

MITP Workshop
Proton Radius Puzzle
Waldthausen Castle,
Mainz, Germany

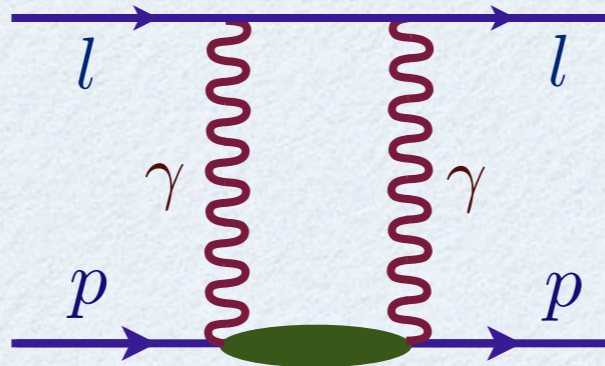
4 June, 2014

*Two-photon exchange corrections
in elastic ep and μp scattering.
Elastic contribution. Dispersive framework*

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Mainz, Germany

Outline

- Motivation
- Previous two-photon exchange (TPE) calculations
- Dispersion relation framework for 2γ corrections
- Results for elastic ep scattering and polarisation transfer
 - Estimations for elastic μp scattering



Proton radius puzzle

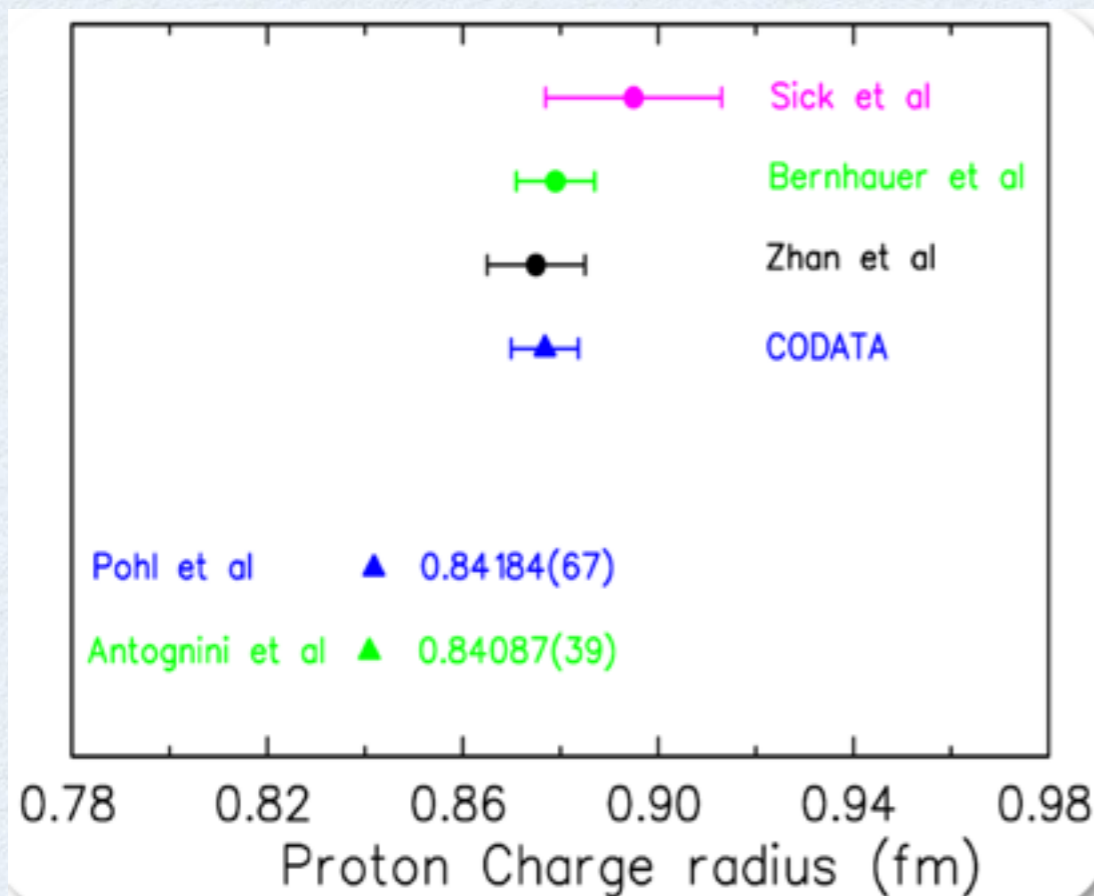


Lamb shift

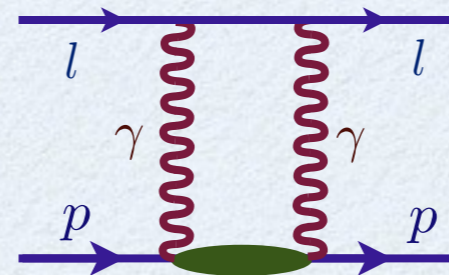
	e hydrogen	$r_E = 0.8768 \pm 0.0069 \text{ fm}$	$r_E = 0.8409 \pm 0.0004 \text{ fm}$
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ep-elastic scattering

		$r_E = 0.870 \pm 0.025 \text{ fm}$	7.7σ difference!
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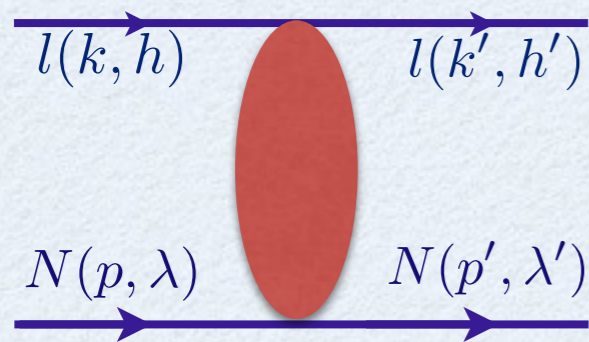


TPE hadronic correction is dominant uncertainty in scattering experiments



$$\sigma^{exp} \equiv \sigma_{1\gamma}(1 + \delta_{soft} + \delta_{2\gamma})$$

Structure amplitudes: TPE correction



$$Q^2 = -(k - k')^2$$

$$s = (p + k)^2$$

$$u = (k - p')^2$$

$$\nu = \frac{s - u}{4}$$

$$\epsilon$$

momentum transfer

crossing symmetric variable

photon polarization parameter

Discrete symmetries



6 structure amplitudes

$$T^{non-flip} = \frac{e^2}{Q^2} \bar{l}(k', h') \gamma_\mu l(k, h) \cdot \bar{N}(p', \lambda') [\mathcal{G}_M(\nu, t) \gamma^\mu - \mathcal{F}_2(\nu, t) \frac{P^\mu}{M} + \mathcal{F}_3(\nu, t) \frac{\hat{K} P^\mu}{M^2}] N(p, \lambda)$$

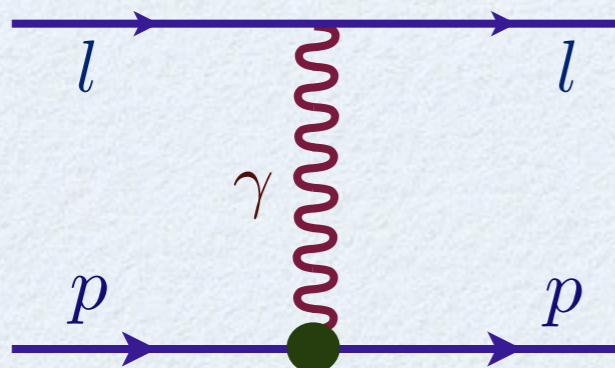
P.A.M. Guichon and M. Vanderhaeghen (2003)

$$T^{flip} \sim \mathcal{F}_4, \mathcal{F}_5, \mathcal{F}_6$$

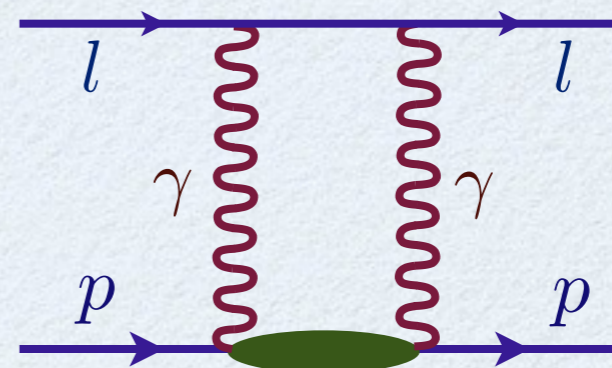
M. Gorchtein, P.A.M. Guichon and M. Vanderhaeghen (2004)

Leading TPE contribution to cross section - interference term

1 photon diagram



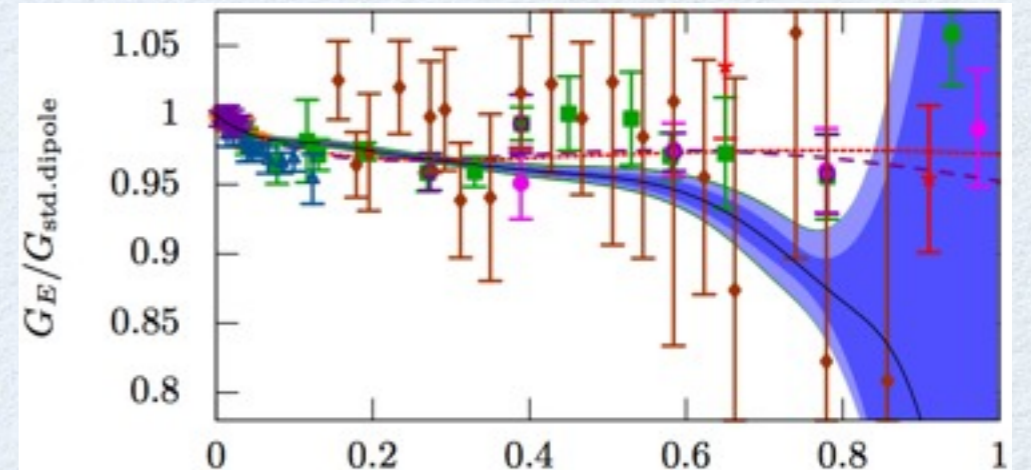
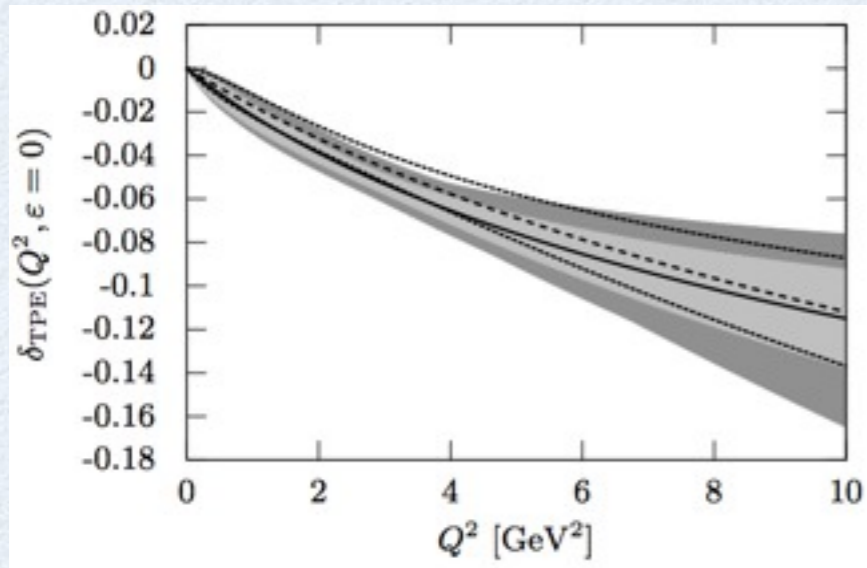
2 photon exchange diagram



$$\delta_{2\gamma} \sim \Re \mathcal{G}_M, \Re \mathcal{F}_2, \Re \mathcal{F}_3, \Re \mathcal{F}_4, \Re \mathcal{F}_5$$

