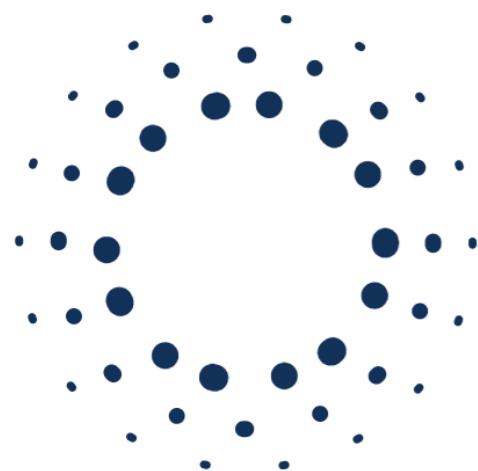


Testing MOdified Gravity (MOG) Theory with the Milky Way

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Tensions in the LCDM paradigm

16/05/2018

Goal

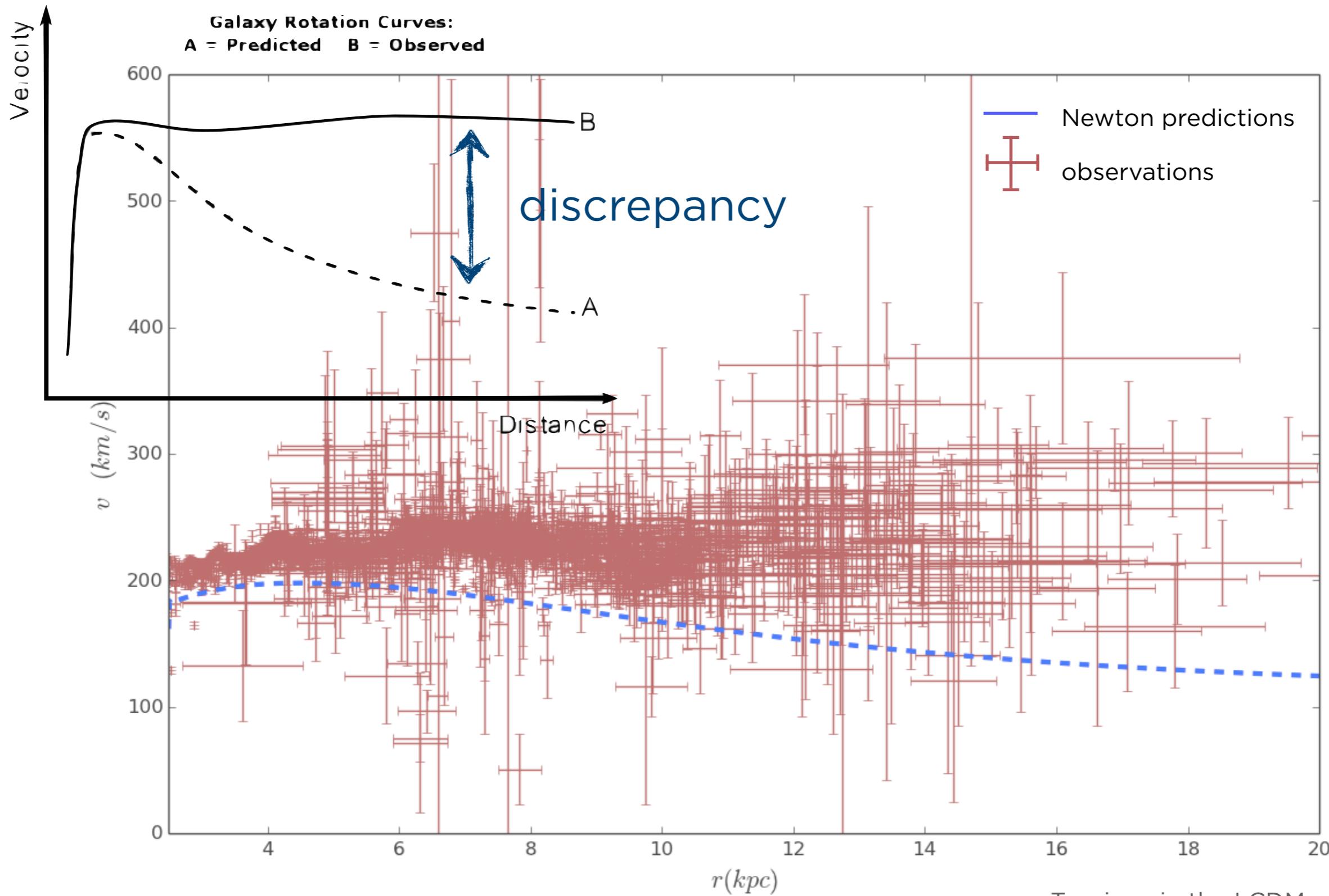
- Test MOdified Gravity (MOG) theory of John Moffat using tracers of the total gravitational potential of the MW
- Methodology already used for testing MOND phenomenology.

