

# Generalized Entropy and higher derivative Gravity

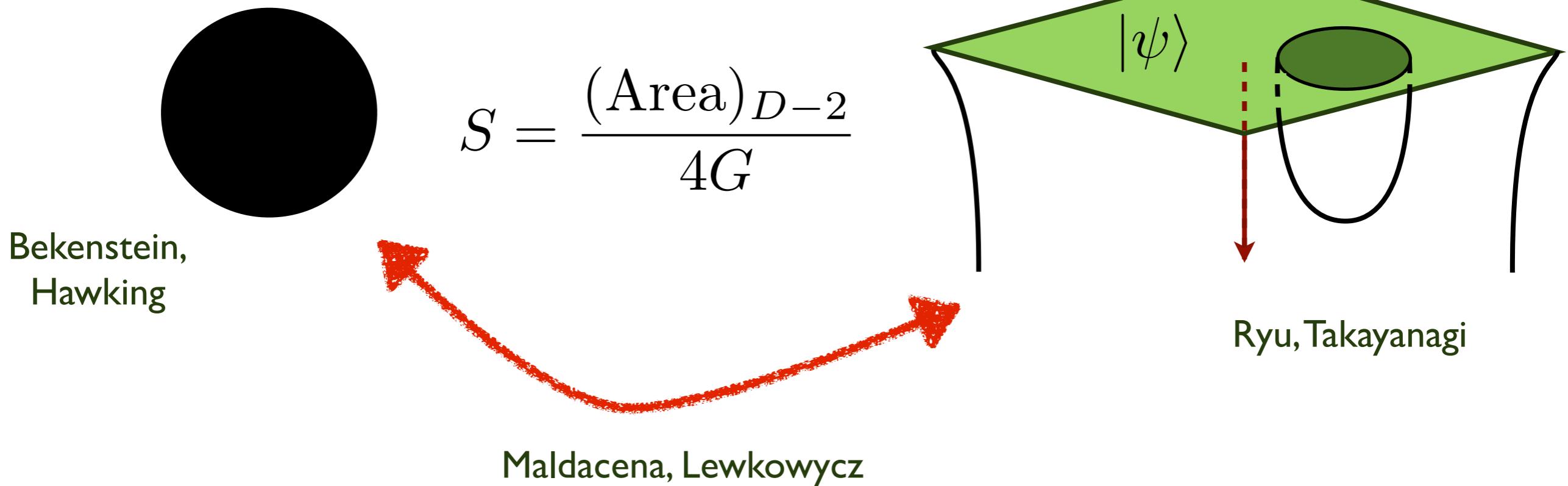
Joan Camps  
DAMTP, Cambridge

based on 1310.6659, see also 1310.5713 by Xi Dong

The String Theory Universe, Mainz, 23 September 2014

# Entropy in General Relativity

$$I = \frac{1}{16\pi G} \int \sqrt{-g} d^D x R$$



# Why higher derivatives

- Such corrections are predicted in String Theory (essential to the great success in microscopic accounts of entropy)

$$I = \frac{1}{16\pi G} \int \sqrt{-g} d^D x \left( R + \alpha' \text{Riem}^2 + \dots \right)$$

- Deformations of AdS/CFT: playground, robustness, universality       $\eta/s$        $c \neq a$
- It is an interesting problem *per se*

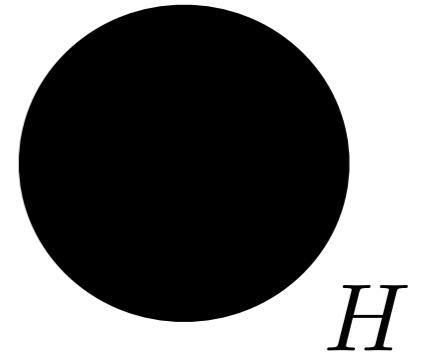
# BH Entropy for higher derivative theories

Wald

$$I = \int \sqrt{-g} d^D x \mathcal{L}(R_{\mu\nu\rho\sigma}, \nabla_\mu, g_{\mu\nu})$$

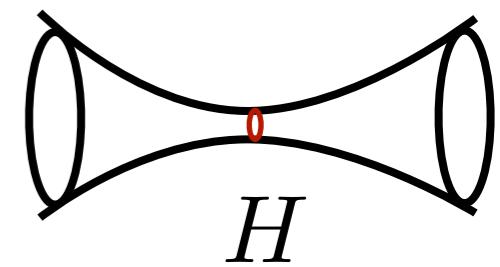
- Black holes enjoy a first law, with entropy

$$S \sim \int_H \sqrt{\gamma} d^{D-2} \sigma \frac{\partial \mathcal{L}}{\partial \text{Riem}}$$



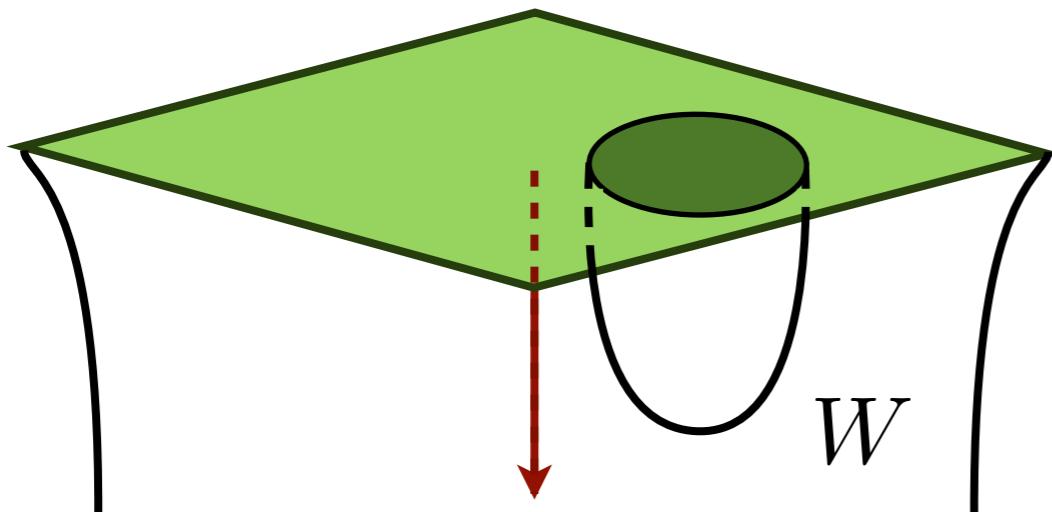
- Not known if they satisfy a second law
- Some proposals for non-stationary

Wald-lyer



# HEE & H Derivatives

- Wald entropy is **not** a good Holographic entanglement entropy (wrong divergences)



$$S \neq \int_W \sqrt{\gamma} d^{D-2} \sigma \frac{\partial \mathcal{L}}{\partial \text{Riem}}$$

