Exceptional Field Theories & Applications

Henning Samtleben

MITP Workshop Holography, Generalized Geometry and Duality 05/2019





exceptional field theory

- dimensional reduction and duality symmetries
- exceptional geometry & tensor hierarchy
- generalized Scherk-Schwarz reductions

applications

- consistent truncations
 - > AdS₅ x S⁵
 - Sp, Hp,q spheres and hyperboloids
 - Sp x Hp,q products of spheres: dyonic gaugings
- supergravity magic
- timelike dualities
- generalized IIB supergravity

based on work with Olaf Hohm, Arnaud Baguet, Hadi Godazgar, Mahdi Godazgar, Hermann Nicolai, Edvard Musaev, Gianluca Inverso, Marc Magro, Emanuel Malek, Mario Trigiante, Valentí Vall Camell





ENS Lyon

exceptional field theory (ExFT)

manifestly duality covariant formulation of maximal supergravity

upon toroidal reduction on T^d, eleven-dimensional supergravity exhibits the global exceptional symmetry group E_{d(d)}

after proper dualisation/reorganisation of the fields



- ExFT : reformulate D=11 supergravity such that E_{d(d)} (or its remnants) becomes manifest <u>before</u> dimensional reduction
- **Example:** $E_{6(6)}$: exceptional field theory



D=5 maximal supergravity (torus reduction of D=11 or IIB)

after proper dualization of the dof's (different for IIA / IIB) the D=5 Lagrangian takes the $E_{6(6)}$ invariant form

$$\mathcal{L} = R + \frac{1}{24} \partial_{\mu} \mathcal{M}_{MN} \partial^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} F_{\mu\nu}{}^{M} F^{\mu\nu N} + \mathcal{L}_{top}$$

 $\begin{array}{l}g_{\mu\nu} : 5 \ge 5 \times 5 \text{ external metric}\\ \mathcal{A}_{\mu}{}^{M} : 27 \text{ vector fields} & \longleftrightarrow 27 \text{ two-form fields} \ \mathcal{B}_{\mu\nu}{}_{M}\\ \mathcal{M}_{MN} : 27 \ge 27 \text{ internal metric, parametrizing } \mathbb{E}_{6(6)}/\mathbb{U}\mathsf{Sp(8)}\end{array}$

with
$$\mathcal{L}_{top} = d_{KMN} F^M \wedge F^N \wedge A^K$$



D=5 maximal supergravity (torus reduction of D=11 or IIB)

after proper dualization of the dof's (different for IIA / IIB) the D=5 Lagrangian takes the $E_{6(6)}$ invariant form

$$\mathcal{L} = R + \frac{1}{24} \partial_{\mu} \mathcal{M}_{MN} \partial^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} F_{\mu\nu}{}^{M} F^{\mu\nu N} + \mathcal{L}_{top}$$

 $g_{\mu\nu}$: 5 x 5 external metric $\mathcal{A}_{\mu}{}^{M}$: 27 vector fields \longleftrightarrow 27 two-form fields $\mathcal{B}_{\mu\nu M}$ \mathcal{M}_{MN} : 27 x 27 internal metric, parametrizing $E_{6(6)}/USp(8)$

exceptional field theory:

- same Kaluza-Klein reorganisation of the higher-dimensional fields
- keeping the dependence on all internal coordinates (non-abelian gauge structure)



exceptional field theory: [Hohm, H.S.]

- same Kaluza-Klein reorganisation of the higher-dimensional fields
- keeping the dependence on all internal coordinates (non-abelian gauge structure)

$$\mathcal{L} = \widehat{R} + \frac{1}{24} \mathcal{D}_{\mu} \mathcal{M}_{MN} \mathcal{D}^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}_{\mu\nu}{}^{M} \mathcal{F}^{\mu\nu N} + \mathcal{L}_{top} - V_{pot}(\mathcal{M}_{MN}, g_{\mu\nu})$$

- internal coordinates enhanced to $E_{6(6)}$ representation: 6 -> 27 subject to the section constraint (covariant restriction down to 6 coordinates)

$$d^{KMN} \partial_M \otimes \partial_N = 0 \begin{cases} d^{KMN} \partial_M \partial_N f = 0 \\ d^{KMN} \partial_M f \partial_N g = 0 \end{cases} \begin{bmatrix} \text{Berman, Godazgar, Perry, West,} \\ \text{Coimbra, Strickland-Constable, Waldram,} \\ \text{Cederwall, Kleinschmidt, Thompson]} \end{cases}$$