B. Famaey & J. Binney,
MNRAS 363 (2005)
[astro-ph/0506723]

S. S. McGaugh, ApJ 683 (2008)
[0804.1314]

F. Iocco, M. Pato & G. Bertone,
Phys. Rev. D92 (2015)
[1505.05181]

Rotation Curve Milky Way



Motivation

- Observations of the dynamics of galaxies reveal a discrepancy between dynamical mass and the mass inferred from luminous matter.

V.C. Rubin +, ApJ 141 (1965)

V.C. Rubin & W.K. Ford Jr., ApJ 159 (1970)

- A proposal for explaining the mismatch is a modification of gravity.

M. Milgrom, ApJ 270 (1983)

J.D. Bekenstein, Phys. Rev. D70 (2004)
[astro-ph/0403694]

J.W. Moffat, JCAP 03 (2006)
[gr-qc/0506021v7]

MOdified Gravity (MOG) Theory

J.W. Moffat, JCAP 03 (2006)
[gr-qc/0506021v7]

■ Scalar-Tensor-Vector Gravity (STVG) theory

Gravitational action:

$$S_G = -\frac{1}{16\pi} \int \frac{1}{G} (R + 2\Lambda) \sqrt{-g} d^4x$$

Massive vector field action:

$$S_\phi = -\frac{1}{4\pi} \int \omega \left[\frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{1}{2} \mu^2 \phi_\mu \phi^\mu + V_\phi(\phi_\mu \phi^\mu) \right] \sqrt{-g} d^4x$$

$$B_{\mu\nu} = \partial_\mu \phi_\nu - \partial_\nu \phi_\mu$$

Scalar fields action:

$$S_S = - \int \frac{1}{G} \left[\frac{1}{2} g^{\alpha\beta} \left(\frac{\nabla_\alpha G \nabla_\beta G}{G^2} + \frac{\nabla_\alpha \mu \nabla_\beta \mu}{\mu^2} \right) + \frac{V_G(G)}{G^2} + \frac{V_\mu(\mu)}{\mu^2} \right] \sqrt{-g} d^4x$$

MOdified Gravity (MOG) Theory

J.W. Moffat, JCAP 03 (2006)
[gr-qc/0506021v7]

- Able to explain data coming from:
 - motion of globular clusters (J.W. Moffat & V.T. Toth, ApJ 680 (2008) [0708.1935])
 - galaxy clusters (e.g. J.W. Moffat & S. Rahvar, MNRAS 441 (2014) [1309.5077])
 - rotation curves (RCs) of spiral and dwarf galaxies (e.g. M.H. Zholideh Haghghi & S. Rahvar, MNRAS 468 (2017) [1609.07851])
 - Bullet Cluster (J.R. Brownstein & J.W. Moffat, MNRAS 382 (2007) [astro-ph/0702146])
- Unable to explain:
 - Bullet Cluster (D. Clowe +, ApJ 648 (2006) [astro-ph/0608407])
 - galaxy clusters (e.g. T.M. Nieuwenhuizen +, MNRS 476 (2018) [1802.04891])

MOdified Gravity (MOG) Theory

J.W. Moffat & S. Rahvar, MNRAS 436 (2013)
[1306.6383]

