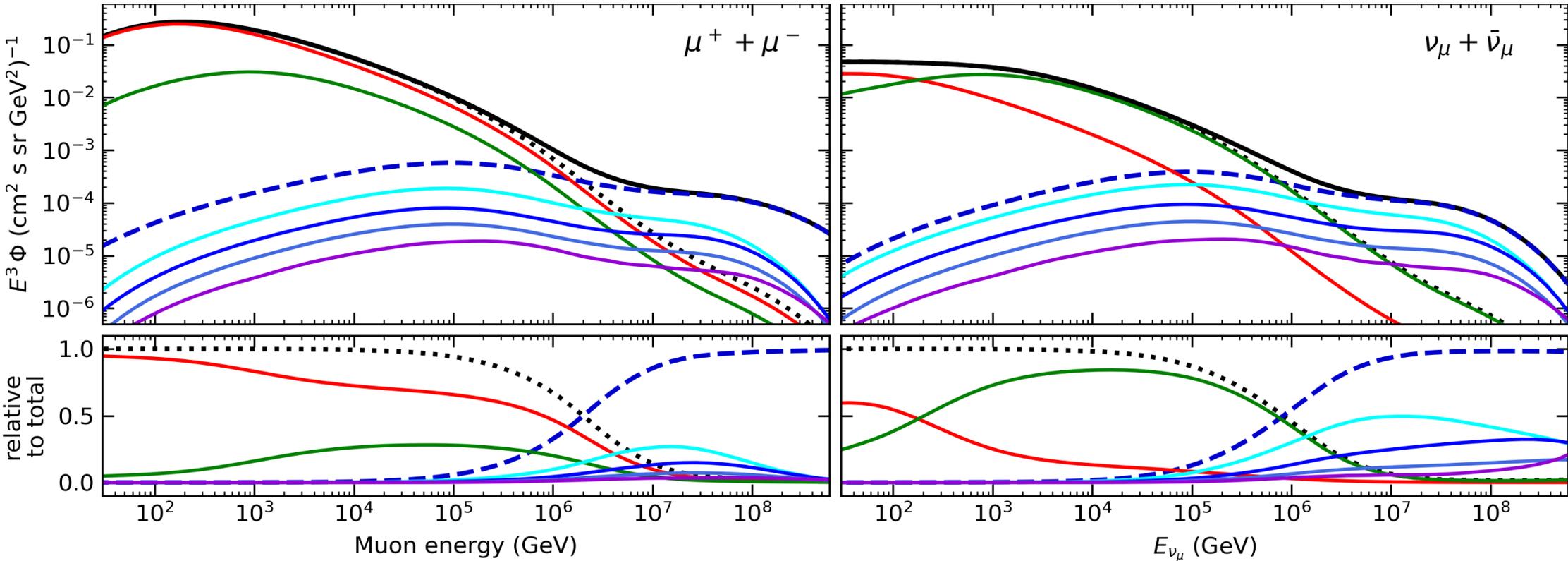
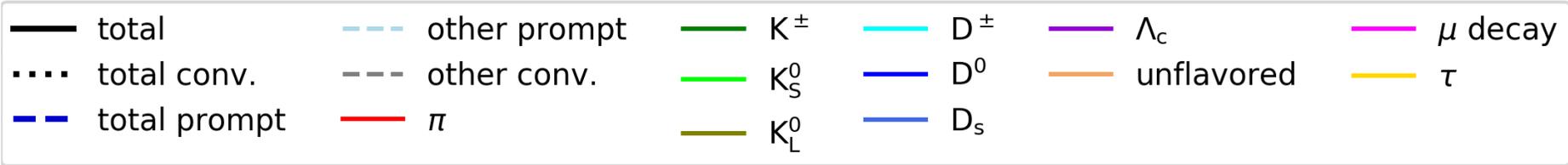
Hadrons contributing to muonic leptons



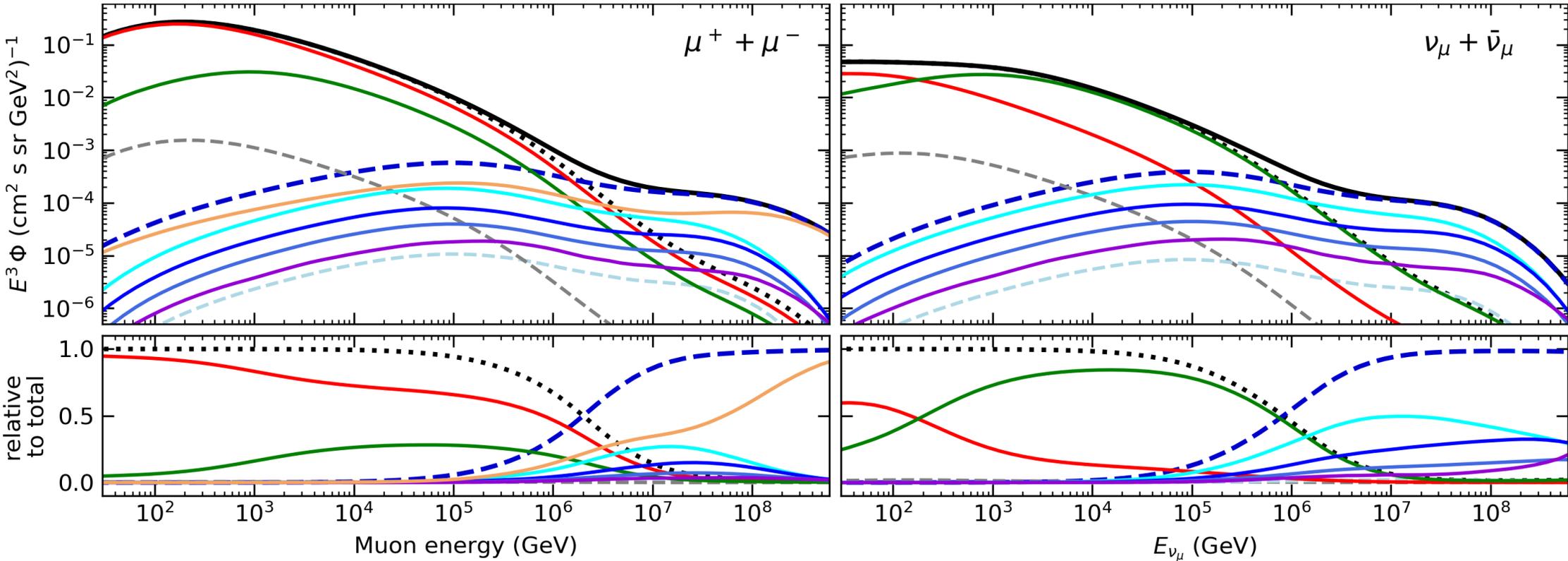
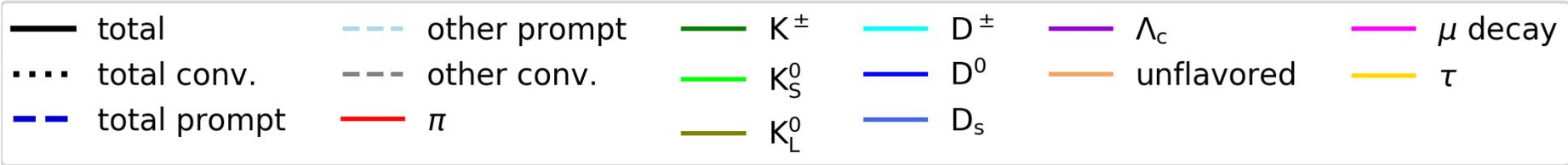
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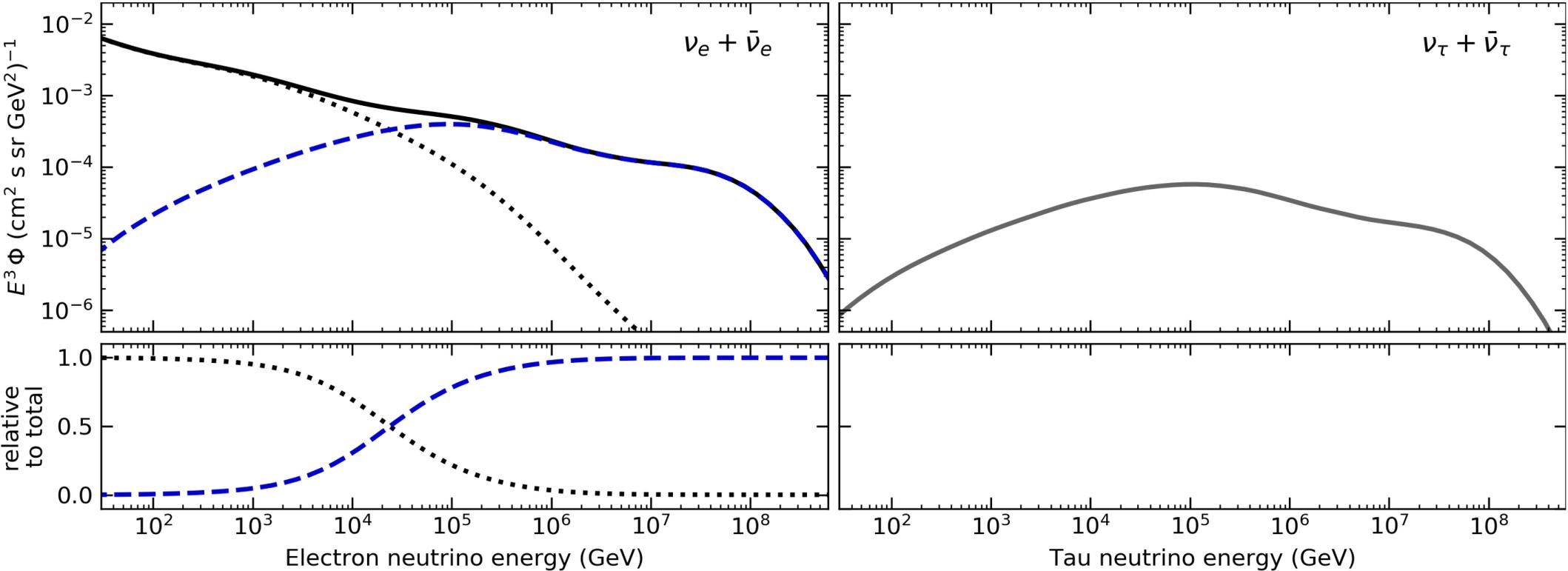
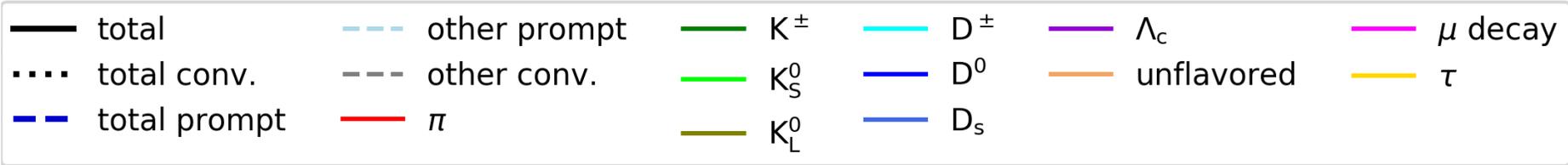


Hadrons contributing to muonic leptons



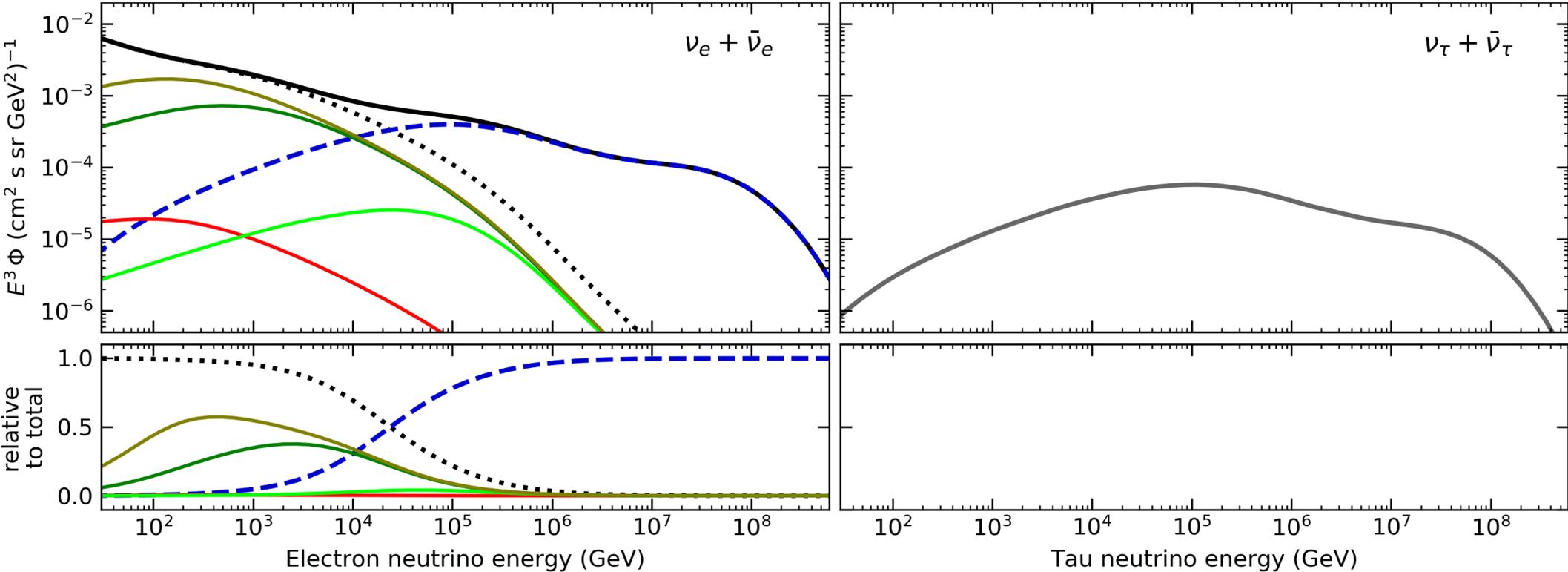
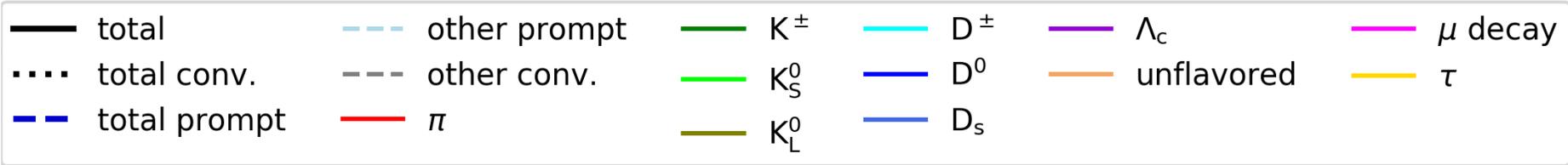
Hadrons contributing to electron and tau neutrinos

arXiv:1806.04140



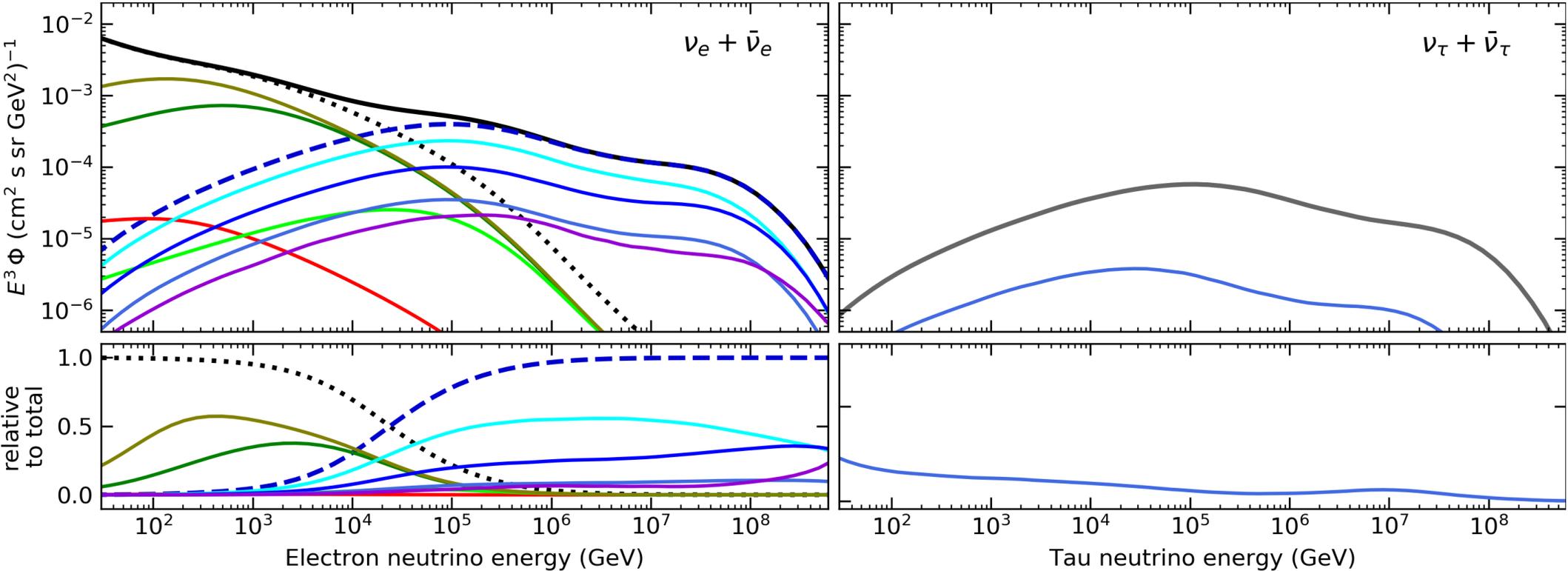
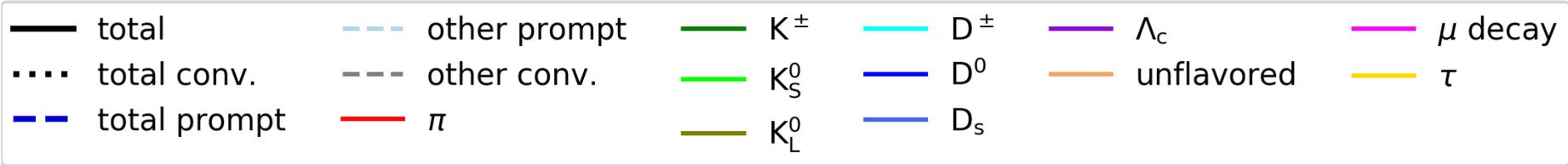
Hadrons contributing to electron and tau neutrinos

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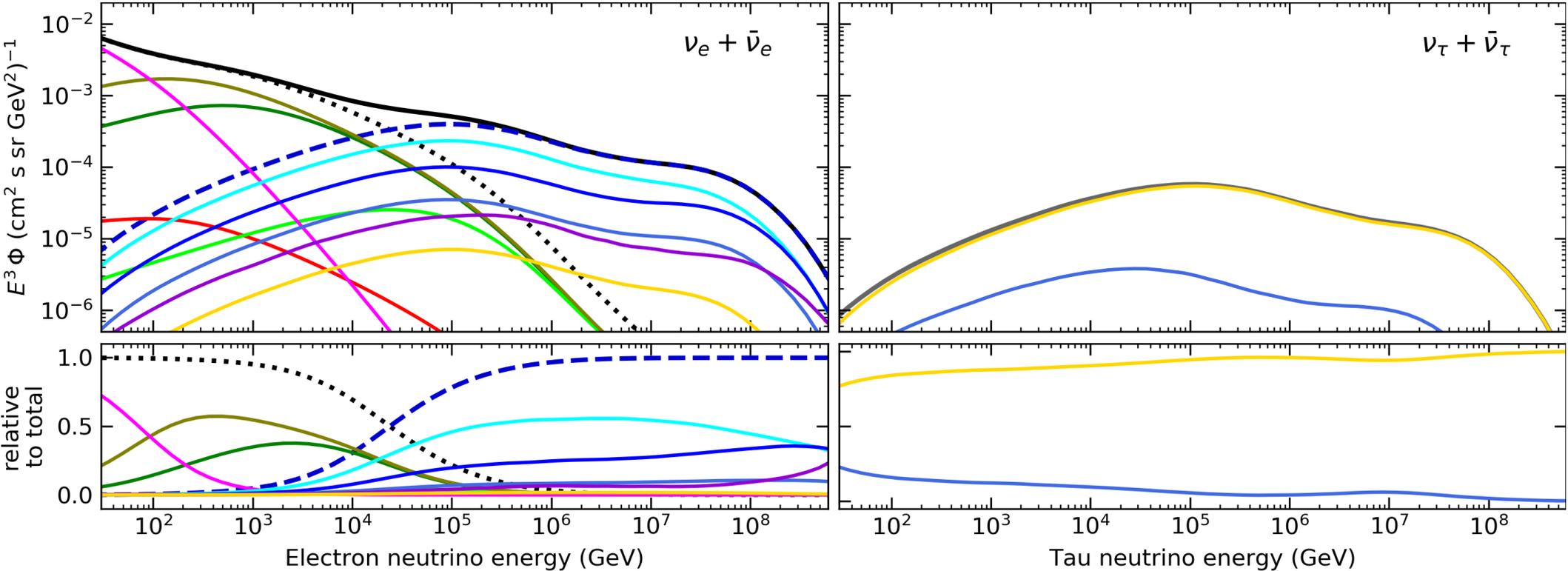
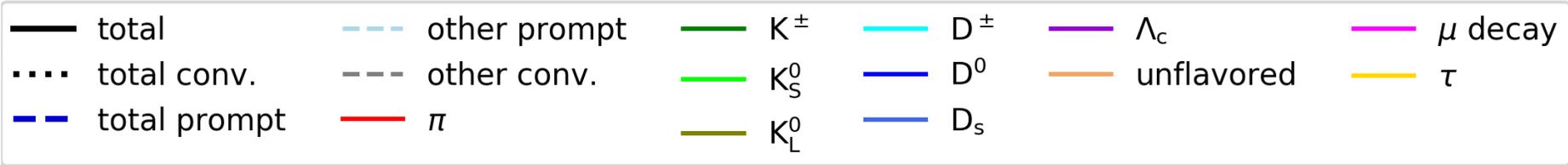
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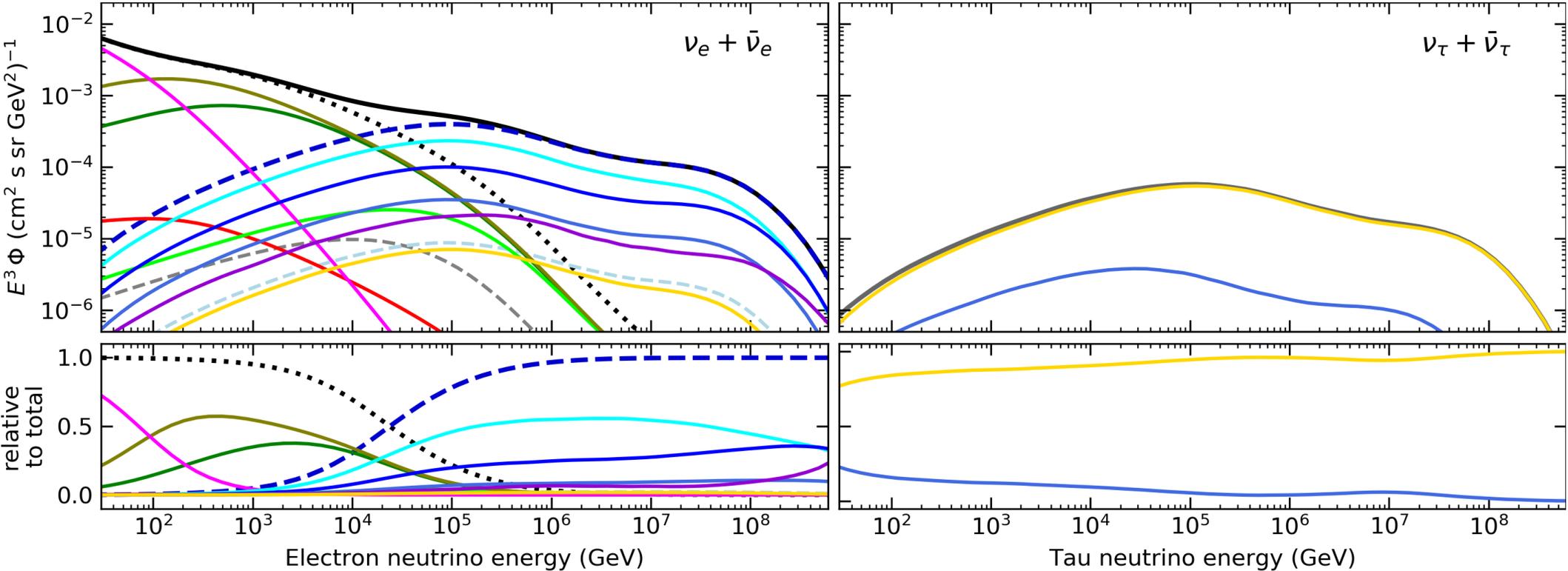
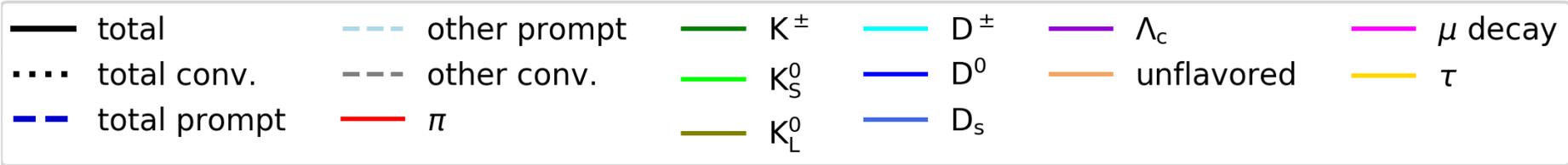
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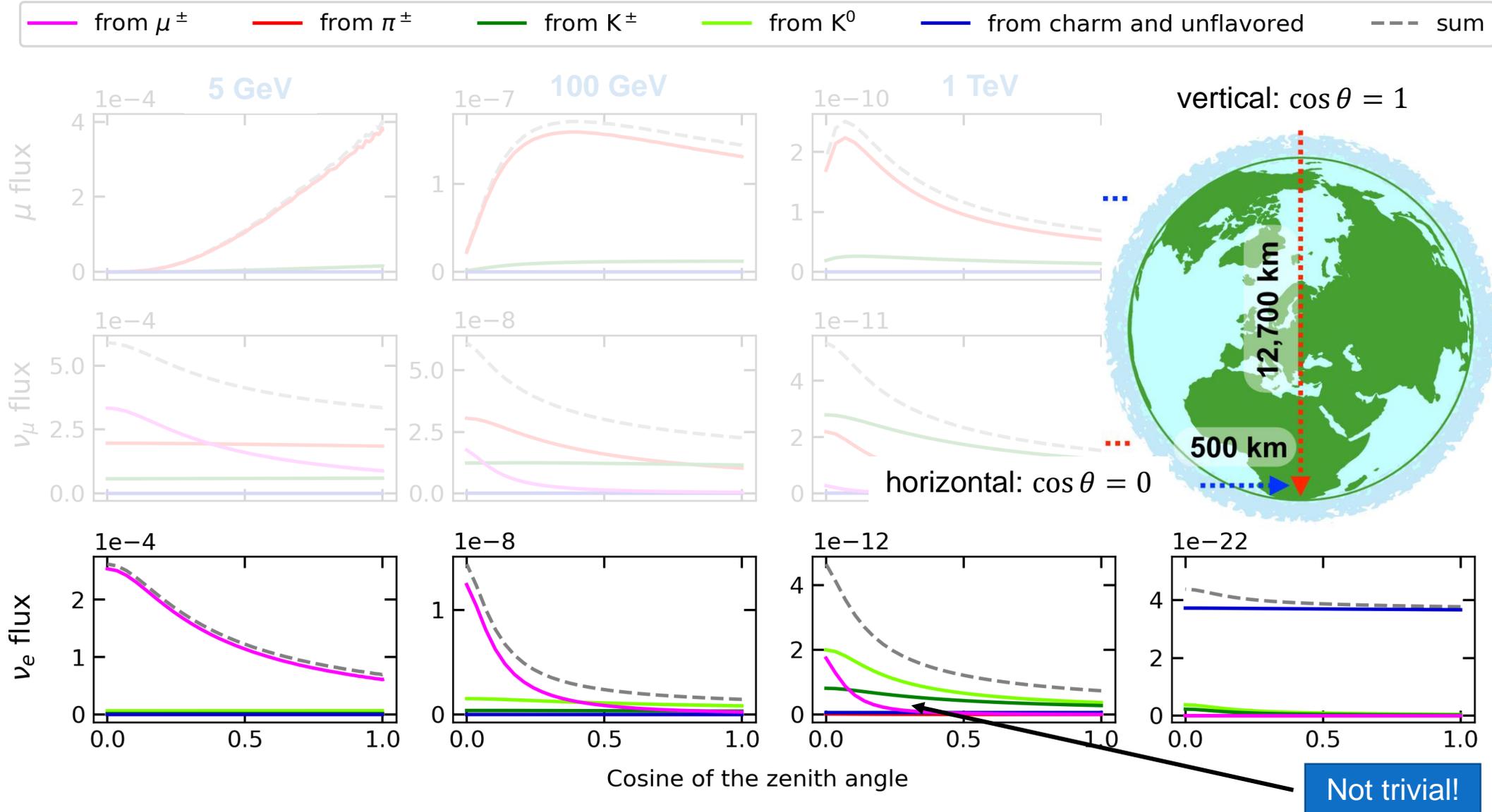


