

Dynamically-coupled partial-waves in $\rho\pi$ scattering and meson-meson operators

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hadspec

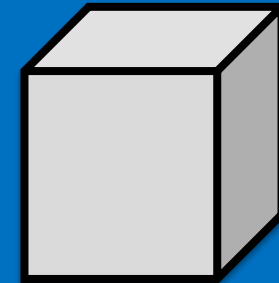
Outline

- Introduction
- Multi-hadron operators
- Isospin-2 $\rho\pi$ scattering with dynamically-coupled partial waves
- Summary

Excited lattice QCD spectroscopy

Finite-volume energy eigenstates from:

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$$




Large bases of interpolating operators and variational method (GEVP)

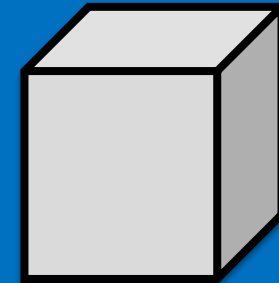
$$C_{ij}(t)v_j^{(n)} = \lambda^{(n)}(t)C_{ij}(t_0)v_j^{(n)}$$

$$\lambda^{(n)}(t) \rightarrow e^{-E_n(t-t_0)} \quad (t \gg t_0)$$

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Lüscher method (and extensions): relate finite-volume energy levels to infinite-volume scattering t -matrix.

Talk by Dave Wilson earlier

Recent review: Briceño, Dudek, Young [arXiv:1706.06223]

Reduced symmetry

Finite cubic lattice, $\vec{p} = \vec{0}$

Broken sym: 3D rotation group \rightarrow octahedral (cubic) group O_h^D

Finite number of *irreps* Λ :

A_1, A_2, T_1, T_2, E

(+ others for
half-integer spin)

Λ^P	J^P
A_1^\pm	$0^\pm, 4^\pm, \dots$
T_1^\pm	$1^\pm, 3^\pm, 4^\pm, \dots$
T_2^\pm	$2^\pm, 3^\pm, 4^\pm, \dots$
E^\pm	$2^\pm, 4^\pm, \dots$
A_2^\pm	$3^\pm, \dots$

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$\vec{p} \neq \vec{0}$

Little group: $LG(\vec{p}) \subset O_h^D \quad \{R \in O_h^D \mid R\vec{p} = \vec{p}\}$

Single-meson operators

[PRL 103, 262001; PR D82, 034508; D84, 074508;
similarly for baryons PR D85, 014507]

Fermion-bilinear 'single-meson' operators:

$$\mathcal{O}^{J^P, M}(\vec{p}, t) = \sum_{\vec{x}} e^{i\vec{p}\cdot\vec{x}} \bar{\psi}(x) \left[\Gamma \times \overleftrightarrow{D} \times \overleftrightarrow{D} \dots \right] \psi(x)$$

Circular basis for D and Γ , couple using SU(2) Clebsch-Gordans $\rightarrow J^P, M$

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'Subduce' operators into lattice irrep (Λ), row (μ):

$$\left[\mathcal{O}_{\Lambda^P, \mu}^{[J]}(\vec{0}) \right]^\dagger = \sum_M \mathcal{S}_{\Lambda, \mu}^{J, M} \left[\mathcal{O}^{J^P, M}(\vec{0}) \right]^\dagger$$

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For $\vec{p} \neq \vec{0}$ 'helicity ops':

$$\left[\mathcal{O}^{J^P, \lambda}(\vec{p}) \right]^\dagger = \sum_M \mathcal{D}_{M\lambda}^{(J)}(R) \left[\mathcal{O}^{J^P, M}(\vec{p}) \right]^\dagger \quad R : (0, 0, |\vec{p}|) \rightarrow \vec{p}$$

Subduce \rightarrow little group irreps

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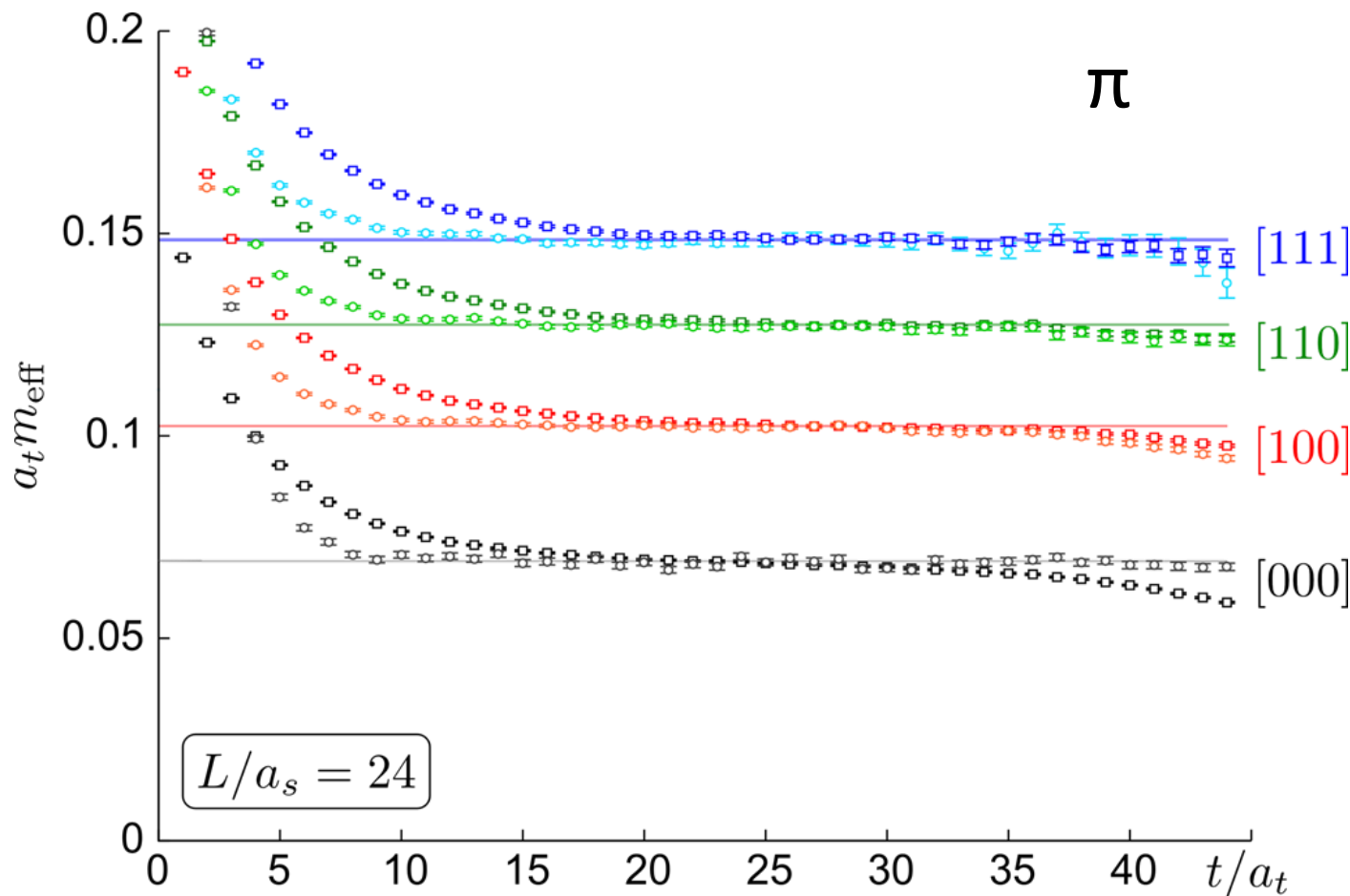
Subduce \rightarrow little group irreps

