# Supplemental Assignments for Collider Phenomenology

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### Lecture 1: Introduction and Basic Formalism

**Exercise 1.1:** A *B*-factor (e. g. KEKB) is designed for asymmetric head-on collisions between a positron beam of energy 3.5 GeV and an electron beam of energy 8 GeV. Find the center-of-mass energy for the *B*-factory. Do you understand why to adopt this design for the energy and for the asymmetry?

**Exercise 1.2:** The dominant decay channel of the top quark is  $t \to W^+b$ . The partial decay width given in terms of the known mass parameters at the leading order is

$$\Gamma_t = \frac{G_F m_t^3}{8\pi\sqrt{2}} (1 - \frac{m_W^2}{m_t^2})^2 (1 + 2\frac{m_W^2}{m_t^2}).$$

Assuming this formula gives its total decay width, estimate the top-quark life-time in units of yocto-second.

If the QCD scale is  $\Lambda_{\rm QCD} \approx 200$  MeV, compare the top-quark life-time with the time scale at which the QCD strong interaction sets in.

Also compare with the *b*-quark life-time, and try to understand the differences between the decays of the two quarks.

(Use the PDG review for the parameters needed.)

Exercise 1.3: (challenging problem) In the "Standard Model" of elementary particle physics, the amplitude for the scattering of the (longitudinally polarized) weak gauge bosons (the force mediator for the nuclear  $\beta$  decay)  $W^+W^+ \rightarrow W^+W^+$  is calculated at high energies to be

$$f(k, \ \theta) = \frac{1}{16\pi k} \left(\frac{-M_H^2}{v^2}\right) \left(\frac{t}{t - M_H^2} + \frac{u}{u - M_H^2}\right)$$

where k is the  $W^+$  momentum in the Center-of-Momentum frame,  $M_H$  is the mass of the Higgs boson, and  $v \approx 250$  GeV is the Higgs vacuum expectation value. The angular-dependent kinematical variables are

$$t = -2k^2(1 - \cos\theta)$$
 and  $u = -2k^2(1 + \cos\theta)$ .

Note that the amplitude is give in the "natural units" where  $c = \hbar = 1$ , and everything is expressed in terms of the energy units electron-volts: 1 GeV =  $10^9$  eV.

(a). Take the high-energy limit  $2k \gg M_H$ , compute the partial wave amplitude  $a_\ell$ . Note that for final state identical particles  $W^+W^+$ , the angular integration should be  $1/2 \int_{-1}^1 d\cos\theta$ .

(b). Impose the partial wave unitarity condition on  $a_{\ell}$  for *s*-wave, determine the bound on the mass of the Higgs boson  $M_H$  (in units of GeV).

(c). If the Higgs boson did not exist in Nature, then the amplitude for the weak gauge boson scattering for  $W^+W^+ \to W^+W^+$  would be expressed by taking the limit  $2k \ll M_H \to \infty$ . Using the same procedure above, determine at what energy scale 2k the Standard Model theory would break down to violate the partial wave unitarity.

(Remark: The "Large Hadron Collider" (LHC) at CERN, Geneva, provides proton-proton collisions at a c.m. energy of 13,000 GeV, which was designed based on the above physics argument. Consequently, we have witnessed the historical discovery of the Higgs boson!)

**Exercise 1.4:** A 125 GeV Higgs boson will have a production cross section of 20 pb at the LHC. How many events per year do you expect to produce for the Higgs boson with a designed LHC luminosity  $10^{33}/\text{cm}^2/\text{s}$ ? With the expected events, why is the Higgs boson so difficult to observe?

#### Lecture 2: Relativistic Kinematics and Phase Space, Collider Detectors

**Exercise 2.1:** Show that the phase space element  $d\vec{p}/2p^0$  is Lorentz invariant.

Exercise 2.2: (challenging problem) A particle of mass M decays to two particles isotropically in its rest frame. What does the momentum distribution look like in a frame in which the particle is moving with a speed  $\beta_z$ ? Compare the result with your expectation for the shape change for a basket ball.

**Exercise 2.3:** Consider a  $2 \to 2$  scattering process  $p_a + p_b \to p_1 + p_2$ . Assume that  $m_a = m_1$  and  $m_b = m_2$ . Show that

$$t = -2p_{cm}^{2}(1 - \cos\theta_{a1}^{*}),$$
  
$$u = -2p_{cm}^{2}(1 + \cos\theta_{a1}^{*}) + \frac{(m_{1}^{2} - m_{2}^{2})^{2}}{s},$$

 $p_{cm} = \lambda^{1/2}(s, m_1^2, m_2^2)/2\sqrt{s}$  is the momentum magnitude in the c.m. frame. Note: t is negative definite;  $t \to 0$  in the collinear limit, that could be singular for massless-exchange. Comment on the u-channel.

