

Heavy-quark Production and Further Recent Developments in EPOS

Tanguy Pierog

Karlsruhe Institute of Technology, Institut für Kernphysik,
Karlsruhe, Germany

K. Werner, Iu. Karpenko, G. Sophys, Subatech, Nantes, France
B. Guiot, FSMT University, Valparaiso, Chile
M. Stefaniak, Warsaw University of Technology, Warsaw, Poland



**Heavy Quark Hadroproduction from Collider to
Astroparticle Physics, MITP, Mainz, Germany**

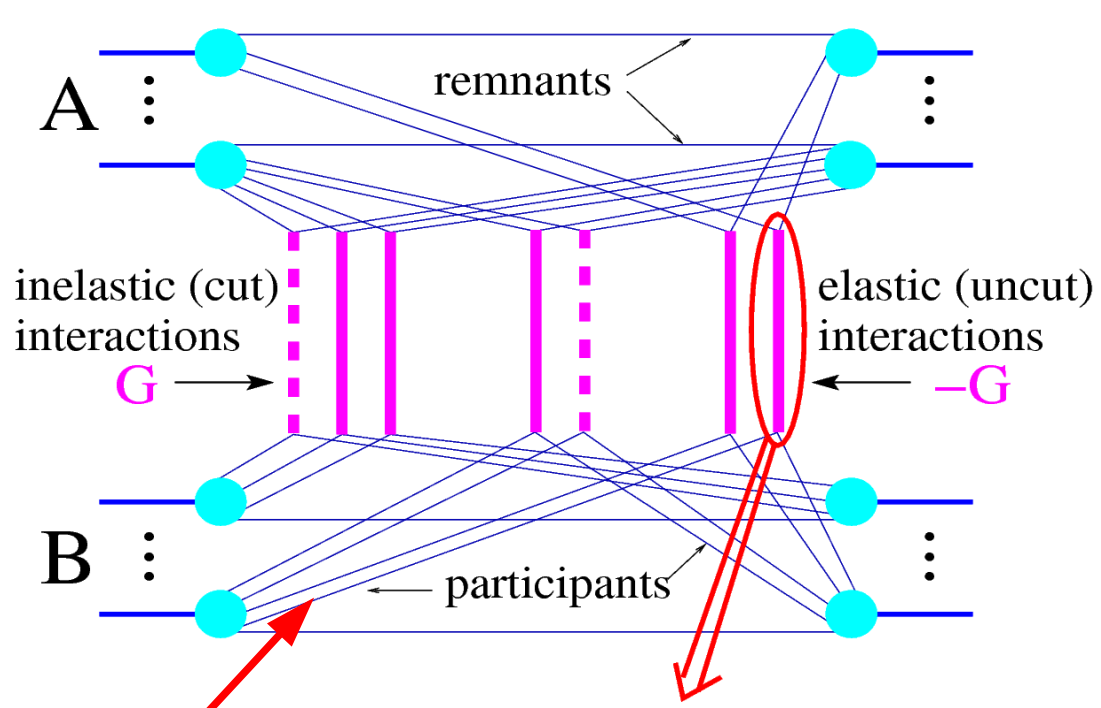
September the 30th 2019

Outline

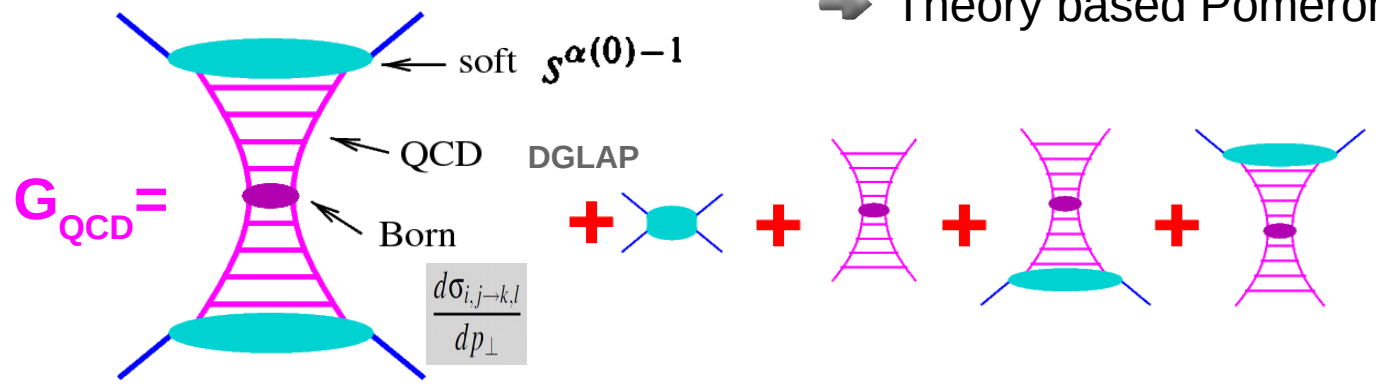
- **EPOS Basic principles**
 - ➔ For EPOS LHC and EPOS 3
- **Heavy Quark (HQ or Q) production**
 - ➔ Soft and perturbative production (EPOS 3)
- **Non-linear perturbative scale**
 - ➔ Evolution of particle production with multiplicity (EPOS 3)
- **Summary**

To reconcile soft and hard observables, EPOS 3 will have **both collective effects and variable non-linear perturbative scale : impact on HQ production.**

Parton-Based Gribov-Regge Theory



Pomeron momentum fraction x^- (x^+ for projectile)

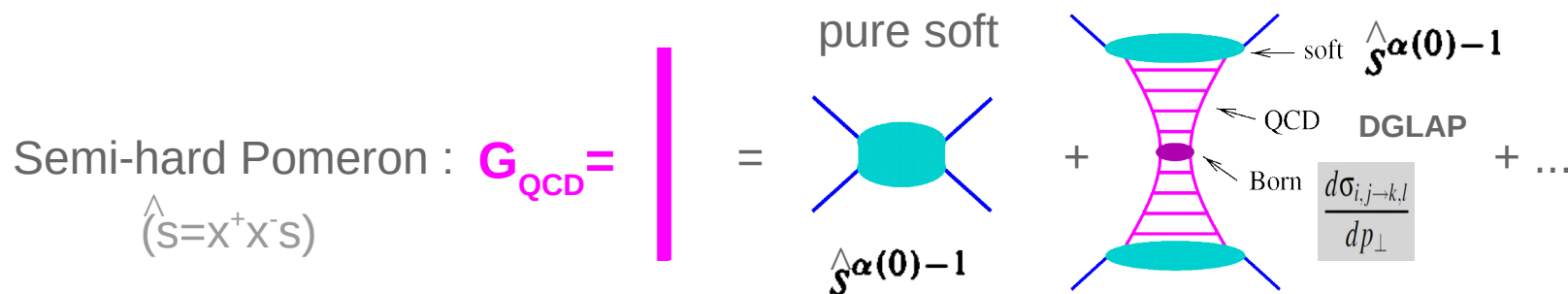


Energy sharing at the cross section level

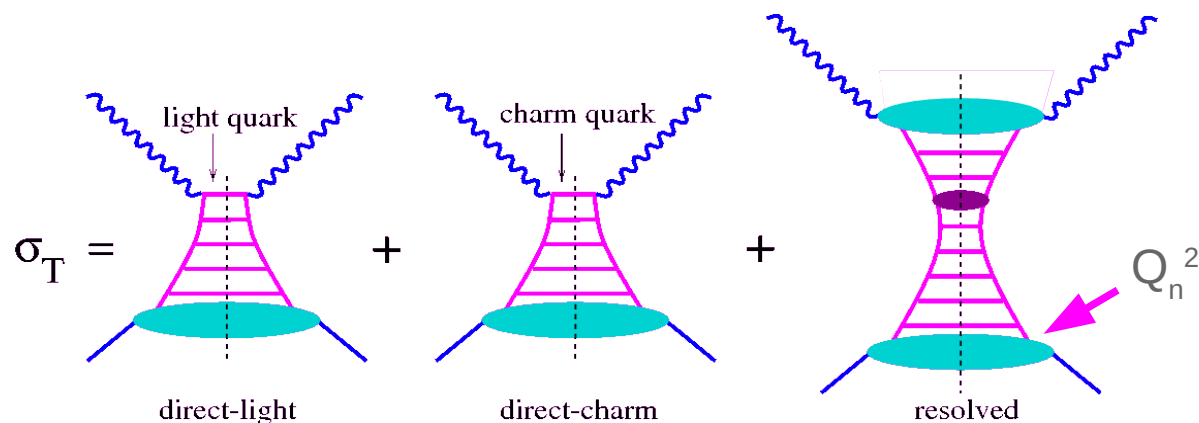
- ➔ Energy shared between cut and uncut diagrams (Pomeron)
- ➔ Reduced number of elementary interactions
- ➔ Generalization to (h)A-B
- ➔ Particle production from momentum fraction matrix (Markov chain metropolis)
- ➔ Theory based Pomeron definition

Parton-based Gribov-Regge Theory, H. J. Drescher, M. Hladik, S. Ostapchenko, T. Pierog, and K. Werner, Phys. Rept. 350 (2001) 93-289;

EPOS : Pomeron Definition



Test of semi-hard Pomeron with DIS: (Parton Distribution Function from HERA)



→ Theory based Pomeron definition

- pQCD based (DGLAP and Born)

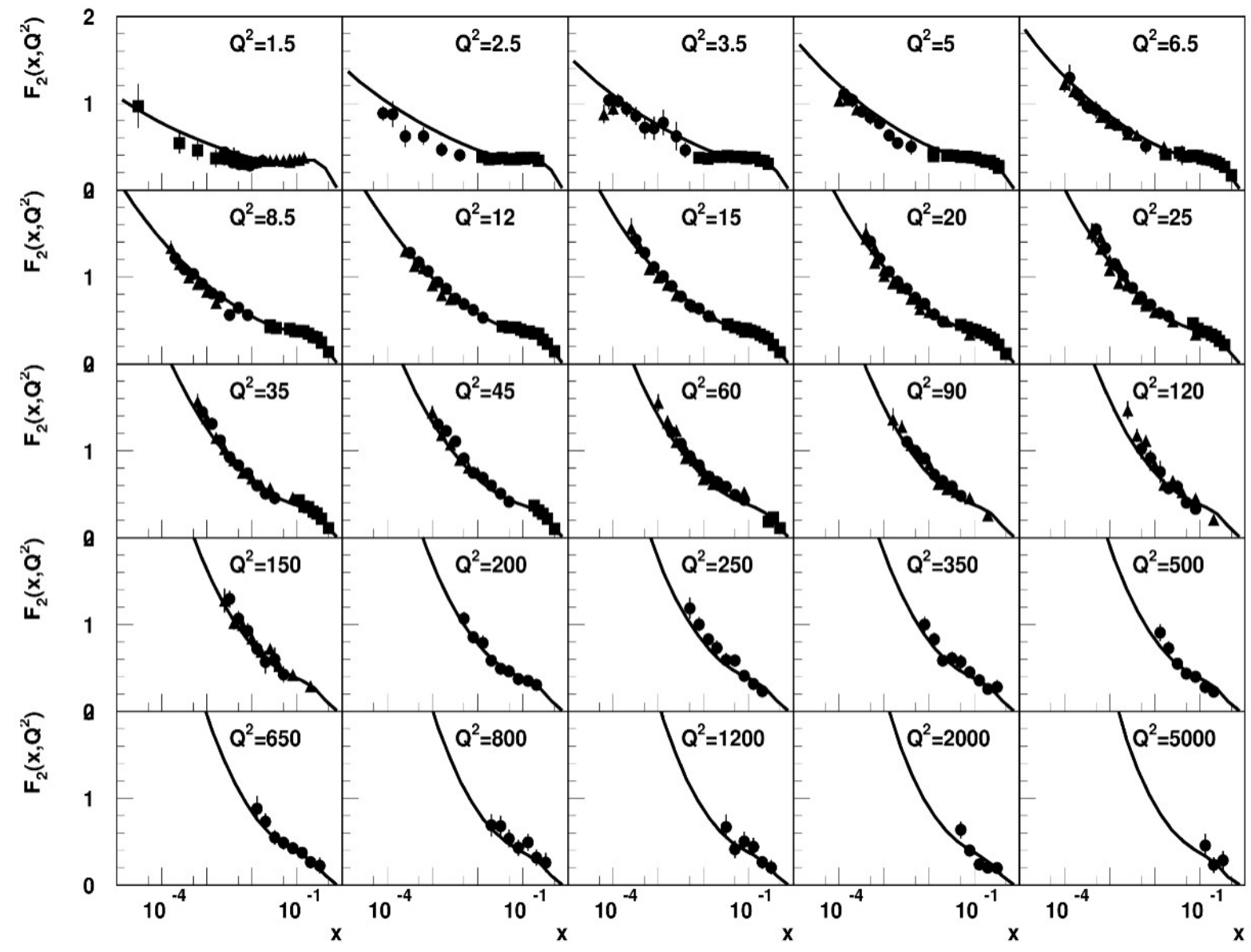
➔ large increase at small x (without saturation)

- External pdf only for valence quark

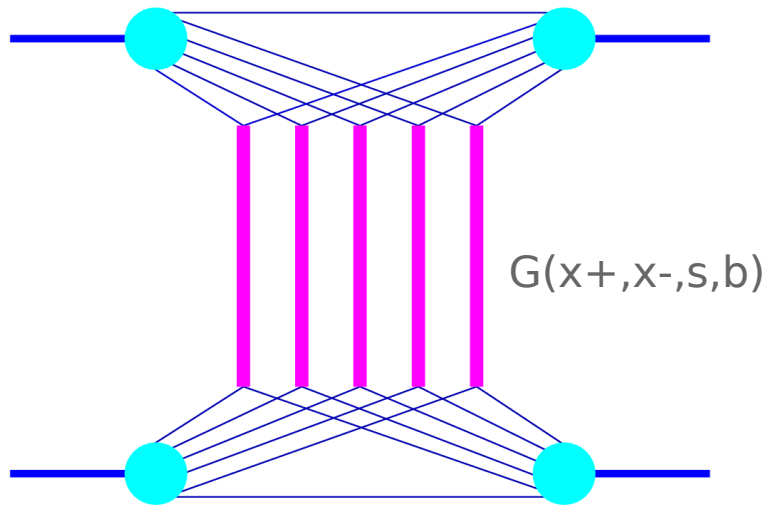
- Minimum non-perturbative scale $Q_n^2 = 2 \text{ GeV}^2$ with soft pre-evolution $s^{\alpha(0)-1}$

- F2 from HERA used to fix parameters for sea quarks and gluons below Q_n^2

EPOS Parton Distribution Function $Q_n^2=2 \text{ GeV}^2$



Cross Section Calculation : EPOS



- ➔ Gribov-Regge but with energy sharing at parton level (Parton Based Gribov Regge Theory)
- ➔ amplitude parameters fixed from QCD and pp cross section (semi-hard Pomeron)
- ➔ cross section calculation take into account interference term

$$\sigma_{\text{ine}}(s) = \int d^2b (1 - \Phi_{\text{pp}}(1, 1, s, b))$$

$$\Phi_{\text{pp}}(x^+, x^-, s, b) = \sum_{l=0}^{\infty} \int dx_1^+ dx_1^- \dots dx_l^+ dx_l^- \left\{ \frac{1}{l!} \prod_{\lambda=1}^l -G(x_\lambda^+, x_\lambda^-, s, b) \right\} \\ \times F_{\text{proj}}\left(x^+ - \sum x_\lambda^+\right) F_{\text{targ}}\left(x^- - \sum x_\lambda^-\right).$$

can not use complex diagram with energy sharing:
non linear effects taken into account as correction of single amplitude G

EPOS – non-linear effects

Well known problem with pQCD based Pomerons

➔ total cross-section too high : MPI required

➔ in EPOS $\langle \text{Pomerons} \rangle$ fixed by b-dep of Pomeron amplitude (slope)

➔ effective coupling introduced to mimic effect of enhanced diagrams and reduce cross-section (screening effect) to get cross-section AND multiplicity right in p-p, p-A and AA

➔ Amplitude G_{eff} no longer fit to G_{QCD}

No effective coupling

$$G_{\text{QCD}} \sim (x_1 x_2)^\beta$$

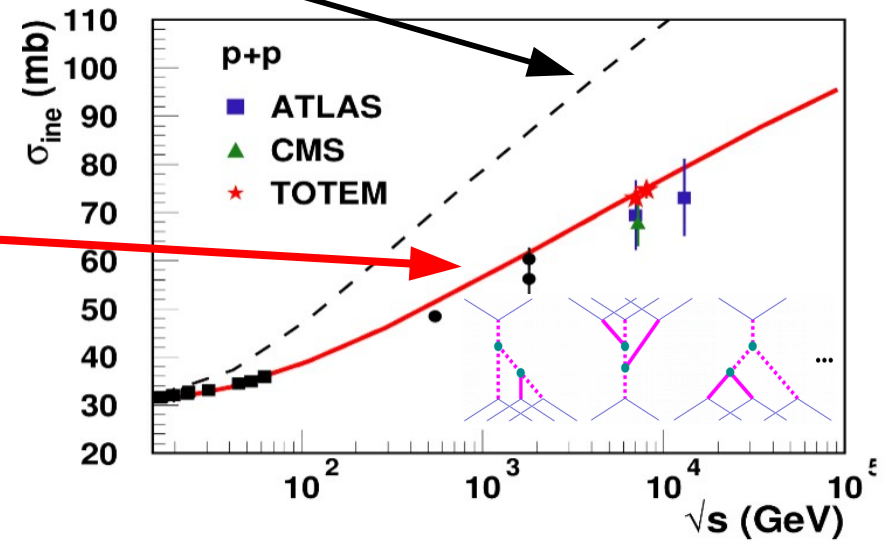
With effective coupling

$$G_{\text{eff}} \sim x_1^\beta x_2^{\beta-\epsilon}$$

Parametrization

$$\epsilon_S = a_S \beta_S Z(s,b,A)$$

$$\epsilon_H = a_H \beta_H Z(s,b,A)$$



Particle Production in EPOS

m number of exchanged elementary interaction per event fixed from elastic amplitude taking into account energy sharing :

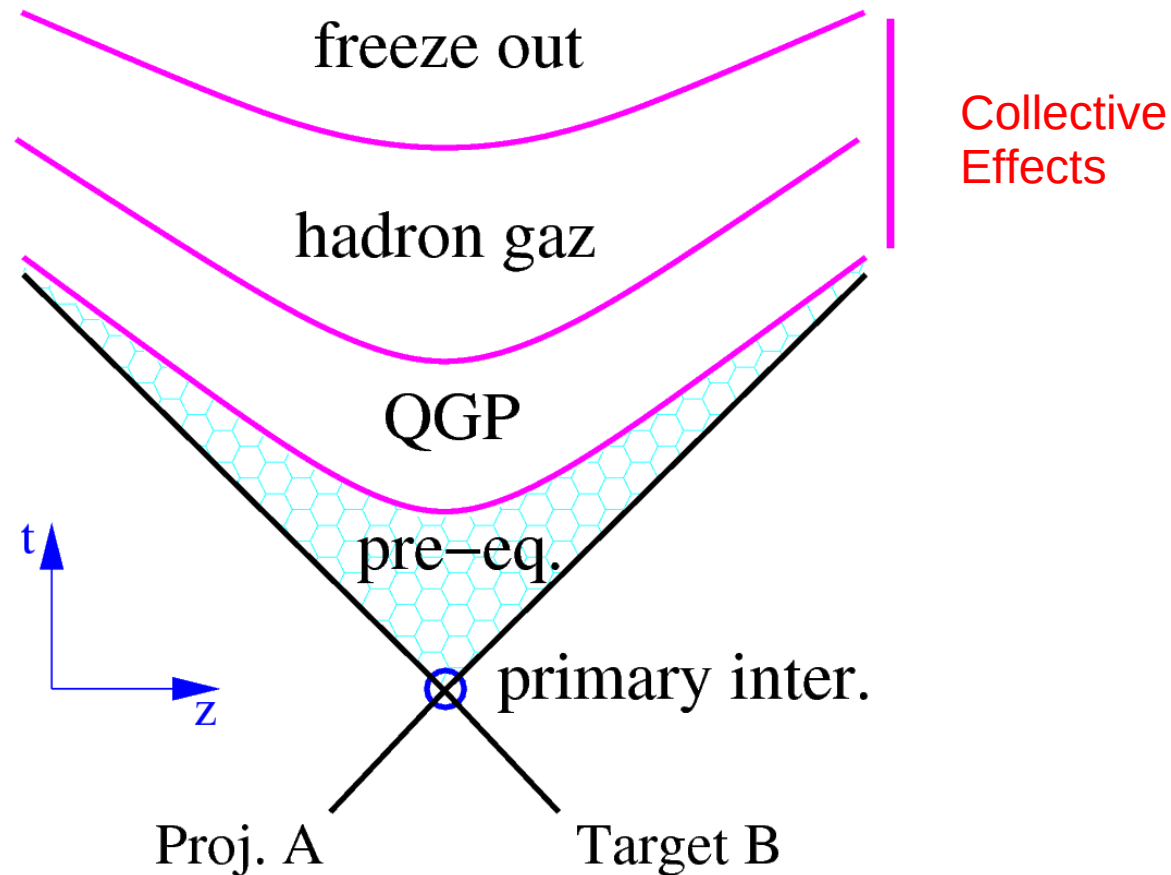
➔ m cut Pomerons from :

$$\Omega_{AB}^{(s,b)}(m, X^+, X^-) = \prod_{k=1}^{AB} \left\{ \frac{1}{m_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, b_k) \right\} \Phi_{AB}(x^{\text{proj}}, x^{\text{targ}}, s, b)$$

- m and X fixed together by a complex Metropolis (Markov chain)
- ➔ 2m “kinky” strings formed from the m elementary interactions
- **energy conservation** : energy fraction of the 2m strings given by X
- ➔ consistent scheme : energy sharing reduce the probability to have large m

To reconcile minimum bias (MB) and underlying events (UE) in pp, we need both collective effects and variable non-linear perturbative scale : impact on HQ production.