Weak field limit:

$$\Phi_{eff}(\vec{x}) = - \int \frac{G_0 \rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3 \vec{x}' + \kappa^2 \int \frac{e^{-\mu |\vec{x} - \vec{x}'|}}{|\vec{x} - \vec{x}'|} \rho(\vec{x}') d^3 \vec{x}'$$



$$G_0 - \kappa^2 = G_N$$

$$\alpha = (G_\infty - G_N)/G_N$$

$$\vec{a} = -\vec{\nabla} \Phi_{eff}$$

$$\vec{a}(\vec{x}) = -G_N \int \frac{\rho(\vec{x}')(\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3} \times \left[1 + \alpha - \alpha e^{-\mu |\vec{x} - \vec{x}'|} (1 + \mu |\vec{x} - \vec{x}'|) \right] d^3 \vec{x}'$$

For $r < 1/\mu$: recover Newtonian gravity

MOdified Gravity (MOG) Theory

Weak field limit:

$$\vec{a}(\vec{x}) = -G_N \int \frac{\rho(\vec{x}')(\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3} \times \left[1 + \alpha - \alpha e^{-\mu|\vec{x} - \vec{x}'|} (1 + \mu|\vec{x} - \vec{x}'|) \right] d^3 \vec{x}'$$

Parameters (α, μ):

- control the strength and the range of the repulsive force and
- can be expressed in terms of the mass of the system

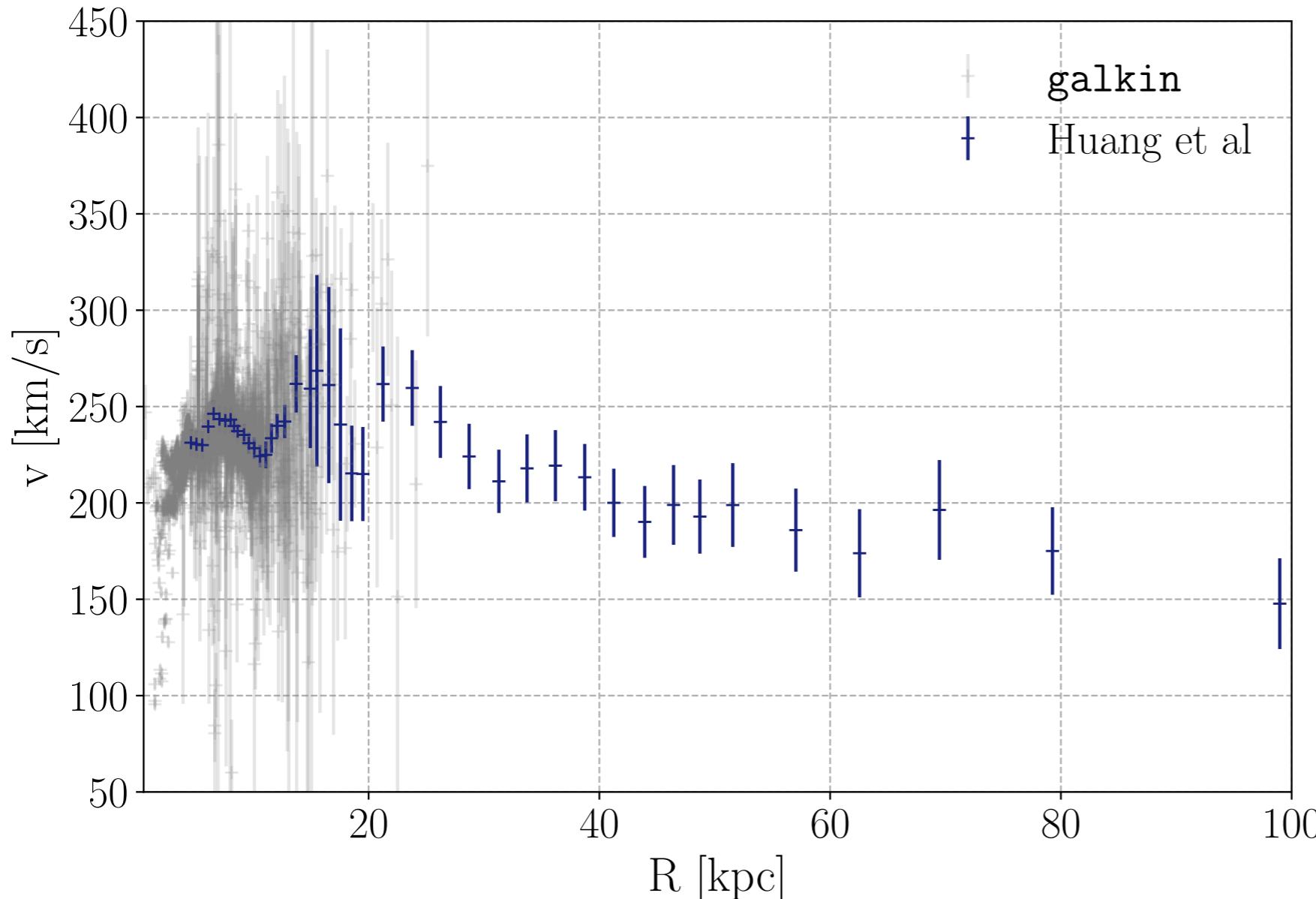
$$\mu = \frac{D}{\sqrt{M}} \quad \alpha = \frac{M}{(\sqrt{M} + E)^2} \left(\frac{G_\infty}{G_N} - 1 \right)$$