- For Lovelock theories, there is a proposal with the correct structure
- That proposal coincides with Wald-Iyer's

de Boer et al  
Myers et al  
Fursaev et al

# Result

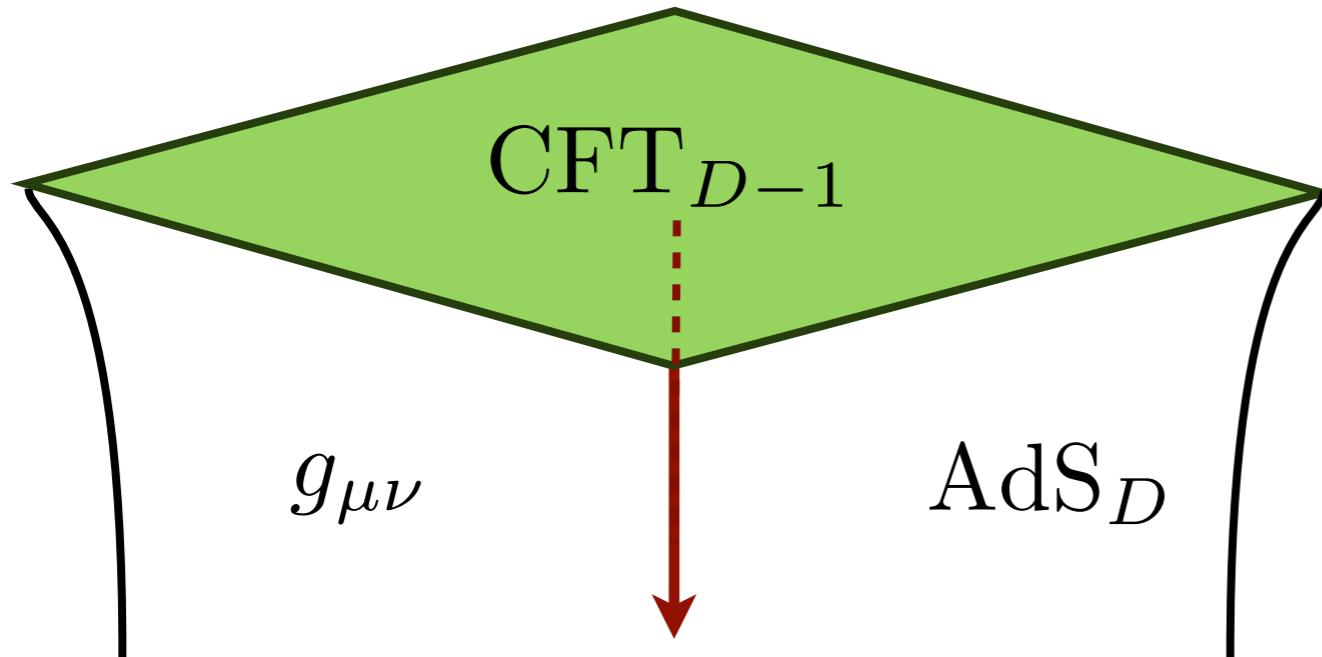
$$S \sim \int_W \sqrt{\gamma} d^{D-2} \sigma \left( \frac{\partial \mathcal{L}}{\partial \text{Riem}} + \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2 \right)$$

- It reduces to Wald's for stationary cases (trivially)
- For Lovelock Gravity, it gives the Jacobson-  
Myers entropy functional
  - de Boer et al
  - Myers et al
  - Fursaev et al
- More generally, disagrees with Wald-Iyer's

# Ingredients

Maldacena, Lewkowycz

- Saddle point approximation to holography



$$Z_{\text{CFT}_{D-1}} \approx e^{-I_E[g_{\mu\nu}]}$$

- Replica trick

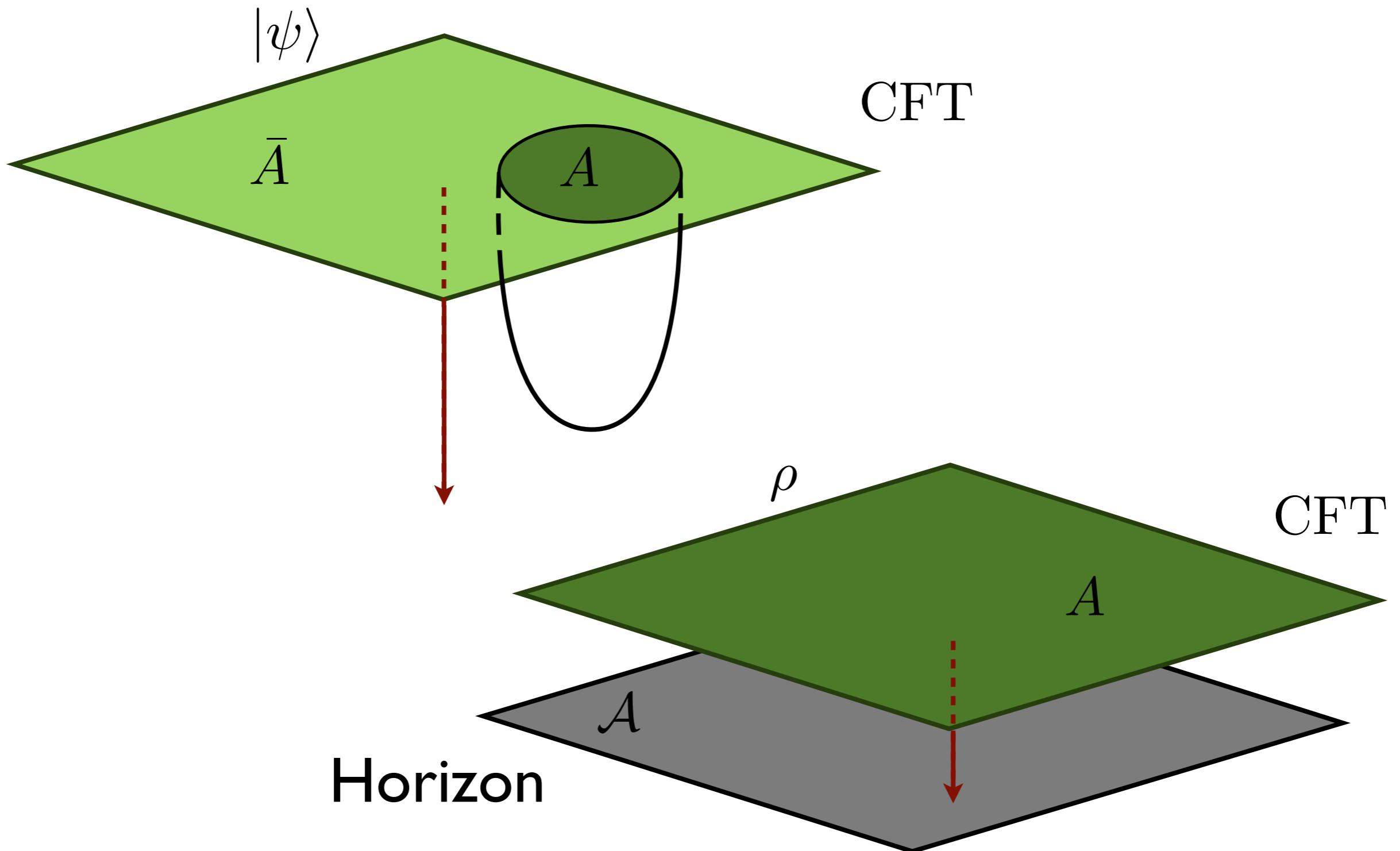
$$\text{tr}(\rho \log \rho) = \lim_{n \rightarrow 1} \frac{\log(\text{tr} \rho^n)}{n - 1}$$

# How to get there

$$S \sim \int_W \sqrt{\gamma} d^{D-2} \sigma \left( \frac{\partial \mathcal{L}}{\partial \text{Riem}} + \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2 \right)$$