Many ops in each channel (up to 3 derivs at rest, 2 at non-zero mom.)

Variationally optimised meson ops

For each \vec{p} , Λ : variational analysis with large basis of fermion-bilinear ops
 n 'th eigenvector \rightarrow optimal operator for n 'th state

$$\Omega^{(n)\dagger} \sim \sum_i v_i^{(n)} O_i^\dagger$$



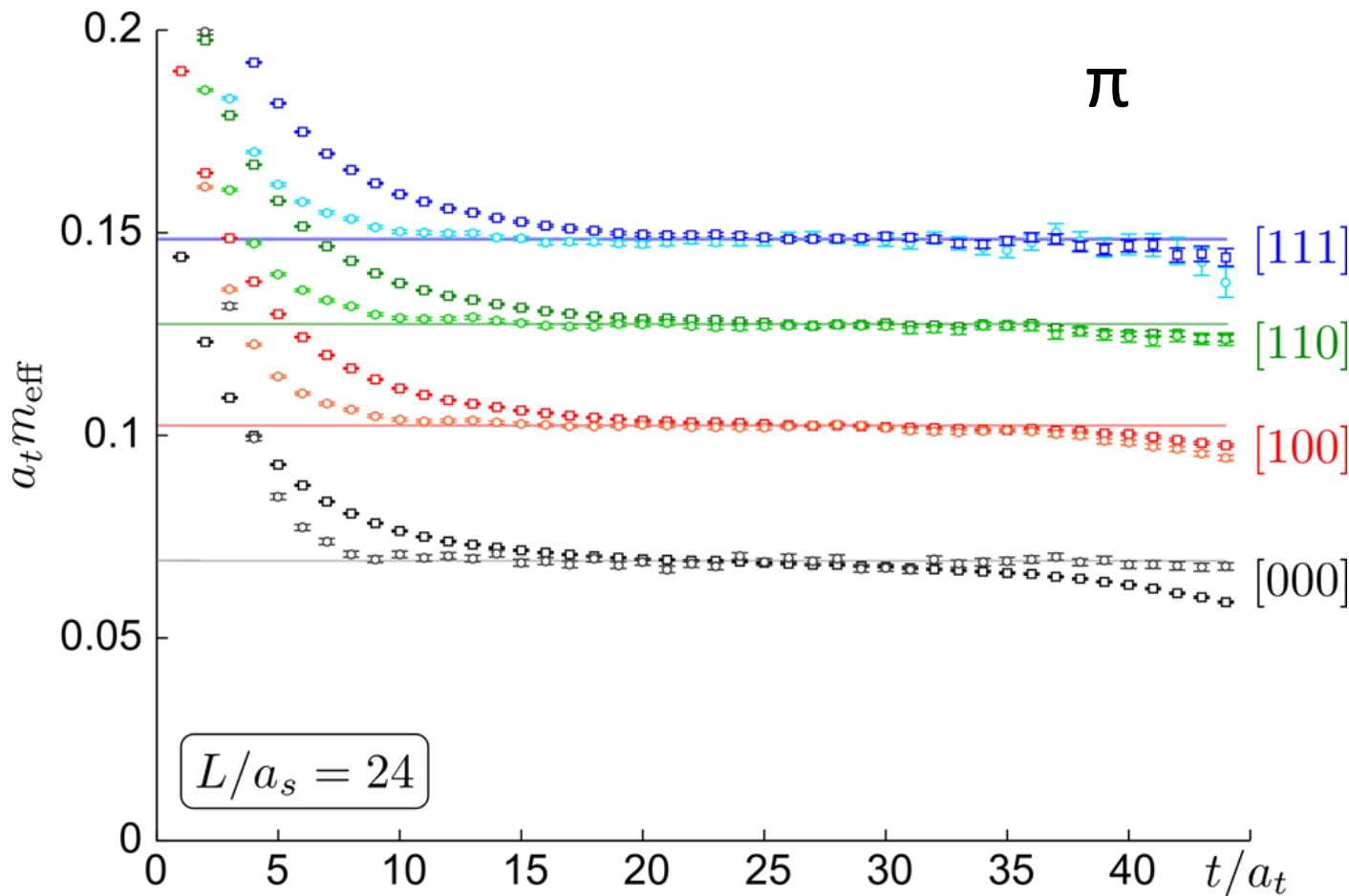
$m_\pi \approx 400 \text{ MeV}$

[PR D86, 34031
(2012)]

