Dynamic AdS/QCD and the Nc,Nf Dependence of Gauge Theory

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New understanding of holographic descriptions of mesons

A stab at predicting the spectra of arbitrary AF gauge theories

Mainz September 2014

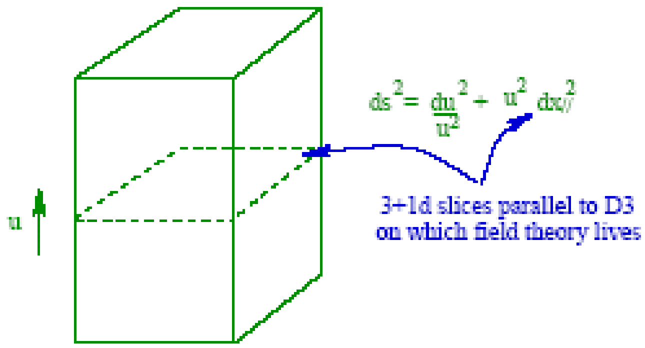
Building Holographic QCD

AdS/CFT Correspondence

Maldacena, Witten...

4d strongly coupled $\mathcal{N}$=4 SYM = IIB strings on $AdS_5 \times S^5$

Pretty well established by this point!



u corresponds to energy (RG) scale in field theory

The SUGRA fields act as sources

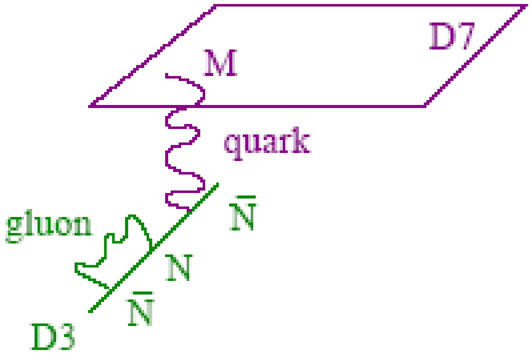
$$\int d^4x \; \Phi_{SUGRA}(u_0)\lambda\lambda$$

eg asymptotic solution ($u \to \infty$) of scalar

$$\varphi \simeq \frac{m}{u} + \frac{\langle\lambda\lambda\rangle}{u^3}$$

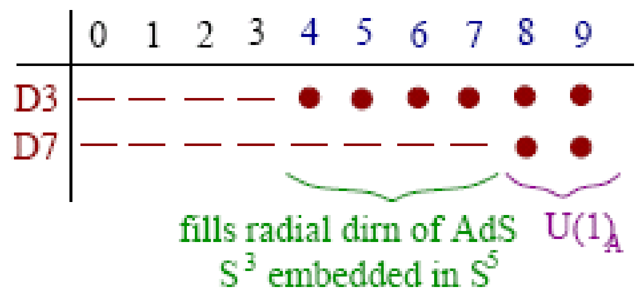
Adding Quarks

Bertolini, DiVecchia...; Polchinski, Grana; Karch, Katz...



Quarks can be introduced via D7 branes in AdS

The brane set up is



We will treat D7 as a probe - quenching in the gauge theory.

Minimize D7 world volume with DBI action

$$S_{D7} = -T_7 \int d\xi^8 \sqrt{P[G_{ab}]}, \qquad P[G_{ab}] = G_{MN} \frac{dx^M}{d\xi^a} \frac{dx^N}{d\xi^b}$$

Mesons in AdS_5

Kruczenski, Myers, Mateos, Winters

The D7 lie flat in AdS. We can consider fluctuations that describe R-chargeless mesons

$$w_6 + i w_5 = d + \delta(\rho) e^{ik.x}$$

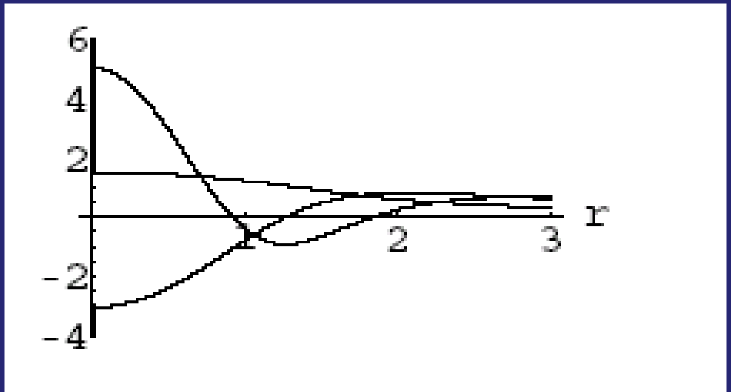
δ satisfies a linearized EoM

$$\partial_\rho^2 \delta + \frac{3}{\rho}\partial_\rho \delta + \frac{M^2}{(\rho^2+1)^2}\delta = 0$$

and the mass spectrum is

$$M = \frac{2d}{R^2}\sqrt{(n+1)(n+2)} \sim \frac{2m}{\sqrt{\lambda_{YM}}}$$

Tightly bound - meson masses suppressed relative to quark mass



Orthonormal wave functions

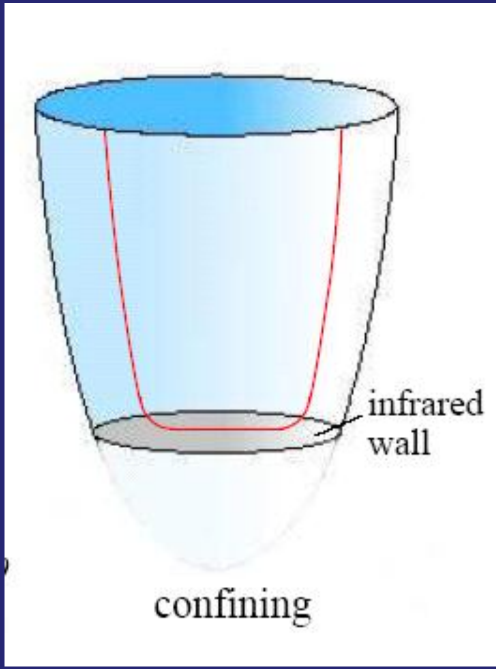
Add Confinement and Chiral Symmetry Breaking

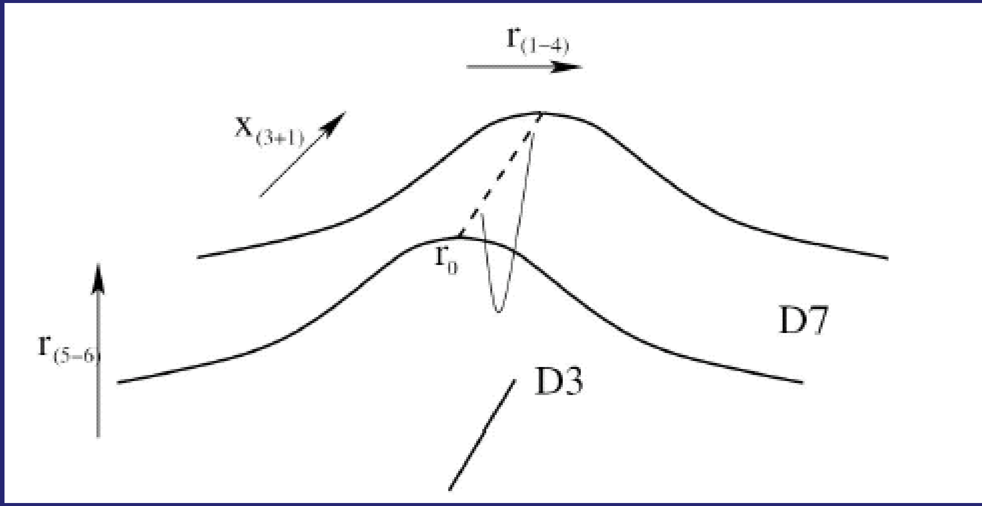
$$ds^2 = \frac{r^2}{R^2} A^2(r) dx_{3+1}^2 + \frac{R^2}{r^2} dr_6^2,$$

$$A(r) = \left(1 - \left(\frac{r_w}{r}\right)^8\right)^{1/4}, \qquad e^{\phi} = \left(\frac{1 + (r_w/r)^4}{1 - (r_w/r)^4}\right)^{\sqrt{3/2}}$$