# Setup

Casini, Huerta, Myers

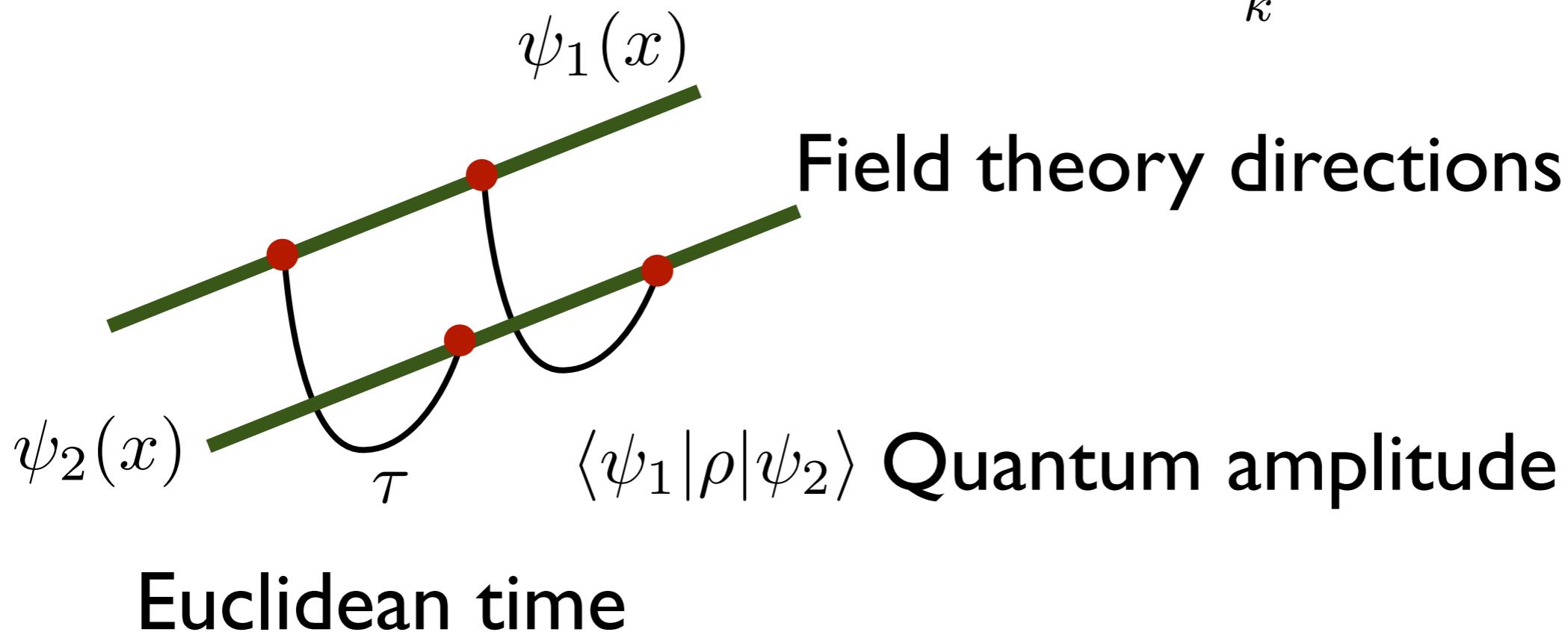


# States considered

- Generated by Euclidean paths integrals

$$\rho = P e^{- \int H(\tau) d\tau}$$

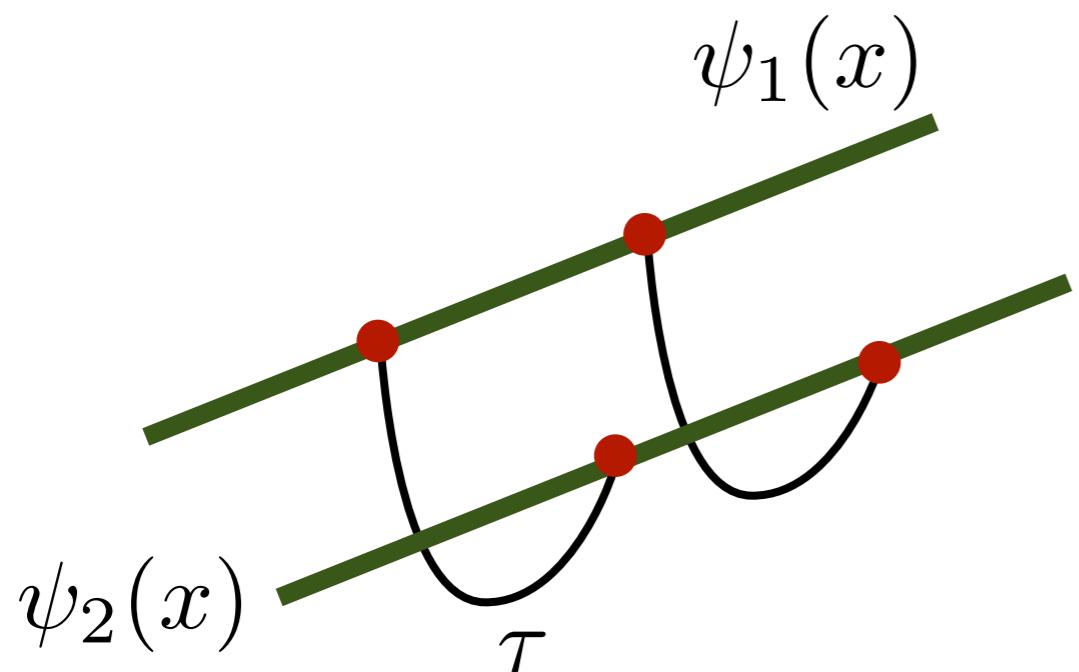
$$\rho_T = \sum_k |k\rangle e^{-E_k \beta} \langle k|$$



