

Experience with EFT fits to Higgs data

Chris Hays, **Gabija Žemaitytė**

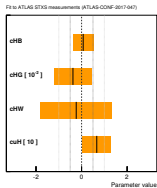
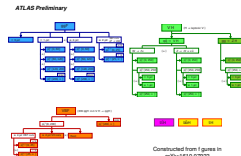
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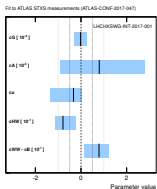
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- Motivation
- ATLAS STXS measurement
- EFT models
- Procedure to relate EFT to measurements
- Fits using HEL and SMEFTsim models
- STXS prospects

Overview map

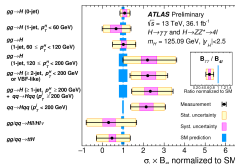


SMEFT



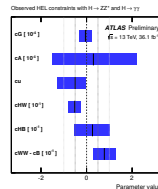
HEL

LHCHXSWG-INT-2017-001



ATLAS-CONF-2017-047

HEL



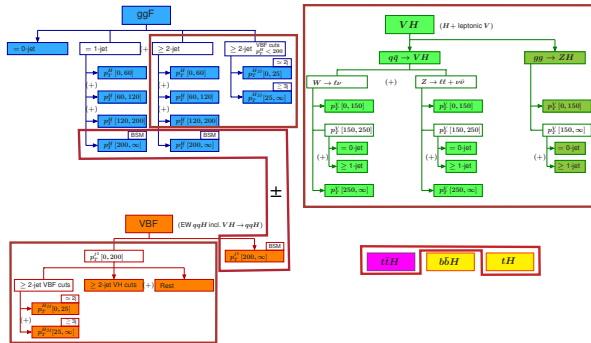
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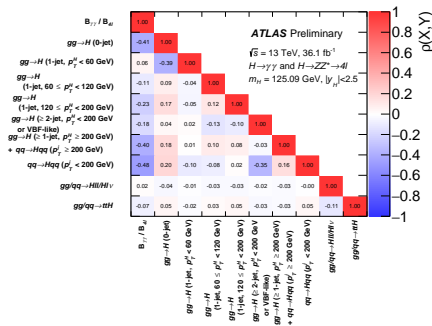
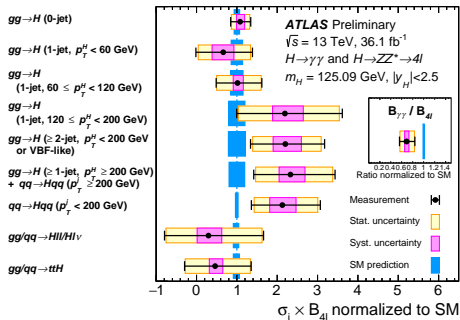
- Prepare experiments for EFT fits by designing strategies and providing tools for fitting Higgs measurements.
 - ▶ Working within the LHC Higgs WG
- Use public EFT implementations to develop parameterizations
- Apply fit procedures to an example set of STXS measurements

- The STXS are an evolution from inclusive production cross section measurements to production cross sections split into a few kinematic regions, which are most sensitive to new physics.

- Measured at the generator level.
- STXS stage 1 binning is on the right. Red boxes show how ATLAS merged the bins in ATLAS-CONF-2017-047.

ATLAS preliminary





■ The value of the difference between $qq \rightarrow Hqq$ $p_T^H \geq 200 \text{ GeV}$ and $gg \rightarrow H$ $p_T^H \geq 200 \text{ GeV}$ is not shown due to the low experimental sensitivity: the fit result is $1.7^{+1.7}_{-1.5} \text{ pb}$ while SM prediction is 0.4 pb .

