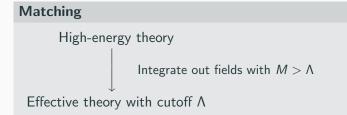
MatchingTools: a Python library for symbolic effective field theory calculations

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Introduction



Removing redundancies

- Group theory identities.
- Integration by parts.
- Field redefinitions (EOMs for the light fields).

From an action:

$$S(\phi,\Phi) = -\int \Phi^{\dagger} \mathcal{D} \Phi + S_{int}(\phi,\Phi)$$

Iterative solution of the EOMs ($1/M_{\Phi}$ expansion):

$$\Phi = \mathcal{D}^{-1} \frac{\delta S_{in}}{\delta \Phi^{\dagger}} (\phi, \Phi) \quad \longrightarrow \quad \Phi_c(\phi)$$

Effective action:

 $S_{\text{eff}}(\phi) = S(\phi, \Phi_c(\phi))$

Complete tree-level dictionary [1711.10391]

Extensions with new particles \longleftrightarrow Dimension-6 SMEFT

We consider a general SM extension with:

- $SU(3) \times SU(2) \times U(1)$ -invariance.
- SM fields + all new ones with contributions to the dimension-6 SMEFT.

and integrate out all the new fields.

Size of the problem:

- 47 new fields (apart from the SM ones).
- The interaction Lagrangian contains hundreds of terms.
- After integration, before simplifying, thousands of terms.

MatchingTools helps reducing the possibility of errors and the time of the calculations.

Outline

1	Overview	Organization, basic tools, matching, extras.
2	Toy example	Integrating out a vector-like quark in a toy model.
3	Application	From general extensions of the SM to the dim-6 SMEFT
4	Links	Repository, installation.
5	Conclusions	Future work, summary.

Overview

matchingtools

.core	The tools for defining a model and the basic simbolic tensor algebraic operations.
.integration	The classes to define heavy fields and the function integrate.
.transformations	The functions apply_rules, simplify,
.output	The class Writer, which provides methods for representing results in plain text or LaTeX
.extras	Package with tools for SMEFT applications.

Basic objects

- Lagrangians: sums of terms (operators).
- Operators: products of tensors with arbitrary index contractions.

Basic operation

Search and replace a pattern in each term. Examples:

- Substitute heavy field by its EOM solution.
- Group theory identities (e.g. $\sigma^a_{ij}\sigma^a_{kl} \rightarrow 2\delta_{il}\delta_{kj} \delta_{ij}\delta_{kl}$)

Tensors (fields) with some indices:

tensor(index_label_1, index_label_2, ...)

Operators (products of tensors):

Op(tensor1(i1, i2, ...), tensor2(j1, ...), ...)

Lagrangians (and other sums of operators) as:

OpSum(operator_1, operator_2, ...)

Repeated index labels

Non-negative integers repeated exactly twice inside each operator to indicate contraction.

$$\mathcal{L}_{example} = \phi_{ij}\psi_{iab}F_{jab} + X_{mn}Y^{nm}$$

Lexample = OpSum(
 Op(phi(0, 1), psi(0, 2, 3), F(1, 2, 3)),
 Op(X(0, 1), Y(1, 0)))

Free index labels

Negative integers. Used in substitution rules: match the ones in the pattern with the ones in the replacement.

$$\sigma^{a}_{ij}\sigma^{a}_{kl} \rightarrow 2\delta_{il}\delta_{kj} - \delta_{ij}\delta_{kl}$$

SU2Fierz = (
 Op(sigma(0, -1, -2), sigma(0, -3, -4)),
 OpSum(
 2 * Op(kdelta(-1, -4), kdelta(-3, -2)),
 - Op(kdelta(-1, -2), kdelta(-3, -4))))

Notation for (covariant) derivatives:

D(vector_index_label, field(i1, i2, ...))

