

Non-global logarithms in jet and isolation cone cross sections

arXiv: 1803.07045 (MB, Thomas Becher, Ding Yu Shao)

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**UNIVERSITÄT
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AEC
ALBERT EINSTEIN CENTER
FOR FUNDAMENTAL PHYSICS

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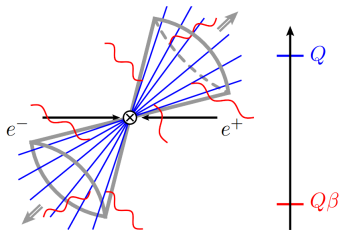
Overview

- 1 Non-Global Logarithms
- 2 Photon isolation cone cross section
- 3 LL resummation of isolated photons
- 4 Summary

Non-Global Logarithms

Non-Global Logarithms

Logarithms appearing in the calculation of non-global observables are called non-global logarithms.



$$\text{terms} \propto \alpha_s^n \ln^n \left(\frac{E_{out}}{E_{in}} \right) = \alpha_s^n \ln^n (\beta)$$

Classical perturbation theory **breaks down** for small β

"Order-By-Order" versus "Logarithm-by-Logarithm"

$$a_i \propto \mathcal{O}(1)$$

$$\text{LO: } \sigma = a_0 + \mathcal{O}(\alpha_s)$$

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$$\text{N}^5\text{LO: } \sigma = a_0 + a_1\alpha_s + a_2\alpha_s^2 + a_3\alpha_s^3 + a_4\alpha_s^4 + a_5\alpha_s^5 + \mathcal{O}(\alpha_s^6)$$

.....

"Order-By-Order" versus "Logarithm-by-Logarithm"

$$L = \ln(\beta)$$

$$\alpha_s L \propto \mathcal{O}(1)$$

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.....

"Order-By-Order" versus "Logarithm-by-Logarithm"

$$\text{LL: } \sigma = a + \sum b_m \alpha_s^m L^m + \mathcal{O}(\alpha_s^n L^{n-1})$$

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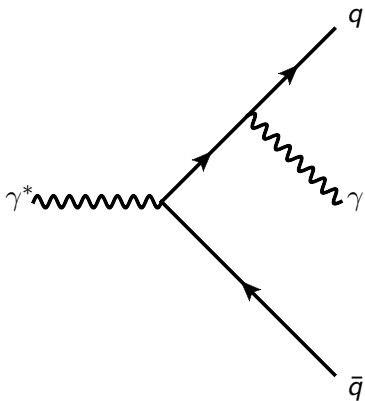
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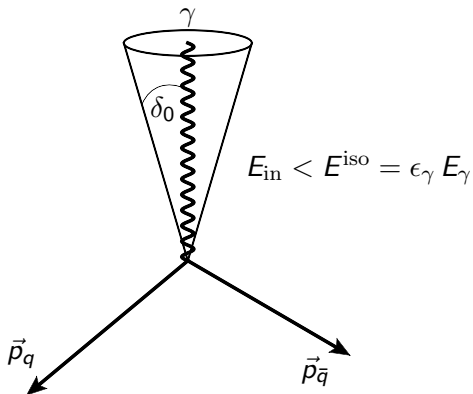
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Leading-Logarithm technique fixes perturbative approach

Photon Production



Isolated Photon

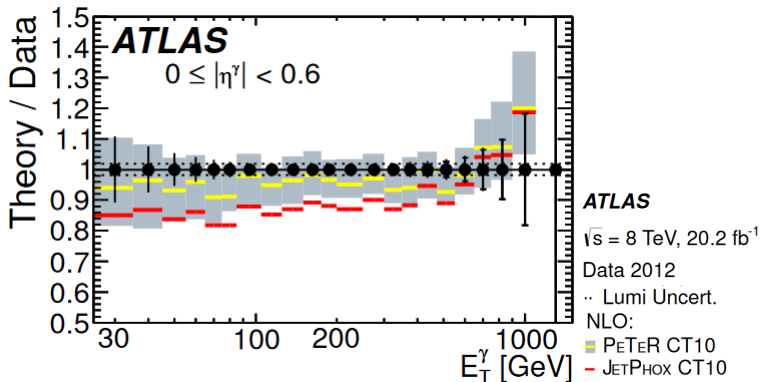


Why should we care?

No experimental measurements possible of direct photon production processes without such cuts!