High Energy Hadronic Interactions

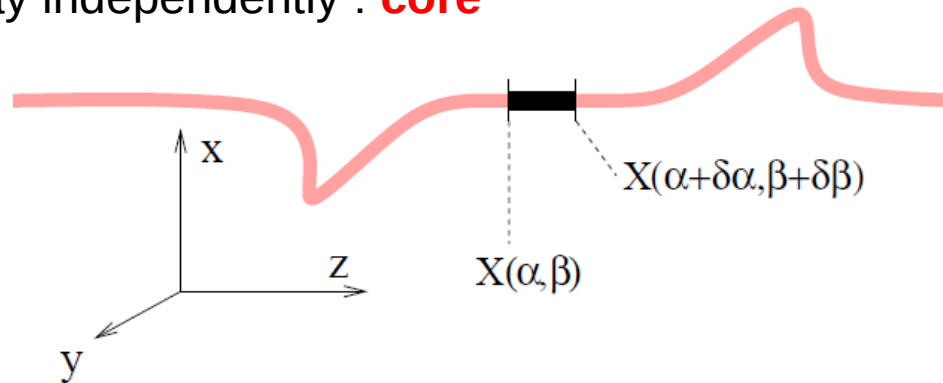
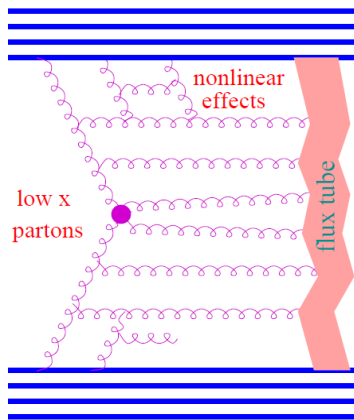


For a complete description of LHC pp data (from minimum-bias to high multiplicity), same process chain as for heavy ion needed.

High Density Core Formation

Heavy ion collisions or high energy proton-proton scattering:

➔ the usual procedure has to be modified, since the density of strings will be so high that they cannot possibly decay independently : **core**



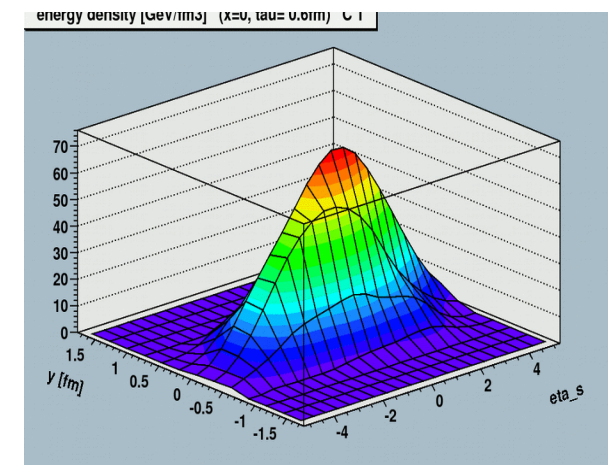
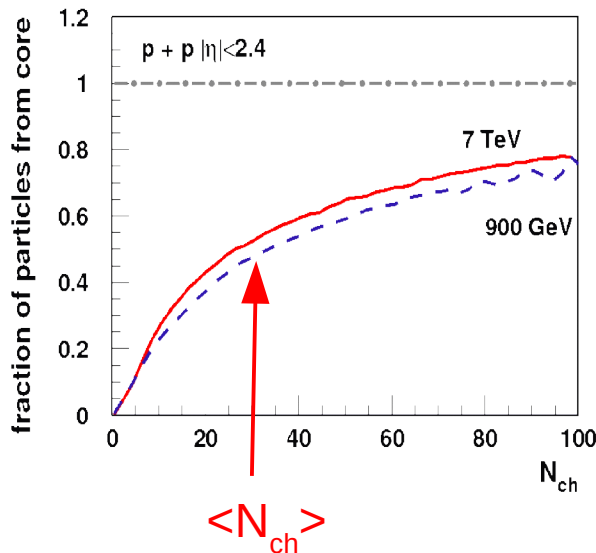
➔ Each string split into a sequence of string segments, corresponding to widths $\delta\alpha$ and $\delta\beta$ in the string parameter space

➔ If energy density from segments high enough

- ◆ segments fused into core
 - full 3D+1 hydro evolution
 - lattice QCD EoS

➔ If low density (corona)

- ◆ segments remain hadrons
 - string fragmentation



EPOS LHC

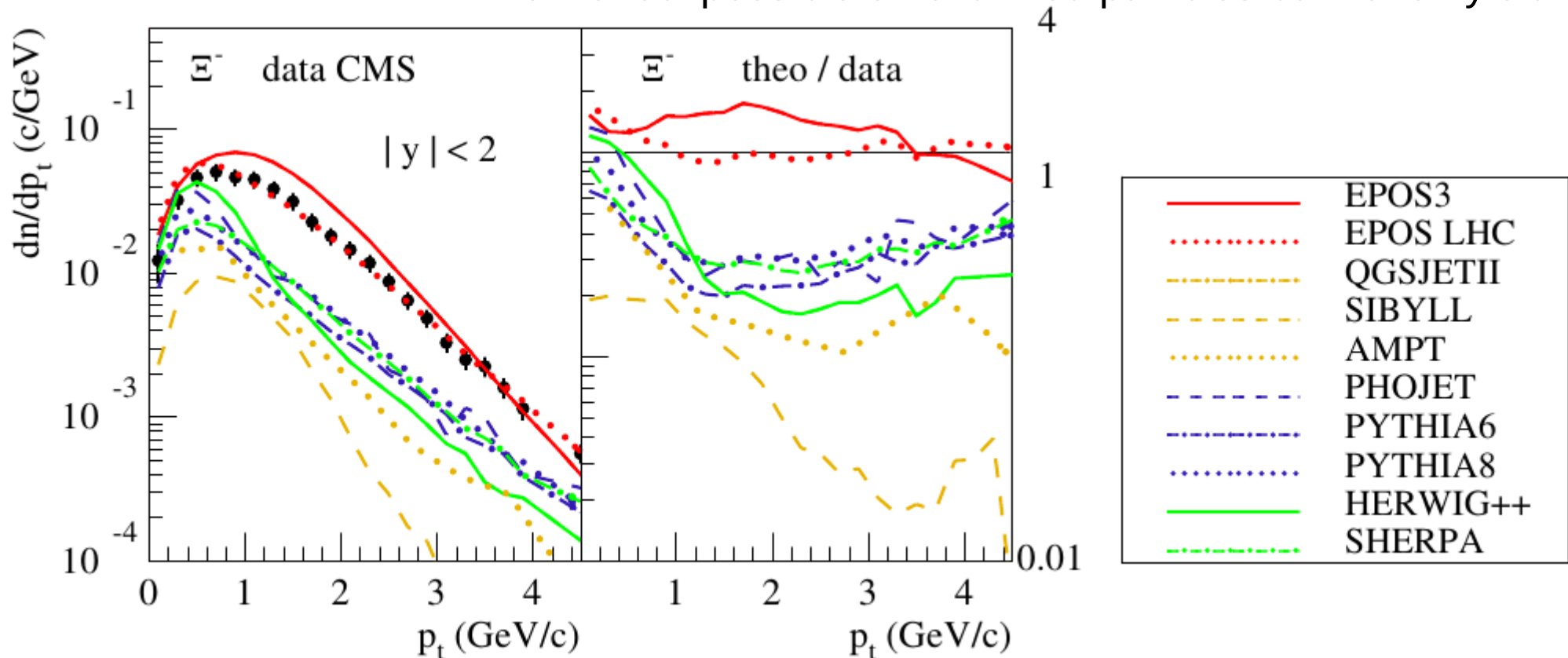
Detailed description can be achieved

➔ identified spectra

➔ p_t behavior driven by collective effects (statistical hadronization + flow)

➔ large effect for multi-strange baryons (yield AND $\langle p_t \rangle$)

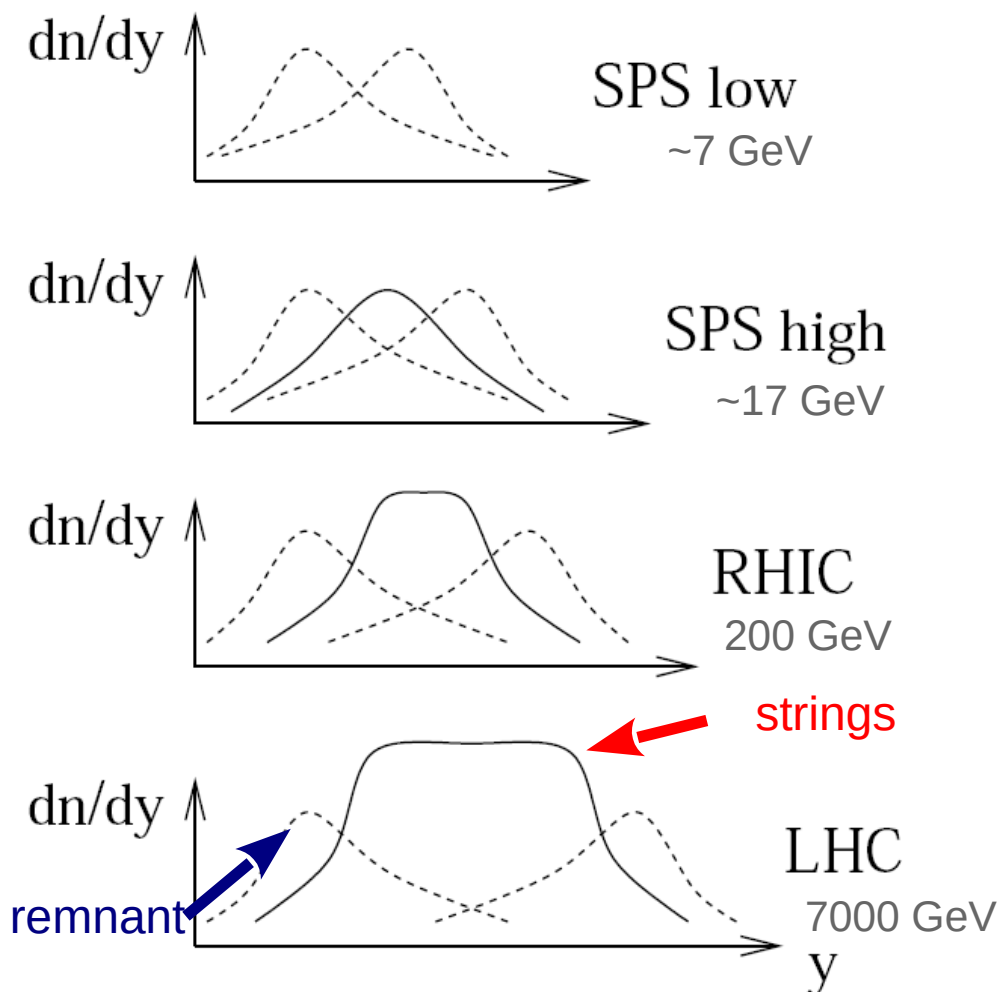
➔ flow effect possible on charmed particles but not on yield ?



Remnants

Forward particles mainly from projectile remnant

Forward hadronization from remnant :

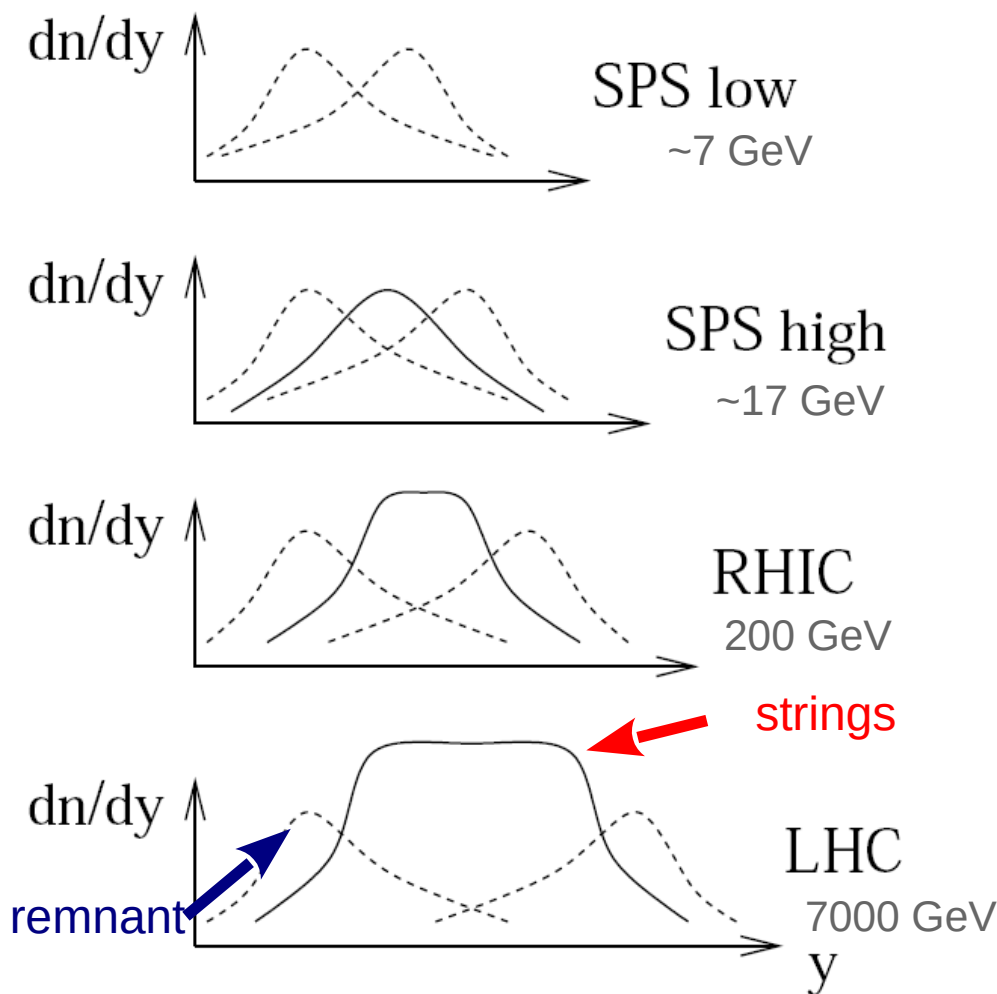


- ➔ At very low energy only particles from remnants
- ➔ At low energy (fixed target experiments) (SPS) strong mixing
- ➔ At intermediate energy (RHIC) mainly string contribution at mid-rapidity with tail of remnants.
- ➔ At high energy (LHC) only strings at mid-rapidity (baryon free)

Remnant considered as universal object : **same behavior at low or high energy**

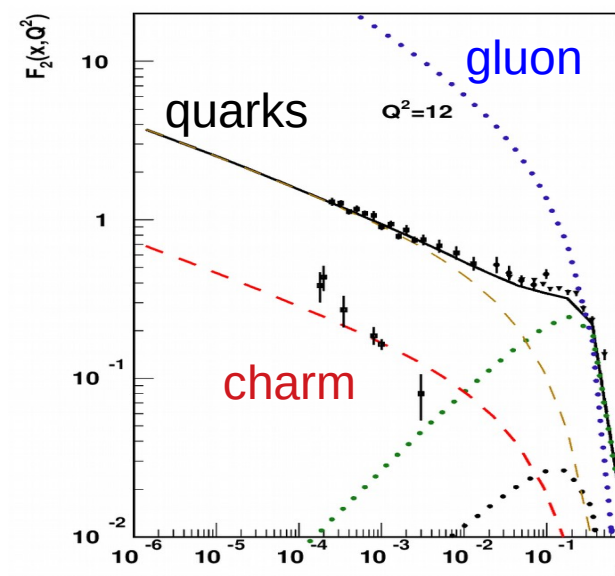
Remnants

Forward particles mainly from projectile remnant



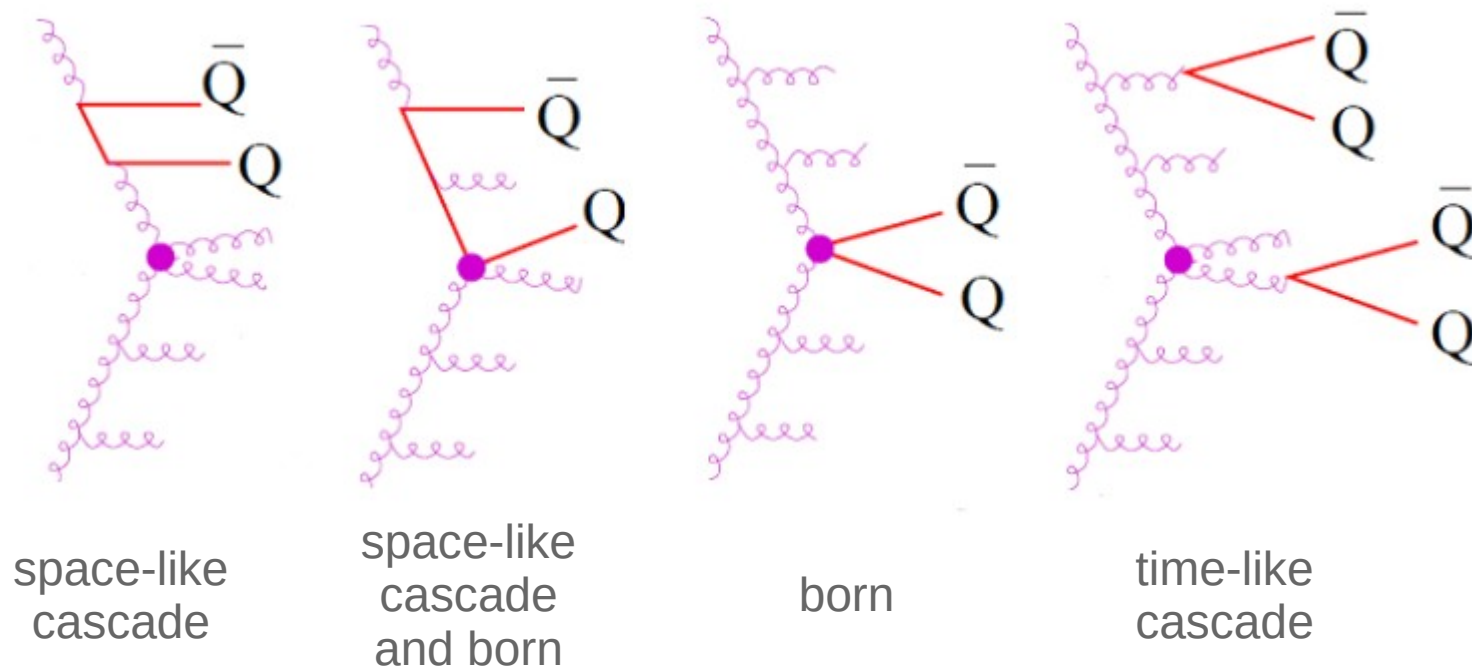
Forward hadronization from remnant :

- ➔ Place for forward charm
- Soft intrinsic charm ?
- Coalescence ?
- ➔ Not yet implemented test with low energy data
- ➔ But charm in PDF to start parton cascade



Heavy Flavor Production

Heavy flavor production included in perturbative ($Q^2 > Q_n^2$) calculation in EPOS 3

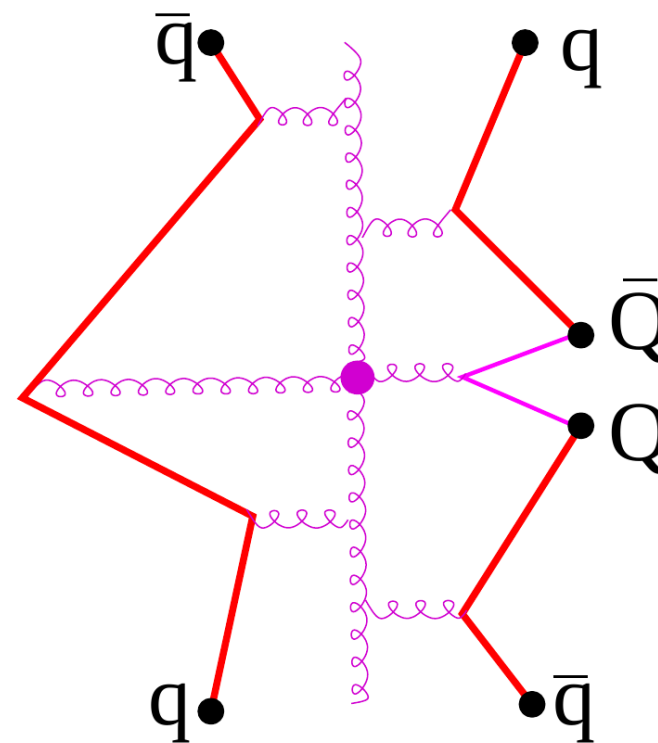


- ➔ “parameter free” : good test of hard Pomeron
- ➔ Heavy quarks (Q) taken as string-end for the hadronization