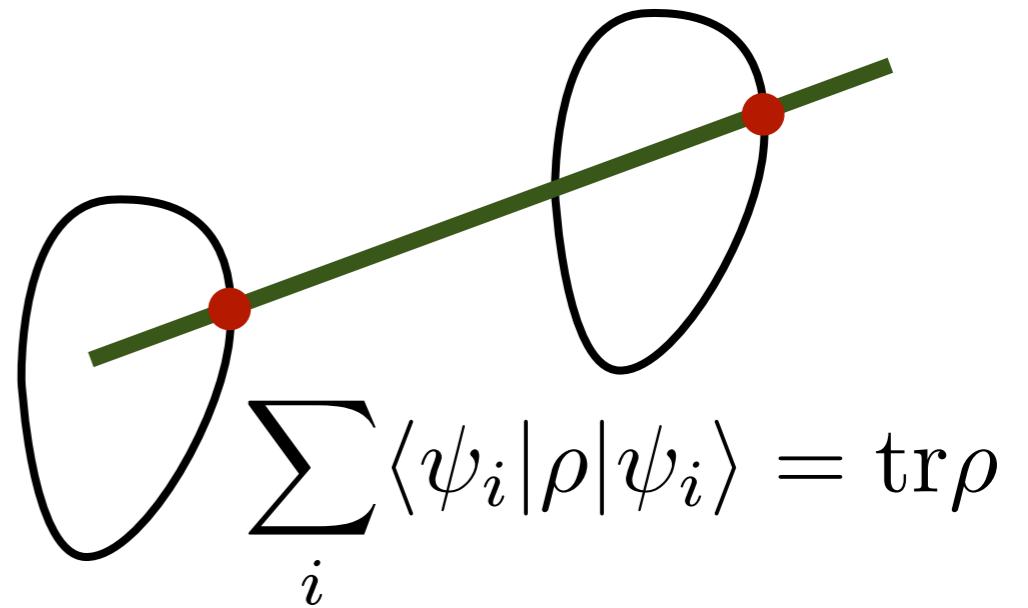
# States considered

- Generated by Euclidean paths integrals

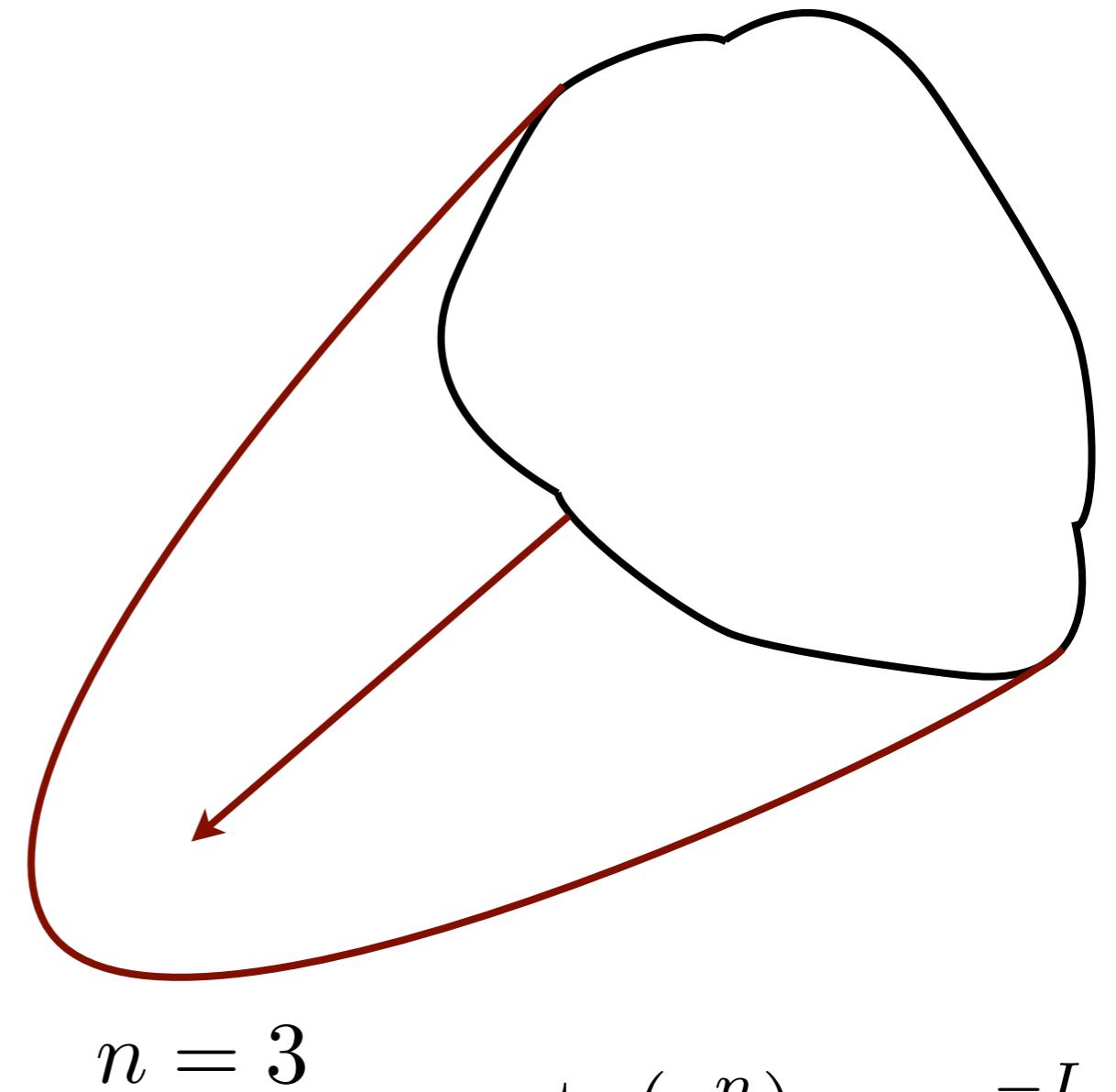
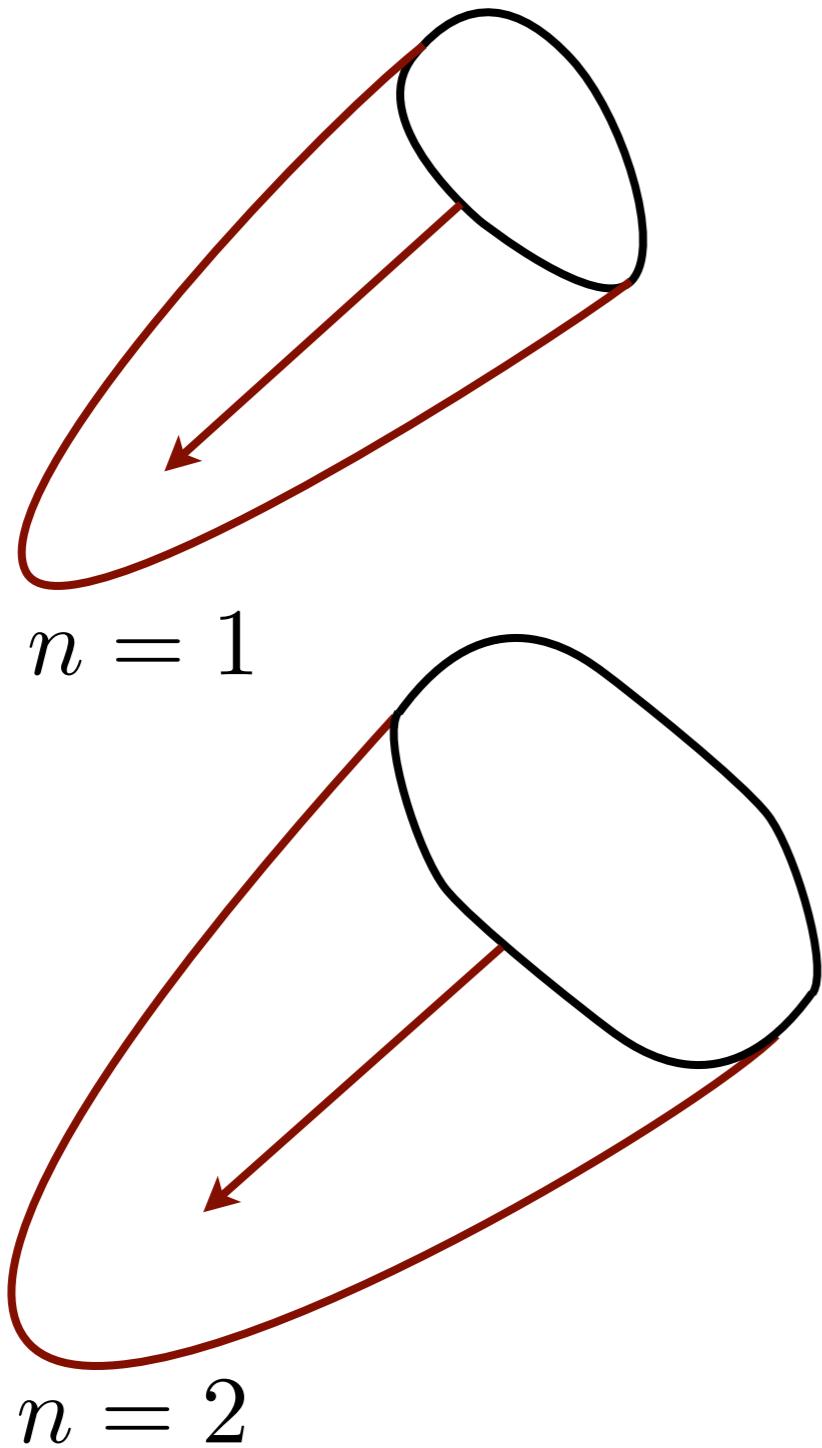
$$\rho = P e^{- \int H(\tau) d\tau}$$



$$\rho_T = \sum_k |k\rangle e^{-E_k \beta} \langle k|$$



# Gravity dual



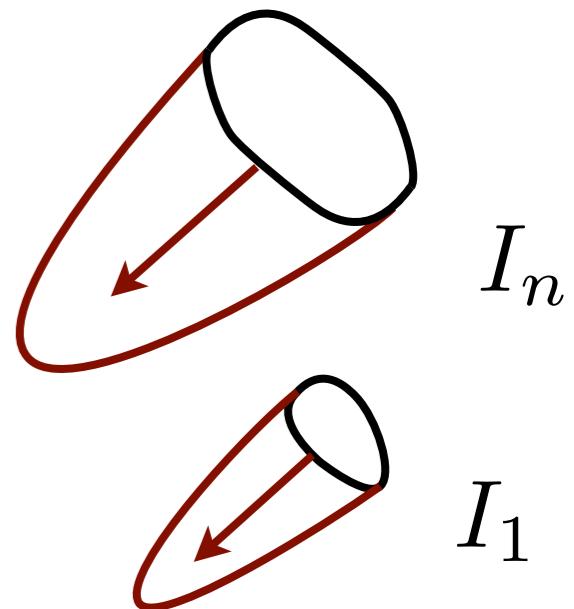
$$\frac{\text{tr}(\rho^n)}{(\text{tr}\rho)^n} \approx \frac{e^{-I_n}}{e^{-nI_1}}$$

# Generalized entropy

Entanglement entropy:  
Replica trick  $S = \lim_{n \rightarrow 1} S_n = -\partial_n \log \frac{\text{tr}(\rho^n)}{(\text{tr}\rho)^n} \Big|_{n=1}$