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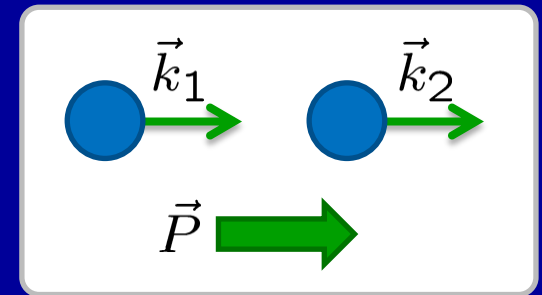


Could have multi-hadron ops

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[PR D86, 34031 (2012)]

Hadron-hadron (or multi-hadron) operators

$$\mathbb{O}_{\Lambda, \mu}^{\dagger}(\vec{P}) = \sum_{\mu_1, \mu_2} \sum_{\vec{k}_1, \vec{k}_2} C(\vec{P}, \Lambda, \mu; \vec{k}_1, \Lambda_1, \mu_1; \vec{k}_2, \Lambda_2, \mu_2) \mathbb{O}_{\Lambda_1 \mu_1}^{\dagger}(\vec{k}_1) \mathbb{O}_{\Lambda_2 \mu_2}^{\dagger}(\vec{k}_2)$$

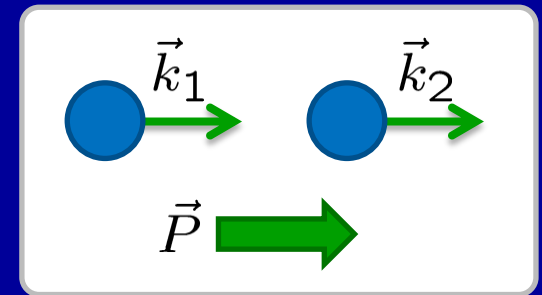


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Could be: simple fermion-bilinears, optimised ops, multi-hadron ops, ...



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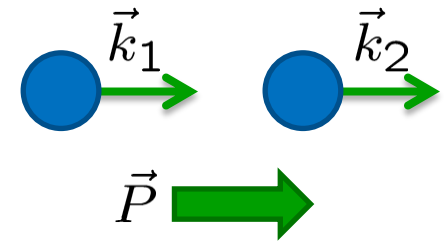
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'Generalised Clebsch-Gordans' $\Lambda_1 \otimes \Lambda_2 \rightarrow \Lambda$

$\Lambda_1 \in \text{LG}(\vec{k}_1)$, $\Lambda_2 \in \text{LG}(\vec{k}_2)$, $\Lambda \in \text{LG}(\vec{P})$

(calculate using induced representation)



[PR D86, 034031 (2012)]

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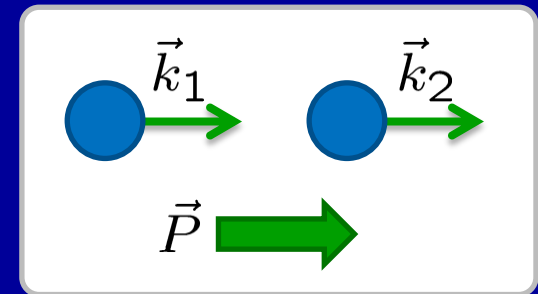
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Sum over all \vec{k}_1, \vec{k}_2 related by allowed lattice rot. such that $\vec{P} = \vec{k}_1 + \vec{k}_2$
 $R \vec{k}_{1,2} \forall R \in \text{LG}(\vec{P})$

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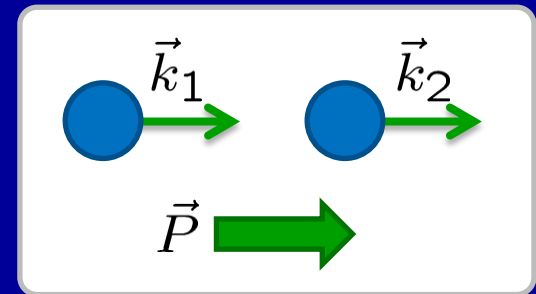
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Why this approach?

- Don't mix different Λ_i
- Can use optimised single-hadron ops
- Can iteratively construct >2 hadron ops

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Compare continuum formulation:

$$\mathbb{O}_{J, M}^{\dagger[S, \ell]} \sim \sum_{\lambda_1 \lambda_2} \int d\hat{p} C(J, \ell, S, M; \vec{p}, S_1, \lambda_1; -\vec{p}, S_2, \lambda_2) \mathbb{O}^{\dagger S_1 \lambda_1}(\vec{p}) \mathbb{O}^{\dagger S_2 \lambda_2}(-\vec{p})$$
$$C = \langle S_1, \lambda_1; S_2, -\lambda_2 | S, \lambda \rangle \langle \ell, 0; S, \lambda | J, \lambda \rangle \mathcal{D}_{M\lambda}^{(J)*}(\hat{p})$$