Dispersion relation framework

2γ corrections



$f(z)$
analyticity

cross section correction

exp. data/phenomenology

$$\Re \mathcal{F}(\nu) = \frac{2\nu}{\pi} \mathcal{P} \int_{\nu_{min}}^{\infty} \frac{\Im \mathcal{F}(\nu' + i0)}{\nu'^2 - \nu^2} d\nu'$$

unitarity

amplitudes: real parts

DR

amplitudes: imaginary parts

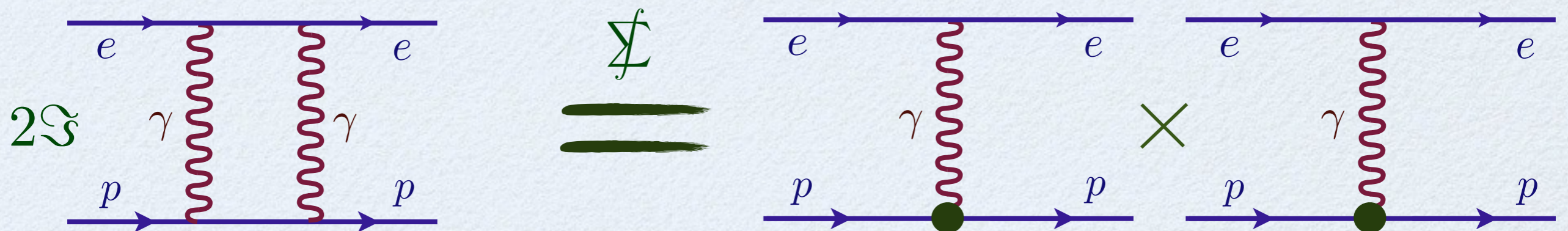
Unitarity relations. Imaginary part

$$S = 1 + iT \quad S^\dagger S = 1 \quad \longrightarrow \quad \Im T_{h'\lambda', h\lambda}$$

only on-shell information is required

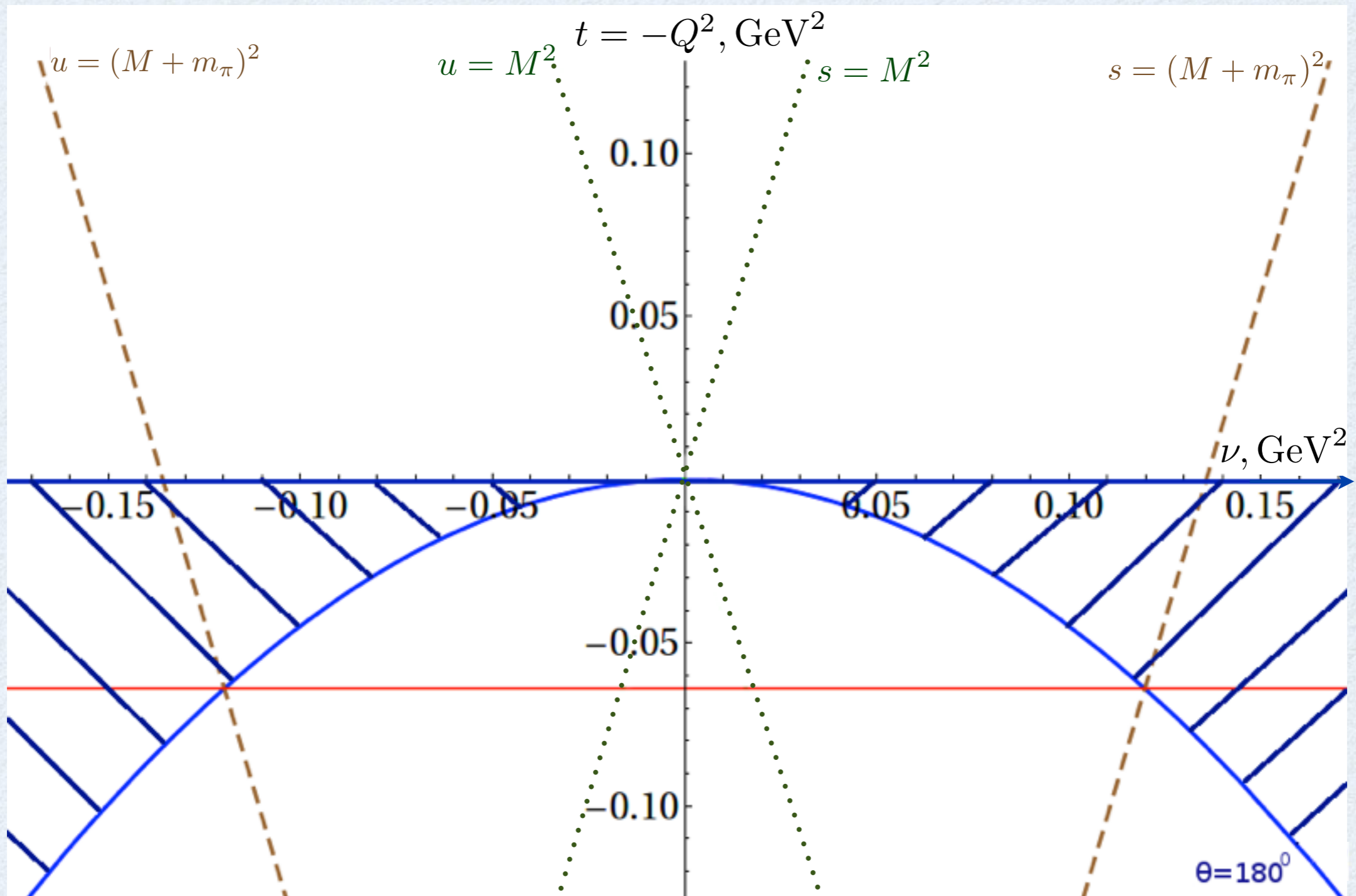
$$2\Im T_{h'\lambda', h\lambda} = \int d\Pi'' T_{h'\lambda', \mu}^+ T_{\mu, h\lambda} (2\pi)^4 \delta^4(k + p - \sum_i q_i)$$

e and N intermediate state



on-shell one-photon amplitudes

Kinematic regions (e-p)



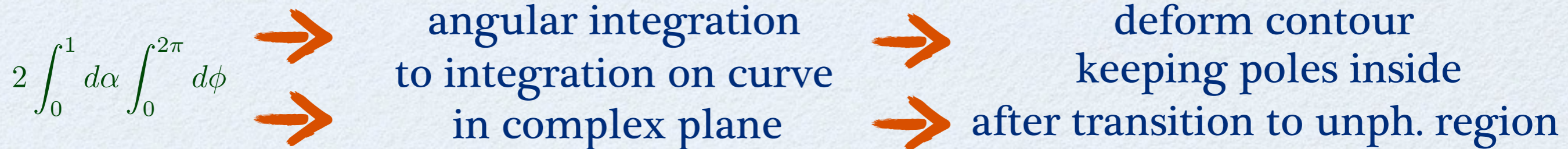
$Q^2 = m_\pi^2 \left(\frac{2M + m_\pi}{M + m_\pi} \right)^2 \sim 3.5 m_\pi^2$ - intersection of phys. region and inelastic threshold

Proton intermediate state is **outside** physical region

Analytical continuation

$\int d\Omega$ symmetric coordinates wrt electron momentum transfer

$$\cos \theta_1 = \sqrt{1 - \alpha^2} (b \cos \phi + c \sin \phi) \quad \cos \theta_2 = \sqrt{1 - \alpha^2} (b \cos \phi - c \sin \phi)$$

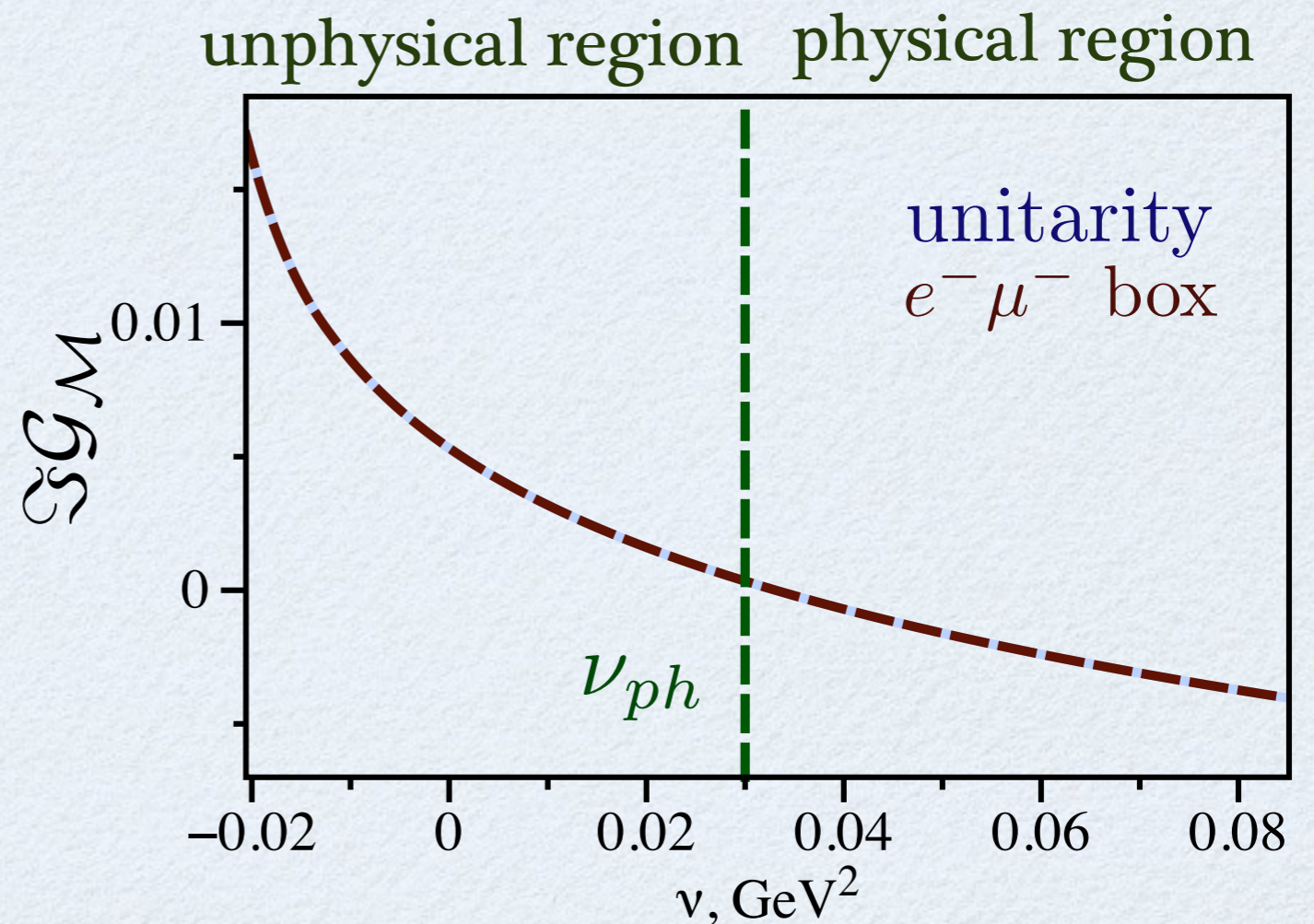


Analytical continuation
reproduces results
in unphysical region

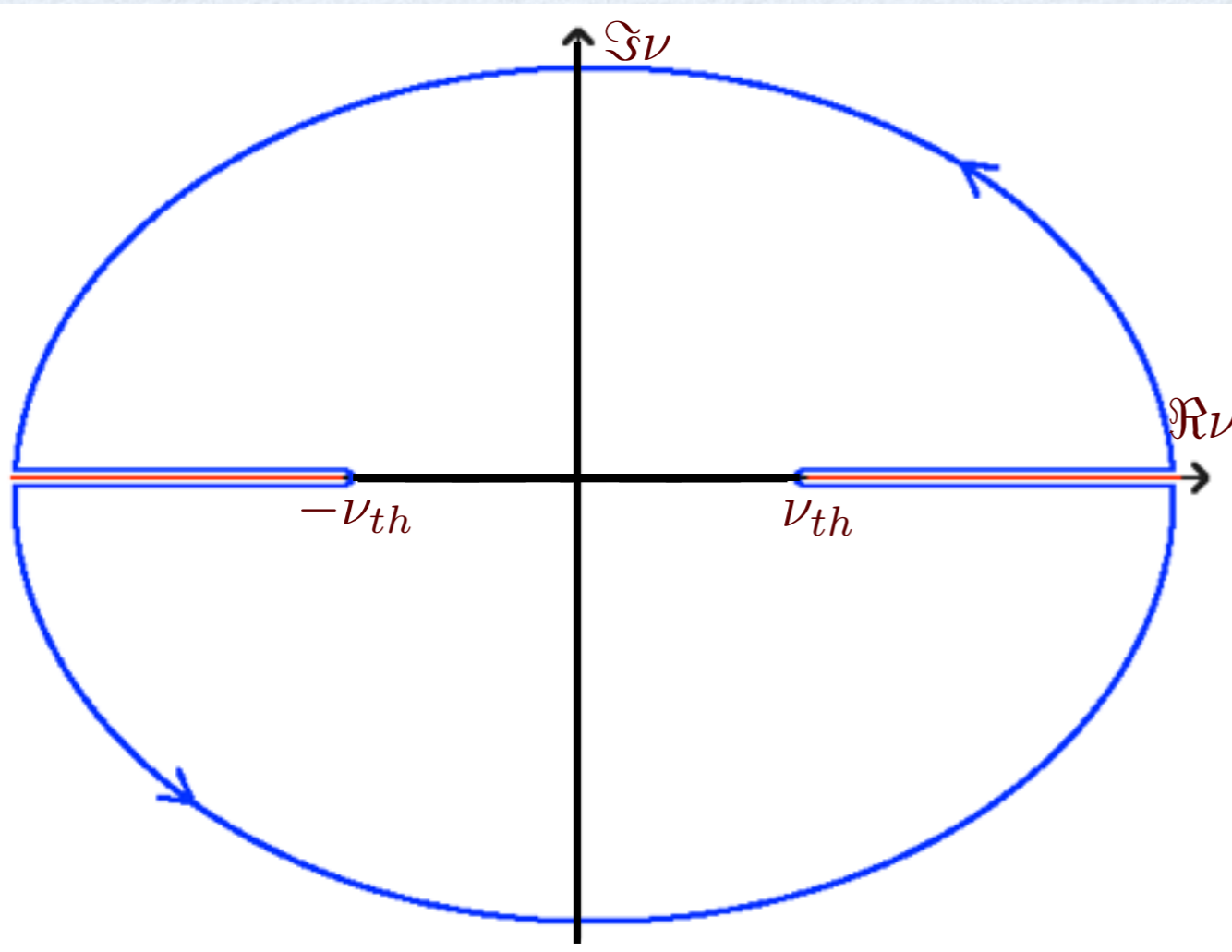
This work

$$Q^2 = 0.1 \text{ GeV}^2$$

$$\nu_{ph} = 0.03 \text{ GeV}^2$$



Fixed-t dispersion relation. Real part



Under crossing $\nu \rightarrow -\nu$
two-photon exchange amplitudes

$\mathcal{G}_M(\nu, t), \mathcal{F}_2(\nu, t) \Leftrightarrow \mathcal{G}_1(\nu, t), \mathcal{G}_2(\nu, t)$ are odd
 $\mathcal{F}_3(\nu, t)$ is even

$$\mathcal{G}_1 = \mathcal{G}_M + \frac{\nu}{M^2} \mathcal{F}_3$$

$$\mathcal{G}_2 = \mathcal{G}_E + \frac{\nu}{M^2} \mathcal{F}_3$$



good
HE behavior

Real part can be reconstructed with DRs

$$\Re \mathcal{G}^{odd}(\nu, t) = \frac{2\nu}{\pi} \mathcal{P} \int_{\nu_{th}}^{\infty} \frac{\Im \mathcal{G}^{odd}(\nu' + i0, t)}{\nu'^2 - \nu^2} d\nu'$$