Gravity saddlepoint

$$\frac{\text{tr}(\rho^n)}{(\text{tr}\rho)^n} \approx \frac{e^{-I_n}}{e^{-nI_1}}$$

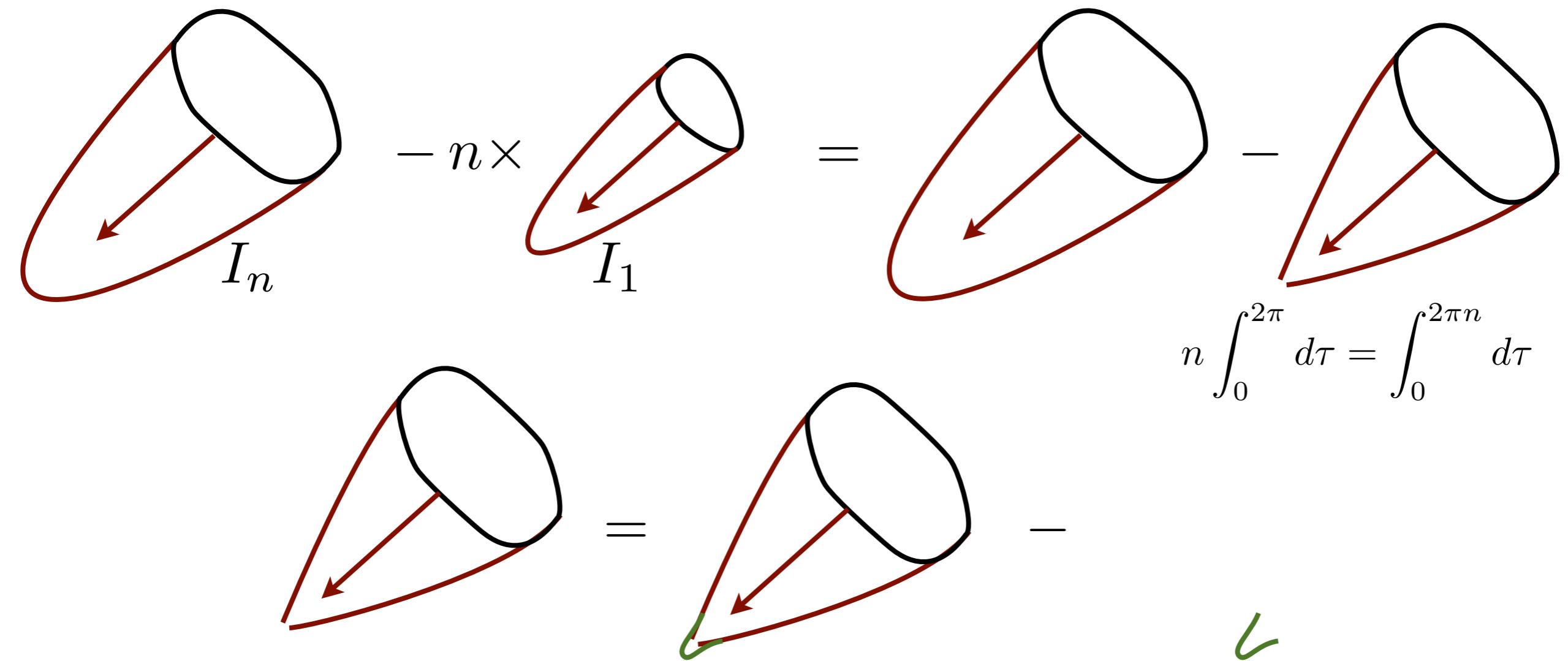


Gravitational entropy

$$S = \partial_n (I_n - nI_1) \Big|_{n=1}$$

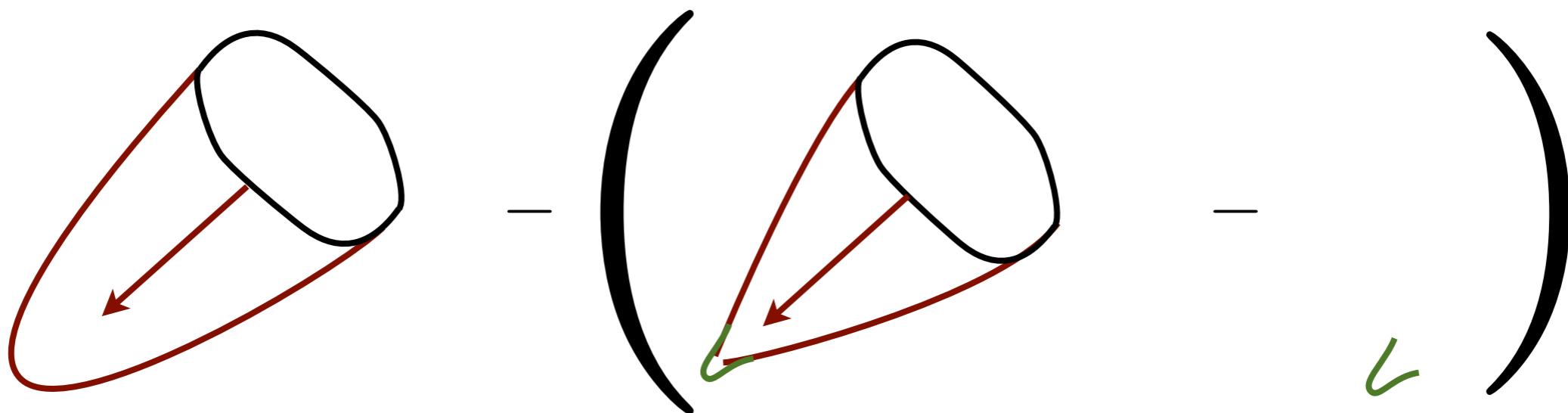
# Replica manipulations

$$S = \partial_n (I_n - nI_1)|_{n=1}$$



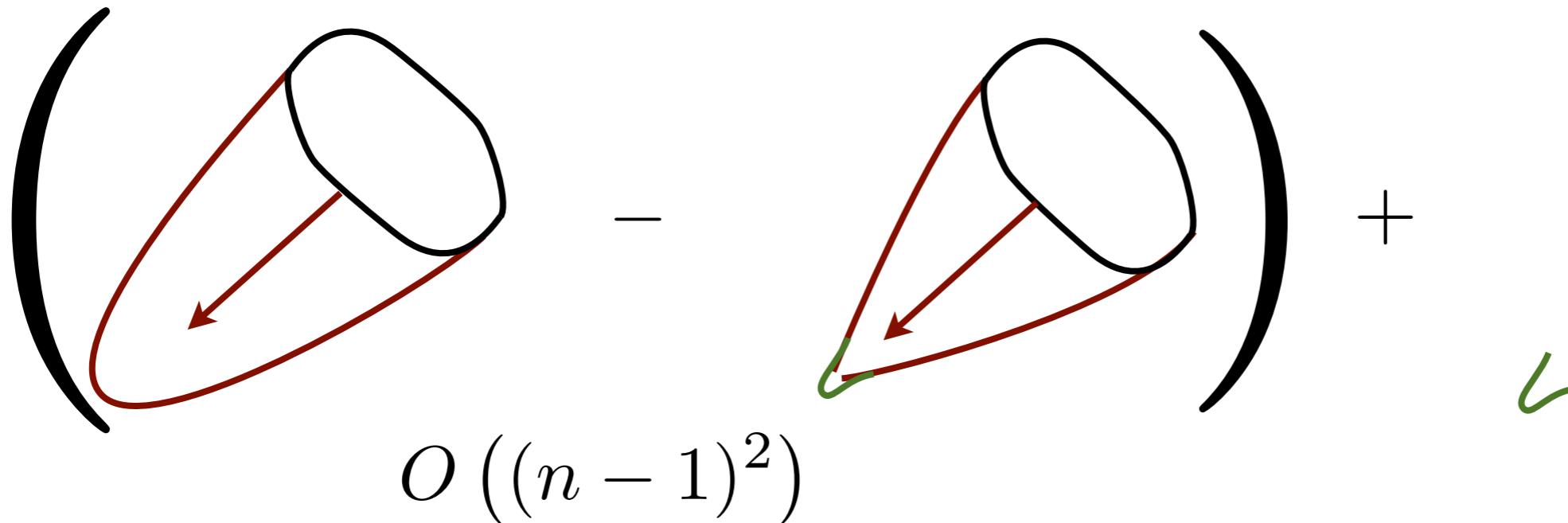
# Replica manipulations

$$S = \partial_n (I_n - nI_1) \Big|_{n=1}$$



