# Minimal dark matter models

with radiative neutrino masses

#### Michael Klasen

#### Institute for Theoretical Physics, University of Münster

#### 18 September 2018





GEFÖRDERT VOM



Radiative seesaw models

Conclusion O

#### References

 D. Restrepo, O. Zapata, C. Yaguna Models with radiative neutrino masses and viable dark matter candidates JHEP 1311 (2013) 011 [1308.3655]

# References

- D. Restrepo, O. Zapata, C. Yaguna Models with radiative neutrino masses and viable dark matter candidates JHEP 1311 (2013) 011 [1308.3655]
- MK, C. Yaguna, J. Ruiz-Alvarez, D. Restrepo, O. Zapata Scalar dark matter and fermion coannihilations in the radiative seesaw model JCAP 1304 (2013) 044 [1302.5298]
- S. Esch, MK, D. Lamprea, C. Yaguna Lepton flavor violation and scalar dark matter in a radiative model of neutrino masses Eur. Phys. J. C 78 (2018) 88 [1602.05137]
- S. Esch, MK, C. Yaguna A singlet doublet dark matter with radiative neutrino masses JHEP (under review) [1804.03384]

Radiative seesaw models

Conclusion O

### Observational evidence for dark matter

MK, M. Pohl, G. Sigl, Prog. Nucl. Part. Phys. 85 (2015) 1 [1507.03800]









#### Galactic rotation curve data and MOND

M. Frandsen, J. Petersen, 1805.10706



Radiative seesaw models

#### Constraints on Primary Black Hole mass and DM fraction

B. Carr, M. Raidal, T. Tenkanen, V. Vaskonen, H. Veermäe, Phys. Rev. D96 (2017) 023514 [1705.05567]



Radiative seesaw models

## Probing the Weakly Interacting Massive Particle paradigm

E. Aprile et al. [Xenon1T], 1805.12562



#### DAMA/LIBRA-phase2 and isospin-violating DM

R. Bernabei et al. [DAMA/LIBRA-phase2], 1805.10486; S. Baum, K. Freese, C. Kelso, 1804.01231



#### Models with radiative neutrino masses and dark matter





Radiative seesaw models

Conclusion O

#### The best studied class: T3

Y. Farzan, S. Pascoli, M.A. Schmidt, JHEP 1010 (2010) 111; M. Aoki, S. Kanemura, K. Yagyu, Phys. Lett. B702 (2011) 355
 E. Ma, Phys. Rev. D73 (2006) 077301; E. Ma, D. Suematsu, Mod. Phys. Lett. A24 (2009) 583

Model	~	Ferm	ionic	Scalar		Exotic charges	# of N'plets
Model	u	DM	DD	DM	DD	Exotic charges	$\pi$ or in pieces
Τ2 Λ	0	21	×	$1_0, 3_2$	√	√	3
13-A	-2	$2_{-1}$	×	3 <sub>0</sub>	~	×	3
T2 B	1,-3	×	×	2 <sub>±1</sub>	~	√	3
13-0	-1	10	$\checkmark$	$2_{\pm 1}$	$\checkmark$	×	2
T3-C	1, -3	$3_{\pm 2}$	×	$2_{\pm 1}$	<ul> <li>✓</li> </ul>	√	3
13-0	-1	3 <sub>0</sub>	$\checkmark$	$2_{\pm 1}$	~	×	2
T3-E	0, -2	$2_{\pm 1}$	×	$3_0, 3_{\pm 2}$	✓	✓	3

Radiative seesaw models

### The best studied class: T3

Y. Farzan, S. Pascoli, M.A. Schmidt, JHEP 1010 (2010) 111; M. Aoki, S. Kanemura, K. Yagyu, Phys. Lett. B702 (2011) 355
 E. Ma, Phys. Rev. D73 (2006) 077301; E. Ma, D. Suematsu, Mod. Phys. Lett. A24 (2009) 583

Model	0	Fermionic		Scala	r	Exotic charges	# of N'plets
Widder		DM	DD	DM	DD	Exotic charges	
Τ3-Δ	0	21	×	1 <sub>0</sub> , 3 <sub>2</sub>	~	√	3
13-4	-2	$2_{-1}$	×	3 <sub>0</sub>	<ul> <li>✓</li> </ul>	×	3
T3 B	1,-3	×	×	2 <sub>±1</sub>	~	√	3
13-0	-1	10	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	×	2
T3-C	1, -3	$3_{\pm 2}$	×	$2_{\pm 1}$	<ul> <li>✓</li> </ul>	√	3
15-0	-1	3 <sub>0</sub>	$\checkmark$	$2_{\pm 1}$	~	×	2
T3-E	0, -2	2+1	×	$3_0, 3_{+2}$	~	✓	3