Example:

$$\mathcal{O} = V^{\mu}\phi_i D_{\mu}\phi_i$$

O = Op(V(0), phi(1), D(0, phi(1)))

Tree-level matching in any Lorentz invariant field theory.

Heavy fields

- Predefined: scalars, Dirac or Majorana fermions and vectors.
- Other kinds of heavy fields can be added by the user.

The user specifies (unrestricted in priciple):

- Order in 1/M for the solution of the EOMs.
- Max. dim. for the operators in the effective Lagrangian.

extras

.SU2	Common tensors and identities for $SU(2)$.
.SU3	Common tensors and identities for $SU(3)$.
.Lorentz	Common tensors and identities for the Lorentz group.
.SM	The SM fields and their equations of motion.
.SM_dim_6_basis	Definition of a basis for the dimension-6 SMEFT [arXiv:1412.8480]. (Other bases will be included soon in other modules)

- 1. Define fields and coupling constants.
- 2. Define the interaction Lagrangian.
- 3. Specify which fields are heavy.
- 4. Integrate out the heavy fields.
- 5. Define substitution rules **rewrite** the effective Lagrangian and apply them.
- [Define LaTeX representation of the coupling constants and Wilson coefficients and output to a .tex file].

Toy example

Consider an extension of the SM with vector-like quark doublet Q of hypercharge 7/6 and interaction Lagrangian:

$$\mathcal{L}_{\rm int} = -(y_Q)_i \bar{Q}_L \phi u_{Ri} + {\rm h.c.}$$

When integrated out, it'll give contributions to the effective Lagrangian as $\sim \bar{u}_R \phi^{\dagger} \not{D}(\phi u_R) \sim \bar{u}_R \phi^{\dagger} (\not{D} \phi) u_R + \bar{u}_R \phi^{\dagger} \phi (\not{D} u_R)$.

We can then use the EOM of u_R to write the result in terms of:

$$\mathcal{O}_{\phi u} = \bar{u}_R \phi^{\dagger}(\not{D}\phi) u_R$$

 $\mathcal{O}_{u\phi} = \bar{q}_L \tilde{\phi} u_R \phi^{\dagger} \phi.$

Definition of the model and matching

```
phi = FieldBuilder('phi', 1, boson)
phic = FieldBuilder('phic', 1, boson)
uR = FieldBuilder('uR', 1.5, fermion)
uRc = FieldBuilder('uRc', 1.5, fermion)
QL = FieldBuilder('QL', 1.5, fermion)
QLc = FieldBuilder('QLc', 1.5, fermion)
QR = FieldBuilder('QR', 1.5, fermion)
QRc = FieldBuilder('ORc', 1.5, fermion)
yQ = TensorBuilder('yQ')
vQc = TensorBuilder('vQc')
Lint = -OpSum(
    Op(vQ(0), QLc(1, 2, 3), phi(3), uR(1, 2, 0)),
    Op(vQc(0), uRc(1, 2, 0), phic(3), QL(1, 2, 3)))
heavy_Q = VectorLikeFermion(
    'Q', 'QL', 'QR', 'QLc', 'QRc', 3, has flavor=False)
Leff = integrate([heavy_Q], Lint, 6)
Leff writer = Writer(Leff, {})
print(Leff writer)
```