(D, E, G_∞) are constants that can be determined by observations

J.W. Moffat & V.T. Toth, Classical and Quantum Gravity 26 (2009)
[0712.1796v5]

Methodology: Rotation Curve

Observed RC: Two different compilations



$$R_0 = 8.34 \text{ kpc}$$
$$V_0 = 239.89 \text{ km/s}$$

$$(U_\odot, V_\odot, W_\odot) = (7.01, 10.13, 4.95) \text{ km/s}$$

Tensions in the LCDM paradigm
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Methodology: Rotation Curve

■ Observed RC: Two different compilations

M. Pato & F. Iocco,
SoftwareX 6 (2017)
[1703.00020]

Y. Huang +,
MNRAS 463 (2016)
[1604.01216]

■ Visible component (star + gas): Large array of observationally inferred 3D density profiles

$$\rho_{\text{bulge}}(x, y, z)$$

$$\rho_{\text{disc}}(r, z)$$

$$\rho_{\text{gas}}(x, y, z)$$

$$\Phi_{\text{bulge}}(x, y, z)$$

$$\Phi_{\text{disc}}(r, z)$$

$$\Phi_{\text{gas}}(x, y, z)$$

$$v^2(r) = \sum_i v_i^2(r)$$

$$v_i^2(r) = r \frac{d\Phi_i}{dr}$$

Data-based
density profiles

Gravitational
potential

Expected
RC

50 kpc

$R_0 = 8 \text{ kpc}$

Most of the galaxy's light comes from stars and gas in the galactic disk and central bulge ...

SUN

Galactic Bulge region

Stellar + gas disc

... but measurements suggest that most of the mass lies unseen in the spherical halo that surrounds the entire disk.

The visible Milky Way

Bulge distribution: $\rho_b(x, y, z) = \bar{\rho}_b f(x, y, z)$

$f(x, y, z)$	Bar angle [°]	x ₀ :y ₀ :z ₀	Reference
e^{-r}	25	2.8 : 1.4 : 1	K.Z. Stanek + (1996) [G2]
$e^{-r_s^2/2}$	24	3.6 : 1.5 : 1	K.Z. Stanek + (1996) [E2]
$e^{-r_s^2/2} + r_a^{-1.85} e^{-r_a}$	20	3.7 : 1.5 : 1	H. Zhao (1996)
$e^{-r_s^2}/(1 + r_s)^{1.8}$	20	2.6 : 0.8 : 1	N. Bissantz & O. Gerhard (2002)
$\text{sech}^2(-r_s) + e^{-r_s}$	13	3.7 : 1.3 : 1	A.C. Robin + (2012)
$e^{-r_s^2}/(1 + r_s)^{1.8}$	15	3.2 : 2.2 : 1	E. Vanhollebeke + (2013)

Normalisation $\bar{\rho}_b$

$$\langle \tau \rangle = 2.17_{-0.38}^{+0.47} \times 10^{-6} \quad (\ell, b) = (1.50^\circ, -2.68^\circ)$$

P. Popowski +, ApJ 631 (2005)
[astrop-ph/0410319]

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Stellar disc distribution: $\rho_d(r, z) = \bar{\rho}_d f(r, z)$