Potential in T3-B with  $\alpha = -1$  ("Inert Doublet Model"):

$$\begin{aligned} -\mathcal{L}_{H_{1},H_{2}} &= \mu_{1}^{2}H_{1}^{\dagger}H_{1} + \mu_{2}^{2}H_{2}^{\dagger}H_{2} + \frac{\lambda_{1}}{2}\left(H_{1}^{\dagger}H_{1}\right)^{2} + \frac{\lambda_{2}}{2}\left(H_{2}^{\dagger}H_{2}\right)^{2} \\ &+ \lambda_{3}\left(H_{1}^{\dagger}H_{1}\right)\left(H_{2}^{\dagger}H_{2}\right) + \lambda_{4}\left(H_{1}^{\dagger}H_{2}\right)\left(H_{2}^{\dagger}H_{1}\right) + \frac{\lambda_{5}}{2}\left(H_{1}^{\dagger}H_{2}\right)^{2} \\ &+ \text{ h.c.} \end{aligned}$$

#### Radiative corrections to the Inert Doublet Model

MK, J. Ruiz-Alvarez, C. Yaguna, Phys. Rev. D87 (2013) 075025 [1302.1657]



Radiative seesaw models

Conclusion O

#### The best studied class: T3

Y. Farzan, S. Pascoli, M.A. Schmidt, JHEP 1010 (2010) 111; M. Aoki, S. Kanemura, K. Yagyu, Phys. Lett. B702 (2011) 355
 E. Ma, Phys. Rev. D73 (2006) 077301; E. Ma, D. Suematsu, Mod. Phys. Lett. A24 (2009) 583

Model	0	Fermionic		Scalar		Exotic charges	# of N'plets	
Widder		DM	DD	DM	DD	Exotic charges	# Of N piets	
Τ3_Δ	0	21	×	1 <sub>0</sub> , 3 <sub>2</sub>	<ul> <li>✓</li> </ul>	√	3	
13-4	-2	$2_{-1}$	×	3 <sub>0</sub>	~	×	3	
T3 B	1,-3	×	×	2 <sub>±1</sub>	~	√	3	
13-0	-1	10	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	×	2	
T3.C	1, -3	$3 \pm 2$	×	$2_{\pm 1}$	√	√	3	
13-0	-1	3 <sub>0</sub>	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	×	2	
T3-E	0, -2	$2_{\pm 1}$	×	$3_0, 3_{\pm 2}$	~	$\checkmark$	3	

Potential in T3-B with  $\alpha = -1$  ("Scotogenic Model"):

$$\begin{aligned} -\mathcal{L}_{H_{1},H_{2}} &= \mu_{1}^{2}H_{1}^{\dagger}H_{1} + \mu_{2}^{2}H_{2}^{\dagger}H_{2} + \frac{\lambda_{1}}{2}\left(H_{1}^{\dagger}H_{1}\right)^{2} + \frac{\lambda_{2}}{2}\left(H_{2}^{\dagger}H_{2}\right)^{2} \\ &+ \lambda_{3}\left(H_{1}^{\dagger}H_{1}\right)\left(H_{2}^{\dagger}H_{2}\right) + \lambda_{4}\left(H_{1}^{\dagger}H_{2}\right)\left(H_{2}^{\dagger}H_{1}\right) + \frac{\lambda_{5}}{2}\left(H_{1}^{\dagger}H_{2}\right)^{2} \\ &+ \text{ h.c.} \end{aligned}$$

Radiative seesaw models

### The best studied class: T3

Y. Farzan, S. Pascoli, M.A. Schmidt, JHEP 1010 (2010) 111; M. Aoki, S. Kanemura, K. Yagyu, Phys. Lett. B702 (2011) 355
 E. Ma, Phys. Rev. D73 (2006) 077301; E. Ma, D. Suematsu, Mod. Phys. Lett. A24 (2009) 583

Model	0	Fermionic		Scalar		Exotic charges	# of N'plets
Iniodei		DM	DD	DM	DD	Exotic charges	# of it piets
T3-4	0	21	×	$1_0, 3_2$	✓	√	3
13-4	-2	2-1	×	3 <sub>0</sub>	~	×	3
T3 B	1,-3	×	×	2 <sub>±1</sub>	~	√	3
13-0	-1	10	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	×	2
T3.C	1, -3	3 <sub>±2</sub>	×	$2_{\pm 1}$	√	√	3
13-0	-1	30	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	×	2
T3-E	0, -2	$2_{\pm 1}$	×	$3_0, 3_{\pm 2}$	<ul> <li>✓</li> </ul>	√	3