Dilaton Flow Geometry: Gubser, Sfetsos

Here, this is just a simple, back reacted, repulsive, hard wall….





Babington et al, Ghoroku..

Erdmenger,
NE,Kirsch,
Threlfall
0711.4467

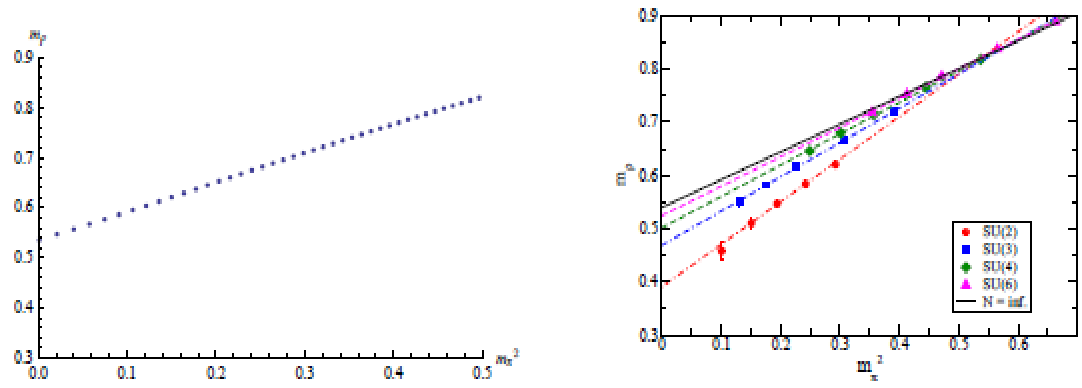


Figure 6.5: A plot of m_ρ vs m_π^2 in the Constable-Myers background on the left (we thank Andrew Tedder for generating this plot). Lattice data [20] (preliminary, quenched and at finite spacing) for the same quantity is also shown on the right.

These models live at very large coupling – QCD is at intermediate coupling strength…

The running coupling is rather different (UV strong fixed point) and there are lots of extra super-partners…

Success beyond caricature seems surprising… how does one improve towards QCD?

Holographic Models for QCD in the Veneziano Limit
Matti Jarvinen, Elias Kiritsis
1112.1261

We model the qq condensate by a scalar in "AdS"…

Breitenlohner-Freedman Bound

A scalar in AdS is stable until

msq < - 4

ie Δ < 2

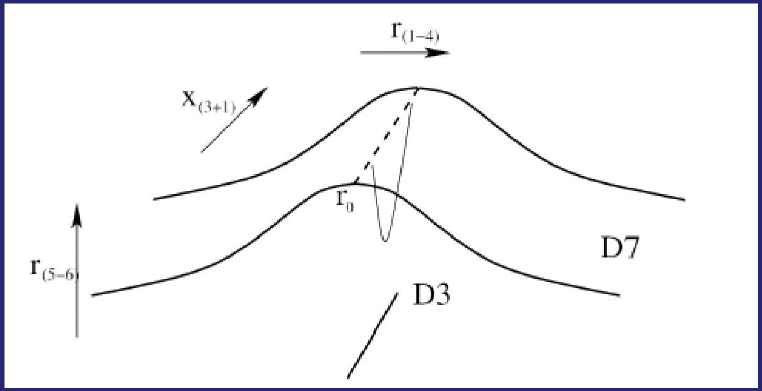
$$m^2 = \Delta(\Delta - 4)$$

A hard prediction that matches gap equation-ology for the on-set of chiral symmetry breaking…

Holographic models work for QCD mesons because they are describing the running of the anomalous dimension of qq rather than the running of the coupling…

D3/ Probe D7 Model

Alvares, NE, Kim, 1204.2474



λ

$$S_{D7} = -T \int d^4x d\rho \, \rho^3 e^{\phi} \sqrt{1 + (\partial_\rho L)^2}$$

$$S = \int \mathrm{d}\rho \lambda(r) \rho^3 \sqrt{1 + L'^2}$$

We expand for small L

$$S = \int \mathrm{d}\rho \left(\frac{1}{2} \lambda(r) \Big|_{L=0} \rho^3 L'^2 + \rho^3 \frac{d\lambda}{dL^2}\Big|_{L=0} L^2 \right)$$

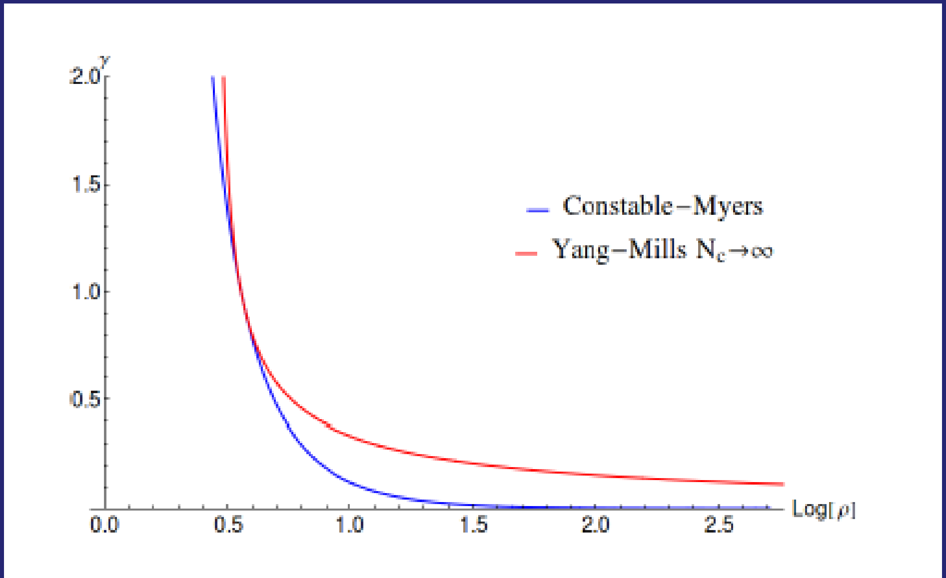
we can now make a coordinate transformation

$$\lambda(\rho)\rho^3 \frac{d}{d\rho} = \tilde{\rho}^3 \frac{d}{d\tilde{\rho}}, \qquad \tilde{\rho} = \sqrt{\frac{1}{2} \frac{1}{\int_\rho^\infty \frac{d\rho}{\lambda \rho^3}}}$$

$$L = \tilde{\rho}\phi$$

$$S = \int \mathrm{d}\tilde{\rho} \frac{1}{2} \left(\tilde{\rho}^5 \phi'^2 - 3\tilde{\rho}^3 \phi^2 \right) + \int \mathrm{d}\tilde{\rho} \frac{1}{2} \lambda \frac{\rho^5}{\tilde{\rho}} \frac{d\lambda}{d\rho} \phi^2$$