- Use Madgraph with Feynrules models HEL and SMEFTsim
- Dimension 6 operators, at LO.
- HEL - used as effective Lagrangian:
 - ▶ SILH basis with no four-fermion operators and assuming flavour-universal couplings

A. Alloul, B. Fuks, V. Sanz

- SMEFTsim - complete Warsaw basis implementation:
 - ▶ includes all four-fermion couplings and flavour dependence: 2499 parameters
 - ▶ use $U(3)^5$ symmetry assumption: 76 parameters
 - ▶ two sets of equivalent models for validation

I. Brivio, Y. Jiang, M. Trott

- Parameterize each STXS production bin, partial decay width, and total width:

- Derive interference-only and quadratic-only terms with dedicated MadGraph options:

$$\frac{\sigma}{\sigma_{SM}} = 1 + \sum_i A_i \bar{c}_i + \sum_{ij} B_{ij} \bar{c}_i \bar{c}_j \quad (1)$$

- Include dependence on width:

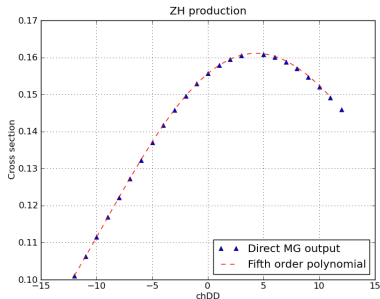
$$\sigma = \sigma_{SM}(\Delta w_p) + A_i(\Delta w_p) c_i \quad (2)$$

$$A(\Delta w) = A(0) + a_p \Delta w_p + b_{pq} \Delta w_p \Delta w_q + \dots \quad (3)$$

$$w = w_{SM} + k_i c_i + t_{ij} c_i c_j \quad (4)$$

$$\Rightarrow \sigma = \sigma_{SM}(0) + a_p k_i^p c_i + A_i(0) c_i \quad (5)$$

Left: derivation of 5th-order dependence on Wilson coefficients, and the linearized dependence [5].

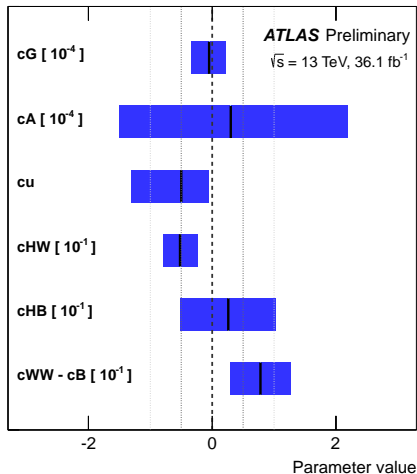


- Fit to an effective Lagrangian (includes quadratic dependence on Wilson coefficients)
- ATLAS: fit to STXS full stage 1 binning, external to ATLAS: fit to merged version of stage 1.
- Operators aimed to be constrained:

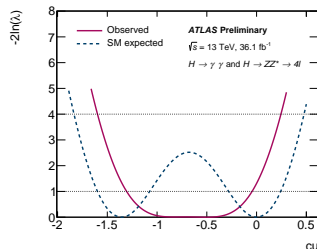
Operator	Expression	HEL coefficient	Vertices
\mathcal{O}_g	$ H ^2 G_{\mu\nu}^A G^{A\mu\nu}$	$cG = \frac{m_W^2}{g_s^2} \bar{c}_g$	Hgg
\mathcal{O}_γ	$ H ^2 B_{\mu\nu} B^{\mu\nu}$	$cA = \frac{m_W^2}{g'^2} \bar{c}_\gamma$	$H\gamma\gamma, HZZ$
\mathcal{O}_u	$y_u H ^2 \bar{u}_l H u_R + \text{h.c.}$	$cu = v^2 \bar{c}_u$	$Ht\bar{t}$
\mathcal{O}_{HW}	$i (D^\mu H)^\dagger \sigma^a (D^\nu H) W_{\mu\nu}^a$	$cHW = \frac{m_W^2}{g} \bar{c}_{HW}$	HW, HZZ
\mathcal{O}_{HB}	$i (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$	$cHB = \frac{m_W^2}{g'} \bar{c}_{HB}$	HZZ
\mathcal{O}_W	$i (H^\dagger \sigma^a D^\mu H) D^\nu W_{\mu\nu}^a$	$cWW = \frac{m_W^2}{g} \bar{c}_W$	HW, HZZ
\mathcal{O}_B	$i (H^\dagger D^\mu H) \partial^\nu B_{\mu\nu}$	$cB = \frac{m_W^2}{g'} \bar{c}_B$	HZZ