Henning Samtleben

 $\mathcal{L} = R + \frac{1}{24} \partial_{\mu} \mathcal{M}_{MN} \partial^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} F_{\mu\nu}{}^{M} F^{\mu\nu N}$

 $+ \mathcal{L}_{top}$

$$\mathcal{L} = \widehat{R} + \frac{1}{24} \mathcal{D}_{\mu} \mathcal{M}_{MN} \mathcal{D}^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}_{\mu\nu}{}^{M} \mathcal{F}^{\mu\nu N} + \mathcal{L}_{top} - V_{pot}(\mathcal{M}_{MN}, g_{\mu\nu})$$

dependence on internal coordinates induces non-abelian gauge structure

 $\mathcal{D}_{\mu} = \partial_{\mu} - \mathcal{L}_{A_{\mu}} \qquad \qquad \mathcal{L}_{\Lambda} V^{M} = \Lambda^{N} \partial_{N} V^{M} - \kappa \left[\partial_{N} \Lambda^{M} \right]_{\text{adj}} V^{N}$ generalized diffeomorphisms
[Coimbra, Strickland-Constable, Waldram]

non-associativity of generalized diffeomorphisms induces modified YM field strengths

$$\mathcal{F}_{\mu\nu}{}^{M} \equiv 2\partial_{\left[\mu\right]}A_{\nu}{}^{M} - \left[A_{\mu}, A_{\nu}\right]_{\mathrm{E}}^{M} + 10\,d^{MNK}\partial_{K}B_{\mu\nu N}$$

two-form field equations compensated by the topological term

$$S_{\text{top}} = \int d^{27}Y \int_{\mathcal{M}_6} \left(d_{MNK} \mathcal{F}^M \wedge \mathcal{F}^N \wedge \mathcal{F}^K - 40 \, d^{MNK} \mathcal{H}_M \wedge \partial_N \mathcal{H}_K \right)$$

boundary term of a six-dimensional bulk



ENS Lyon

E₆₍₆₎ : exceptional field theory

$$\mathcal{L} = \widehat{R} + \frac{1}{24} \mathcal{D}_{\mu} \mathcal{M}_{MN} \mathcal{D}^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}_{\mu\nu}{}^{M} \mathcal{F}^{\mu\nu N} + \mathcal{L}_{top} - V_{pot}(\mathcal{M}_{MN}, g_{\mu\nu})$$

"potential"

$$V_{\text{pot}} = \frac{1}{24} \mathcal{M}^{MN} \partial_M \mathcal{M}^{KL} \left(12 \partial_L \mathcal{M}_{NK} - \partial_N \mathcal{M}_{KL} \right) - \frac{1}{2} g^{-1} \partial_M g \partial_N \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}^{MN} g^{-1} \partial_M g g^{-1} \partial_N g - \frac{1}{4} \mathcal{M}^{MN} \partial_M g^{\mu\nu} \partial_N g_{\mu\nu}$$

- invariant under generalized diffeomorphisms

- generalised (internal) curvature scalar



E₆₍₆₎ : exceptional field theory

$$\mathcal{L} = \widehat{R} + \frac{1}{24} \mathcal{D}_{\mu} \mathcal{M}_{MN} \mathcal{D}^{\mu} \mathcal{M}^{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}_{\mu\nu}{}^{M} \mathcal{F}^{\mu\nu N} + \mathcal{L}_{top} - V_{pot}(\mathcal{M}_{MN}, g_{\mu\nu})$$

- unique action with generalized diffeomorphism invariance in all 5+27 coordinates (modulo section condition)
- section condition is solved upon breaking $E_{6(6)} \rightarrow SL(5)$ $d^{KMN} \partial_M \otimes \partial_N = 0$ $27 \longrightarrow 1 + 5' + 10 + 5 + 5' + 1$

5' : reproduces full IIA supergravity (additional singlet 1:
$$D=11$$
 supergravity)

