

Direct Solutions of IBP Systems & Conjugate Polynomials

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@ Mathematics of Linear Relations between Feynman Integrals
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[1804.00131] & work in progress

Integrals

$$I^{(L)}[Poly(\{\ell_i\}, \{k_j\})] = \int \prod_{i=1}^L \frac{d^D \ell_i}{(2\pi)^D} \frac{Poly}{\prod_{j=1}^N d_j}$$

- Many integrals on line 1, mostly differing in *Poly*
- Polylogarithmic content determined by denoms
- Expect linear relations & a small number of master integrals

Integration by Parts

Chetyrkin & Tkachov (1981)

- Technique for finding all linear relations
 - Sufficient: Petukhov & Smirnov (2011)

$$0 = \int \prod_{i=1}^L \frac{d^D \ell_i}{(2\pi)^D} \frac{\partial}{\partial \ell_a^\mu} \frac{v^\mu \text{Poly}}{\prod_{j=1}^N d_j}$$

- No boundary term thanks to dimensional regularization
- Choose all possible vectors v^μ , sufficient number of *Polys* of sufficient dimension; system will close
- Solve system by Gaussian elimination

Laporta (2000)

- Nontrivial beyond one loop because of *irreducible numerators*
- Implementations: AIR [Anastasiou & Lazopoulos (2004)]; FIRE [Smirnov (2008–15)]; Reduze [Studerus & van Manteuffel (2010–12)]; Kira [Maierhoefer, Usovitch, & Uwer (2017)]; ...
- Non-Laporta: LiteRed [Lee (2012)]; Azurite [Georgoudis, Larsen, Zhang (2016)]

Generic Reduction

$I[\text{General Irreducible}]$

$$= \sum_{j \in \text{basis}} c_j(s_{ij}, \epsilon) I_j + \text{simpler graphs}$$

Problems

- Introduces unwanted integrals
 - In intermediate stages, lots of integrals with doubled propagators which never arise from Feynman diagrams & ultimately cancel

$$\frac{\partial}{\partial \ell_i^\mu} \frac{1}{d_j(\ell_i)} = - \frac{1}{d_j^2(\ell_i)} \frac{\partial}{\partial \ell_i^\mu} d_j(\ell_i)$$

- Need to solve huge systems of equations
 - Careful ordering, programming, disk management, ...
- Hard to do arbitrary powers

General parameter, not fixed value

$$I[\text{Irred}^n] = \sum_{j \in \text{basis}} c^{(n)}(s_{ij}, \epsilon) I_j + \text{simpler graphs}$$

Resolve these problems

IBP-Generating Vectors

Gluza, Kajda, & DAK (2011)

- Basic idea: choose *special* vectors s.t.

$$v_i^\mu \frac{\partial}{\partial \ell_i^\mu} d_j \propto d_j \quad \forall j$$

This eliminates new doubled propagators

$$\begin{aligned} v_i^\mu \frac{\partial}{\partial \ell_i^\mu} \frac{1}{d_j} &= -\frac{1}{d_j^2} v_i^\mu \frac{\partial}{\partial \ell_i^\mu} d_j \\ &= -\frac{p(\ell_i)}{d_j} \end{aligned}$$

Generating-Vector Equations

- Vectors come in L -tuples; focusing on two loops, we have pairs (v_1, v_2) , with equations

$$\sum_{i=1}^2 v_i^\mu \frac{\partial}{\partial \ell_i^\mu} d_j = c_j d_j \quad \forall j$$

where each c_j is a polynomial in the variables

$$\{\ell_1^2, \ell_1 \cdot \ell_2, \ell_2^2, \{\ell_1 \cdot b\}_{b \in B}, \{\ell_2 \cdot b\}_{b \in B}, s_{12}\}$$

with B a basis of external momenta; coefficients of terms in the polynomials are rational functions of $\chi_{ij} \equiv s_{ij}/s_{12}$

- Parametrize the vectors

$$v_i^\mu = c_i^{\ell_1} \ell_1^\mu + c_i^{\ell_2} \ell_2^\mu + \sum_{b \in B} c_i^b b^\mu$$

Generating-Vector Equations

- Organize the c s into a row vector \tilde{c}
- Doing all the algebra leads to an equation

$$\tilde{c} E = 0$$

where each column corresponds to a different denominator

This is a syzygy equation.