Hadrons contributing to electron and tau neutrinos

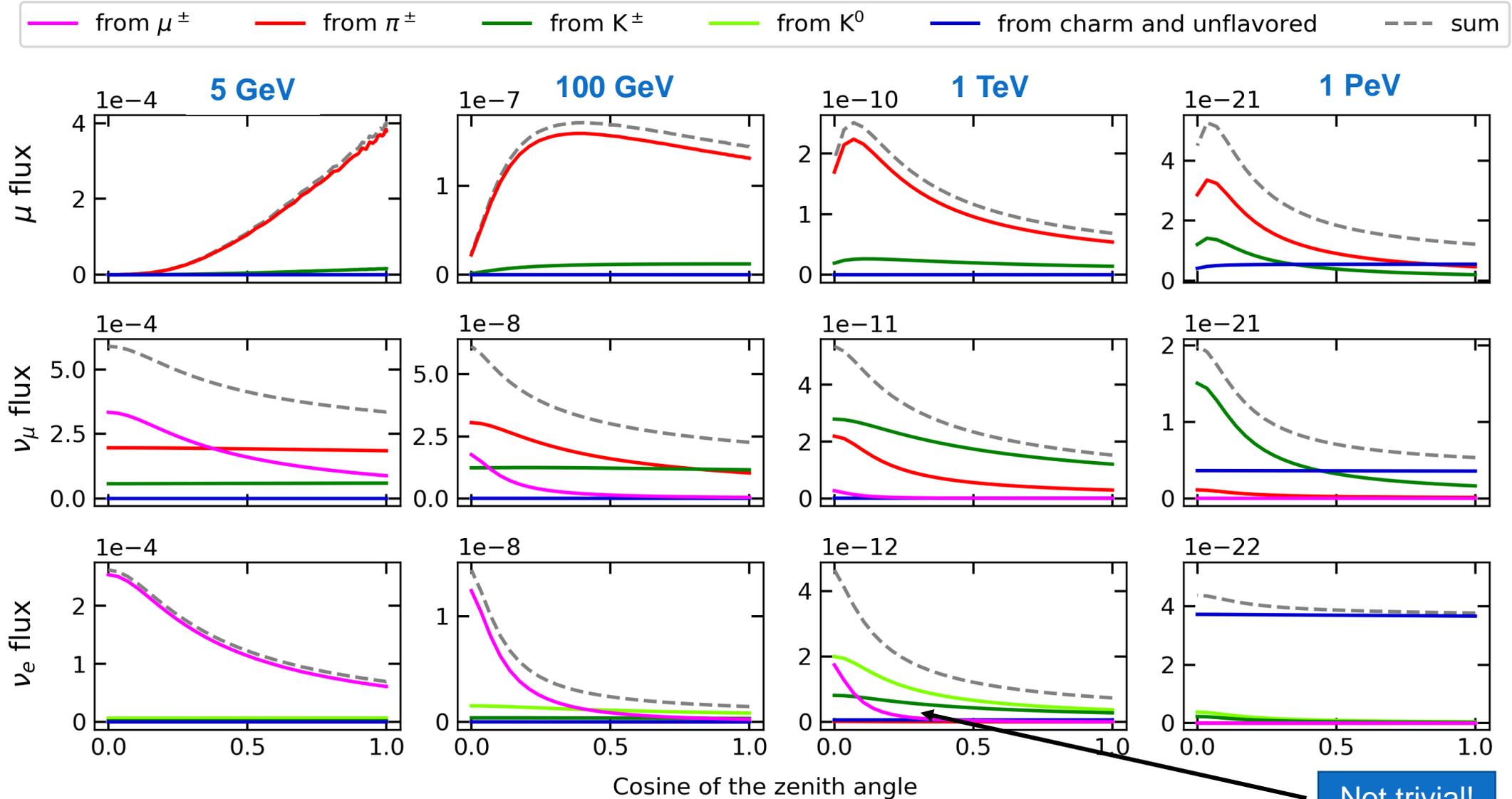
arXiv:1806.04140



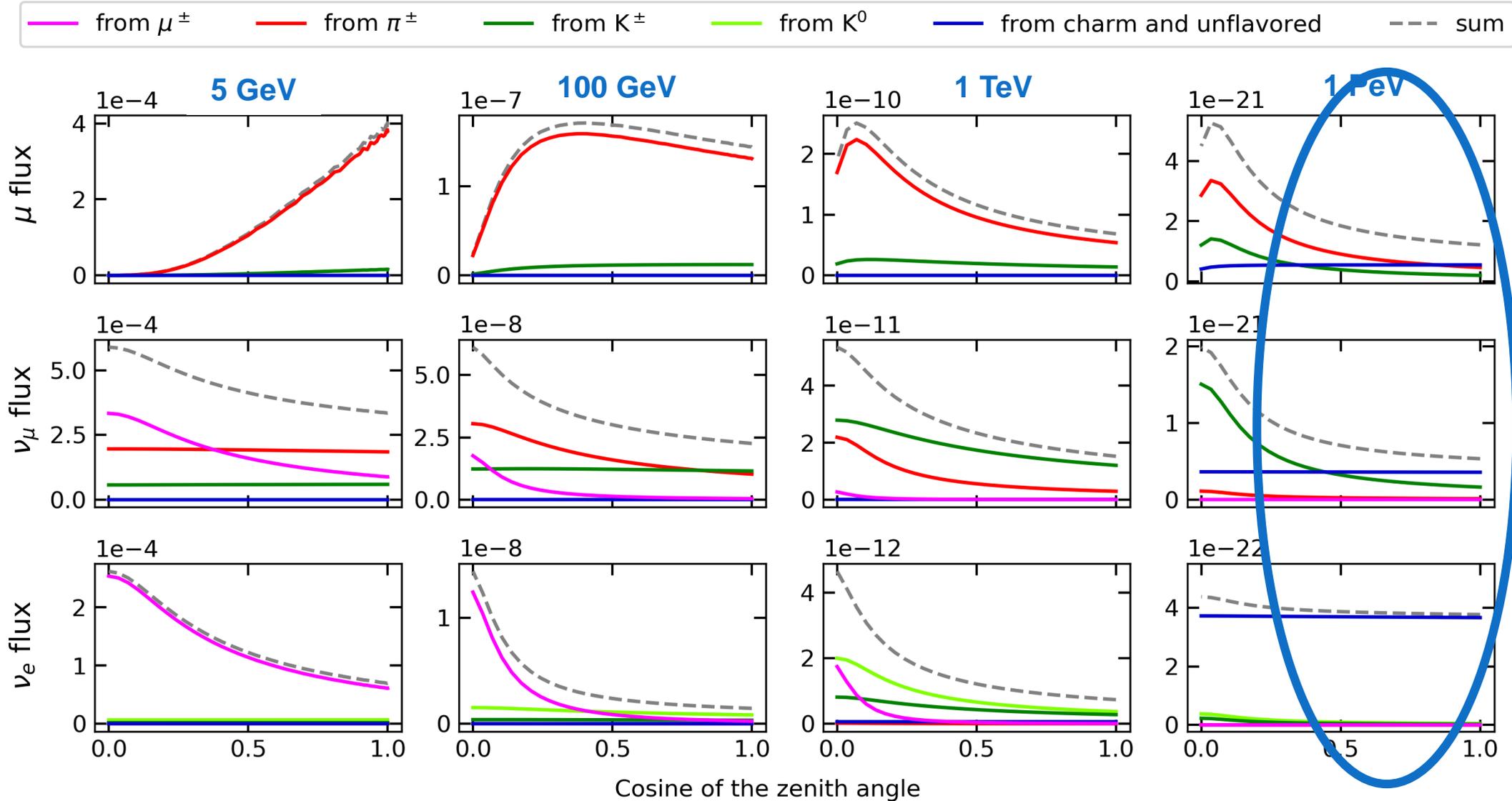
Different hadronic components shape the zenith distribution



Different hadronic components shape the zenith distribution

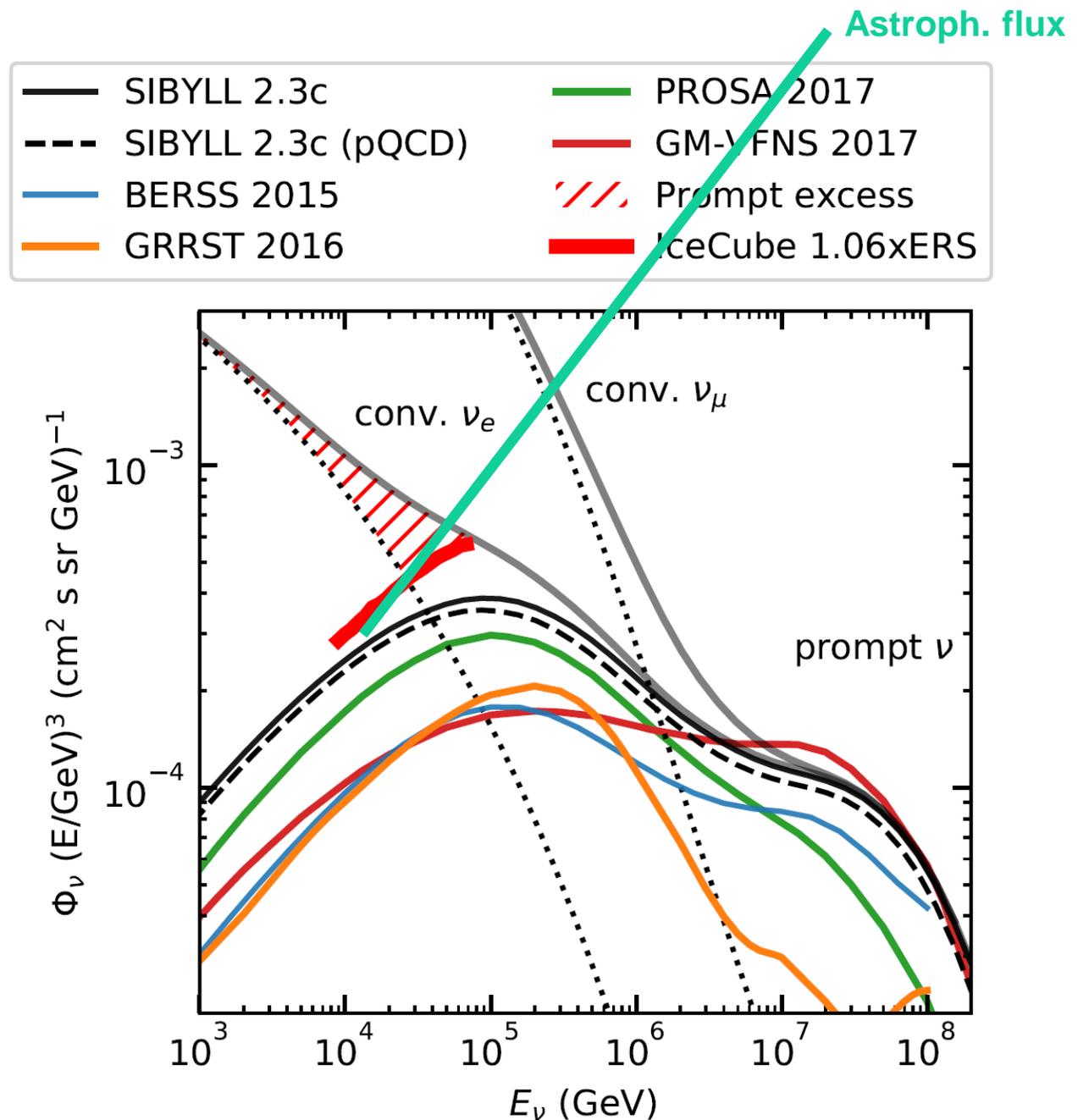


Different hadronic components shape the zenith distribution



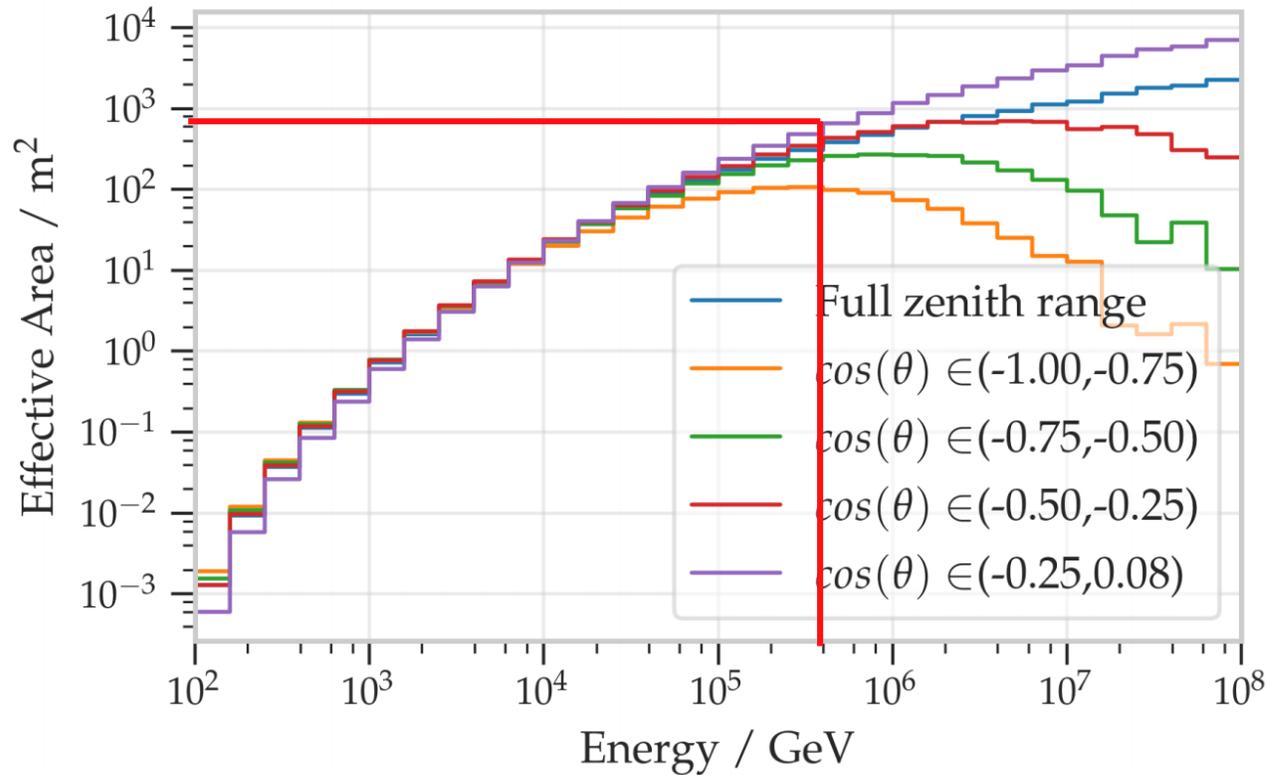
Prompt ν_μ and ν_e fluxes

- Prompt **muon neutrino flux** **buried** under conventional and astrophysical flux
- This depends on the zenith angle (Tom's talk)
- Clear excess over conventional background in electron neutrinos (hatched area)



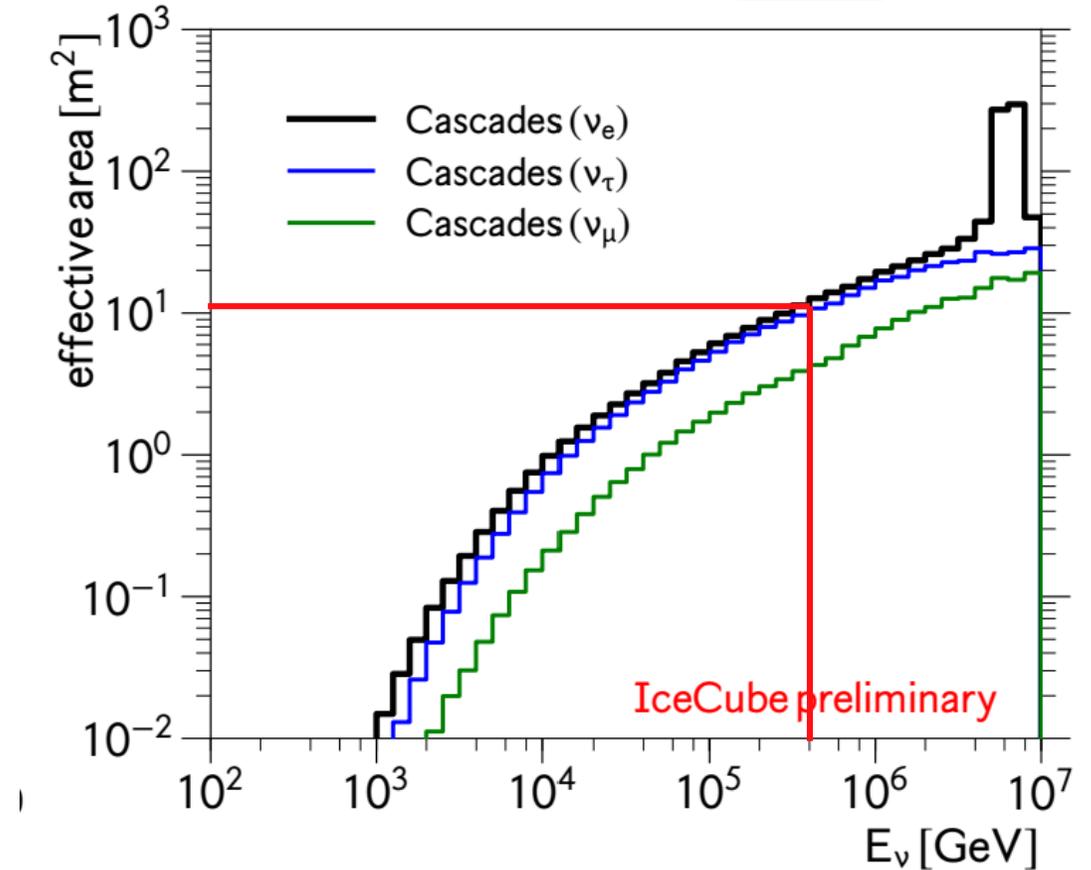
Tracks vs. Cascades effective areas

IceCube, Diffuse NuMu, 10 years



J. Stettner (IceCube Collaboration),
ICRC, arxiv: 1908.09551

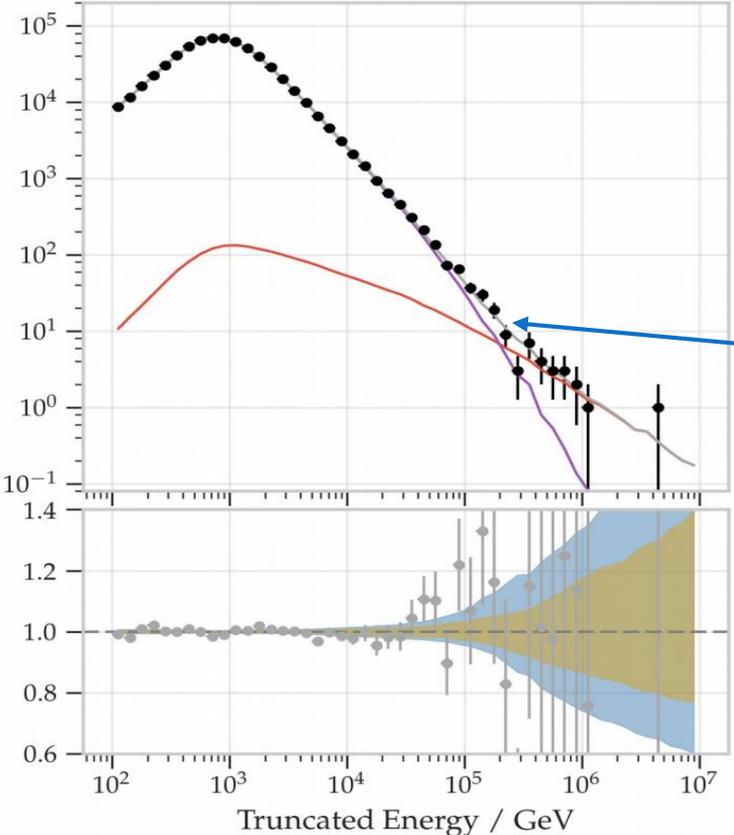
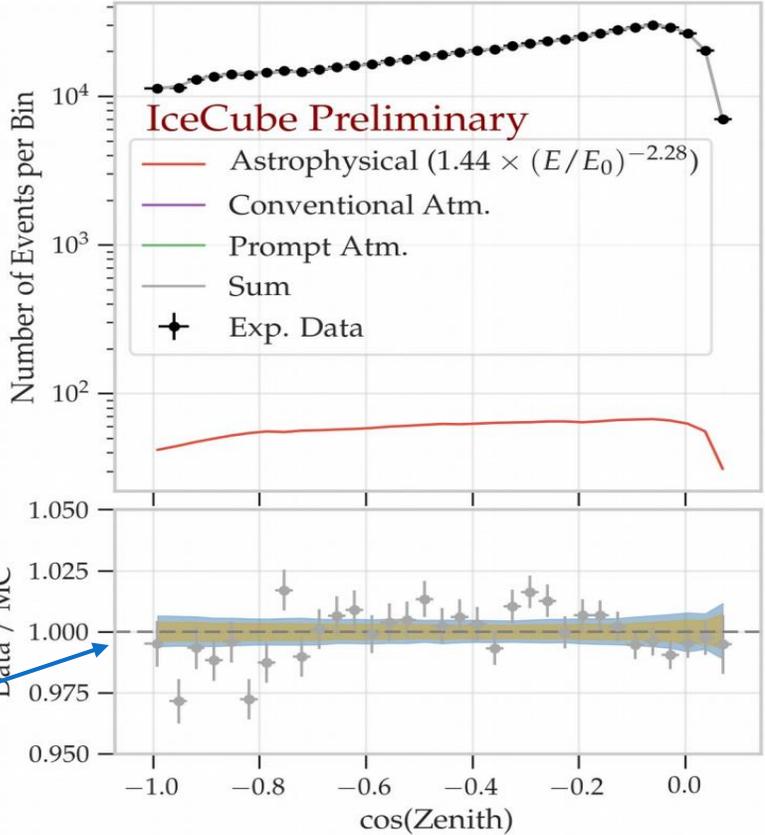
IceCube, Cascades, 4 years



H. Niederhausen (IceCube Collaboration)
PoS(ICRC2017)968

IceCube tracks state-of-the-art, latest data

Energy and Zenith Distributions

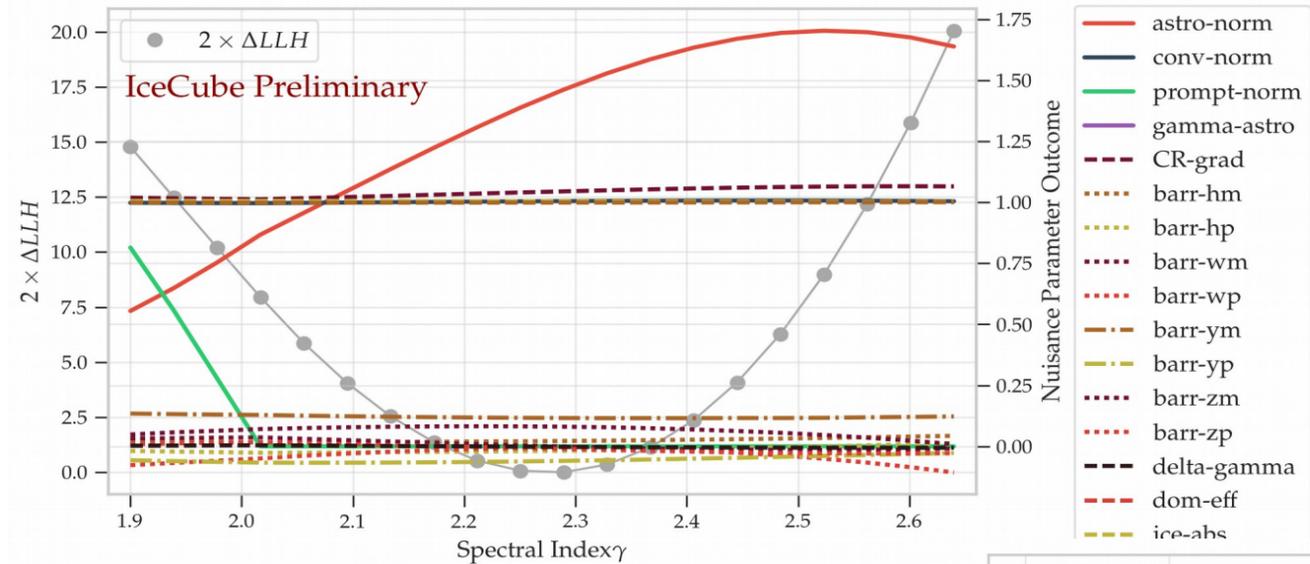


And there, but data currently lower?