Hadronization

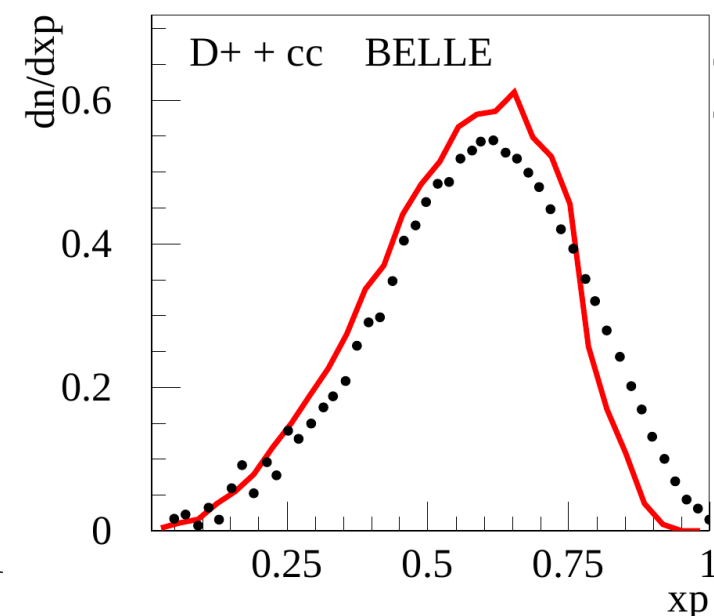
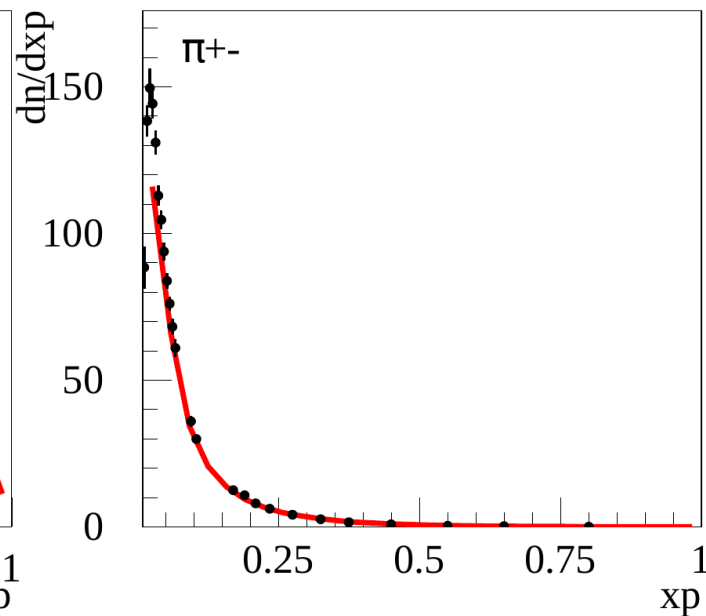
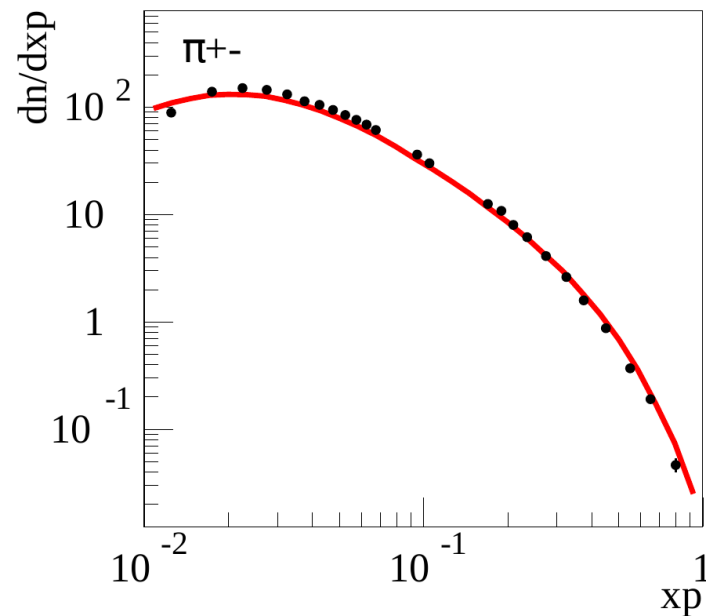
Heavy Quark String

- ➔ 2 “kinky” strings from one Pomeron (=parton ladder)
- ➔ HQ always used as string ends
- ➔ Hadronization tested with e^+e^- data



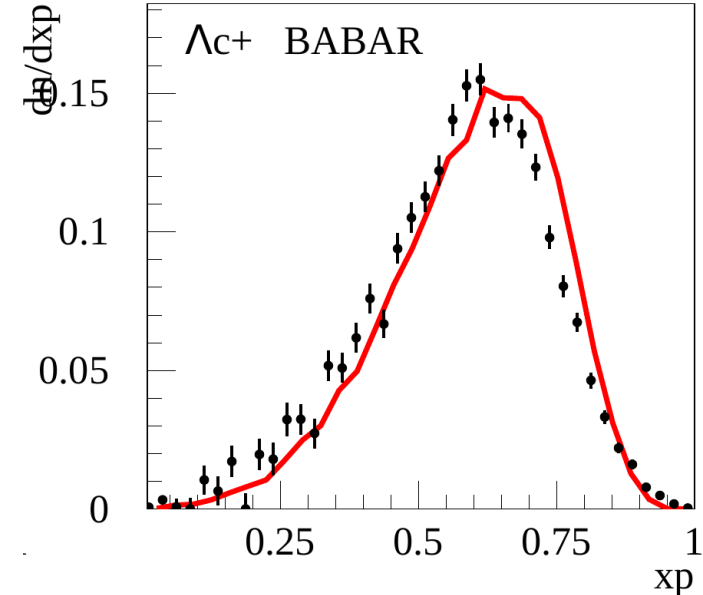
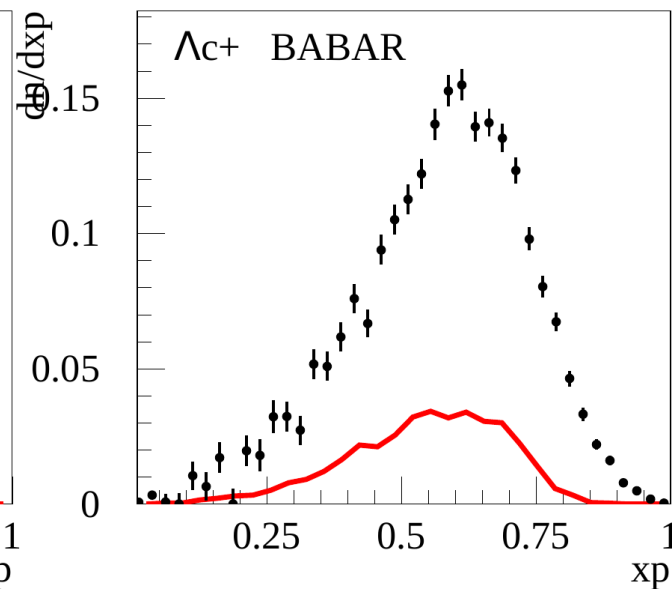
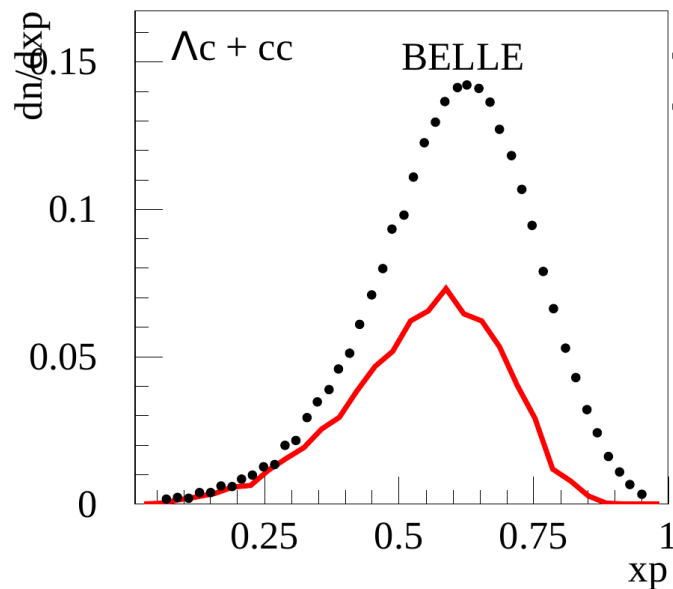
Test of Hadronization with e+e- Data : Mesons

- ➔ LEP data used to fix light hadron hadronization (area law used for string fragmentation, not Lund model)
 - ◆ Free parameters = effective mass of quarks and diquarks
 - ◆ String tension
- ➔ $x_p = E/E_{\max}$ distribution much harder for charmed meson : **HQ from string ends only** (not produced during fragmentation : only 3 (light) flavors)

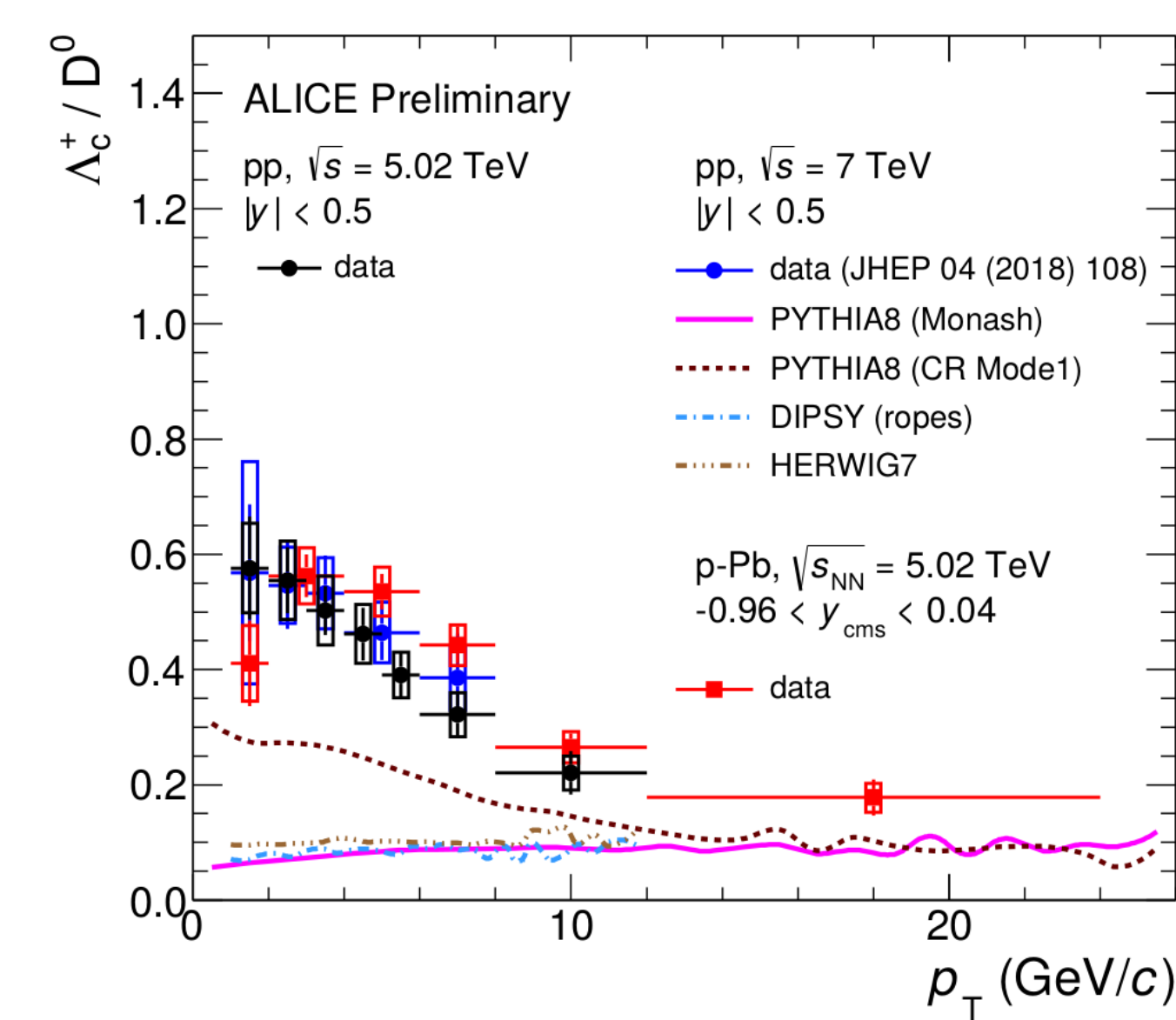
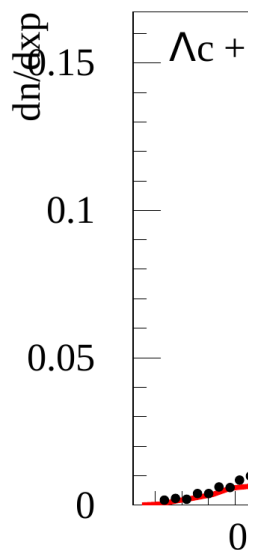


Test of Hadronization with e+e- Data : Baryons

- ➔ LEP data used to fix light hadron hadronization (area law used for string fragmentation, not Lund model)
 - ◆ Free parameters = effective mass of quarks and diquarks
 - ◆ String tension
- ➔ $x_p = E/E_{\max}$ distribution much harder for charmed meson : HQ from string ends only (not produced during fragmentation : only 3 (light) flavors)
- ➔ **Charmed baryon production underestimated using diquark mass suppression**
 - ◆ Correct description of data if diquark mass is not used in diquark break probability ($\sim \times 10$ compared to light quarks only)



Test of Hadronization with e+e- Data : Baryons



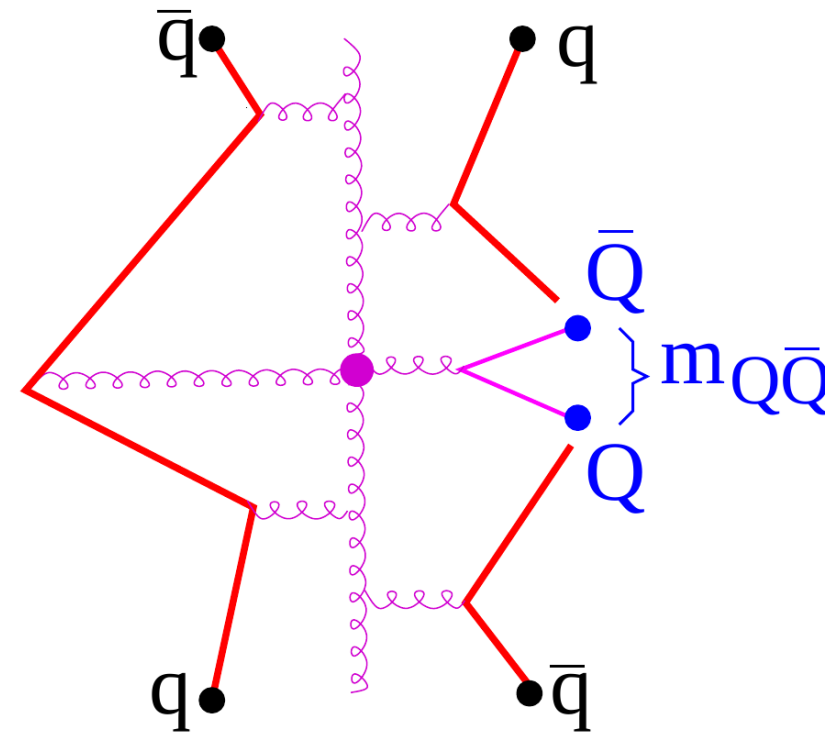
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Hadronization

Heavy Quark String

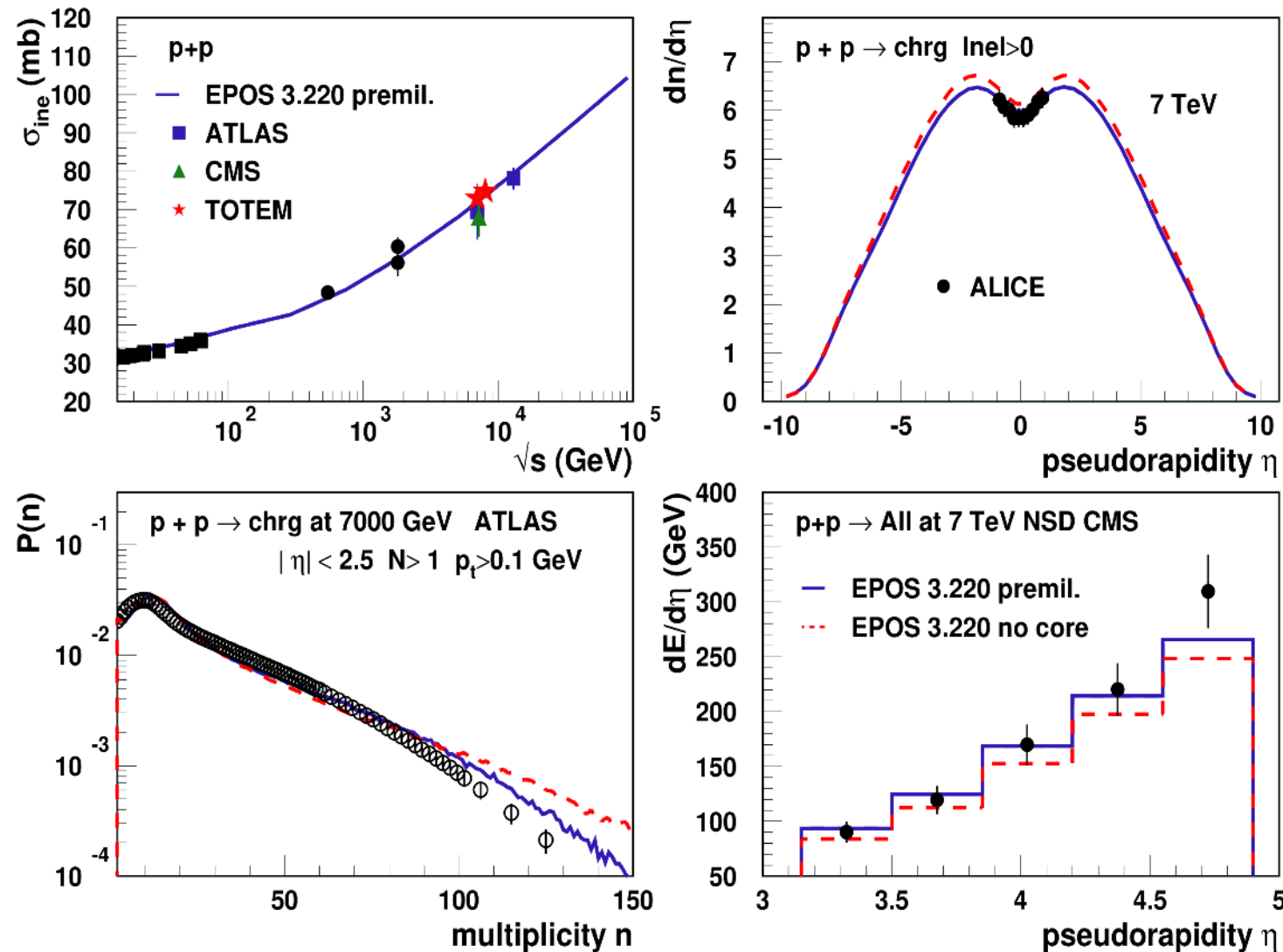
- ➔ 2 “kinky” strings from one Pomeron (=parton ladder)
- ➔ HQ always used as string ends
- ➔ Hadronization tested with e^+e^- data
- ➔ Quarkonium produced in born or TLC by pair production
 - Compute mass of the $Q\bar{Q}$ system
 - If $2m_Q < m_{Q\bar{Q}} < 2m_{D(B)}$

then $Q\bar{Q}$ =quarkonium
with probability $w_{Q\bar{Q} \rightarrow J/\Psi(\Upsilon)}$
(color evaporation model)



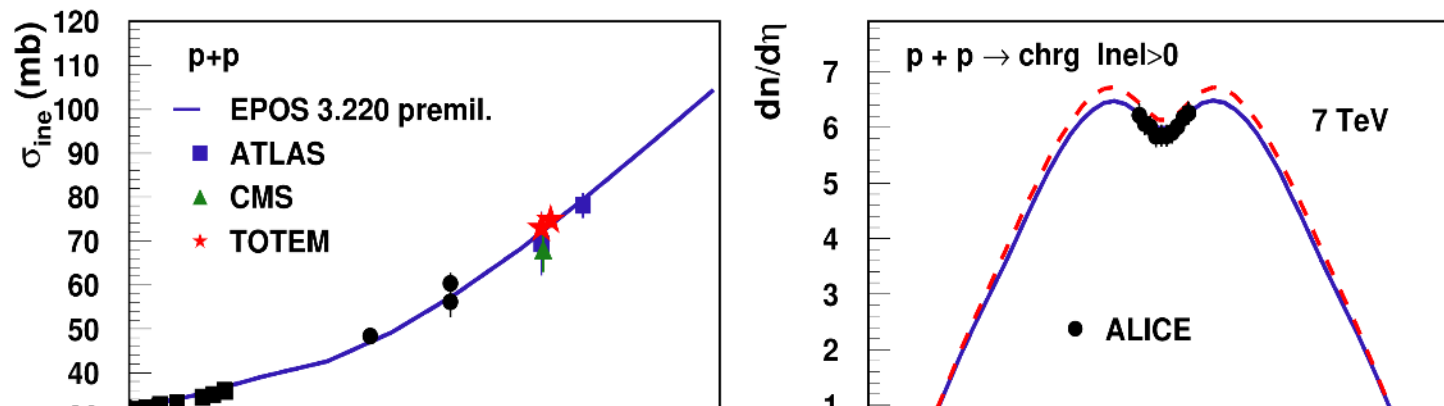
Preliminary Results : With/out Core

Excellent results for minimum bias soft physics

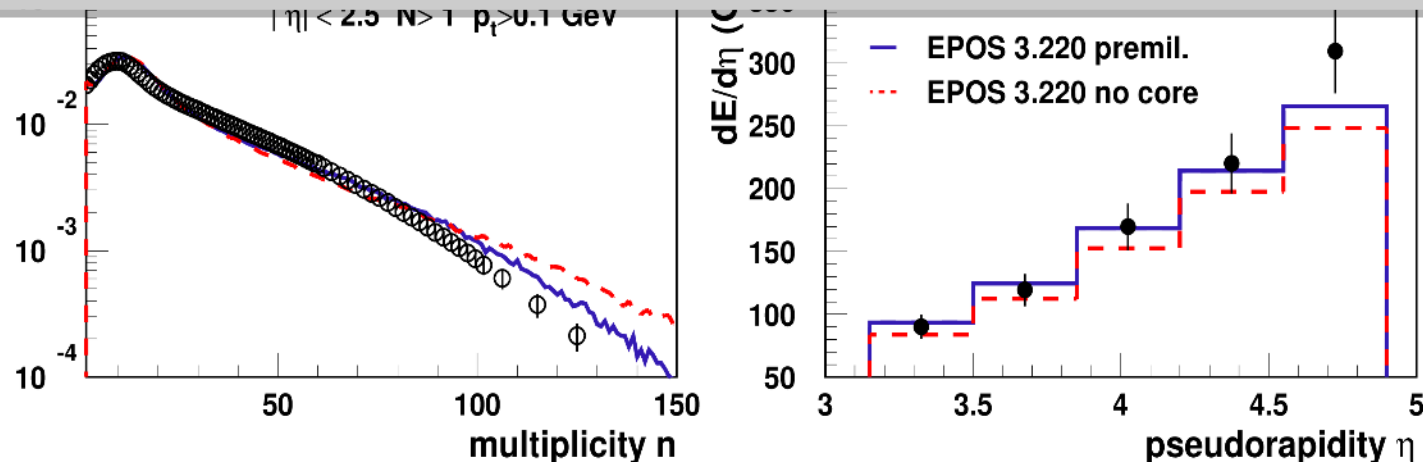


Preliminary Results : With/out Core

Excellent results for minimum bias soft physics



Nice but what to do with harder scales ? Jets, UE, etc ...
Can we recover G_{QCD} ?