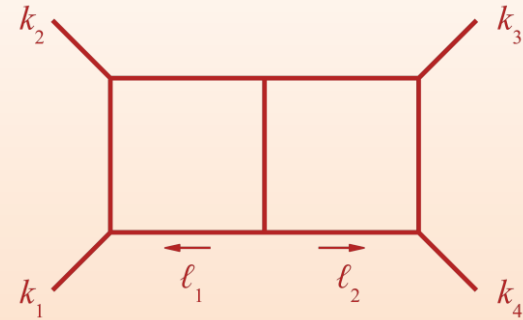
Solving

- We *cannot* solve this by linear algebra without using ansätze for the polynomials, because that would give rational and not polynomial solutions
- Find all syzygies
- Find independent basis set
 - Vector pairs that differ only by reducible invariants are not considered independent
- Standard problem in computational algebraic geometry
- Method uses Gröbner bases of tuples

[Linear algebra approach *Schabinger (2012), von Manteuffel (2018)*]

Example 1

- All-massless planar double box
- 3 vector pairs



$$v_{1;1}^{\mu} = -2k_4 \cdot \ell_1 k_1^{\mu} + \ell_1^2 k_2^{\mu} + (2k_1 \cdot \ell_1 - \ell_1^2) k_4^{\mu} - (2k_2 \cdot \ell_1 - 2k_4 \cdot \ell_1 - s_{12}) \ell_1^{\mu}$$

$$v_{1;2}^{\mu} = -2k_4 \cdot \ell_2 k_1^{\mu} - \ell_2^2 k_2^{\mu} + (2k_1 \cdot \ell_2 + \ell_2^2) k_4^{\mu} - (2k_4 \cdot \ell_2 - 2k_2 \cdot \ell_2 - s_{12}) \ell_2^{\mu}$$

other two a bit longer & of engineering dimension 5

$$0 = c_0[V_1, V_2] - c_1 V_1 - c_2 V_2 - c_3 V_3 + \text{purely reducible},$$

Define a compact notation for derivatives $\partial_A v_A \equiv \sum_{i=1}^2 \frac{\partial}{\partial \ell_i^{\mu}} v_i^{\mu}$

Compact Variables

- Use variables to simplify structure: r, u are purely reducible; t are irreducible

$$r_{11} = \ell_1^2 ,$$

$$r_{12} = \ell_1 \cdot \ell_2 ,$$

$$r_{22} = \ell_2^2 ,$$

$$u_{11} = \ell_1 \cdot k_1 ,$$

$$u_{12} = \ell_1 \cdot k_2 - s_{12}/2 ,$$

$$u_{23} = \ell_2 \cdot k_3 - s_{12}/2 ,$$

$$u_{24} = \ell_2 \cdot k_4 ,$$

$$t_{14} = \ell_1 \cdot k_4 ,$$

$$t_{21} = \ell_2 \cdot k_1$$

Conjugate Polynomials

- Goal: diagonalize systems a priori
- Make choices of the *Poly* to obtain IBPs for target numerators

$$0 = \int \prod_{i=1}^L \frac{d^D \ell_i}{(2\pi)^D} \partial_A \frac{v_A \text{Poly}}{\prod_{j=1}^N d_j}$$

- General expression for numerator

ME: $\sum_{r=1}^{n_v} \text{Poly}_r \text{Denom} \partial_A \frac{v_{rA}}{\text{Denom}} + \sum_{r=1}^{n_v} v_{rA} \partial_A \text{Poly}_r.$