**Exercise 2.4:** (challenging problem) A particle of mass M decays to three particles  $M \rightarrow abc$ . Show that the phase space element can be expressed as

$$dPS_{3} = \frac{1}{2^{7}\pi^{3}} M^{2} dx_{a} dx_{b}.$$
$$x_{i} = \frac{2E_{i}}{M}, \ (i = a, b, c, \ \sum_{i} x_{i} = 2).$$

where the integration limits for  $m_a = m_b = m_c = 0$  are

$$0 \le x_a \le 1, \quad 1 - x_a \le x_b \le 1.$$

Note: For the decay in the M-rest frame, three of the four angular variables can be trivially integrated out (ignoring the spins of the particles).

**Exercise 2.5:** For a  $\pi^0$ ,  $\mu^-$ , or a  $\tau^-$  respectively, calculate its decay length if the particle has an energy E = 10 GeV.

#### Lecture 3: Lepton Colliders

Exercise 3.1: For a resonant production  $e^+e^- \rightarrow V^*$  with a mass  $M_V$  and total width  $\Gamma_V$ , derive the Breit-Wigner formula (If you find it too challenging for the calculation, you may skip this part and move on to the next line.)

$$\sigma(e^+e^- \to V^* \to X) = \frac{4\pi(2j+1)\Gamma(V \to e^+e^-)\Gamma(V \to X)}{(s - M_V^2)^2 + \Gamma_V^2 M_V^2} \frac{s}{M_V^2},$$

Consider a beam energy spread  $\Delta$  in Gaussian distribution

$$\frac{dL}{d\sqrt{\hat{s}}} = \frac{1}{\sqrt{2\pi} \ \Delta} \exp[\frac{-(\sqrt{\hat{s}} - \sqrt{s})^2}{2\Delta^2}],$$

obtain the appropriate cross section formulas for (a)  $\Delta \ll \Gamma_V$  (resonance line-shape) and (b)  $\Delta \gg \Gamma_V$  (narrow-width approximation).

**Exercise 3.2:** An event was identified to have a  $\mu^+$  and a  $\mu^-$  along with some missing energy. What can you say about the kinematics of the system of the missing particles? Consider for both an  $e^+e^-$  and a hadron collider.

Exercise 3.3 (challenging problem): Derive the Weizsäcker-Williams spectrum for a photon with an energy xE off an electron with an energy E

$$P_{\gamma/e}(x) \approx \frac{\alpha}{2\pi} \frac{1 + (1-x)^2}{x} \ln \frac{E^2}{m_e^2}.$$

Note that this procedure is the direct analog to deriving the DGLAP  $q \to q'g$  splitting in QCD.

## Lecture 4: Hadron Colliders

**Exercise 4.1:** For a four-momentum  $p \equiv p^{\mu} = (E, \vec{p})$ , define

$$E_T = \sqrt{p_T^2 + m^2}, \quad p_T^2 = p_x^2 + p_y^2, \quad y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$
  
then show  $p^{\mu} = (E_T \cosh y, \ p_T \cos \phi, \ p_T \sin \phi, \ E_T \sinh y),$   
and,  $\frac{d^3 \vec{p}}{E} = p_T dp_T d\phi \ dy = E_T dE_T d\phi \ dy.$ 

Due to the random boost between the Lab-frame (O) and the c.m. frame (O') for every event,

$$y' = \frac{1}{2} \ln \frac{E' + p'_z}{E' - p'_z} = \frac{1}{2} \ln \frac{(1 - \beta_{cm})(E + p_z)}{(1 + \beta_{cm})(E - p_z)} = y - y_{cm}$$

where  $\beta_{cm}$  and  $y_{cm}$  are the speed and rapidity of the c.m. frame w.r.t. the lab frame.

In the massless limit, the rapidity y defines the pseudo-rapidity:

$$y \to \eta = \frac{1}{2} \ln \frac{1 + \cos \theta}{1 - \cos \theta} = \ln \cot \frac{\theta}{2}.$$

**Exercise 4.2:** For a two-body massless final state with an invariant mass squared s, show that

$$\frac{d\hat{\sigma}}{dp_T} = \frac{4p_T}{s\sqrt{1-4p_T^2/s}} \; \frac{d\hat{\sigma}}{d\cos\theta^*}$$

where  $p_T = p \sin \theta^*$  is the transverse momentum and  $\theta^*$  is the polar angle in the c.m. frame. Comment on the apparent singularity at  $p_T^2 = s/4$ .

**Exercise 4.3**: What would be needed to construct a CP-odd observable in the hadron collider environment, such as the LHC? Try to construct one (anyone and for any process of your choice).

Reference: Genuine CP-odd Observables at the LHC, by Tao Han and Yingchuan Li: Phys.Lett. B683 (2010) 278-281, e-Print: arXiv:0911.2933.