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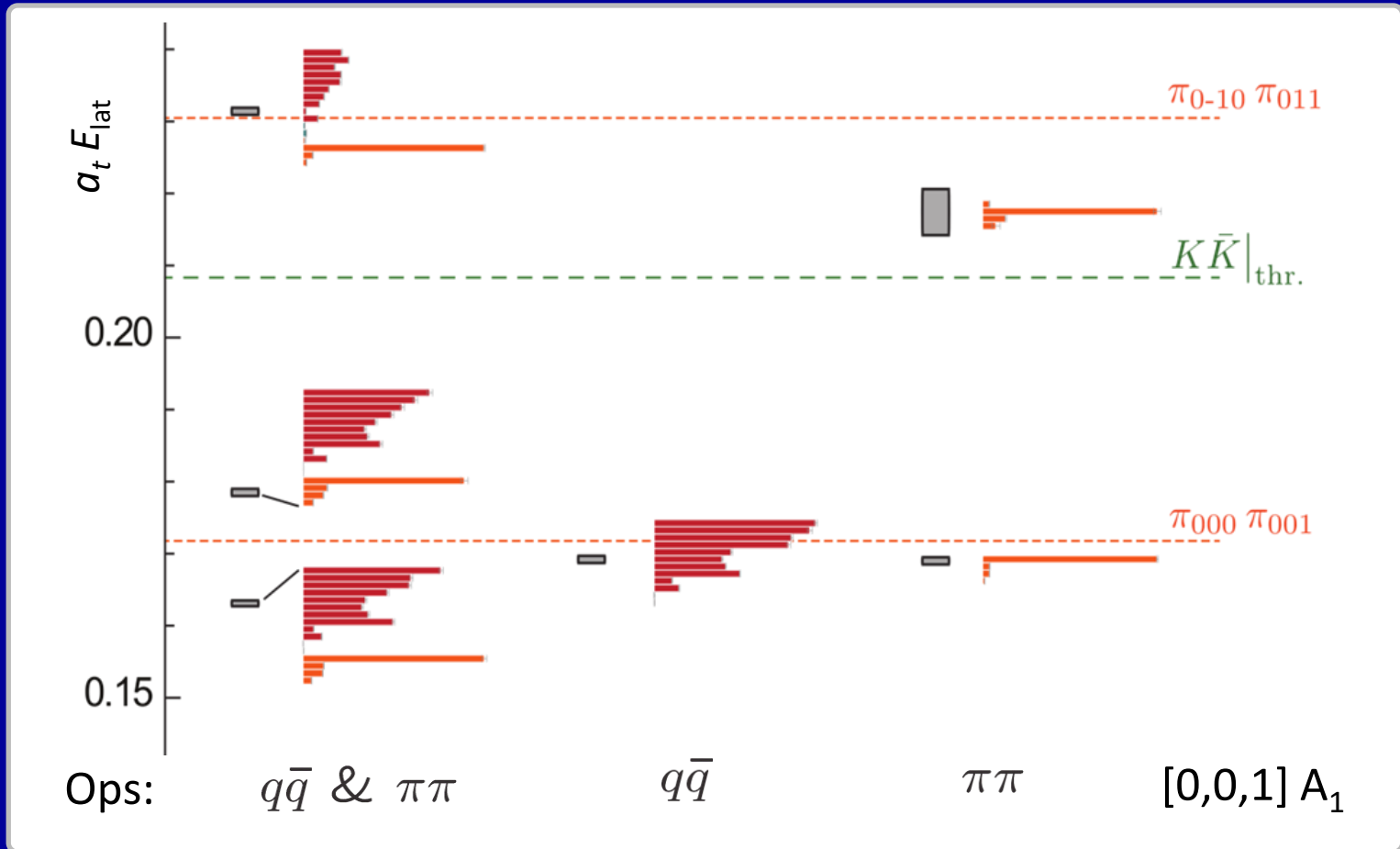
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Operators and energy levels

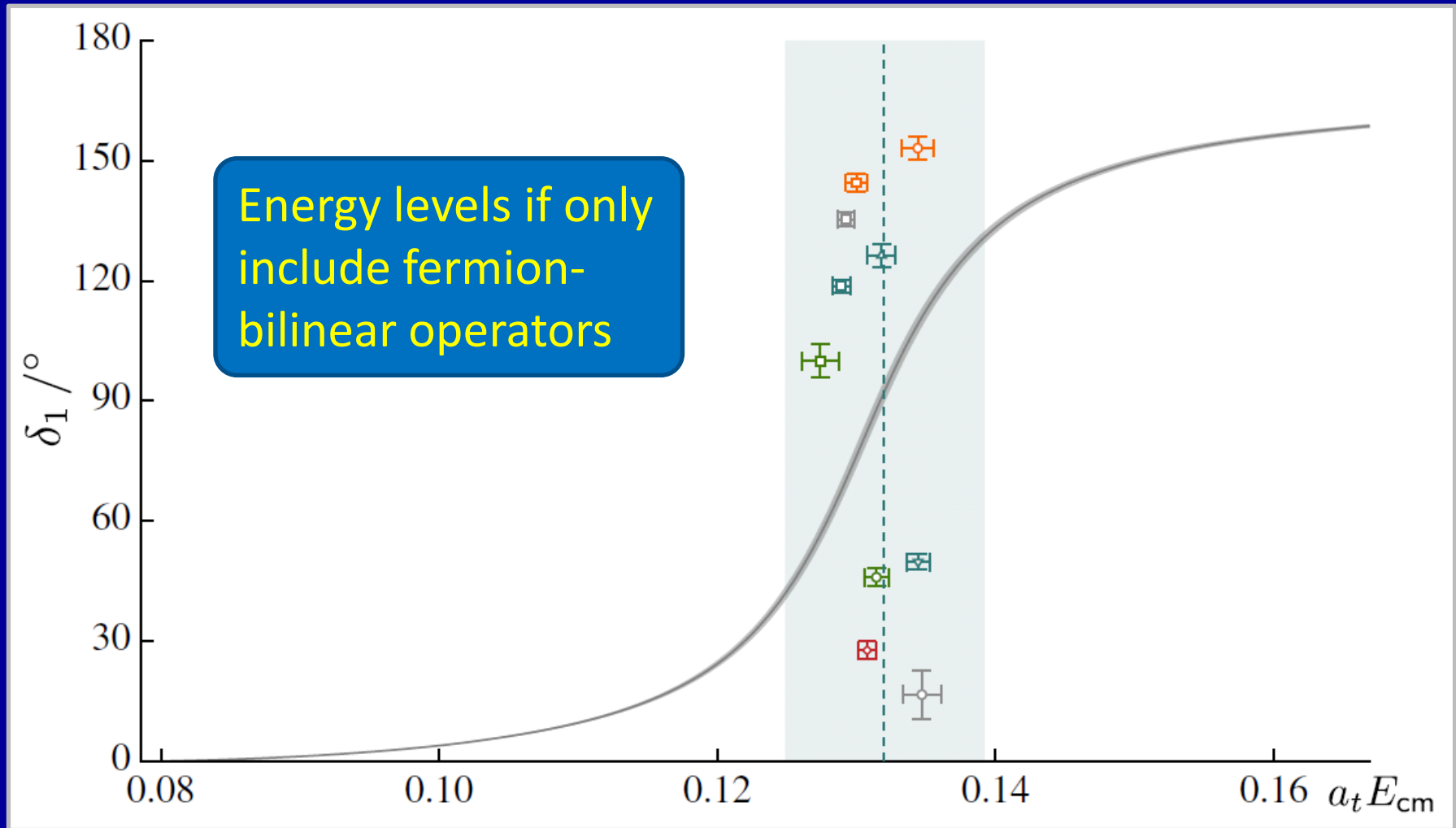
$l=1$ $\pi\pi$ scattering (ρ resonance)