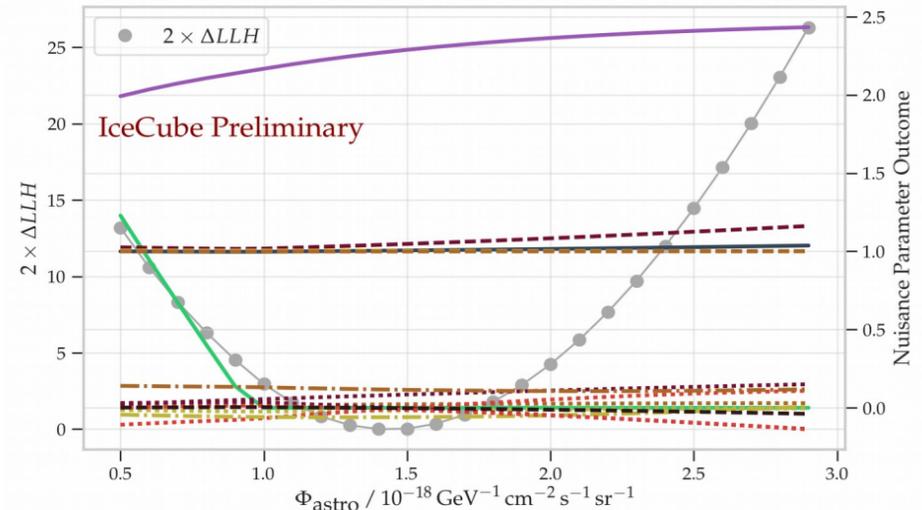
Excess of a high-energy component clearly visible
 More complex models for the astrophysical component are currently being tested

Prompt flux anti-correlated with astrophysical flux

Astrophysical Parameters and Systematics



- One-dimensional profile likelihood scans
- The two astrophysical parameters are correlated (as expected)
 - Very little correlation with other parameters
 - Only for much harder or smaller astroph. fluxes, a prompt flux is needed to fill the gaps (light green line)



Impact of prompt flux on astrophysical interpretation

- Best-fit astroph. Flux: $\frac{d\phi_{\nu+\bar{\nu}}}{dE} = (1.44^{+0.25}_{-0.24}) \left(\frac{E}{100 \text{ TeV}} \right)^{-2.28^{+0.08}_{-0.09}} \cdot 10^{-18} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$
(no prompt)

- What if a prompt flux is present ?
Repeated the fit with prompt normalization fixed to 1 x baseline prediction
 - Astro. normalization decreases
 - Spectral index hardens slightly

$$\phi_0 = 1.17 / \text{std. units}$$

$$\gamma_{\text{astro}} = 2.24$$

- Best-fit astroph. parameters, if prompt=1.0:

→ Astrophysical flux remains necessary to explain the experimental data

Challenge to advance beyond a single power-law fit

- Two power-law fit not statistically significant

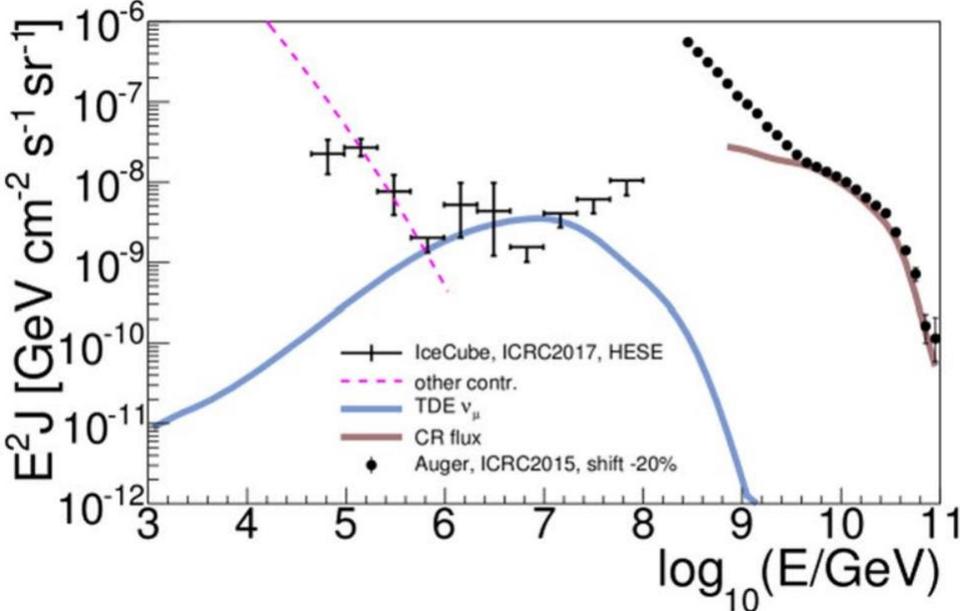
$$\frac{d\Phi_{6\nu}}{dE} = \left[\Phi_{\text{astro1}} \left(\frac{E_\nu}{100\text{TeV}} \right)^{-\gamma_{\text{astro1}}} + \Phi_{\text{astro2}} \left(\frac{E_\nu}{100\text{TeV}} \right)^{-\gamma_{\text{astro2}}} \right] \cdot 10^{-18} [\text{GeV}^{-1}\text{cm}^{-2}\text{s}^{-1}\text{sr}^{-1}].$$

- Large uncertainty/choice in assumed shape of astrophysical flux

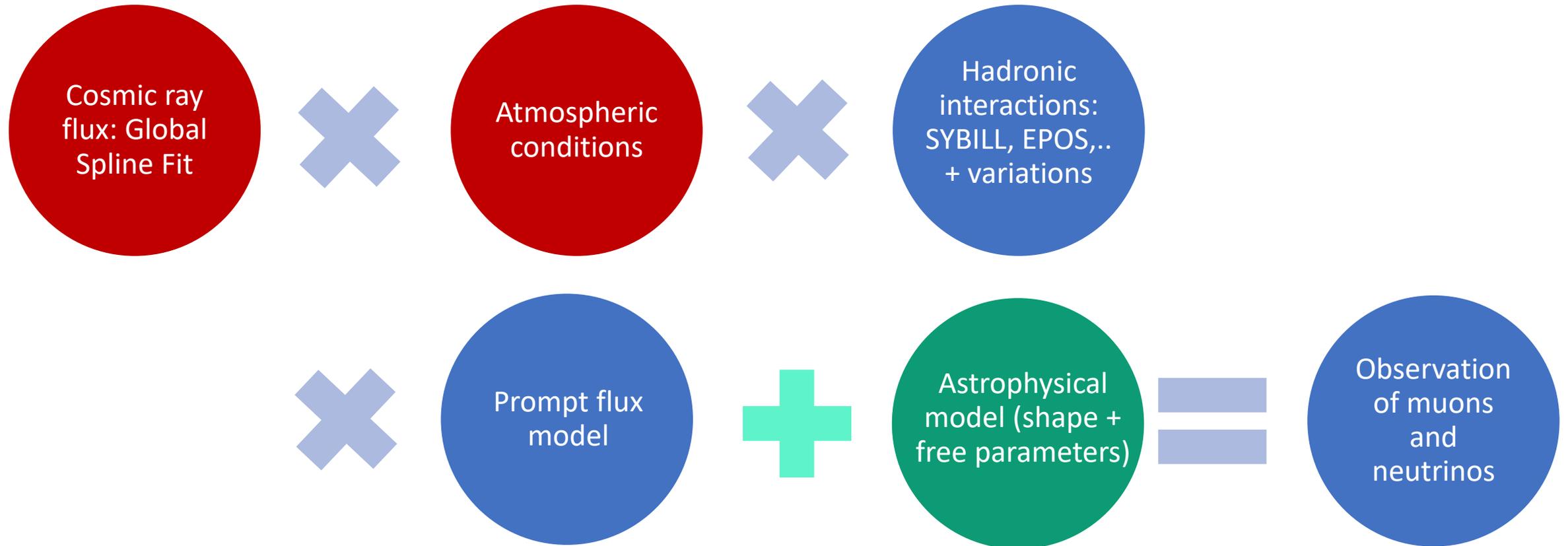
Model	Model only Bayes factor	Model + SPL Bayes factor	Most-likely SPL γ_{astro}	Most-likely SPL Φ_{astro}
Stecker [26]	4.32×10^{-13}	1.45×10^{-10}	$3.97^{+0.54}_{-0.47}$	$4.08^{+1.8}_{-1.13}$
Fang et al. [27]	0.281	0.248	$3.83^{+0.81}_{-0.5}$	$2.56^{+1.28}_{-1.44}$
Kimura et al. (B1) [28]	4.84×10^{-6}	8.38×10^{-7}	$4.5^{+0.5}_{-0.67}$	$0.98^{+1.04}_{-0.98}$
Kimura et al. (B4) [28]	3.44×10^{-4}	0.666	$2.43^{+0.31}_{-0.26}$	$1.39^{+1.18}_{-0.77}$
Kimura et al. (two component) [28]	1.73×10^{-4}	6.12×10^{-6}	$4.15^{+0.84}_{-0.73}$	$0.0^{+0.69}_{-0}$
Padovani et al. [29]	6.20×10^{-11}	3.32×10^{-7}	$3.59^{+0.59}_{-0.34}$	$4.97^{+1.68}_{-1.46}$
Senno et al. [30]	0.256	3.52	$3.67^{+0.57}_{-0.62}$	$3.36^{+1.56}_{-1.34}$
Bartos et al. [31]	1.15×10^{-14}	2.81×10^{-16}	$4.25^{+0.75}_{-0.83}$	$0.0^{+0.49}_{-0}$
Tavecchio et al. [32]	0.0730	1.04	$3.88^{+0.65}_{-0.49}$	$3.7^{+1.39}_{-1.48}$
Biehl et al. [33]	8.66×10^{-7}	0.362	$3.35^{+0.4}_{-0.38}$	$5.09^{+2.07}_{-1.03}$

Austin Schneider (IceCube Collaboration), arXiv:1907.11266

Example for astroph. model:
Biehl et al., Nat. Sci. Rep. 2018



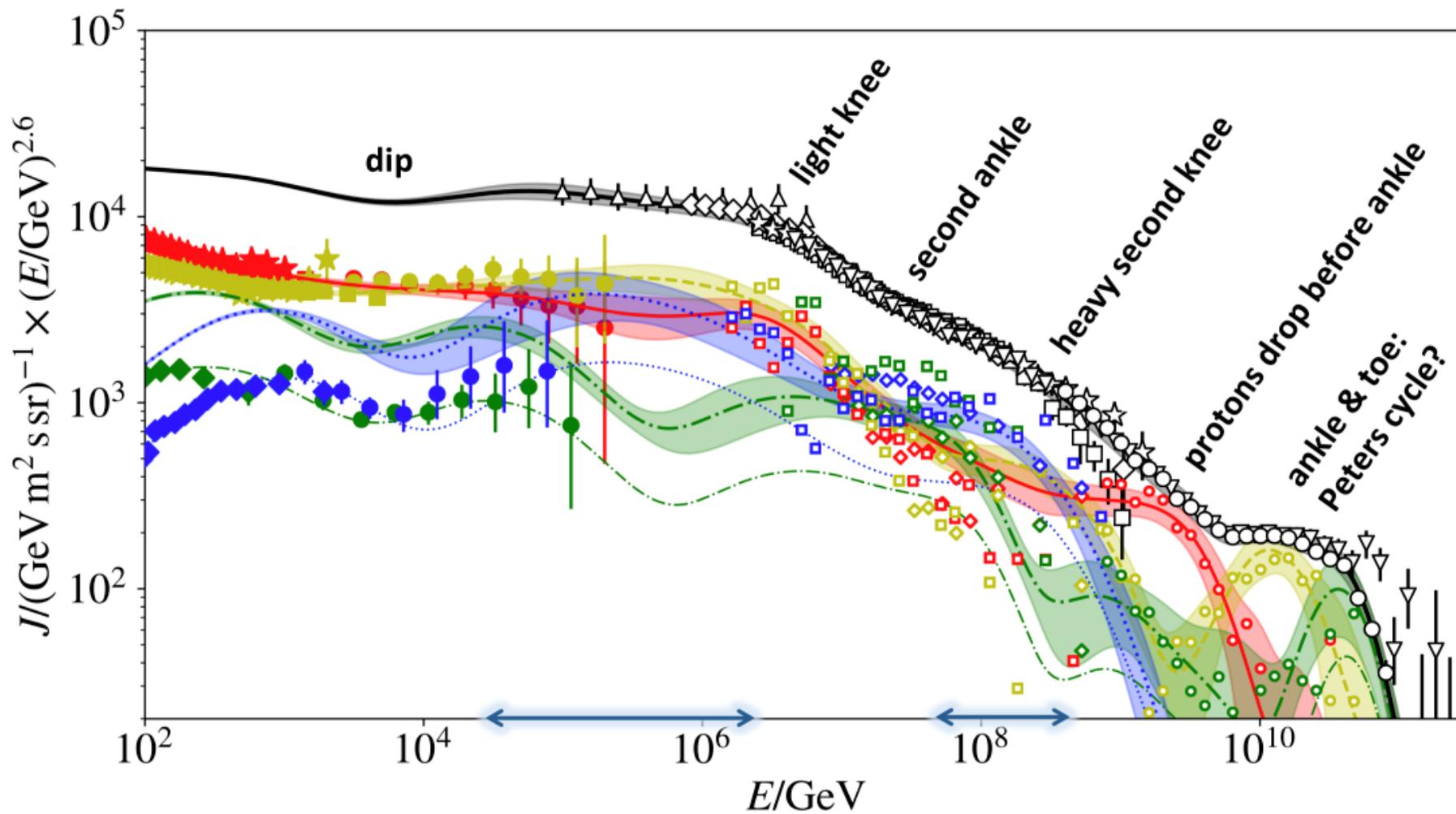
Correct parameterization of uncertainties is a chance to measure prompt fluxes and constrain charm



Circle means central prediction + error band.
Much better: free parameters + covariance matrix.

The Global Spline Fit

Dembinski, AF, Engel, Gaisser, Stanev
PoS(ICRC2017)533 & in prep.

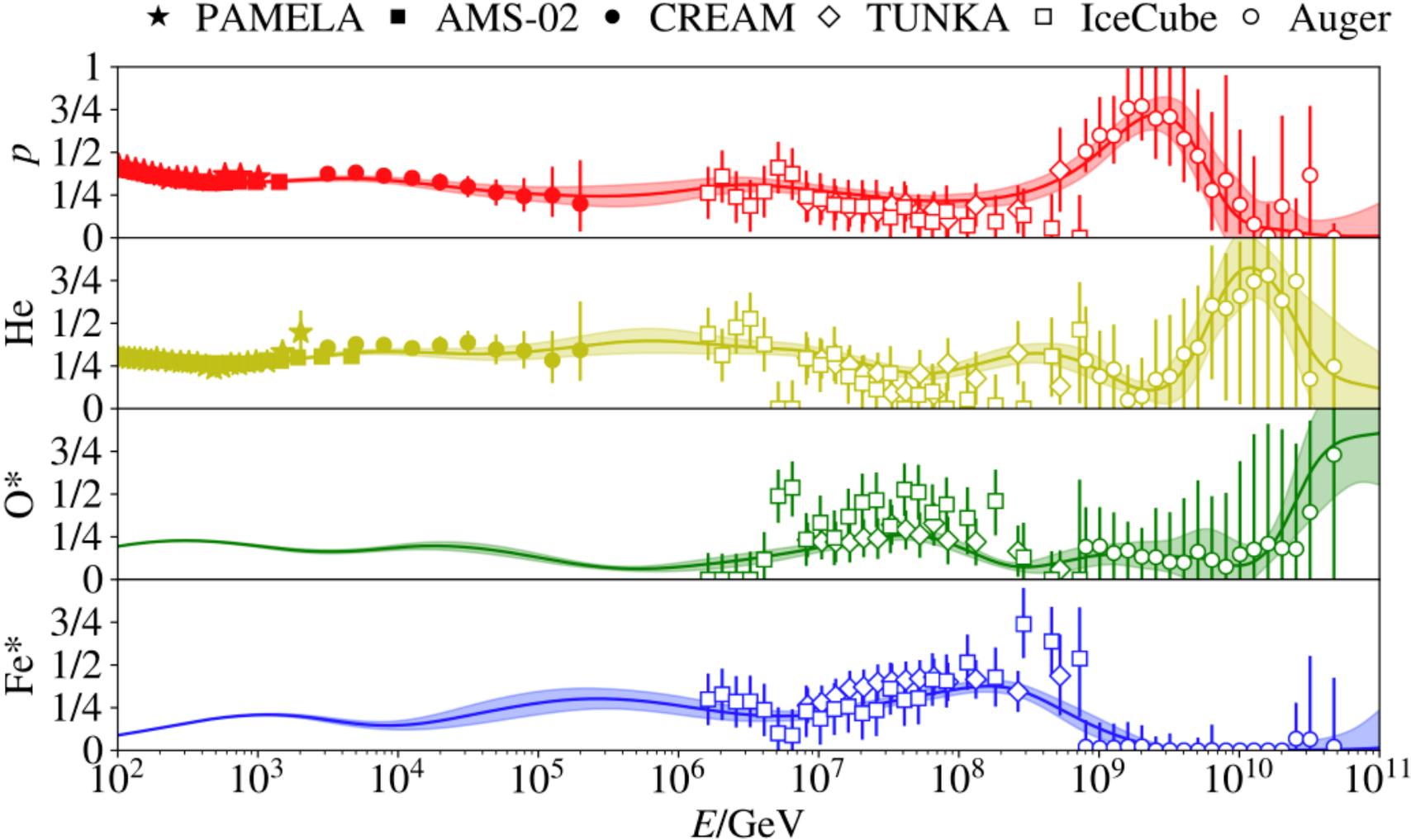


More composition data needed

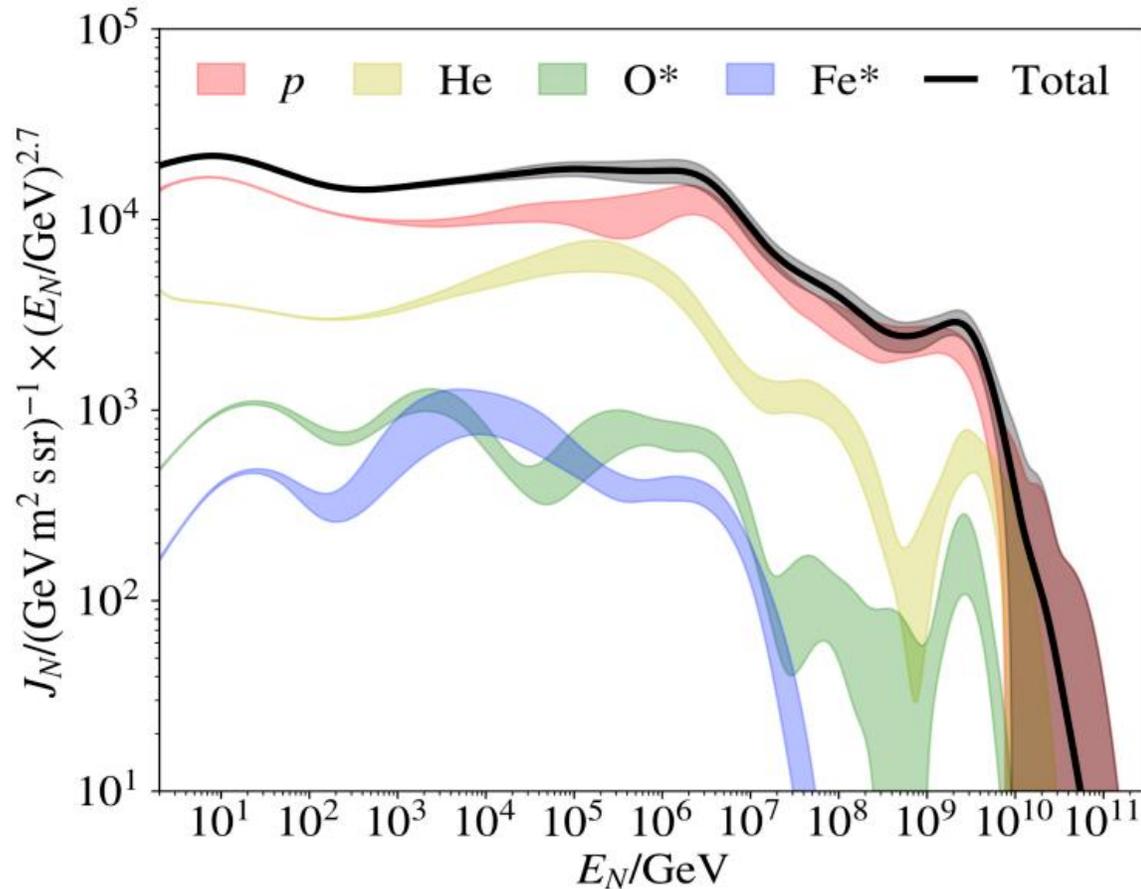
Fitted composition data

Dembinski, AF, Engel, Gaisser, Stanev
PoS(ICRC2017)533 & in prep.

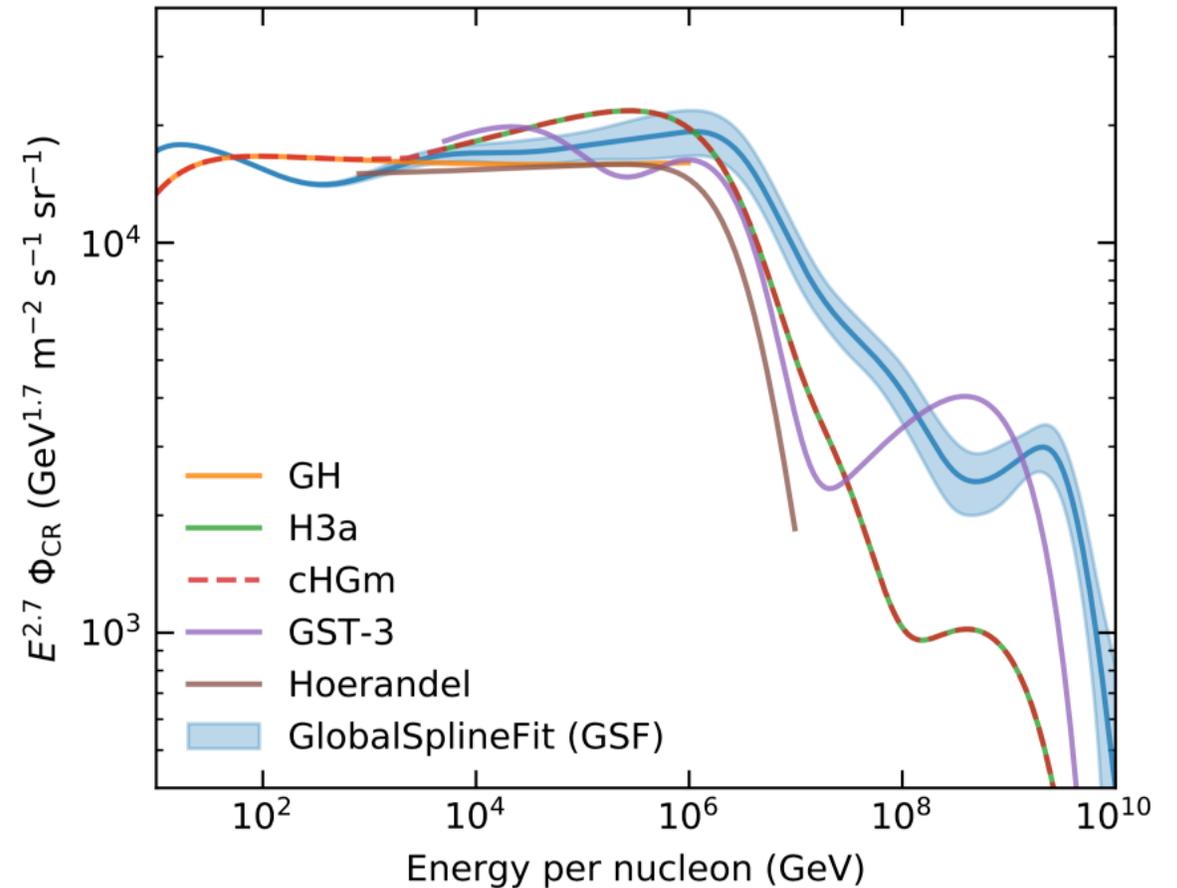
4-mass group experiments



Derived result: nucleon flux



Dominated by proton flux. Details of sub-leading elements not important.