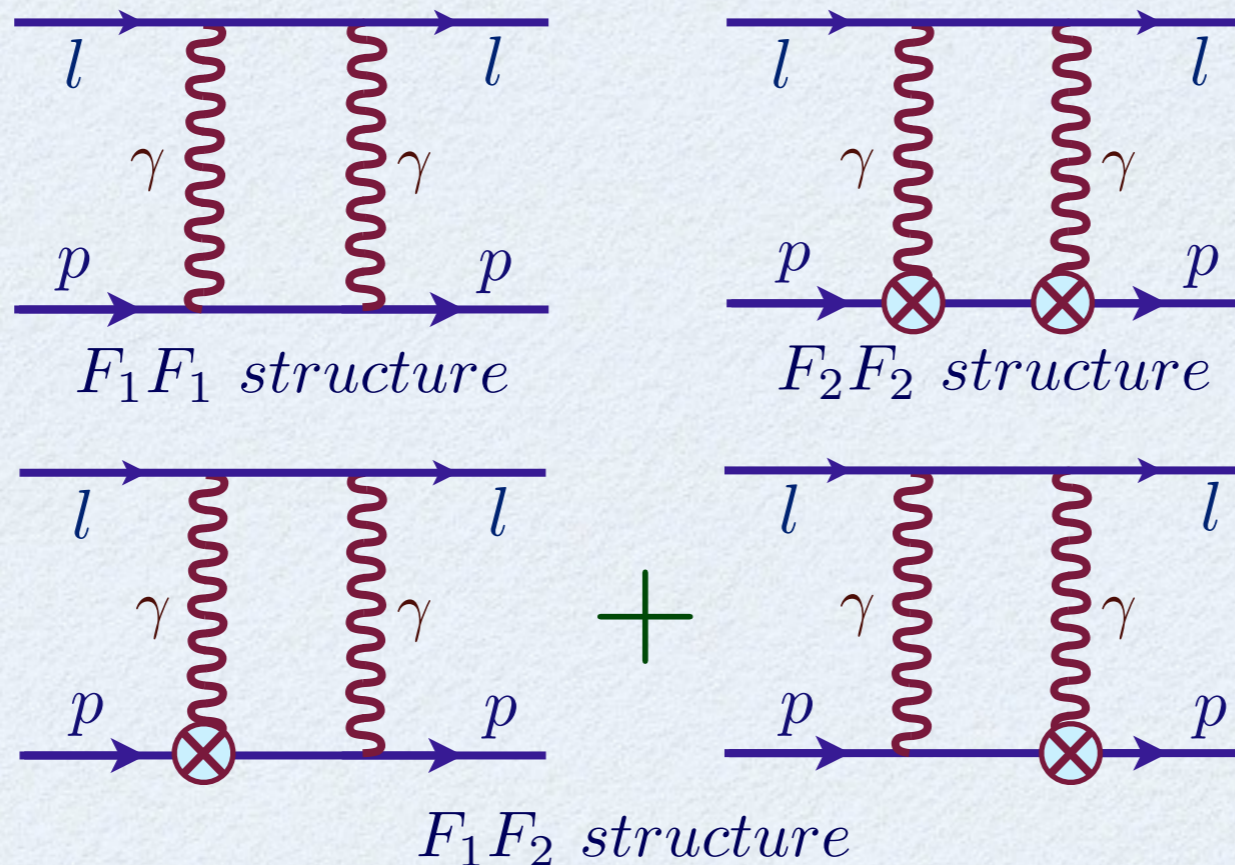
$$\Re \mathcal{G}^{even}(\nu, t) = \frac{2}{\pi} \mathcal{P} \int_{\nu_{th}}^{\infty} \nu' \frac{\Im \mathcal{G}^{even}(\nu' + i0, t)}{\nu'^2 - \nu^2} d\nu'$$

Hadronic model

The one-photon exchange **on-shell** vertex

$$\Gamma^\mu(Q^2) = \gamma^\mu F_1(Q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2M} F_2(Q^2)$$

P. G. Blunden, W. Melnitchouk, and J. A. Tjon (2003)



IR divergencies
are subtracted

L.C. Maximon and J. A. Tjon (2000)

Point-like couplings



$F_1 = 1$

$F_2 = \mu_p - 1$

Dipole FFs for G_M, G_E



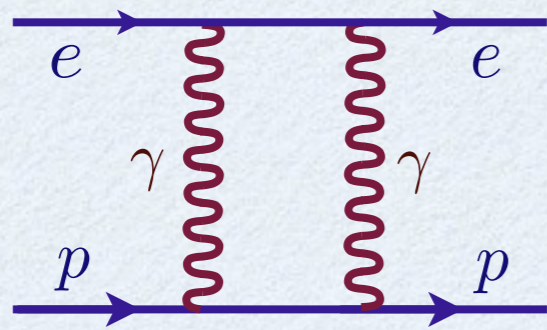
$G_E = F_1 - \tau F_2$

$G_M = F_1 + F_2$

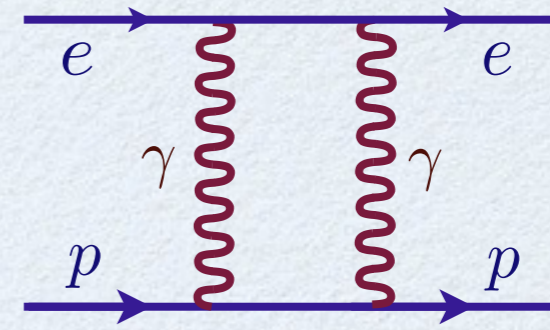
$$\left(\tau = \frac{Q^2}{4M^2}\right)$$

Small momentum transfer limit (e-p)

Feshbach correction - scattering correction in Dirac theory (HE)



$F_1 F_1$ structure



Coulomb photons

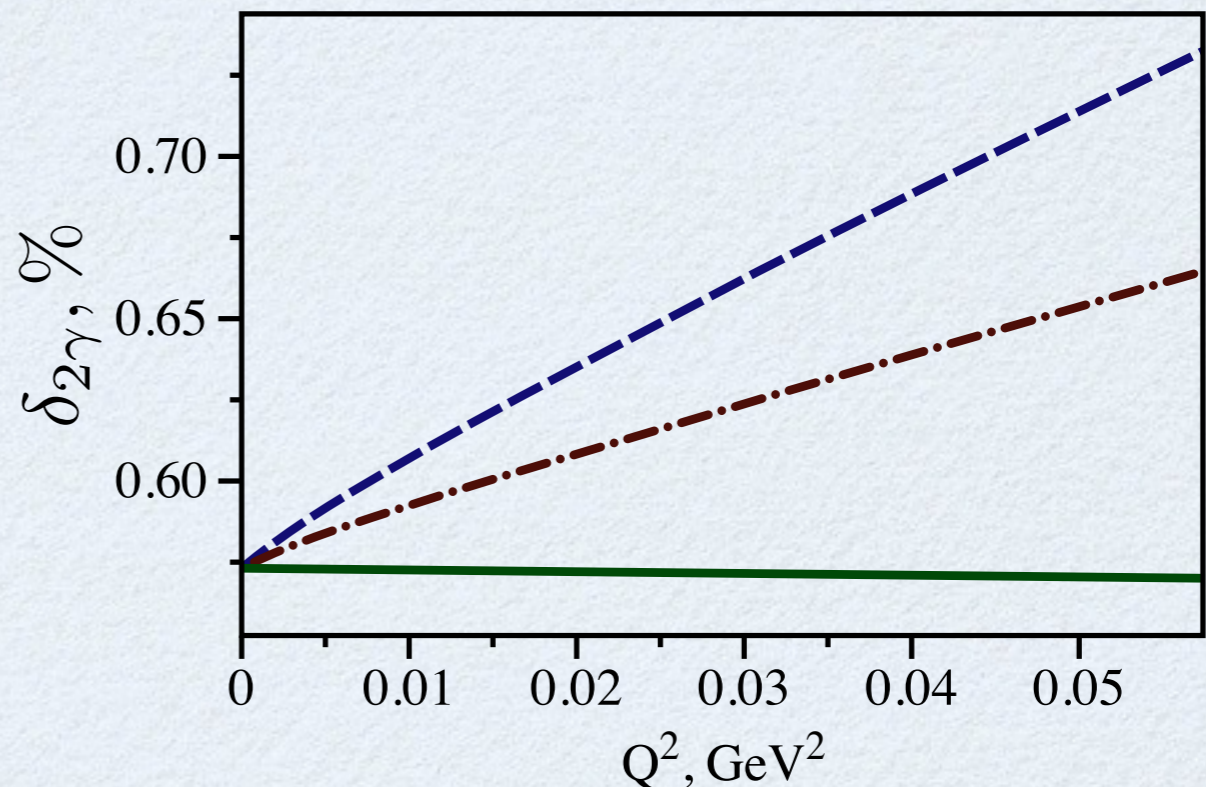
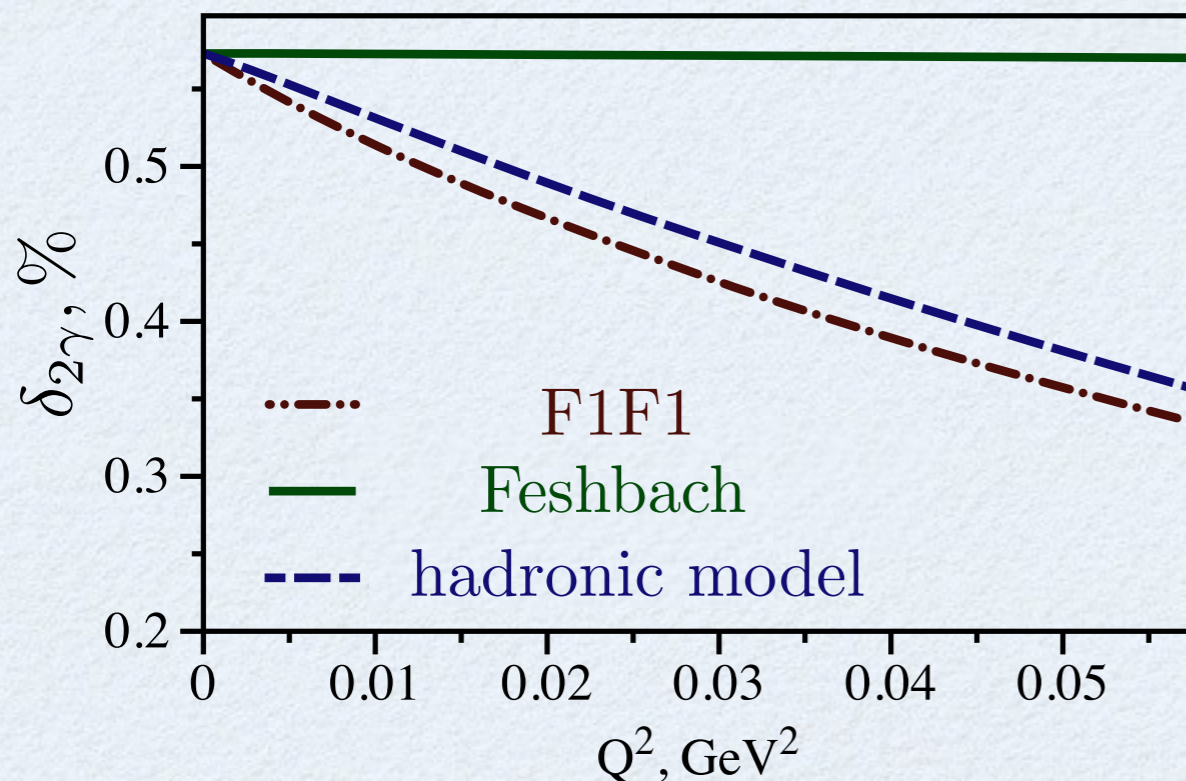
$$\delta_F = \pi\alpha \frac{\sqrt{1-\epsilon}}{\sqrt{1-\epsilon} + \sqrt{1+\epsilon}} = \pi\alpha \frac{\sin(\frac{\theta}{2}) - \sin(\frac{\theta}{2})^2}{\cos(\frac{\theta}{2})^2}$$