This is the action of a scalar in AdS with a mass squared of -3 + ρ dependent correction from the gradient of λ



Top-down probe-brane models of QCD are just AdS/QCD with the background providing a running γ…

Dynamic AdS/QCD

Timo Alho, NE, KimmoTuominen
1307.4896

$$S = \int d^4x\, d\rho\, \mathrm{Tr}\, \rho^3 \left[\frac{1}{\rho^2 + |X|^2} |DX|^2 + \frac{\Delta m^2}{\rho^2} |X|^2 + \frac{1}{2\kappa^2} (F_V^2 + F_A^2) \right]$$

$$X = L(\rho)\, e^{2i\pi^a T^a}.$$

$$ds^2 = \frac{d\rho^2}{(\rho^2 + |X|^2)} + (\rho^2 + |X|^2) dx^2,$$

D7 probe action in AdS expanded to quadratic order

X is now a dynamical field whose solution will determine the condensate as a function of m

We use the top-down IR boundary condition on mass-shell: $X'(\rho=X) = 0$

X enters into the AdS metric to cut off the radial scale at the value of m or the condensate – no hard wall

Dynamic AdS/QCD

The gauge DYNAMICS is input through Δm

$$\Delta m^2 = -2\gamma = -\frac{3(N_c^2-1)}{2N_c\pi}\alpha$$

$$m^2 = \Delta(\Delta - 4)$$

For example try to (very naively) input the two loop QCD running of the coupling…

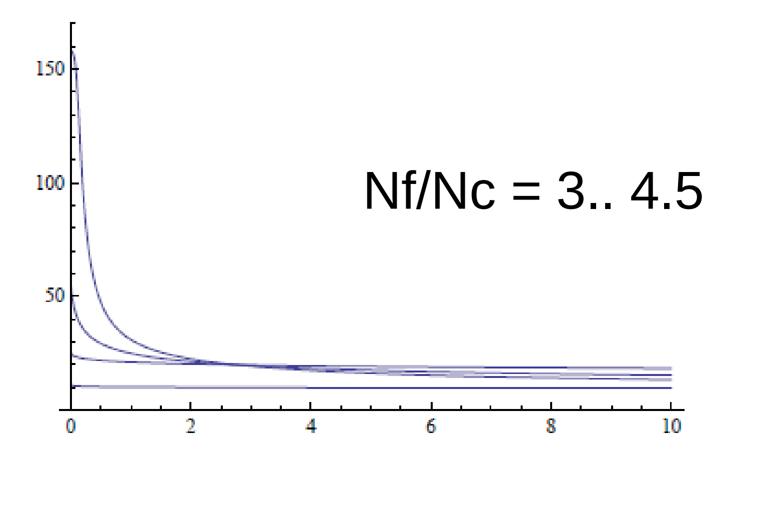
$$\beta(g) = -\frac{g^3}{(4\pi)^2}\left\{\frac{11}{3}N_c - \frac{2}{3}N_f\right\} - \frac{g^5}{(4\pi)^4}\left\{\frac{34}{3}N_c^2 - \frac{N_f}{N_c}\left[\frac{13}{3}N_c^2 - 1\right]\right\} + \cdots$$

Using the 't Hooft coupling, and setting $\frac{N_f}{N_c} \to x$ we obtain

$$\lambda \equiv g^2 N_c \quad , \quad \dot{\lambda} = -b_0\lambda^2 + b_1\lambda^3 + \mathcal{O}(\lambda^4)$$

with

$$b_0 = \frac{2}{3}\frac{(11-2x)}{(4\pi)^2} \quad , \quad \frac{b_1}{b_0^2} = -\frac{3}{2}\frac{(34-13x)}{(11-2x)^2}$$