$f(r, z)$		Scale-length [kpc]	Scale-height [kpc]	Reference
$e^{-r} \operatorname{sech}^2(z)$	thin	2.75	0.27 $\eta(r)$	C. Han & A. Gould (2003)
	thick	2.75	0.44 $\eta(r)$	
$e^{-r} e^{- z }$ $e^{-r} e^{- z }$ $(r^2 + z^2)^{-2.77/2}$	thin	2.6	0.30	M. Juric + (2008)
	thick	3.6	0.90	
	halo			
$e^{-r} e^{- z }$ $e^{-r} e^{- z }$ $(r^2 + z^2)^{-2.75/2}$	thin	2.75	0.25	J. T. A. De Jong + (2010)
	thick	4.1	0.75	
	halo			
$e^{-r} e^{- z }$ $e^{-r} e^{- z }$	thin	2.75	0.25	S. Calchi Novati & L. Mancini (2011)
	thick	4.1	0.75	
$e^{-r} e^{- z }$	single	2.15	0.4	J. Bovy & H.W. Rix (2013)

Normalisation $\bar{\rho}_d$

$$\Sigma_*(R_0) = 38 \pm 4 \text{ M}_\odot \text{pc}^{-2}$$

J. Bovy & H.W. Rix, ApJ 779
(2013) [1309.0809]