W. A. McKinley and H. Feshbach (1948)

dipole FFs

$\epsilon = 0.8$

point-like coupling

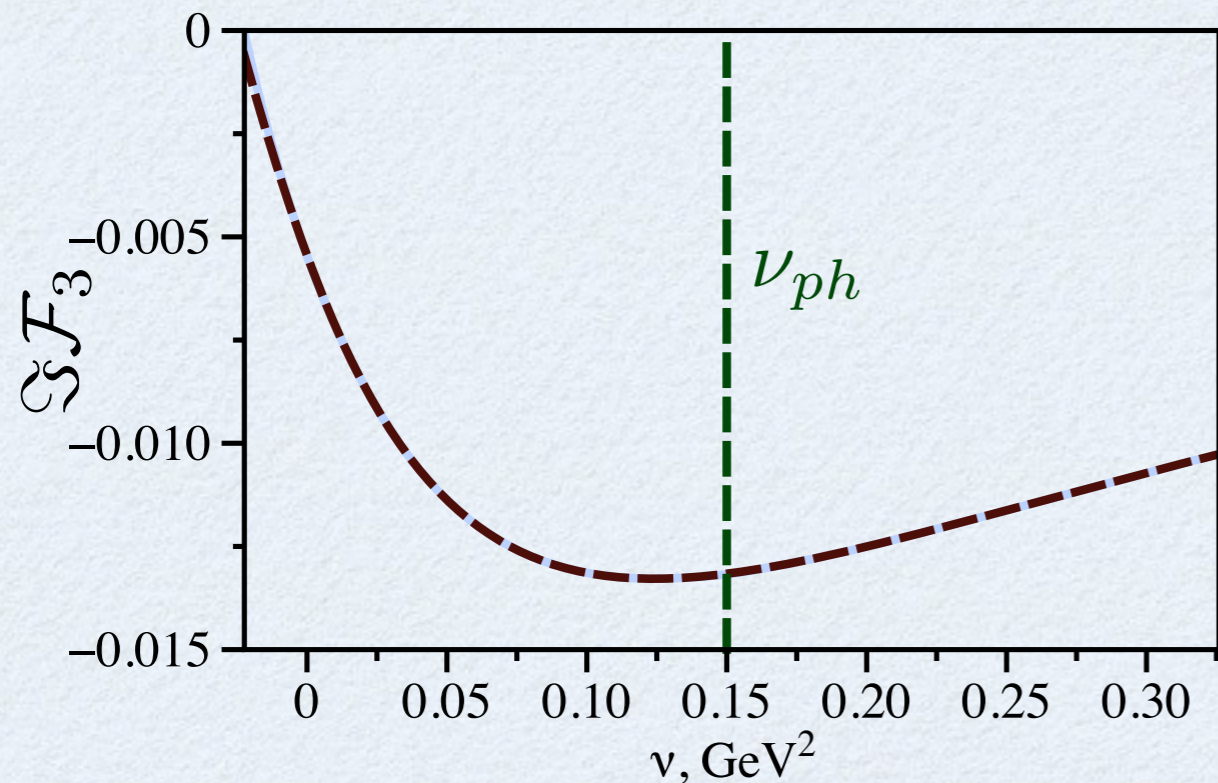
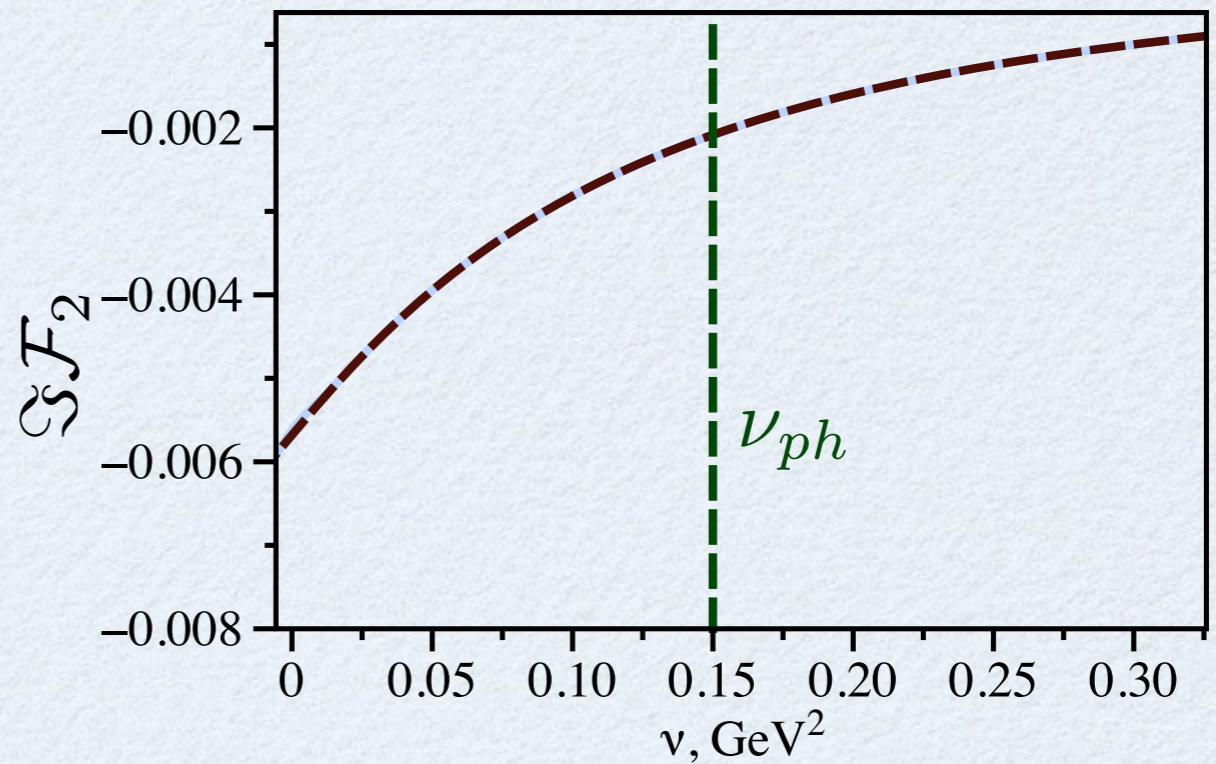
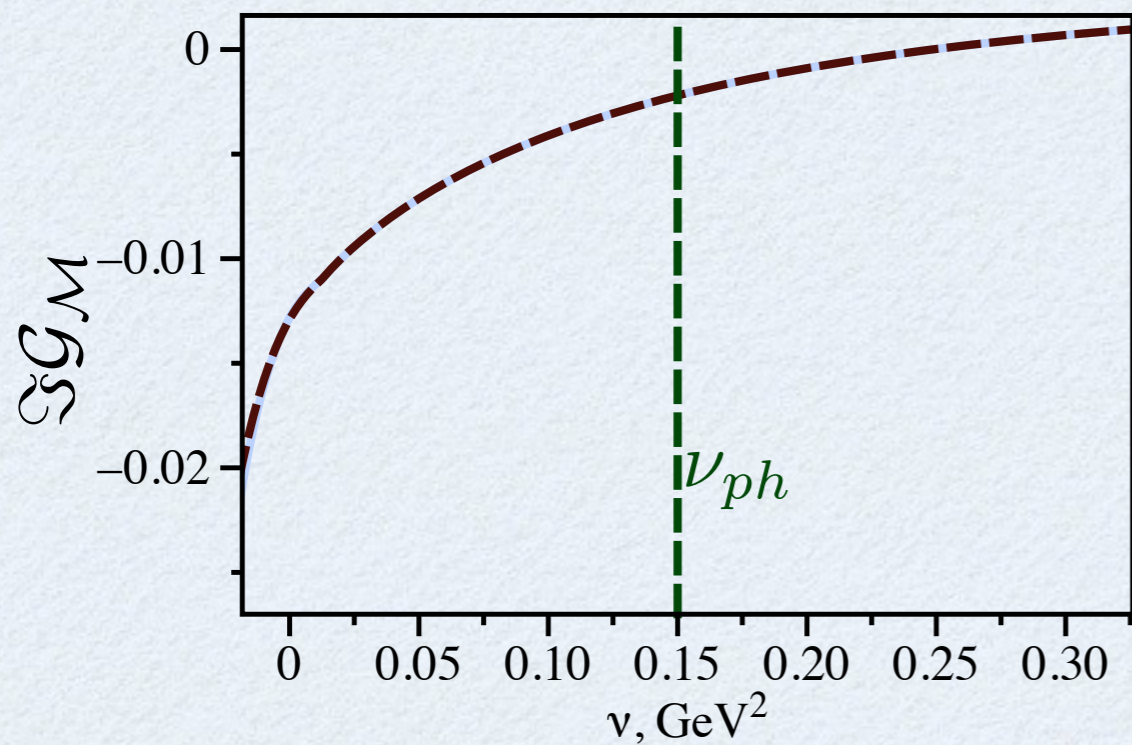