- Other parameters are set to zero based on their constraints from other measurements.
- Full HEL equations documented in [LHCHSWG-INT-2017-001](#), [WG2 twiki](#)

Observed HEL constraints with $H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$

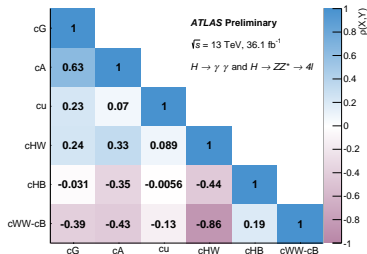
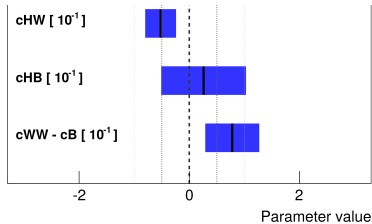


- Quoted central values are those closest to the SM predictions.
- c_u 's observed likelihood has a small local maximum between the two solutions, leading to an asymmetric uncertainty.



Profiled likelihood scan of c_u parameter.

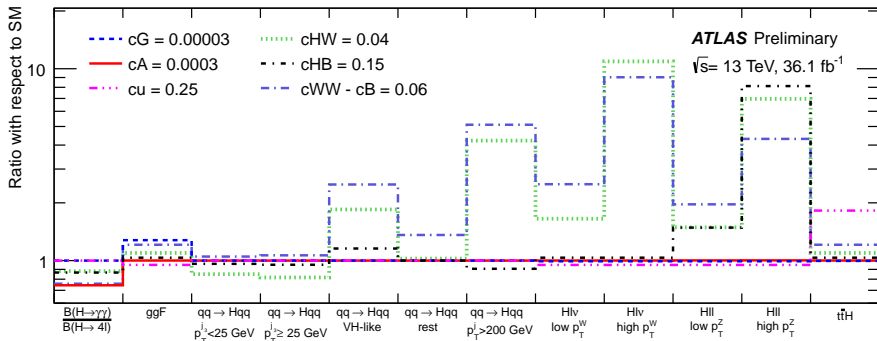
Results: cHW and cWW-cB



- c_{HW} and c_{WW-cB} are strongly anticorrelated.
- Part of the reason - similar impact on STXS bins.

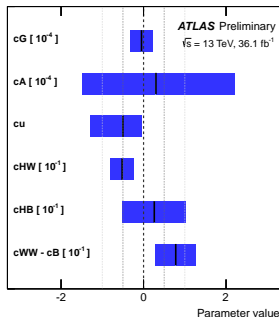
Results: impact plot

- Impact plot: values of STXS regions relative to the SM, for $\approx +1\sigma$ expected SM values.
- cHW and cWW-cB have similar impact.
- Different WH and ZH profile show sensitivity to cHB.

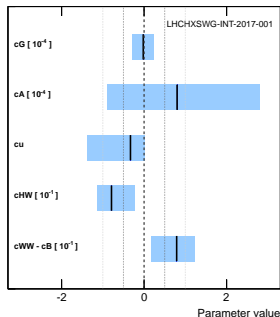


- Later replicated study using STXS measurements.
- Fit to public results loses sensitivity to one parameter (c_{HB} , c_{HW} or c_{WW-cB}) due to WH and ZH merging.
- Other parameters have similar constraints in both studies.