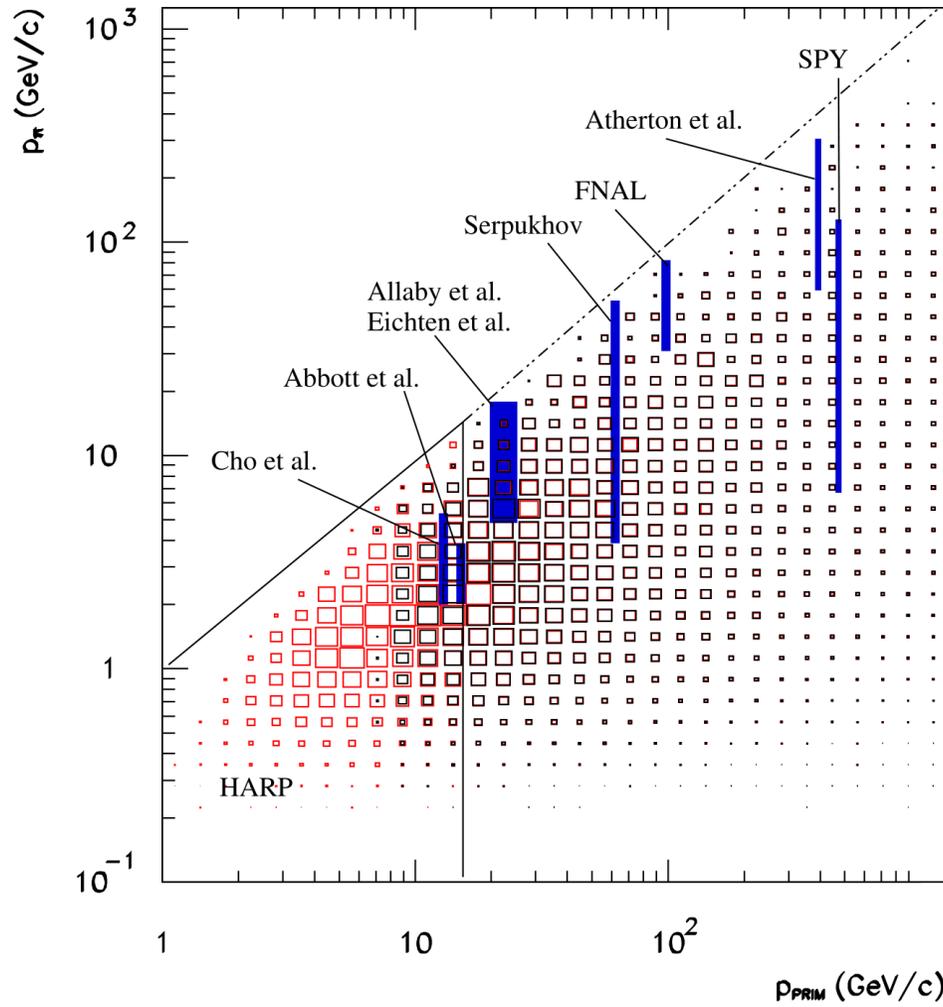


Harder spectrum at the knee due to lighter composition as assumed by 3-population models

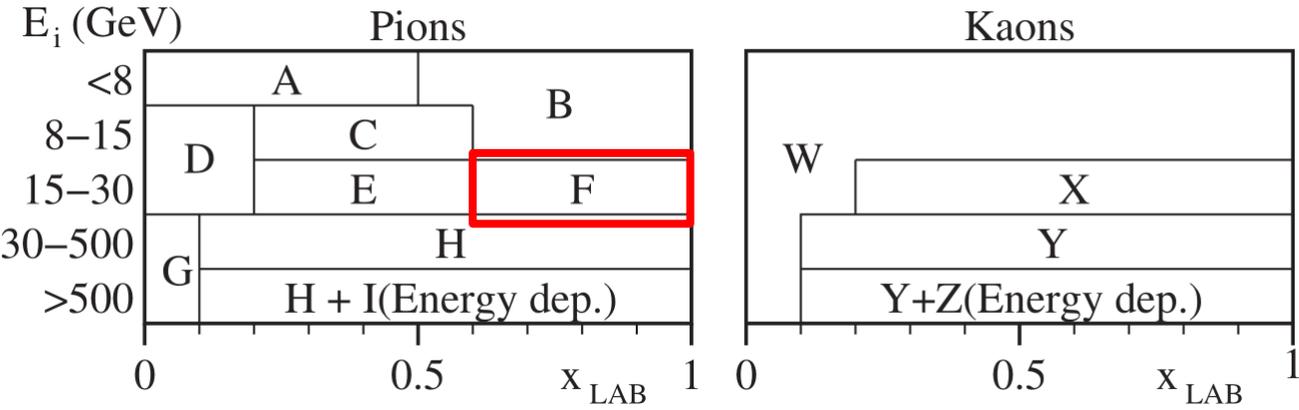
Hadronic uncertainties in new IceCube analyses

- “*Uncertainties in atmospheric neutrino fluxes*”, G. D. Barr, S. Robbins, T. K. Gaisser, and T. Stanev, Phys. Rev. D 74, 094009 (2006) (extensive discussion also in Sanuki et al. PRD 75 (2007))
- Cut phase-space in regions/slices in E_{lab} and x_{lab} and **assign** uncertainty to each slice (uncorrelated)
- Uncertainty assigned by hand and not derived from data. Assignment based on availability of data, not how well the model [TARGET2.1] describes it
- Many “free” parameters with unclear correlations

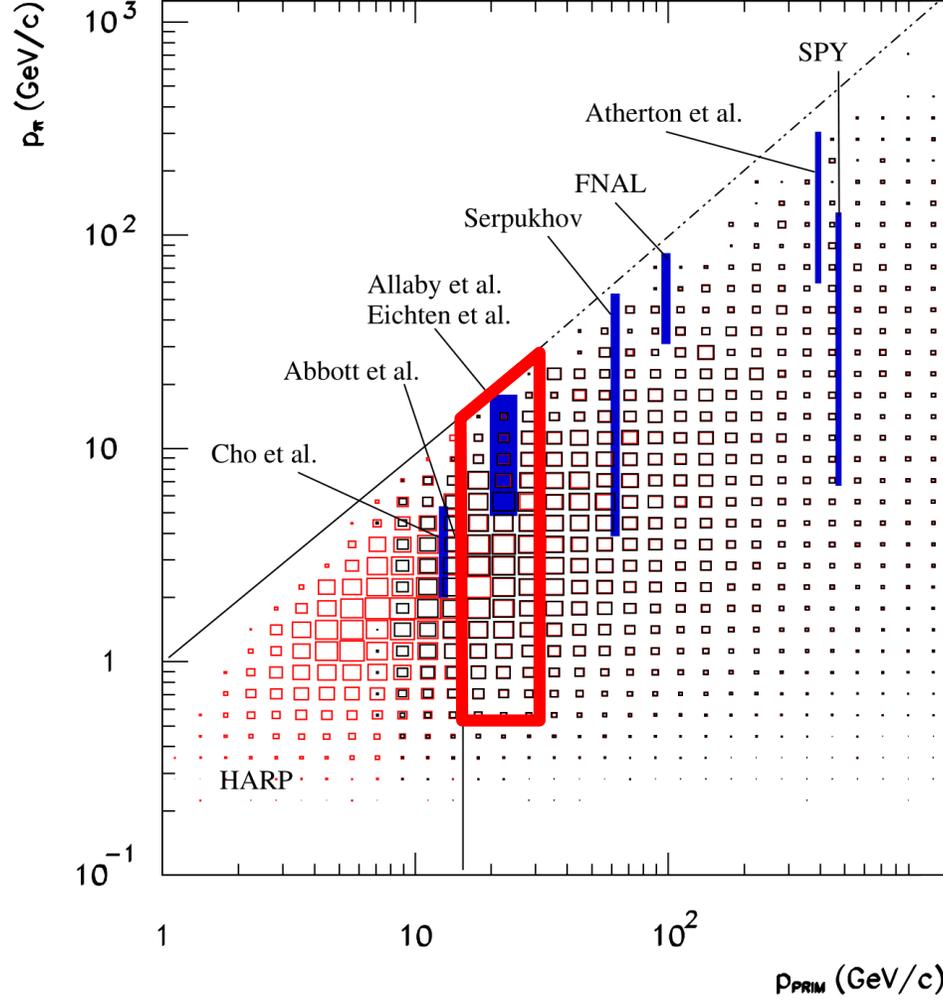
PHYSICAL REVIEW D 74, 094009 (2006)



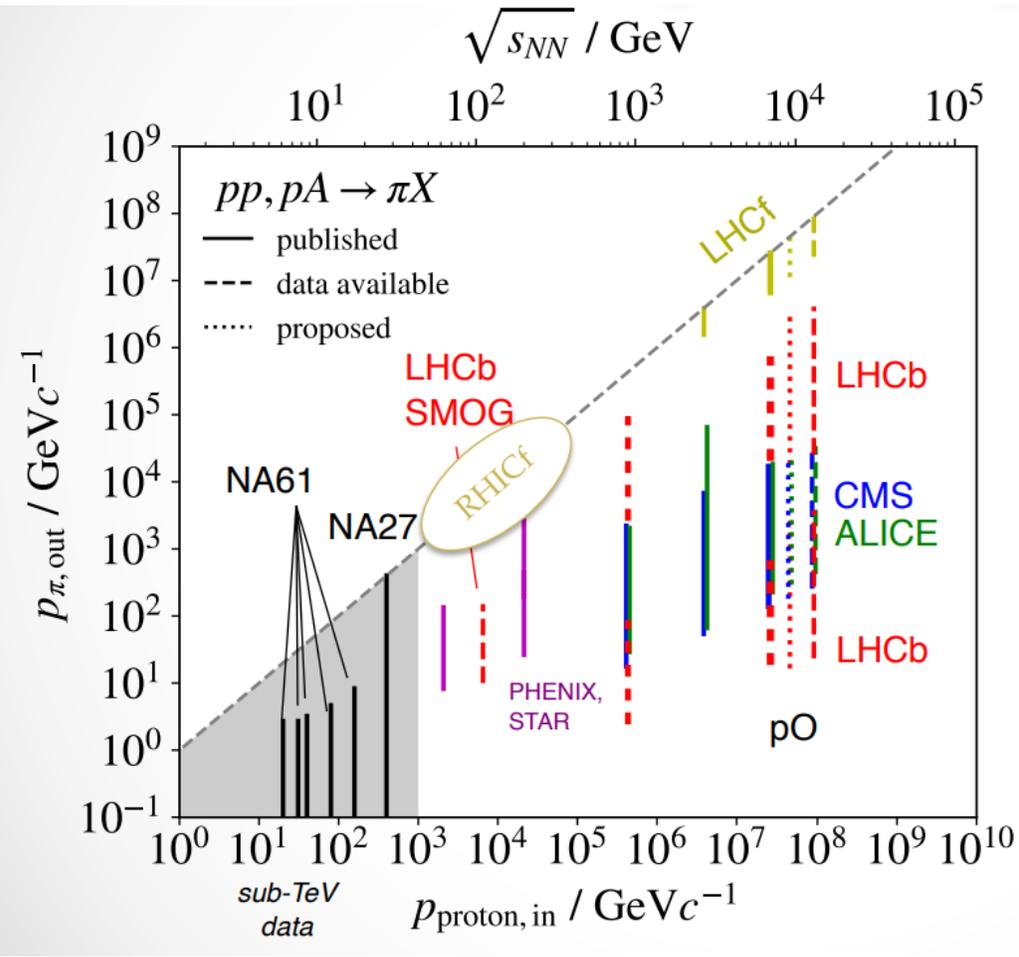
Phase space regions



PHYSICAL REVIEW D **74**, 094009 (2006)



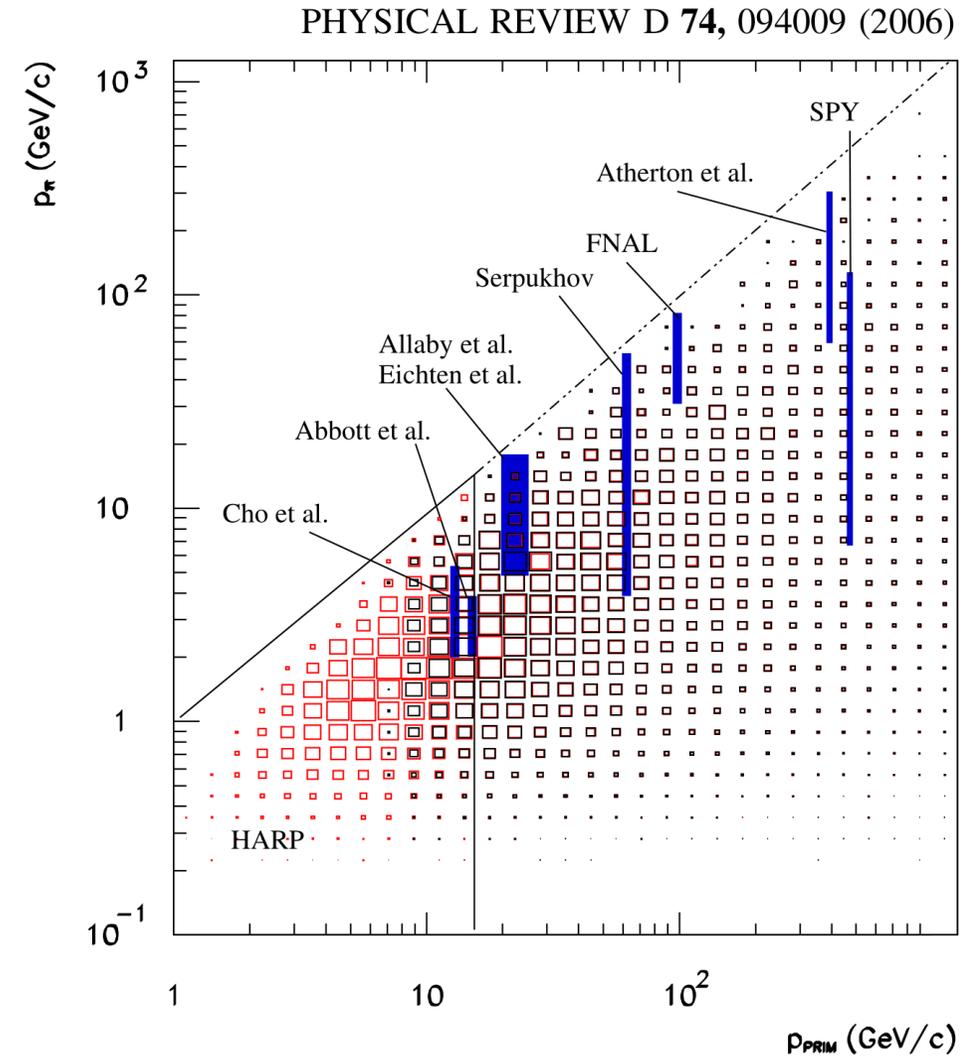
2019 coverage of phase-space



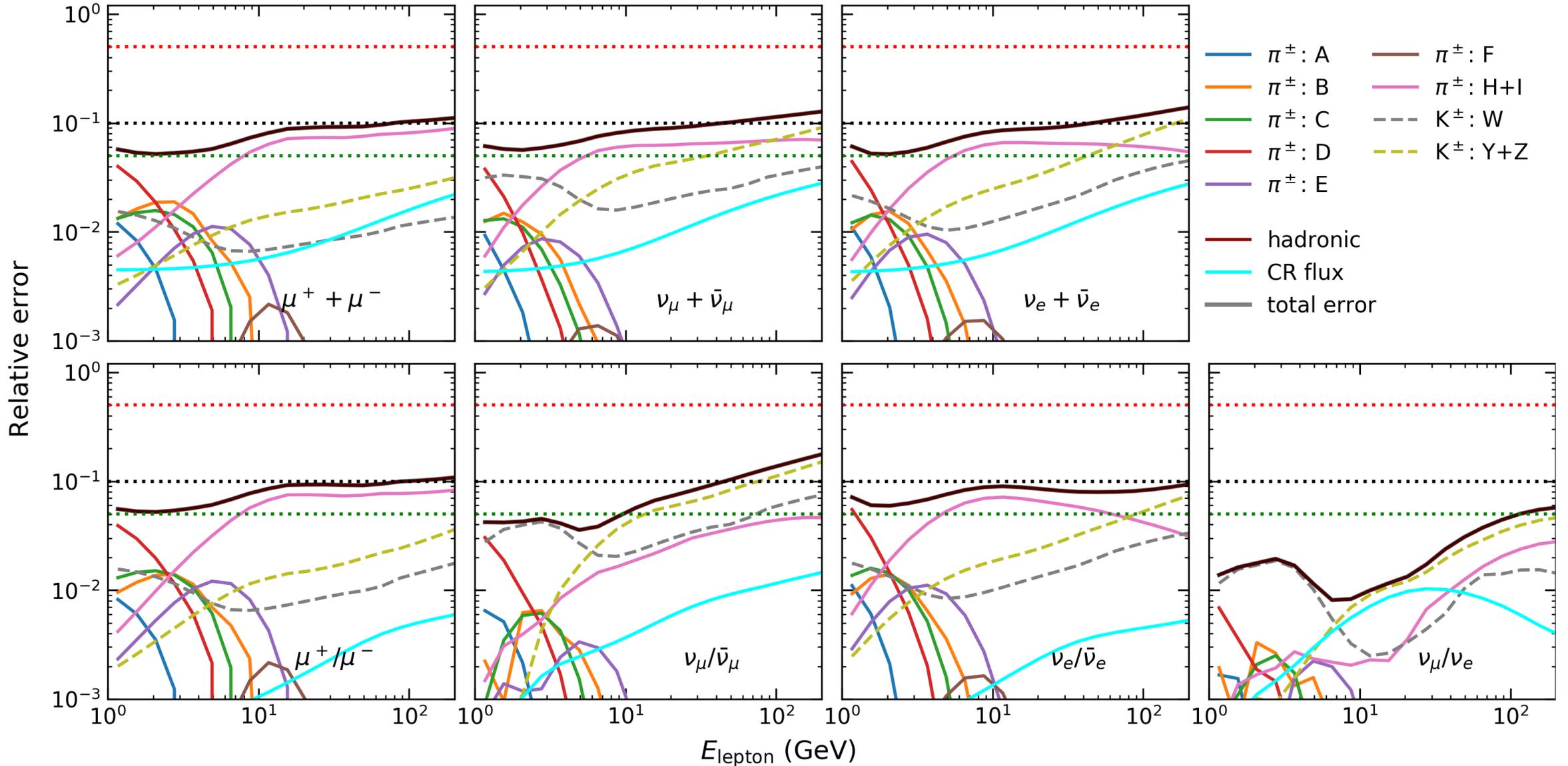
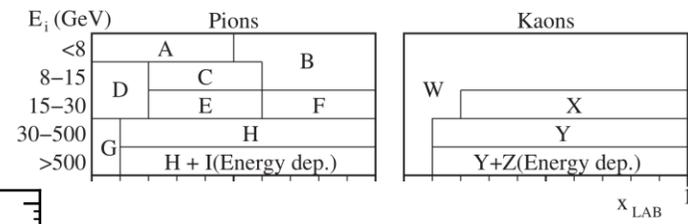
H. Dembinski,
ISVHECRI 2018

Phase space of air shower interactions as covered by various experiments (beam-beam collisions transformed to equivalent fixed-target system)

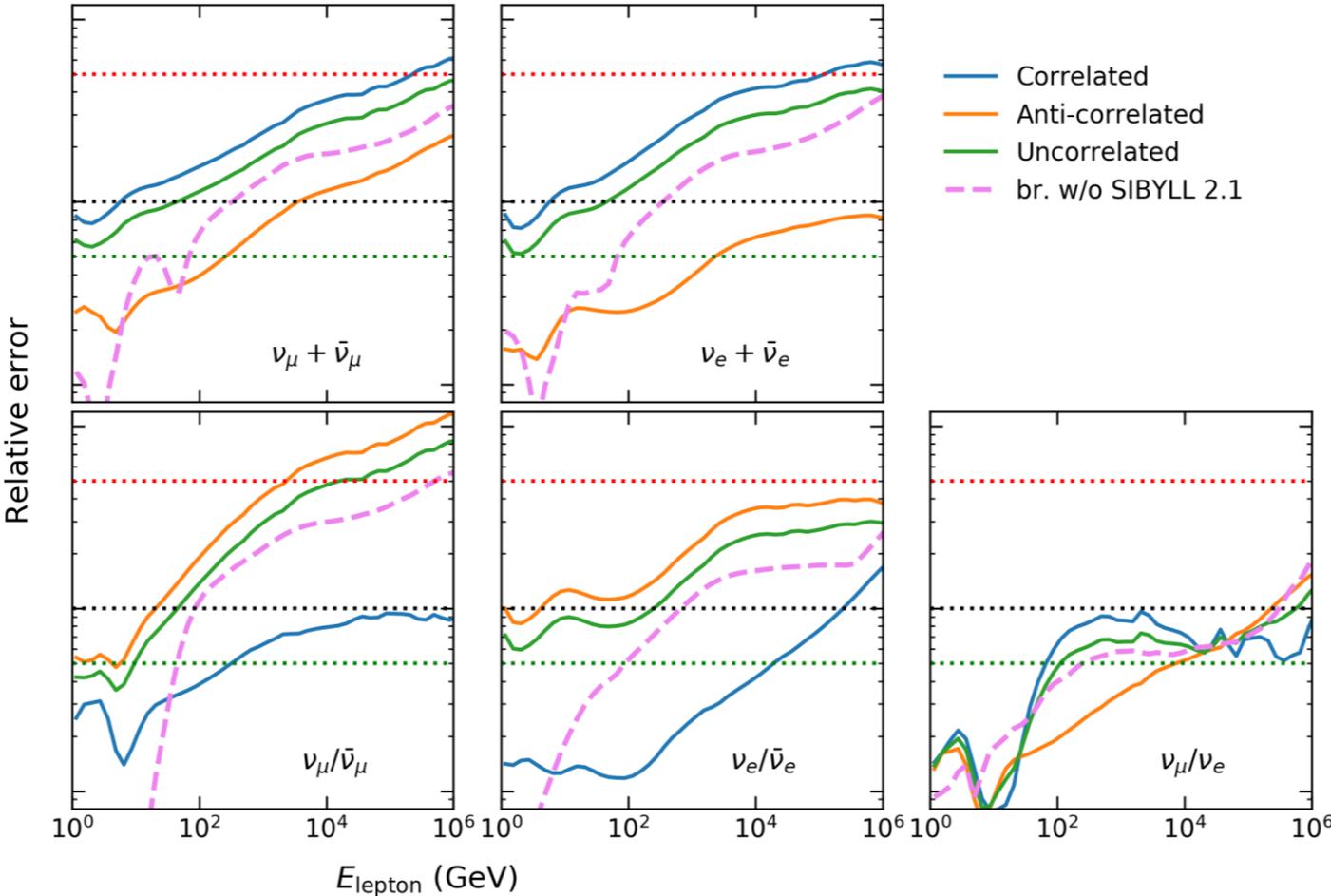
LHCb could significantly increase coverage



Contribution of individual “Barr groups”