The only free parameters are Nc, Nf, m, Λ

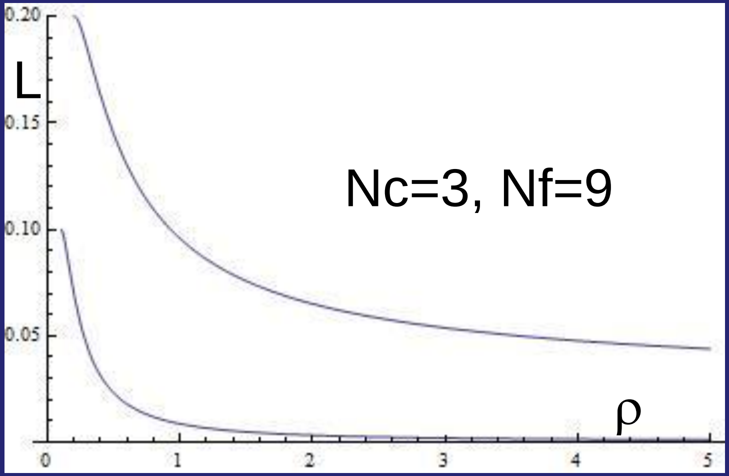
Formation of the Chiral Condensate

We solve for the vacuum configuration of L

$$\partial_\rho[\rho^3 \partial_\rho L] - \rho \Delta m^2 L = 0\,.$$

Shoot out with

$L'(\rho = L) = 0$

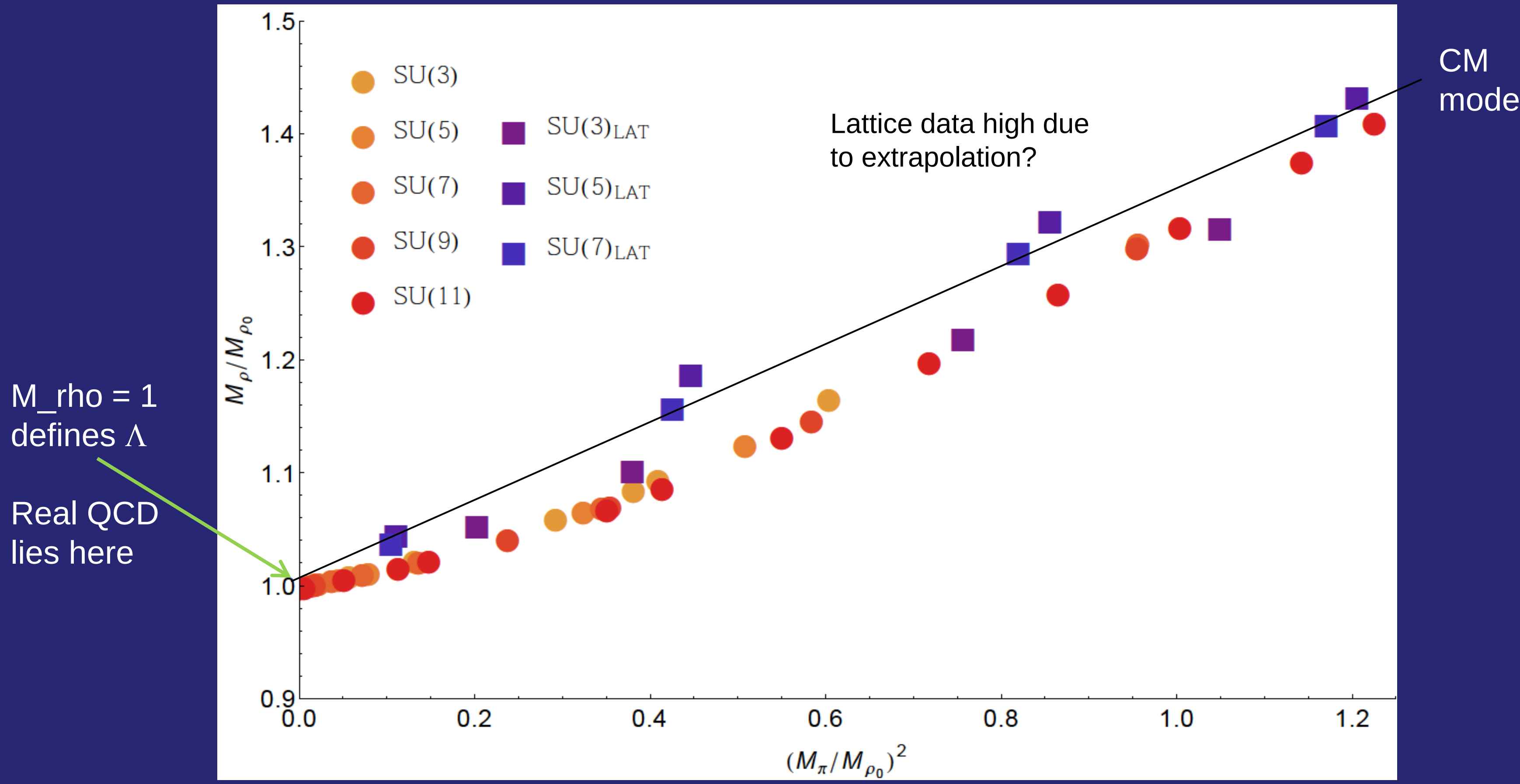


Read off m and qq in the UV…

Now solve for meson masses by looking at linearized fluctuations about this vacuum…

SU(Nc) gauge + 3 quarks

NE, Erdmenger & Mark Scott

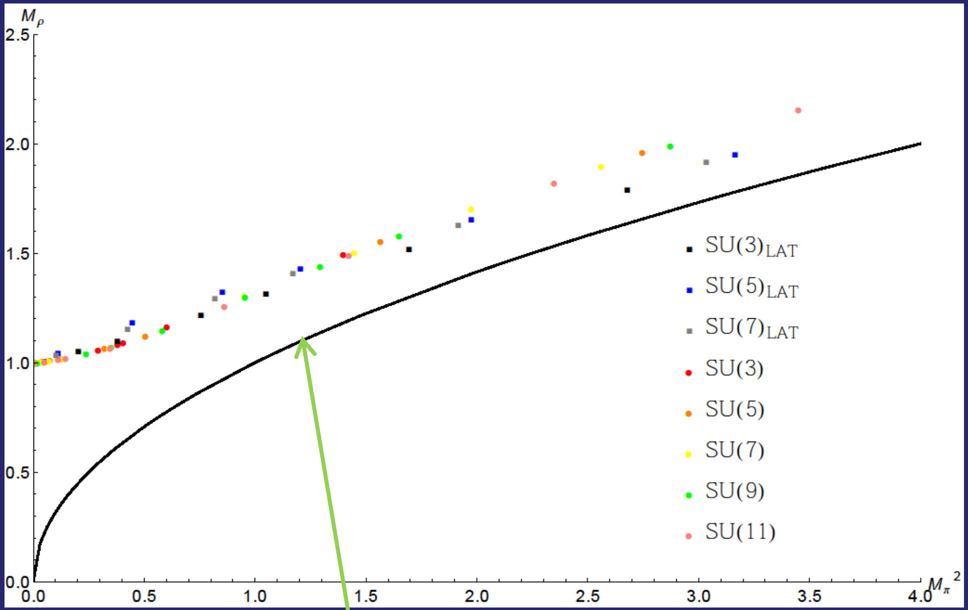


There is very little Nc dependence – basically quenched…

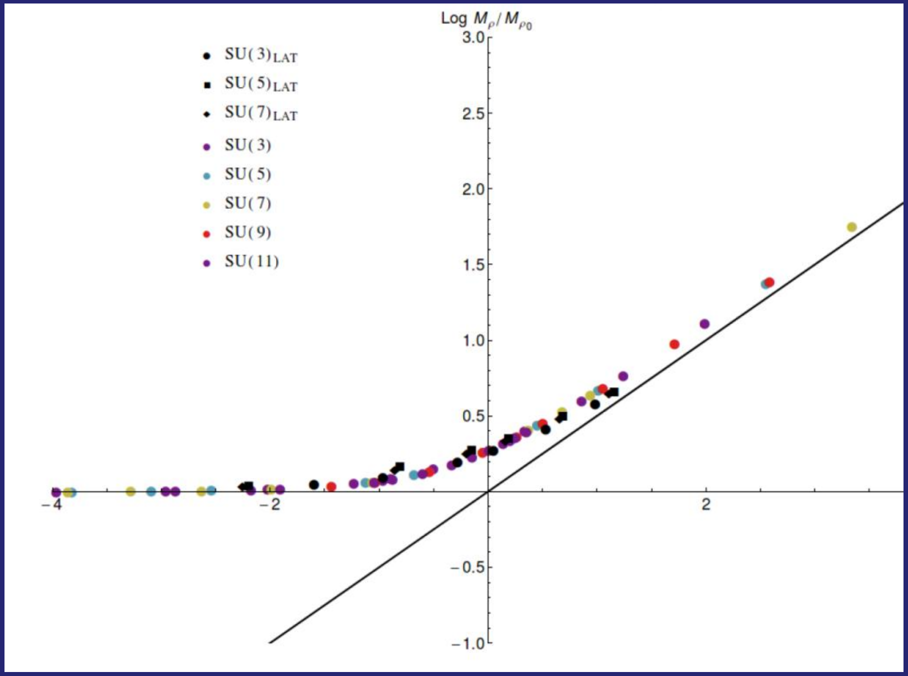
Hence comparison to quenched lattice data (Bali et al… arXiv1304.4437)

All of these models lie within 10% on any point....