Observed HEL constraints with $H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$



Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



Left: ATLAS-CONF-2017-047
ATLAS

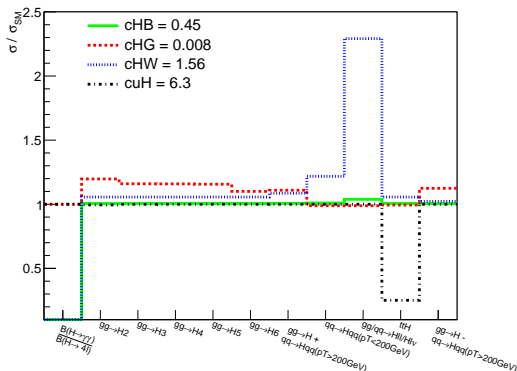
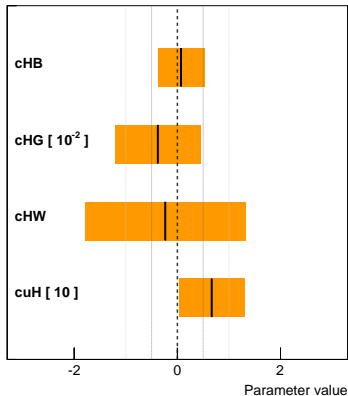
Right: LHCHXSWG-INT-2017-001
C. Hays, V. Sanz,
G. Žemaitytė

- Warsaw basis, with $\Lambda = 1$ TeV.
- Use linear interference terms only.
- External to ATLAS: fit to merged version of STXS stage 1.
- Operators aimed to be constrained:

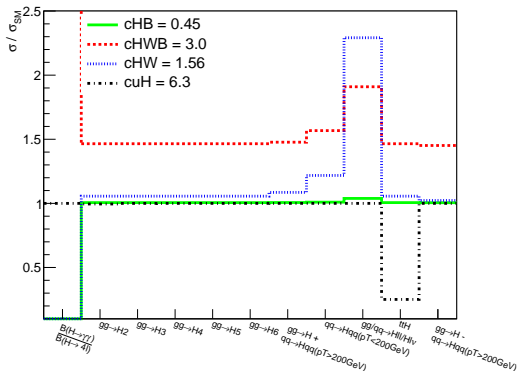
Operator	Expression	SMEFT coefficient	Production/ decay
\mathcal{O}_{HG}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	cHG	ggf ,
\mathcal{O}_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	cHW	$H\gamma\gamma$, HZZ , WH , ZH , VBF
\mathcal{O}_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	cHB	$H\gamma\gamma$, HZZ , ZH , VBF
\mathcal{O}_{HWB}	$H^\dagger H \tau^I W_{\mu\nu}^I B^{\mu\nu}$	$cHWB$	$H\gamma\gamma$, HZZ , ZH , VBF
\mathcal{O}_{uH}	$H^\dagger H \bar{q}_p u_r \tilde{H}$	cuH	ttH , ggf

- Fit to ATLAS public STXS measurements using SMEFTsim.
- Right: values of merged STXS regions for $\approx +1\sigma$ data values. Each parameter has a distinct profile.

Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)

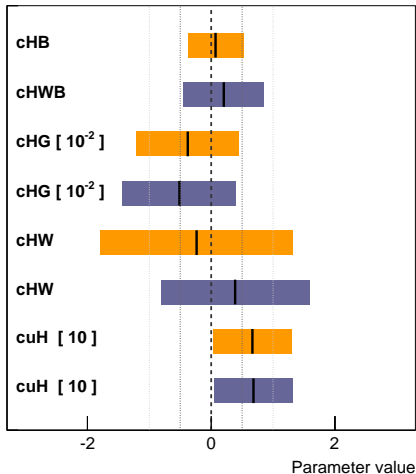


- Based on impact plot that includes $c_{HWB}(=3.0)$, we can see what other four parameter subsets can be fitted.



Replacing cHB by cHWB

Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



- The plot shows two different four parameter fits:
 - ▶ Orange: cG, cHW, cuH, cHB
 - ▶ Violet: cG, cHW, cuH, cHWB
- Similar constraints on cHG and cuH.

- Individual experimental measurements sensitive to isolated parameter set at LO.
- Ideally: each measurement provides the multidimensional likelihood for this parameter set.
- To a good approximation: parameters highly constrained by external measurements can be fixed.
 - ▶ Not a replacement for a rigorous fit but can be useful for studies.
- Fits that constrain the reduced parameter set could quote the dependence on the other parameter values.