Hadronic model vs. dispersion relations

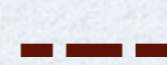
- Imaginary parts

Amplitudes imaginary parts

Dipole form of G_M, G_E



$F_1 F_1$ structure



hadronic model



unitarity relations

$$Q^2 = 0.1 \text{ GeV}^2$$

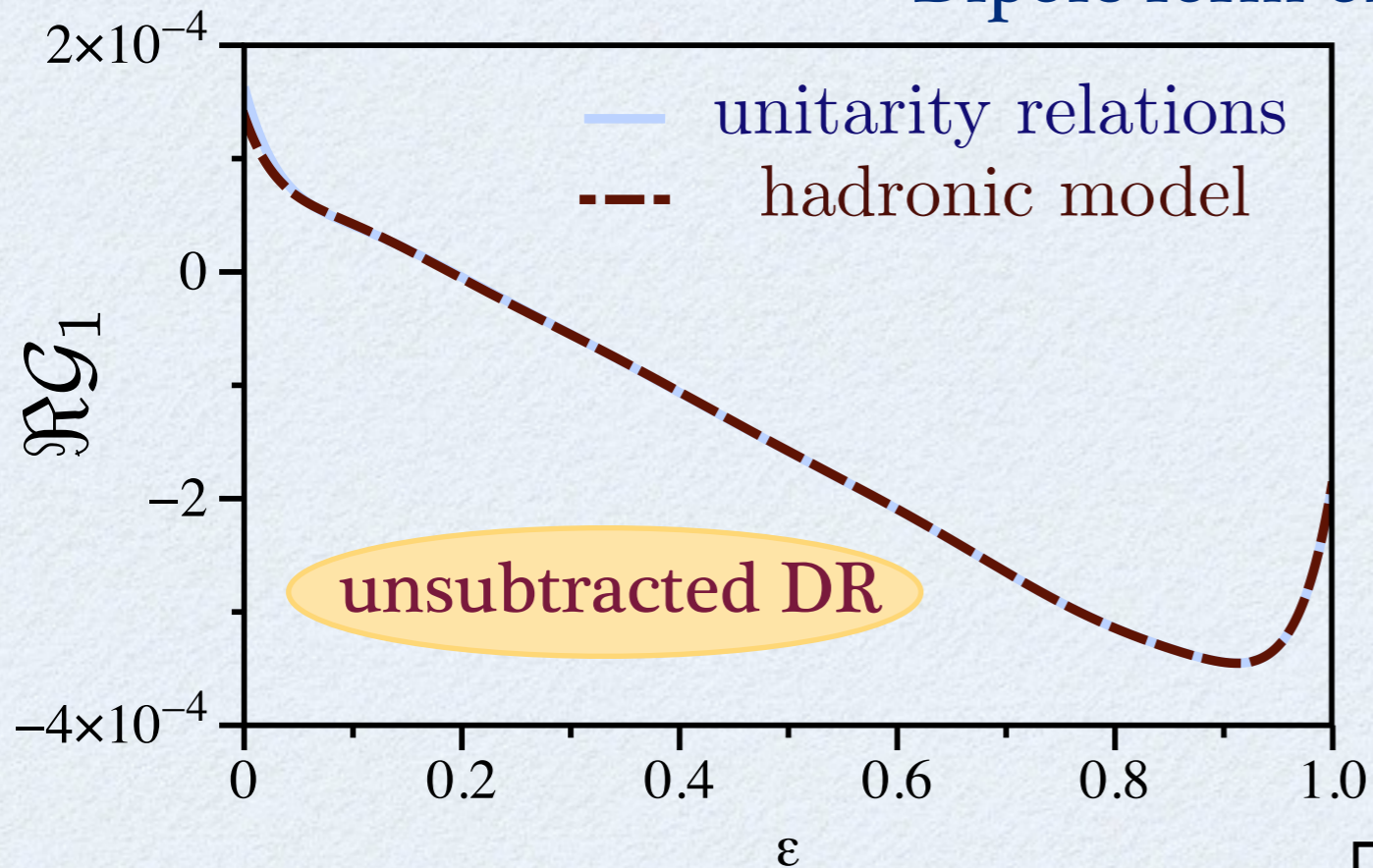
$$\nu_{ph} = 0.15 \text{ GeV}^2$$

Hadronic model vs. dispersion relations

- Imaginary parts are the same
- Real parts

Amplitudes real parts

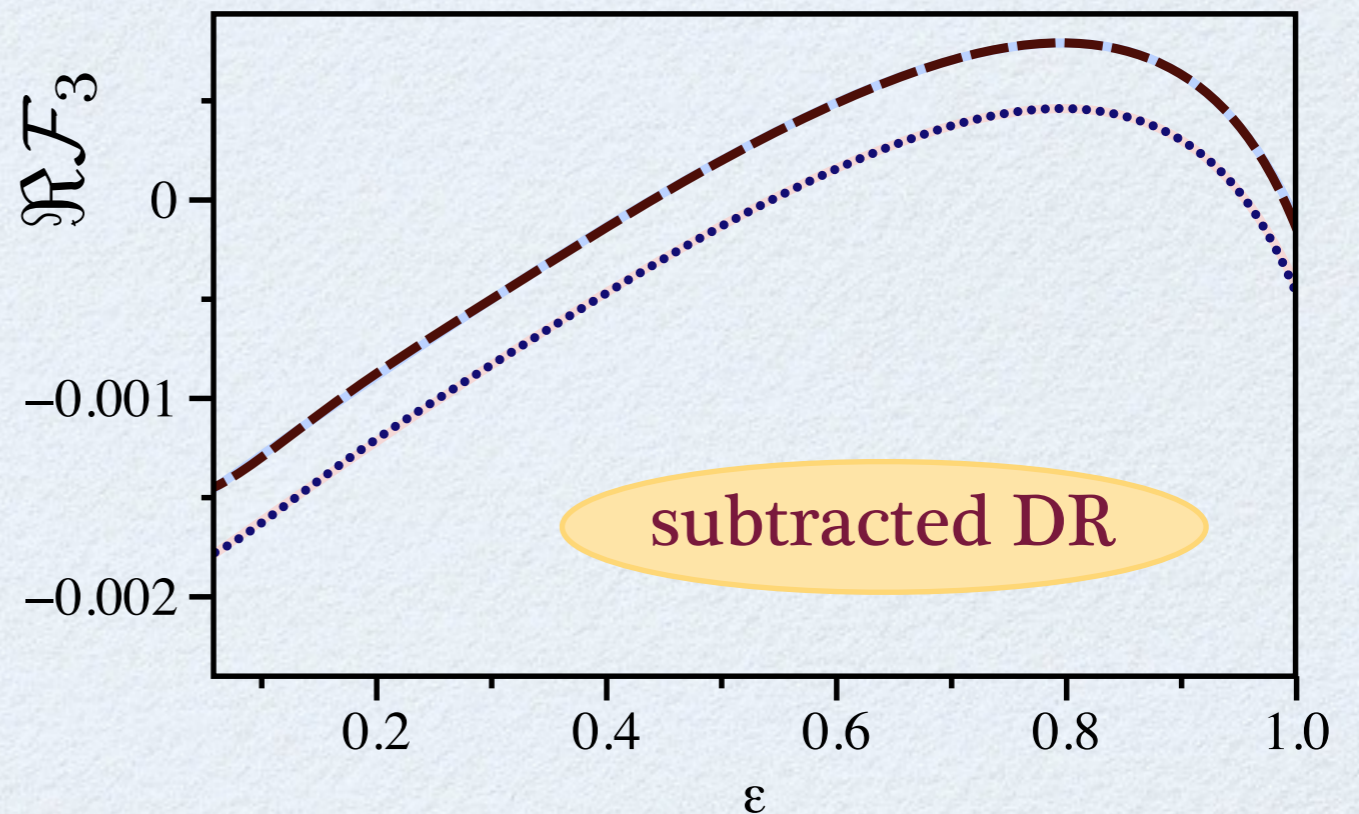
Dipole form of G_M, G_E



$F_2 F_2$ structure

$$Q^2 = 0.1 \text{ GeV}^2$$

- exact
- - - DR($\nu_0 = 2 \text{ GeV}^2$)
- exact
- DR($\nu_0 = 1 \text{ GeV}^2$)