SU(Nc) gauge + 3 quarks



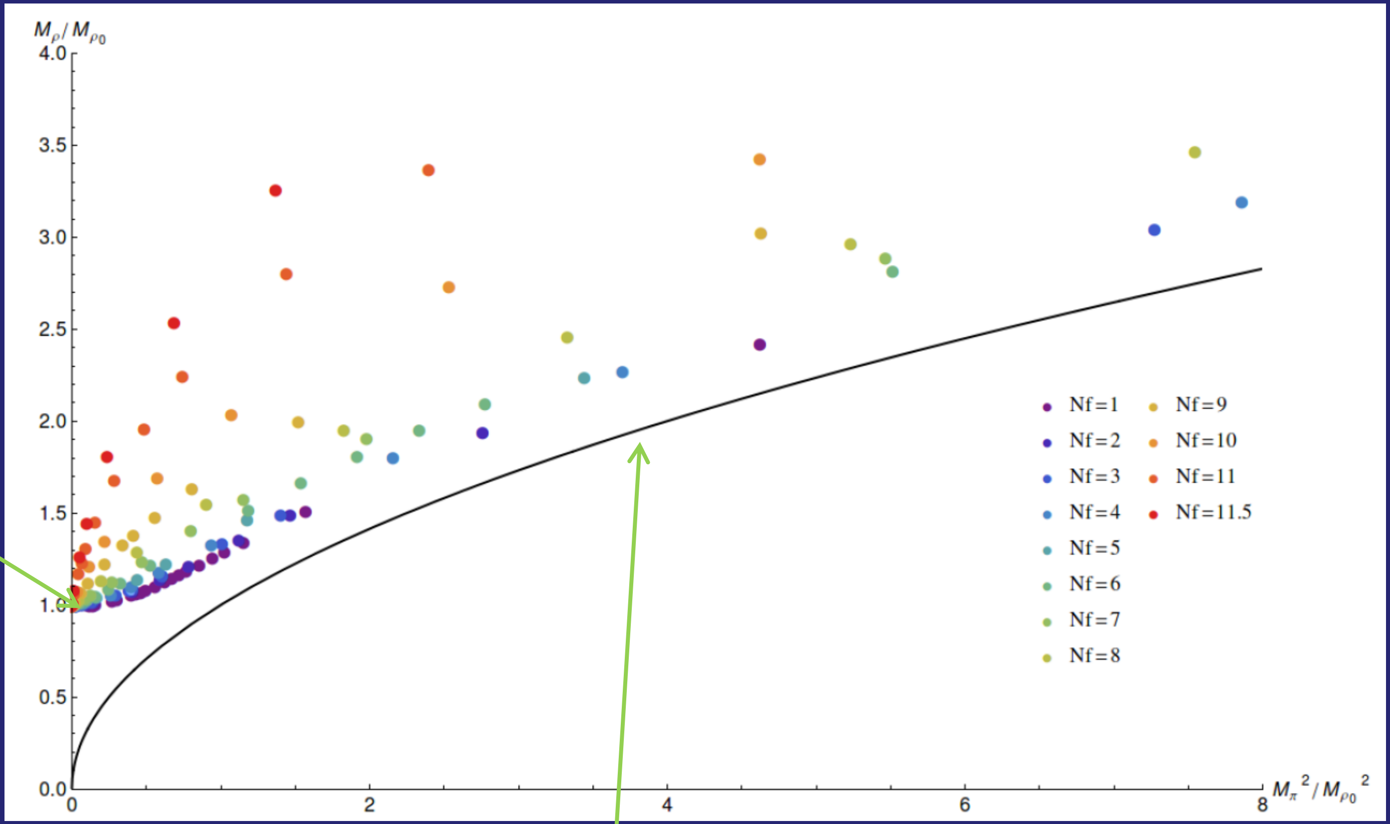
M_rho = M_pi



SU(3) gauge theory + Nf quarks

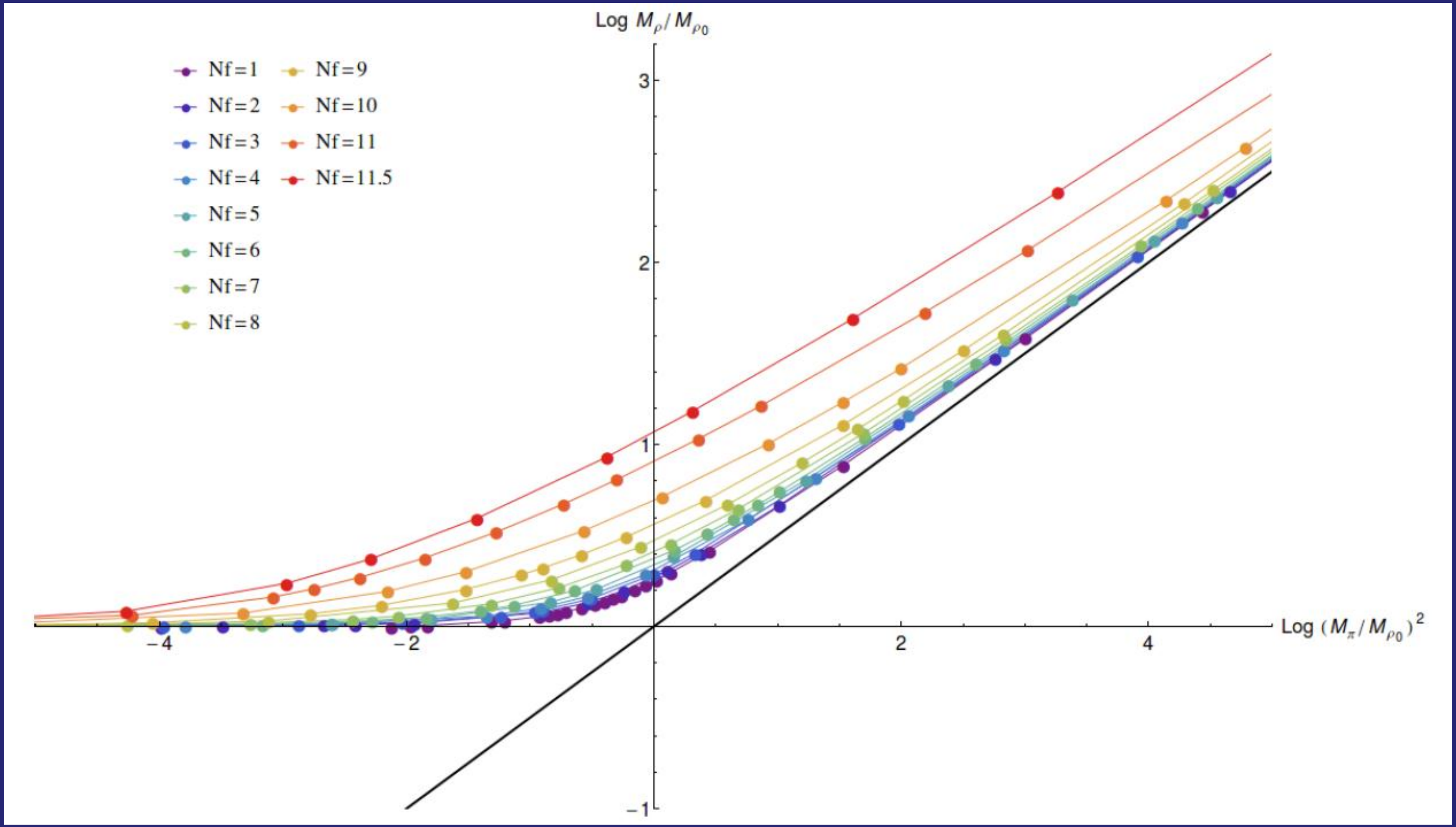
M_rho = 1 defines Λ

Real QCD lies here



M_rho = M_pi

SU(3) gauge theory + Nf quarks



We do see new behaviour as Nf heads towards 12….