The likelihood L is expressed as:

$$-2 \log L = (Ac - D)^T M (Ac - D) \quad (6)$$

where A is the EFT equation matrix, c is the Wilson coefficient vector, D is the vector of data deviation from SM, M - the inverse of covariance matrix.

Let $A'c'$ be EFT matrix and (externally determined) Wilson vector of not fitted Wilson parameters (i.e. the ones that we usually set to 0), then likelihood is:

$$-2 \log L = (Ac + A'c' - D)^T M (Ac + A'c' - D) \quad (7)$$

When we assume that c' is constant vector, then $A'c'$ can be interpreted as corrections to data deviations from SM:

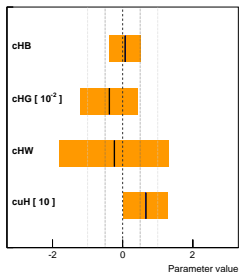
$$D \longrightarrow D - A'c' \quad (8)$$

then solution to (7):

$$c = (A^T M A)^{-1} A^T M (D - A'c') = (A^T M A)^{-1} A^T M D - (A^T M A)^{-1} A^T M A'c' \quad (9)$$

$(A^T M A)^{-1} A^T M D$ - central value, $-(A^T M A)^{-1} A^T M A'$ - prefactors for unfitted coefficients.

Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



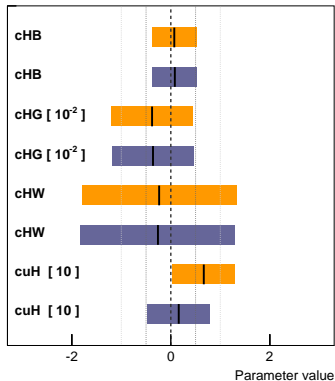
$$cHB = 0.071 + 0.68 \cdot cHWB - 1.5 \cdot cH\beta + ..$$

$$cHG = -0.0038 + 0.1 \cdot cH\beta - 0.1 \cdot cHq\beta + ..$$

$$cHW = -0.24 - 0.19 \cdot cHbox - 0.16 \cdot cHDD + ..$$

$$cuHAbs = 6.7 + 0.84 \cdot cHbox + 0.5 \cdot cHDD + ..$$

Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)

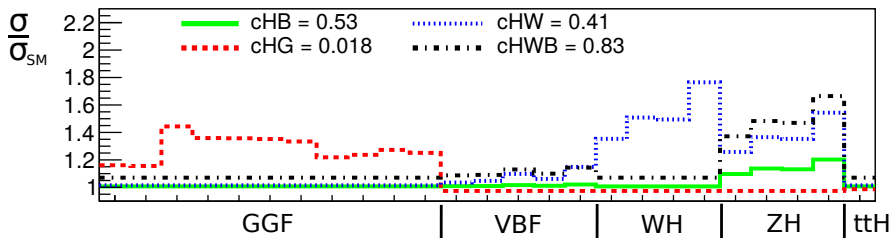


- External constraints give an indication of uncertainties:

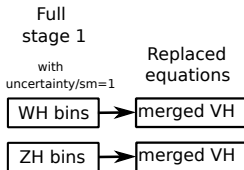
- ▶ **Orange**: other parameters set to 0 vs **violet**: not fitted parameters set to values from a global fit from *1803.03252, J. Ellis, C.W. Murphy, V. Sanz, T. You.*
- ▶ Central values shift insignificantly.