Hadronic model vs. dispersion relations

- Imaginary parts are the same
- Real parts are the same for

all F_1F_1 amplitudes

all F_1F_2 amplitudes

F_2F_2 amplitudes

$$\mathcal{G}_M \quad \mathcal{F}_2 \quad \mathcal{F}_3$$

$$\mathcal{G}_M \quad \mathcal{F}_2 \quad \mathcal{F}_3$$

$$\mathcal{F}_2 \quad \mathcal{G}_M + \frac{\nu}{M^2} \mathcal{F}_3$$

Fixed-t subtracted dispersion relation works

$$F_2F_2 \text{ amplitudes} \quad \mathcal{G}_M \quad \mathcal{F}_3$$

Hadronic model vs. dispersion relations

- Imaginary parts are the same
- Real parts are the same for

all F_1F_1 amplitudes

all F_1F_2 amplitudes

F_2F_2 amplitudes

$$\mathcal{G}_M \quad \mathcal{F}_2 \quad \mathcal{F}_3$$

$$\mathcal{G}_M \quad \mathcal{F}_2 \quad \mathcal{F}_3$$

$$\mathcal{F}_2 \quad \mathcal{G}_M + \frac{\nu}{M^2} \mathcal{F}_3$$

Fixed-t subtracted dispersion relation works

F_2F_2 amplitudes $\mathcal{G}_M \quad \mathcal{F}_3$

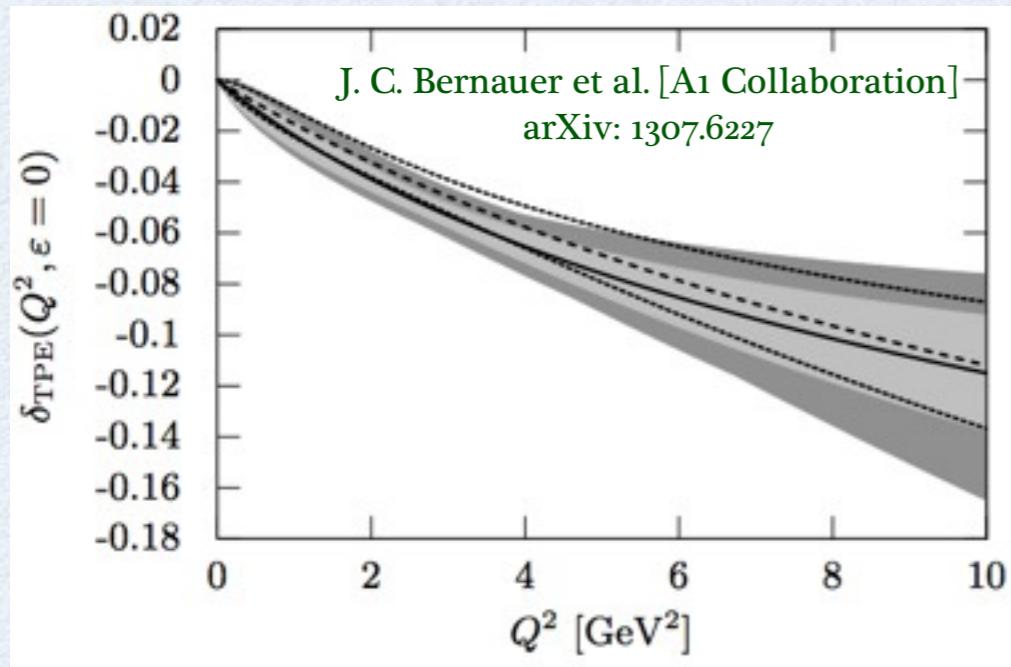
- Calculation based on DR

This work

- for amplitudes $\mathcal{G}_1, \mathcal{G}_2$ unsubtracted DR can be used

- for amplitude \mathcal{F}_3 subtracted DR should be used

2 γ in e-p elastic scattering



$$\delta_{2\gamma} = \delta_F + \delta_{TPE}$$

$$\delta_{TPE} = -(1 - \epsilon) a \ln(bQ^2 + 1)$$

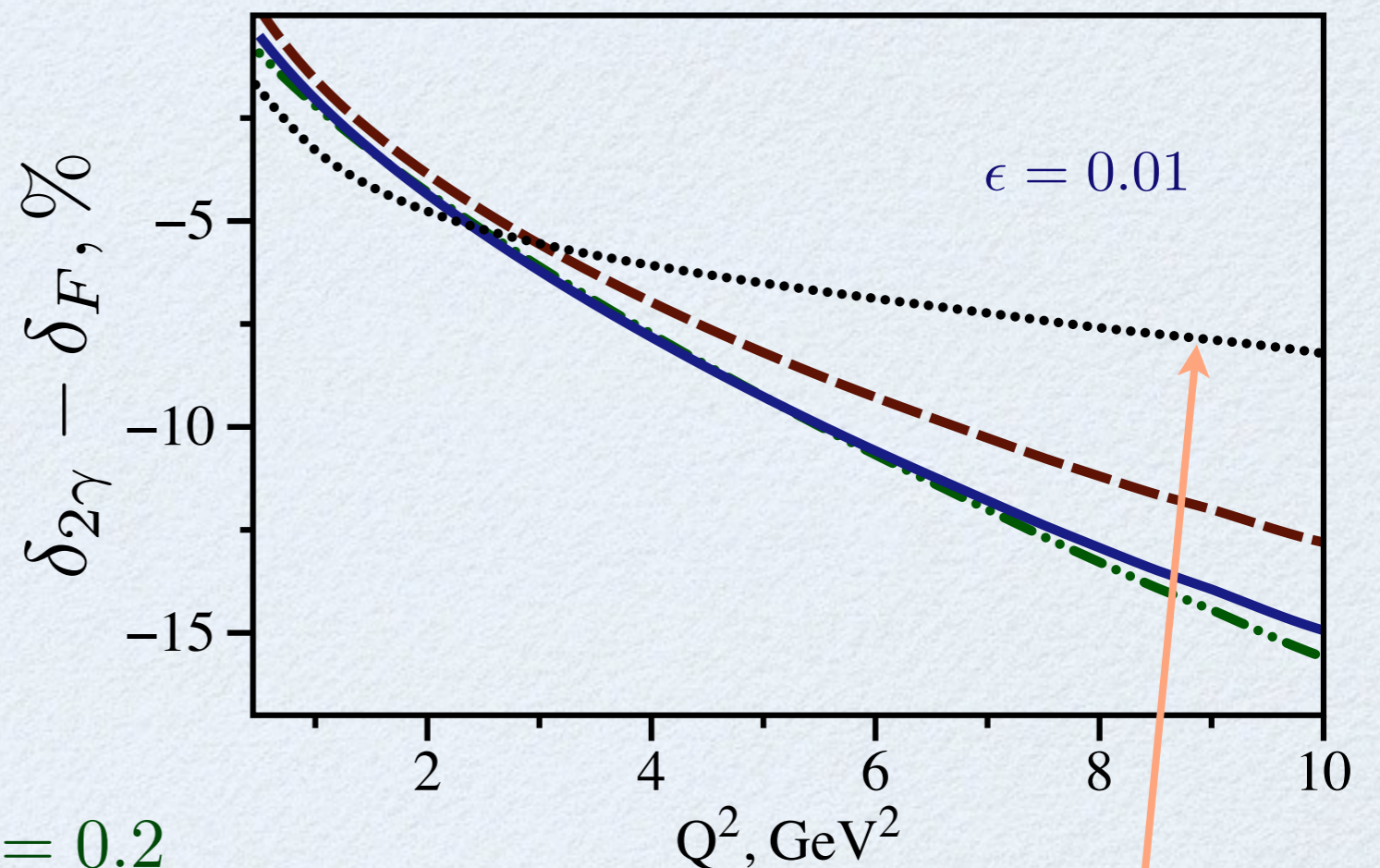
subtraction points

- $\epsilon_0 = 0.2$
- $\epsilon_0 = 0.5$
- $\epsilon_0 = 0.8$

subtracted DR

subtraction point $\Re \mathcal{F}_3^{F_2 F_2}(\nu_0, Q^2)$

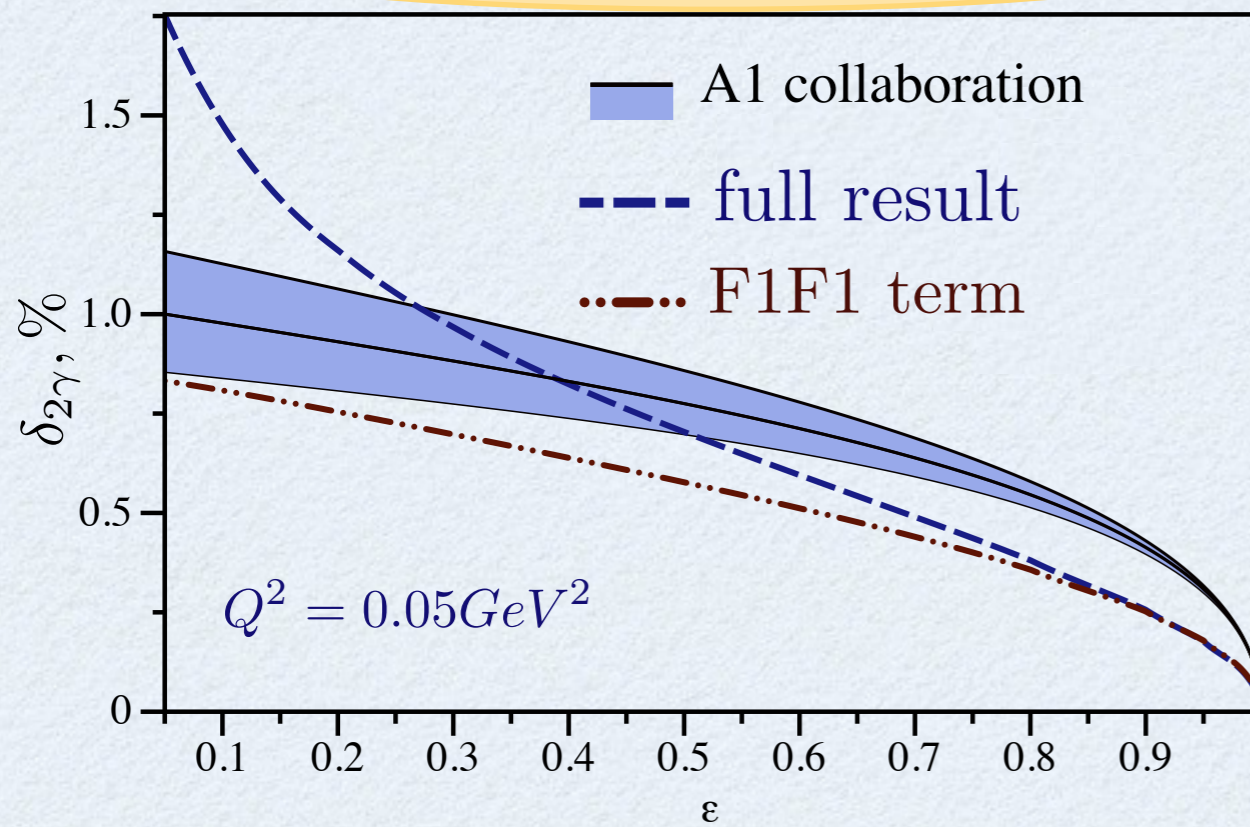
fixed from $\delta_{2\gamma}(\nu_0, Q^2)$ data



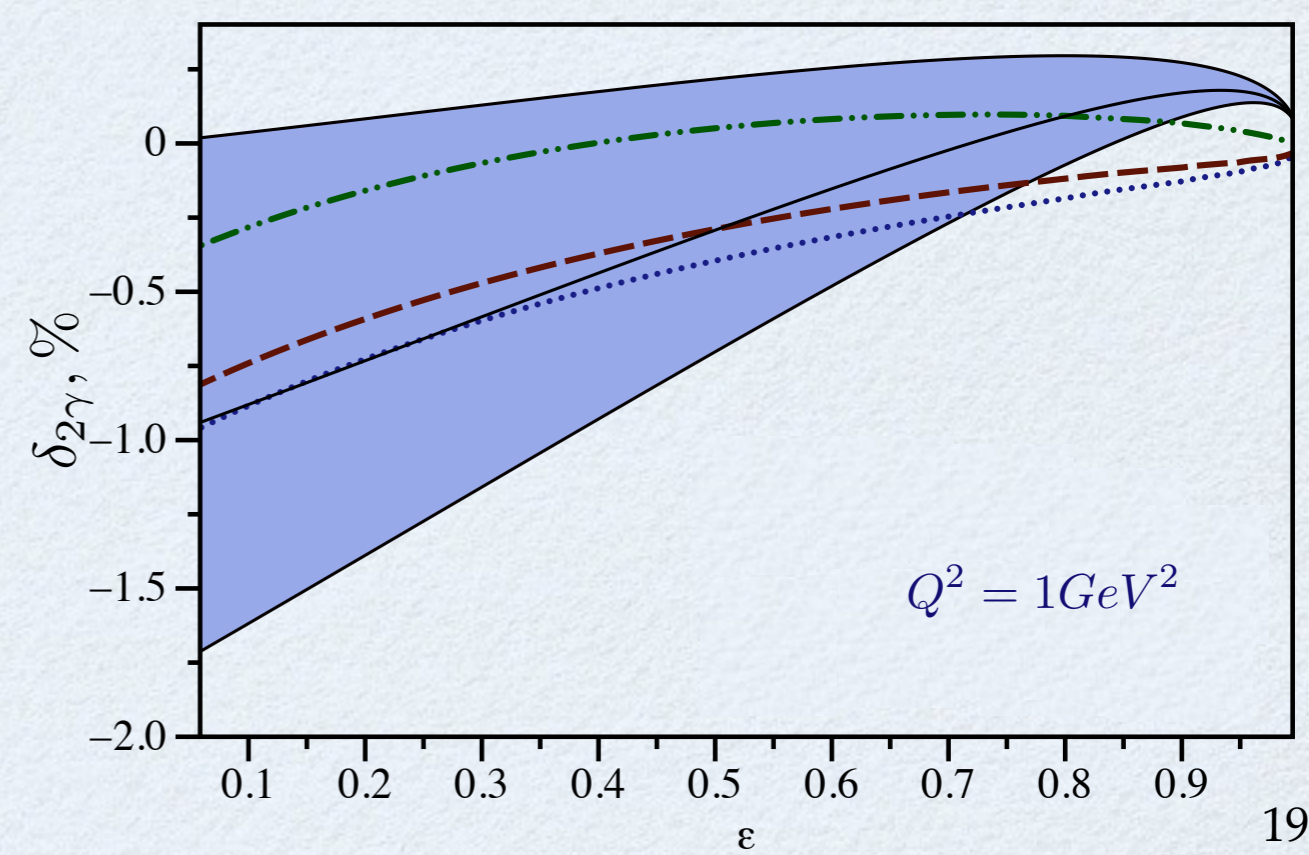
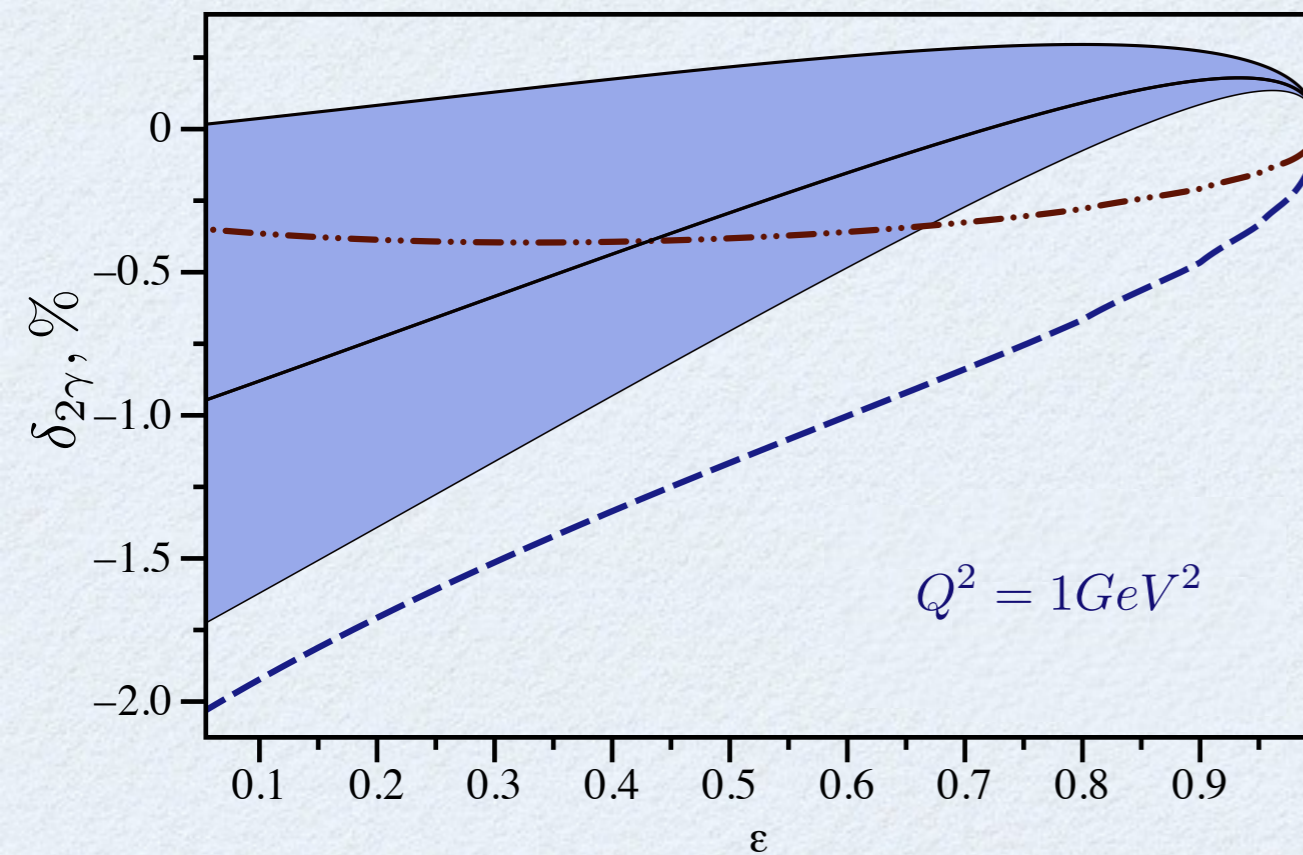
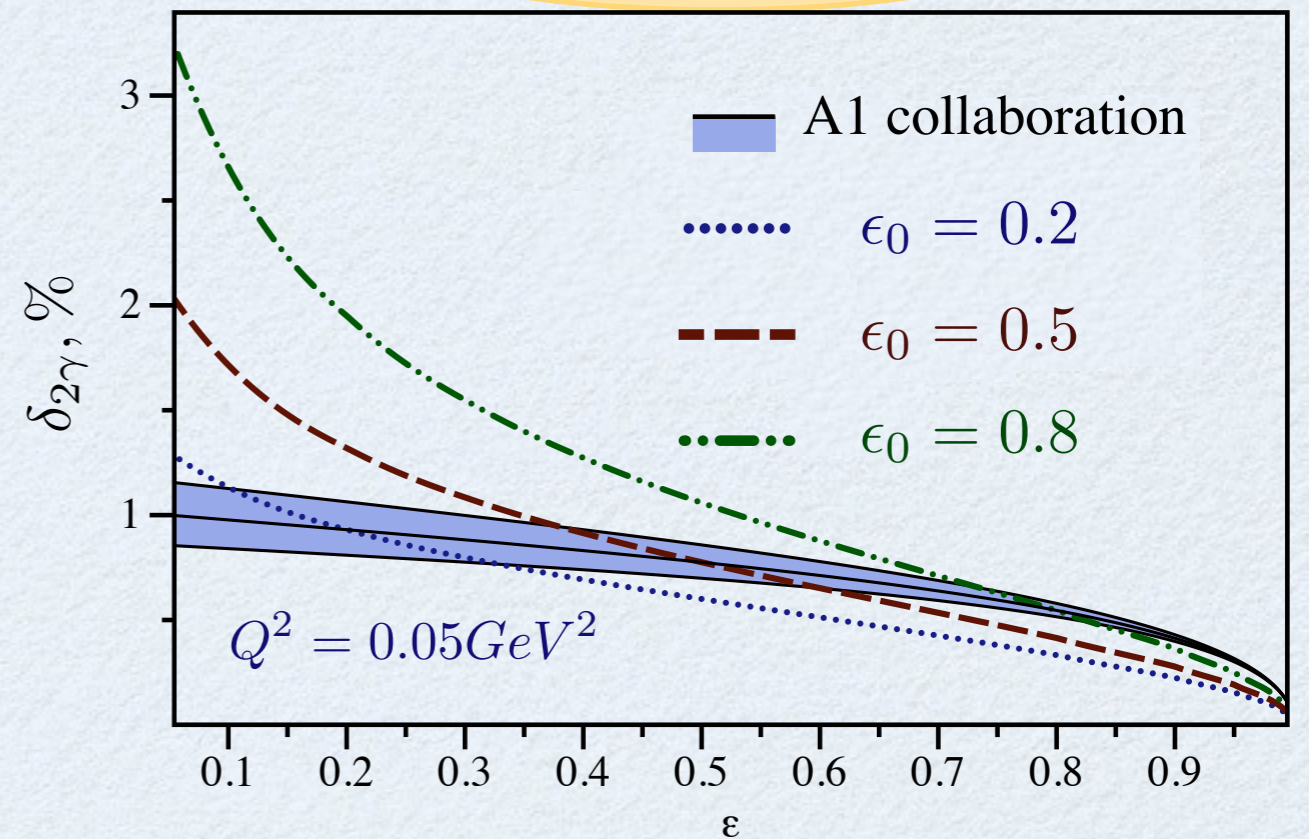
box diagram model calculation