Rewriting operators and defining a basis

```
isigma2 = TensorBuilder("isigma2")
yu = TensorBuilder("yu"); yu_dagger = TensorBuilder("yu_dagger")
qL = FieldBuilder("qL", 1.5, fermion); qLc = FieldBuilder("qLc", 1.5, fermion)
rules_uR_eom = [
    (Op(sigma4(0, -1, 1), D(0, uR(1, -2, -3))),
     OpSum(Op(yu_dagger(-3, 0), isigma2(1, 2), phi(2), qL(-1, -2, 1, 0)))),
    (Op(sigma4(0, 1, -1), D(0, uRc(1, -2, -3))),
     OpSum(Op(yu(0, -3), isigma2(1, 2), phic(2), qLc(-1, -2, 1, 0))))]
Ophiu = flavor_tensor_op("Ophiu"); Ophiuc = flavor_tensor_op("Ophiuc")
Ouphi = flavor_tensor_op("Ouphi"); Ouphic = flavor_tensor_op("Ouphic")
rules ops = [
    (Dp(uRc(0, 1, -1), phic(2), sigma4(3, 0, 4), D(3, phi(2)), uR(4, 1, -2)),
     OpSum(Ophiu(-1, -2))),
    (Dp(uRc(0, 1, -2), D(3, phic(2)), sigma4(3, 0, 4), phi(2), uR(4, 1, -1)),
     OpSum(Ophiuc(-1, -2))),
    (Op(qLc(0, 1, 2, -1), isigma2(2, 3), phic(3), uR(0, 1, -2), phic(4), phi(4)),
     OpSum(Ouphi(-1, -2))),
    (Op(uRc(0, 1, -2), qL(0, 1, 2, -1), isigma2(2, 3), phi(3), phic(4), phi(4)),
     OpSum(Ouphic(-1, -2)))]
Lfinal = apply rules(Leff, rules uR eom + rules ops, 1)
```

```
Lfinal_writer = Writer(Lfinal, ["Ophiu", "Ophiuc", "Ouphi", "Ouphic"])
print(Lfinal_writer)
latex_structures = {
    "yQ": "(y_Q)_{}", "yQc": "(y_Q)^*_{}",
    "yu": "(y_u)_{{{}}", "yQc": "(y_Q)^*_{}",
    "yQ": "(y_u)_{{{}}", "yu_dagger": "(y_u)^\dagger_{{{}}",
    "yQ": "M_Q"}
latex_ops = {
    "Ophiu": r"(C_{{\phi u}})_{{{}}",
    "Ophiuc": r"(C_{{\phi u}})^*_{{{}}",
    "Ouphi": r"(C_{{\phi u}})^*_{{}}",
    "Ouphi": r"(C_{{\u03c6}},
    "yl])^*_{{{}}",
    "Ouphic": r"(C_{{u \phi}})_{{}}",
    "Ouphic": r"(C_{{u \phi}})^*_{{}}",
    "Uuphic": r"(C_{{}}"(u \phi))^*_{{}}",
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```

Output

The final LaTeX output of the program is:

$$(C_{\phi u})_{ij} = + \frac{i(y_Q)_i^*(y_Q)_j}{2(M_Q)^2}$$
$$(C_{\phi u})_{ij}^* = -\frac{i(y_Q)_i(y_Q)_j^*}{2(M_Q)^2}$$
$$(C_{u\phi})_{ij} = -\frac{i(y_u)_{ik}(y_Q)_j(y_Q)_k^*}{2(M_Q)^2}$$
$$(C_{u\phi})_{ij}^* = +\frac{i(y_u)_{ki}^*(y_Q)_j^*(y_Q)_k}{2(M_Q)^2}$$

An application to a complex case

Complete tree-level dictionary [1711.10391]

Extensions with new particles \longleftrightarrow Dimension-6 SMEFT

Size:

- $\bullet~\sim 50$ multiplets to integrate out.
- 1000 10000 terms in some intermediate steps.

MatchingTools takes **less than a minute** in a i5 to do the complete calculation.

UV/IR dictionary: bottom-up

$$\begin{split} (C_{le})_{ijkl} &= -\frac{(y_{\varphi}^{e})_{rkl}^{*}(y_{\varphi}^{e})_{rkj}}{2M_{\varphi_{r}}^{2}} - \frac{(g_{B}^{e})_{rkl}(g_{B}^{l})_{rij}}{M_{B_{r}}^{2}} + \frac{(g_{L_{3}})_{rki}^{*}(g_{L_{3}})_{rlj}}{M_{L_{3r}}^{2}} \\ &- \frac{\tilde{y}_{\ell_{1}}^{e*}(\delta_{L_{1}\varphi})_{sr}(\gamma_{L_{1}})_{s}^{*}(y_{\varphi}^{e})_{rkj}}{2M_{\varphi_{r}}^{2}M_{L_{1s}}^{2}} - \frac{\tilde{y}_{\ell_{3}}^{e}(\delta_{L_{1}\varphi})_{sr}^{*}(\gamma_{L_{1}})_{s}(y_{\varphi}^{e})_{rli}}{2M_{\varphi_{r}}^{2}M_{L_{1s}}^{2}} \\ &- \frac{\tilde{y}_{\ell_{3}}^{e}\tilde{y}_{l}^{i*}(\delta_{L_{1}\varphi})_{is}(\gamma_{L_{1}})_{t}(\delta_{L_{1}\varphi})_{rs}(\gamma_{L_{1}})_{r}^{*}}{2M_{L_{1r}}^{2}M_{Z_{1r}}^{2}} \\ &+ \frac{1}{f} \left\{ \frac{\tilde{y}_{\ell_{1}}^{e*}(\tilde{y}_{L}^{eD})_{rkj}(\gamma_{L_{1}})_{r}^{*}}{4M_{L_{1r}}^{2}} + \frac{\tilde{y}_{\ell_{1}}^{e*}(\tilde{y}_{L}^{DD})_{rkj}(\gamma_{L_{1}})_{r}^{*}}{4M_{L_{1r}}^{2}} \\ &+ \frac{\tilde{y}_{\ell_{3}}^{e}(\tilde{y}_{L}^{eDD})_{rkj}(\gamma_{L_{1}})_{r}}{4M_{L_{1r}}^{2}} + \frac{\tilde{y}_{\ell_{3}}^{e*}(\tilde{y}_{L}^{DD})_{rkj}(\gamma_{L_{1}})_{r}}{4M_{L_{1r}}^{2}} \right\}, \end{split}$$

$$(C_{ld})_{ijkl} = -\frac{(y_{\Pi_{1}})_{rjk}^{*}(y_{\Pi_{1}})_{ril}}{2M_{\Pi_{1r}}^{2}} - \frac{(g_{B}^{d})_{rkl}(g_{B}^{l})_{rij}}{M_{B_{r}}^{2}} + \frac{(g_{Q_{3}}^{d})_{rki}^{*}(g_{Q_{3}}^{d})_{rlj}}{M_{Q_{0r}}^{2}}, \end{cases}$$

. . .

UV/IR dictionary: top-down

Fields	Operators
S	$\mathcal{O}_{\phi4}, \mathcal{O}_{\phi}, \mathcal{O}_{\phi\Box}, \mathcal{O}_{\phiB}, \mathcal{O}_{\phi\tilde{B}}, \mathcal{O}_{\phiW}, \mathcal{O}_{\phi\tilde{W}}, \mathcal{O}_{\phi G}, \mathcal{O}_{\phi\tilde{G}}, \mathcal{O}_{e\phi}, \mathcal{O}_{d\phi}, \mathcal{O}_{u\phi}$
\mathcal{S}_1	\mathcal{O}_{ll}
S_2	\mathcal{O}_{ee}
φ	$\mathcal{O}_{le}, \mathcal{O}_{qu}^{(1)}, \mathcal{O}_{qu}^{(8)}, \mathcal{O}_{qd}^{(1)}, \mathcal{O}_{qd}^{(8)}, \mathcal{O}_{ledq}, \mathcal{O}_{quqd}^{(1)}, \mathcal{O}_{lequ}^{(1)}, \mathcal{O}_{\phi}, \mathcal{O}_{e\phi}, \mathcal{O}_{d\phi}, \mathcal{O}_{u\phi}$