EPOS LHC : Fixed $Q_0^2 = Q_n^2$ (old)

- Excellent results for soft physics

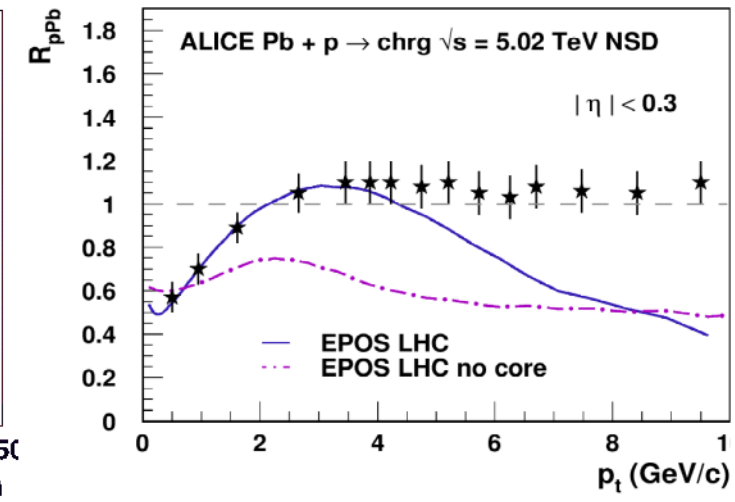
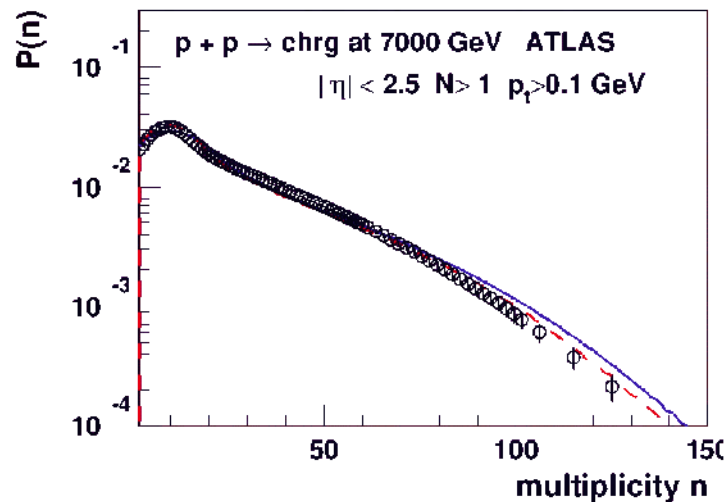
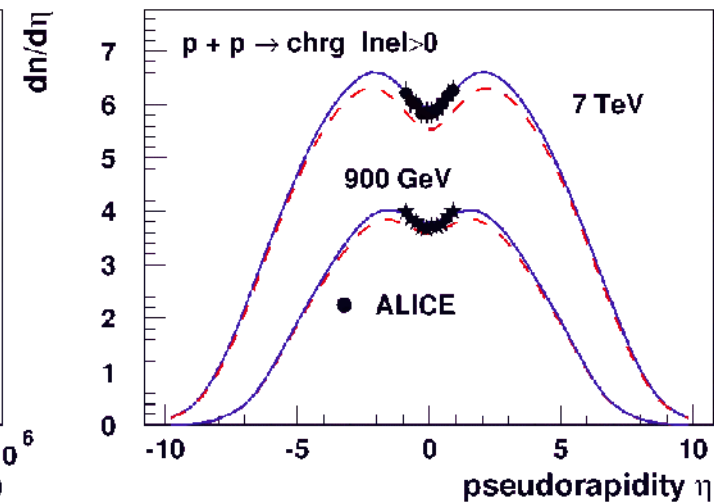
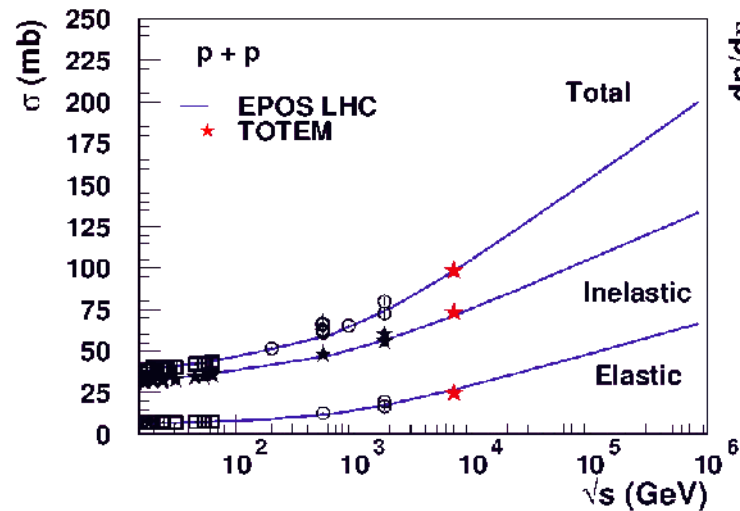
- ➔ cross-section, multiplicity, etc ...

- Problem for hard processes

- ➔ Minimum scale for DGLAP = $Q_0^2 = Q_n^2$

- ➔ lack of high p_t

- ➔ no binary scaling for pA or AB



Since Q_0^2 is fixed both low and high p_t are suppressed: in contradiction with data.

Jet Production in EPOS

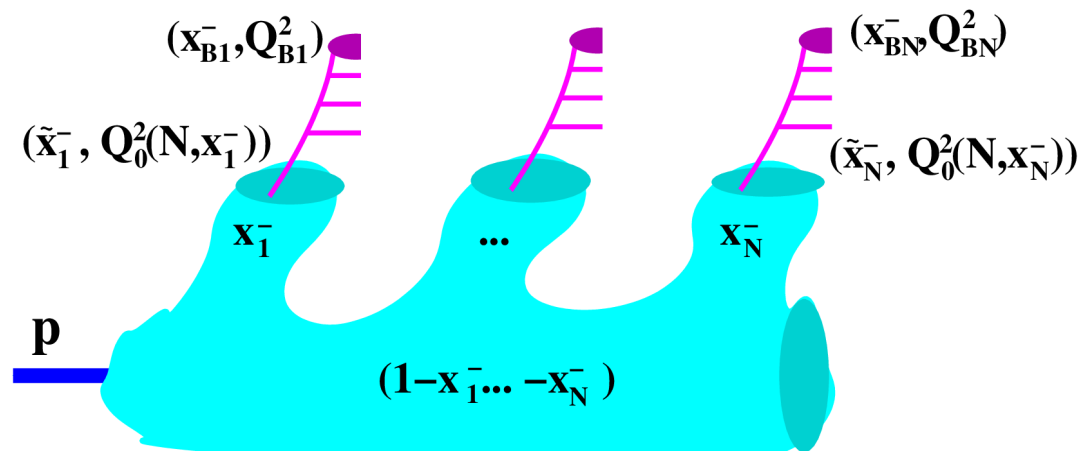
m number of exchanged elementary interaction per event fixed from elastic amplitude taking into account energy sharing :

→ m cut Pomerons from :

$$\Omega_{AB}^{(s,b)}(m, X^+, X^-) = \prod_{k=1}^{AB} \left\{ \frac{1}{m_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, b_k) \right\} \Phi_{AB}(x^{\text{proj}}, x^{\text{targ}}, s, b)$$

■ m(=N) and $X = \{x_1, \dots, x_N\}$ fixed together by a complex Metropolis (Markov chain)

→ if $G_{\text{eff}} = G_{\text{QCD}}$ then each hard elementary interaction will be a minijet



EPOS as an N-pdf generator (event-by-event) if $Q_0^2(N, x)$ could be determined !

Non-linear Perturbative Scale Q_0^2

Model property : AGK cancellation

$$\begin{aligned} \frac{dn_{\text{Pom}}^{h_1 h_2}}{dx^+ dx^-}(x^+, x^-, s, b) &= \frac{dn_{\text{Pom}}^{(1)h_1 h_2}}{dx^-}(x^+, x^-, s, b) \\ &= G_{\text{eff}}(x^+, x^-, s, b) F_{\text{remn}}^{h_1}(1-x^+) F_{\text{remn}}^{h_2}(1-x^-) \end{aligned}$$

Assumption : factorization should be satisfied at large Q^2

→ satisfied if: $\langle N_{\text{hard}} \rangle G_{\text{QCD}}(x, b, Q_0^2) = G_{\text{eff}}(s, x, b, A)$

→ different non-linear perturbative scale event-by-event and even Pomeron-by-Pomeron depending on momentum fraction x

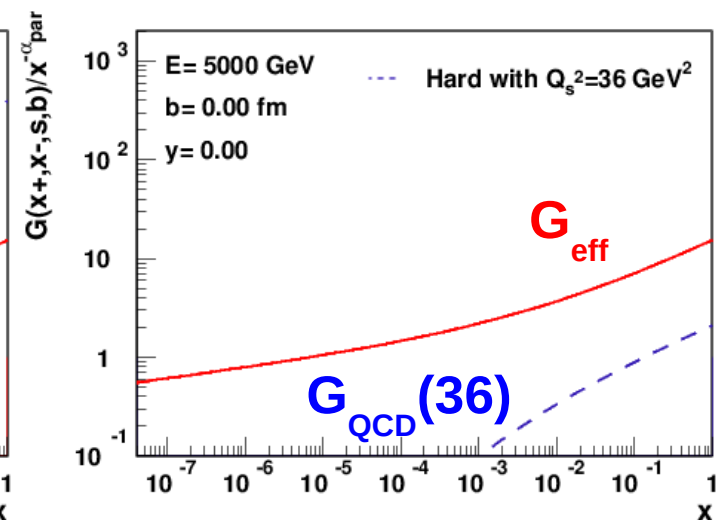
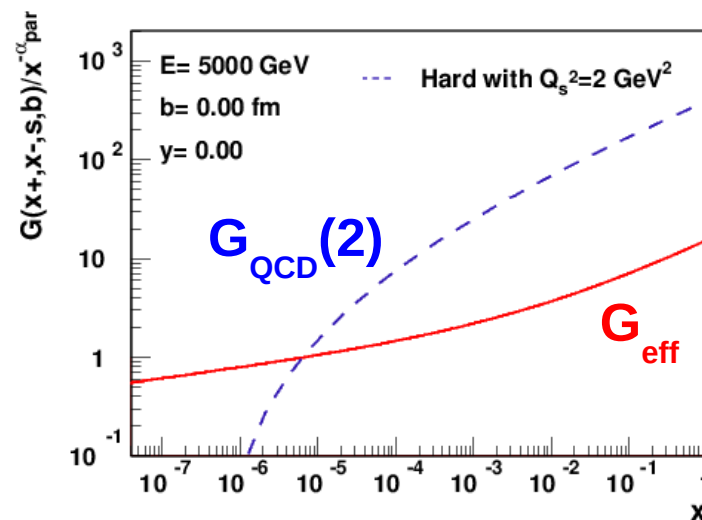
Matching amplitude

→ $G_{\text{eff}} \leq G_{\text{QCD}}(Q_0^2 = Q_n^2)$

→ increase Q_0^2 until

$G_{\text{eff}} = \langle N_{\text{hard}} \rangle G_{\text{QCD}}(Q_0^2)$
for each parton scattering

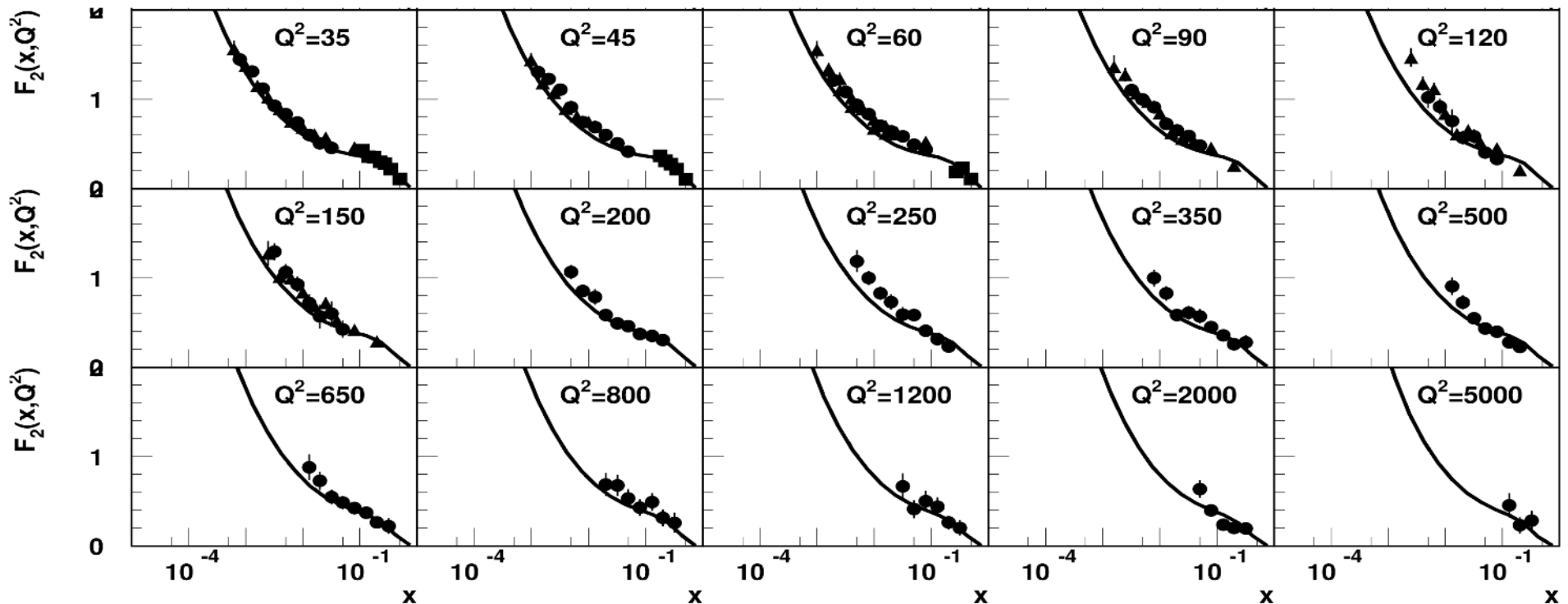
→ for $Q^2 \gg Q_0^2(x, b)$
factorization holds



EPOS Parton Distribution Function $Q_0^2=30 \text{ GeV}^2$

● Larger Q_0^2

- ➔ partons which can be treated perturbatively (linear DGLAP evolution) and independently have already a large virtuality
- ➔ soft preevolution changed to get the same parton distribution than with Q_n^2 (perturbative but non-linear evolution)
- ➔ PDF for $Q^2 > Q_0^2$ independent of Q_0^2

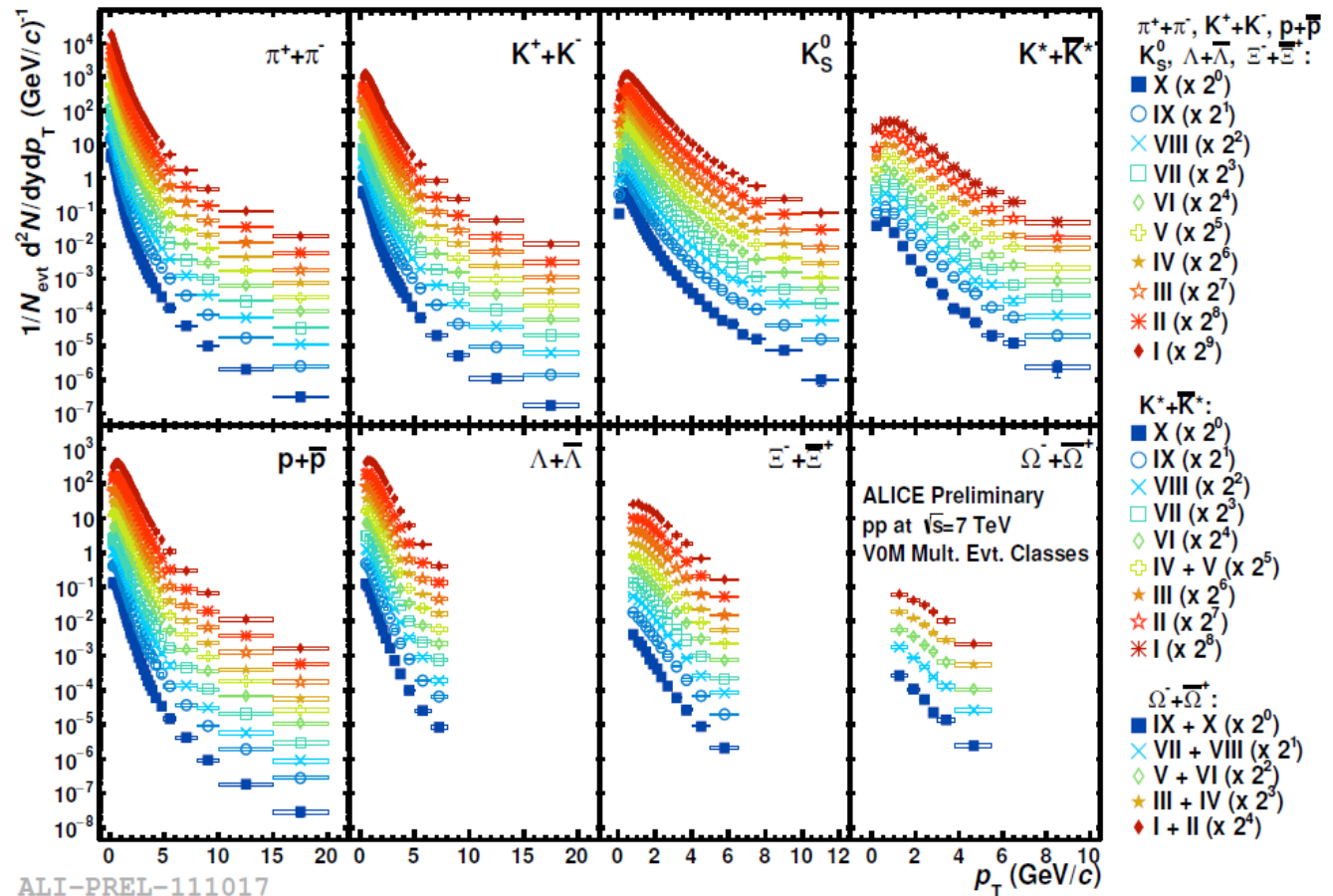
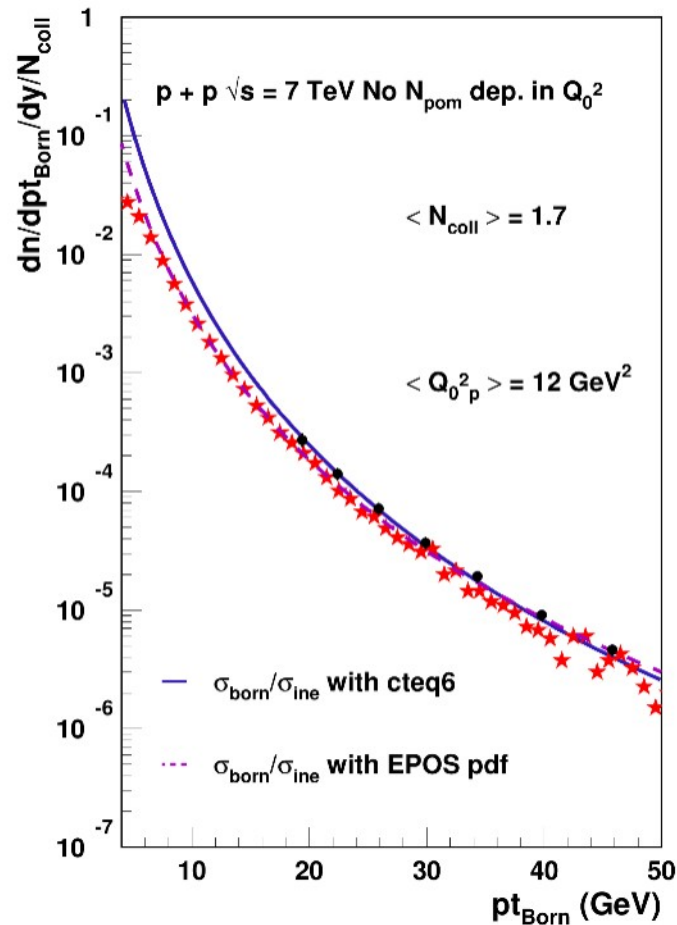