The Conformal Window

SU(Nc) gauge theory with Nf fundamental quarks

	No AF
Nf=11/2 Nc	________________
	CFT
Nf = 4 Nc	________________
	χ SB

$$\beta(g) = -\frac{g^3}{(4\pi)^2}\left\{\frac{11}{3}N_c - \frac{2}{3}N_f\right\} - \frac{g^5}{(4\pi)^4}\left\{\frac{34}{3}N_c^2 - \frac{N_f}{N_c}\left[\frac{13}{3}N_c^2 - 1\right]\right\} + \cdots$$

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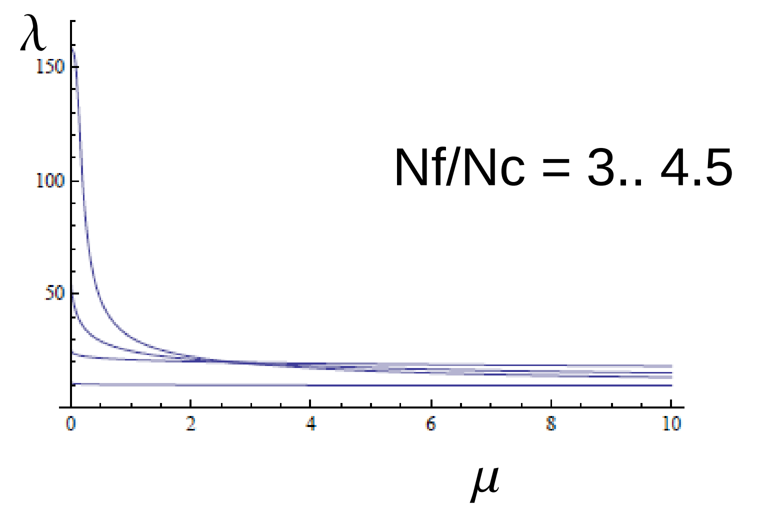
with

$$b_0 = \frac{2}{3}\frac{(11-2x)}{(4\pi)^2} \quad , \quad \frac{b_1}{b_0^2} = -\frac{3}{2}\frac{(34-13x)}{(11-2x)^2}$$

$m\,\bar{q}q$

$$\gamma_m^{(1)} = \mu\frac{d\ln m_q}{d\mu} = \frac{3(N_c^2-1)}{4N_c\pi}\alpha$$



If critical γ = 1…. Nf/Nc ~ 4

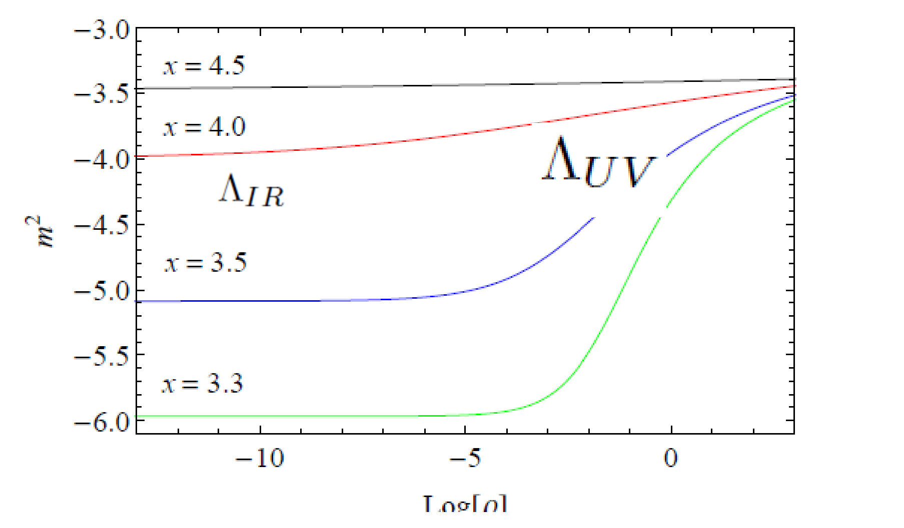
Yamawaki,Appelquist, Terning, Sannino,…

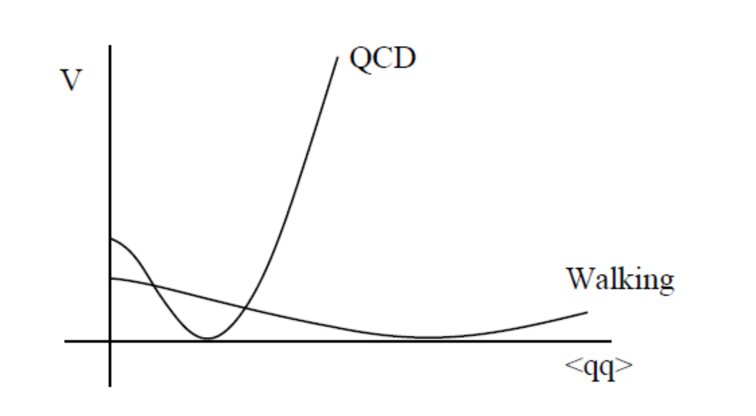
Walking Dynamics Holdom

Just above the CW regime theories have an enhanced UV quark condensate

$$\langle \bar{q}q \rangle_{UV} \sim \Lambda_{UV} \langle \bar{q}q \rangle_{IR} \sim \Lambda_{UV} \Lambda_{IR}^2$$