$m_\pi \approx 400$ MeV [PR D87, 034505 (2013)]

Operators and energy levels

$l=1$ $\pi\pi$ scattering (ρ resonance)



$m_\pi \approx 236$ MeV [PR D92, 094502 (2015)]

$\rho\pi$ isospin-2 scattering

[Woss, CT, Dudek, Edwards,
Wilson, JHEP 1807, 043 (2018)]

Different partial waves with same J^P can mix dynamically

E.g. $J^P = 1^+$ (${}^3S_1, {}^3D_1$), $J^P = 2^-$ (${}^3P_2, {}^3F_2$)

$$\mathbf{t} = \begin{bmatrix} t({}^3S_1|{}^3S_1) & t({}^3S_1|{}^3D_1) \\ t({}^3S_1|{}^3D_1) & t({}^3D_1|{}^3D_1) \end{bmatrix}$$

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Finite-volume lattice QCD
calculations: reduced symmetry
→ additional 'mixing'

$[0,0,0]$	T_1^+
$J^+(^3\ell_J)$	$1^+ \begin{pmatrix} ^3S_1 \\ ^3D_1 \end{pmatrix}$
	$3^+ \begin{pmatrix} ^3D_3 \\ ^3G_3 \end{pmatrix}$
	$4^+ \begin{pmatrix} ^3G_4 \end{pmatrix}$

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Finite-volume lattice QCD calculations: reduced symmetry
 \rightarrow additional 'mixing'

$$\mathbf{t} = \begin{bmatrix} t(^3S_1|{}^3S_1) & t(^3S_1|{}^3D_1) & 0 \\ t(^3S_1|{}^3D_1) & t(^3D_1|{}^3D_1) & 0 \\ 0 & 0 & t(^3D_3|{}^3D_3) \end{bmatrix}$$

$[0,0,0]$	T_1^+
$J^+({}^3\ell_J)$	$1^+ \begin{pmatrix} {}^3S_1 \\ {}^3D_1 \end{pmatrix}$
	$3^+ \begin{pmatrix} {}^3D_3 \\ {}^3G_3 \end{pmatrix}$
	$4^+ ({}^3G_4)$

$[00n] \Lambda$	A_1	A_2	E
		$0^- \left({}^3P_0 \right)$	
		$1^+ \left(\begin{matrix} {}^3S_1 \\ {}^3D_1 \end{matrix} \right)$	$1^+ \left(\begin{matrix} {}^3S_1 \\ {}^3D_1 \end{matrix} \right)$
	$1^- \left({}^3P_1 \right)$		$1^- \left({}^3P_1 \right)$
	$2^+ \left({}^3D_2 \right)$		$2^+ \left({}^3D_2 \right)$
$J^P \left({}^3\ell_J \right)$		$2^- \left(\begin{matrix} {}^3P_2 \\ {}^3F_2 \end{matrix} \right)$	$2^- \left(\begin{matrix} {}^3P_2 \\ {}^3F_2 \end{matrix} \right)$
		$3^+ \left(\begin{matrix} {}^3D_3 \\ {}^3G_3 \end{matrix} \right)$	$3^+ \left(\begin{matrix} {}^3D_3 \\ {}^3G_3 \end{matrix} \right)_{[2]}$
	$3^- \left({}^3F_3 \right)$		$3^- \left({}^3F_3 \right)_{[2]}$
	$4^- \left(\begin{matrix} {}^3F_4 \\ {}^3H_4 \end{matrix} \right)$	$4^- \left(\begin{matrix} {}^3F_4 \\ {}^3H_4 \end{matrix} \right)_{[2]}$	$4^- \left(\begin{matrix} {}^3F_4 \\ {}^3H_4 \end{matrix} \right)_{[2]}$

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

Calc. at $SU(3)_F$ symmetric point ($m_u=m_d=m_s$)

$$m_\pi \approx 700 \text{ MeV}$$

$$m_\rho \approx 1020 \text{ MeV (stable)}$$

Large bases of $\rho\pi$ $SU(3)_F$ ops:

$$\text{'optimised' } \rho \text{ and } \pi \text{ ops } \sim \bar{\psi} \Gamma D \dots \psi$$

Use 'distillation' [PR D80 054506 (2009)]

One lattice spacing, 2 volumes ($L \approx 2, 3 \text{ fm}$)

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

$$\begin{aligned}\rho[000] &\rightarrow T_1^-, & \pi[000] &\rightarrow A_1^- \\ [000] T_1^- \times [000] A_1^- &\rightarrow [000] T_1^+\end{aligned}$$

$$1 T_1^+$$

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

$$\begin{aligned} \rho[000] &\rightarrow T_1^-, \quad \pi[000] \rightarrow A_1^- \\ [000] T_1^- \times [000] A_1^- &\rightarrow [000] T_1^+ \end{aligned}$$

1 T_1^+

$$\begin{aligned} \rho[001] \lambda=0 &\rightarrow A_1, \quad \lambda=\pm 1 \rightarrow E_2, \quad \pi[001] \rightarrow A_2 \\ [001] A_1 \times [001] A_2 &\rightarrow [000] T_1^+, A_1^-, E^- \\ [001] E_2 \times [001] A_2 &\rightarrow [000] T_1^+, T_2^+, T_1^-, T_2^- \end{aligned}$$

2 T_1^+

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

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1 T_1^+

$$\begin{aligned} \rho[001] \lambda=0 &\rightarrow A_1, \quad \lambda=\pm 1 \rightarrow E_2, \quad \pi[001] \rightarrow A_2 \\ [001] A_1 \times [001] A_2 &\rightarrow [000] T_1^+, A_1^-, E^- \\ [001] E_2 \times [001] A_2 &\rightarrow [000] T_1^+, T_2^+, T_1^-, T_2^- \end{aligned}$$