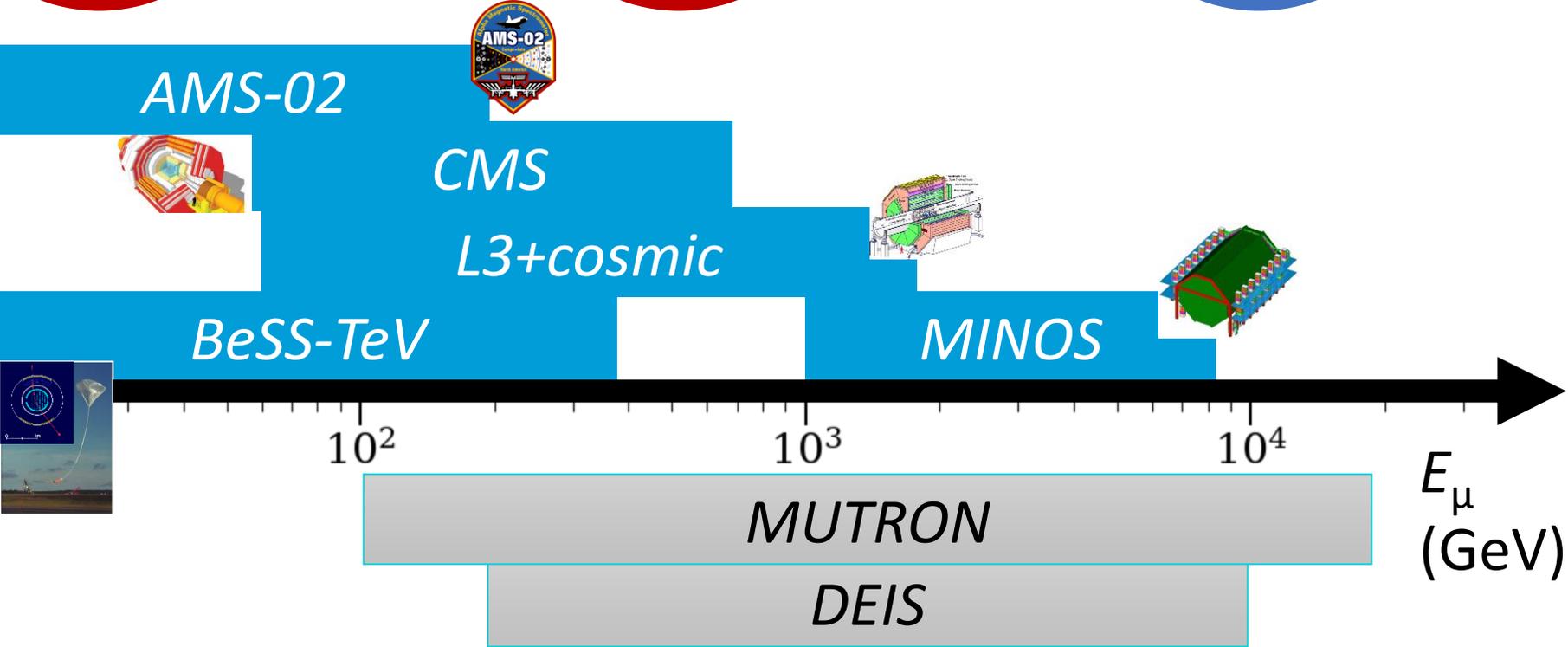
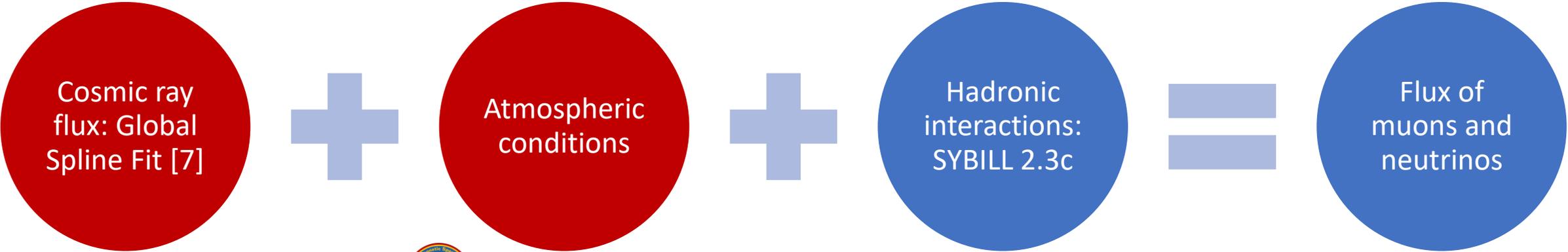
Correlations between phase-space patches unclear



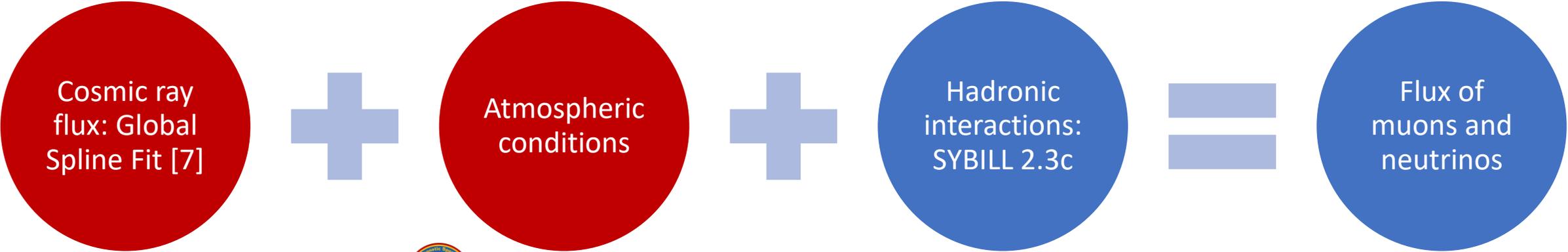
Examples	For one “Barr” - parameters
symmetric	$\rho\pi^+\uparrow \ n\pi^+\uparrow \ \rho\pi^-\uparrow \ n\pi^-\uparrow$
asymmetric	$\rho\pi^+\uparrow \ n\pi^+\uparrow \ \rho\pi^-\downarrow \ n\pi^-\downarrow$
uncorrelated	$\rho\pi^+\uparrow \ n\pi^+0 \ \rho\pi^-0 \ n\pi^-0$

- The production of charged secondaries is physically not independent
- It is very difficult to extract this information from hadronic interaction models directly

Calibration of ν uncertainties with “global fit” to μ data



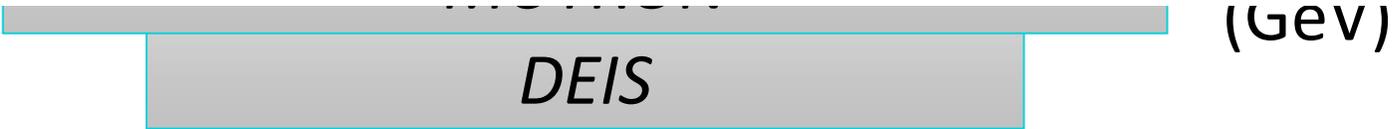
Calibration of ν uncertainties with “global fit” to μ data



AMS-02



Experiment	Energy (GeV)	Measurements	Reported unit	Location	Altitude	Zenith range
AMS-02	0.1-2500	Flux & charge ratio	rigidity	28.57°N , 80.65° W	5 m (sea level)	
BESS-TeV	0.6-400	Flux	momentum	36.2°N, 140.1°W	30 m	0-25.8°
CMS	5-1000	Charge ratio	momentum	46.31°N, 6.071°E	420 m	$p \cos \theta_z$
L3+C	20-3000	Flux & charge ratio	momentum	46.25°N, 6.02°E	450 m	0-58°
MINOS	1000-7000	Charge ratio	total energy	47.82°N, 92.24°W	5 m (sea level)	unfolded
OPERA	891-7079	Charge ratio	total energy	42.42°N, 13.51°E	5 m (sea level)	$E \cos \theta^*$



How we did it

- New version the cascade code MCEq with improved accuracy at low E
- Cut secondary particle phase-space according to parameters B_i from Barr et al.
- Generate database of fluxes $\Phi(E_\mu)$ and Jacobians

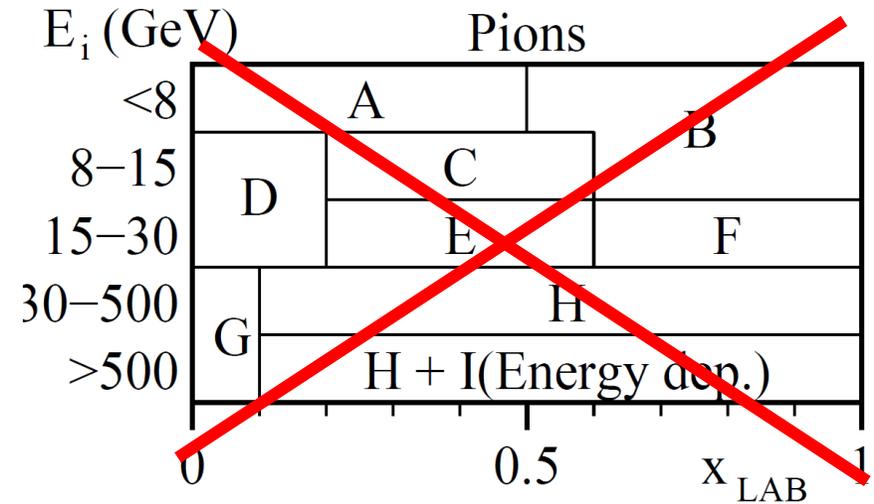
$$\frac{\partial \Phi(E_\mu)}{\partial \mathcal{B}_i} = \frac{\Phi(E_\mu, \mathcal{B}_i = 1 + \delta) - \Phi(E_\mu, \mathcal{B}_i = 1 - \delta)}{2\delta}$$

- Fluxes with modifications to B_i can be quickly evaluated in the fit:

$$\Phi(E_\mu, \mathcal{B}_a, \mathcal{B}_b, \dots) = \Phi(E_\mu) + \sum_i \mathcal{B}_i \frac{\partial \Phi(E_\mu)}{\partial \mathcal{B}_i}$$

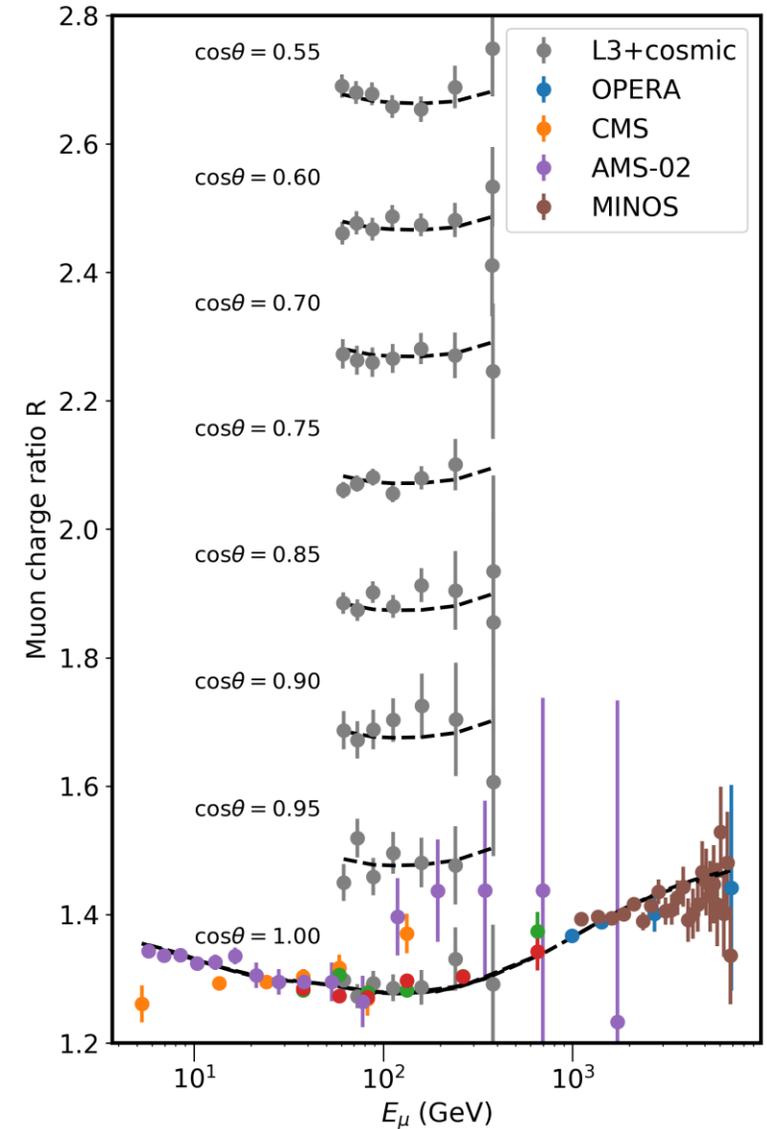
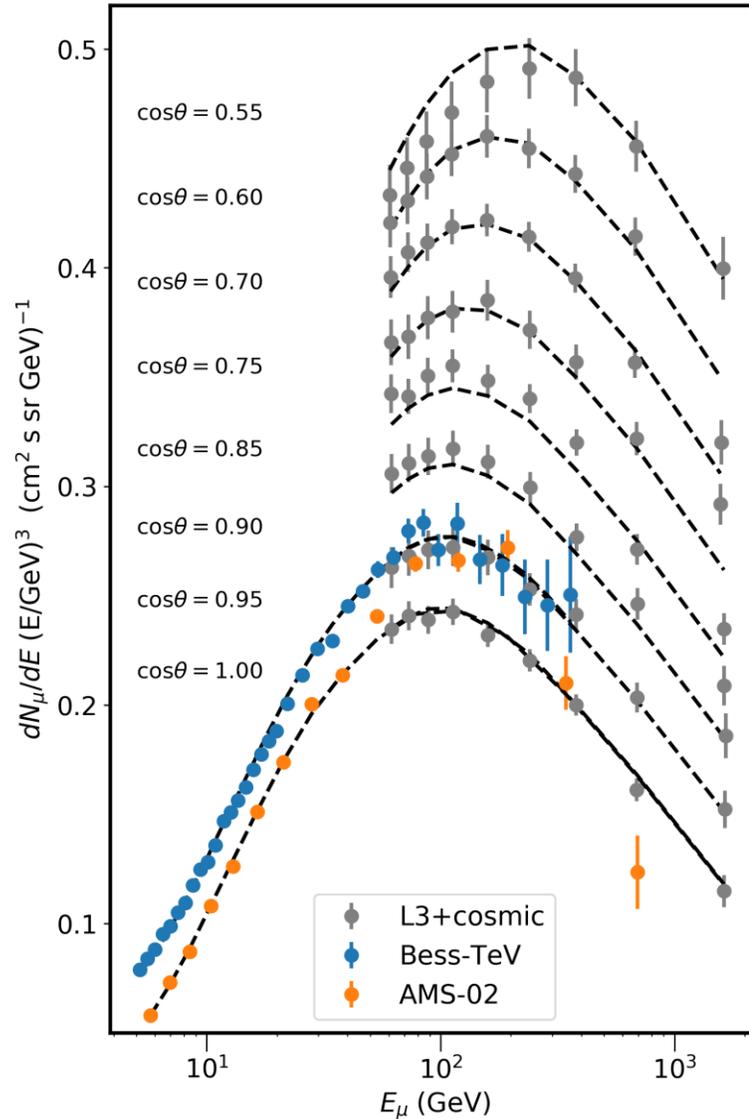
What we found

- Original attempt was to use the parameterization of Barr et al.
 1. Found data to be insensitive
 2. Too many correlations
 3. Impossible to constrain
- Simplified to four parameters
 - Yields of each meson species
 - Global, energy-independent scales
 - Enough to describe data



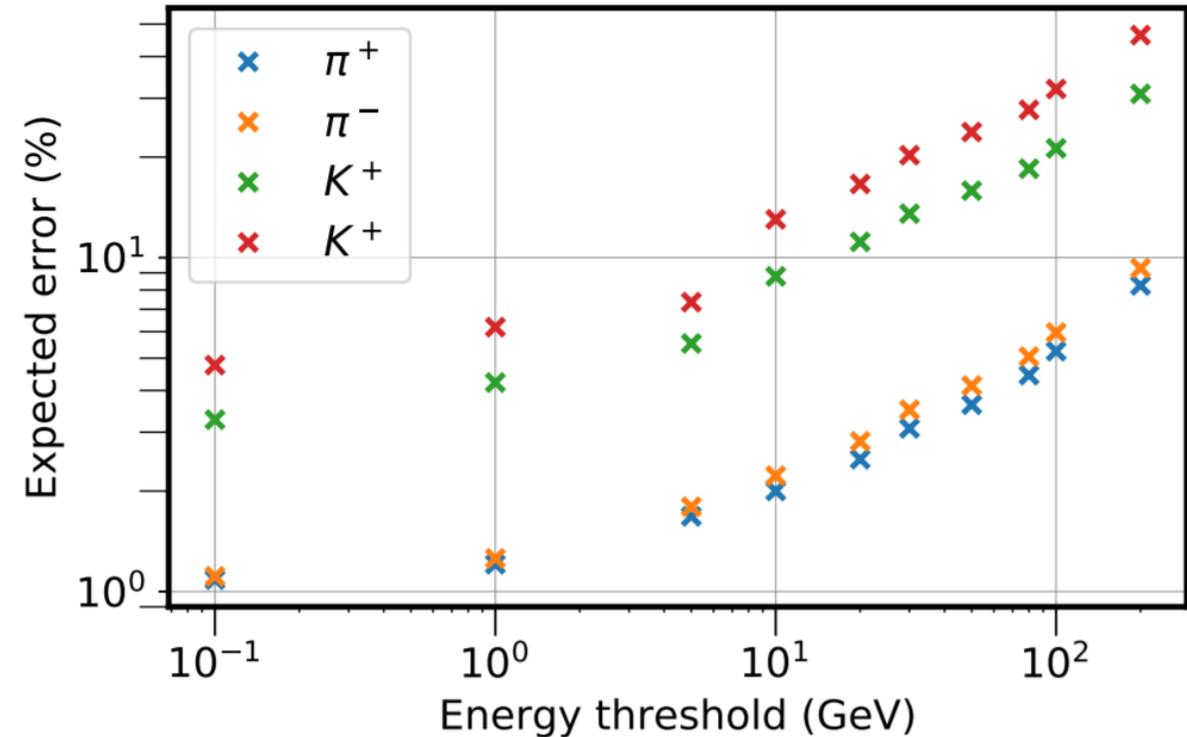
Fit results

- Some experiments are hard to fit regardless of modifications
- Possible systematic effects not reported
- L3+C – previously “the reference dataset” – is not as good as we thought
- We will include more data and CR flux uncertainties in the next iteration and report later this year
- **This approach limited to < few TeV ☹ lepton energy**



Impact of energy threshold for the fit

- High energy data less sensitive
- This is because the features in the muon spectrum are smooth
- and fit variables become strongly correlated
- More angles are needed
- We're investigating horizontal and high-altitude balloon data

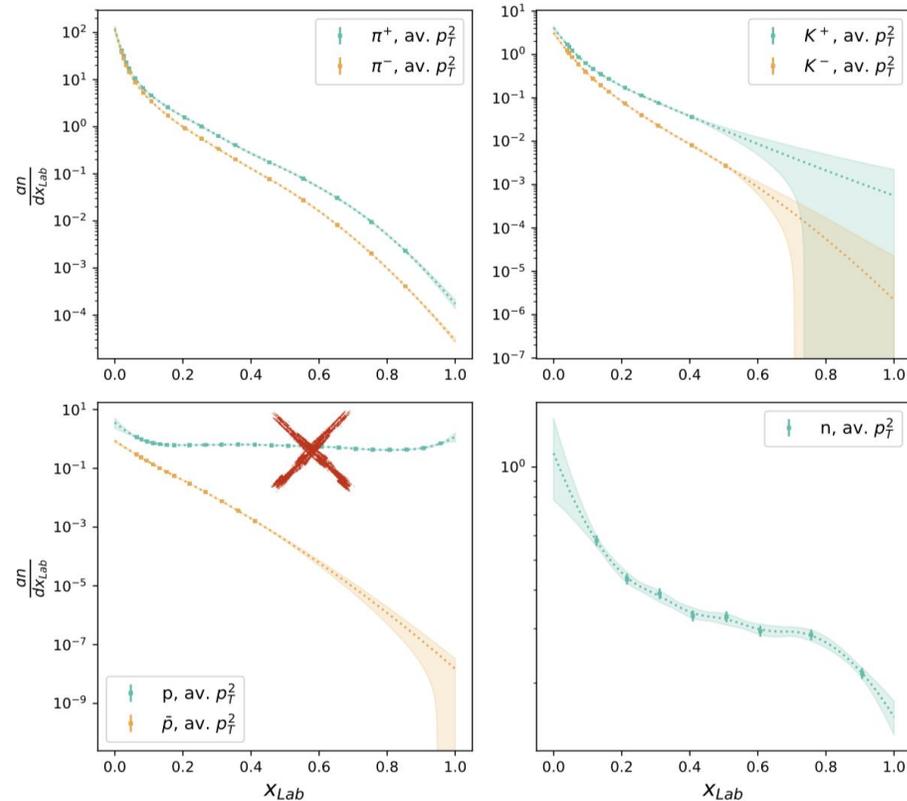


Alternative under investigation: data-driven inclusive interaction model

"The SHIn-project"

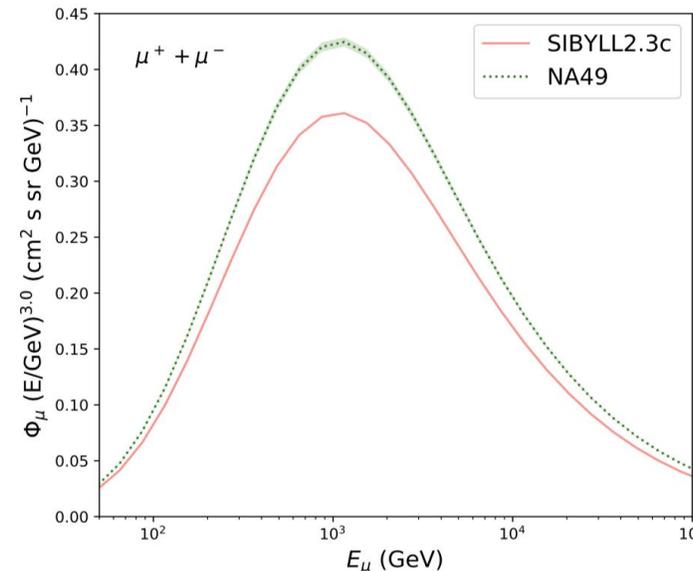


NA49 pp data (158GeV)



$p+p \rightarrow \text{part}+X$

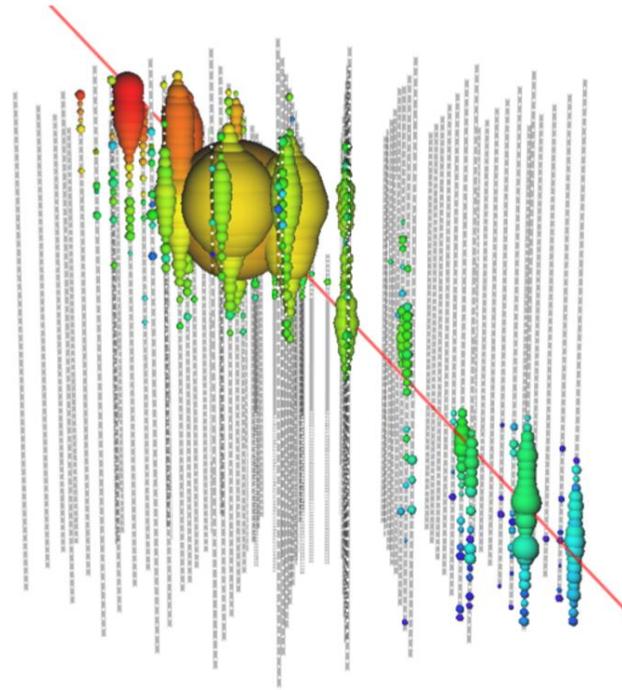
Experiment	Interaction	E_p [GeV]	yields
NA49	pp	158	$\pi^\pm, K^\pm, \bar{p}, n$
NA49	pC	158	π^\pm, \bar{p}, n
NA61/SHINE	pC	31	$\pi^\pm, K^\pm, K_S^0, \Lambda$
NA61/SHINE	pp	20, 31, 40, 80, 158	π^\pm, K^\pm, \bar{p}
NA61/SHINE	π^-C	158, 350	ρ^0, ω, K^{*0}
NA61/SHINE (upcoming)	π^-C	158, 350	π^\pm, K^\pm, \bar{p}



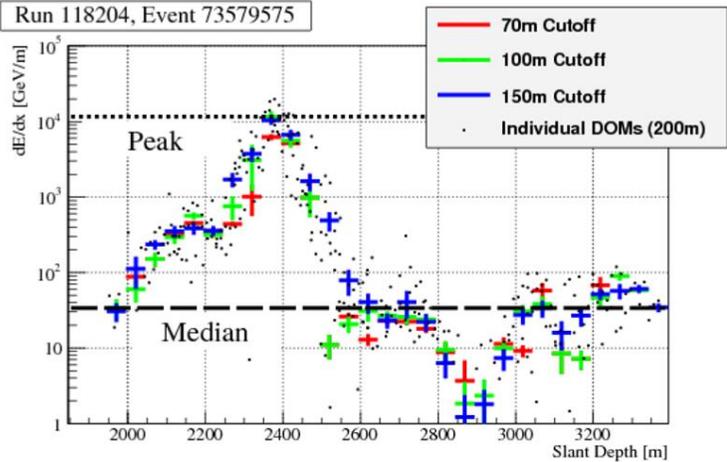
Matthias Huber

High-energy muons in IceCube

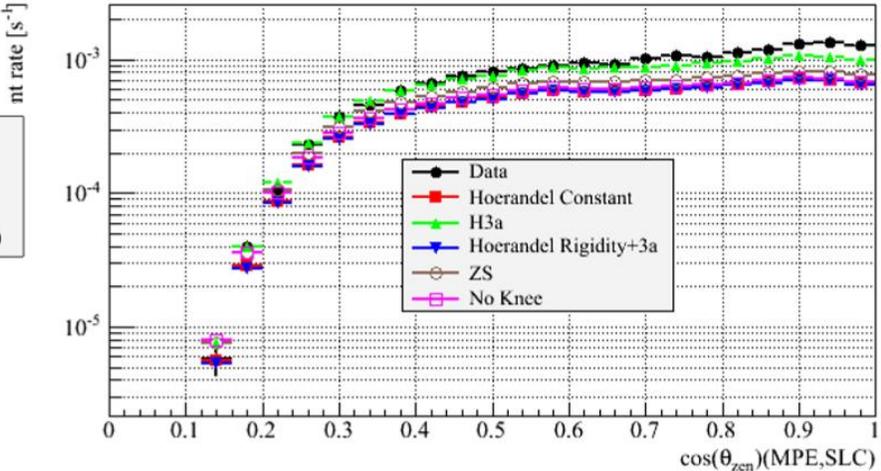
See talk by D. Soldin at UHECR 2018



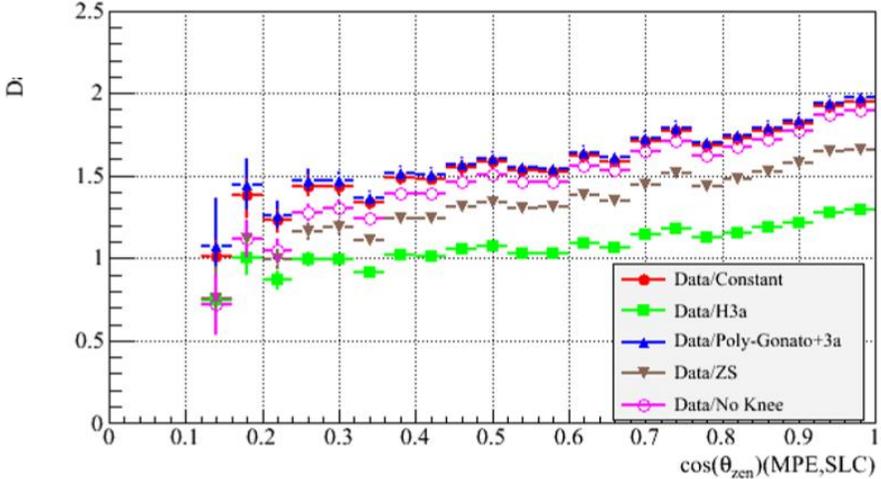
Energy loss pattern



Angular distribution has tilt

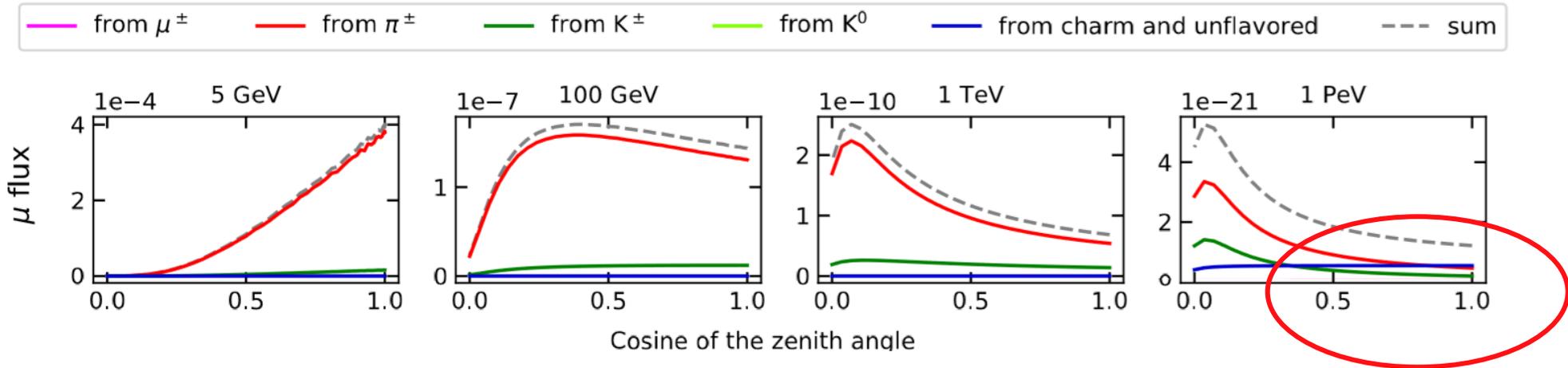


- Energy loss pattern sensitive to muon multiplicity and energy
- Angular distribution of muon bundles always tilted compared to MC
- Neither hadronic interactions nor cosmic ray models can explain this
- Maybe unknown detector systematics