Potential in T3-B with  $\alpha = -1$  ("Scotogenic Model"):

$$\begin{aligned} -\mathcal{L}_{H_1,H_2} &= \mu_1^2 H_1^{\dagger} H_1 + \mu_2^2 H_2^{\dagger} H_2 + \frac{\lambda_1}{2} \left( H_1^{\dagger} H_1 \right)^2 + \frac{\lambda_2}{2} \left( H_2^{\dagger} H_2 \right)^2 \\ &+ \lambda_3 \left( H_1^{\dagger} H_1 \right) \left( H_2^{\dagger} H_2 \right) + \lambda_4 \left( H_1^{\dagger} H_2 \right) \left( H_2^{\dagger} H_1 \right) + \frac{\lambda_5}{2} \left( H_1^{\dagger} H_2 \right)^2 \\ &+ \text{h.c.} \\ \mathcal{L}_N &= h_{\alpha i} \bar{\ell}_{\alpha} H_2^{\dagger} P_R N_i + \text{h.c.} \end{aligned}$$

#### Coannihilations in the Scotogenic Model

MK, C. Yaguna, J. Ruiz-Alvarez, D. Restrepo, O. Zapata, JCAP 1304 (2013) 044



Radiative seesaw models

Conclusion O

#### Class T1-3

S. Fraser, E. Ma, O. Popov, Phys. Lett. B737 (2014) 280; D. Restrepo et al., Phys. Rev. D92 (2015) 013005;
 S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C78 (2018) 88; J. Fiaschi, MK, S. May, in preparation

Model	~	Fermionic		Sca	lar	Exotic charges	# of N'plets
Widder	L u	DM	DD	DM	DD	Exotic charges	# Of N piets
T1-3-A	0	$1_0, 2_{\pm 1}$	~	10	$\checkmark$	×	3
T1-3-B	0	$1_0, 2_{\pm 1}$	~	30	~	×	3
T1-3-C	±1	$1_0, 2_{\pm 1}$	√	21	$\checkmark$	×	4
T13D	1	2 <sub>1</sub> , 3 <sub>0</sub>	~	21	$\checkmark$	×	4
11-5-0	$^{-1}$	$1_0, 2_{-1}, 3_{-2}$	$\checkmark$	$2_{-1}$	$\checkmark$	$\checkmark$	4
T1-3-F	±1	$2_{\pm 1}, 3_0, 3_{\pm 2}$	√	$2_{\pm 1}$	$\checkmark$	√	4
T1-3-G	0	$2_{\pm 1}, 3_0$	~	10	$\checkmark$	×	3
T1-3-H	0	$2_{\pm 1}, 3_0$	$\checkmark$	30	$\checkmark$	×	3

Conclusion O

### Class T1-3

S. Fraser, E. Ma, O. Popov, Phys. Lett. B737 (2014) 280; D. Restrepo et al., Phys. Rev. D92 (2015) 013005;
 S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C78 (2018) 88; J. Fiaschi, MK, S. May, in preparation

Model	0	Fermionic		Sca	ılar	Exotic charges	# of N'plets
Widder		DM	DD	DM	DD	Exotic charges	# Of N piets
T1-3-A	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-B	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	30	<ul> <li>✓</li> </ul>	×	3
T1-3-C	±1	$1_0, 2_{\pm 1}$	✓	21	<ul> <li>✓</li> </ul>	×	4
T1.3.D	1	2 <sub>1</sub> , 3 <sub>0</sub>	<ul> <li>✓</li> </ul>	21	$\checkmark$	×	4
11-5-0	-1	$1_0, 2_{-1}, 3_{-2}$	$\checkmark$	2_1	$\checkmark$	$\checkmark$	4
T1-3-F	±1	$2_{\pm 1}, 3_0, 3_{\pm 2}$	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	<ul> <li>✓</li> </ul>	√	4
T1-3-G	0	$2_{\pm 1}, 3_0$	<ul> <li>✓</li> </ul>	10	~	×	3
T1-3-H	0	$2_{\pm 1}, 3_0$	~	30	$\checkmark$	×	3