2 T_1^+

$$\begin{aligned} \rho[011] \lambda=0 &\rightarrow A_1, \quad \lambda=\pm 1 \rightarrow B_1, B_2, \quad \pi[011] \rightarrow A_2 \\ [011] A_1 \times [011] A_2 &\rightarrow [000] T_1^+, \dots \\ [011] B_1 \times [011] A_2 &\rightarrow [000] T_1^+, \dots \\ [011] B_2 \times [011] A_2 &\rightarrow [000] T_1^+, \dots \end{aligned}$$

3 T_1^+

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

$$\rho[000] \rightarrow T_1^-, \quad \pi[000] \rightarrow A_1^-$$
$$[000] T_1^- \times [000] A_1^- \rightarrow [000] T_1^+$$

1 T_1^+

$$\rho[001] \lambda=0 \rightarrow A_1, \quad \lambda=\pm 1 \rightarrow E_2, \quad \pi[001] \rightarrow A_2$$
$$[001] A_1 \times [001] A_2 \rightarrow [000] T_1^+, A_1^-, E^-$$
$$[001] E_2 \times [001] A_2 \rightarrow [000] T_1^+, T_2^+, T_1^-, T_2^-$$

2 T_1^+

$$\rho[011] \lambda=0 \rightarrow A_1, \quad \lambda=\pm 1 \rightarrow B_1, B_2, \quad \pi[011] \rightarrow A_2$$
$$[011] A_1 \times [011] A_2 \rightarrow [000] T_1^+, \dots$$
$$[011] B_1 \times [011] A_2 \rightarrow [000] T_1^+, \dots$$
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3 T_1^+

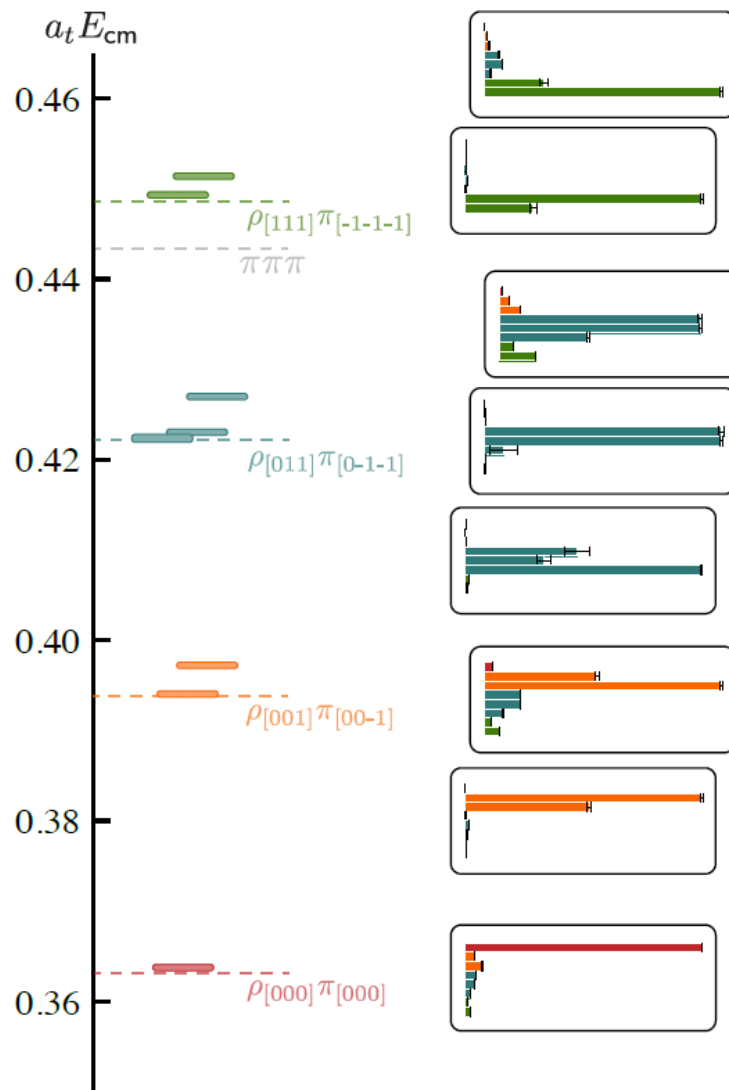
Overall non-zero momentum

$$[001] A_1 \times [011] A_2 \rightarrow [001] E_2, \dots$$
$$[001] E_2 \times [011] A_2 \rightarrow [001] E_2, E_2, \dots$$