$$f_\pi \sim \Lambda_{IR}$$

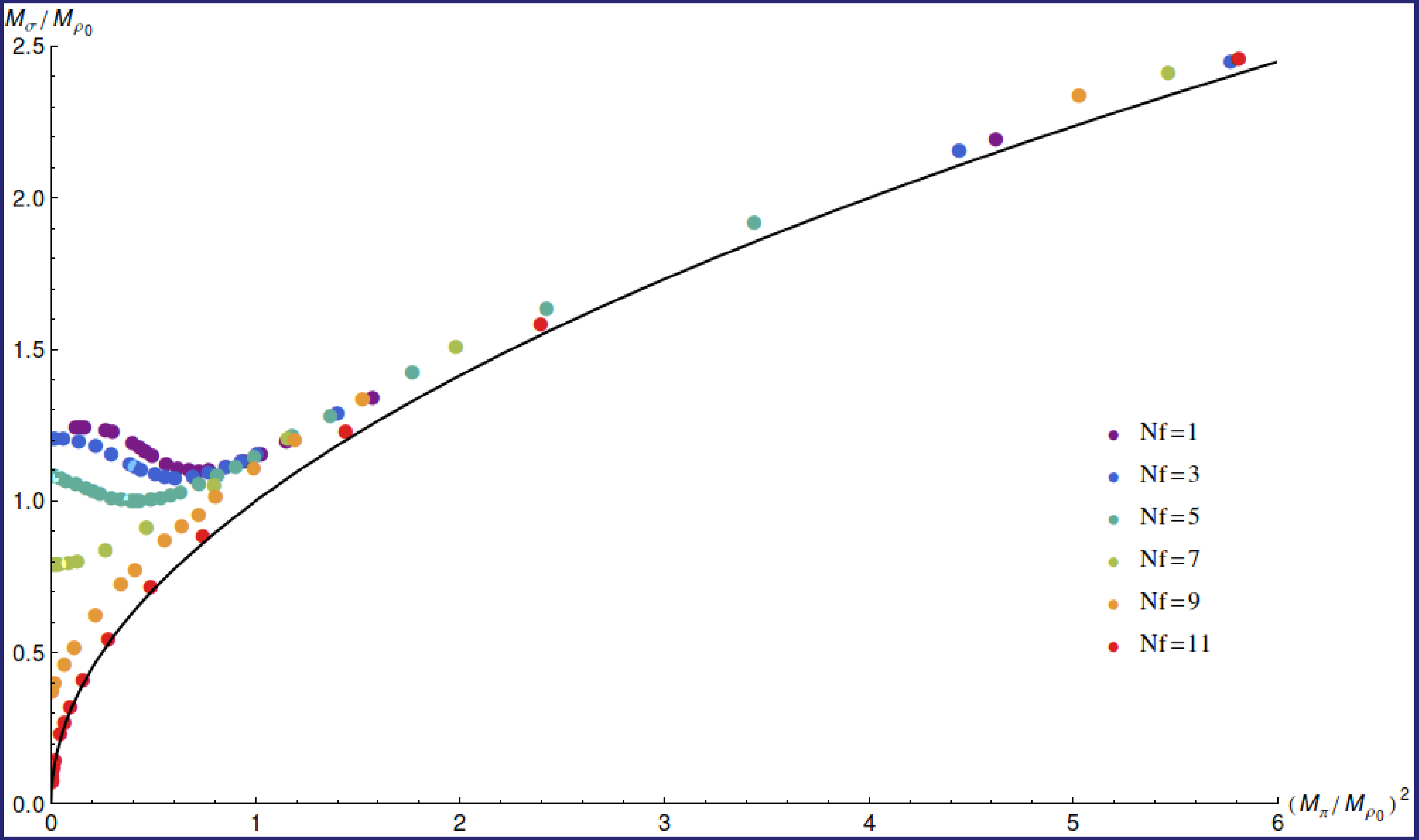




- Is the sigma particle light – a techni-dilaton?
- Is the higgs such a technicolor state?

SU(3) gauge theory + Nf quarks

The QCD point is not right for the f0(500) but about right for the f0(980) – is the f0(500) odd eg a molecule ???



We indeed see a light sigma relative to the rho…

Alho, NE, Tuominen 1307.4896

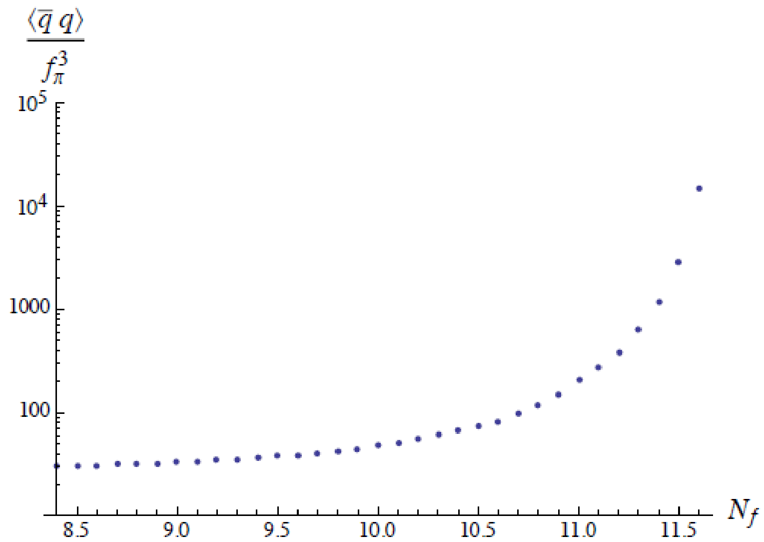


FIG. 5: The quark condensate normalized by f_π^3 vs N_f.

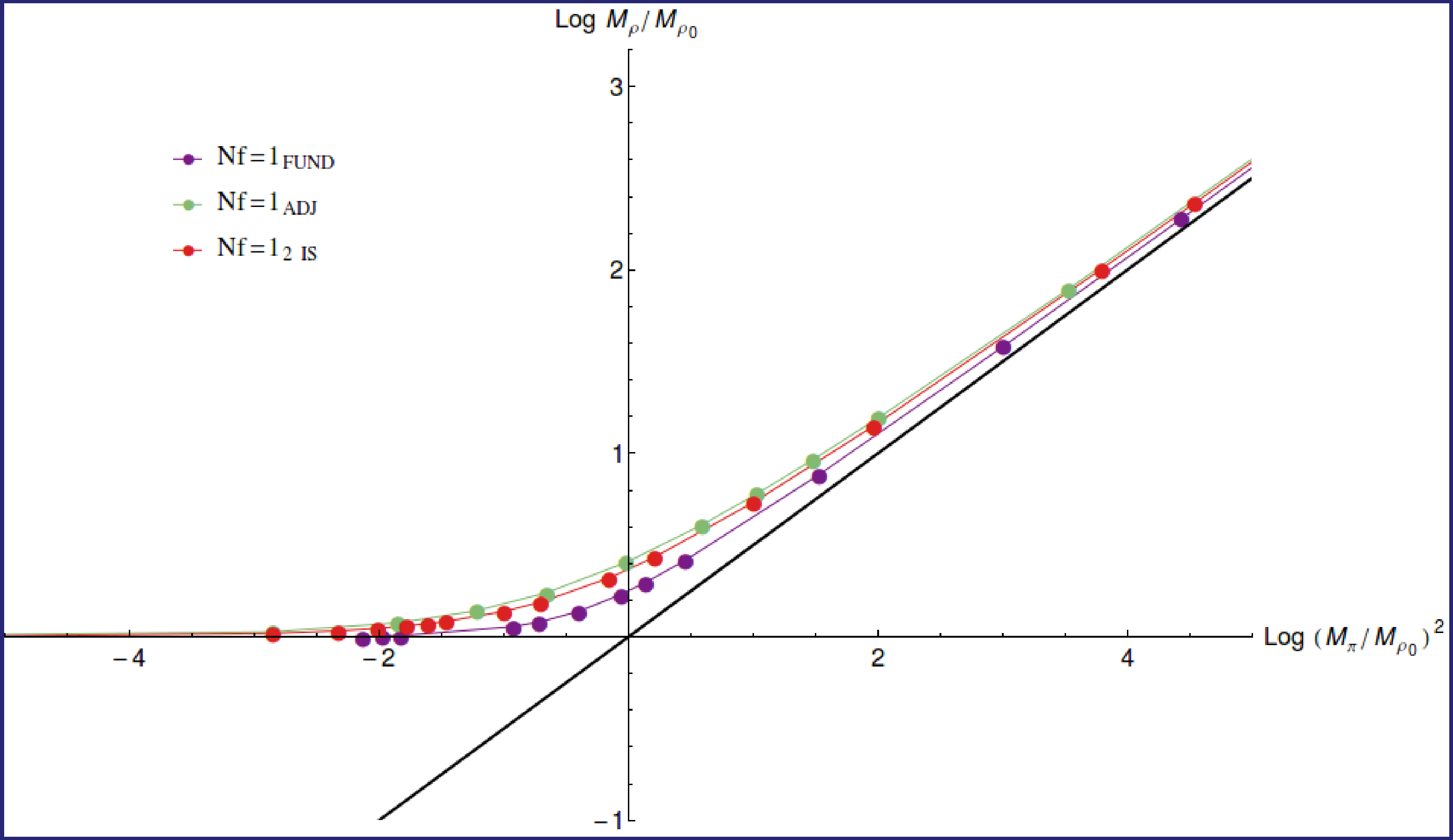
qq enhancement reproduced….