5 : reproduces full IIB supergravity (preserves an SL(2))

together with proper dictionary of ExFT fields into IIA/IIB







- upon solving the section constraint: reformulation of the original theories
- IIA and IIB supergravity accommodated in the same framework





makes the symmetry enhancement after torus reduction manifest



ENS Lyon



also allows a compact description of complicated reductions



ENS Lyon



also allows a compact description of complicated reductions



ENS Lyon

consistent truncation on AdS₅ x S⁵





- AdS₅ x S⁵ : maximal supersymmetric solution of IIB
- fluctuations around the background: D=5 gauged supergravity
- explicit reduction formulas: highly non-trivial

$$ds^{2} = \Delta^{-2/3}(x,y) g_{\mu\nu}(x) dx^{\mu} dx^{\nu} + G_{mn}(x,y) \left(dy^{m} + \mathcal{K}_{[ab]}{}^{m}(y) A_{\mu}^{ab}(x) dx^{\mu} \right) \left(dy^{n} + \mathcal{K}_{[cd]}{}^{n}(y) A_{\nu}^{cd}(x) dx^{\nu} \right)$$
$$G^{mn}(x,y) = \Delta^{2/3}(x,y) \mathcal{K}_{[ab]}{}^{m}(y) \mathcal{K}_{[cd]}{}^{n}(y) M^{ab,cd}(x)$$

 $\begin{array}{ll} D=11 \mbox{ on } AdS_4 \ge S^7: \mbox{ [de Wit, Nicolai] 1987} \\ D=11 \mbox{ on } AdS_7 \ge S^4: \mbox{ [Nastase, van Nieuwenhuizen, Vaman] 1999} \\ \mbox{ IIB on } AdS_5 \ge S^5: \mbox{ shown for various sub-sectors...} \end{array}$





consistent truncation via generalized Scherk-Schwarz ansatz in ExFT

in terms of an E₆₍₆₎-valued twist matrix $U_M{}^N(Y)$ and scale factor ho(Y)





- consistent truncation via generalized Scherk-Schwarz ansatz in ExFT in terms of an E₆₍₆₎—valued twist matrix $U_M{}^N(Y)$ and scale factor $\rho(Y)$
- **system of consistency equations** $[(U^{-1})_M{}^P(U^{-1})_N{}^L \partial_P U_L{}^K]_{351} \stackrel{!}{=} \rho X_{MN}{}^K$
- generalized parallelizability
- no general classification of its solutions (Lie algebras vs Leibniz algebras)





twist matrix $U \in SL(6)$ associated to SO(6)

full reduction formulas of IIB on AdS₅ x S⁵

background AdS₅ x S⁵

 $U = \begin{pmatrix} g^{-1/2} \partial_i \mathcal{Y}^A \\ \mathcal{Y}^A - 2\zeta^i \partial_i \mathcal{Y}^A \end{pmatrix} \in \mathrm{SL}(6)$

in terms of sphere harmonics and the fields of D=5 maximal supergravity



ENS Lyon

e.g. metric (standard Kaluza-Klein form)

$$ds^{2} = \Delta^{-2/3}(x,y) g_{\mu\nu}(x) dx^{\mu} dx^{\nu} + G_{mn}(x,y) \left(dy^{m} + \mathcal{K}_{[ab]}{}^{m}(y) A_{\mu}^{ab}(x) dx^{\mu} \right) \left(dy^{n} + \mathcal{K}_{[cd]}{}^{n}(y) A_{\nu}^{cd}(x) dx^{\nu} \right)$$
$$G^{mn}(x,y) = \Delta^{2/3}(x,y) \mathcal{K}_{[ab]}{}^{m}(y) \mathcal{K}_{[cd]}{}^{n}(y) M^{ab,cd}(x)$$

e.g. 4-form (after reconstructing all components, in Kaluza-Klein basis)

proves the consistent truncation of IIB on AdS₅ x S⁵



 $+\frac{1}{40}\sqrt{2}A_{[\mu}{}^{ab}A_{\nu}{}^{cd}A_{\rho}{}^{ef}A_{\sigma}{}^{gh}A_{\tau]}{}^{ij}\varepsilon_{abcegi}\eta_{df}\eta_{hj}$



twist matrix $U \in SL(6)$ associated to SO(6)