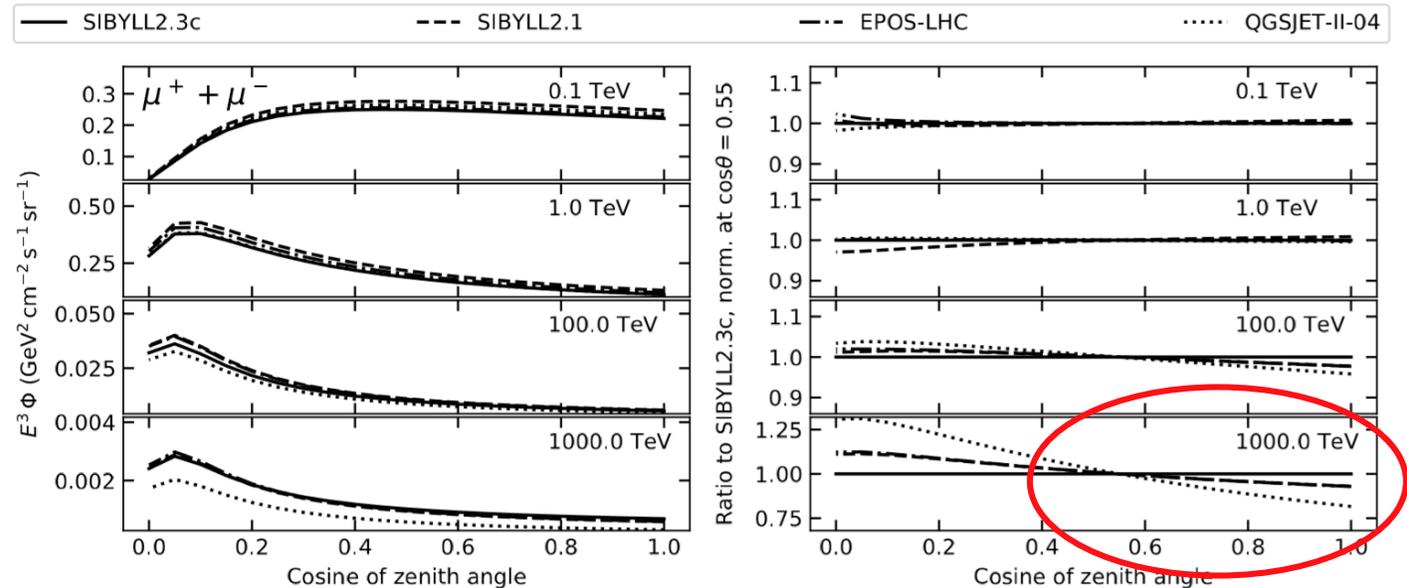


IceCube, Astropart. Phys. 78 (2016)

Possible origin from prompt muons (not necessarily charm)



- First **indication** that zenith shift is a signature of **prompt muons**
- But! **Other processes** contribute
 - Muon pair production $\gamma \rightarrow \mu^+ \mu^-$
 - Quasi-elastic vector meson production $\gamma + A \rightarrow (\rho, \omega, \phi) + A' \rightarrow \text{hadrons} + X$
- The “electromagnetic” processes might be under better control and thus would allow for constraining the contribution of charm (50%)
- **Tilt in the zenith distribution is simple to measure**



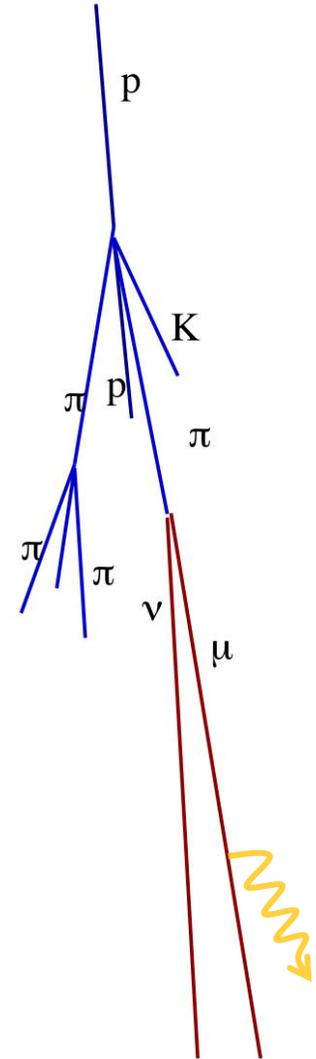
Conclusions and future path

- Current atmospheric neutrino detectors cover **9 orders of magnitude** in energy (MeV-PeV) → **challenge** for modeling!
- **High-precision** (and high-performance) scheme to take **uncertainties self-consistently** into account **is available** through MCEq (<https://github.com/afedynitch/MCEq>), but good/realistic parameterizations are under investigation
- **Data-driven techniques** can constrain some of the expected/known uncertainties, but is this a consistent scheme? What do we learn?
- Measurement of the prompt flux unlikely to happen in a model-independent way for the present generation of detectors
- Help creating a better likelihood for the (global) fit
- Build KM3Net

Transport equations (hadronic cascade equations)

System of coupled non-linear PDE for each particle species h :

$$\begin{aligned} \frac{d\Phi_h(E, X)}{dX} = & - \frac{\Phi_h(E, X)}{\lambda_{\text{int},h}(E)} && \text{Interactions with air} \\ & - \frac{\Phi_h(E, X)}{\lambda_{\text{dec},h}(E, X)} && \text{Decays} \\ & - \frac{\partial}{\partial E} (\mu(E)\Phi_h(E, X)) && \text{Continuous losses} \\ & + \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{int},k}(E_k)} && \text{Re-injection from interactions} \\ & + \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}^{\text{dec}}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{dec},k}(E_k, X)} && \text{Re-injection from decays} \end{aligned}$$



$$X(h_0) = \int_0^{h_0} d\ell \rho_{\text{air}}(\ell)$$

Transport equations (hadronic cascade equations)

System of coupled non-linear PDE for each particle species h :

$$\frac{d\Phi_h(E, X)}{dX} = - \frac{\Phi_h(E, X)}{\lambda_{\text{int},h}(E)} \quad \text{cosmic ray physics}$$

$$- \frac{\Phi_h(E, X)}{\lambda_{\text{dec},h}(E, X)} \quad \text{Interactions with air}$$

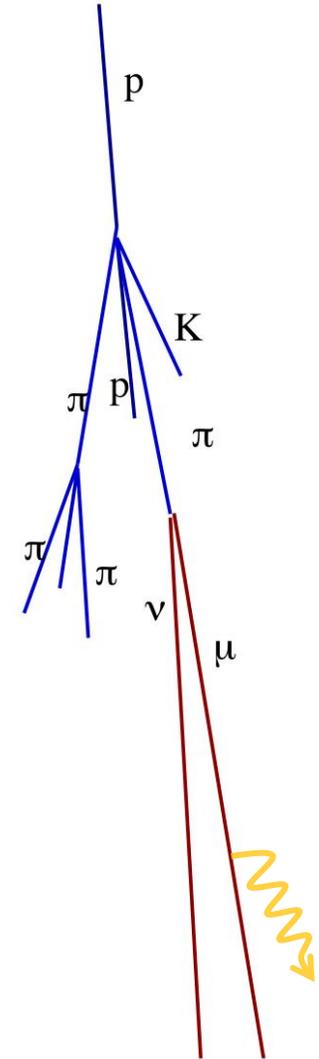
$$- \frac{\partial}{\partial E} (\mu(E) \Phi_h(E, X)) \quad \text{Decays}$$

$$+ \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{int},k}(E_k)} \quad \text{Continuous losses}$$

$$+ \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}^{\text{dec}}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{dec},k}(E_k, X)} \quad \text{particle physics}$$

Re-injection from interactions

Re-injection from decays



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cosmic ray physics

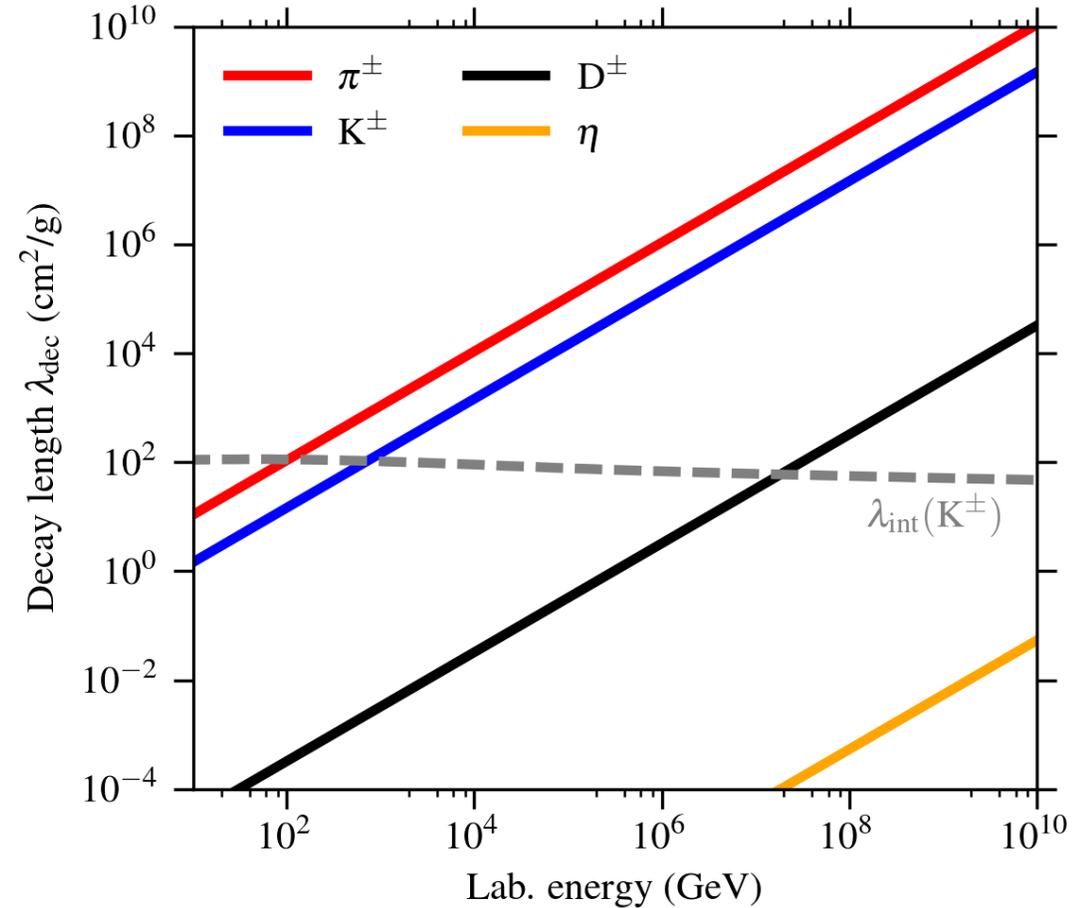
Interactions with air

Decays

atmospheric physics

Continuous losses

particle physics



$$X(h_0) = \int_0^{h_0} dl \rho_{\text{air}}(l)$$

Transport equations (hadronic cascade equations)

System of coupled non-linear PDE for each particle species h :

$$\frac{d\Phi_h(E, X)}{dX} = - \frac{\Phi_h(E, X)}{\lambda_{\text{int},h}(E)} - \frac{\Phi_h(E, X)}{\lambda_{\text{dec},h}(E, X)} - \frac{\partial}{\partial E} (\mu(E)\Phi_h(E, X)) + \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{int},k}(E_k)} + \sum_k \int_E^\infty dE_k \frac{dN_{k(E_k) \rightarrow h(E)}^{\text{dec}}}{dE} \frac{\Phi_k(E_k, X)}{\lambda_{\text{dec},k}(E_k, X)}$$

cosmic ray physics

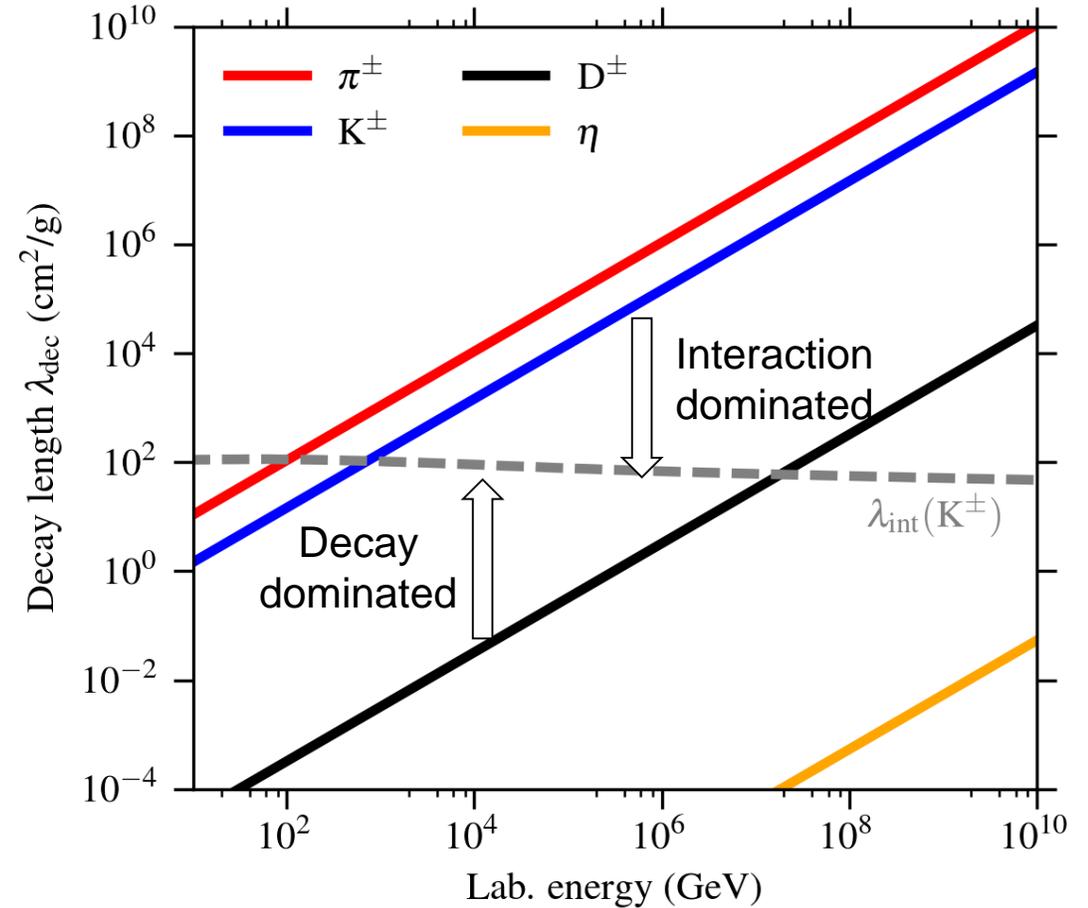
Interactions with air

Decays

atmospheric physics

Continuous losses

particle physics



$$X(h_0) = \int_0^{h_0} dl \rho_{\text{air}}(l)$$

MCEq: Matrix Cascade Equations

$$\begin{aligned} \frac{d\Phi_h(E, X)}{dX} = & - \frac{\Phi_h(E, X)}{\lambda_{\text{int},h}(E)} \\ & - \frac{\Phi_h(E, X)}{\lambda_{\text{dec},h}(E, X)} \\ & - \frac{\partial}{\partial E}(\mu(E)\Phi_h(E, X)) \\ & + \sum_{\ell} \int_E^{\infty} dE_{\ell} \frac{dN_{\ell(E_{\ell}) \rightarrow h(E)}}{dE} \frac{\Phi_{\ell}(E_{\ell}, X)}{\lambda_{\text{int},\ell}(E_{\ell})} \\ & + \sum_{\ell} \int_E^{\infty} dE_{\ell} \frac{dN_{\ell(E_{\ell}) \rightarrow h(E)}^{\text{dec}}}{dE} \frac{\Phi_{\ell}(E_{\ell}, X)}{\lambda_{\text{dec},\ell}(E_{\ell}, X)} \end{aligned}$$



$$\begin{aligned} \frac{d\Phi_{E_i}^h}{dX} = & - \frac{\Phi_{E_i}^h}{\lambda_{\text{int},E_i}^h} \\ & - \frac{\Phi_{E_i}^h}{\lambda_{\text{dec},E_i}^h(X)} \\ & - \vec{\nabla}_i(\mu_{E_i}^h \Phi_{E_i}^h) \\ & + \sum_{E_k \geq E_i}^{E_N} \sum_{\ell} \frac{C_{\ell(E_k) \rightarrow h(E_i)}}{\lambda_{\text{int},E_k}^{\ell}} \Phi_{E_k}^{\ell} \\ & + \sum_{E_k \geq E_i}^{E_N} \sum_{\ell} \frac{d_{\ell(E_k) \rightarrow h(E_i)}}{\lambda_{\text{dec},E_k}^{\ell}(X)} \Phi_{E_k}^{\ell} \end{aligned}$$

State (or flux) vector

$$\vec{\Phi} = \left(\vec{\Phi}^{\text{p}} \quad \vec{\Phi}^{\text{n}} \quad \vec{\Phi}^{\pi^+} \quad \dots \quad \vec{\Phi}^{\bar{\nu}_{\mu}} \quad \dots \right)^T$$

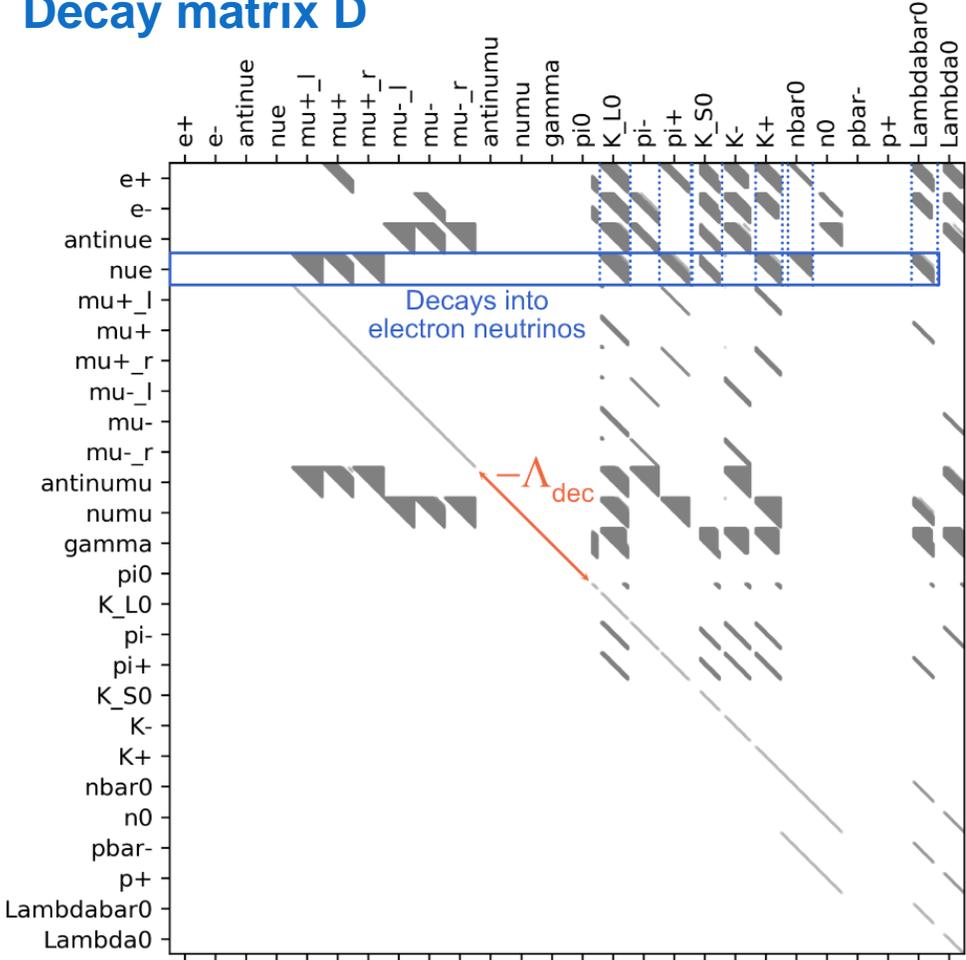
$$\vec{\Phi}^{\text{p}} = \left(\Phi_{E_0}^{\text{p}} \quad \Phi_{E_1}^{\text{p}} \quad \dots \quad \Phi_{E_N}^{\text{p}} \right)^T$$

“Matrix form”

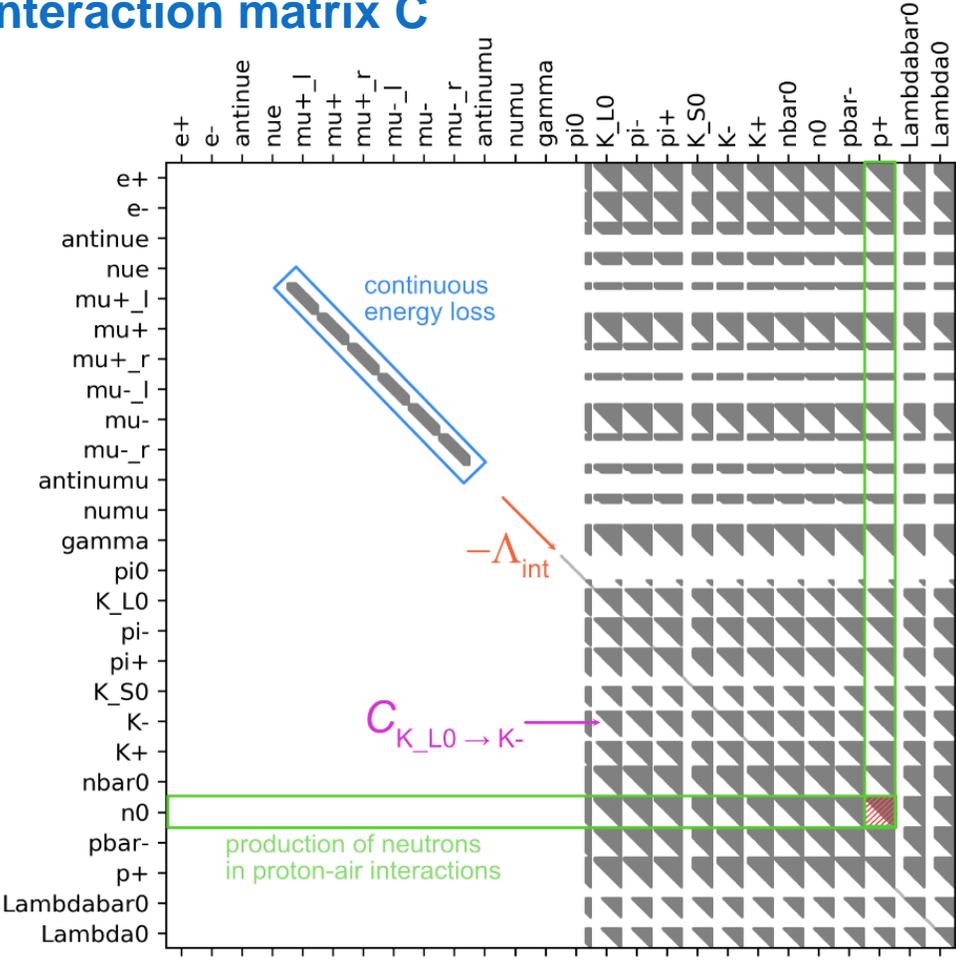
$$\begin{aligned} \frac{d}{dX} \vec{\Phi} = & - \vec{\nabla}_E (\text{diag}(\vec{\mu}) \vec{\Phi}) + (-\mathbf{1} + \mathbf{C}) \Lambda_{\text{int}} \vec{\Phi} \\ & + \frac{1}{\rho(X)} (-\mathbf{1} + \mathbf{D}) \Lambda_{\text{dec}} \vec{\Phi} \end{aligned}$$