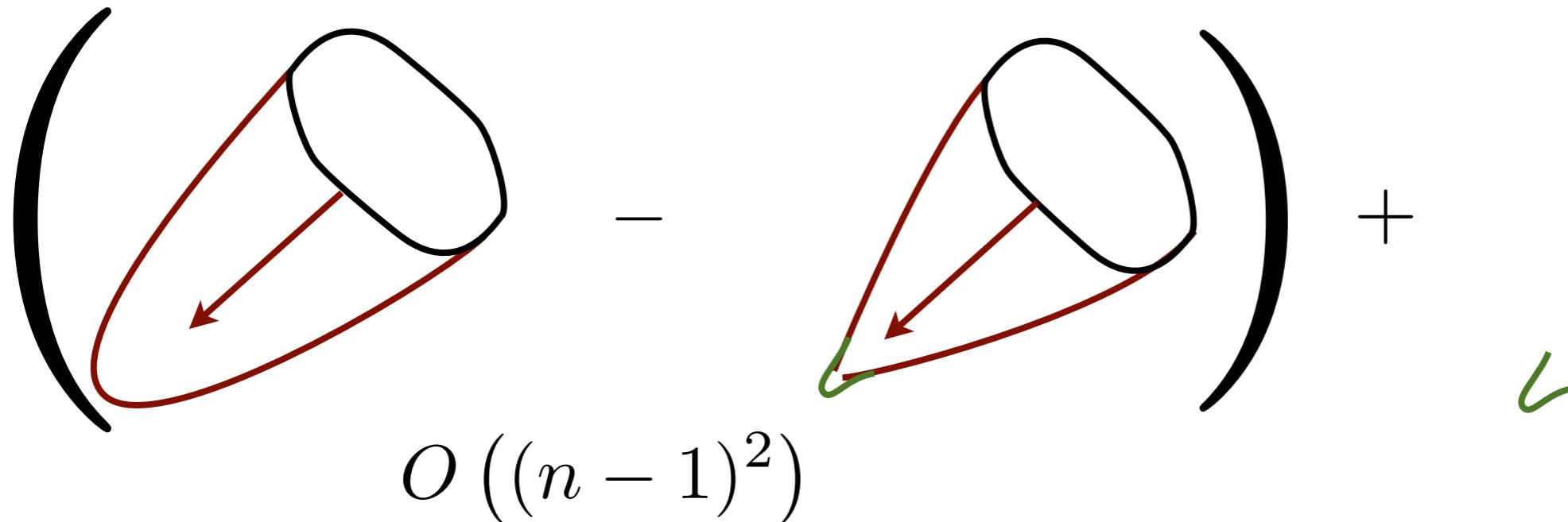
# Replica manipulations

$$S = \partial_n (I_n - nI_1)|_{n=1}$$



# Replica manipulations

$$S = \partial_n (I_n - nI_1)|_{n=1}$$



$$S = \partial_n (I[\mathcal{L}])|_{n=1}$$

# Recap

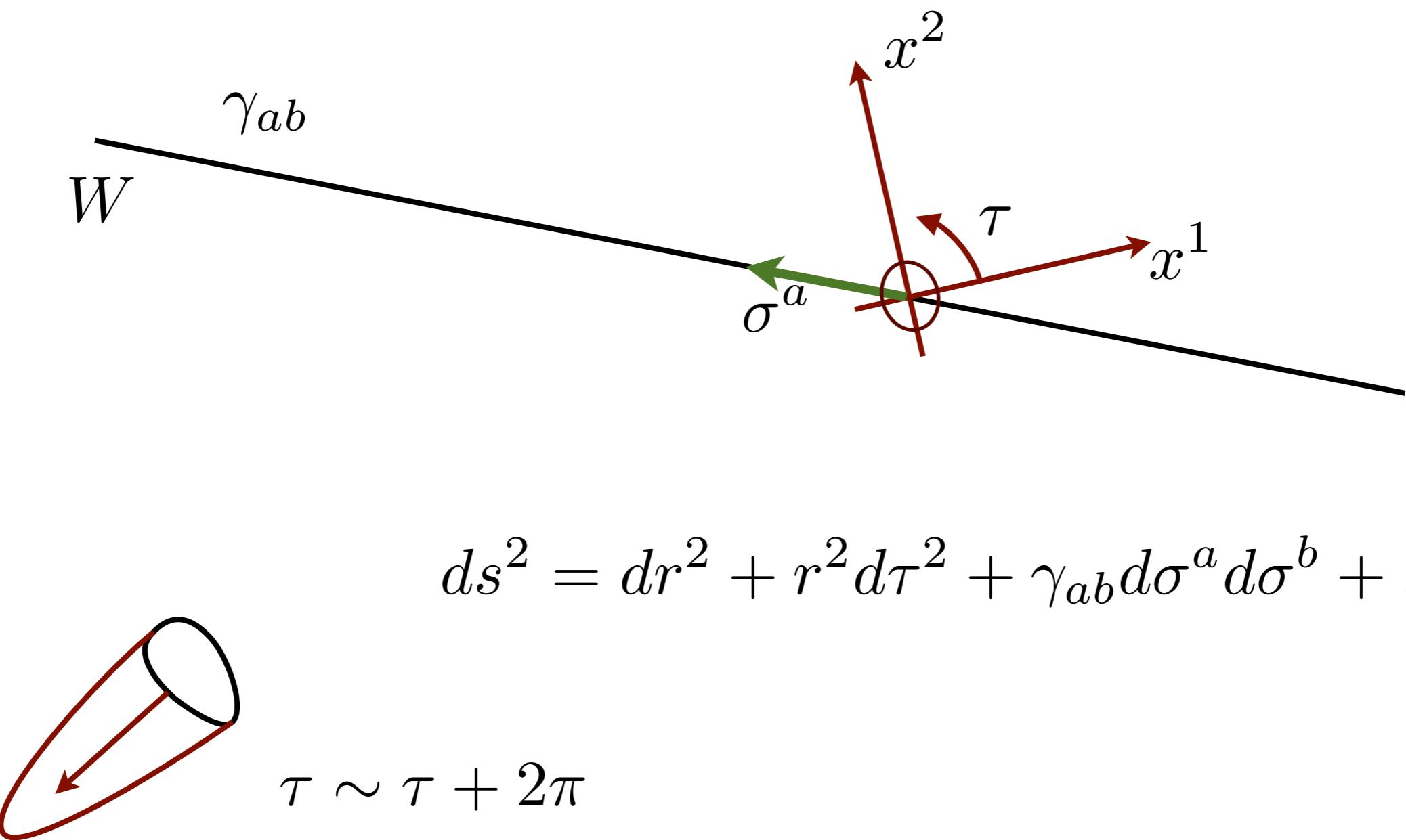
- Replica trick: EE and loops in eucl. time
- Assume ‘holography’  $Z_{\text{FT}} \approx e^{-I_n}$
- Entanglement entropy becomes:

$$S = \partial_n (I[\mathcal{L}])|_{n=1}$$

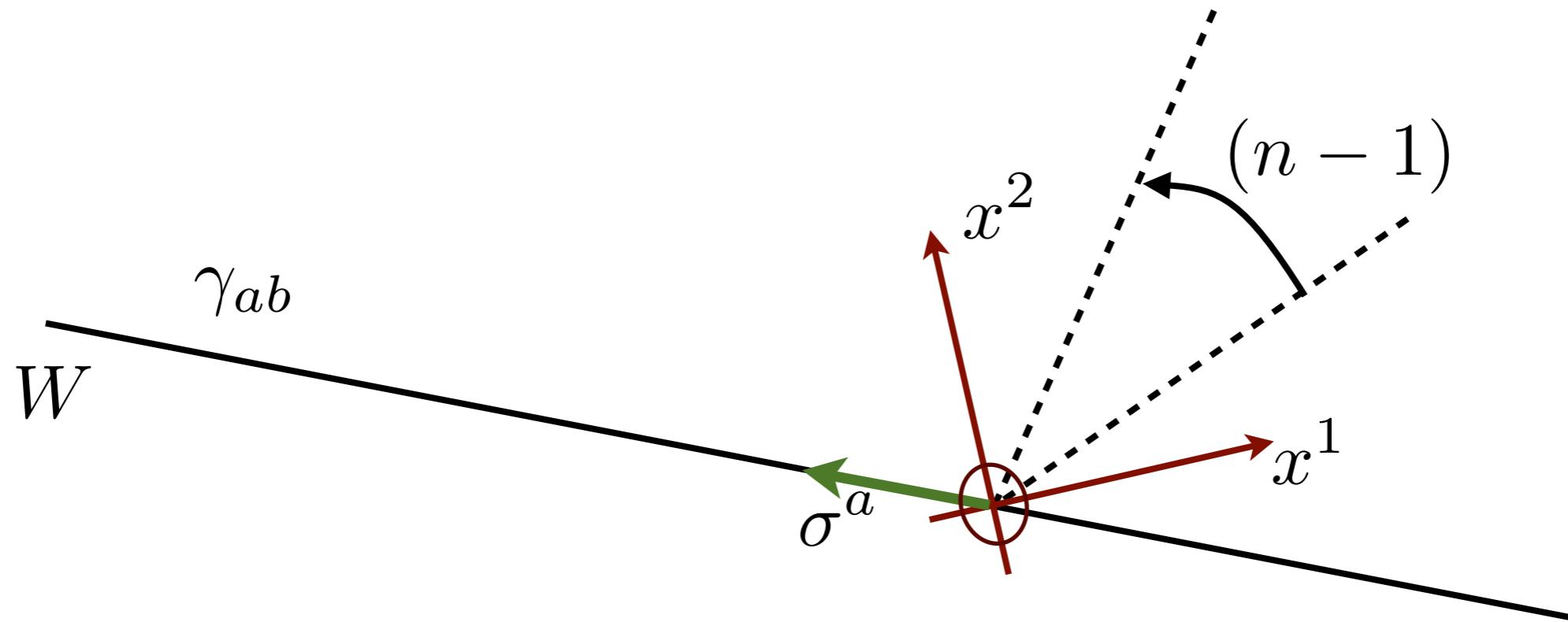
# The calculation

$$S = \partial_n (I[\mathcal{L}])|_{n=1}$$

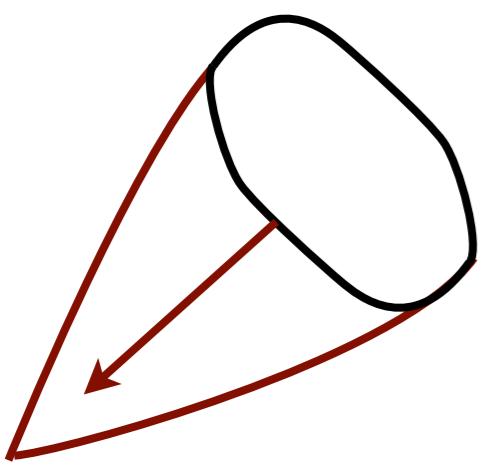
# Adapted coordinates



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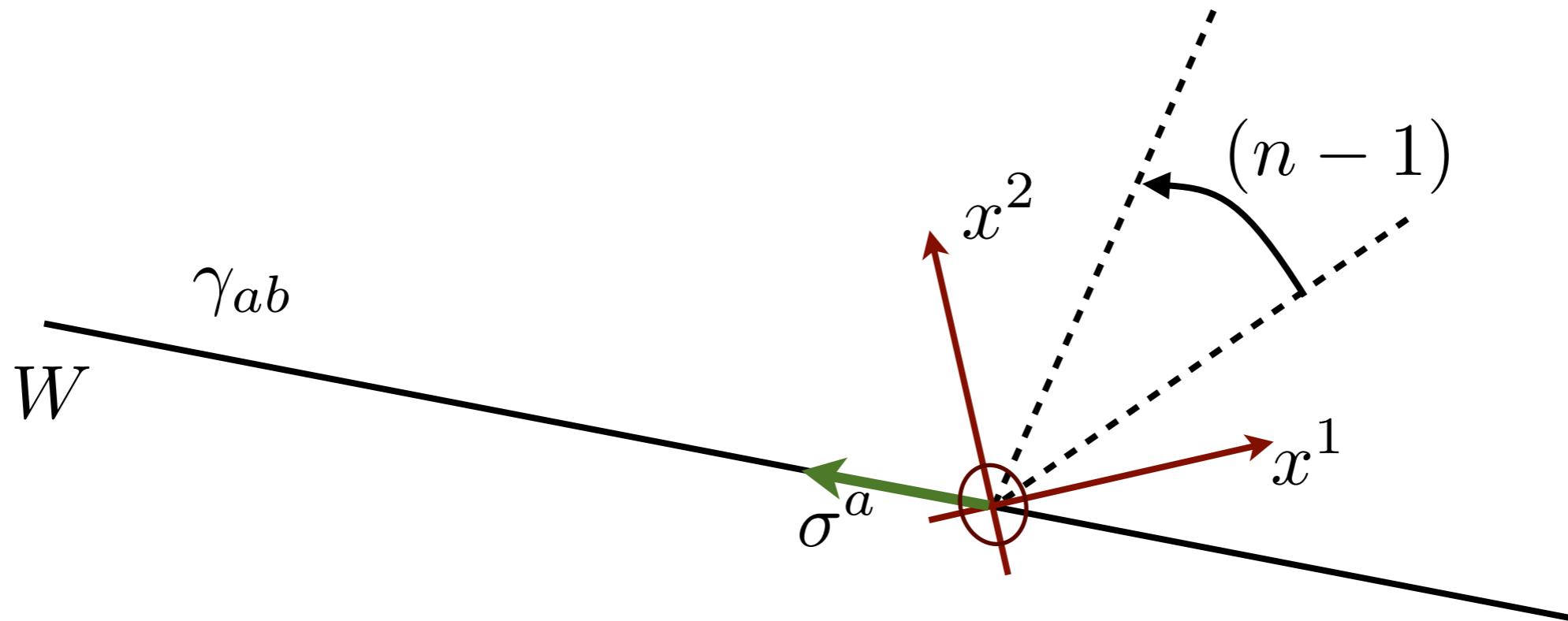