Ξ	$\mathcal{O}_{\phi 4}, \mathcal{O}_{\phi}, \mathcal{O}_{\phi D}, \mathcal{O}_{\phi \Box}, \mathcal{O}_{\phi WB}, \mathcal{O}_{\phi W\tilde{B}}, \mathcal{O}_{e\phi}, \mathcal{O}_{d\phi}, \mathcal{O}_{u\phi}$
Ξ_1	$\mathcal{O}_{\phi4}, \mathcal{O}_5, \mathcal{O}_{ll}, \mathcal{O}_{\phi}, \mathcal{O}_{\phi D}, \mathcal{O}_{\phi \Box}, \mathcal{O}_{e\phi}, \mathcal{O}_{d\phi}, \mathcal{O}_{u\phi}$
Θ_1	\mathcal{O}_{ϕ}
Θ_3	\mathcal{O}_{ϕ}
ω_1	$\mathcal{O}_{qq}^{(1)}, \mathcal{O}_{qq}^{(3)}, \mathcal{O}_{lq}^{(1)}, \mathcal{O}_{lq}^{(3)}, \mathcal{O}_{eu}, \mathcal{O}_{ud}^{(1)}, \mathcal{O}_{ud}^{(8)}, \mathcal{O}_{quad}^{(1)}, \mathcal{O}_{quad}^{(8)},$
	$\mathcal{O}_{lequ}^{(1)} \mathcal{O}_{lequ}^{(3)}, \mathcal{O}_{duq}, \mathcal{O}_{qqu}, \mathcal{O}_{qqq}, \mathcal{O}_{duu}$
ω_2	\mathcal{O}_{dd}
ω_4	$\mathcal{O}_{uu},\mathcal{O}_{ed},\mathcal{O}_{duu}$
Π_1	\mathcal{O}_{ld}
Π_7	$\mathcal{O}_{lu}, \mathcal{O}_{qe}, \mathcal{O}_{lequ}^{(1)}, \mathcal{O}_{lequ}^{(3)}$
ς	$\mathcal{O}_{qq}^{(1)},\mathcal{O}_{qq}^{(3)},\mathcal{O}_{lq}^{(1)},\mathcal{O}_{lq}^{(3)},\mathcal{O}_{qqq}$
Ω_1	$\begin{array}{l} \mathcal{O}_{lu}, \mathcal{O}_{qe}, \mathcal{O}_{lu}^{(1)}, \mathcal{O}_{qq}^{(3)}, \\ \mathcal{O}_{ql}^{(1)}, \mathcal{O}_{ql}^{(3)}, \mathcal{O}_{ll}^{(1)}, \mathcal{O}_{ld}^{(3)}, \mathcal{O}_{qqq} \\ \mathcal{O}_{ql}^{(2)}, \mathcal{O}_{ql}^{(3)}, \mathcal{O}_{ul}^{(1)}, \mathcal{O}_{ul}^{(3)}, \mathcal{O}_{uld}^{(1)}, \mathcal{O}_{qqd}^{(8)} \end{array}$
Ω_2	\mathcal{O}_{dd}
Ω_4	\mathcal{O}_{uu}
Υ	$\mathcal{O}_{qq}^{(1)},\mathcal{O}_{qq}^{(3)}$
Φ	$\mathcal{O}_{qu}^{(1)},\mathcal{O}_{qu}^{(8)},\mathcal{O}_{qd}^{(1)},\mathcal{O}_{qd}^{(8)},\mathcal{O}_{quqd}^{(8)}$

+ fermions, vectors

Links and installation

GitHub repository:

https://github.com/jccriado/matchingtools

Available at PyPI:

pip install matchingtools

Documentation:

http://matchingtools.readthedocs.io/en/latest/ arXiv:1710.06445

Future work and conclusions

Future work:

- Include more application-specific tools in extras.
- Connection with other software.

With MatchingTools we can automatize the process of:

- Tree-level matching.
- Rewritting the effective Lagrangian in a chosen basis.

This lets us reduce the possibility of errors and the time it takes to do the calculations.