Potential in T1-3-B with  $\alpha = 0$ :

$$\begin{aligned} -\mathcal{L}_{Y} &= (\lambda_{1})^{ij} (H^{\dagger} H) \operatorname{Tr}(\phi_{i} \phi_{j}) + (\lambda_{3})^{ijkm} \operatorname{Tr}(\phi_{i} \phi_{j} \phi_{k} \phi_{m}) \\ &+ (\lambda_{4} (H^{\dagger} \psi') \Psi + \text{h.c.}) + (\lambda_{5} (H \psi) \Psi + \text{h.c.}) \\ &+ ((\lambda_{6})^{ij} L_{i} \phi_{j} \psi' + \text{h.c.}) \end{aligned}$$

Conclusion O

### Class T1-3

S. Fraser, E. Ma, O. Popov, Phys. Lett. B737 (2014) 280; D. Restrepo et al., Phys. Rev. D92 (2015) 013005; S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C78 (2018) 88; J. Fiaschi, MK, S. May, in preparation

Model	0	Fermionic		Sca	lar	Exotic charges	# of N'plets
Widder		DM	DD	DM	DD	Exotic charges	# of N piets
T1-3-A	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-B	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	30	$\checkmark$	×	3
T1-3-C	±1	$1_0, 2_{\pm 1}$	✓	21	$\checkmark$	×	4
T1.3.D	1	2 <sub>1</sub> , 3 <sub>0</sub>	<ul> <li>✓</li> </ul>	21	$\checkmark$	×	4
11-3-0	$^{-1}$	$1_0, 2_{-1}, 3_{-2}$	~	2-1	$\checkmark$	√	4
T1-3-F	±1	$2_{\pm 1}, 3_0, 3_{\pm 2}$	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	√	4
T1-3-G	0	$2_{\pm 1}, 3_0$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-H	0	$2_{\pm 1}, 3_0$	~	30	$\checkmark$	×	3

Potential in T1-3-B with  $\alpha = 0$ :

$$\begin{aligned} -\mathcal{L}_{Y} &= (\lambda_{1})^{ij} (H^{\dagger}H) \operatorname{Tr}(\phi_{i}\phi_{j}) + (\lambda_{3})^{ijkm} \operatorname{Tr}(\phi_{i}\phi_{j}\phi_{k}\phi_{m}) \\ &+ (\lambda_{4}(H^{\dagger}\psi')\Psi + \mathrm{h.c.}) + (\lambda_{5}(H\psi)\Psi + \mathrm{h.c.}) \\ &+ ((\lambda_{6})^{ij}L_{i}\phi_{j}\psi' + \mathrm{h.c.}) \end{aligned}$$

Neutrino mass matrix for two generations of scalars:

$$\frac{1}{32\pi^2} \left( A_1 \begin{pmatrix} (\lambda_6^{e1})^2 & \lambda_6^{e1}\lambda_6^{\mu1} & \lambda_6^{e1}\lambda_6^{\tau1} \\ \lambda_6^{\mu1}\lambda_6^{\tau1} & (\lambda_6^{\mu1})^2 & \lambda_6^{\mu1}\lambda_6^{\tau1} \\ \lambda_6^{e1}\lambda_6^{e1} & \lambda_6^{\mu1}\lambda_6^{\tau1} & (\lambda_6^{\tau1})^2 \end{pmatrix} + A_2 \begin{pmatrix} (\lambda_6^{e2})^2 & \lambda_6^{e2}\lambda_6^{\mu2} & \lambda_6^{e2}\lambda_6^{\tau2} \\ \lambda_6^{\mu2}\lambda_6^{\tau2} & (\lambda_6^{\mu2})^2 & \lambda_6^{\mu2}\lambda_6^{\tau2} \\ \lambda_6^{e2}\lambda_6^{\mu2}\lambda_6^{\pi2} & (\lambda_6^{\tau2})^2 \end{pmatrix} \right)_{13}$$

13 / 24

Radiative seesaw models

Conclusion 0

#### Neutrino masses in T1-3-B

J. Fiaschi, MK, S. May, in preparation



Radiative seesaw models

Conclusion O

#### Relic density in T1-3-B

J. Fiaschi, MK, S. May, in preparation



Conclusion O

#### Class T1-3

S. Fraser, E. Ma, O. Popov, Phys. Lett. B737 (2014) 280; D. Restrepo et al., Phys. Rev. D92 (2015) 013005;
 S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C78 (2018) 88; J. Fiaschi, MK, S. May, in preparation