3 E_2

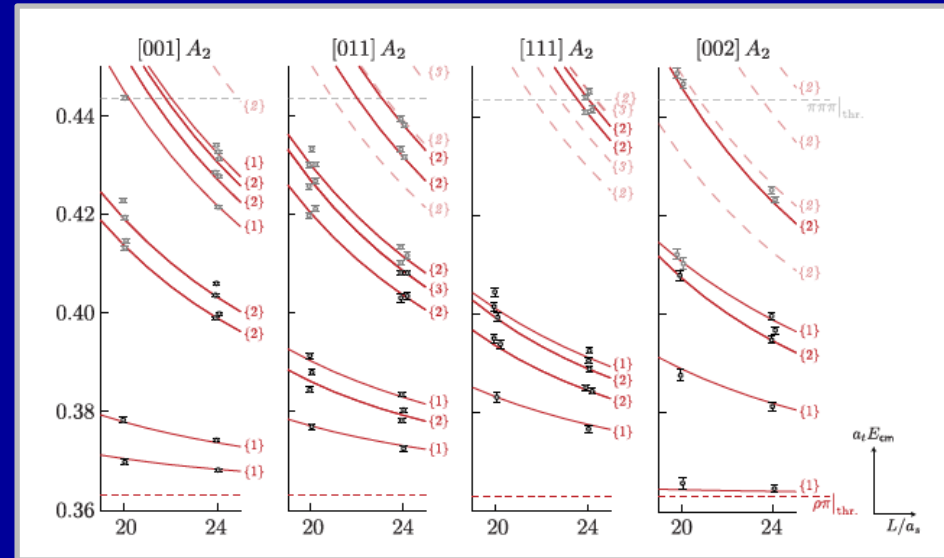
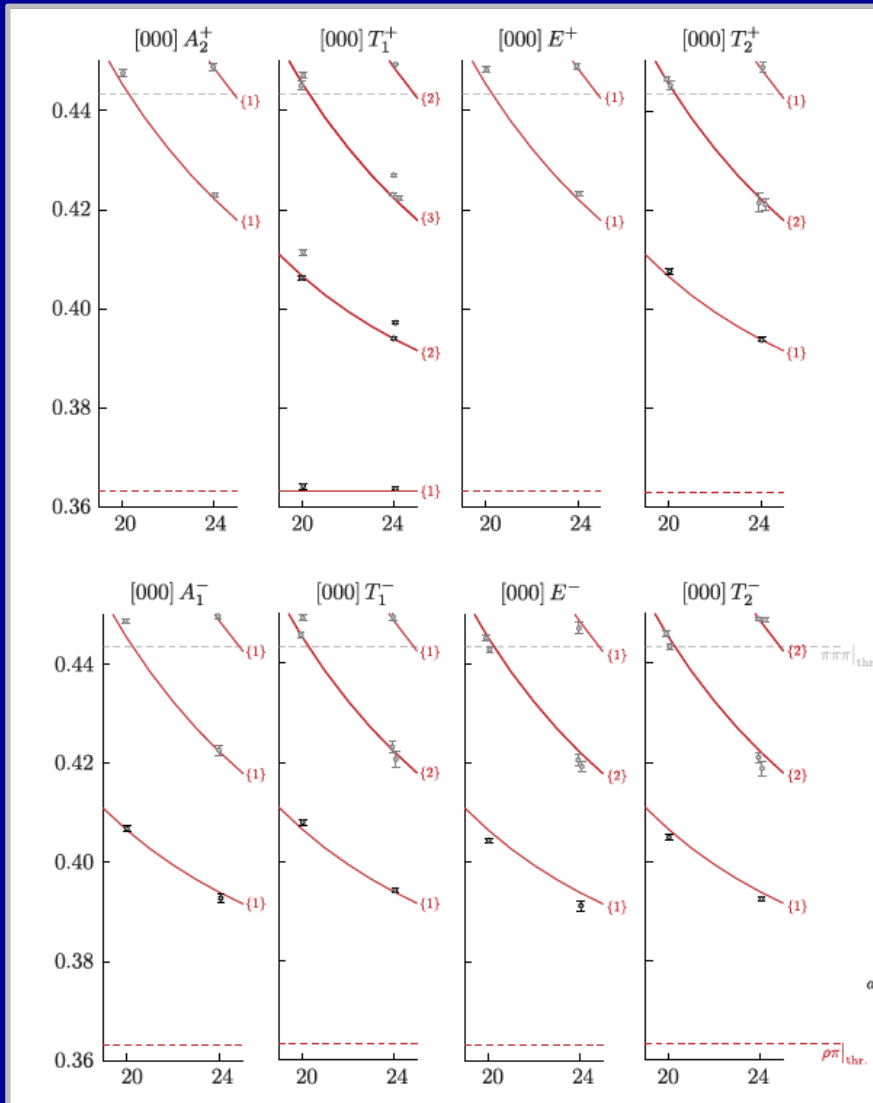
$[000] T_1^+ (1^+, 3^+, \dots)$

ops
$\rho_{[000]}\pi_{[000]}$
$\{2\} \rho_{[001]}\pi_{[00-1]}$
$\{3\} \rho_{[011]}\pi_{[0-1-1]}$
$\{2\} \rho_{[111]}\pi_{[-1-1-1]}$



$\rho\pi$ isospin-2 scattering – spectra

[JHEP 1807, 043 (2018)]

