

Coset Representation of Newton-Cartan Geometry

Kevin T. Grosvenor

Niels Bohr Institute, Copenhagen University



Johannes Gutenberg Universität, Mainz
March 14, 2018

Collaborators

Homogeneous Nonrelativistic Geometries as Coset Spaces,
[arXiv:1712.03980] with

Jelle Hartong

ITP and Delta ITP, University of Amsterdam

Cynthia Keeler

Arizona State University

Niels A. Obers

Niels Bohr Institute, Copenhagen University

Outline

Motivation

Relativistic Example: Minkowski and (A)dS

Non-relativistic Examples

Bargmann

Newton-Hooke

Schrödinger

TNC with $\mathbb{R} \times SO(3)$ isometries

Outlook

Motivation

- ▶ Probes of geometry (geodesics, Killing vectors, etc.)
- ▶ Nonrelativistic gravity (e.g. HLG)
- ▶ Nonrelativistic holography
- ▶ Condensed matter
- ▶ Strings

Relativistic Example: Minkowski and (A)dS

\mathfrak{g} = Poincaré, \mathfrak{h} = Lorentz, \mathfrak{m} = translations

$$[M_{ab}, M_{cd}] = \eta_{ac}M_{bd} \pm \dots \text{ and } [M_{ab}, P_c] = \eta_{ac}P_b - \eta_{bc}P_a$$

Relativistic Example: Minkowski and (A)dS

\mathfrak{g} = Poincaré, \mathfrak{h} = Lorentz, \mathfrak{m} = translations

$$[M_{ab}, M_{cd}] = \eta_{ac}M_{bd} \pm \dots \text{ and } [M_{ab}, P_c] = \eta_{ac}P_b - \eta_{bc}P_a$$

Associated structure constants: $f_{\mathfrak{h}\mathfrak{m}}^{\mathfrak{m}} = f_{[cd]a}^b = \eta_{ca}\delta_d^b - \eta_{da}\delta_c^b$

Solution to ad- \mathfrak{h} inv. symm. bil. form $f_{\mathfrak{h}(a}^c\Omega_{b)c} = 0$ is $\Omega_{ab} \propto \eta_{ab}$

Relativistic Example: Minkowski and (A)dS

\mathfrak{g} = Poincaré, \mathfrak{h} = Lorentz, \mathfrak{m} = translations

$$[M_{ab}, M_{cd}] = \eta_{ac}M_{bd} \pm \dots \text{ and } [M_{ab}, P_c] = \eta_{ac}P_b - \eta_{bc}P_a$$

Associated structure constants: $f_{\mathfrak{h}\mathfrak{m}}^{\mathfrak{m}} = f_{[cd]a}^b = \eta_{ca}\delta_d^b - \eta_{da}\delta_c^b$

Solution to ad- \mathfrak{h} inv. symm. bil. form $f_{\mathfrak{h}(a}^c\Omega_{b)c} = 0$ is $\Omega_{ab} \propto \eta_{ab}$

$$\mathcal{G} \ni \underbrace{g = \exp(P_a \delta_\mu^a x^\mu)}_{\text{Lie group element}} \implies \underbrace{g^{-1}dg = P_a \delta_\mu^a dx^\mu}_{\text{Maurer-Cartan form}} \implies \underbrace{e^a = \delta_\mu^a dx^\mu}_{\text{vielbein}}$$

Metric on $\mathcal{M} = \mathcal{G}/\mathcal{H}$ is $ds^2 = \Omega_{ab}e^ae^b = \eta_{\mu\nu}dx^\mu dx^\nu$

Relativistic Example: Minkowski and (A)dS

\mathfrak{g} = Poincaré, \mathfrak{h} = Lorentz, \mathfrak{m} = translations

$$[M_{ab}, M_{cd}] = \eta_{ac}M_{bd} \pm \dots \text{ and } [M_{ab}, P_c] = \eta_{ac}P_b - \eta_{bc}P_a$$