Universal prefactor

Requires more than
Gröbner bases

- Universal prefactor for v_1

$$\text{Denom} \partial_A \frac{v_{1A}}{\text{Denom}} = -2 \epsilon (t_{14} - t_{21} - u_{12} - u_{23} - 2 u_{24})$$

- Simplest example: $Poly = 1$

$$\begin{aligned} 0 &= -2\epsilon I_{\text{DB}}[t_{14} - t_{21} - u_{12} - u_{23} - 2 u_{24}] \\ &= 2\epsilon I_{\text{DB}}[t_{21} - t_{14} + \text{purely reducible}] \\ &= 2\epsilon I_{\text{DB}}[t_{21} - t_{14}] + \text{simpler integrals} \end{aligned}$$

- Here, the simpler integrals cancel, and we can solve for $I_{\text{DB}}[t_{14}]$ in terms of one of the masters $I_{\text{DB}}[t_{21}]$

Motivation

- Multiply the first vector pair by

$$a_1 t_{14} + a_2 t_{21}$$

and the second pair by $b_1(1 + \chi)$, where $\chi = t/s$

We obtain the IBP

$$\begin{aligned} 0 = & I_{\text{DB}} \left[-\frac{1}{2} b_1 \chi \epsilon s_{12}^2 + \frac{1}{2} (-a_1 \chi + 2 b_1 \epsilon) s_{12} t_{14} \right. \\ & + a_1 (1 - 2 \epsilon) t_{14}^2 + \frac{1}{2} (a_2 \chi + b_1 \chi + 4 b_1 \epsilon) s_{12} t_{21} \\ & + 2 (a_1 - a_2 + b_1) \epsilon t_{14} t_{21} - (a_2 + b_1) (1 - 2 \epsilon) t_{21}^2 \\ & \left. + \text{purely reducible} \right]. \end{aligned}$$

If we make the ‘canonical’ choices

$$(a_1, a_2, b_1) = (1, 0, 0); (0, 1, 0); (0, 0, 1)$$

we obtain three equations which still need to be diagonalized

Example 2

- If instead we make the choice

$$a_1 = \frac{1}{(1 - 2\epsilon)} , \quad a_2 = \frac{1}{2(1 - 2\epsilon)} , \quad b_1 = -\frac{1}{2(1 - 2\epsilon)} ,$$

we obtain the direct IBP for $I_{\text{DB}}[t_{14}^2]$

$$I_{\text{DB}}[t_{14}^2] = \frac{(\chi + 3\epsilon)}{2(1 - 2\epsilon)} s_{12} I_{\text{DB}}[t_{21}] - \frac{\chi \epsilon}{4(1 - 2\epsilon)} s_{12}^2 I_{\text{DB}}[1] \\ + \text{simpler integrals}$$

We can obtain similar direct equations for the other quadratic irreducibles

Finding Master Integrals

- Write out all monomials of a given degree in a vector; for example, degree 3

$$\begin{pmatrix} t_{14}^3 \\ t_{14}^2 t_{21} \\ t_{14} t_{21}^2 \\ t_{21}^3 \\ t_{14}^2 s_{12} \\ t_{14} t_{21} s_{12} \\ t_{21}^2 s_{12} \\ t_{14} s_{12}^2 \\ t_{21} s_{12}^2 \\ s_{12}^3 \end{pmatrix}$$

- One such monomial vector for each IBP-generating tuple of vectors, with different degrees

Finding Master Integrals

- Independently take *Poly* to be each entry
 - Compute numerator using the Master Equation
 - Set reducibles to zero
- ⇒ Reduced numerators, which we can regard as a linear transformation of irreducibles
- Organize these into a matrix
 - Columns correspond to monomials
 - Rows correspond to IBP equations
 - Independent reductions \Leftrightarrow range of the matrix
 - Redundant reductions \Leftrightarrow master integrals \Leftrightarrow kernel of matrix