$$ds^2 = dr^2 + r^2 d\tau^2 + \gamma_{ab} d\sigma^a d\sigma^b + \dots$$



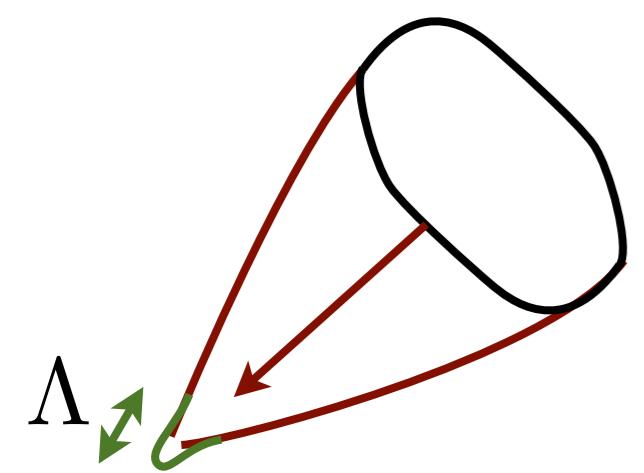
$$\tau \sim \tau + 2\pi n$$

# Adapted coordinates

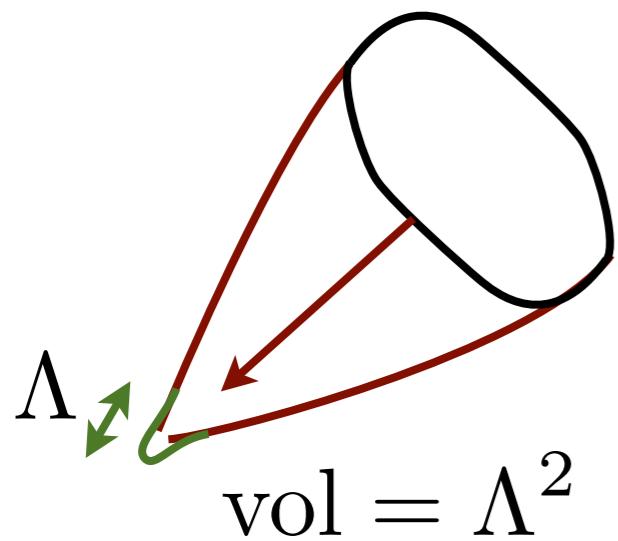


$$ds^2 = dr^2 + r^2 d\tau^2 + \gamma_{ab} d\sigma^a d\sigma^b + \dots$$

$$\tau \sim \tau + 2\pi n$$



# Entropy



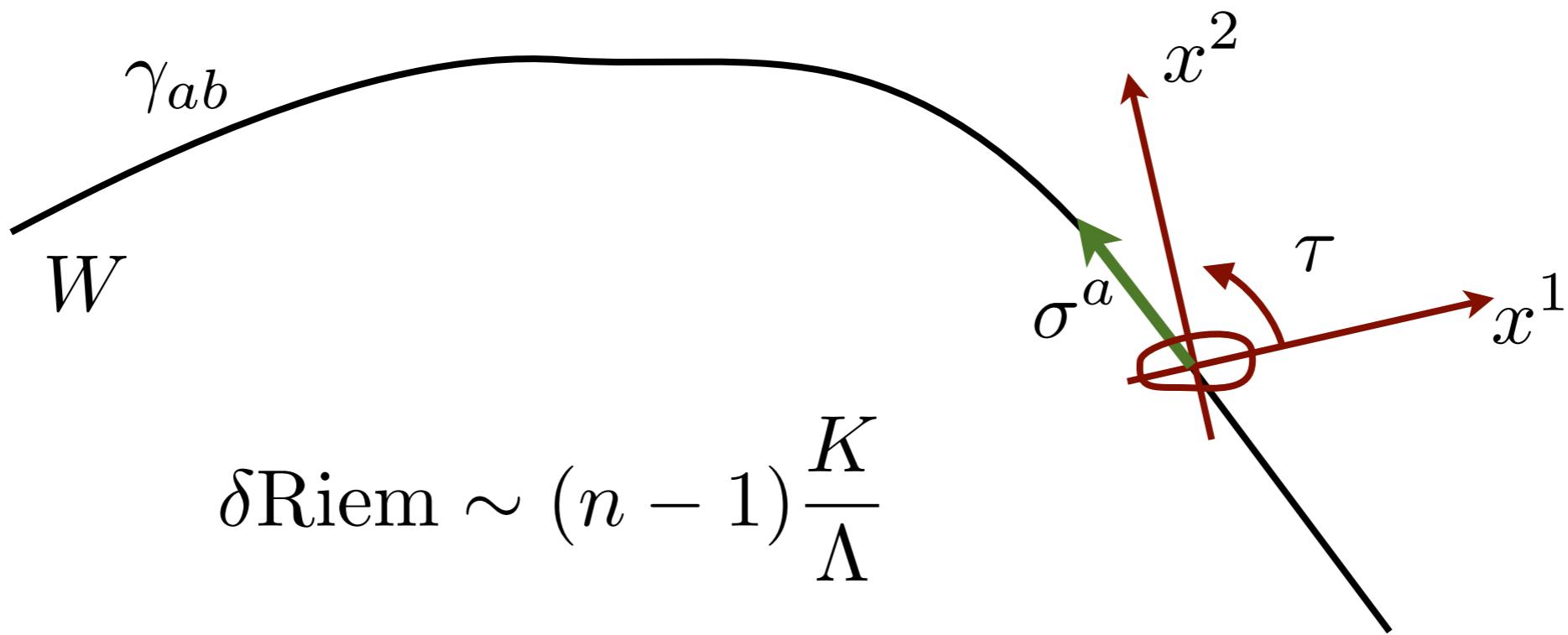
$$I = \int d^D x \mathcal{L}[\text{Riem}]$$

$\Lambda^2 \int \sqrt{\gamma} d^{D-2} \sigma$

$\delta \text{Riem} \sim \frac{n-1}{\Lambda^2}$

$$S = \partial_n (I[\text{L}])|_{n=1} = \int_W \sqrt{\gamma} d^{D-2} \sigma \frac{\partial \mathcal{L}}{\partial \text{Riem}}$$

# Extrinsic curvature



$$\delta \text{Riem} \sim (n - 1) \frac{K}{\Lambda}$$

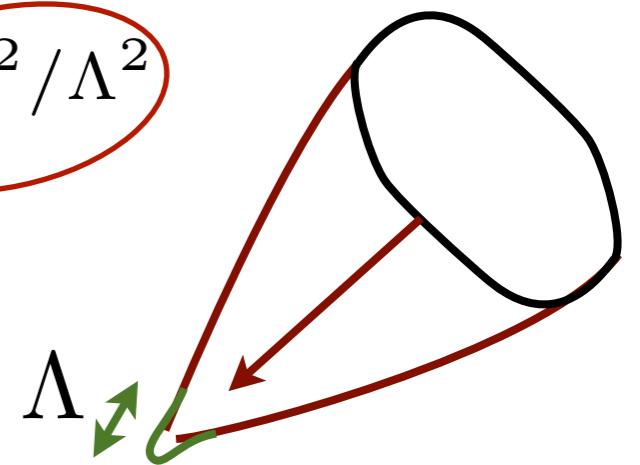
$$\Lambda \rightarrow 0$$

$$\delta S = \partial_n (\delta I[\mathcal{L}])|_{n=1} \sim \int_W \sqrt{\gamma} d^{D-2} \sigma \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2$$

# The subtlety

$$S = \partial_n (I[\mathcal{L}])|_{n=1}$$

$$\lim_{n \rightarrow 1} \partial_n \int_0^\infty dr (n-1)^2 r^{2n-3} e^{-r^2/\Lambda^2}$$



$$\lim_{n \rightarrow 1} \partial_n \left( (n-1)^2 \frac{\Gamma(n-1)}{2} \right) = \frac{1}{2}$$

$$\delta S \sim \int_W \sqrt{\gamma} d^{D-2} \sigma \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2$$

# Comments on the new formula

$$S \sim \int_W \sqrt{\gamma} d^{D-2} \sigma \left( \frac{\partial \mathcal{L}}{\partial \text{Riem}} + \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2 \right)$$

- It reduces to Wald's for stationary cases (trivially)
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# Final remarks & Open questions

- First principles derivation of Ryu-Takayanagi
- Euclidean space essential: Lorentzian?
- Non-stationary entropy? Second law?
- Nothing we did relies on replica symmetry;  
Where should we evaluate the entropy?

Cheng et al; Bhattacharyya et al; Dong; Erdmenger et al

# Extra slides

# On the location of $W$

- For Einstein Gravity, the LM construction locates  $W$  in GR by demanding regularity of a replica-symmetric ansatz, to  $O(n - 1)$
- A similar argument is not fully successful for more general theories  
Cheng et al; Bhattacharyya et al; Dong; Erdmenger et al
- That may hint to  
Replica Symmetry Breaking

JC & William Kelly. In progress