Associated structure constants: $f_{\mathfrak{h}\mathfrak{m}}^{\mathfrak{m}} = f_{[cd]a}^b = \eta_{ca}\delta_d^b - \eta_{da}\delta_c^b$

Solution to ad- \mathfrak{h} inv. symm. bil. form $f_{\mathfrak{h}(a}{}^c\Omega_{b)c} = 0$ is $\Omega_{ab} \propto \eta_{ab}$

$$\mathcal{G} \ni \underbrace{g = \exp(P_a \delta_\mu^a x^\mu)}_{\text{Lie group element}} \implies \underbrace{g^{-1}dg = P_a \delta_\mu^a dx^\mu}_{\text{Maurer-Cartan form}} \implies \underbrace{e^a = \delta_\mu^a dx^\mu}_{\text{vielbein}}$$

Metric on $\mathcal{M} = \mathcal{G}/\mathcal{H}$ is $ds^2 = \Omega_{ab}e^a e^b = \eta_{\mu\nu}dx^\mu dx^\nu$

With $[P_a, P_b] = \lambda M_{ab}$, this gives (A)dS

Relativistic Example: Minkowski and (A)dS

\mathfrak{g} = Poincaré, \mathfrak{h} = Lorentz, \mathfrak{m} = translations

$$[M_{ab}, M_{cd}] = \eta_{ac}M_{bd} \pm \dots \text{ and } [M_{ab}, P_c] = \eta_{ac}P_b - \eta_{bc}P_a$$

Associated structure constants: $f_{\mathfrak{h}\mathfrak{m}}^{\mathfrak{m}} = f_{[cd]a}^b = \eta_{ca}\delta_d^b - \eta_{da}\delta_c^b$

Solution to ad- \mathfrak{h} inv. symm. bil. form $f_{\mathfrak{h}(a}^c\Omega_{b)c} = 0$ is $\Omega_{ab} \propto \eta_{ab}$

$$\mathcal{G} \ni \underbrace{g = \exp(P_a \delta_\mu^a x^\mu)}_{\text{Lie group element}} \implies \underbrace{g^{-1}dg = P_a \delta_\mu^a dx^\mu}_{\text{Maurer-Cartan form}} \implies \underbrace{e^a = \delta_\mu^a dx^\mu}_{\text{vielbein}}$$

Metric on $\mathcal{M} = \mathcal{G}/\mathcal{H}$ is $ds^2 = \Omega_{ab}e^ae^b = \eta_{\mu\nu}dx^\mu dx^\nu$

With $[P_a, P_b] = \lambda M_{ab}$, this gives (A)dS

There is also Ω^{ab} such that $\Omega^{c(a}f_{bc}^b) = 0$, but Ω_{ab} and Ω^{ab} are inverses.

Bargmann

cf. 1407.7730 [Brauner, Endlich, Monin, Penco]

$\mathfrak{g} = \text{Barg}$, $\mathfrak{h} = \{J, G_i, N\}$, $\mathfrak{m} = \{H, P_i\}$ and $g = e^{Ht} e^{P_i x^i}$

$\Omega_{ab} e^a e^b = -dt^2$ (degenerate!)

But, contravariant version $\Omega^{c(a} f_{lc}{}^{b)} = 0$ gives $\Omega^{ab} e_a e_b = \delta^{ij} \partial_i \partial_j$

$$\tau = dt, \quad h^{\mu\nu} \partial_\mu \partial_\nu = \delta^{ij} \partial_i \partial_j, \quad m = 0.$$

Loc. tan. space trans.: $g \rightarrow gh$ with $h \in \mathcal{H}$

From Relativistic to Non-relativistic

Two derivations of Bargmann coset from relativistic starting point:

- Inönü-Wigner contraction

$$\begin{array}{ccc} \text{Poinc} \times U(1) & \xrightarrow{c \rightarrow \infty} & \text{Barg} \\ \downarrow \text{coset} & & \downarrow \text{coset} \\ \Omega_{ab} = \begin{pmatrix} -1 & 0 \\ 0 & \frac{1}{c^2} \end{pmatrix} & \xrightarrow{c \rightarrow \infty} & \Omega_{ab} = \begin{pmatrix} -1 & 0 \\ 0 & 0 \end{pmatrix} \end{array}$$

From Relativistic to Non-relativistic

Two derivations of Bargmann coset from relativistic starting point:

- ▶ Inönü-Wigner contraction

$$\begin{array}{ccc} \text{Poinc} \times U(1) & \xrightarrow{c \rightarrow \infty} & \text{Barg} \\ \downarrow \text{coset} & & \downarrow \text{coset} \\ \Omega_{ab} = \begin{pmatrix} -1 & 0 \\ 0 & \frac{1}{c^2} \end{pmatrix} & \xrightarrow{c \rightarrow \infty} & \Omega_{ab} = \begin{pmatrix} -1 & 0 \\ 0 & 0 \end{pmatrix} \end{array}$$

- ▶ Null-reduction ($N \in \mathfrak{m}$, with associated coordinate u)

$$ds^2 = 2 \underbrace{dt}_{\tau} \left(du - \underbrace{dt}_{m} \right) + \underbrace{\delta_{ij} dx^i dx^j}_{h_{\mu\nu} dx^\mu dx^\nu}$$

Newton-Hooke

Add commutator $[H, P_i] = \Lambda G_i$

$$\begin{aligned}\tau &= dt \\ h^{\mu\nu} &= \delta^{ij} \delta_i^\mu \delta_j^\nu \\ m &= -\frac{1}{2} \Lambda x^2 dt\end{aligned}$$

\Rightarrow

$$\begin{aligned}\lambda^i &= -\sqrt{\Lambda} x^i \\ \sigma &= \frac{1}{2} \sqrt{\Lambda} x^2 \\ t' &= t \\ x'^i &= e^{-\sqrt{\Lambda} t} x^i\end{aligned}$$

\Rightarrow

$$\begin{aligned}\tau' &= dt' \\ h'^{\mu\nu} &= \\ &e^{-2\sqrt{\Lambda} t'} \delta^{ij} \delta_i^\mu \delta_j^\nu \\ m' &= 0\end{aligned}$$

Newton-Hooke

Add commutator $[H, P_i] = \Lambda G_i$

$$\begin{aligned}\tau &= dt \\ h^{\mu\nu} &= \delta^{ij} \delta_i^\mu \delta_j^\nu \\ m &= -\frac{1}{2} \Lambda x^2 dt\end{aligned}$$

\Rightarrow

$$\begin{aligned}\lambda^i &= -\sqrt{\Lambda} x^i \\ \sigma &= \frac{1}{2} \sqrt{\Lambda} x^2 \\ t' &= t \\ x'^i &= e^{-\sqrt{\Lambda} t} x^i\end{aligned}$$

\Rightarrow

$$\begin{aligned}\tau' &= dt' \\ h'^{\mu\nu} &= \\ &e^{-2\sqrt{\Lambda} t'} \delta^{ij} \delta_i^\mu \delta_j^\nu \\ m' &= 0\end{aligned}$$

Can also derive via contraction or null-reduction.