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Gas distribution:

$$\rho_g(x, y, z) = \rho_{\text{H}_2}(x, y, z) + \rho_{\text{H}_\text{I}}(x, y, z) + \rho_{\text{H}_{\text{II}}}(x, y, z)$$

Components		Range	Reference
molecular ring	H_2	$r = 3 - 20 \text{ kpc}$	K. Ferrière (1998)
cold, warm	HI		
warm, hot	HII		
CMZ, disc	H_2	$r = 0.01 - 3 \text{ kpc}$	K. Ferrière + (2007)
CMZ, disc	HI		
warm, hot, very hot	HII		

Uncertainties

CO-to- H_2 factor: $X_{\text{CO}}(r > 3 \text{ kpc}) = (5.0 \pm 2.5) \times 10^{19} \text{ cm}^{-2} \text{ K}^{-1} \text{ km}^{-1} \text{s}$

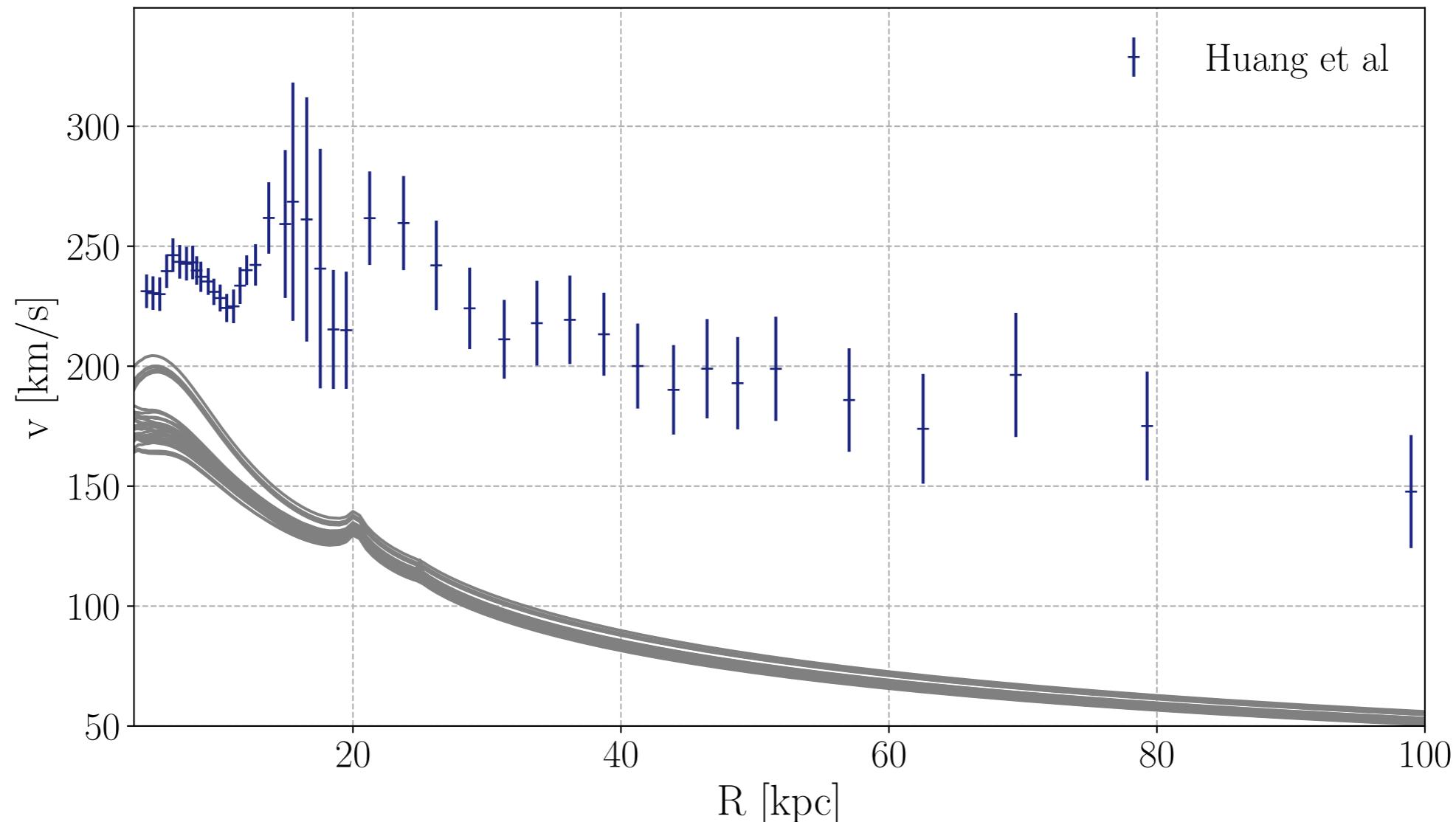
$$X_{\text{CO}}(r < 3 \text{ kpc}) = (1.9 \pm 1.4) \times 10^{20} \text{ cm}^{-2} \text{ K}^{-1} \text{ km}^{-1} \text{s}$$