Masters Example: Double Box

- Take first vector pair, with degree-0 vector

$$M_1 = \begin{pmatrix} -2\epsilon & 2\epsilon & 0 \end{pmatrix}$$

with columns corresponding to t_{14}, t_{21}, s_{12}

Corresponding IBP is

$$0 = I_{\text{DB}} \left[M_1 \begin{pmatrix} t_{14} \\ t_{21} \\ s_{12} \end{pmatrix} \right] + \text{simpler topologies}$$

Kernel is two-dimensional, generated by

$$\begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

Masters Example

- Use non-trivial IBP to reduce first vector to get $I_{\text{DB}}[t_{21}]$ and $I_{\text{DB}}[1]$ as masters
- Might worry that we're missing equations
- Try polynomials of higher degree: degree-1 for 1st pair, degree-0 for 2nd & 3rd pairs

$$M_2 = \begin{pmatrix} 2(1 - 2\epsilon) & 4\epsilon & 0 & -\chi & 0 & 0 \\ 0 & -4\epsilon & -2(1 - 2\epsilon) & 0 & \chi_{14} & 0 \\ 0 & 0 & 0 & -2\epsilon & 2\epsilon & 0 \\ 0 & \frac{2\epsilon}{1+\chi} & -\frac{1-2\epsilon}{1+\chi_{14}} & \frac{\epsilon}{1+\chi} & \frac{\chi+4\epsilon}{2(1+\chi)} & -\frac{\chi\epsilon}{2(1+\chi)} \\ 0 & -8\epsilon & 0 & -6\epsilon & 0 & \chi\epsilon \end{pmatrix}$$

with columns corresponding to $t_{14}^2, t_{14}t_{21}, t_{21}^2, s_{12}t_{14}, s_{12}t_{21}, s_{12}^2$

Masters Example

- Again a kernel of dimension 2, with same masters as before
- Repeat with one higher dimension \Rightarrow same result
- Pedestrian procedure: assume convergence
- More complete solution will presumably require D -algebras
- Should also provide connections to algebraic geometry

Doubled Propagators

- IBP-generating vectors ensure that no doubled propagators are newly introduced
- What about those already present?
- IBP-generating vectors can be used with them too: ensure no higher powers are introduced
- For double box, usual integrals require only first two vector pairs
- Doubled propagators also need third pair, but reduce to usual masters (no new masters)

General Powers

- Multiply the 1st vector pair by

$$a_1 t_{14}^{n-1} + a_2 t_{14}^{n-2} t_{21} + a_3 t_{14}^{n-2} s_{12}$$

and the 2nd pair by

$$b_1(1 + \chi)t_{14}^{n-2}$$

- Feed it through the Master Equation to obtain the IBP

$$\begin{aligned} 0 = I_{\text{DB}} & \left[a_1 (1 + 2\epsilon - n) t_{14}^n \right. \\ & - (2a_1\epsilon + (b_1 - a_2)(2 + 2\epsilon - n)) t_{14}^{n-1} t_{21} \\ & + (a_2 + b_1)(1 - 2\epsilon) t_{14}^{n-2} t_{21}^2 \\ & \left. + \text{lower irreducible degree} \right] \\ & + \text{simpler integrals} \end{aligned}$$

General Powers

- Choose

$$a_1 = \frac{1}{n-1-2\epsilon},$$

$$a_2 = -\frac{\epsilon}{(n-1-2\epsilon)(n-2-2\epsilon)},$$

$$a_3 = \frac{1}{2(n-1-2\epsilon)},$$

$$b_1 = \frac{\epsilon}{(n-1-2\epsilon)(n-2-2\epsilon)}$$

to obtain an IBP equation for $I_{\text{DB}}[t_{14}^n]$

$$0 = I_{\text{DB}} \left[t_{14}^n + \frac{(2+(n-1)\chi+3\epsilon-n)}{2(1+2\epsilon-n)} t_{14}^{n-1} s_{12} - \frac{\chi(2+\epsilon-n)}{4(1+2\epsilon-n)} t_{14}^{n-2} s_{12}^2 \right]$$

