

Vector polarimetry at MAMI

Measurements of tensor correlation coefficients in
 e^- bremsstrahlung processes

Fabian Nillius

Institut für Kernphysik
Johannes Gutenberg-Universität, Mainz

February 13, 2014

This work was supported by the BMBF Verbundforschung
„Spinoptimierung“.

1 Introduction

- How to measure polarisation
- The correlation tensor
- Measuring principle

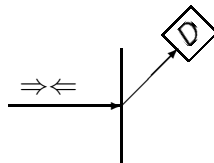
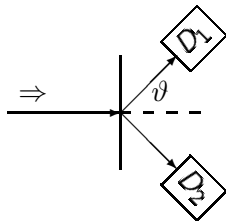
2 Setup and measurements

- Experimental setup
- Energy resolved measurement of the longitudinal polarisation via C_{32}
- First energy resolved measurement of the transverse polarisation via C_{12}

3 Summary & Outlook

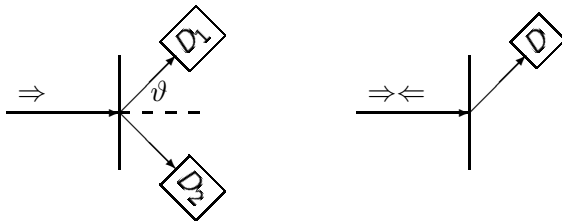
Polarisation is measured by an asymmetry:

- either $A = \frac{N_1 - N_2}{N_1 + N_2} = S_\vartheta \cdot P$ with a pair of two detectors, D_1 and D_2
- or by switching the polarisation and receiving different counting rates N^\Rightarrow and N^\Leftarrow in a single detector



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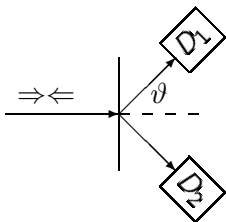
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Why not use both techniques?

Measuring a polarisation with the super ratio:

- using two detectors
- switching the polarisation
- measuring four rates N_1^{\Rightarrow} , N_2^{\Rightarrow} , N_1^{\Leftarrow} and N_2^{\Leftarrow} .



$$\mathcal{A} = \frac{1 - \sqrt{\frac{N_1^{\Rightarrow} \cdot N_2^{\Leftarrow}}{N_1^{\Leftarrow} \cdot N_2^{\Rightarrow}}}}{1 + \sqrt{\frac{N_1^{\Rightarrow} \cdot N_2^{\Leftarrow}}{N_1^{\Leftarrow} \cdot N_2^{\Rightarrow}}}}$$

$$N_m = N_t \cdot f$$

Tseng and Pratt¹ introduced a correlation tensor between the incoming electrons and the emitted photons $C_{\mu\nu}(Z, \vartheta, E)$.

- The first μ index represents the polarisation of the electrons:

$\mu = 0$: intensity

$\mu = 1$: transverse in the emission plane

$\mu = 2$: perpendicular to the emission plane

$\mu = 3$: longitudinal

- The second ν index describes the properties of the emitted photons

$\nu = 0$: intensity

$\nu = 1$: linear polarisation diagonal to emission plane
(45° & 135°)

$\nu = 2$: circular polarisation

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plane (x & y)

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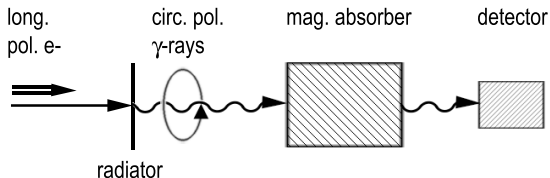
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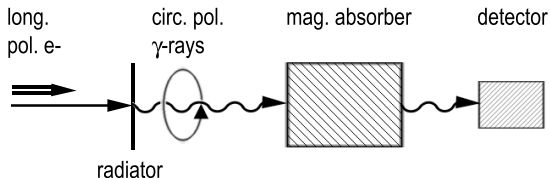
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This design is also known as a compton absorption polarimeter.

$$A_{full}(E_\gamma) = \frac{e^{-\sigma_p Z N_A \rho d / B} - e^{+\sigma_p Z N_A \rho d / B}}{e^{-\sigma_p Z N_A \rho d / B} + e^{+\sigma_p Z N_A \rho d / B}} = -\tanh\left(\frac{\sigma_p Z N_A \rho d}{B}\right)$$

$$\begin{aligned} A(E_\gamma) &= -P_e C_{32}(E_\gamma) \tanh\left(\sigma_p(E_\gamma) \frac{2N_A \rho d}{B}\right) \\ &\approx -P_e C_{32}(E_\gamma) \sigma_p(E_\gamma) \frac{2N_A \rho d}{B} \end{aligned}$$