Model	0	Fermionic	Scalar		Exotic charges	# of N'plets	
Widder		DM	DD	DM	DD	Exotic charges	# Of N piets
T1-3-A	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-B	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	30	$\checkmark$	×	3
T1-3-C	±1	$1_0, 2_{\pm 1}$	✓	21	$\checkmark$	×	4
T1.3.D	1	2 <sub>1</sub> , 3 <sub>0</sub>	<ul> <li>✓</li> </ul>	21	$\checkmark$	×	4
11-3-0	$^{-1}$	$1_0, 2_{-1}, 3_{-2}$	~	$2_{-1}$	$\checkmark$	√	4
T1-3-F	±1	$2_{\pm 1}, 3_0, 3_{\pm 2}$	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	√	4
T1-3-G	0	$2_{\pm 1}, 3_0$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-H	0	$2_{\pm 1}, 3_0$	~	30	$\checkmark$	×	3

Potential in T1-3-A with  $\alpha = 0$ :

$$\mathcal{L}_{Y} = \alpha_{ij}\overline{\psi'}\nu_{jL}^{c}\phi_{i} + \alpha_{ij}\overline{E'}e_{jL}^{c}\phi_{i} + \frac{\beta_{1}}{\sqrt{2}}\overline{\psi}^{c}Sh + \frac{\beta_{2}}{\sqrt{2}}\overline{\psi'}S^{c}h + \text{h.c.}$$

#### Class T1-3

S. Fraser, E. Ma, O. Popov, Phys. Lett. B737 (2014) 280; D. Restrepo et al., Phys. Rev. D92 (2015) 013005;
 S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C78 (2018) 88; J. Fiaschi, MK, S. May, in preparation

Model	0	Fermionic		Scalar		Exotic charges	# of N'plets
Widder		DM	DD	DM	DD	Exotic charges	# Of N piets
T1-3-A	0	$1_0, 2_{\pm 1}$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-B	0	$1_0, 2_{\pm 1}$	~	30	~	×	3
T1-3-C	±1	$1_0, 2_{\pm 1}$	✓	21	$\checkmark$	×	4
T1.3.D	1	2 <sub>1</sub> , 3 <sub>0</sub>	<ul> <li>✓</li> </ul>	21	$\checkmark$	×	4
11-5-0	-1	$1_0, 2_{-1}, 3_{-2}$	$\checkmark$	$2_{-1}$	$\checkmark$	$\checkmark$	4
T1-3-F	±1	$2_{\pm 1}, 3_0, 3_{\pm 2}$	<ul> <li>✓</li> </ul>	$2_{\pm 1}$	$\checkmark$	√	4
T1-3-G	0	$2_{\pm 1}, 3_0$	<ul> <li>✓</li> </ul>	10	$\checkmark$	×	3
T1-3-H	0	$2_{\pm 1}, 3_0$	~	30	~	×	3

Potential in T1-3-A with  $\alpha = 0$ :

$$\mathcal{L}_{Y} = \alpha_{ij}\overline{\psi'}\nu_{jL}^{c}\phi_{i} + \alpha_{ij}\overline{E'}e_{jL}^{c}\phi_{i} + \frac{\beta_{1}}{\sqrt{2}}\overline{\psi}^{c}Sh + \frac{\beta_{2}}{\sqrt{2}}\overline{\psi'}S^{c}h + \text{h.c.}$$

Lepton flavor violation:



16 / 24

Conclusion O

#### Lepton flavor violation in T1-3-A

S. Esch, MK, D. Lamprea, C. Yaguna, Eur. Phys. J. C 78 (2018) 88



#### Class T1-2

S. Esch, MK, C. Yaguna, 1804.03384, JHEP (under review); R. Longas, D. Portillo, D. Restrepo, O. Zapata, JHEP 1603 (2016)

Model	~	Fermionic		Scalar		Exotic charges	# of N'plets
Woder	α	DM	DD	DM	DD	Exotic charges	# Of N piets
T1-2-A	0	$1_0, 2_1$	<ul> <li>✓</li> </ul>	$1_0, 2_1$	$\checkmark$	×	4
11-2-7	-2	2_1	×	2-1	~	×	4
T1.2 B	0	$1_0, 2_1$	<ul> <li>✓</li> </ul>	2 <sub>1</sub> , 3 <sub>0</sub>	$\checkmark$	×	4
11-2-0	-2	2_1	×	$2_{-1}, 3_{-2}$	$\checkmark$	√	4
T1.2 D	1	$2_1, 3_2$	×	21	$\checkmark$	√	4
11-2-0	-1	$2_{-1}, 3_0$	<ul> <li>✓</li> </ul>	$1_0, 2_{-1}$	$\checkmark$	×	4
T1.2 F	1	$2_1, 3_2$	×	$2_1, 3_2$	$\checkmark$	√	4
11-2-1	-1	$2_{-1}, 3_0$	<ul> <li>✓</li> </ul>	$2_{-1}, 3_0$	$\checkmark$	×	4