+ simpler integrals

n could even be non-integer (with additional masters)

Recurrence Relations

- This is a recurrence relation: $w_n \equiv s_{12}^{-n} I_D[t_{14}^n]$

$$4(1+2\epsilon-n)w_n + 2(2+(n-1)\chi+3\epsilon-n)w_{n-1} - \chi(2+\epsilon-n)w_{n-2} \stackrel{\cdot}{=} 0$$

- How can we solve it?
- First thought: use *Mathematica*
- Solution in terms of **DifferenceRoot** objects
 - Possibly efficient, but useless analytically

Solving: Generating Function

- Instead, build a generating function

$$\sum_{n=0}^{\infty} Rec_n x^n$$

- Substitute

$$\sum_{n=0}^{\infty} c_n a_{n+r} x^n \rightarrow x^{-r} \left(\sum_{n=0}^{\infty} c_{n-r} a_n x^n - x^{-r} \sum_{n=0}^{r-1} c_{n-r} a_n x^n \right)$$

- Replace $\sum_{n=0}^{\infty} n^p a_n x^n \rightarrow D_x^p f(x)$ ($D_x \equiv x \partial_x$)

to obtain a differential equation; differentiate it again to make it homogeneous

Solving: Differential Equation

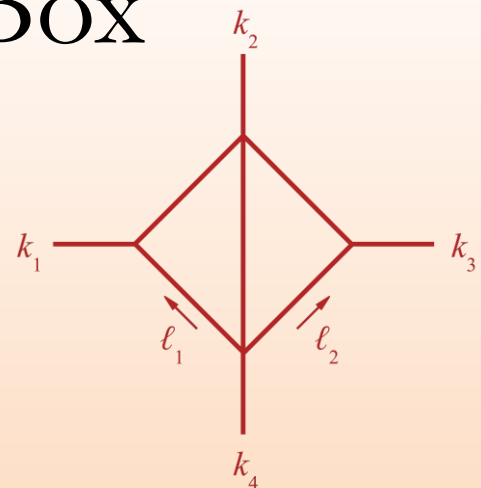
- Solve differential equation
- Boundary conditions given by $x \rightarrow 0$ behavior as well as master integrals
- Extract n^{th} power in x to obtain desired coefficient

Example: Slashed Box

- One master integral $I_{\text{SB}}[1]$

- Four irreducible invariants

- Seven pairs of generating vectors ($\check{r}_{12} = \ell_1 \cdot \ell_2 + t_{14} + t_{24}$, $\check{u}_{23} = \ell_2 \cdot k_3$)



$$v_{1;1}^\mu = -k_1^\mu r_{11} - k_2^\mu r_{11} - \ell_1^\mu (s_{12} - 2t_{12} - 2u_{11}),$$

$$v_{1;2}^\mu = k_1^\mu (r_{22} + 2t_{24}) + k_2^\mu (r_{22} + 2t_{24}) + 2\ell_2^\mu (t_{24} + \check{u}_{23}) + 2k_4^\mu (t_{24} + \check{u}_{23}),$$

$$v_{2;1}^\mu = \frac{1}{2} k_2^\mu r_{11} + \frac{1}{2} k_4^\mu r_{11} - k_1^\mu (\check{r}_{12} - t_{14} - t_{24}) + \ell_2^\mu u_{11}$$

$$+ \frac{1}{2} \ell_1^\mu (s_{12} + \chi s_{12} - 2t_{12} - 2t_{14} - 2t_{22} - 2t_{24} - 2\check{u}_{23}),$$

$$v_{2;2}^\mu = \frac{1}{2} k_2^\mu (2\check{r}_{12} + r_{22} - 2t_{14} - 2t_{24}) + \frac{1}{2} k_4^\mu (2\check{r}_{12} + r_{22} - 2t_{14} - 2t_{24}) + \ell_1^\mu \check{u}_{23}$$

$$+ k_1^\mu (\check{r}_{12} - t_{14} - t_{24}) + \ell_2^\mu (\frac{1}{2}s_{12} + \frac{1}{2}\chi s_{12} - t_{12} - t_{14} - t_{22} - t_{24} - u_{11})$$