K. Ferriere +, ApJ 467
(2007) [astro-ph/0702532]

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Baryonic morphologies

$$\begin{array}{ccc} \rho_{\text{bulge}}(x, y, z) & & \Phi_{\text{bulge}}(x, y, z) \\ \rho_{\text{disc}}(r, z) & \xrightarrow{\quad} & \Phi_{\text{disc}}(r, z) \\ \rho_{\text{gas}}(x, y, z) & & \Phi_{\text{gas}}(x, y, z) \end{array} \quad \begin{aligned} v^2(r) &= \sum_i v_i^2(r) \\ v_i^2(r) &= r \frac{d\Phi_i}{dr} \end{aligned}$$



MOG (α , μ) parameters

■ Rotation curve of the MW

$$(\alpha, \mu)^{\text{MW}} = (15.01, 0.0313 \text{ kpc}^{-1})$$

J.W. Moffat & V.T. Toth, Phys. Revm D91 (2015)
[1411.6701]

■ We use equations

$$\mu = \frac{D}{\sqrt{M}} \quad \alpha = \frac{M}{(\sqrt{M} + E)^2} \left(\frac{G_\infty}{G_N} - 1 \right)$$

to obtain $(\alpha, \mu)^c$ self-consistently with our MW mass determination

$$D = 6.25 \text{ M}_\odot \text{pc}^{-1}$$

$$E = 25 \times 10^3 \text{ M}_\odot^{1/2}$$

$$G_\infty = 20 G_N$$

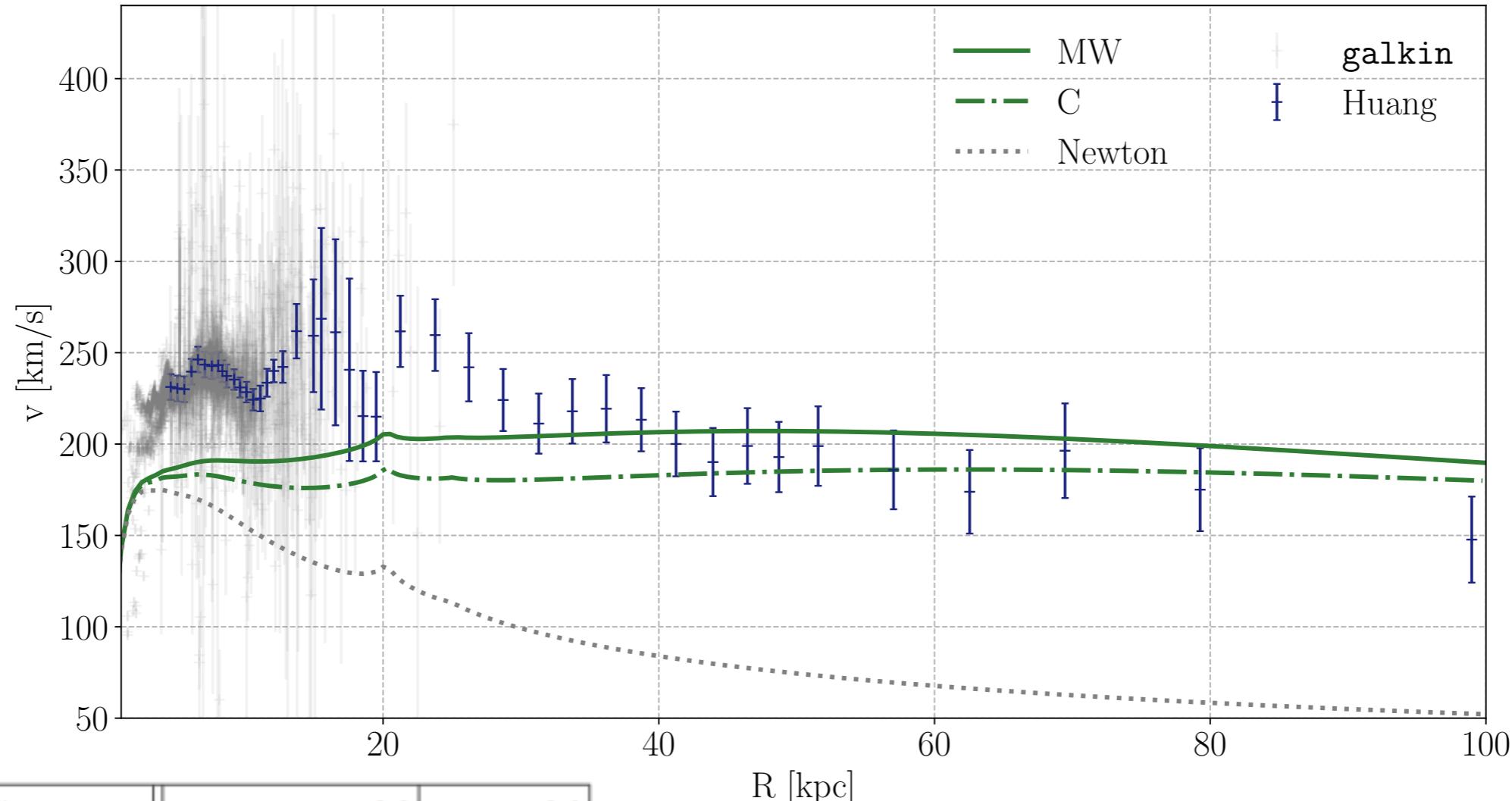
J.W. Moffat & V.T. Toth, Classical & Quantum Gravity 26
(2009) [0712.1796]

Representative Morphology

K.Z. Stanek + (1996) [E2]