2γ in e^-p elastic scattering

box diagram model calculation



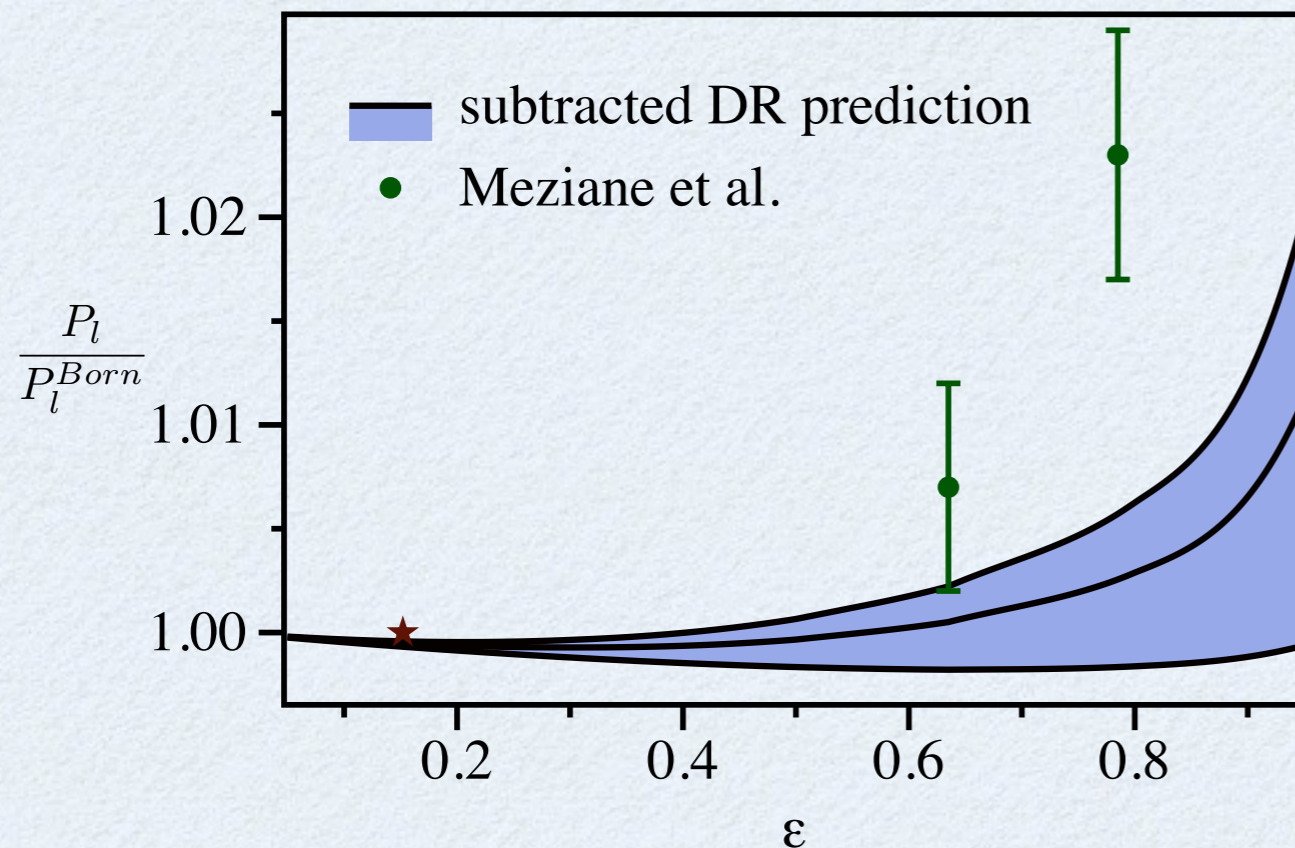
subtracted DR



Polarization transfer observables (e-p)

Experimental points at $Q^2 \sim 2.5 \text{ GeV}^2$

M. Meziane et al. [GEp2gamma Collaboration] (2011)



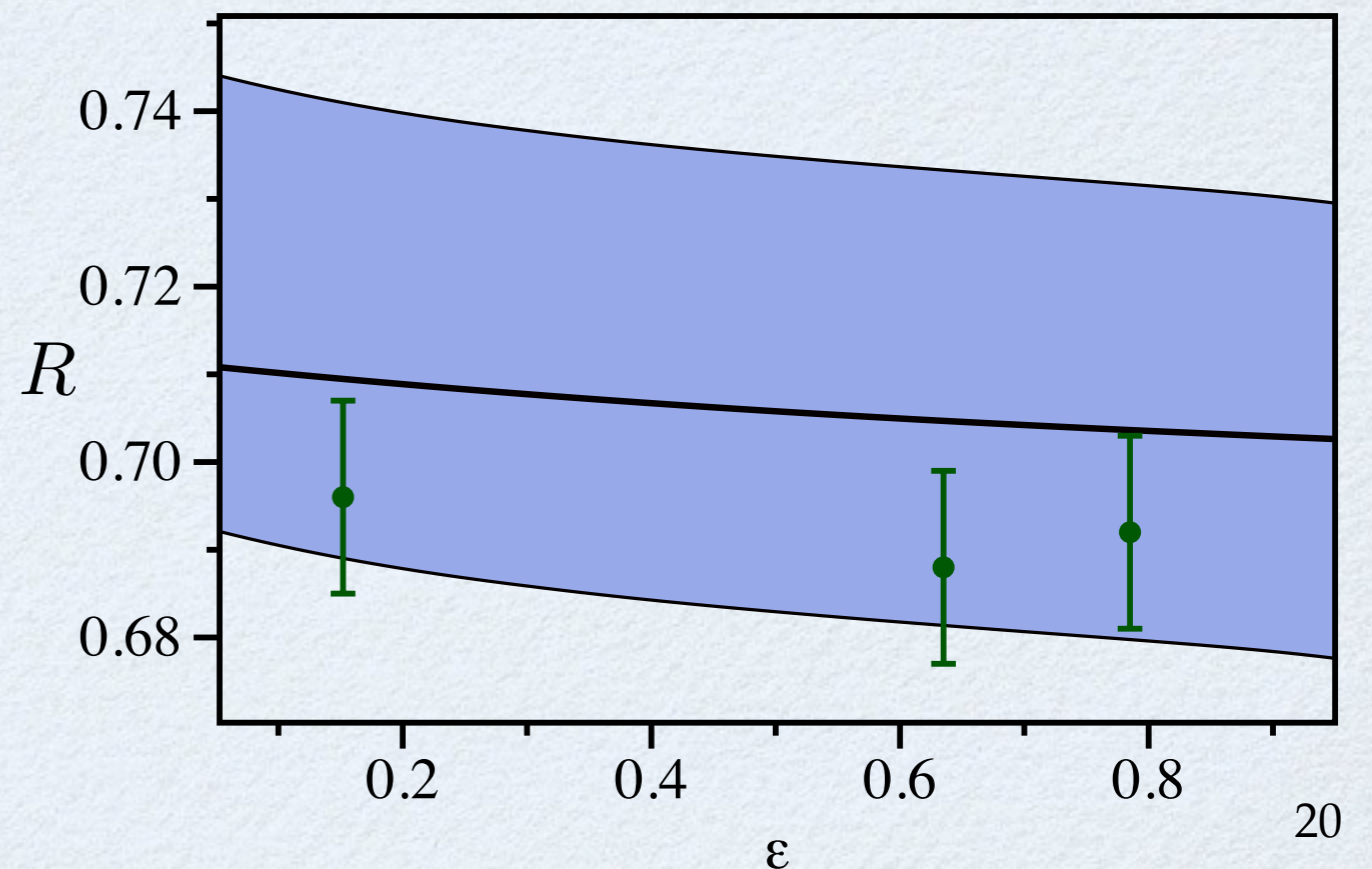
Longitudinal polarization transfer

$$\frac{P_l}{P_l^{Born}} - 1 \sim \Re \mathcal{G}_M, \Re \mathcal{F}_2, \Re \mathcal{F}_3$$

Longitudinal to transverse polarization transfer

$$R = -\mu_p \sqrt{\frac{(1+\epsilon)\tau}{2\epsilon}} \frac{P_t}{P_l}$$

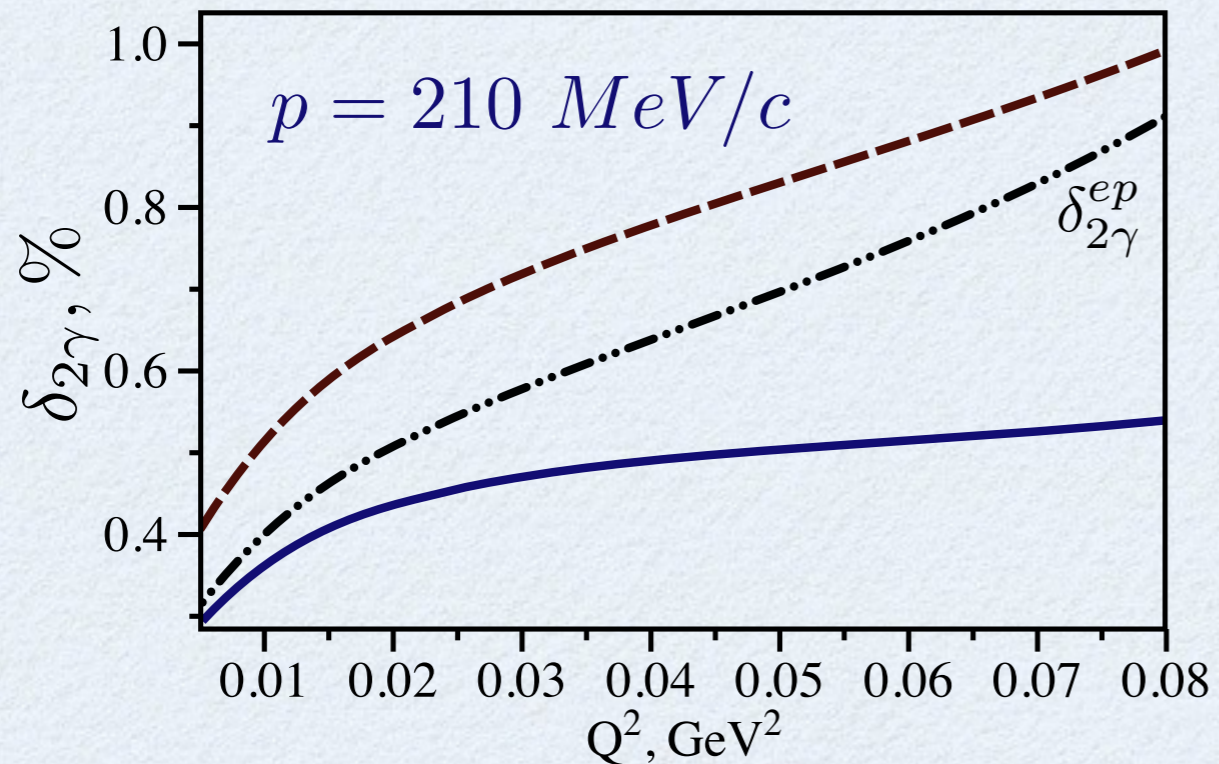
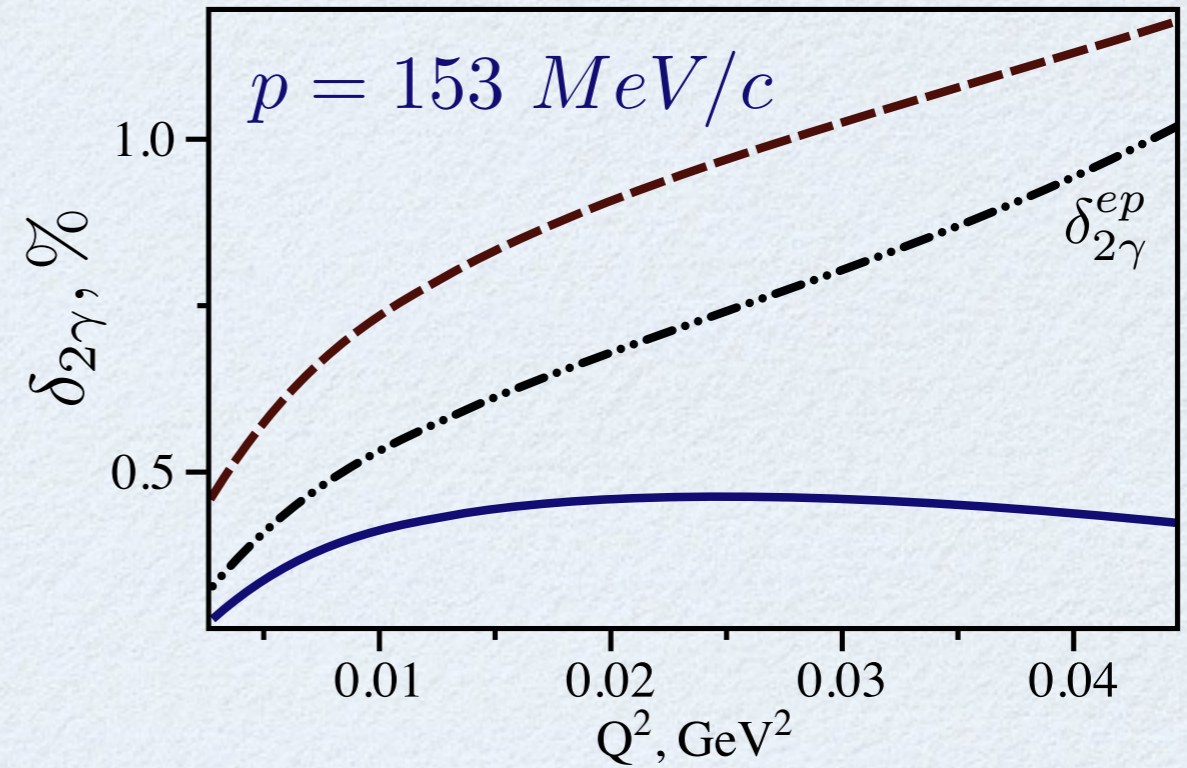
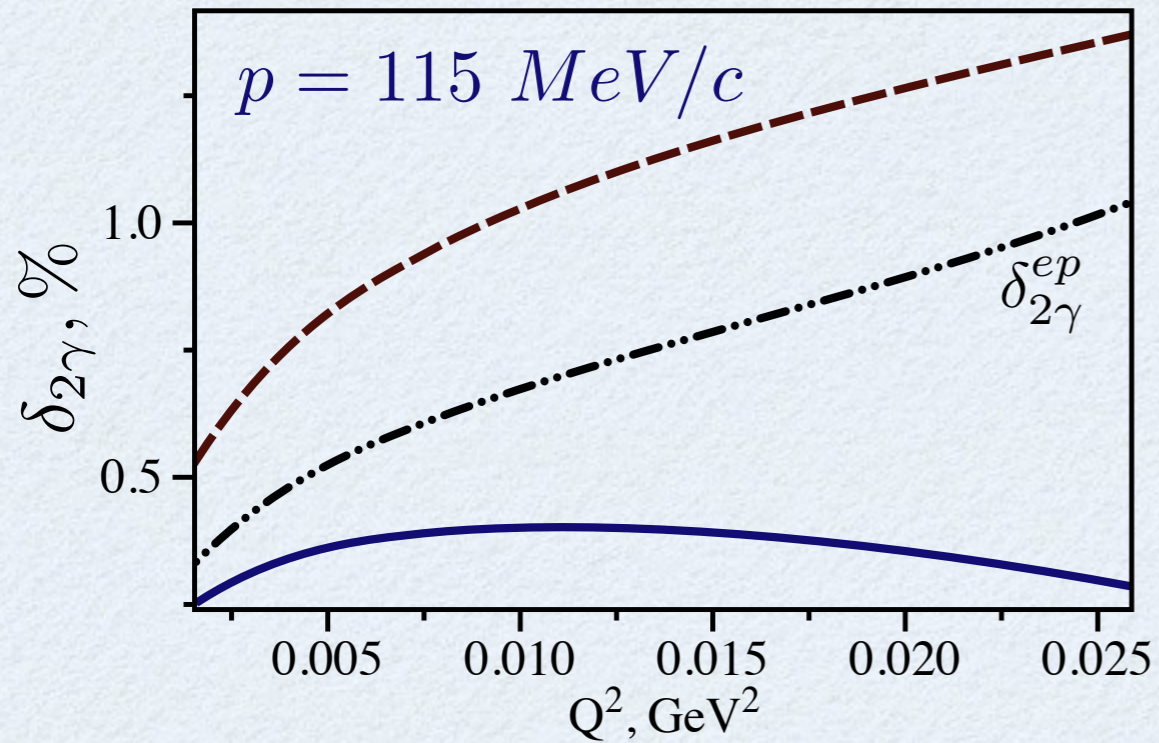
$$R - \mu_p \frac{G_E}{G_M} \sim \Re \mathcal{G}_M, \Re \mathcal{F}_2, \Re \mathcal{F}_3$$