- $E_{ATLAS}^{iso} = 4.8 \text{ GeV} + 4.2 \times 10^{-3} \times E_T^\gamma$
- $\delta_0^{ATLAS} \approx 0.42 \approx 0.13 \pi \approx 25^\circ$

ATLAS Results



arXiv: 1605.03495

Factorization formula

Factorization formula for e^+e^- -Annihilation

$$\frac{d\sigma(\epsilon_\gamma, \delta_0)}{dE_\gamma} = \sum_{m=2}^{\infty} \langle \mathcal{H}_{\gamma+m}(\{\underline{n}\}, E_\gamma, Q, \delta_0) \otimes \mathcal{S}_m(\{\underline{n}\}, \epsilon_\gamma, E_\gamma, \delta_0) \rangle$$

- $\mathcal{H}_{\gamma+m}$ "Hard function": Squared amplitude of hard process (consisting of m partons and the isolated photon)
- $\{\underline{n}\}$ Directions of the m hard partons ($\{\underline{n}\} = \{n_1, \dots, n_m\}$)
- \mathcal{S}_m "Soft function": soft radiation (inside cone)
- \otimes Integration over $\{\underline{n}\}$
- $\langle \dots \rangle$ Sum over Colors

Renormalization group equation

$\mathcal{H}_{\gamma+m}$ fulfill renormalization group equations

$$\frac{d}{d \ln \mu} \mathcal{H}_{\gamma+m}(\{\underline{n}\}, Q, \mu) = - \sum_{l=2}^m \mathcal{H}_l(\{\underline{n}\}, Q, \mu) \Gamma_{lm}^H(\{\underline{n}\}, Q, \mu)$$

- Compute $\mathcal{H}_{\gamma+m}$ at characteristic high scale (such as $\mu_h = E_\gamma$)
- Evolve $\mathcal{H}_{\gamma+m}$ to the scale of \mathcal{S}_m using RGE
- Compute \mathcal{S}_m at characteristic low scale (such as $\mu_s = \epsilon_\gamma E_\gamma$)

Ingredients for Leading Logarithmic accuracy (large N_c)

$\mathcal{H}_{\gamma+m}(\mu_h \approx E_\gamma)$ at LL

$$\begin{aligned}\mathcal{H}_{\gamma+2}(\mu_h \approx E_\gamma) &= \sigma_0 + \mathcal{O}(\alpha_s) \\ \mathcal{H}_{\gamma+m}(\mu_h \approx E_\gamma) &= \mathcal{O}(\alpha_s) \quad \forall m \neq 2\end{aligned}$$

Anomalous Dimension at LL

$$\Gamma^{(1)} = \begin{pmatrix} \mathbf{V}_2 & \mathbf{R}_2 & 0 & 0 & \dots \\ 0 & \mathbf{V}_3 & \mathbf{R}_3 & 0 & \dots \\ 0 & 0 & \mathbf{V}_4 & \mathbf{R}_4 & \dots \\ 0 & 0 & 0 & \mathbf{V}_5 & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$\mathcal{S}_m(\mu_s \approx \epsilon_\gamma E_\gamma)$ at LL

$$\mathcal{S}_m(\mu_s \approx \epsilon_\gamma E_\gamma) = 1 + \mathcal{O}(\alpha_s)$$

What does the anomalous Dimension represent?

Formal:

$$R_m = 4N_c \sum_i W_{i,i+1}^{m+1} \Theta_{OUT}(n_{m+1})$$

$$V_m = -4N_c \sum_i \int \frac{d\Omega(n_{m+1})}{4\pi} W_{i,i+1}^{m+1}$$

$$W_{i,j}^k = \frac{n_i \cdot n_j}{(n_i \cdot n_k)(n_k \cdot n_j)}$$

What does the anomalous Dimension represent?

Picture:

$$R_m \left[\begin{array}{c} 1 \\ \diagdown \quad \diagup \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \right] = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} + \dots + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array}$$

$$2V_m \left[\begin{array}{c} 1 \\ \diagdown \quad \diagup \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \right] = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} + \dots + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array} \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \\ \vdots \\ \text{---} \\ \text{---} \\ m \end{array}$$

Resummation by RGE at Leading Logarithmic accuracy

Renormalization group evolution

$$\frac{d}{d \ln \mu} \mathcal{H}_{\gamma+m}(\{\underline{n}\}, Q, \mu) = - \sum_{l=2}^m \mathcal{H}_l(\{\underline{n}\}, Q, \mu) \Gamma_{lm}^H(\{\underline{n}\}, Q, \mu)$$

Renormalization group evolution at leading log

$$\frac{d}{dt} \mathcal{H}_m(t) = \mathcal{H}_m(t) \mathbf{V}_m + \mathcal{H}_{m-1}(t) \mathbf{R}_{m-1}$$

$$t = \frac{1}{2\beta_0} \ln \frac{\alpha(\mu_s)}{\alpha(\mu_h)} = \frac{\alpha_s}{4\pi} \ln \frac{\mu_h}{\mu_s} + \mathcal{O}(\alpha_s^2)$$