Jet and p_t

Check with pp data at 7 TeV

➔ inclusive jet cross section: **OK**

➔ transverse momentum for different centrality bins



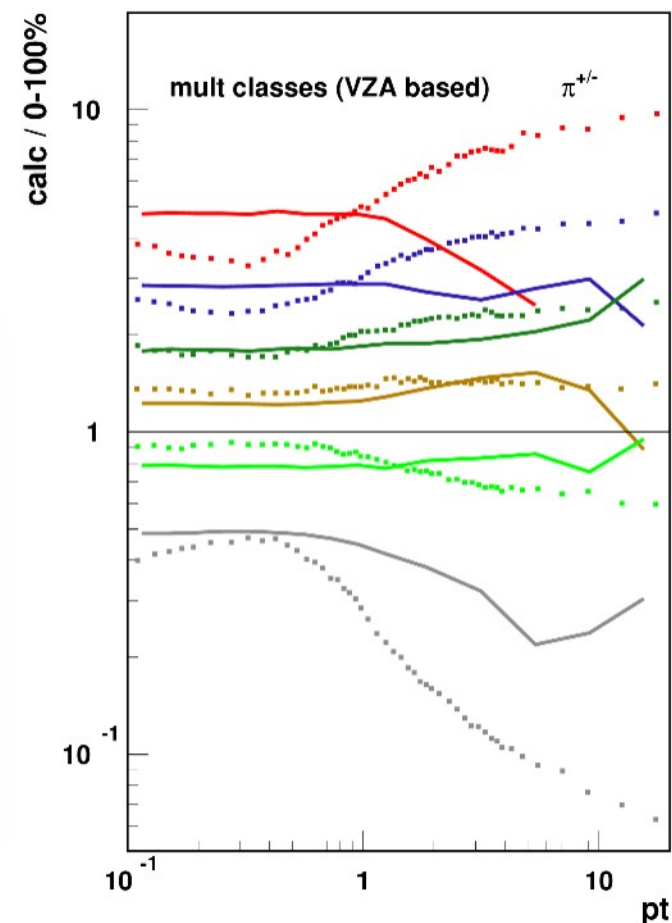
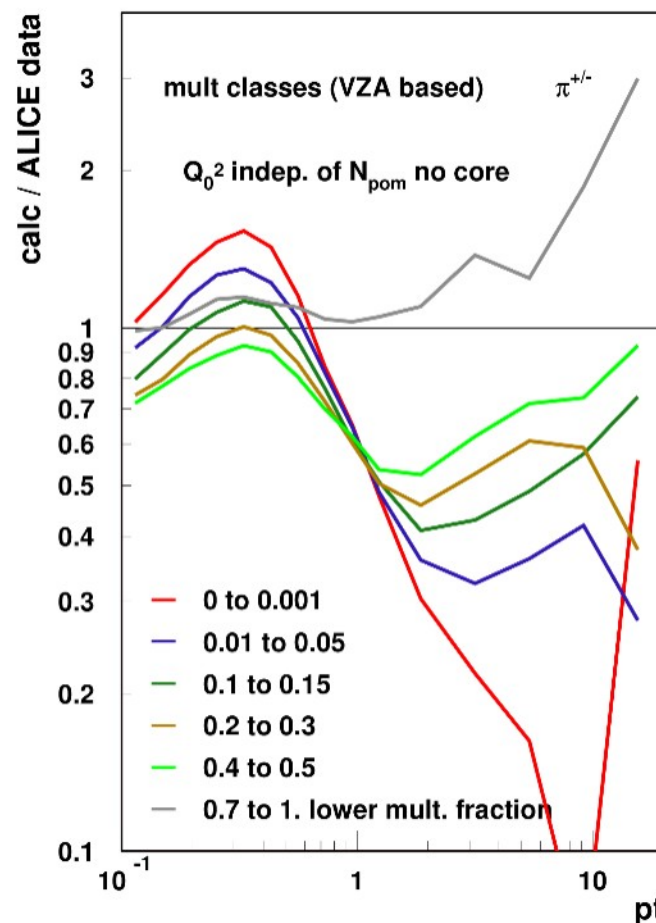
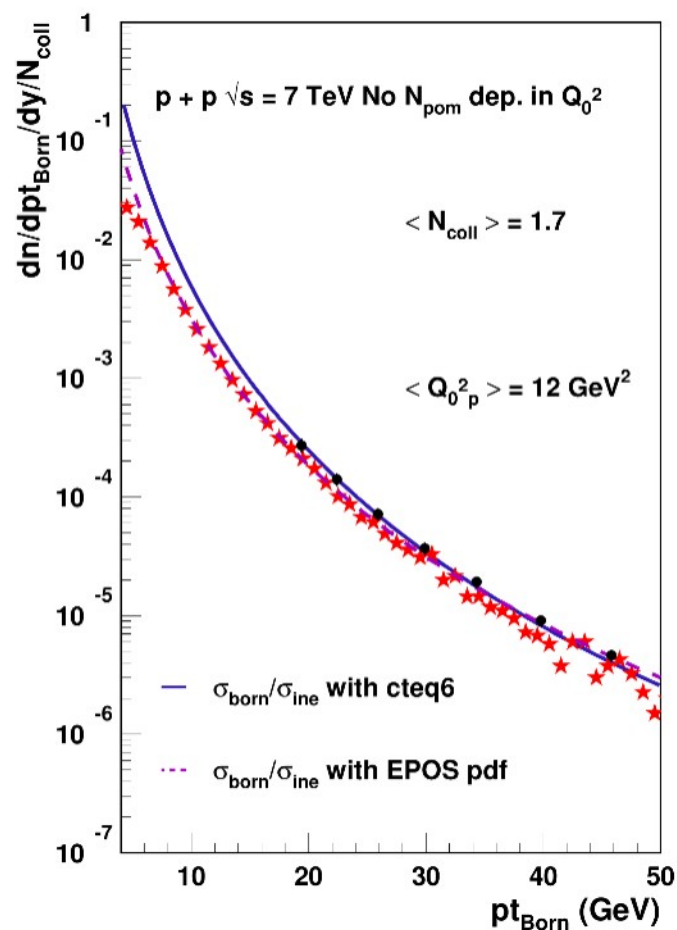
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Jet and p_t

Check with pp data at 7 TeV

- ➔ inclusive jet cross section: **OK**
- ➔ transverse momentum for different centrality bins: **NO**

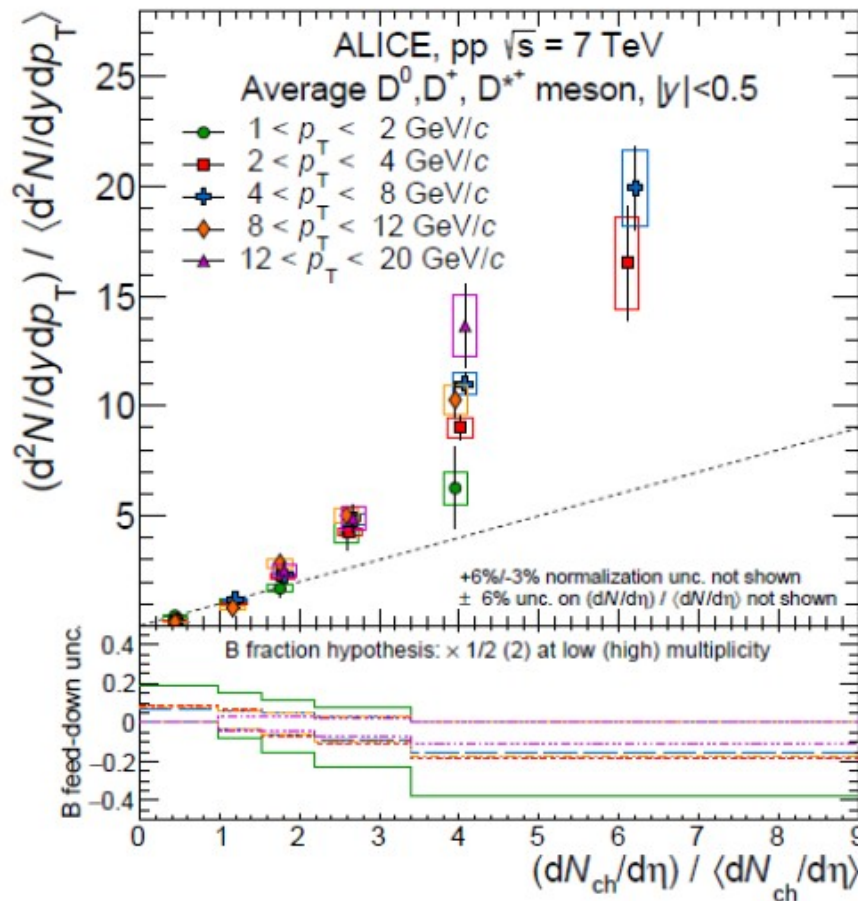
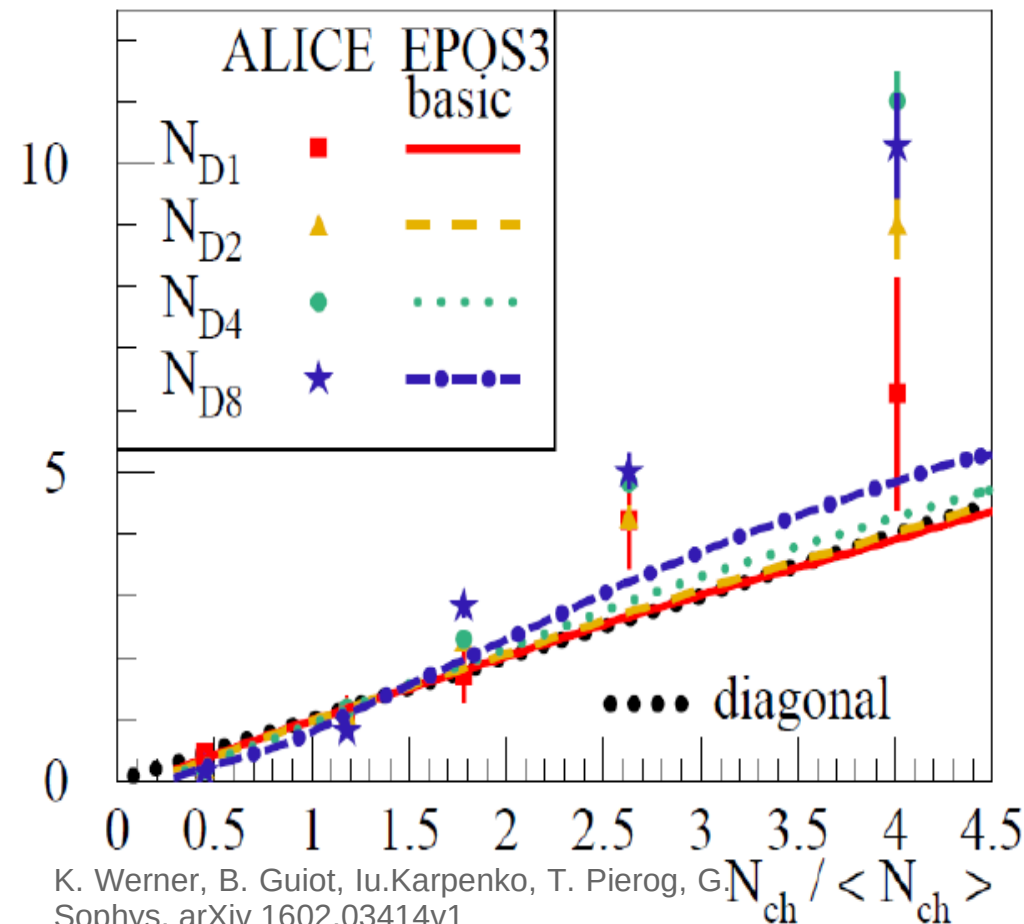
Same slope for all multiplicities while data not flat



Charm Production

Similar behavior observed in D meson but presented in a different way

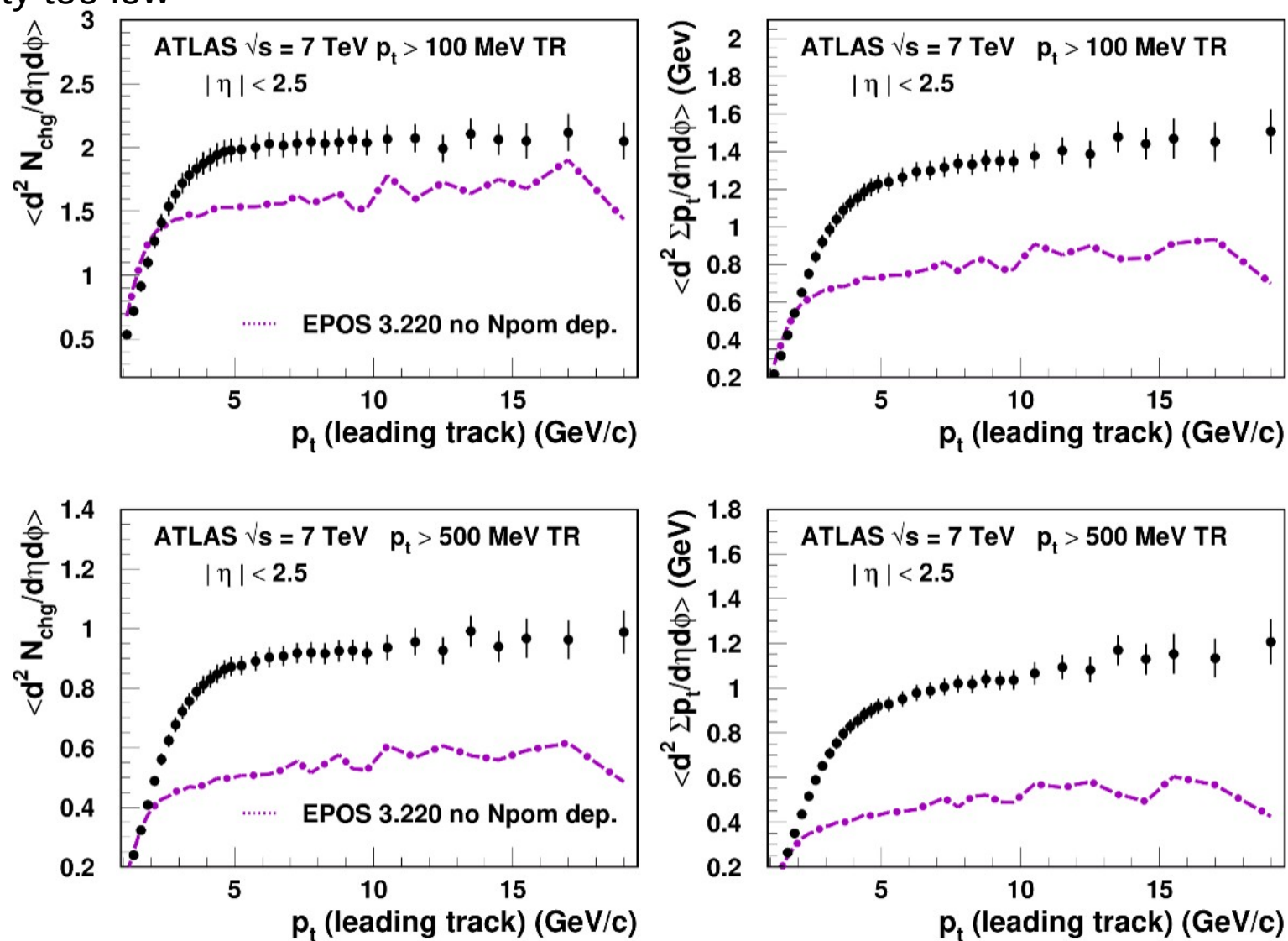
- ➔ more than linear increase of charm production and larger in higher pt bin = hardening of pt spectra with particle multiplicity
- ➔ small increase due to fluctuations observed in EPOS 3 but not sufficient to reproduce data


 $\frac{\langle N \rangle}{N}$


Underlying Events

Check with pp data at 7 TeV

➔ activity too low

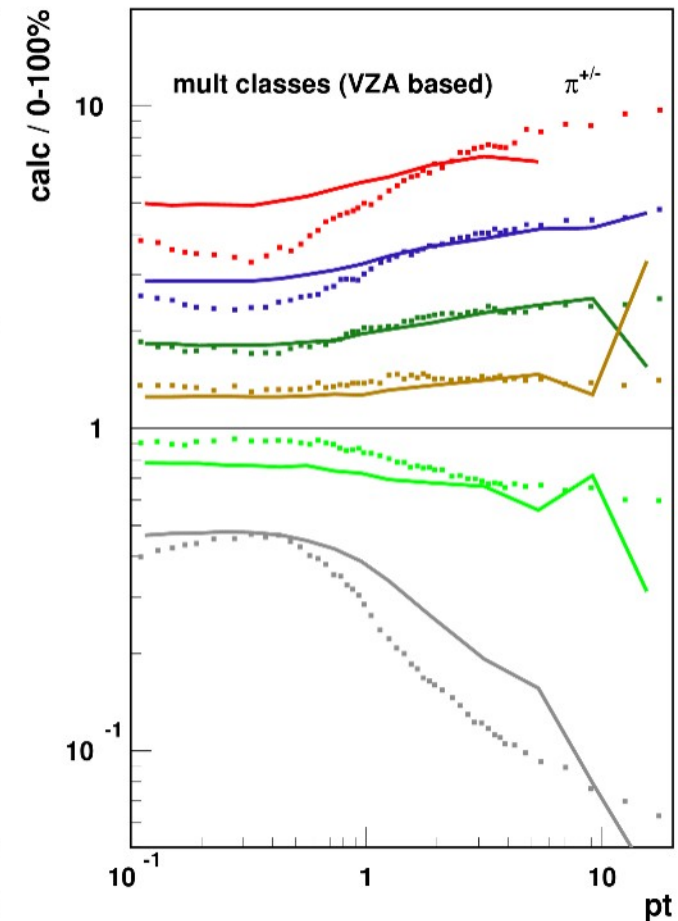
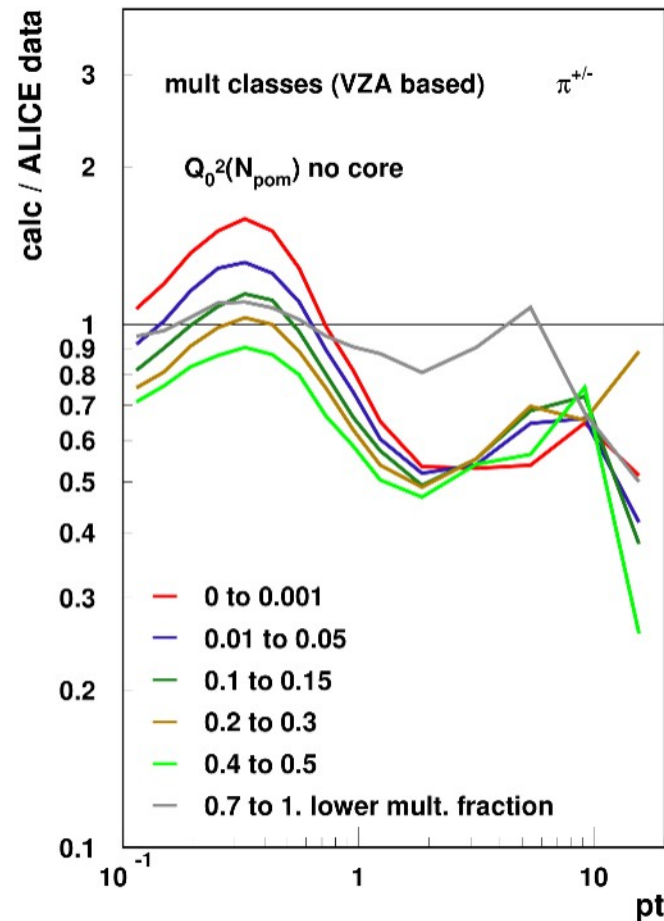
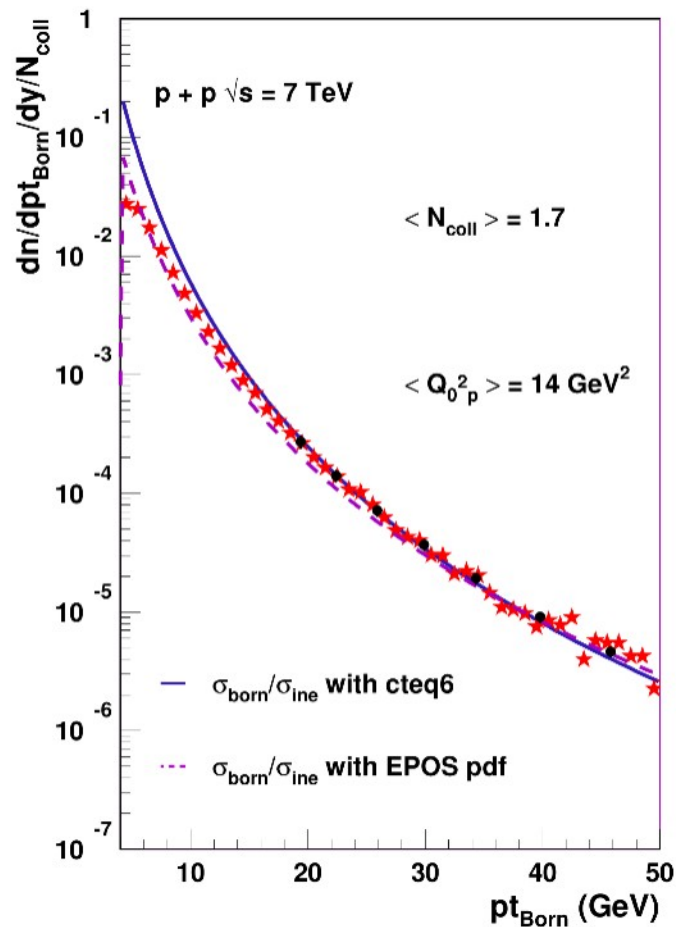


Effect of MPI on Q_0^2

Is it possible to introduce the number of parton scattering N_{hard} in Q_0^2 ?