σ_p = cross section sensitive to polarisation

Z = atomic number

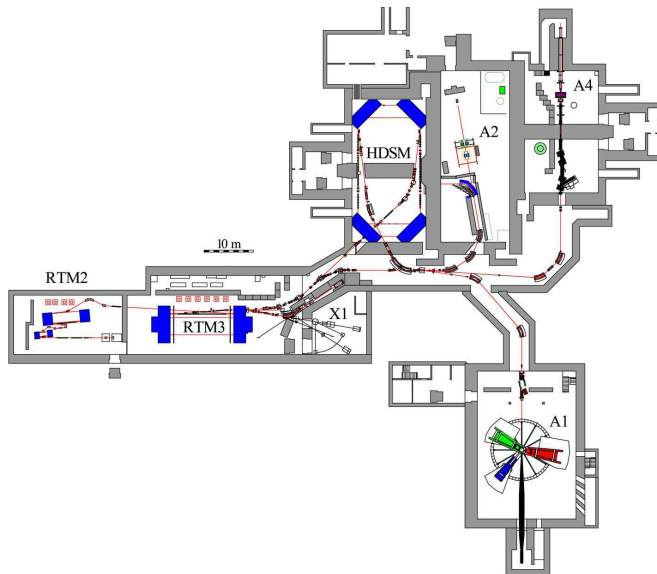
B = baryon number

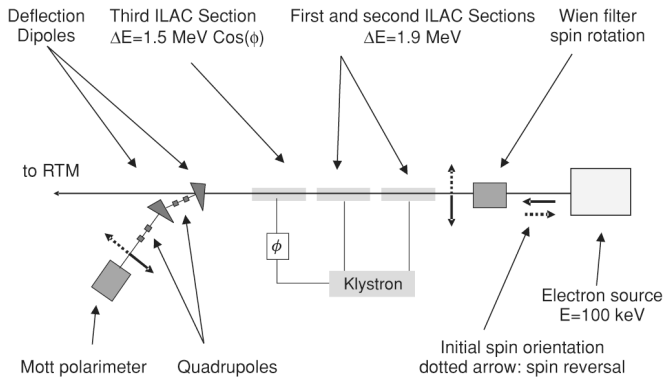
P_e = polarisation of the electron beam

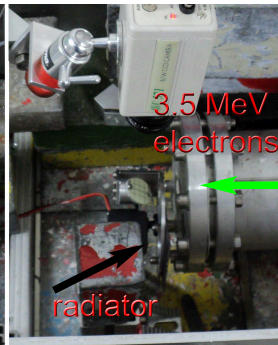
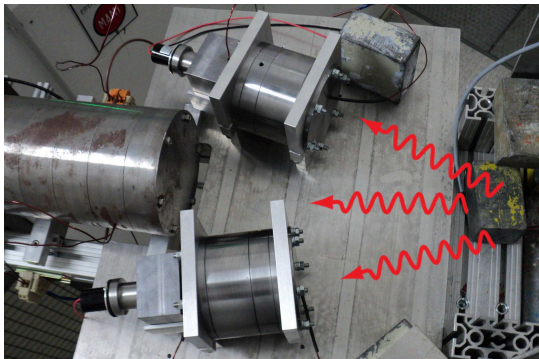
N_A = Avogadro number

d = absorber length

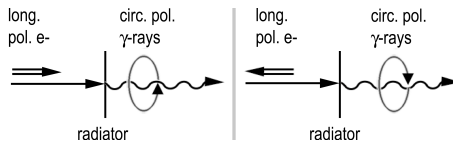
ρ = density





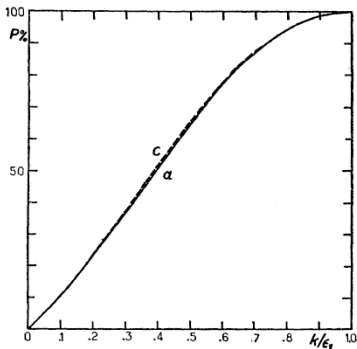


C_{32} describes the correlation between the longitudinal polarisation of the electrons and the circular polarisation of the outgoing photon.



C_{32}

The circular polarisation of bremsstrahlung from longitudinal polarised electrons for relativistic electrons was calculated by H. Olsen and L. C. Maximon² for 50 MeV ($\gamma = 101$).