- background AdS₅ x S⁵
- full reduction formulas of IIB on AdS₅ x S⁵

 $U = \begin{pmatrix} g^{-1/2} \partial_i \mathcal{Y}^A \\ \mathcal{Y}^A - 2 \zeta^i \partial_i \mathcal{Y}^A \end{pmatrix} \in \mathrm{SL}(6)$

ENS Lyon



proves the consistent truncation of IIB on AdS₅ x S⁵



similar: twist matrix $U \in SL(6)$ associated to SO(p,q) and CSO(p,q,r) built from sphere harmonics on SO(p,q)/SO(p-1,q)

$$U = \begin{pmatrix} g^{-1/2} \partial_i \mathcal{Y}^A \\ \mathcal{Y}^A - 2 \zeta^i \partial_i \mathcal{Y}^A \end{pmatrix} \in \mathrm{SL}(6)$$



hyperboloid compactifications



SO(p,q) and CSO(p,q,r)

built from sphere harmonics on SO(p,q)/SO(p-1,q)



hyperboloid compactifications



similar: twist matrix $U \in SL(6)$ associated to SO(p,q) and CSO(p,q,r)

- background: (warped) hyperboloids [Hull, Warner] [Baron, Dall'Agata]
- in general no IIB solutions, still consistent truncations!





ENS Lyon

Consistent truncations with smaller isometry groups [Inverso, HS, Trigiante, Malek] products of spheres and hyperboloids $S^p \times S^q$, $S^p \times H^q$ specific D=4 construction, based on electric/magnetic split of internal coordinates inducing dyonic gaugings $(SO(p,q) \times SO(p',q')) \ltimes N$ [Dall'Agata, Inverso]





consistent truncations with smaller isometry groups [Inverso, HS, Trigiante, Malek] products of spheres and hyperboloids $S^p \times S^q$, $S^p \times H^q$ specific D=4 construction, based on electric/magnetic split of internal coordinates inducing dyonic gaugings $(SO(p,q) \times SO(p',q')) \ltimes N$ [Dall'Agata, Inverso]



- consistent truncations with less supersymmetry via DFT [Baguet, Malek, Pope, HS, Sarioglu]
 - > S³ reduction of the bosonic string
 - > more general: Pauli reduction of the bosonic string on group manifold G
 - > example G = SO*(4), type II uplift of D=4 Minkowski vacua
 - > $AdS_3 \times S^3$ reductions from 6D supergravity, N=(1,1) and N=(2,0) w tensor-multiplets



- consistent truncations with less supersymmetry via DFT [Baguet, Malek, Pope, HS, Sarioglu]
 - > S³ reduction of the bosonic string
 - > more general: Pauli reduction of the bosonic string on group manifold G
 - > example G = SO*(4), type II uplift of D=4 Minkowski vacua
 - > $AdS_3 \times S^3$ reductions from 6D supergravity, N=(1,1) and N=(2,0) w tensor-multiplets
- consistent truncations with less supersymmetry in ExFT (in type II sugra) [Malek] [Malek, HS, Vall Camell] embedding of half-maximal supergravity into ExFT construction and classification of supersymmetric AdS vacua half-maximal supersymmetric AdS vacua induce consistent truncations around

 \rightarrow [Emanuel's talk]



- ExFT for all finite-dimensional exceptional groups $E_{d(d)}$, d<9 [Hohm, HS] [Abzalov, Bakhmatov, Musaev, Hohm, Wang, Berman, Blair, Malek, Rudolph] based on the different splits external/internal coordinates $\{x^{\mu}, y^{m}\} \longrightarrow \{x^{\mu}, Y^{M}\}$
- ExFT embedding of massive IIA theory

[Ciceri, Guarino, Inverso] [Cassani, de Felice, Petrini, Strickland-Constable, Waldram]

- by deformations of ExFT
- by Scherk-Schwarz reduction violating the section conditions
- -> more general theme: consistent theories from reductions violating section constraints
- ExFT embedding of 'generalized IIB' theory

[Baguet, Magro, HS]