#### 162 [1511.01873]

-

Radiative seesaw models

#### Class T1-2

S. Esch, MK, C. Yaguna, 1804.03384, JHEP (under review); R. Longas, D. Portillo, D. Restrepo, O. Zapata, JHEP 1603 (2016)

Model	a	Fermio	nic	Scalar		Evotic charges	# of N'plets
woder	u u	DM	DD	DM	DD	Exotic charges	# Of N piecs
Τ1-2-Δ	0	$1_0, 2_1$	√	$1_0, 2_1$	$\checkmark$	×	4
11-2-7	-2	2_1	×	2-1	~	×	4
T1.2 B	0	$1_0, 2_1$	√	$2_1, 3_0$	~	×	4
11-2-0	-2	2_1	×	$2_{-1}, 3_{-2}$	~	$\checkmark$	4
T1.2 D	1	$2_1, 3_2$	×	21	~	√	4
11-2-0	$^{-1}$	$2_{-1}, 3_0$	$\checkmark$	$1_0, 2_{-1}$	$\checkmark$	×	4
T1.2 F	1	$2_1, 3_2$	×	$2_1, 3_2$	$\checkmark$	$\checkmark$	4
1 1-2-1	-1	$2_{-1}, 3_0$	$\checkmark$	$2_{-1}, 3_0$	$\checkmark$	×	4

#### 162 [1511.01873]

Potential in T1-2-A with  $\alpha = 0$ :

$$\begin{aligned} -\mathcal{L}_{\text{scalar}} &= \frac{1}{2}\lambda_{S}\phi_{S}^{2}|H|^{2} + \lambda_{D}|\phi_{D}|^{2}|H|^{2} + \lambda_{D}'|\phi_{D}^{\dagger}H|^{2} \\ &+ \frac{1}{2}\lambda_{D}''\left[\left(\phi_{D}^{\dagger}H\right)^{2} + \text{h.c.}\right] + A\left[\phi_{D}^{\dagger}H\phi_{S} + \text{h.c.}\right] \\ -\mathcal{L}_{\text{fermion}} &= y_{1}\psi_{D_{1}}H\psi_{S} + y_{2}\psi_{D_{2}}H^{\dagger}\psi_{S} + \text{h.c.} \\ -\mathcal{L}_{\text{lepton}} &= g_{1i}L_{i}\phi_{S}\psi_{D_{2}} + g_{2i}L_{i}\phi_{D}\psi_{S} + \text{h.c.} \end{aligned}$$

Radiative seesaw models

#### Direct detection and lepton flavor violation in T1-2-A

S. Esch, MK, C. Yaguna, 1804.03384, JHEP (under review)



Radiative seesaw models

#### Direct detection and lepton flavor violation in T1-2-A

S. Esch, MK, C. Yaguna, 1804.03384, JHEP (under review)



Radiative seesaw models

Conclusion O

#### Class T1-1

C. Boehm, Y. Farzan, T. Hambye, S. Palomares-Ruiz, S Pascoli, Phys. Rev. D77 (2008) 043516;

Model	a	Ferm	ionic	Scalar		Exotic charges	# of N'plets
Woder	α	DM	DD	DM	DD	Exotic charges	# Of N piets
T1-1-A	±2	×	×	$2_{\pm 1}$	$\checkmark$	√	4
11-1-4	0	10	$\checkmark$	$1_0, 2_{\pm 1}$	$\checkmark$	×	3
T1.1.B	±2	$3_{\pm 2}$	×	$2_{\pm 1}$	$\checkmark$	$\checkmark$	4
11-1-0	0	3 <sub>0</sub>	$\checkmark$	$1_0, 2_{\pm 1}$	$\checkmark$	×	3
T1-1-C	±1	$2_{\pm 1}$	×	$1_0, 2_{\pm 1}$	$\checkmark$	$\checkmark$	4
T11D	1	21	×	$1_0, 2_1, 3_2$	$\checkmark$	$\checkmark$	4
11-1-D	$^{-1}$	$2_{-1}$	×	$2_{-1}, 3_0$	$\checkmark$	×	4
T1-1-F	±1	$2_{\pm 1}$	×	$2_{\pm 1}, 3_0, 3_{\pm 2}$	~	$\checkmark$	4
T11C	±2	×	×	$2_{\pm 1}, 3_{\pm 2}$	$\checkmark$	$\checkmark$	4
11-1-0	0	10	~	$2_{\pm 1}, 3_0$	$\checkmark$	×	3
Т1.1.Н	±2	$3_{\pm 2}$	×	$2_{\pm 1}, 3_{\pm 2}$	$\checkmark$	$\checkmark$	4
1 1-1-11	0	30	$\checkmark$	$2_{\pm 1}, 3_0$	$\checkmark$	Х	3