Wilson	$c'_i = 0$	Best-fit c'_i	$c'_i \pm 1\sigma$ (c_i furthest from SM)
cHB	0.071	0.081	0.089
cHW	-0.24	-0.26	-0.28
cHG	-0.0038	-0.0036	-0.0035
cuH	6.7	4.2	1.6

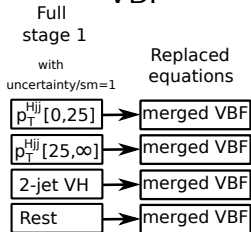
- Let's assume that having full STXS stage 1 we measure bins with uncertainty/sm ≈ 1 .
- ggf bins are not that interesting at LO, VH and VBF bins have potential to constrain 5 parameters (cHW, cHB, cHWB, cHG, cuH):



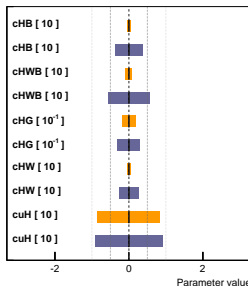
VH



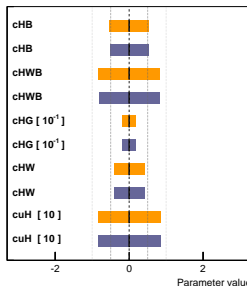
VBF



Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)

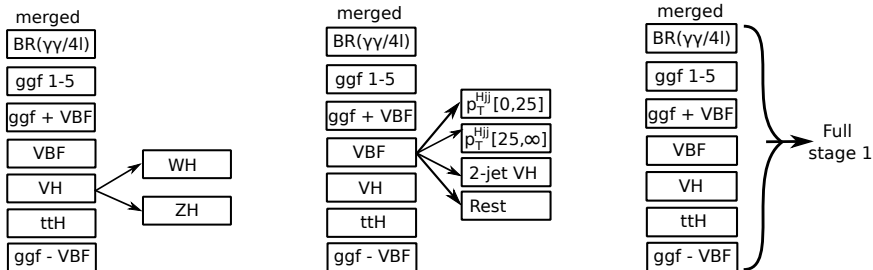


Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



- Assuming uncertainty/sm ≈ 1 .
- Plots: orange - full stage 1; left violet: WH, ZH equations replaced by merged VH; right violet: VBF equations are replaced by VBF merged.
- VBF granularity is not important - results have not changed much.
- Using VH merged equations give much larger uncertainties.

- ATLAS public STXS measurement is sensitive to 4 parameters from SMEFTsim, how can we recover the 5th?
- To predict what ATLAS internal would give - split bins assuming uncertainties are $\sim \frac{1}{\sqrt{N}}$.
- None of the splittings below gave sensitivity to the 5th parameter - full granularity in ATLAS probably would not be sensitive.



- With current ATLAS STXS measurement we can constrain:
 - ▶ 6 HEL parameters inside ATLAS, 5 - external to ATLAS.
 - ▶ 4 SMEFTsim parameters external to ATLAS.
- Equations with unfit parameters can be used to estimate their impact.
- Full Higgs STXS stage 1 $H \rightarrow 4l$ and $H \rightarrow \gamma\gamma$ measurements have the capacity to constrain 5 SMEFT coefficients, and for the improvement the most important productions are WH and ZH.

THANK YOU

BACKUP

$$\begin{aligned} c_{HB} = & 0.071 + 1.8e-08 \cdot c_H + 0.053 \cdot c_{Hbox} + 0.045 \cdot c_{HDD} + 0.68 \cdot \\ & c_{HWB} + 9.4e-10 \cdot c_{eHAbs} - 0.0001 \cdot c_{eWAbs} - 3.1e-07 \cdot c_{eBAbs} + \\ & 0.039 \cdot c_{Hl1} - 1.5 \cdot c_{Hl3} - 0.015 \cdot c_{He} + 0.062 \cdot c_{ll} + 1.2 \cdot c_{ll1} - 0.00016 \cdot \\ & c_{uWAbs} + 2.6e-06 \cdot c_{uBAbs} - 8.3e-05 \cdot c_{dWAbs} - 4.5e-05 \cdot \\ & c_{dBAbs} + 0.01 \cdot c_{Hq1} + 2 \cdot c_{Hq3} + 0.075 \cdot c_{Hu} - 0.034 \cdot c_{Hd} + 3.1e- \\ & 08 \cdot c_{dHAbs} - 3e-05 \cdot c_{uGAbs} - 2.9e-09 \cdot c_{dGAbs} - 6.3e-13 \cdot \\ & c_{HudAbs} + 0.0022 \cdot c_G + 4.1e-17 \cdot c_{quqd1Abs} - 1.4e-17 \cdot c_{quqd8Abs} \end{aligned}$$