- background from $\eta\text{-}deformed\ \mathrm{AdS}_5\ \mathrm{x}\ \mathrm{S}^5$ sigma model
- T-dual of IIA with non-isometric dilaton



magical framework

			I					hi	gher-d	imens	sior	nal orig	gin	of	3D	coset spa	aces
	11	×				[Cremmer Julia I 1 Pone]											
	10	$\mathbb{R} \mid A_1 \mid \times$															
9		$\ A_1 \times \mathbb{R}$		\mathbb{R}												LKeurent	jes
	8	$A_2 \times A_1$		$A_1 \times \mathbb{R}$	A ₂	A ₁											
	7	A ₄		$A_2 \times$	\mathbb{R}	$A_1 \times \mathbb{R}$	\mathbb{R}		×								
	6	D ₅		$A_3 \times A_3$	$A_1 A_1^2 \times \mathbb{R}$		\mathbb{R}^2	A_1^2	\mathbb{R}								
	5	5 E ₆		A ₅		A_2^2	A_1^2	$ imes \mathbb{R}$	$A_1 \times \mathbb{R}$	A ₁							
	4	1	E ₇	D ₆		A ₅		$\times A_1$	$A_2 \times \mathbb{R}$	$A_1 \times \mathbb{R}$	A ₂	\mathbb{R}	×				
	3	3 E ₈		E ₇		E ₆	D ₅		A ₄	$A_2 \times A_1$		$A_1 \times \mathbb{R}$	A ₁	\mathbb{R}	×		
	Dr	8		7	6		5		4	3		2	1		0		
				1		I	I		1	1		1	I				
	Il embedded within Equa EVET: $E_{\alpha(n)} \longrightarrow SI(D-2) \times SI(9-r) \times U_{D}$																
$\mathbb{L}_{\delta(\delta)} = \mathbb{L}_{\delta(\delta)} = $																	
											K						

solution of section constraint:

$$\left\{ \begin{array}{c} \mathbf{Y}^{M} \end{array} \right\} \longrightarrow \left\{ \underbrace{Y^{i}_{j}}_{j}, \underbrace{Y^{a}_{b}}_{b} \right\} \longrightarrow \left\{ y^{i} \equiv Y^{i}_{0} \right\}$$

SL(D-2) SL(9-r)