Sparse matrix structure

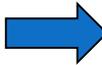
Decay matrix D



Interaction matrix C



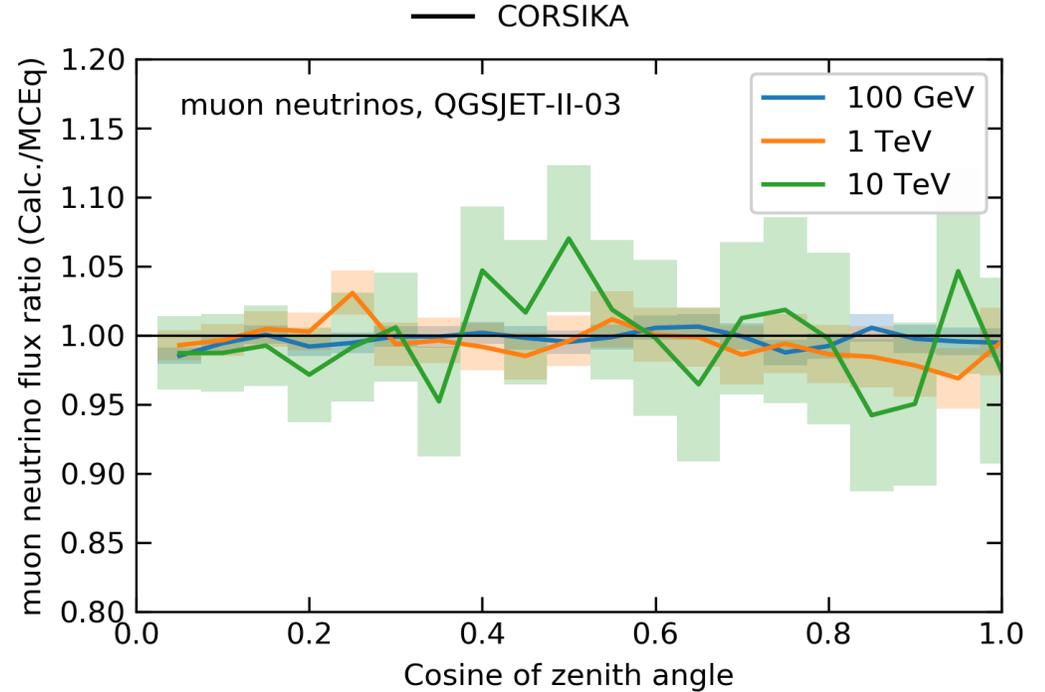
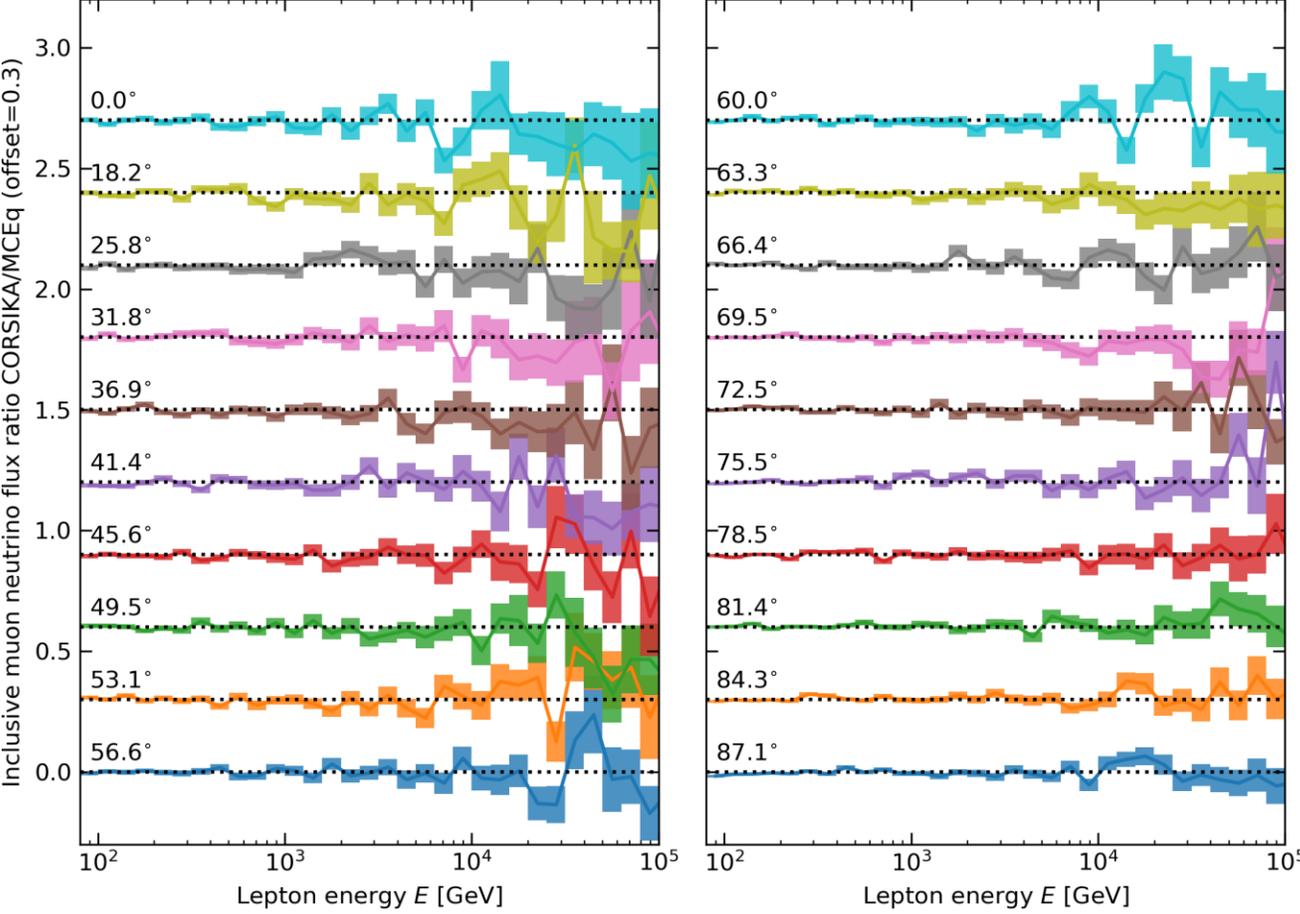
matrices are sparse



high performance

MCEq vs (thinned) CORSIKA calculation in 1D

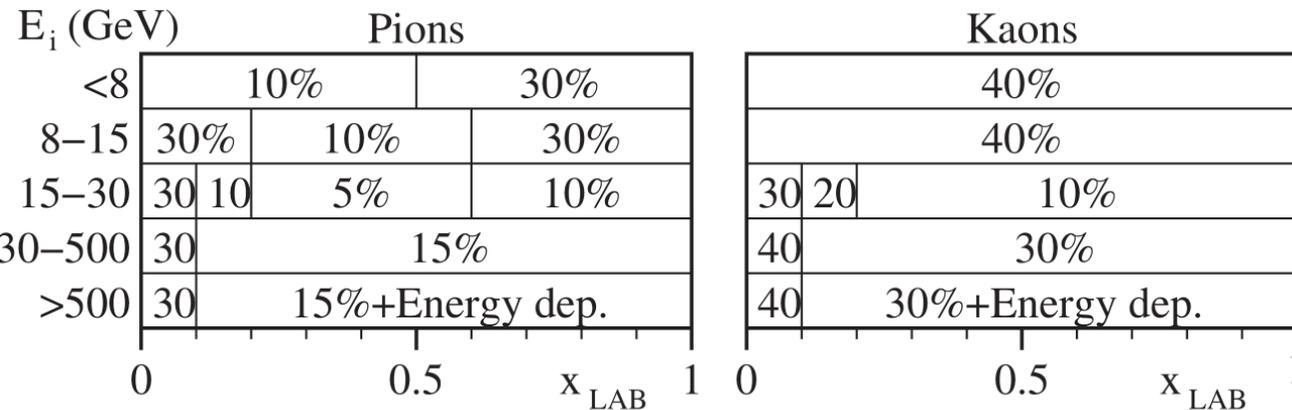
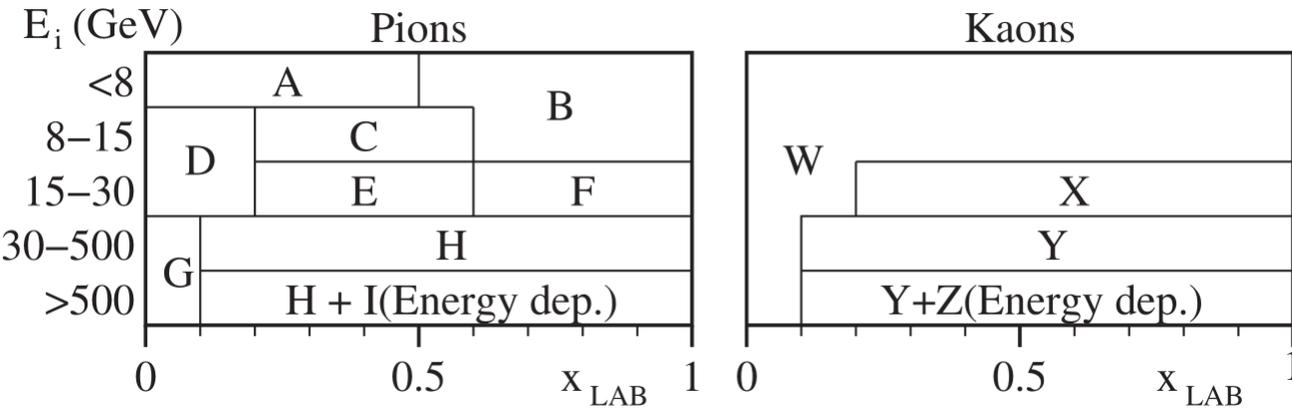
Inclusive muon neutrino flux ratio CORSIKA/MCEQ. QGSJET-II-03 + H3a.



> BSD licensed @ <https://github.com/afedynitch/MCEq>

MCEq-based implementation

“Barr regions”



- Compute partial derivatives wrt. phase-space regions (Taylor expansion), i.e. $\frac{\partial \Phi_\nu}{\partial W}$
- No correlations between phase-space regions (as in Barr et al.) or add. correlations

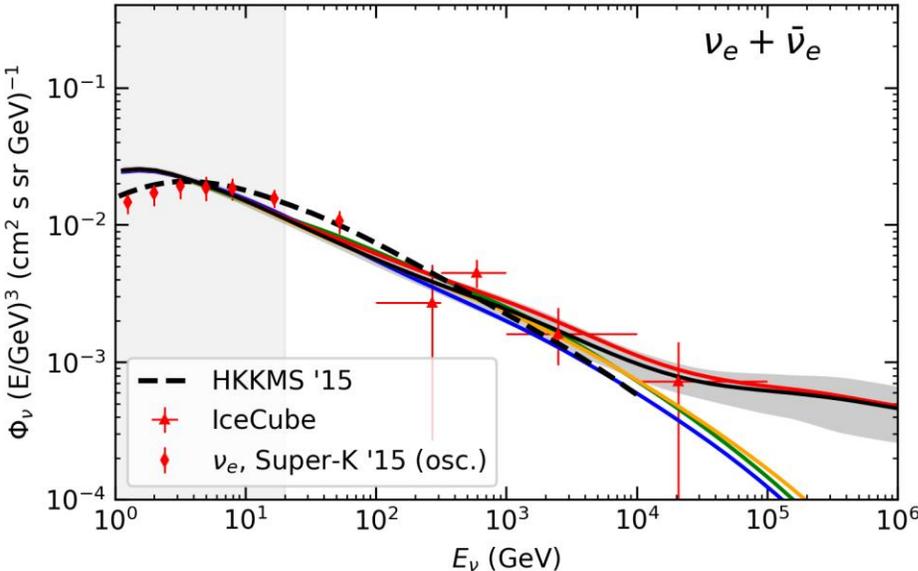
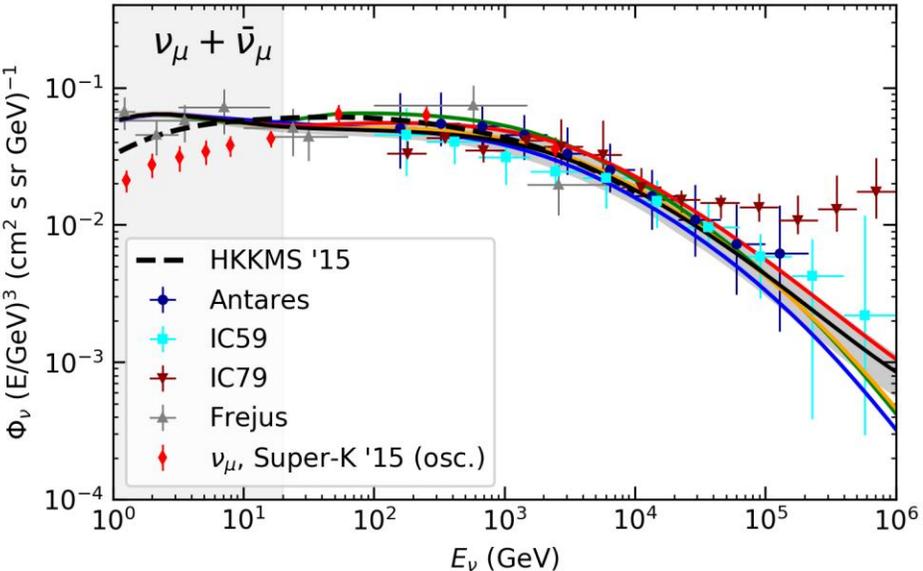
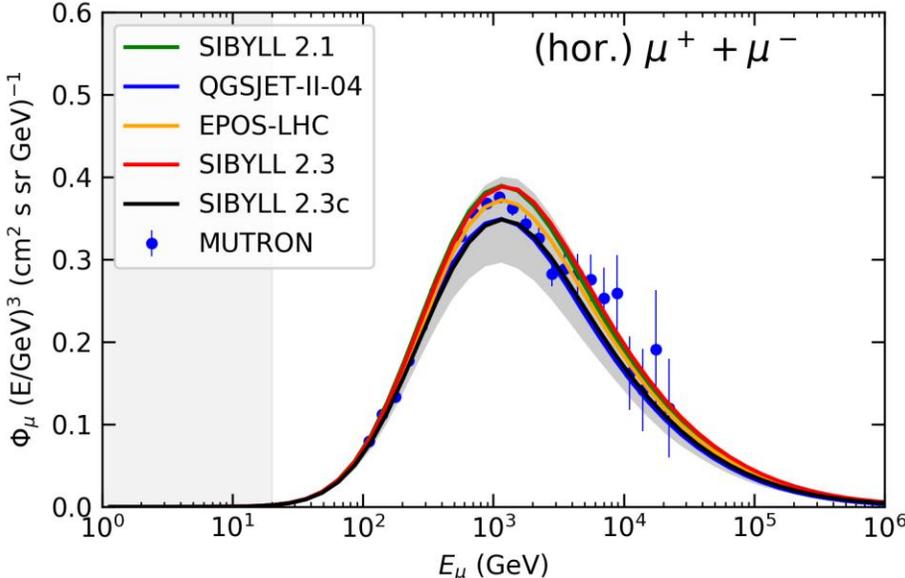
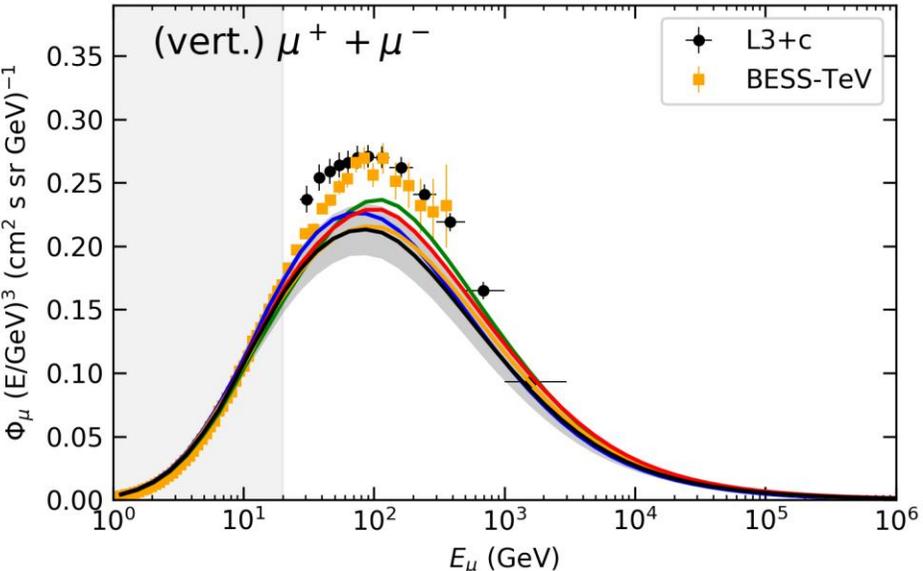
Elements of Jacobian (numerical)

$$J_{E_{ij}} = \frac{\partial \Phi_\nu(E_i)}{\partial p} = \frac{\Phi_\nu(\delta p_j+) - \Phi_\nu(\delta p_j-)}{2\delta p_j}$$

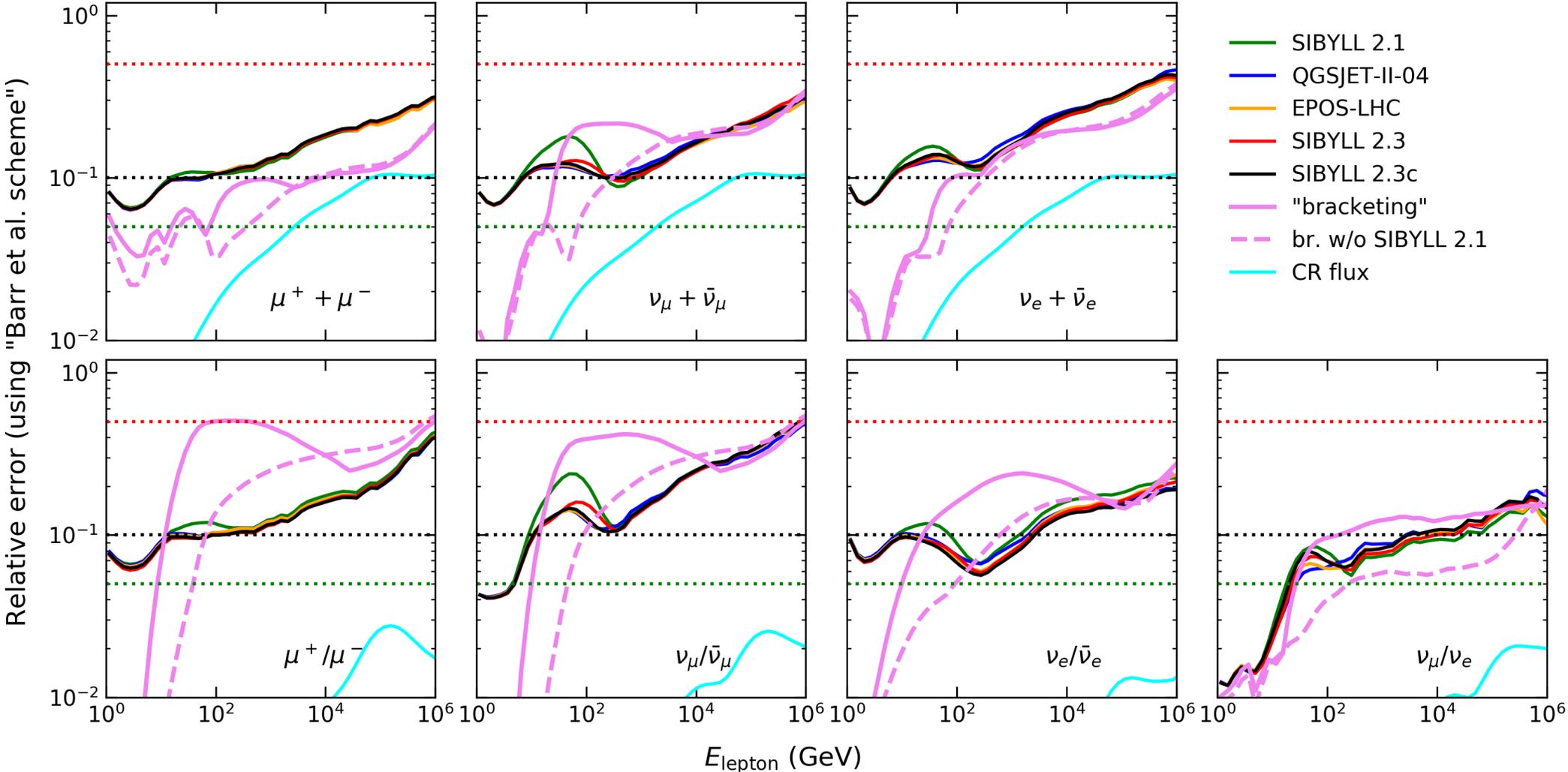
Error propagation

$$\text{cov}[\Phi_\nu(E_i), \Phi_\nu(E_j)] = \sum_{mn} J_{E_i m} J_{E_j n} \text{cov}[p_m, p_l]$$

Computation of error bands through error propagation

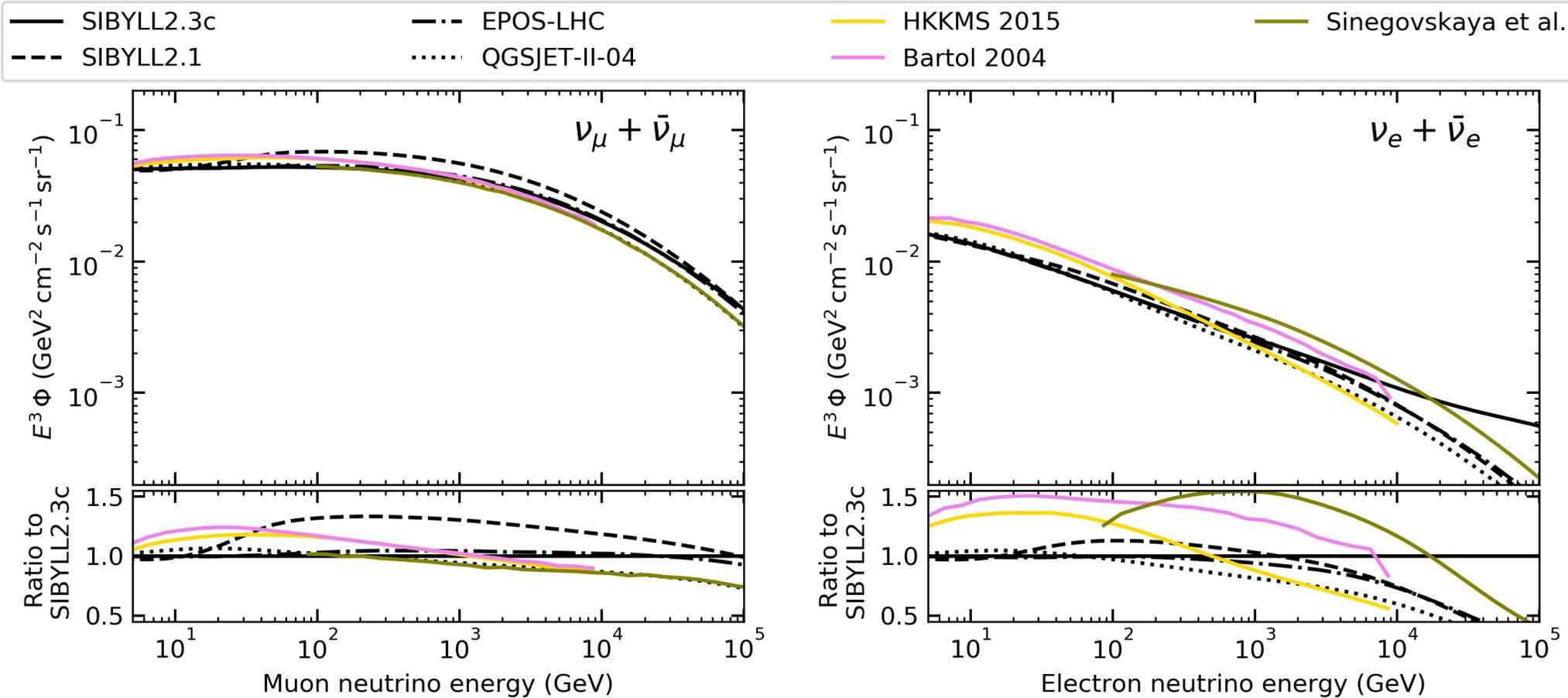


Computation of error bands through error propagation



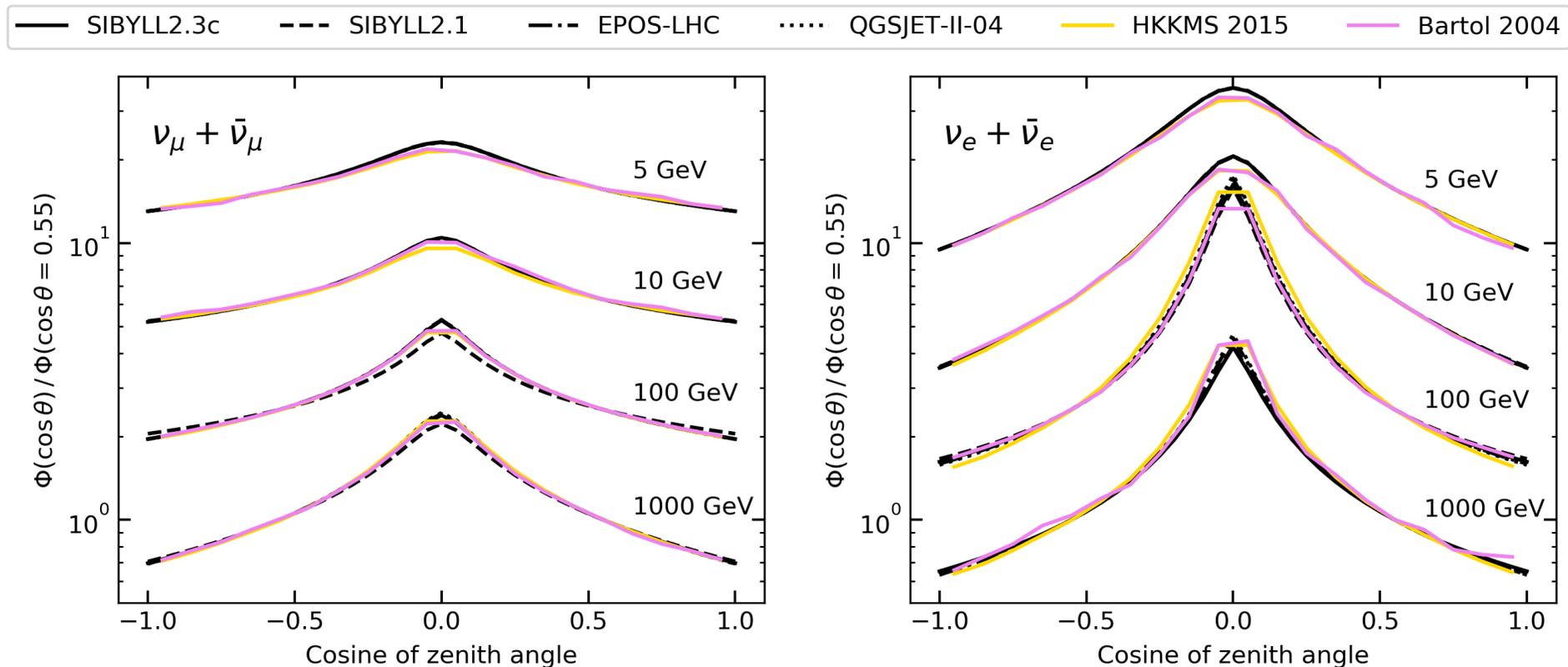
MCEq vs. traditional calculations

HKMS: M. Honda et al., PRD 92 (2015)
 Bartol: G. Barr et al., PRD 70 (2004)
 Sinogovskaya et al. PRD 91 (2015)
 MCEq: AF, R. Engel in prep.