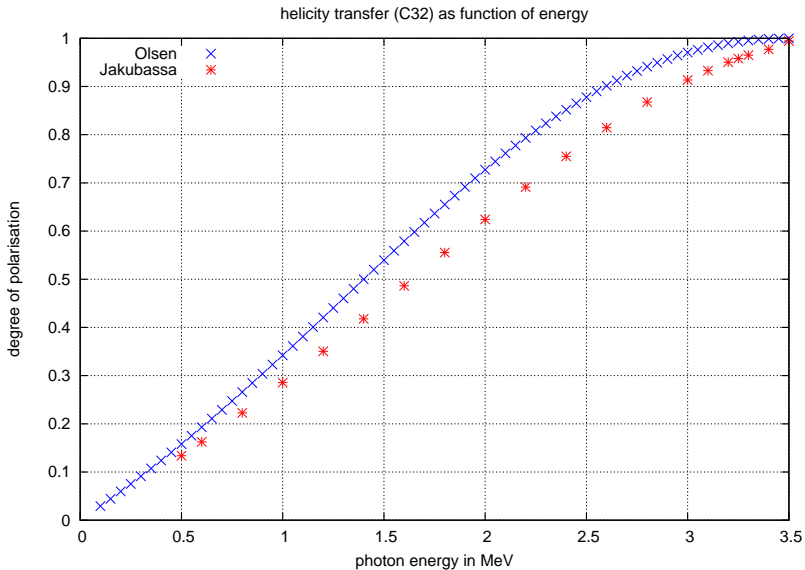


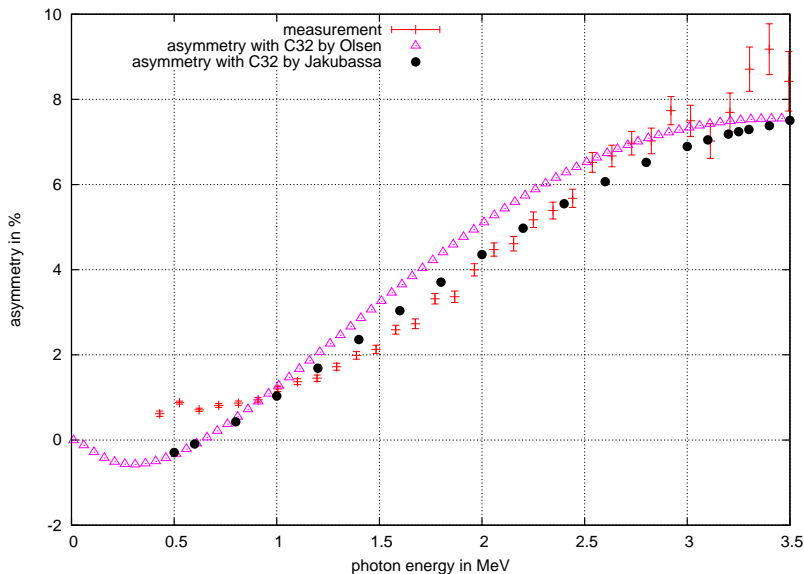
But what does this look like for

$$\gamma \neq 1$$

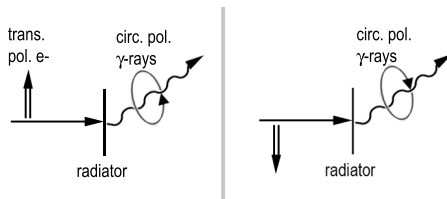
e. g. 8?

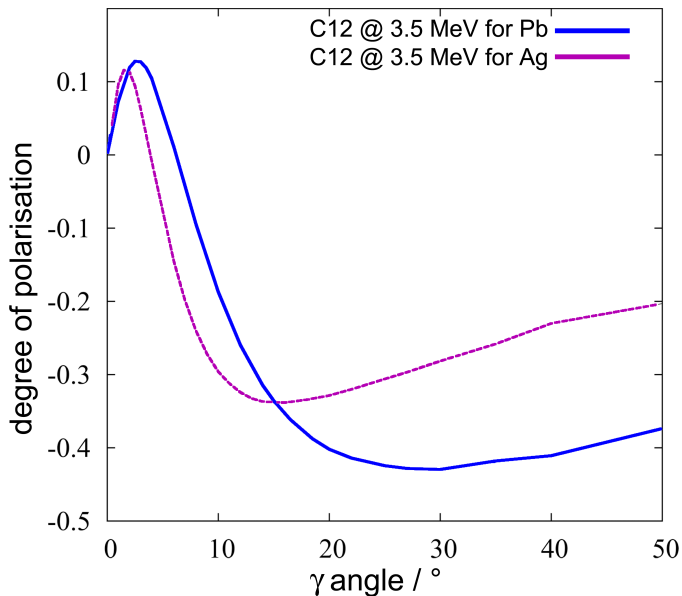
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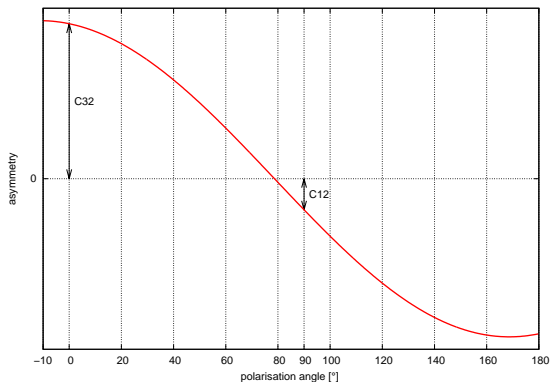