Solution

$$\mathcal{H}_m(t) = \mathcal{H}_m(t_0) e^{(t-t_0)\mathbf{V}_m} + \int_{t_0}^t dt' \mathcal{H}_{m-1}(t') \mathbf{R}_{m-1} e^{-(t-t')\mathbf{V}_m}$$

Resummation by RGE at Leading Logarithmic accuracy

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$$\mathcal{H}_m(t) = \mathcal{H}_m(t_0) e^{(t-t_0)\mathbf{V}_m} + \int_{t_0}^t dt' \mathcal{H}_{m-1}(t') \mathbf{R}_{m-1} e^{-(t-t')\mathbf{V}_m}$$

Iterative procedure

$$\begin{aligned}\mathcal{H}_2(t) &= \mathcal{H}_2(0) e^{t\mathbf{V}_2} \\ \mathcal{H}_3(t) &= \int_0^t dt' \mathcal{H}_2(t') \mathbf{R}_2 e^{(t-t')\mathbf{V}_3} \\ \mathcal{H}_4(t) &= \int_0^t dt' \mathcal{H}_3(t') \mathbf{R}_3 e^{(t-t')\mathbf{V}_4} \\ \mathcal{H}_5(t) &= \dots\end{aligned}$$

Cross Section

$$\sigma(t) = \mathcal{H}_2(t) + \int \frac{d\Omega_1}{4\pi} \mathcal{H}_3(t) + \int \frac{d\Omega_1}{4\pi} \frac{d\Omega_2}{4\pi} \mathcal{H}_4(t) + \dots$$

Resummation by RGE at Leading Logarithmic accuracy

$$\sigma(t) = \mathcal{H}_2 + \mathcal{H}_3 + \mathcal{H}_4^{(1)} + \mathcal{H}_4^{(2)} + \dots$$

The diagram illustrates the resummation of the cross section $\sigma(t)$ at leading logarithmic accuracy. It shows a series of terms in the expansion:

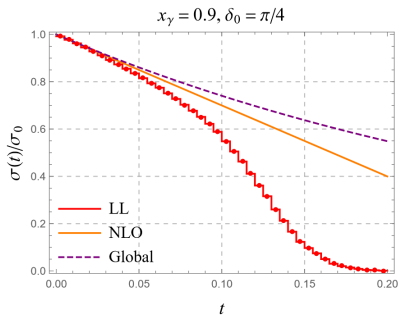
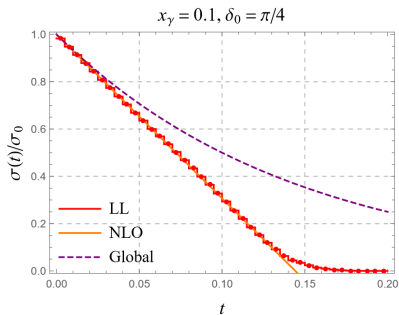
- \mathcal{H}_2 : A single vertex V_2 with external lines n_1 and n_2 .
- \mathcal{H}_3 : Two vertices V_2 and V_3 connected by a radiator R_2 , with external lines n_1, n_2, n_3 .
- $\mathcal{H}_4^{(1)}$: Four vertices $V_2, R_2, V_3, R_3^{(1)}, V_4^{(1)}$ connected in a chain, with external lines n_1, n_2, n_3, n_4 .
- $\mathcal{H}_4^{(2)}$: Four vertices $V_2, R_2, V_3, R_3^{(2)}, V_4^{(2)}$ connected in a chain, with external lines n_1, n_2, n_3, n_4 .

The vertices V_i are represented by yellow ovals, and the radiators R_i are represented by green ovals. The expansion continues with higher-order terms indicated by the ellipsis.

Monte Carlo!

- generate a lot of $\mathcal{H}_{[\gamma+]_2}(t = 0)$ (using event generator [Madgraph])
- for every event
 - calculate \mathbf{V}_2
 - generate $\Delta t'$
 - generate \mathbf{R}_2 and calculate \mathbf{V}_3
 - calculate $\mathcal{H}_3(\Delta t')$
 - repeat for $\mathcal{H}_4(\Delta t' + \Delta t'') \dots$
 - restart the "showering", as soon as we emit into the cone

$$e^+ e^- \rightarrow \gamma + X$$



Summary

- we resum non-global logarithms at leading logarithmic accuracy using a parton shower;
- flexible implementation by using Madgraph, Code can be easily adapted to different non-global observables.

- small-angle observables not (yet) accounted for;
- going to NLL: higher-order corrections to matching coefficients and anomalous dimension;
- $N_c = 3$ difficult to achieve.



Thank you!