together with truncation to SL(9-r) singlets



magical framework

п т			hig	gher-d	imensior	nal orig	gin of	3D	coset space	
×						[Cre	mmer	TI	ilia Lu Pone	
$\mathbb{R} \mid A_1$	×									
$A_1 \times \mathbb{R}$	\mathbb{R}	[Keurentjes								
$A_2 \times A_1$	$A_1 \times \mathbb{R} \mid A_2$	A ₁								
A ₄	$A_2 \times \mathbb{R}$	$A_1 \times \mathbb{R}$	\mathbb{R}	×						
D ₅	$A_3 \times A_1$	$A_1^2 \times \mathbb{R}$	\mathbb{R}^2 A_1^2	\mathbb{R}						
E ₆	A ₅	A ₂ ²	$A_1^2 \times \mathbb{R}$	$A_1 \times \mathbb{R}$	A ₁					
E ₇	D ₆	A ₅	$A_3 \times A_1$	$A_2 \times \mathbb{R}$	$A_1 \times \mathbb{R} \mid A_2$	\mathbb{R}	×			
E ₈	E ₇	E ₆	D ₅	A ₄	$A_2 \times A_1$	$A_1 \times \mathbb{R}$	$ A_1 \mathbb{R}$	×		
8	7	6	5	4	3	2	1	0		
	$ \begin{vmatrix} \times \\ \mathbb{R} & A_1 \\ A_1 \times \mathbb{R} \\ A_2 \times A_1 \\ A_4 \\ D_5 \\ E_6 \\ E_7 \\ E_8 \\ 8 \\ \end{vmatrix} $	$ \begin{vmatrix} \times \\ \mathbb{R} & A_1 \\ A_1 \times \mathbb{R} \\ A_2 \times A_1 \\ A_2 \times A_1 \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ D_5 \\ A_3 \times A_1 \\ E_6 \\ E_7 \\ E_7 \\ B_6 \\ E_8 \\ E_7 \\ B_7 \\ A_5 \\ E_7 \\ A_5 \\ A_5 \\ E_7 \\ A_7 \\$	$ \begin{vmatrix} \times \\ \mathbb{R} & A_1 \\ A_1 \times \mathbb{R} \\ A_2 \times A_1 \\ A_4 \\ A_2 \times \mathbb{R} \\ A_4 \\ A_2 \times \mathbb{R} \\ A_4 \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_2 \times \mathbb{R} \\ A_1 \times \mathbb{R} \\ A_2 \times$	$\begin{array}{c c c c c c c c c c c c c c c c c c c $	× \mathbb{R} A_1 × $A_1 \times \mathbb{R}$ \mathbb{R} $A_2 \times A_1$ $A_1 \times \mathbb{R}$ A_2 A_4 $A_2 \times \mathbb{R}$ $A_1 \times \mathbb{R}$ \mathbb{R} D_5 $A_3 \times A_1$ $A_1^2 \times \mathbb{R}$ \mathbb{R}^2 A_6 A_5 A_2^2 $A_1^2 \times \mathbb{R}$ A_6 A_5 A_2^2 $A_1^2 \times \mathbb{R}$ A_7 D_6 A_5 $A_3 \times A_1$ $A_2 \times \mathbb{R}$ B_7 B_6 D_5 A_4 A_7 A_6 A_5	higher-dimension \times \mathbb{R} $A_1 \times \mathbb{R}$ $A_1 \times \mathbb{R}$ $A_2 \times A_1$ $A_1 \times \mathbb{R}$ $A_2 \times A_1$ $A_1 \times \mathbb{R}$ A_4 $A_2 \times \mathbb{R}$ $A_1 \times \mathbb{R}$ \mathbb{R} X \mathbb{R} \mathbb{R} \mathbb{R} \mathbb{R} \mathbb{R} $A_1 \times \mathbb{R}$ \mathbb{R} </th <th>$\begin{array}{ c c c c c c c c c c c c c c c c c c c$</th> <th>$\begin{array}{ c c c c c c c c c c c c c c c c c c c$</th> <th>higher-dimensional origin of 3D$\times$$\mathbb{R}$$A_1 \times \mathbb{R}$$\mathbb{R}$$A_1 \times \mathbb{R}$$\mathbb{R}$$\mathbb{R}$$\mathbb{Cremmer, Jx}$$A_2 \times A_1$$A_1 \times \mathbb{R}$$A_2$$A_1$$A_4$$A_2 \times \mathbb{R}$$A_1 \times \mathbb{R}$$\mathbb{R}$$\Delta_5$$A_3 \times A_1$$A_1^2 \times \mathbb{R}$$\mathbb{R}^2$$B_6$$A_5$$A_2^2$$A_1^2 \times \mathbb{R}$$A_1 \times \mathbb{R}$$A_7$$D_6$$A_5$$A_3 \times A_1$$A_2 \times \mathbb{R}$$A_4$$A_2 \times A_1$$A_1 \times \mathbb{R}$$A_1$$\mathbb{R}$</th>	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	higher-dimensional origin of 3D \times \mathbb{R} $A_1 \times \mathbb{R}$ \mathbb{R} $A_1 \times \mathbb{R}$ \mathbb{R} \mathbb{R} $\mathbb{Cremmer, Jx}$ $A_2 \times A_1$ $A_1 \times \mathbb{R}$ A_2 A_1 A_4 $A_2 \times \mathbb{R}$ $A_1 \times \mathbb{R}$ \mathbb{R} Δ_5 $A_3 \times A_1$ $A_1^2 \times \mathbb{R}$ \mathbb{R}^2 B_6 A_5 A_2^2 $A_1^2 \times \mathbb{R}$ $A_1 \times \mathbb{R}$ A_7 D_6 A_5 $A_3 \times A_1$ $A_2 \times \mathbb{R}$ A_4 $A_2 \times A_1$ $A_1 \times \mathbb{R}$ A_1 \mathbb{R}	

all embedded within $E_{8(8)}$ ExFT:

 $E_{8(8)} \longrightarrow SL(D-2) \times SL(9-r) \times U_{D,r}$

other magic triangles and pyramids ...