$$\begin{aligned} c_{HW} = & -0.24 - 6.4e-08 \cdot c_H - 0.19 \cdot c_{Hbox} - 0.16 \cdot c_{HDD} - 0.51 \cdot \\ & c_{HWB} - 3.4e-09 \cdot c_{eHAbs} + 0.00036 \cdot c_{eWAbs} + 1e-06 \cdot c_{eBAbs} - \\ & 0.15 \cdot c_{Hl1} + 5.1 \cdot c_{Hl3} + 0.051 \cdot c_{He} - 0.22 \cdot c_{ll} - 4 \cdot c_{ll1} + 0.00056 \cdot \\ & c_{uWAbs} - 9e-06 \cdot c_{uBAbs} + 0.00029 \cdot c_{dWAbs} + 0.00015 \cdot \\ & c_{dBAbs} - 0.036 \cdot c_{Hq1} - 7 \cdot c_{Hq3} - 0.26 \cdot c_{Hu} + 0.12 \cdot c_{Hd} - 8.8e- \\ & 08 \cdot c_{dHAbs} + 8.5e-05 \cdot c_{uGAbs} + 1e-08 \cdot c_{dGAbs} + 2.2e-12 \cdot \\ & c_{HudAbs} - 0.0072 \cdot c_G - 1.2e-16 \cdot c_{quqd1Abs} + 4e-17 \cdot c_{quqd8Abs} \end{aligned}$$

$$\begin{aligned} c_{HG} = & -0.0038 - 1.4e-09 \cdot c_H + 0.00033 \cdot c_{Hbox} - 0.0041 \cdot c_{HDD} - 0.0057 \cdot \\ & c_{HWB} - 7.5e-11 \cdot c_{eHAbs} + 7.9e-06 \cdot c_{eWAbs} + 2.5e-08 \cdot c_{eBAbs} - \\ & 0.0038 \cdot c_{Hl1} + 0.1 \cdot c_{Hl3} + 0.0009 \cdot c_{He} - 0.0034 \cdot c_{ll} - 0.083 \cdot c_{ll1} + \\ & 1.2e-05 \cdot c_{cuWAbs} - 1.9e-07 \cdot c_{cuBAbs} + 6.7e-06 \cdot c_{cdWAbs} + 3.6e- \\ & 06 \cdot c_{cdBAbs} - 0.0027 \cdot c_{Hq1} - 0.1 \cdot c_{Hq3} - 0.0012 \cdot c_{Hu} + 0.00095 \cdot c_{Hd} + \\ & 2e-07 \cdot c_{cdHAbs} - 0.00019 \cdot c_{cuGAbs} + 3.2e-10 \cdot c_{cdGAbs} + 1.4e-14 \cdot \\ & c_{HudAbs} - 0.00029 \cdot c_G + 2.6e-16 \cdot c_{quqd1Abs} - 8.8e-17 \cdot c_{quqd8Abs} \end{aligned}$$

$cuHAbs =$

$$6.7 + 2.6e^{-07} \cdot cH + 0.84 \cdot cHbox + 0.5 \cdot cHDD + 1 \cdot cHWB + 1.4e^{-08} \cdot ceHAbs - 0.0014 \cdot ceWAbs - 4.5e^{-06} \cdot ceBAbs + 0.68 \cdot cHI1 - 20 \cdot cHI3 - 0.16 \cdot cHe + 0.89 \cdot cII + 15 \cdot cII1 - 0.0022 \cdot cuWAbs + 3.4e^{-05} \cdot cuBAbs - 0.0012 \cdot cdWAbs - 0.00066 \cdot cdBAbs + 0.49 \cdot cHq1 + 19 \cdot cHq3 + 0.22 \cdot cHu - 0.18 \cdot cHd - 1.7e^{-06} \cdot cdHAbs - 7.1 \cdot cuGAbs - 4.1e^{-08} \cdot cdGAbs - 1.6e^{-12} \cdot cHudAbs + 1.3 \cdot cG + 9.9e^{-12} \cdot cquqd1Abs - 3.4e^{-12} \cdot cquqd8Abs$$