Y. Farzan, Phys. Rev. D 80 (2009) 073009

Radiative seesaw models

Conclusion O

#### Class T1-1

C. Boehm, Y. Farzan, T. Hambye, S. Palomares-Ruiz, S Pascoli, Phys. Rev. D77 (2008) 043516;

Model	α	Fermionic		Scalar		Exotic charges	# of N'plets	
		DM	DD	DM	DD	Exotic charges	# Of N piets	
T1-1-A	±2	×	×	$2_{\pm 1}$	$\checkmark$	$\checkmark$	4	
	0	10	$\checkmark$	$1_0, 2_{\pm 1}$	$\checkmark$	×	3	
T1-1-B	±2	$3_{\pm 2}$	×	$2_{\pm 1}$	~	$\checkmark$	4	
	0	3 <sub>0</sub>	$\checkmark$	$1_0, 2_{\pm 1}$	$\checkmark$	×	3	
T1-1-C	$\pm 1$	$2_{\pm 1}$	×	$1_0, 2_{\pm 1}$	$\checkmark$	$\checkmark$	4	
T1-1-D	1	21	×	$1_0, 2_1, 3_2$	~	$\checkmark$	4	
	$^{-1}$	$2_{-1}$	×	$2_{-1}, 3_0$	~	×	4	
T1-1-F	$\pm 1$	$2_{\pm 1}$	×	$2_{\pm 1}, 3_0, 3_{\pm 2}$	$\checkmark$	$\checkmark$	4	
T1-1-G	±2	×	×	$2_{\pm 1}, 3_{\pm 2}$	~	$\checkmark$	4	
	0	10	$\checkmark$	$2_{\pm 1}, 3_0$	$\checkmark$	×	3	
T1-1-H	±2	$3_{\pm 2}$	×	$2_{\pm 1}, 3_{\pm 2}$	$\checkmark$	~	4	
	0	3 <sub>0</sub>	$\checkmark$	$2_{\pm 1}, 3_0$	$\checkmark$	X	3	

Y. Farzan, Phys. Rev. D 80 (2009) 073009

Potential in T1-1-A with  $\alpha = 0$  ("SLIM Model"):

$$-\mathcal{L} = \lambda_1 |H^T(i\sigma_2)\Phi|^2 + \operatorname{Re}[\lambda_2(H^T(i\sigma_2)\phi)^2] + \lambda_3\eta^2 H^{\dagger}H + \lambda_4 \Phi^{\dagger}\Phi H^{\dagger}H + \frac{\lambda_1'}{2}(\Phi^{\dagger}\Phi)^2 + \frac{\lambda_2'}{2}\eta^4 + \lambda_3'\eta^2 \Phi^{\dagger}\Phi + g_{i\alpha}\bar{N}_i\Phi^{\dagger}L_{\alpha}$$

#### Scalar as Light as MeV - solution to missing satellites?

Y. Farzan, Phys. Rev. D 80 (2009) 073009; A. Arhrib, C. Boehm, E. Ma, T.C. Yuan, JCAP 1604 (2016) 049

Constraints (scalar or fermion DM):

- Unitarity, vacuum stability, non-tachyonic masses
- LEP invisible Z decay,  $\Delta S$  and  $\Delta T$ ; LHC H decay

#### Scalar as LIght as MeV - solution to missing satellites?

Y. Farzan, Phys. Rev. D 80 (2009) 073009; A. Arhrib, C. Boehm, E. Ma, T.C. Yuan, JCAP 1604 (2016) 049

Constraints (scalar or fermion DM):

- Unitarity, vacuum stability, non-tachyonic masses
- LEP invisible Z decay,  $\Delta S$  and  $\Delta T$ ; LHC H decay



#### Scalar as LIght as MeV - solution to missing satellites?

Y. Farzan, Phys. Rev. D 80 (2009) 073009; A. Arhrib, C. Boehm, E. Ma, T.C. Yuan, JCAP 1604 (2016) 049

Constraints (scalar or fermion DM):

- Unitarity, vacuum stability, non-tachyonic masses
- LEP invisible Z decay,  $\Delta S$  and  $\Delta T$ ; LHC H decay