⋮

Example: Slashed Box

- Study $\check{w}_n \equiv s_{12}^{-n} I_{\text{SB}}[t_{12}^n]$

- Recurrence

$$0 \doteq (1+n) \check{w}_n + 2 \left((3+2\chi)(1+\epsilon) - (\chi+2)(3+n) \right) \check{w}_{n+1} \\ + 4(1+\chi)(n+2-3\epsilon) \check{w}_{n+2}$$

- Differential equation

$$0 = -2f(x) + 2 \left(2(2-\epsilon) + 2(1+\chi)(1-2\epsilon) - 5x \right) f'(x) \\ - \left(4(1+\chi)((2-\epsilon)(1-x) - 2\epsilon) + x(7x-10+2\epsilon) \right) f''(x) \\ - (x-2)x(x-2(1+\chi)) f^{(3)}(x)$$

Example: Slashed Box

- Raw solution to differential equation

$$\begin{aligned}
 f(x) = & \frac{x^{3\epsilon} c_1}{(2-x)^\epsilon (2(1+\chi) - x)^{1+2\epsilon}} \\
 & - \frac{2^{3\epsilon} (1+\chi)^{2\epsilon} c_2}{6\epsilon (2-x)^\epsilon (2(1+\chi) - x)^{1+2\epsilon}} \\
 & \times F_1(-3\epsilon, -\epsilon, -2\epsilon, 1-3\epsilon; \frac{x}{2}, \frac{x}{2(1+\chi)}) \\
 & + \frac{2^{3\epsilon} (1+\chi)^{2\epsilon} x (c_2 + 2c_3)}{4(1-3\epsilon) (2-x)^\epsilon (2(1+\chi) - x)^{1+2\epsilon}} \\
 & \times F_1(1-3\epsilon, 1-\epsilon, -2\epsilon, 2-3\epsilon; \frac{x}{2}, \frac{x}{2(1+\chi)})
 \end{aligned}$$

Appell function

- $x \rightarrow 0$ behavior forces $c_1 = 0$
- $f(0) = \tilde{w}_0$ fixes c_2
- $f'(0) = \tilde{w}_1 = c \tilde{w}_0$ fixes $c_3 = 0$

Example: Slashed Box

- Solution to differential equation with desired boundary behavior

$$\begin{aligned} f(x) = & - \frac{3 \, 2^{3 \epsilon} (1 + \chi)^{1+2 \epsilon} \epsilon x}{(1 - 3 \epsilon) (2 - x)^\epsilon (2(1 + \chi) - x)^{1+2 \epsilon}} \\ & \times F_1 \left(1 - 3 \epsilon, 1 - \epsilon, -2 \epsilon, 2 - 3 \epsilon; \frac{x}{2}, \frac{x}{2(1 + \chi)} \right) \check{w}_0 \\ & + \frac{2^{1+3 \epsilon} (1 + \chi)^{1+2 \epsilon}}{(2 - x)^\epsilon (2(1 + \chi) - x)^{1+2 \epsilon}} \\ & \times F_1 \left(-3 \epsilon, -\epsilon, -2 \epsilon, 1 - 3 \epsilon; \frac{x}{2}, \frac{x}{2(1 + \chi)} \right) \check{w}_0 \end{aligned}$$