Correlation matrix:

	cHB	cHW	cHG	cuHAbs
cHB	1	-1	0.05	0.01
cHW	-1	1	-0.04	-0.01
cHG	0.05	-0.04	1	-0.21
cuHAbs	0.01	-0.01	-0.21	1

- Most efficient derivation of equation is using substitution:

$$\sigma = \sigma_{SM}(\Delta w_p) + A_i(\Delta w_p)c_i + B_{ij}(\Delta w_p)c_i c_j \quad (10)$$

$$A(\Delta w) = A(0) + a_p \Delta w_p + b_{pq} \Delta w_p \Delta w_q \quad (11)$$

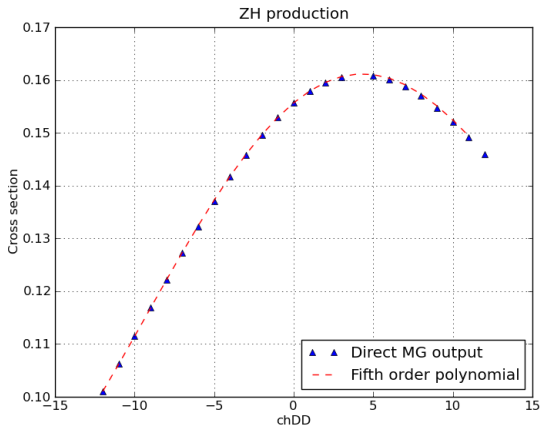
$$w = w_{SM} + k_i c_i + t_{ij} c_i c_j \quad (12)$$

$$\sigma_{SM}(\Delta w) = \sigma_{SM}(0) + r_p \Delta w_p + s_{pq} \Delta w_p \Delta w_q \quad (13)$$

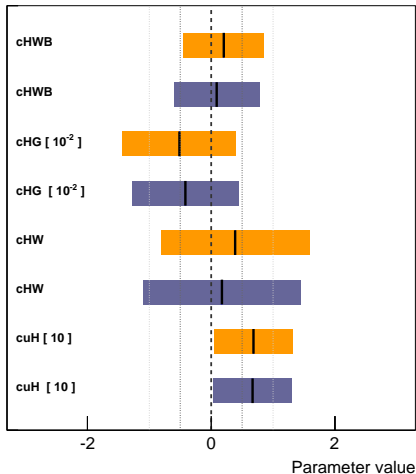
Then 10 becomes:

$$\begin{aligned} \sigma &= \sigma_{SM}(0) \left(1 + \frac{r_p}{\sigma_{SM}(0)} \Delta w_p \right) + A_i(0) c_i \\ &= \sigma_{SM}(0) \left(1 + \frac{r_p}{\sigma_{SM}(0)} k_i^p c_i \right) + A_i(0) c_i \\ &= \sigma_{SM}(0) + r_p k_i^p c_i + A_i(0) c_i \end{aligned} \quad (14)$$

- Validation for ZH production for chDD coefficient. Red dashed - from the equations.
- Advantage: we have only four widths which need to be calculated once.



Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



- Only cHWB (from our 5 param set) affects particle widths.
- Orange: width non constant, violet: width constant SM value.

- Substituted results from 1803.03252 (Ellis et al.)
- There Higgs STXS measurement were included, so substituted results should be consistent.

Wilson	Other=0	1σ closer to SM	Central values	1σ further from SM	Other=1
cHB	0.071	0.073	0.081	0.089	2.7
cHW	-0.24	-0.24	-0.26	-0.28	-7.5
cHG	-0.0038	-0.0038	-0.0036	-0.0035	-0.11
cuH	6.7	6.8	4.2	1.6	19