Schrödinger

cf. 0903.4245 [Schäfer-Nameki, Yamazaki, Yoshida]

$\mathfrak{h} = \{J, G_i, N + 2K\}$ and $\mathfrak{m} = \{N, H, P_i, D\}$

$$ds^2 = 2 \frac{dt}{r^2} \left(du - \frac{\alpha}{2} \frac{dt}{r^2} \right) + \frac{dr^2 + dx^2}{r^2} \xrightarrow{\text{n-red}}$$

$$\begin{aligned} \tau &= \frac{dt}{r^2} \\ h^{\mu\nu} &= r^2 (\delta_r^\mu \delta_r^\nu + \delta_i^\mu \delta_i^\nu) \\ m &= \frac{\alpha}{2} \frac{dt}{r^2} \end{aligned}$$

Schrödinger

cf. 0903.4245 [Schäfer-Nameki, Yamazaki, Yoshida]

$$\mathfrak{h} = \{J, G_i, N + 2K\} \text{ and } \mathfrak{m} = \{N, H, P_i, D\}$$

$$ds^2 = 2 \frac{dt}{r^2} \left(du - \frac{\alpha}{2} \frac{dt}{r^2} \right) + \frac{dr^2 + dx^2}{r^2} \xrightarrow{\text{n-red}}$$

$$\begin{aligned} \tau &= \frac{dt}{r^2} \\ h^{\mu\nu} &= r^2 (\delta_r^\mu \delta_r^\nu + \delta_i^\mu \delta_i^\nu) \\ m &= \frac{\alpha}{2} \frac{dt}{r^2} \end{aligned}$$

Can get this TTNC via $\mathfrak{h} = \{J, G_i, N, K\}$ and $\mathfrak{m} = \{H - \frac{\alpha}{2} N, P_i, D\}$.

TNC with $\mathbb{R} \times SO(3)$ isometries

$$\mathfrak{g} = \underbrace{\mathfrak{so}(3)}_{J_a} \times \underbrace{\mathfrak{u}(1)}_H \times \underbrace{\mathfrak{u}(1)}_N \text{ and rescale } J \equiv J_3 \text{ and } P_{i=1,2} = \frac{1}{R} J_i$$

$$\mathfrak{h} = \{J + aN + bH, N\} \text{ and } \mathfrak{m} = \{H, P_i\}$$

$$g = e^{H\nu} e^{P_1\alpha_1} e^{P_2\alpha_2} \text{ and redefine } \alpha_1 = R\phi \text{ and } \alpha_2 = R\left(\eta - \frac{\pi}{2}\right).$$

Ω_{ab} is diagonal and, in general, non-degenerate.

Choose $\text{sign}\Omega_{ab} = (1, 0, 0)$ and $\text{sign}\Omega^{ab} = (0, 1, 1)$.

TNC with $\mathbb{R} \times SO(3)$ isometries

$$\mathfrak{g} = \underbrace{\mathfrak{so}(3)}_{J_a} \times \underbrace{\mathfrak{u}(1)}_H \times \underbrace{\mathfrak{u}(1)}_N \text{ and rescale } J \equiv J_3 \text{ and } P_{i=1,2} = \frac{1}{R} J_i$$

$$\mathfrak{h} = \{J + aN + bH, N\} \text{ and } \mathfrak{m} = \{H, P_i\}$$

$$g = e^{Hv} e^{P_1 \alpha_1} e^{P_2 \alpha_2} \text{ and redefine } \alpha_1 = R\phi \text{ and } \alpha_2 = R\left(\eta - \frac{\pi}{2}\right).$$

Ω_{ab} is diagonal and, in general, non-degenerate.

Choose $\text{sign}\Omega_{ab} = (1, 0, 0)$ and $\text{sign}\Omega^{ab} = (0, 1, 1)$.

$$\begin{aligned} \tau &= dv + b \cos \eta d\phi \\ h_{\mu\nu} dx^\mu dx^\nu &= \\ R^2 (d\eta^2 + \sin^2 \eta d\phi^2) \\ m &= a \cos \eta d\phi \end{aligned}$$

Outlook

Classification problems (algebras to coset spaces)

Quotient out by discrete subgroups (e.g., analog of BTZ)

Solutions of HLG or CS, cf. 1712.05794 [Hartong, Lei, Obers, Oling]

Nonrelativistic dual for Schrödinger field theories (or Lifshitz?)

Entanglement entropy and Ryu-Takayanagi.