→ $\langle N_{\text{hard}} \rangle G_{\text{QCD}}(x,b,Q_0^2) = G_{\text{eff}}(s,x,b,A)$ on average but for each event we can define :

$$\rightarrow N_{\text{hard}} G_{\text{QCD}}(x,b,Q_0^2) = G_{\text{eff}}(s,x,b,A)$$



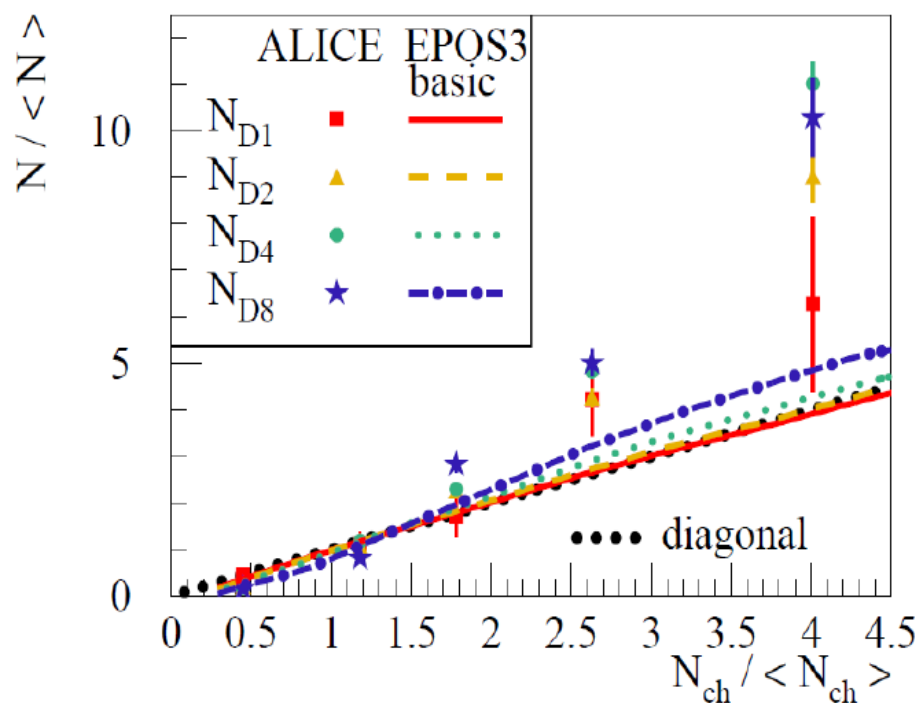
Effect on Heavy Flavor Production

Similar behavior observed in D meson but presented in a different way

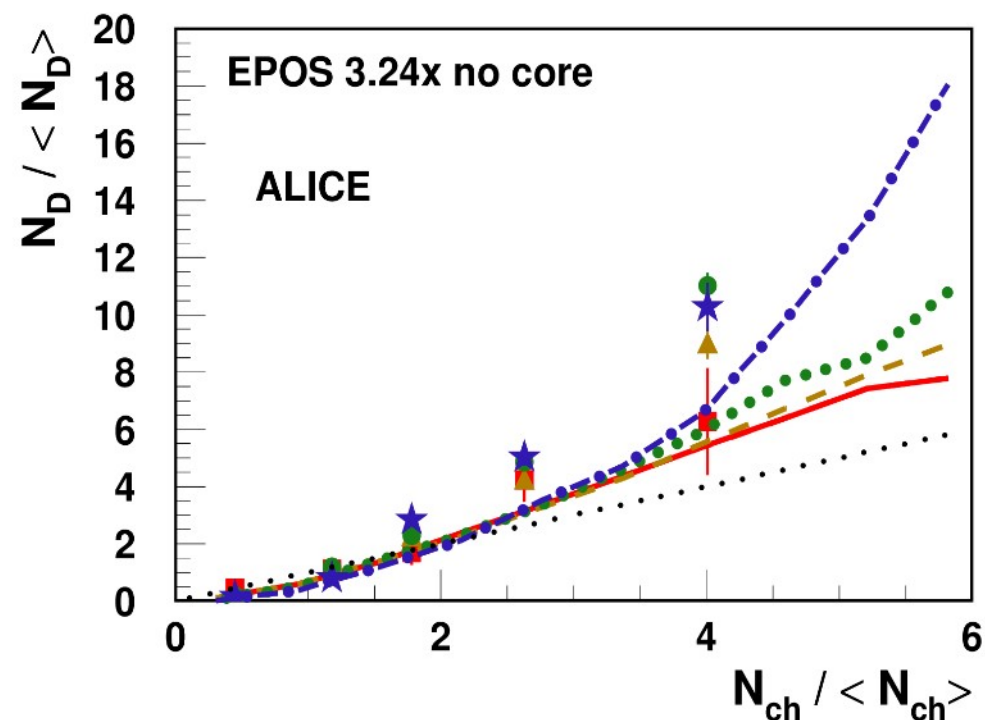
➔ increase of Q_0^2 with multiplicity imply a non linear increase of charm production as function of multiplicity

➔ strong effect but still not enough compared to data

➔ room for reduction of multiplicity due to collective effect (core)



Q_0^2 with $\langle N_{hard} \rangle$ dep.

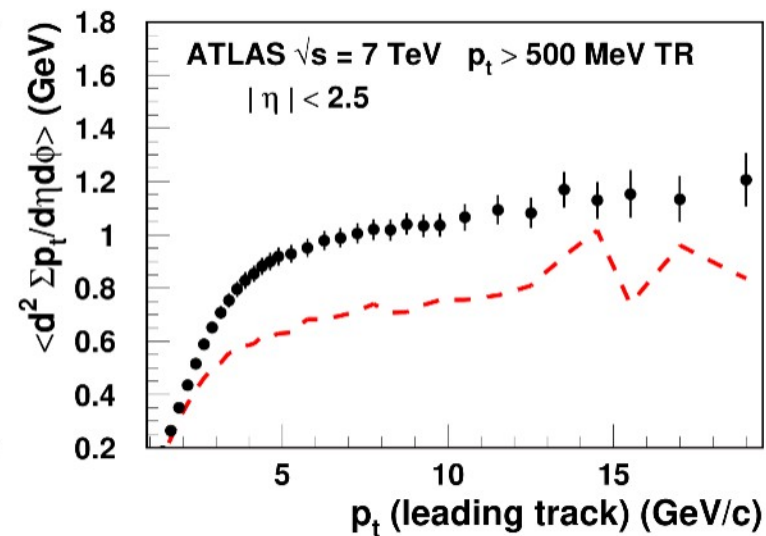
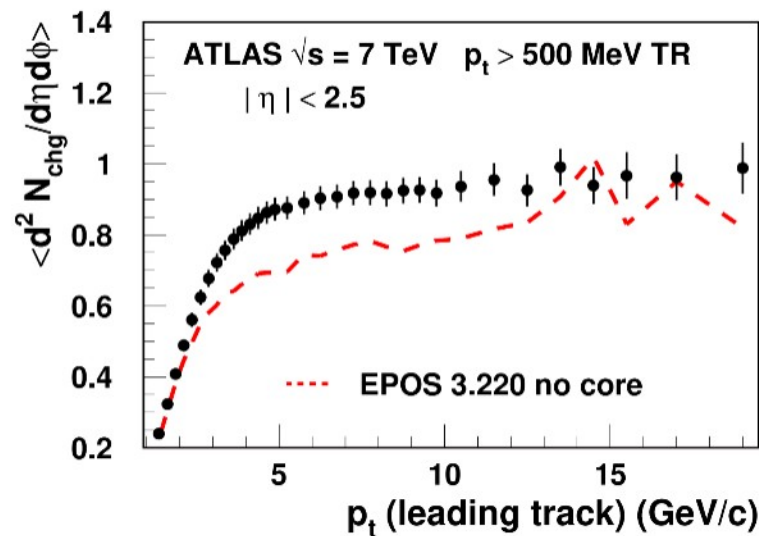
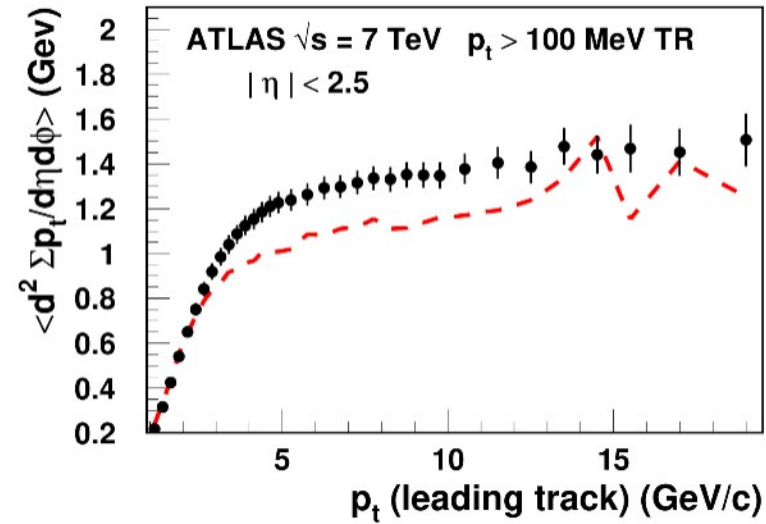
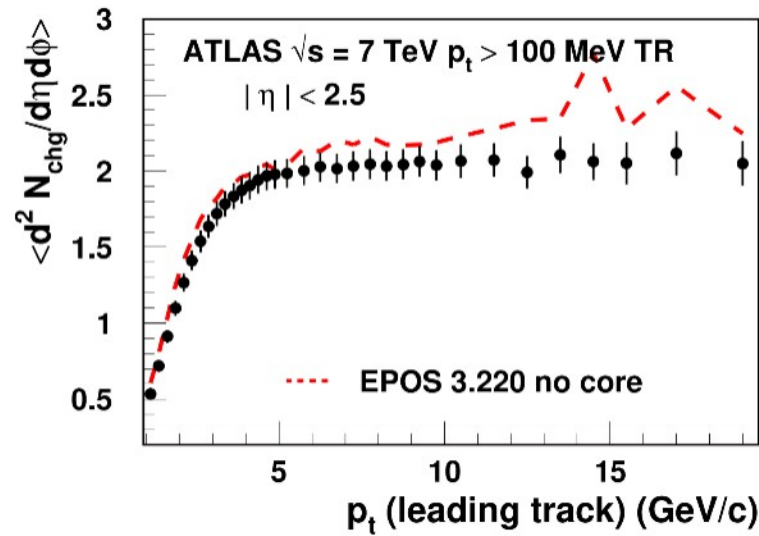


Q_0^2 with N_{hard} dep.

Underlying Events

Check with pp data at 7 TeV

➔ activity too high ?



Nuclear Interactions

Factorization holds independently of centrality

Once normalized by the number of binary collisions and inelastic cross-section, hard parton production (large Q^2) similar in pp or nuclear collisions.

EPOS 3

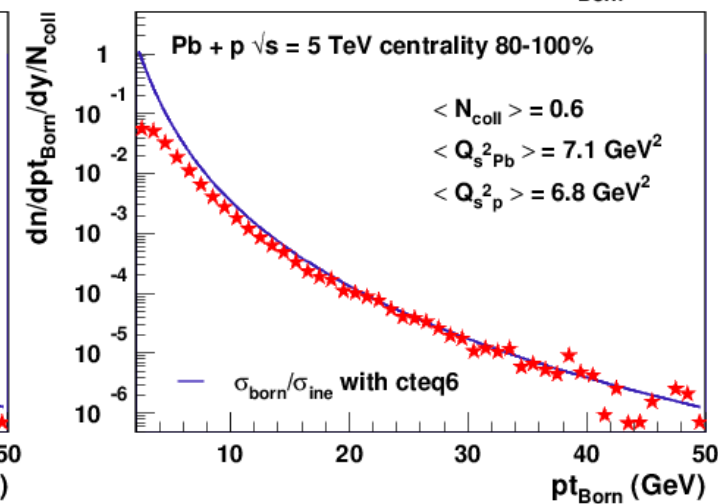
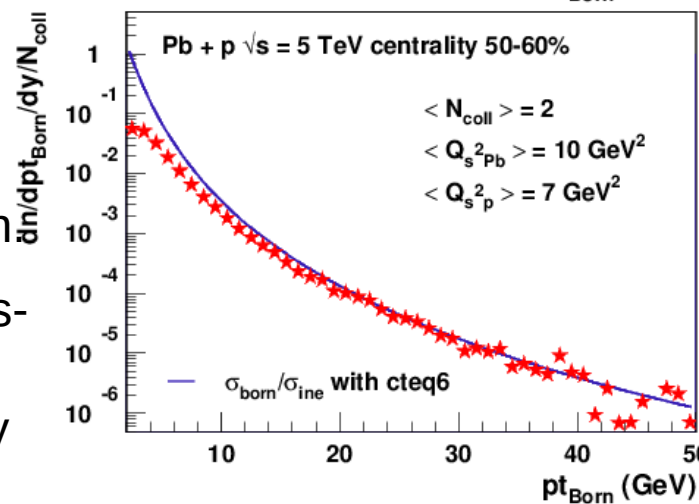
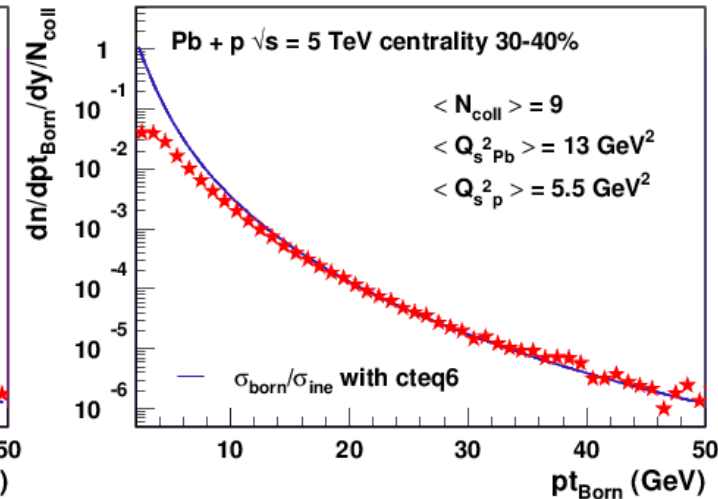
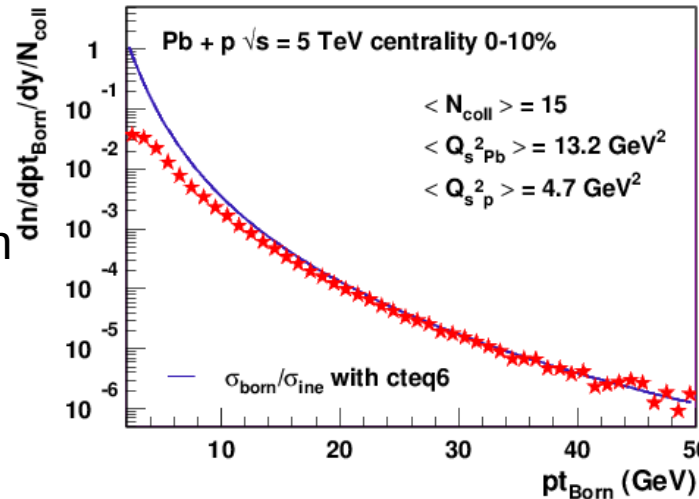
extend N_{hard} to take into account connections with other nucleons ($\sim N_{bin}$)

Define Q_0^2 such that

$$(\sum N_{hard}) G_{QCD}(x,b,Q_0^2) = G_{eff}(s,x,b,A)$$

to produce ISR and born process in hard Pomeron

Scaling of inclusive cross-section if N_{hard} and N_{soft} ($N_{pom} = N_{hard} + N_{soft}$) properly determined

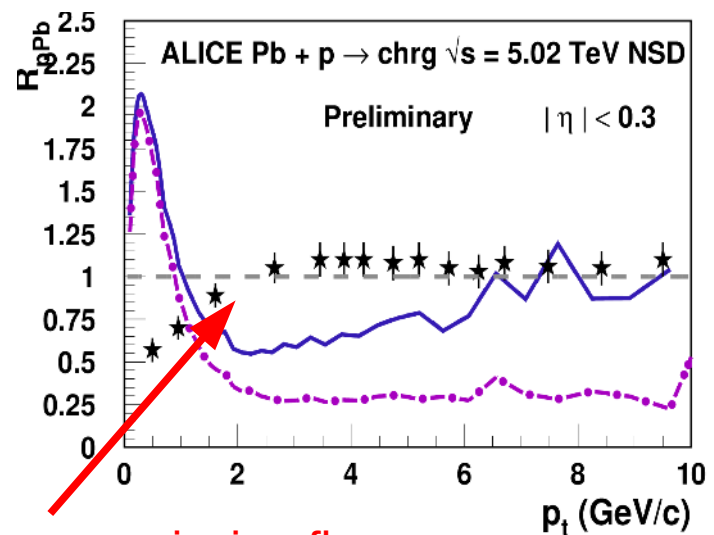
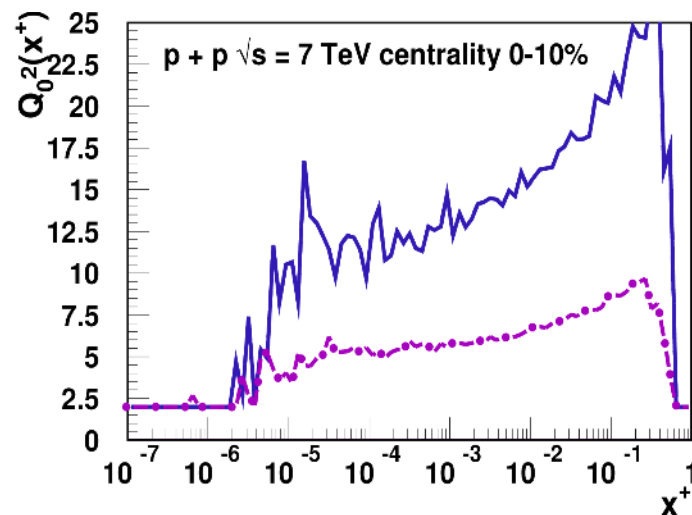
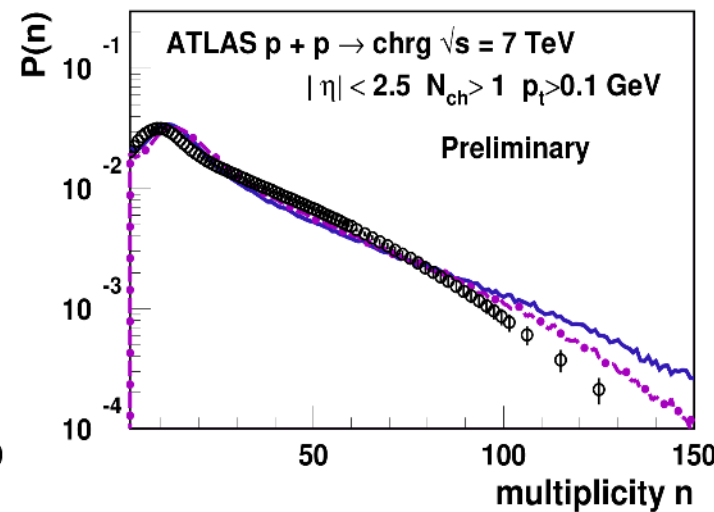
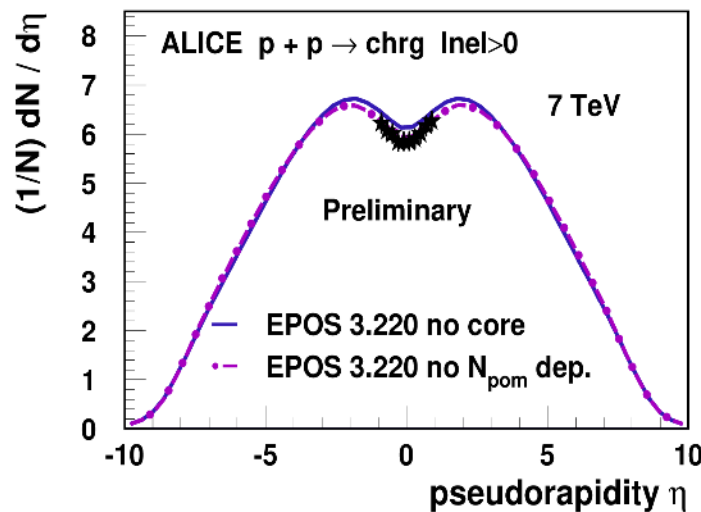


Preliminary Results : Without Core

- Overestimate multiplicity to take into account the effect of hydro
 - ➔ change in multiplicity by changing Q_0^2 definition only in the tail (as expected)

- Problem solved for hard processes
 - ➔ complete factorization
 - ➔ binary scaling for nuclear scattering simply by adding collision from all nucleons in N_{pom}

Same process to scale Q_0^2 in pp, pA and AA gives factorization and binary scaling.