Cf FCNC problem in TC

Pick your theory…

Input (assumed) running of γ for any theory you fancy…

eg SU(3) gauge +

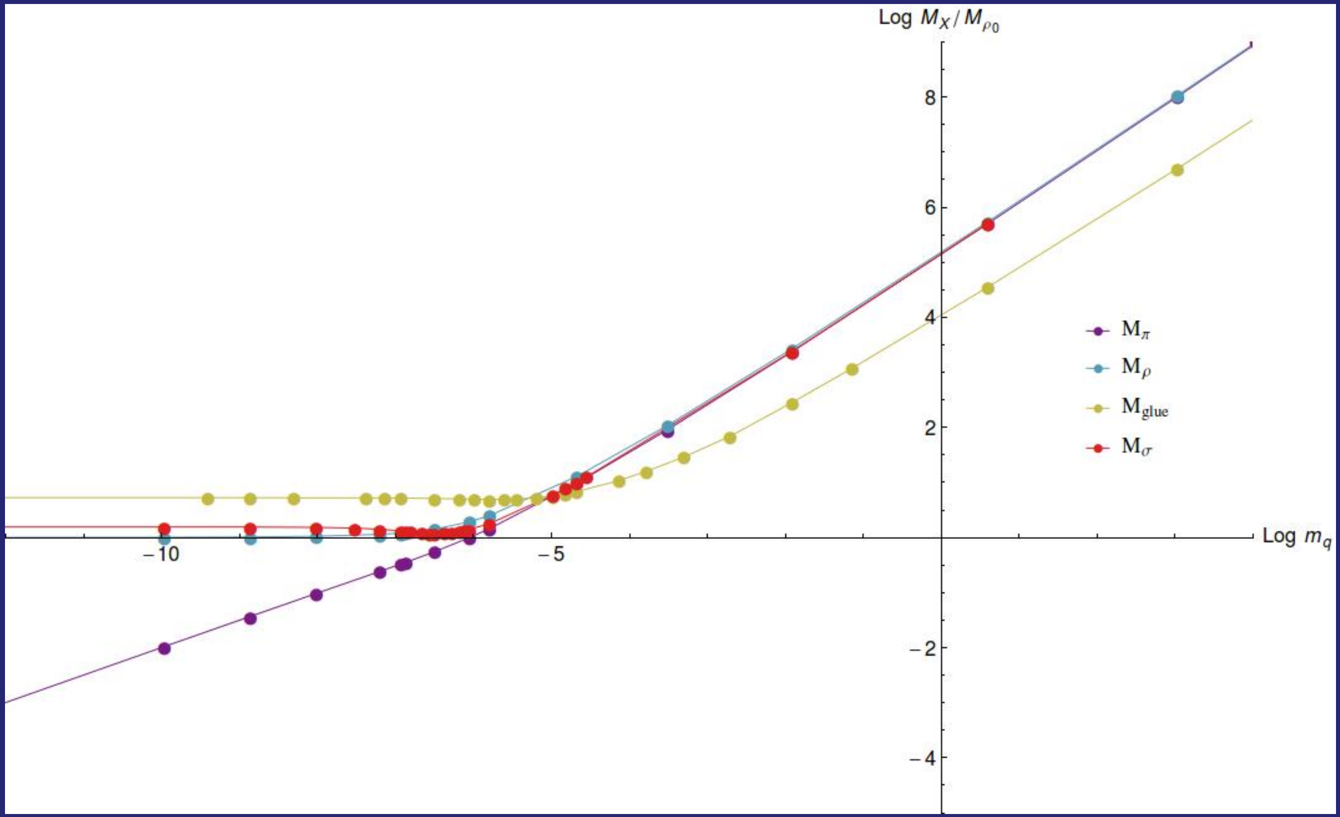


Glueballs

The model concentrates on the quark states… we're not trying to describe the running of the coupling or Tr F^2… talks later today will try to include that extra scalar correctly… however for us roughly…

Find the dynamical IR quark mass… below that scale run the coupling as pure YM.. Find the IR pole… multiply by 8 and that's the glueball mass!

SU(3)
Nf=3

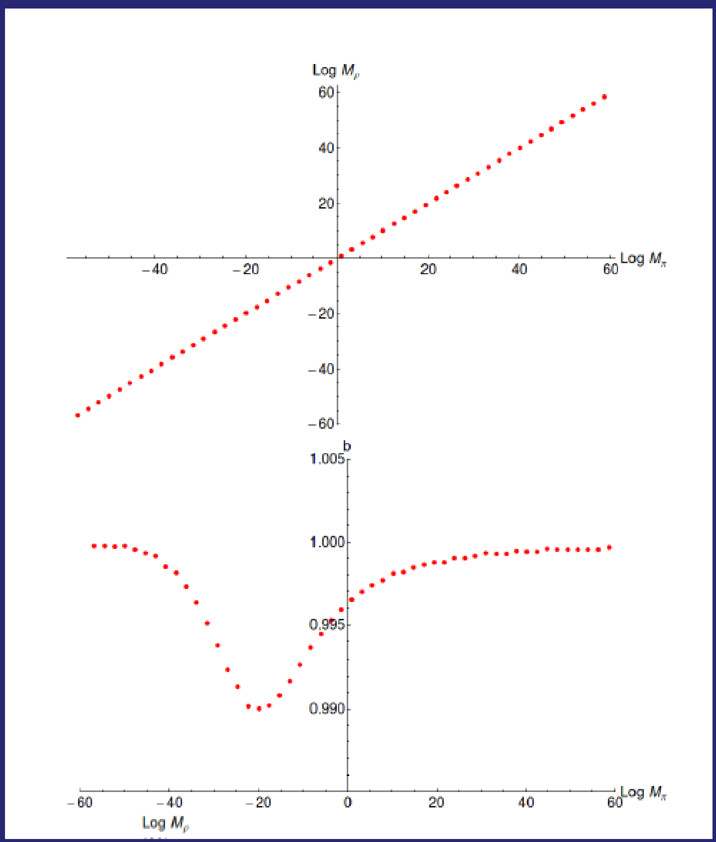


Study the Conformal Window

NE, Scott 1405.5373

For Nf > 4 Nc there is no dynamical scale so study the theory at non-zero quark mass…

eg SU(3) gauge + 13 flavours



$$M_\rho \propto M_\pi^b$$

b=1 for the conformal IR and UV regimes but deviates at the intermediate running scale…

Conclusions

Holographic models of QCD continue to improve… running of γ is crucial…one gets a good description of the lowest lying spectra at better than 10% and you can see generic behaviours with Nc Nf easily...

Holography is a remarkably simple method to get a ball park answer for behaviour… but it still can't be systematically improved…

Backreacted flavours remain a very interesting project (Barcelona ESF workshop in November)

On going work: T μ phase structure… pomeron physics… enlarging to the full QCD spectra…