ATLAS (ICHEP 2018):  $R_{\gamma\gamma} = 1.08 \pm 0.08$ ,  $R_{\gamma Z} = 2.7^{+4.6}_{-4.5} < 6.6$ 

Radiative seesaw models

Conclusion 0

#### Models allowing for gauge coupling unification

C. Hagedorn, T. Ohlsson, S. Riad, M.A. Schmidt, JHEP 1609 (2016) 111 [1605.03986]

Model	m	P1	P2	P3	P4	$\Lambda~({\rm GeV})$	$\alpha^{-1}(\Lambda)$	$\begin{array}{c} \frac{\Delta \log_{10}(\Lambda)}{\log_{10}(\Lambda)} \\ \begin{pmatrix} \% \end{pmatrix} \end{array}$	$\frac{\Delta \alpha^{-1}}{\alpha^{-1}}$ (%)
T1-1-D	$1 \\ -1$	$(1, 2, \frac{1}{2})_S$ $(1, 2, -\frac{1}{2})_S$	$(1, 1, 0)_S$ $(1, 1, -1)_S$	$(1, 2, \frac{1}{2})_F$ $(1, 2, -\frac{1}{2})_F$	$(1,3,1)_S$ $(1,3,0)_S$	$\begin{array}{c} 1.3 \cdot 10^{13} \\ 3.1 \cdot 10^{13} \end{array}$	38.4 38.2	7.7 3.2	3.9 1.7
T1-2-A	0	$(1, 1, 0)_F$	$(1, 2, \frac{1}{2})_S$	$(1, 1, 0)_S$	$(1,2,\frac{1}{2})_F$	$5.3\cdot 10^{13}$	39.4	4.1	2.9
T1-2-B	$0 \\ -2$	$(1, 1, 0)_F$ $(1, 1, -1)_F$	$(1, 2, \frac{1}{2})_S$ $(1, 2, -\frac{1}{2})_S$	$(1, 3, 0)_S$ $(1, 3, -1)_S$	$(1, 2, \frac{1}{2})_F$ $(1, 2, -\frac{1}{2})_F$	$\begin{array}{c} 4.6 \cdot 10^{13} \\ 3.2 \cdot 10^{12} \end{array}$	$38.4 \\ 35.9$	$5.6 \\ 0.54$	2.9 0.28
T1-3-A	0	$(1, 1, 0)_F$	$(1,2,\frac{1}{2})_F$	$(1, 1, 0)_S$	$(1,2,-\tfrac{1}{2})_F$	$2.8\cdot 10^{13}$	37.7	6.5	3.3
Т3-А	$0 \\ -2$	$(1, 1, 0)_S$ $(1, 1, -1)_S$	$(1,3,1)_S$ $(1,3,0)_S$	$(1, 2, \frac{1}{2})_F$ $(1, 2, -\frac{1}{2})_F$	-	$\begin{array}{c} 1.6 \cdot 10^{13} \\ 4.0 \cdot 10^{13} \end{array}$	37.3 38.7	4.4 0.21	2.3 0.11
T1-3-A	0	$(1, 1, 0)_F$	$(1,2,\frac{1}{2})_F$	$2 (1,1,0)_S$	-	$6.9\cdot 10^{13}$	39.8	7.4	4.0
T1-3-B	0	$(1, 1, 0)_F$	$(1,2,\frac{1}{2})_F$	$2 \ (1,3,0)_S$	-	$5.7\cdot 10^{13}$	38.9	2.5	1.3



## Conclusion

Minimal dark matter models:

- Bottum-up, small parameter space, often completely testable
- Best motivated, if they solve also other SM problems
- Intriguing: Connection to Higgs, neutrino masses, unification



# Conclusion

Minimal dark matter models:

- Bottum-up, small parameter space, often completely testable
- Best motivated, if they solve also other SM problems
- Intriguing: Connection to Higgs, neutrino masses, unification
- Many models (here only 1-loop), larger parameter spaces
- Coannihilation decouples relic density from direct detection
- Lepton flavor violation provides crucial tests

#### Conclusion

# Conclusion

Minimal dark matter models:

- Bottum-up, small parameter space, often completely testable
- Best motivated, if they solve also other SM problems
- Intriguing: Connection to Higgs, neutrino masses, unification
- Many models (here only 1-loop), larger parameter spaces
- Coannihilation decouples relic density from direct detection
- Lepton flavor violation provides crucial tests
- Many other constraints (indirect detection, LHC)
- Work in progress ...