μ^-p experiments estimates

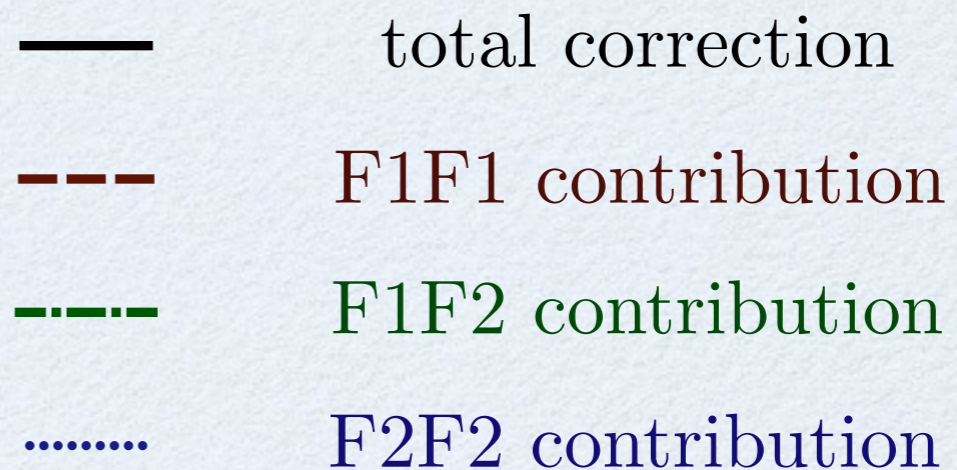
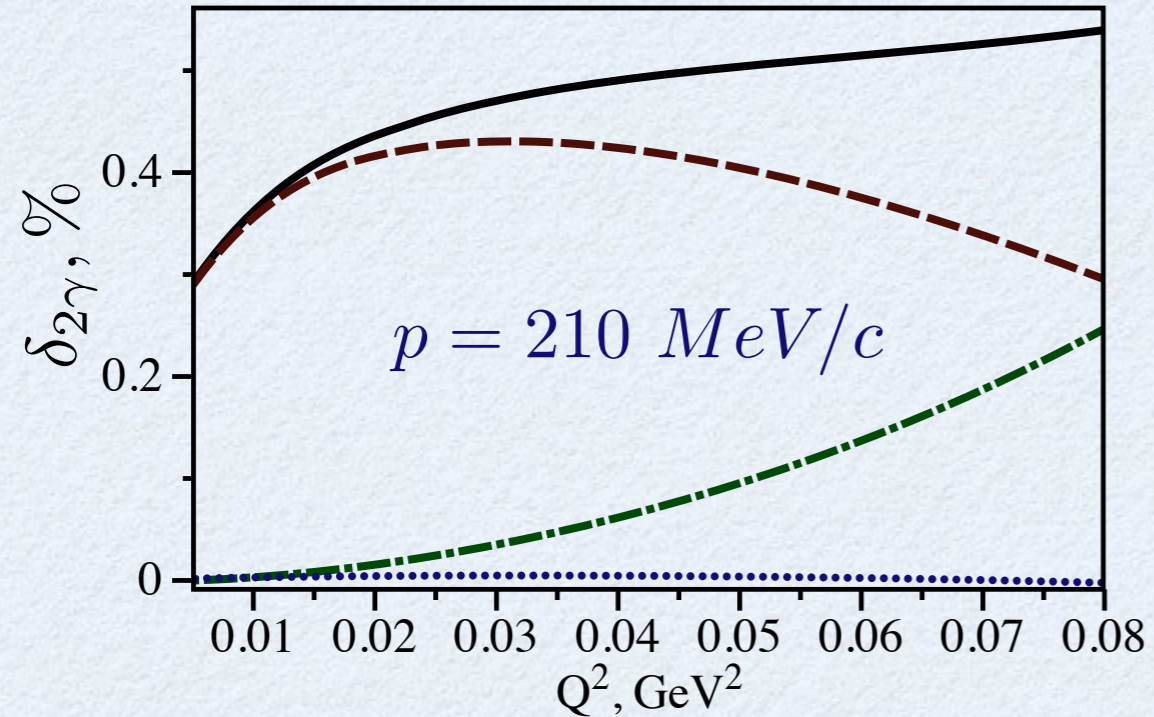
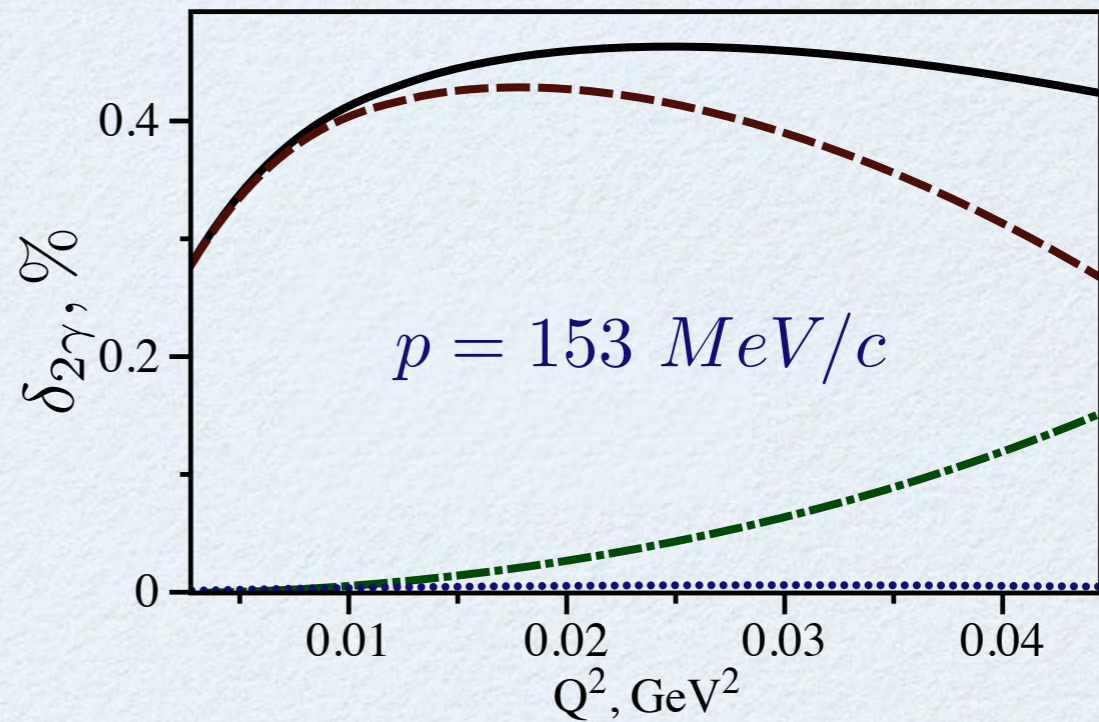
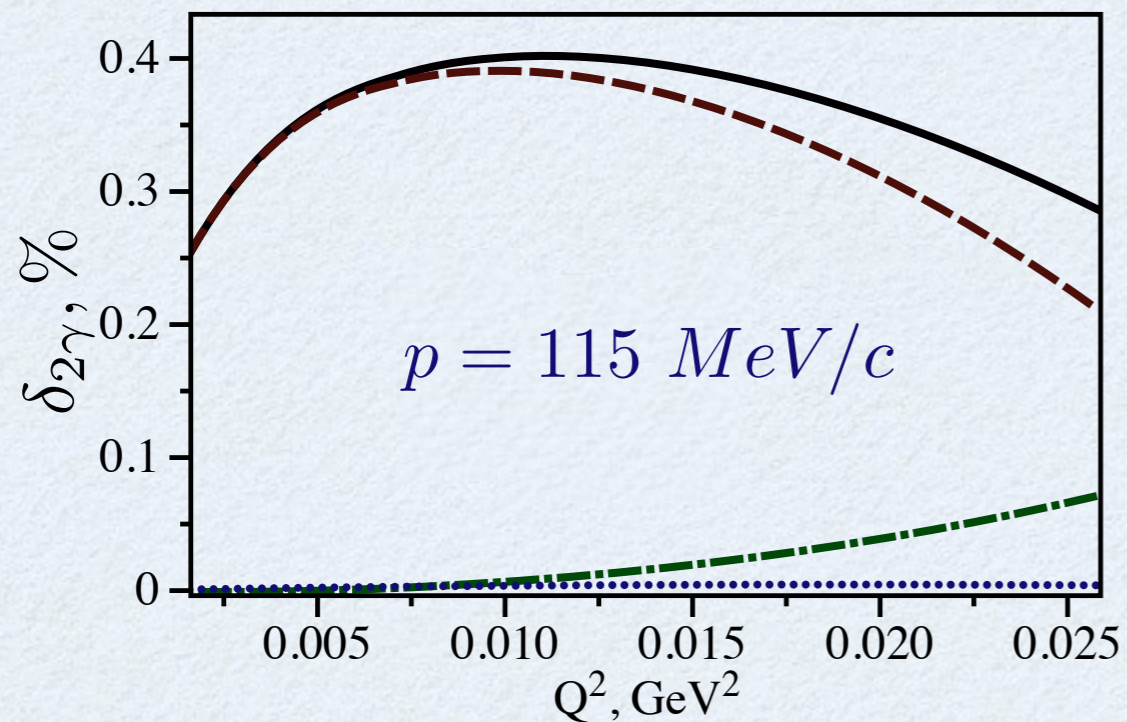
TPE correction

$$\delta_{2\gamma} \sim \Re\mathcal{G}_M, \Re\mathcal{F}_2, \Re\mathcal{F}_3, \Re\mathcal{F}_4, \Re\mathcal{F}_5$$



- μ^-p scattering
- - - $\mu^-p, \mathcal{F}_4 = \mathcal{F}_5 = 0$
- · - · - e^-p scattering

μ^-p experiments estimates



F1F1 contribution dominates

DRs approach requires two subtraction points

Conclusions and outlook

- Dispersion relations framework was developed
 - DR for ep scattering require 1 exp. point as input
 - DR for μp scattering require 2 exp. points as input
- DR checked vs. hadronic model calculation (ep):
F1F1, F1F2: agreement F2F2 : on-shell model violates DR
- Theoretical estimates for elastic (ep and μp) cross section and polarization transfer observables were made
- Next step: inclusion of inelastic intermediate states (πN)

Thanks for your attention !!!