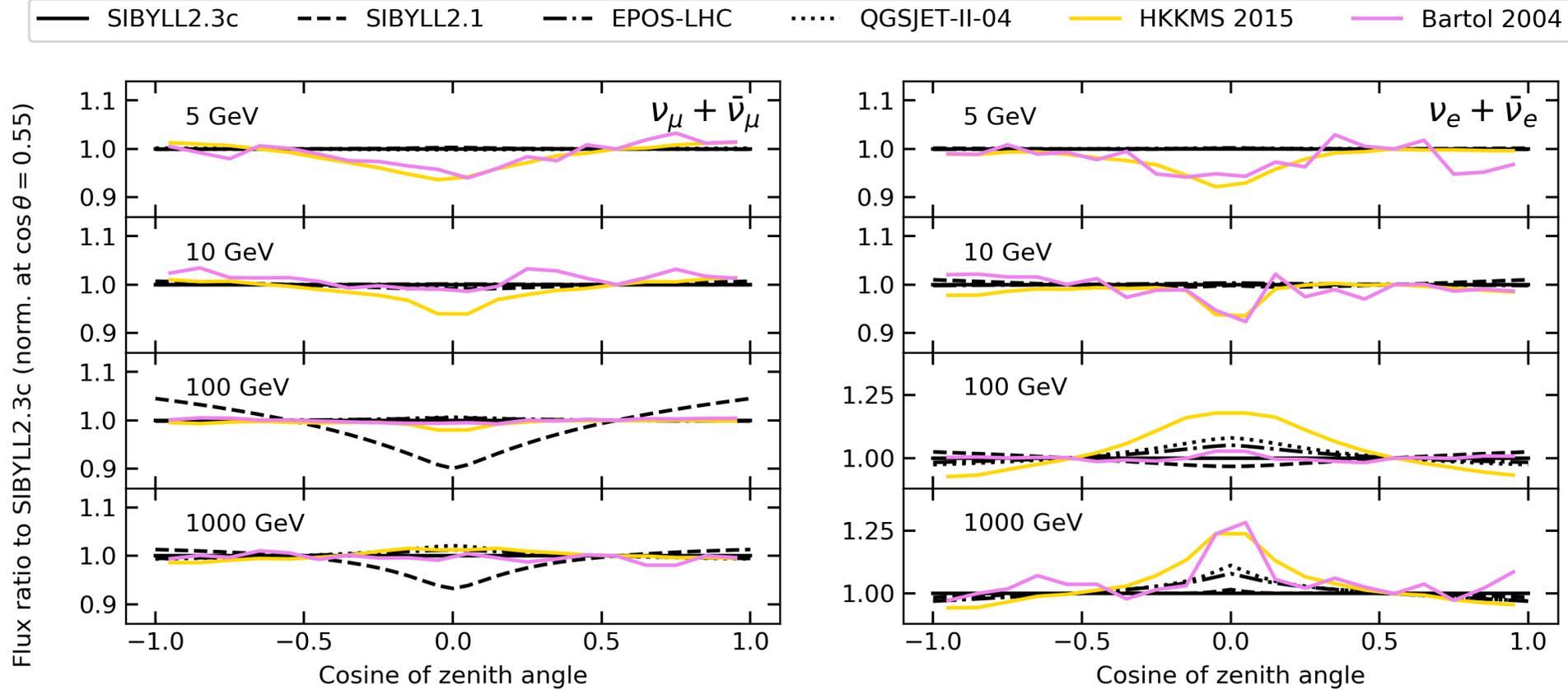
- Old 2002 (GH) primary model for HKMS and Bartol, H3a for the rest
- Data can not discriminate between calculations
- Shown are zenith and azimuth averages

Hadronic model dependence of zenith distributions



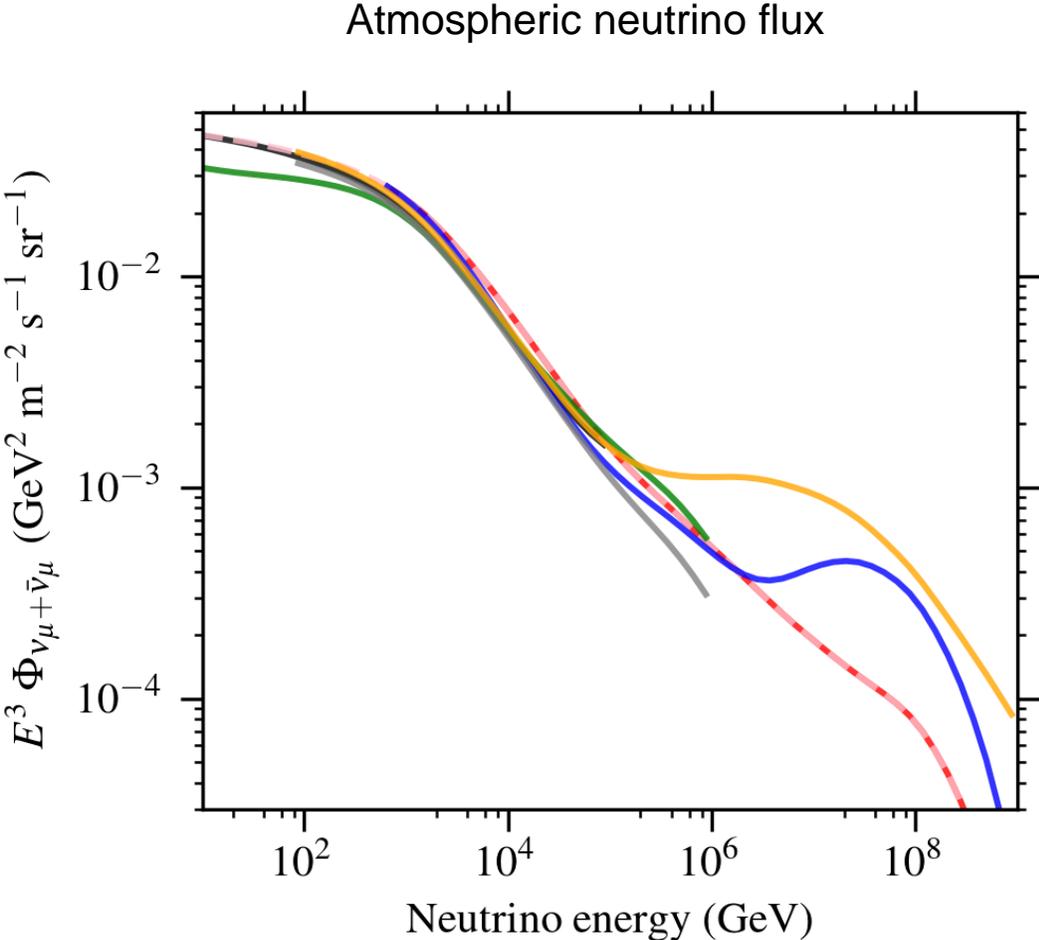
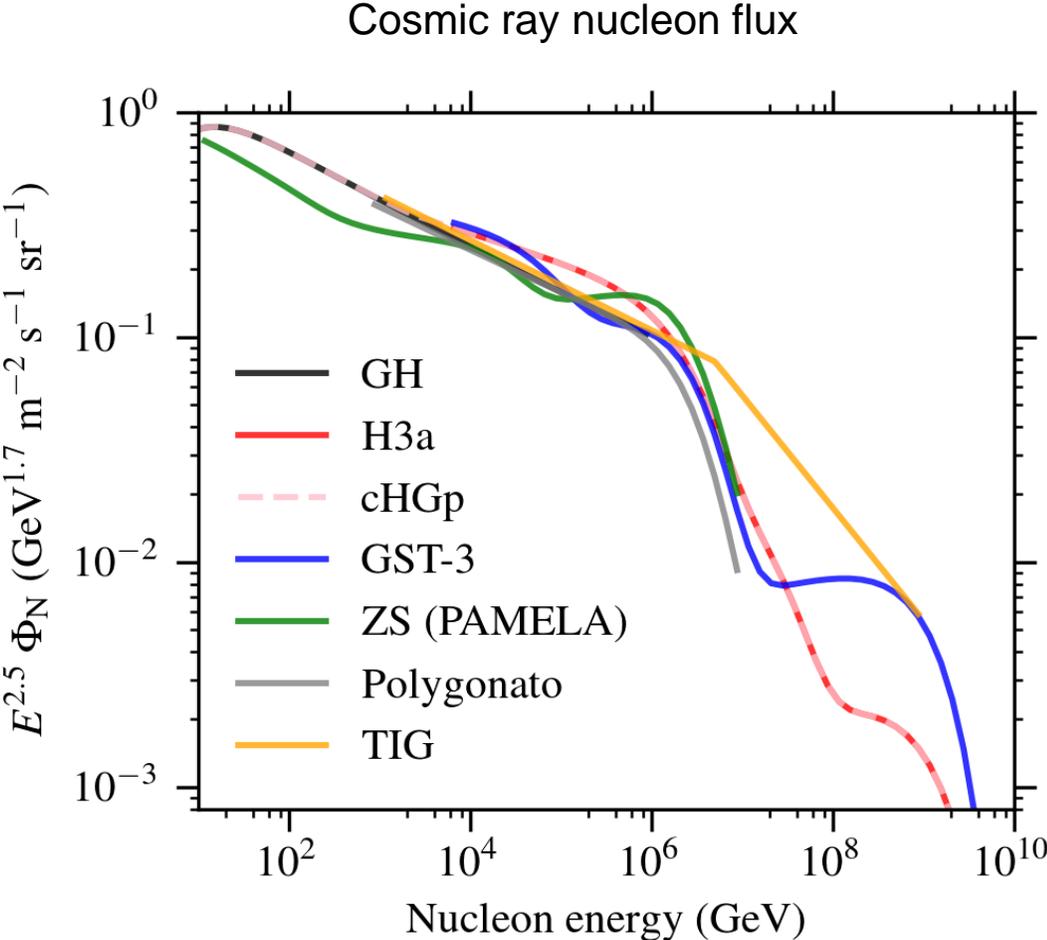
- Good agreement above tens of GeV for muon neutrinos
- Some tension between calculations at the horizon in electron neutrinos
- Affected by K/Pi, K^+/K^0_L ratios

Hadronic model dependence of zenith distributions



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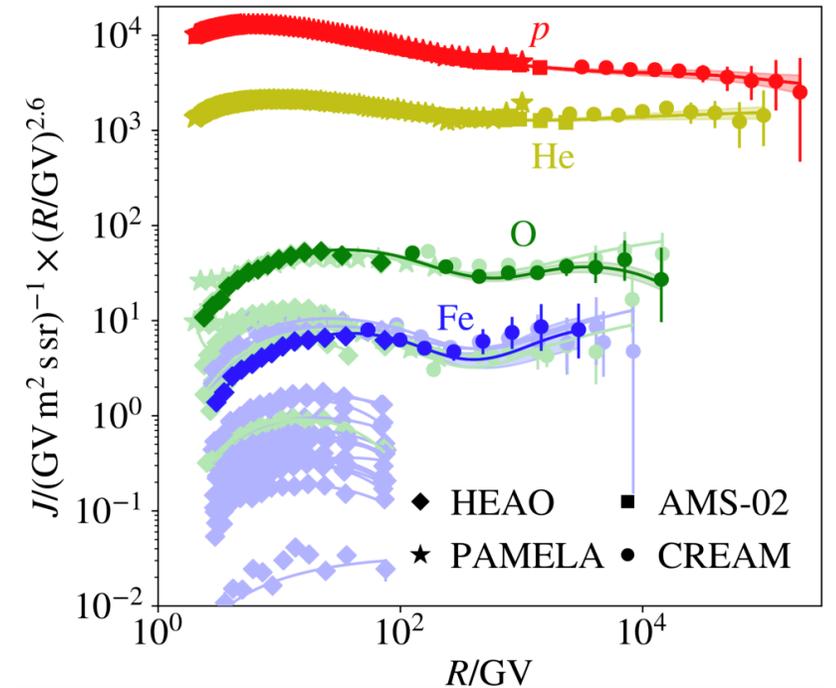
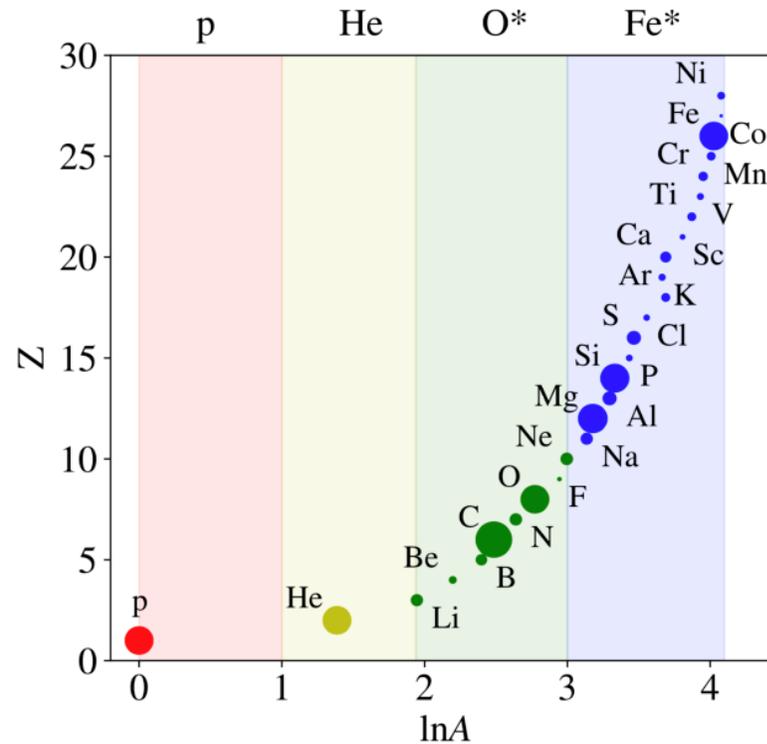
Cosmic ray flux uncertainties – ‘bracketing’ overestimates



Global Spline Fit – fit to direct & indirect observations

H. Dembinski, AF, T. Gaisser
PoS(ICRC2017)533

- Fit **four** independent mass groups, which cover equal ranges in $\ln A$:
proton (p), helium (He), oxygen group (O*), and iron group (Fe*)
- Assumption: this holds **at all energies**
- One leading element L per group described by smooth spline curve
- Other elements j in a group kept in constant ratio: $J_j(R)/J_L(R) = \text{const.}$

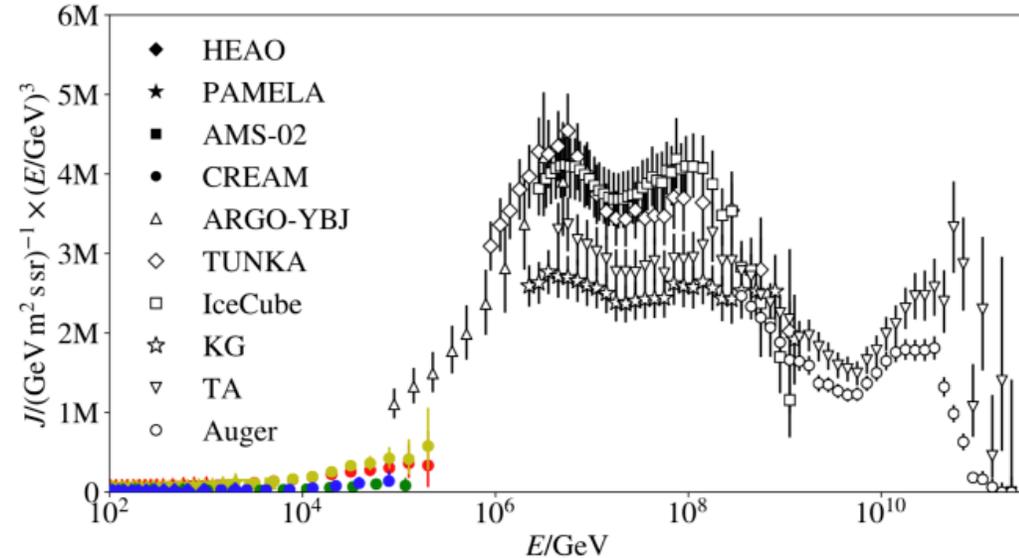
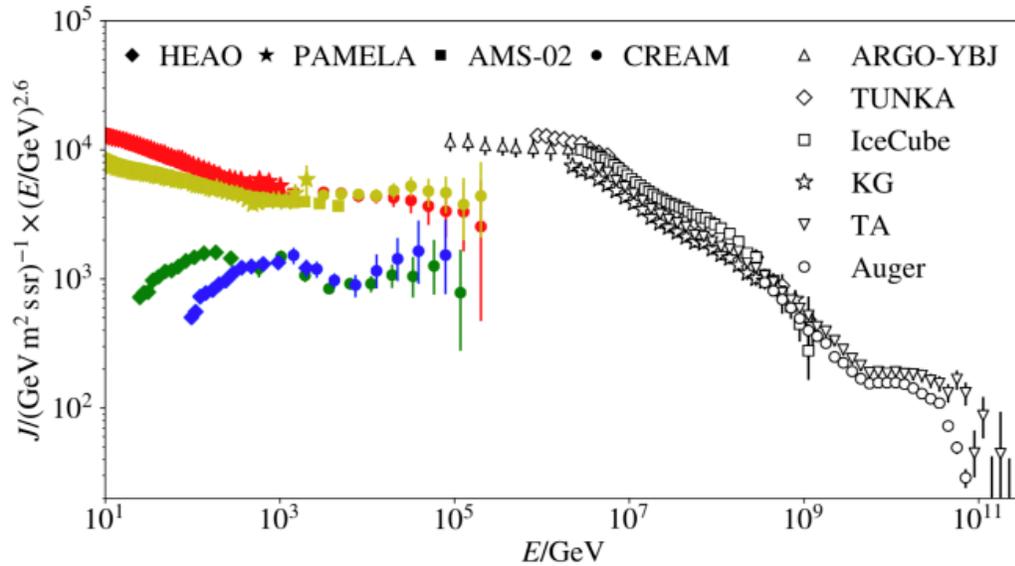


Mass sensitivity of air-shower experiments is $\sim \ln A$

Handling energy-scale uncertainty

H. Dembinski, AF, T. Gaisser
PoS(ICRC2017)533

Original data



- The determination of **energy scale in air-shower experiments is uncertain**
- This is caused by inconsistencies of **hadronic interaction models**
- Fit adjusts energy scales **within systematic uncertainties** of the experiment

$$\tilde{J}(\tilde{E}) = J(E) \frac{dE}{d\tilde{E}} = J \left(\frac{\tilde{E}}{1 + z_E} \right) \frac{1}{1 + z_E}$$

Flux distortion caused by energy-scale offset z_E

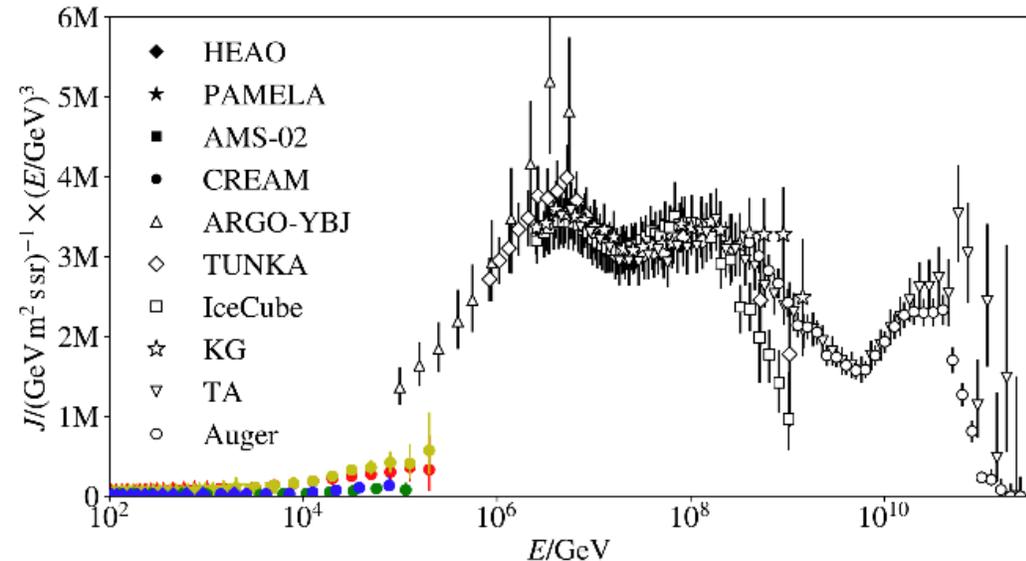
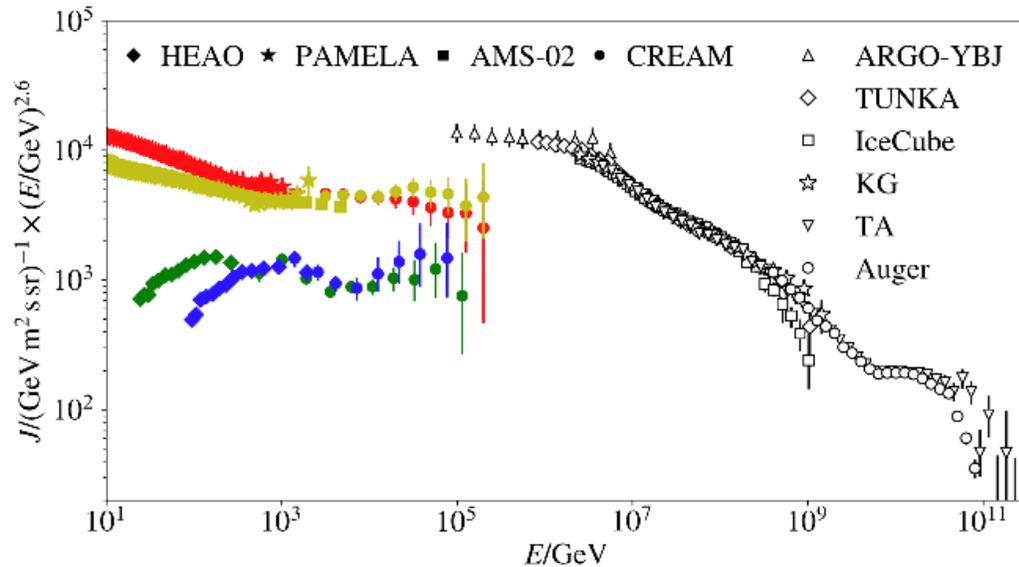
$$S = \sum_i z_i^2 + \sum_j \left(\frac{z_{Ej}}{(\sigma[E]/E)_j} \right)^2$$

Flux residuals Energy-scale offset residuals

Handling energy-scale uncertainty

H. Dembinski, AF, T. Gaisser
PoS(ICRC2017)533

Adjusted data



- The determination of **energy scale in air-shower experiments is uncertain**
- This is caused by inconsistencies of **hadronic interaction models**
- Fit adjusts energy scales **within systematic uncertainties** of the experiment

$$\tilde{J}(\tilde{E}) = J(E) \frac{dE}{d\tilde{E}} = J \left(\frac{\tilde{E}}{1 + z_E} \right) \frac{1}{1 + z_E}$$

Flux distortion caused by energy-scale offset z_E

$$S = \sum_i z_i^2 + \sum_j \left(\frac{z_{Ej}}{(\sigma[E]/E)_j} \right)^2$$

Flux residuals Energy-scale offset residuals