no core = missing flow

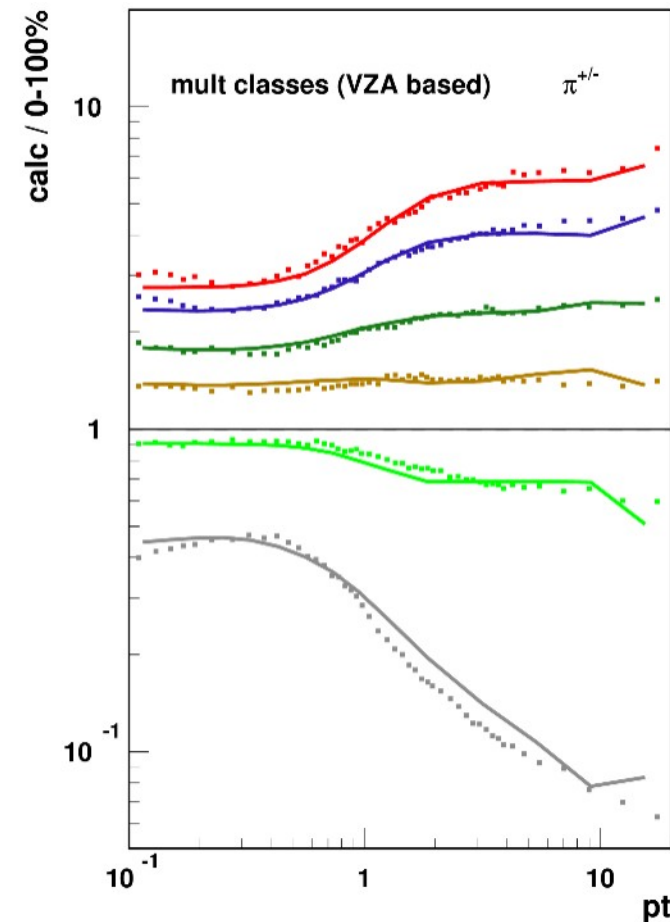
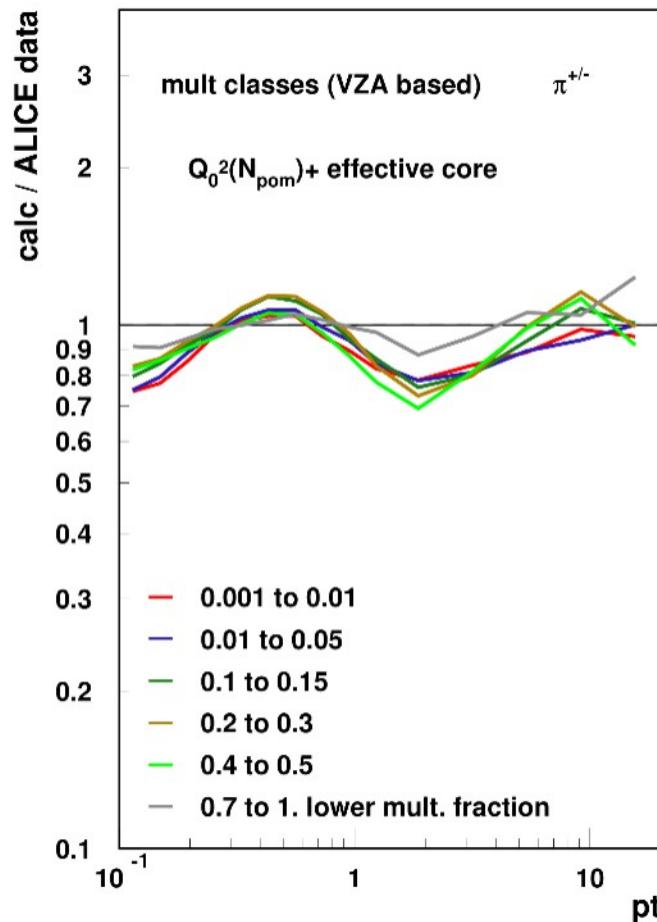
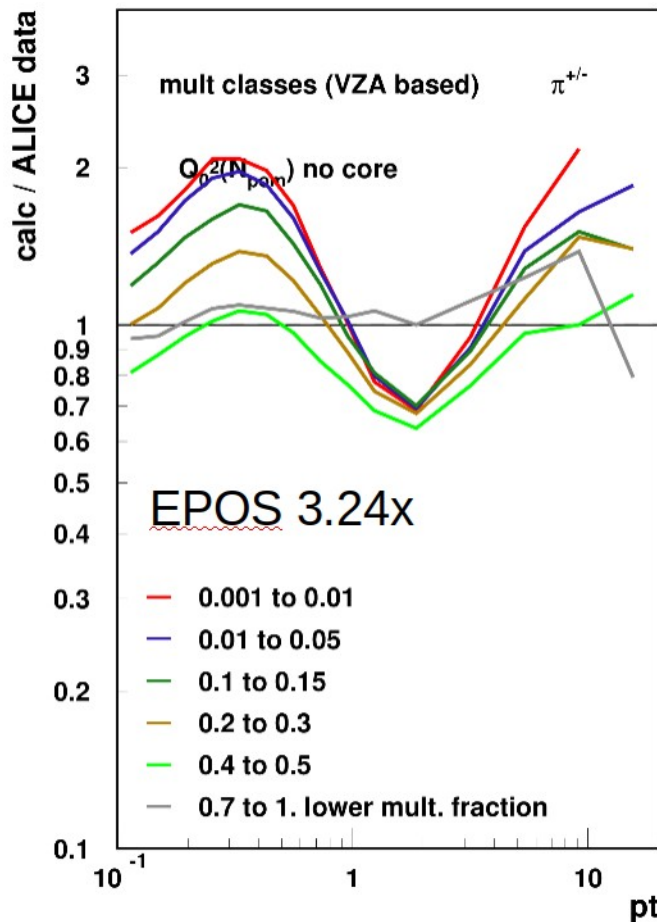
Core Effect vs Centrality

Flow change the shape of transverse momentum distribution

➔ effect on high p_t due to parton energy loss (same as in HI) ... to be confirmed !

➔ test of real hydro vs CGC ?

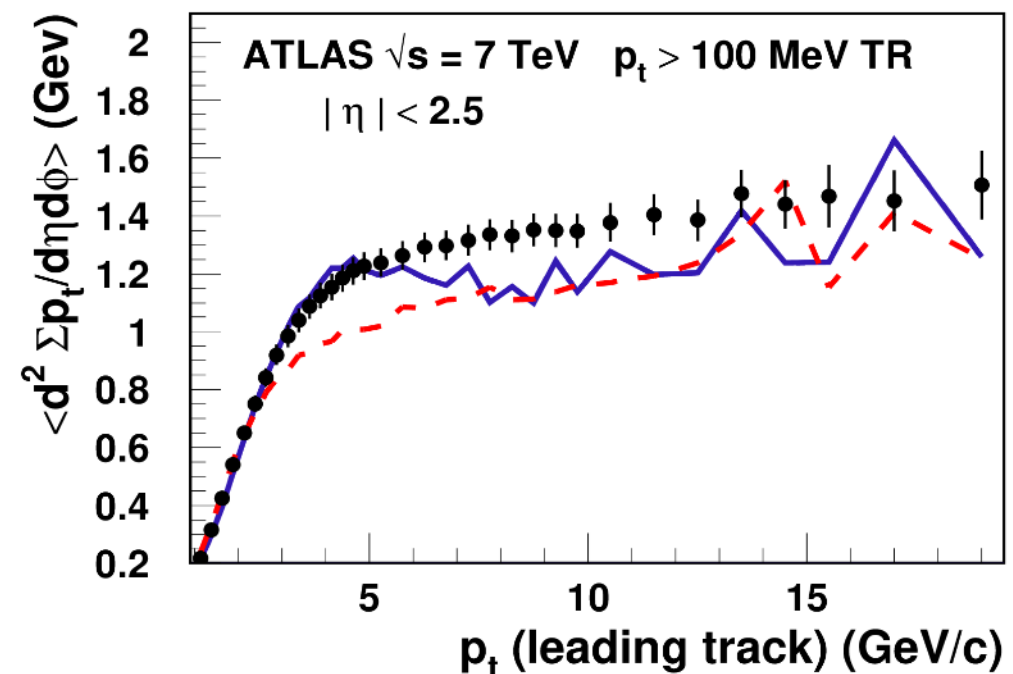
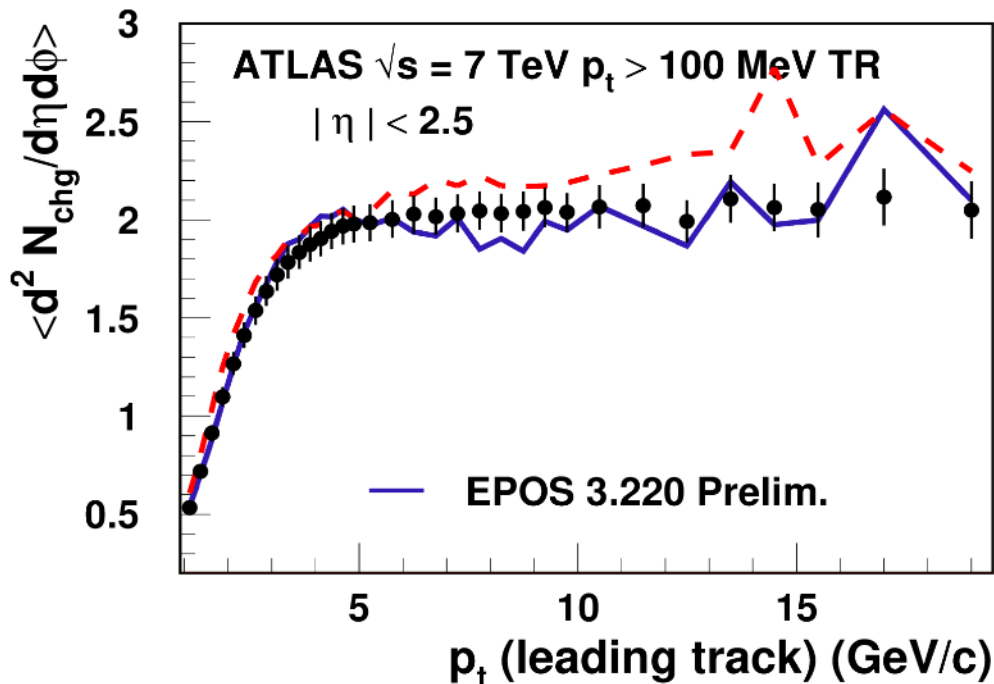
➔ (tension between these data and ATLAS/CMS (absolute scale))



Underlying Events: $p_t > 100 \text{ MeV}/c$

$p_t > 100 \text{ MeV}/c$ particles in TRANS region

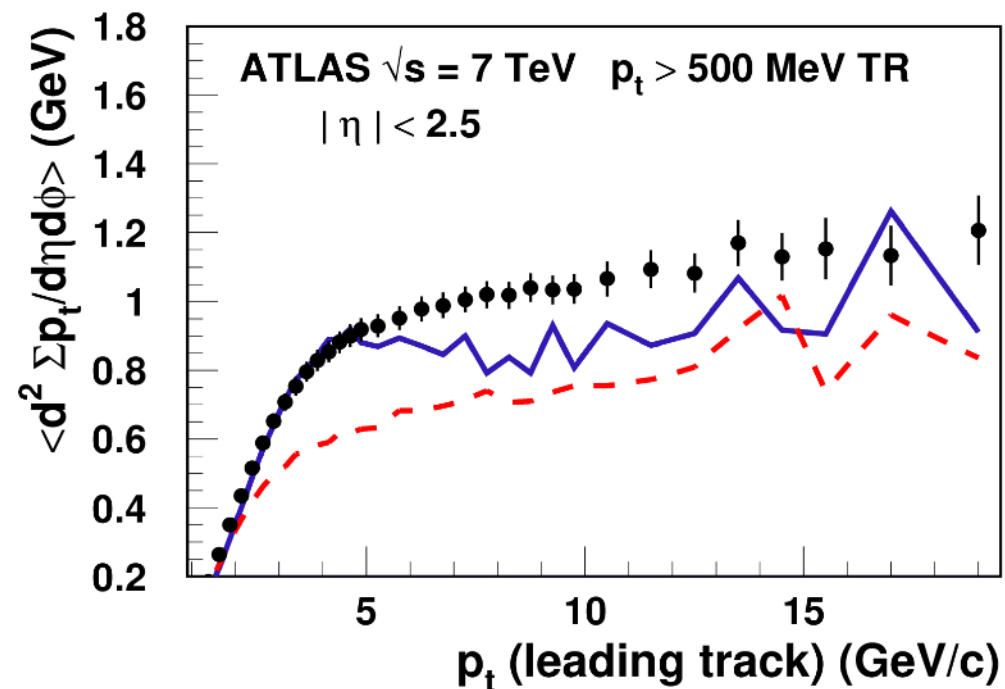
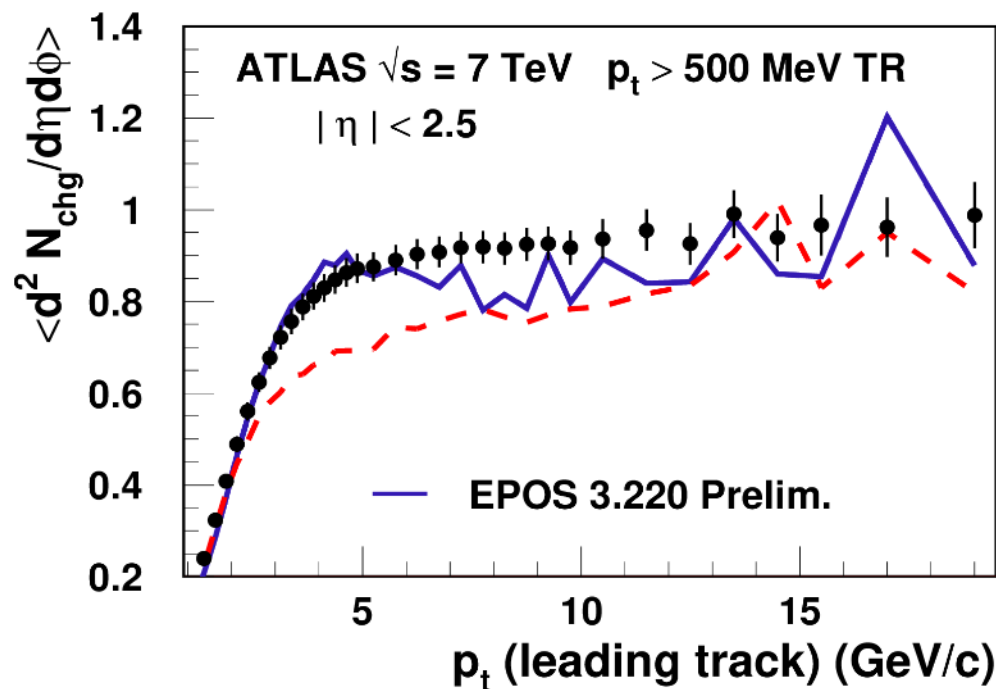
- ➔ without core N_{ch} is large like in MB but energy density is too low for p_t leading $> 2 \text{ GeV}/c$
- ➔ with core the multiplicity is reduced and energy density at intermediate p_t is increased
- ➔ reasonable agreement with data
 - ◆ mean transverse energy still a bit low for high p_t leading track



Underlying Events: $p_t > 500 \text{ MeV}/c$

$p_t > 500 \text{ MeV}/c$ particles in TRANS region

- ➔ without core N_{ch} is too low and energy density is too low
- ➔ with core here both multiplicity and energy density are increased at intermediate p_t
- ➔ reasonable agreement with data
 - ◆ mean transverse energy still a bit low for high p_t leading track
 - ➔ more study needed (sensitivity to FSR, flow orientation, ...)



Underlying Events: Strangeness

Lambda production in UE

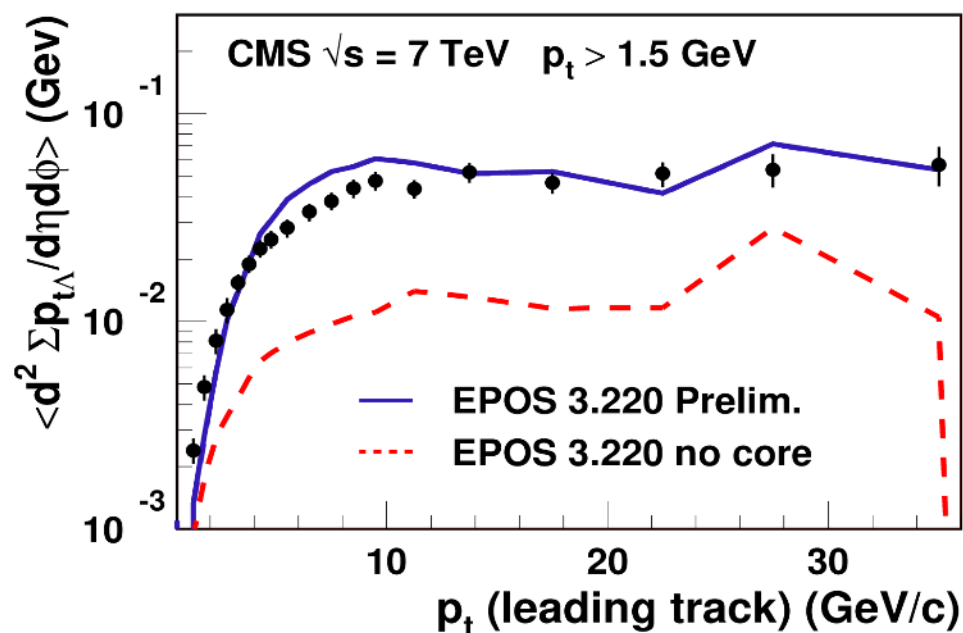
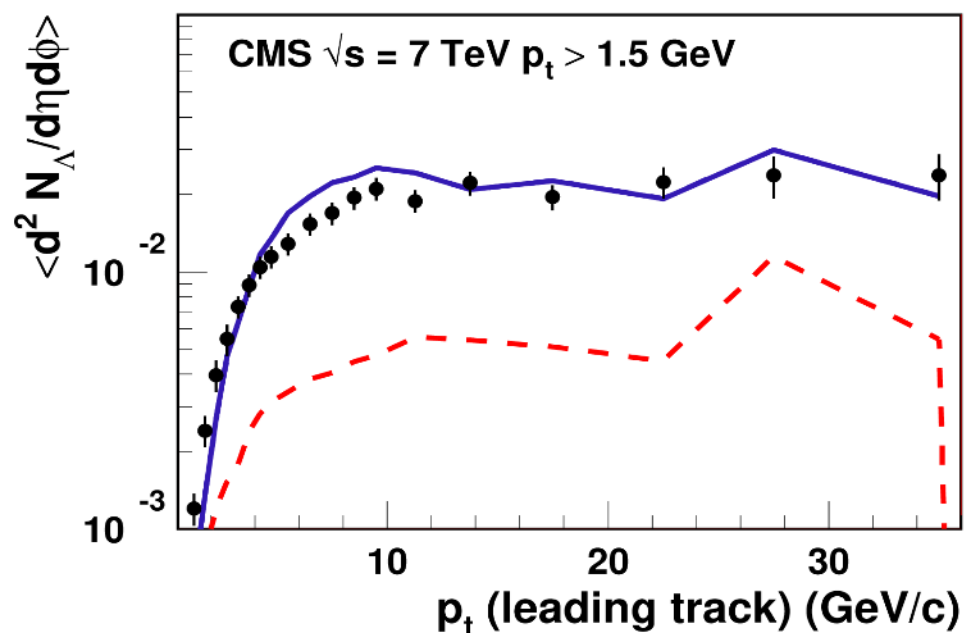
➔ Without core, very low lambda production like for other HEP models

➔ With core (and so hydro), much higher strangeness production

■ statistical hadronization

■ flow effect on transverse energy

➔ very strong effect of collective hadronization in UE for strange baryon production



Summary

To reconcile soft and hard observables, EPOS 3 will have **both collective effects and variable non-linear perturbative scale** : impact on HQ production.