[Julia]

Henning Samtleben



ENS Lyon

FNS de LYO

the ExFT construction mainly relies on the (complex) algebra e_6

▶ different real forms: e.g. $E_{6(6)}/USp(8) \longrightarrow E_{6(-26)}/F_4$:

$$\mathcal{L} \equiv \hat{R} + \frac{1}{24} g^{\mu\nu} \mathcal{D}_{\mu} \mathcal{M}^{MN} \mathcal{D}_{\nu} \mathcal{M}_{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}^{\mu\nu M} \mathcal{F}^{N}_{\mu\nu} + e^{-1} \mathcal{L}_{\text{top}} - V(\mathcal{M}, e).$$

ExFT formally identical

different structure of (real) solutions to the section constraint

uplift to D=6 -----> magic N=2 series



the ExFT construction mainly relies on the (complex) algebra e_6

b different real forms: e.g. $E_{6(6)}/USp(8) \longrightarrow E_{6(-26)}/F_4$:

$$\mathcal{L} \equiv \hat{R} + \frac{1}{24} g^{\mu\nu} \mathcal{D}_{\mu} \mathcal{M}^{MN} \mathcal{D}_{\nu} \mathcal{M}_{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}^{\mu\nu M} \mathcal{F}^{N}_{\mu\nu} + e^{-1} \mathcal{L}_{\text{top}} - V(\mathcal{M}, e).$$

ExFT formally identical

different structure of (real) solutions to the section constraint

uplift to D=6 \longrightarrow magic N=2 series

▶ different coset space: $E_{6(6)}/USp(8) \longrightarrow E_{6(6)}/(USp(4) \times USp(4))$ reduction on torus including time-like circle

$$\mathcal{L} \equiv \hat{R} + \frac{1}{24} g^{\mu\nu} \mathcal{D}_{\mu} \mathcal{M}^{MN} \mathcal{D}_{\nu} \mathcal{M}_{MN} - \frac{1}{4} \mathcal{M}_{MN} \mathcal{F}^{\mu\nu M} \mathcal{F}^{N}_{\mu\nu} + e^{-1} \mathcal{L}_{\text{top}} - V(\mathcal{M}, e).$$

ExFT formally identical, flip of external/internal signatures

$$(4,1) \& (27,0) \iff (5,0) \& (16,11)$$

equivalent description, also captures Hull's *-theories ! SL(5): [Blair,Malek,Park]

Henning Samtleben



ENS Lyon

- unifying framework for brane solutions
 [Berman, Rudolph, Bakhmatov, Kleinschmidt, Musaev, Otsuki, Fernandez-Melgarejo, Kimura, Sakatani]
 1/2 BPS branes from a single ExFT solution, organisation of exotic branes
- orbifolds and orientifolds in ExFT

[Blair, Malek, Thompson]
unified approach in terms of generalized orbifolds (O-folds)

exceptional string sigma model [Arvanitakis, Blair] string sigma model with ExFT background fields

ExFT loop calculations [Bossard, Kleinschmidt]

duality covariant graviton amplitudes

underlying mathematical structures

[Cederwall, Palmkvist][Hohm, Kupriyanov, Lüst, Traube] [Cagnacci, Codina, Marques][Arvanitakis]

 L_{∞} -algebras, Borchers superalgebras, tensor hierarchy algebras

conclusions

exceptional field theory

- based on generalized diffeomorphisms in exceptional geometry
- unique theory with generalized diffeomorphism invariance in all coordinates (modulo section condition)
- upon an explicit solution of the section condition the theory coincides with full D=11 supergravity or full D=10 IIB supergravity
- includes generalized type II supergravities and Hull's II*-theories
- powerful tool for construction & analysis of vacua & consistent truncations



ENS Lyon

conclusions

exceptional field theory

- based on generalized diffeomorphisms in exceptional geometry
- unique theory with generalized diffeomorphism invariance in all coordinates (modulo section condition)
- upon an explicit solution of the section condition the theory coincides with full D=11 supergravity or full D=10 IIB supergravity
- includes generalized type II supergravities and Hull's II*-theories
- powerful tool for construction & analysis of vacua & consistent truncations
- tool for analyzing existing theories
 - or hints towards a more fundamental structure ..?
- weaken / relax section constraints
- \triangleright decrease number of external dimensions \longrightarrow unifying picture