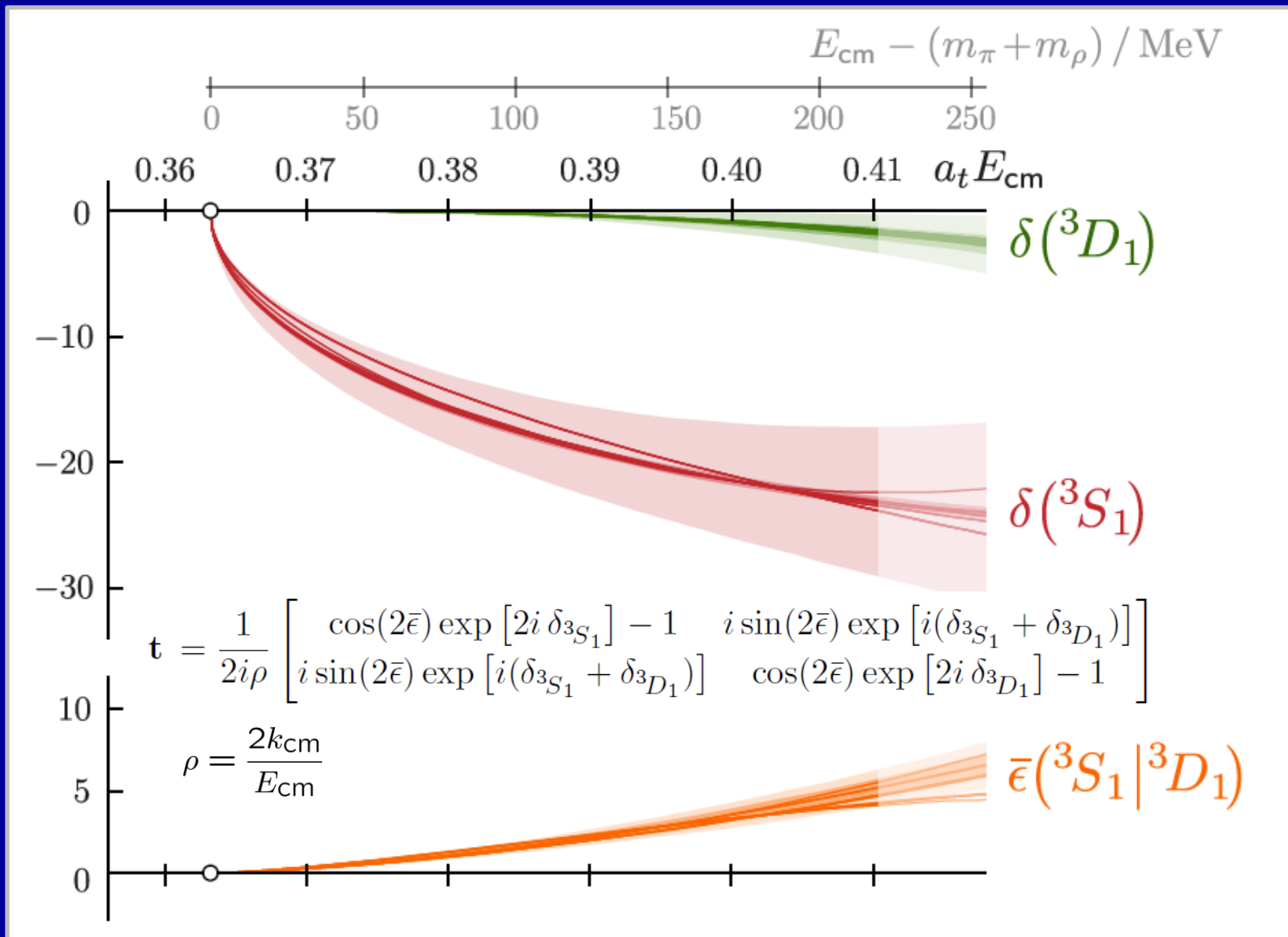


+ others

Used 141 energy levels for $\ell = 0, 1, 2$

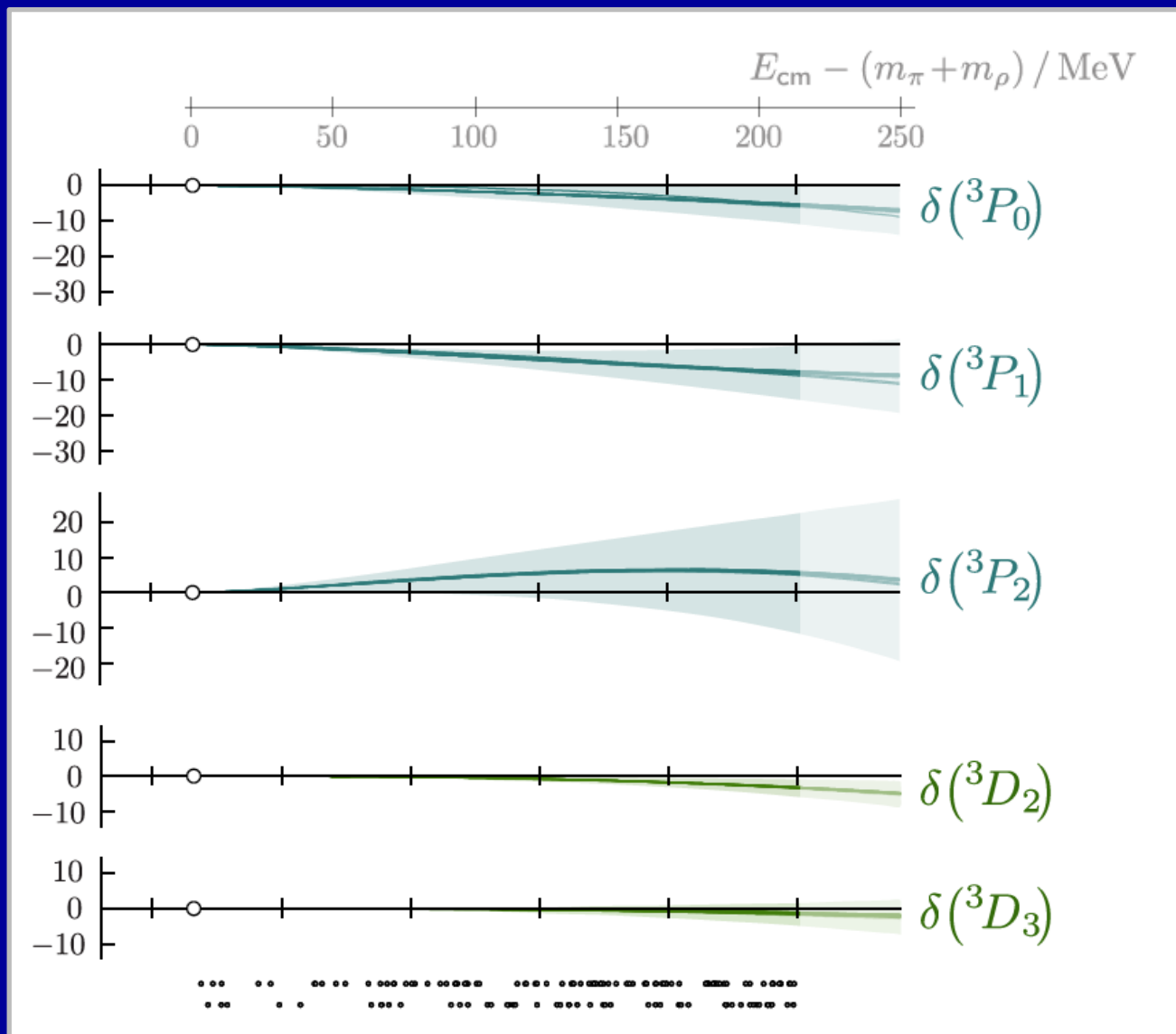
$\rho\pi$ isospin-2 scattering – amplitudes

[JHEP 1807, 043 (2018)]



$\rho\pi$ isospin-2 scattering – amplitudes

[JHEP 1807, 043 (2018)]



Summary

- General multi-hadron operator constructions
- Example of recent progress:
 - $\rho\pi$ isospin-2 scattering with dynamically-coupled $^3S_1, ^3D_1$ partial waves
- Work in progress on other channels and different m_π
- Also matrix elements, transitions

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Extra slides

$\rho\pi$ isospin-2 scattering

[JHEP 1807, 043 (2018)]