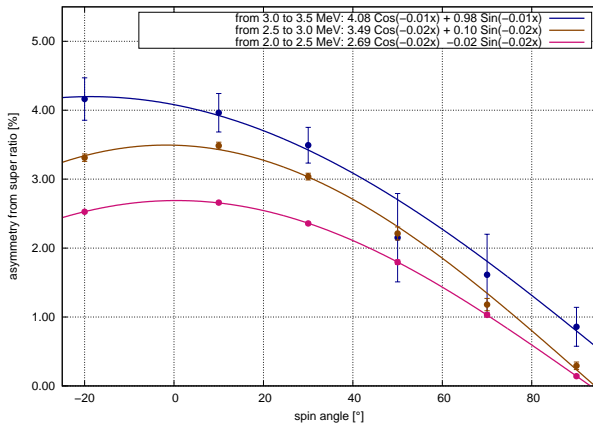
C_{12} describes the correlation between the transverse polarisation of the electrons and the circular polarisation of the outgoing photon.

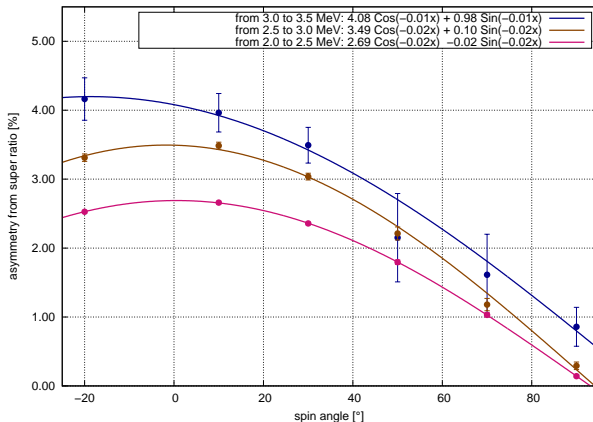






$$A(\phi_{\text{Pol}}) = \chi (C_{32} \cos(\phi_{\text{Pol}}) + C_{12} \sin(\phi_{\text{Pol}}))$$





$$\left(\frac{C_{32}}{C_{12}} \right)_{exp} = 4.163$$

$$\left(\frac{C_{32}}{C_{12}} \right)_{theo} = 4.086$$

Summary

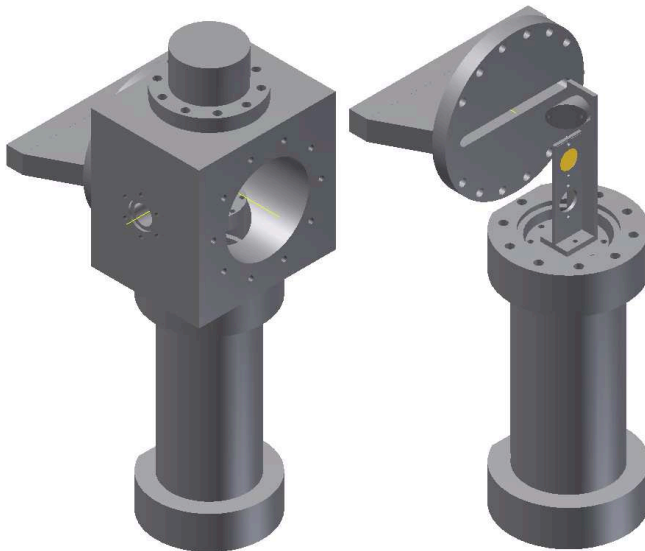
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- difference between theory of Olsen and Jakubassa can be seen

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Outlook

- simulation for analysis of systematic errors in
 - degree of polarisation
 - magnetisation
 - energy resolution
 - background
- measurements for different γ -angles
- improved experimental setup



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