Example: Slashed Box

- Brute-force expansion to extract x^n term yields a triple sum

$$\begin{aligned}
 \check{w}_n = & \frac{6 \epsilon^3 (1 + \chi) \check{w}_0}{2^n \Gamma(1 - 2 \epsilon) \Gamma(1 - \epsilon) \Gamma(1 + \epsilon) \Gamma(1 + 2 \epsilon)} \left[\sum_{n_1=1}^n \frac{1}{(1 + \chi)^{n_1} (n - n_1 + 1 - 3 \epsilon)} \right. \\
 & \times \sum_{n_3=0}^{n-n_1} \frac{\Gamma(n - n_1 - n_3 + 2 - \epsilon) \Gamma(n_3 - 2 \epsilon)}{(1 + \chi)^{n_3} (n - n_1 - n_3 + 1)! n_3!} \\
 & \times \sum_{n_4=0}^{n_1} \frac{(1 + \chi)^{n_4} \Gamma(n_1 - n_4 + 2 \epsilon) \Gamma(n_4 + \epsilon)}{(n_1 - n_4 - 1)! n_4!} \\
 & + \Gamma(1 - \epsilon) (1 + \chi)^{-n-1} \sum_{n_1=1}^{n+1} \frac{\Gamma(n - n_1 + 1 - 2 \epsilon)}{(n - n_1 + 1 - 3 \epsilon) (n - n_1 + 1)!} \\
 & \left. \times \sum_{n_4=0}^{n_1} \frac{(1 + \chi)^{n_4} \Gamma(n_1 - n_4 + 2 \epsilon) \Gamma(n_4 + \epsilon)}{(n_1 - n_4 - 1)! n_4!} \right]
 \end{aligned}$$

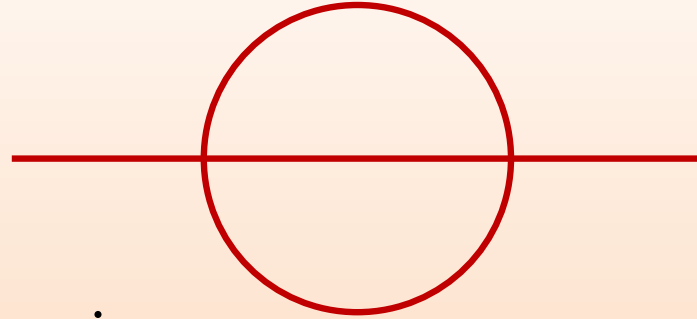
Example: Slashed Box

- A little sprinkling of magic simplifies the expression

$$\check{w}_n = -\frac{2^{1-n}\epsilon\Gamma(n-2\epsilon)\Gamma(1-3\epsilon)}{\Gamma(1-2\epsilon)\Gamma(n+1-3\epsilon)} {}_2F_1(1-\epsilon, -n; 1-n+2\epsilon; (1+\chi)^{-1}) \check{w}_0$$

- Still manifestly rational in χ and ϵ

Sunrise at Noon



- Equal-mass case
- 2 irreducible invariants
- 4 IBP-generating vector pairs
- Naively, four masters; but symmetries & iterated integration reduce this to 2: $\text{Sun}_{2:m}[1]$ and $\text{Sun}_{2:m}[(\ell_1 \cdot K)^2]$
- Conjugate polynomial: degrees (3,2,2,2)
- $y_n \equiv s^{-n} \text{Sun}_{2:m}[(\ell_1 \cdot K)^n]$
- Recurrence

$$\begin{aligned} 0 = & n\tau(1 + \tau)(-1 + 3\tau)y_{n-1} + 2\tau(-2 + \epsilon - 2n + \epsilon\tau + 2n\tau)y_n \\ & + (3 - 2\epsilon + n - 14\tau + 8\epsilon\tau - 6n\tau - 9\tau^2 + 6\epsilon\tau^2 - 3n\tau^2)y_{n+1} \\ & + 2(8 - 5\epsilon + 2n - 6\tau + 3\epsilon\tau - 2n\tau)y_{n+2} - 4(-5 + 3\epsilon - n)y_{n+3} + c_n \text{tadpole} \end{aligned}$$

Summary

- New approach to integration-by-parts systems
- IBP-generating vectors to remove unwanted doubled propagators & block-diagonalize system
- Conjugate polynomials to fully diagonalize, and target specific numerators
- Derive recurrence relations for arbitrary powers of irreducibles