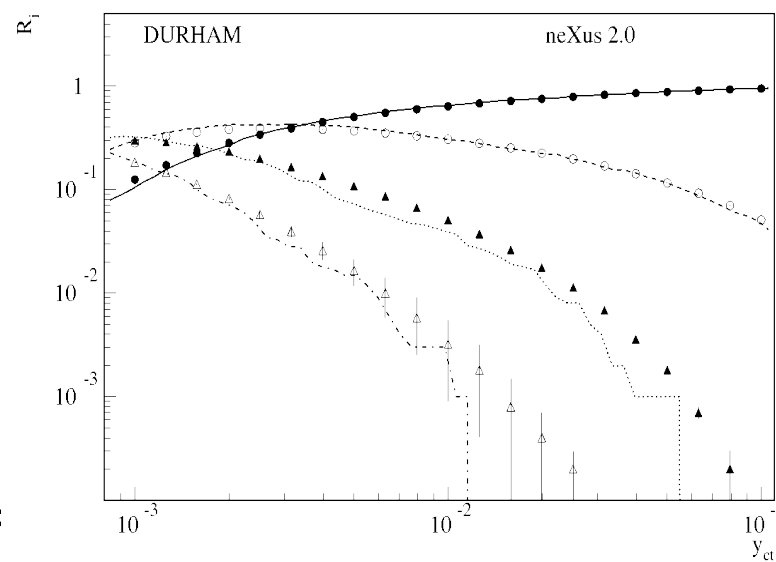
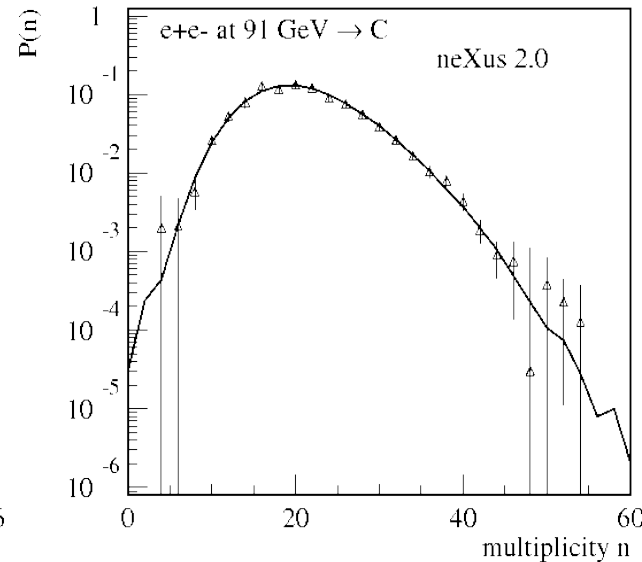
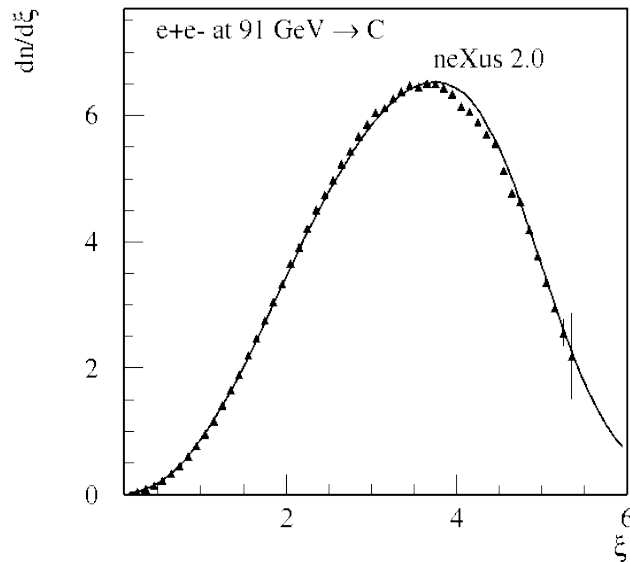
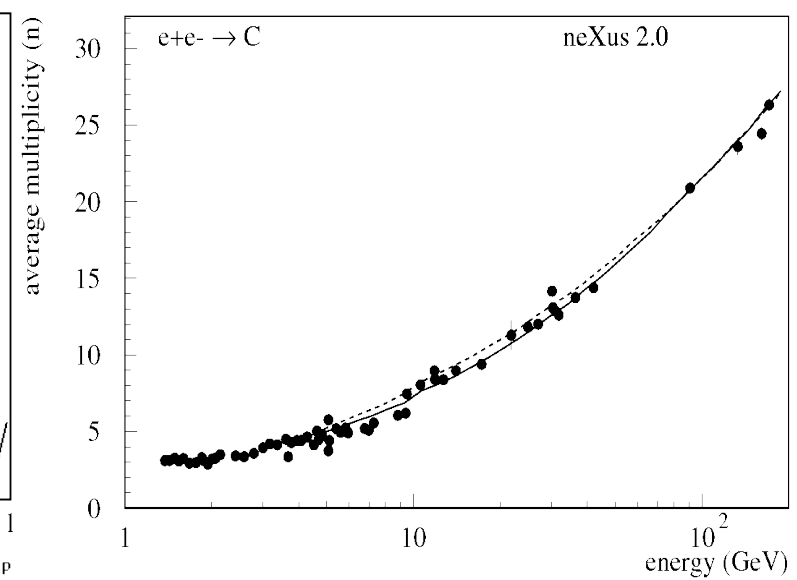
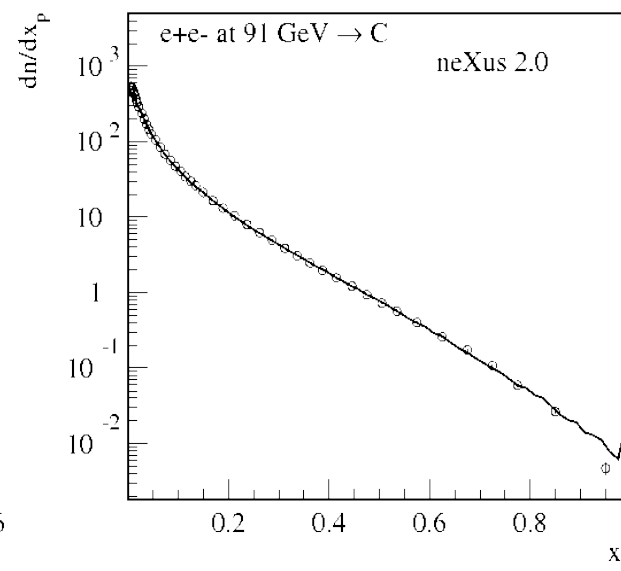
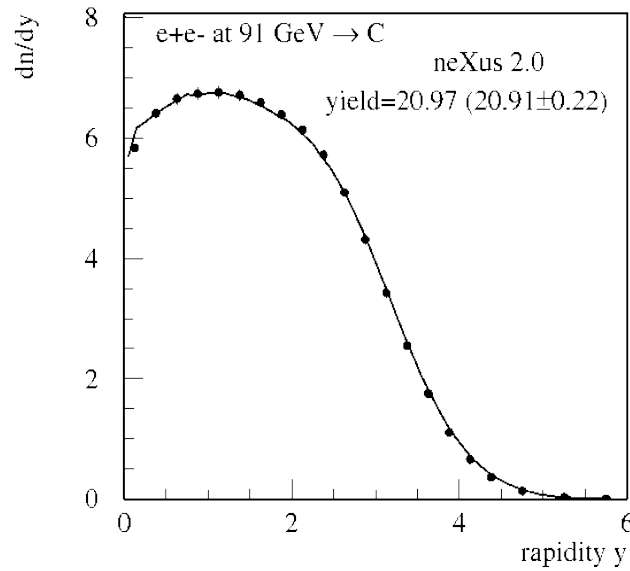
Difficult to describe min bias and hard scale events at the same time

- ➔ UE and heavy flavor require large multiplicity
- ➔ transverse momentum distributions depend on event multiplicity even at high p_t (to high for flow effect)

EPOS 3

- ➔ introduce (non-)linear perturbative scale Q_0^2 **GENERATED** Pomeron-by-Pomeron and dependent on the number of MPI event-by-event.
- ➔ non-linear N-pdf generation coupled with N independent DGLAP evolutions to get N hard partons event-by-event
- ➔ recover factorization and binary scaling for inclusive hard processes above Q_0^2
- ➔ hydro expansion require higher MPI than imposed by multiplicity that reflects on UE and other variables like charm production.
- ➔ improve underlying event description in p-p but real hydro still to be tried for final results

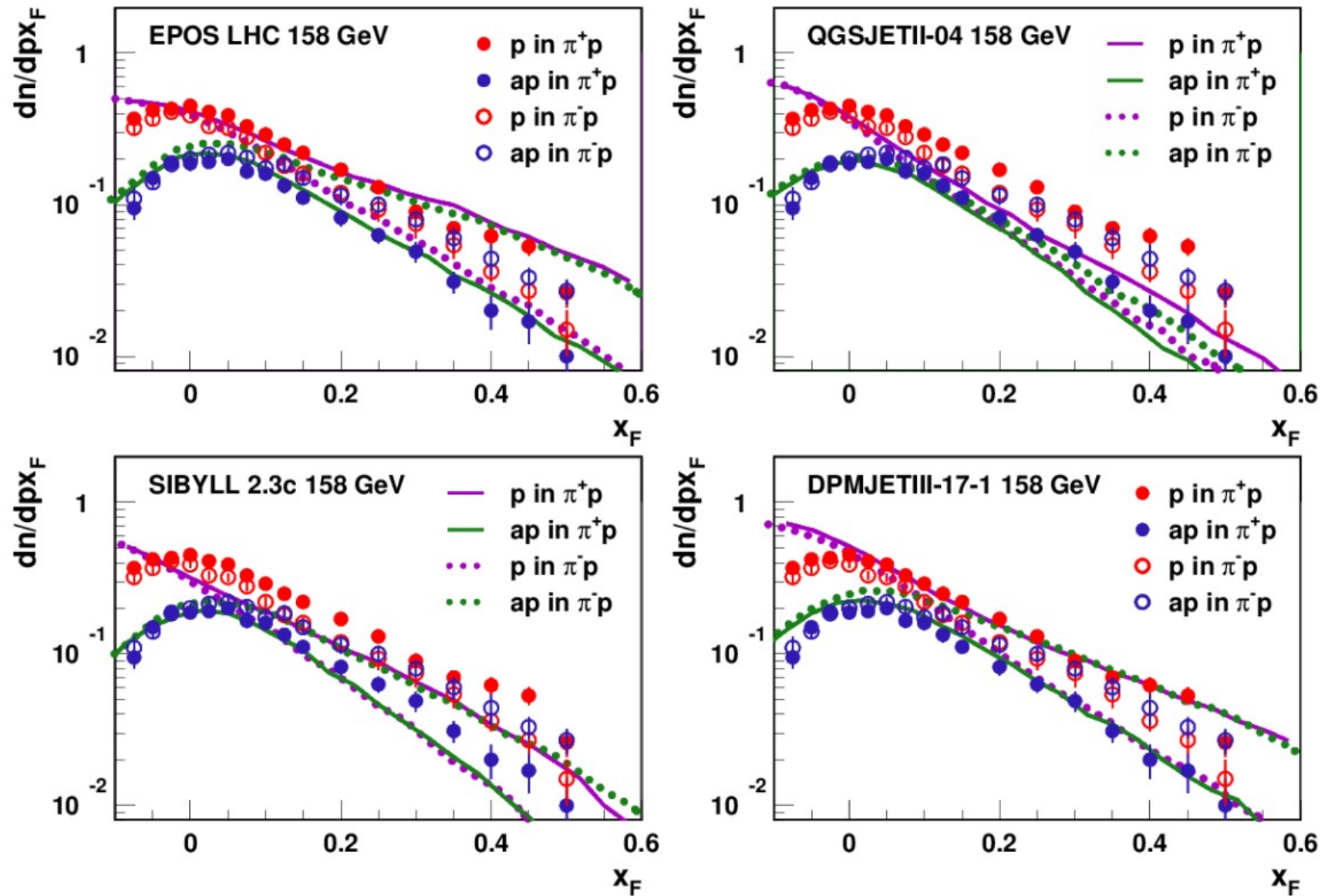
Test of string fragmentation with LEP data



Area law

Baryons in Pion Interactions

Data from NA49 (Gabor Veres PhD) : full picture



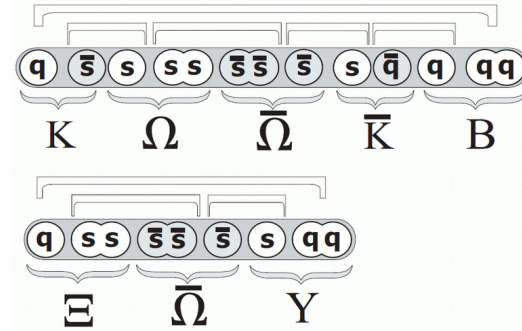
Baryons and Remnants

Parton ladder string ends :

➔ Problem of multi-strange baryons at low energy (Bleicher et al., Phys.Rev.Lett.88:202501,2002)

◆ 2 strings approach :

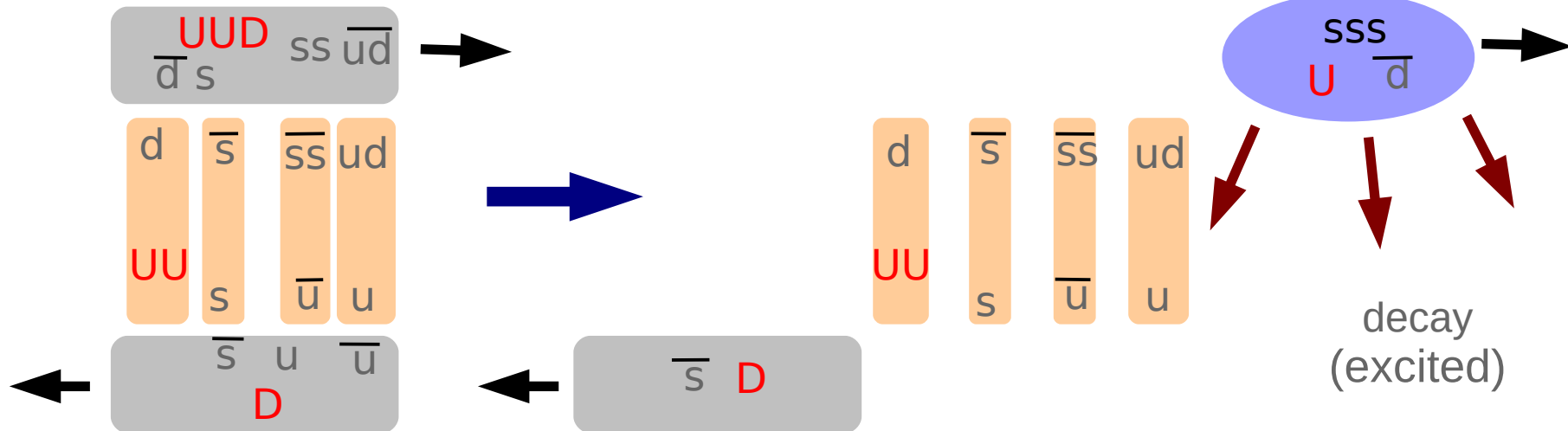
- ➔ $\bar{\Omega} / \Omega$ always > 1
- ➔ But data < 1 (Na49)



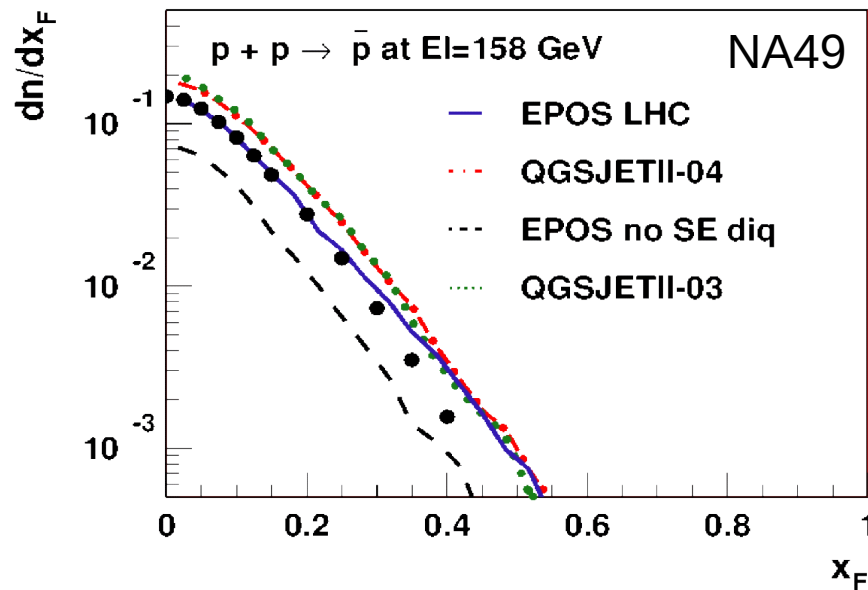
➔ EPOS

- ◆ No “first string” with valence quarks : all strings equivalent
- ◆ Wide range of excited remnants (hadronization via light resonance decay, string fragmentation or heavy quark-bag statistical decay)

➔ $\bar{\Omega} / \Omega$ always < 1

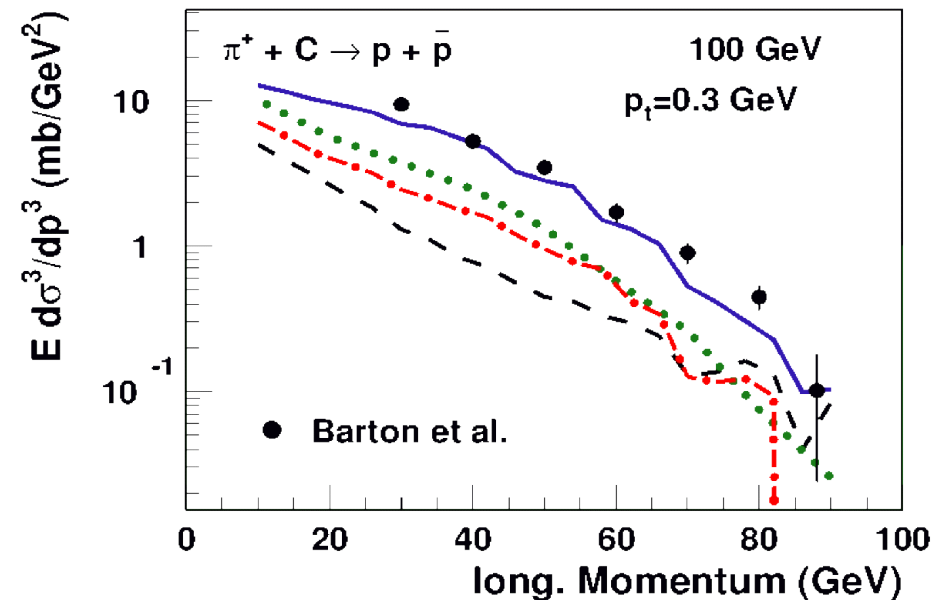
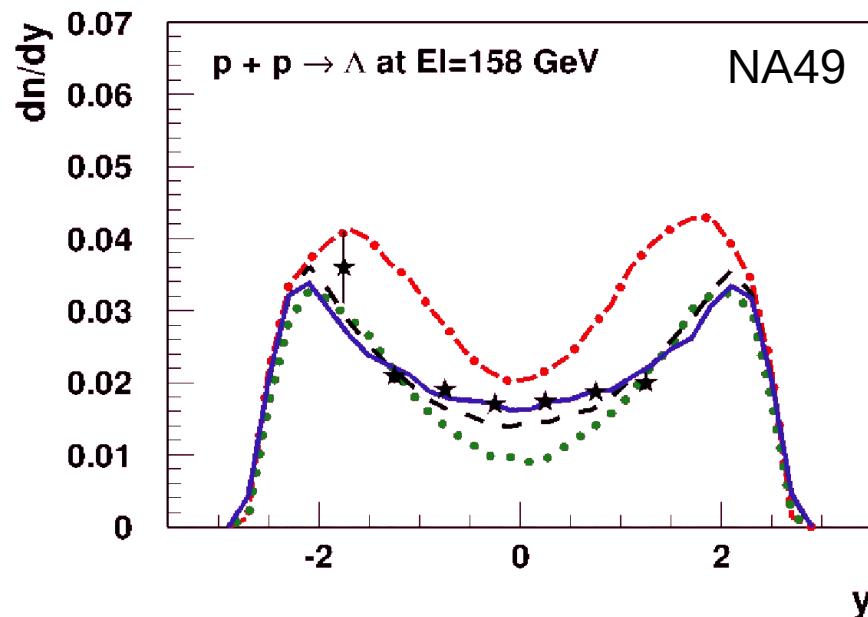


Forward Baryons (low energy)



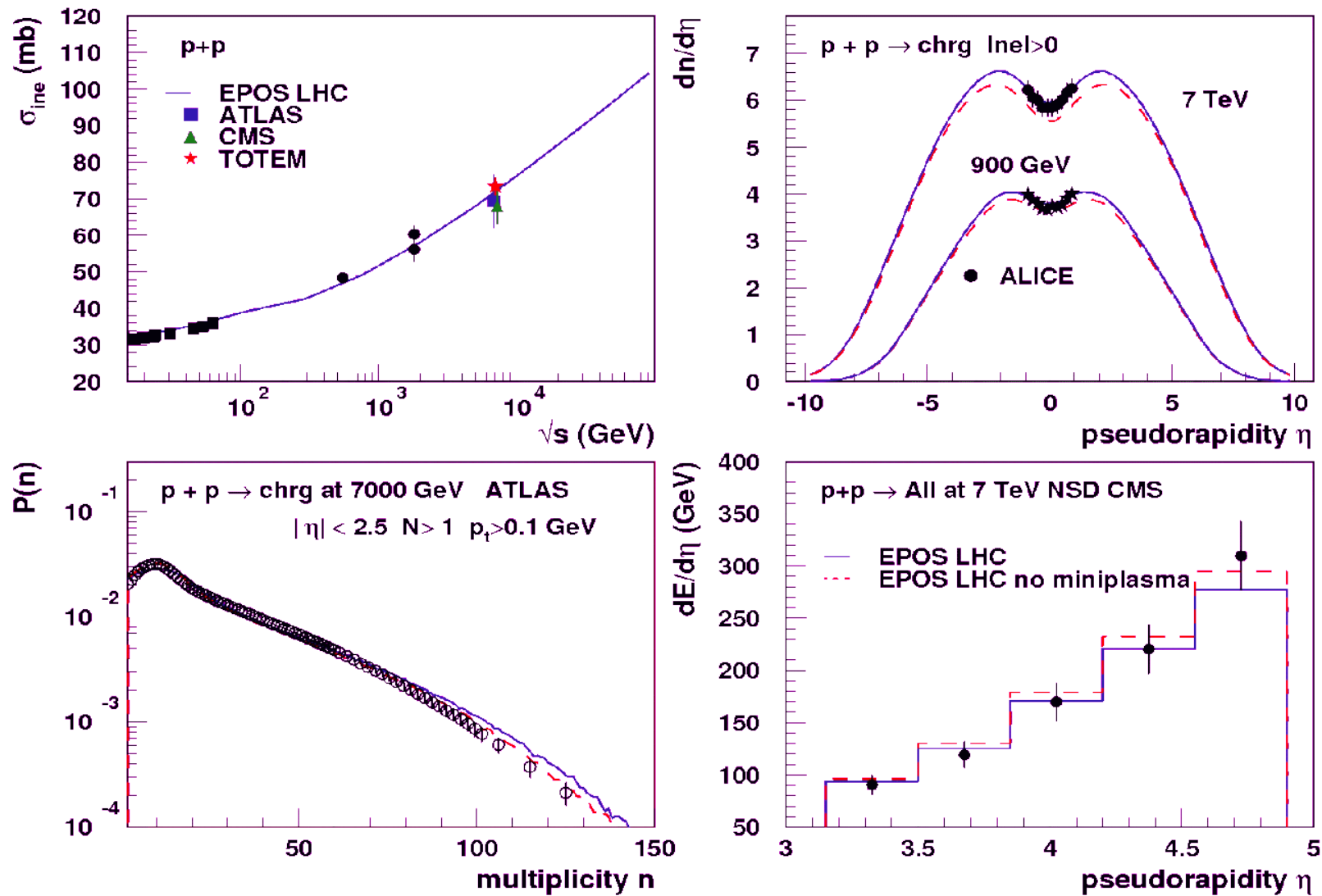
- ➔ Large differences between models
- ➔ Need a new remnant approach for a complete description (EPOS)
- ➔ Problems even at low energy
- ➔ No measurement at high energy !

Without remnant, string fragmentation has to be changed for baryon production



EPOS LHC

Effective flow treatment



EPOS LHC

Detailed description can be achieved

- ➔ identified spectra
- ➔ p_t behavior driven by collective effects (statistical hadronization + flow)