S. Calchi Novati & L. Mancini (2011)

$$M_{\text{MW}} = 6.7^{+0.7}_{-0.6} \times 10^{10} M_{\odot}$$

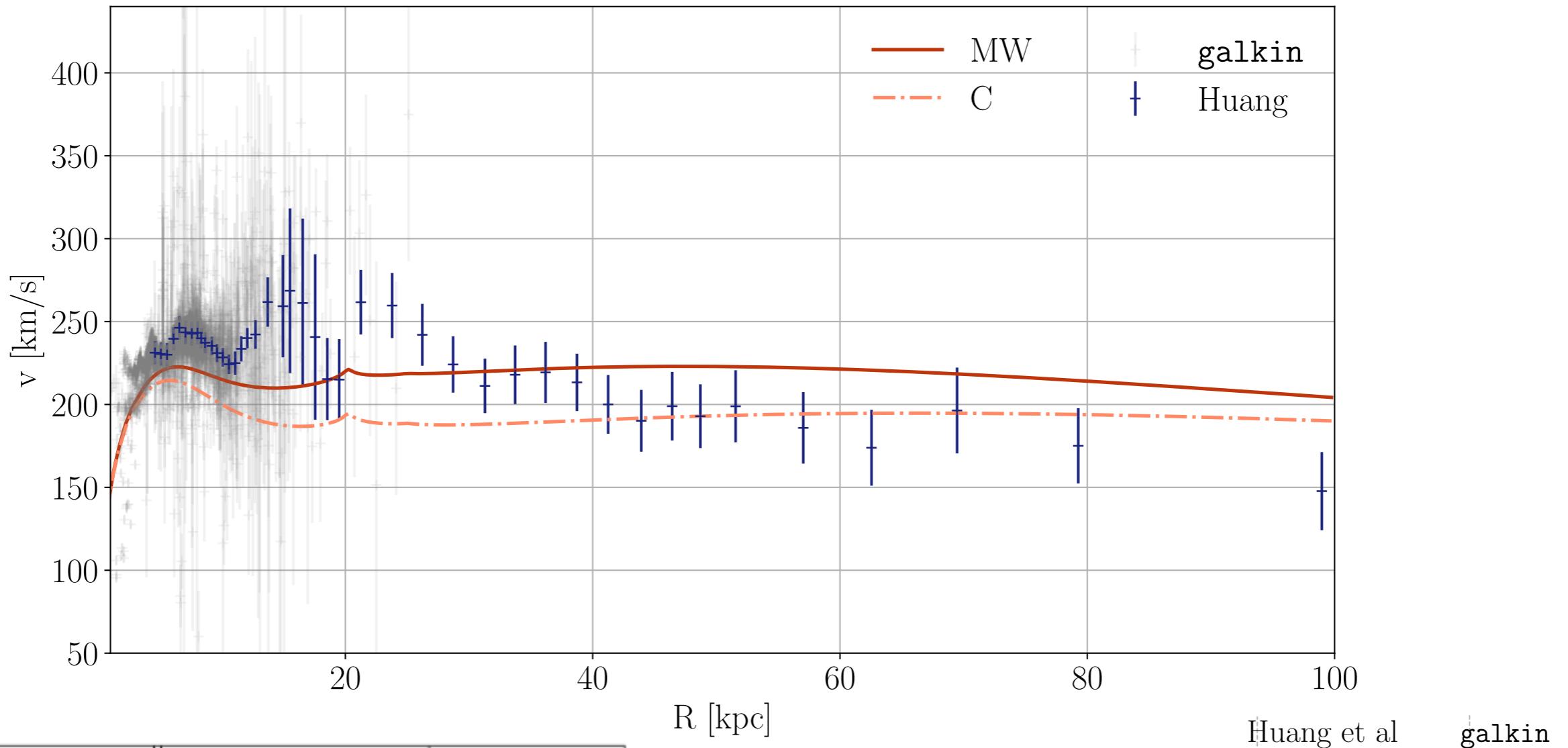


$(\alpha, \mu [\text{kpc}^{-1}])$	Huang et al. [1]	galkin [2]
$(15.01, 3.13 \cdot 10^{-2})^{\text{MW}}$	5.02	8.54
$(15.80, 2.41 \cdot 10^{-2})^{\text{C}}$	9.14	10.61
Newton	32.65	21.84
5σ equivalent $\tilde{\chi}_{5\sigma}^2$	2.41	1.14

Best Performing Morphology

$$M_{\text{MW}} = 7.7^{+0.8}_{-0.7} \times 10^{10} M_{\odot}$$

N. Bissantz & O. Gerhard (2002)
J. Bovy & H.W. Rix (2013)



$(\alpha, \mu [\text{kpc}^{-1}])$	Huang et al. [1]	galkin [2]
$(15.01, 3.13 \cdot 10^{-2})^{\text{MW}}$	1.58	1.98
$(15.98, 2.26 \cdot 10^{-2})^{\text{C}}$	3.82	2.82

2.6 σ discrepancy

5 σ equivalent $\tilde{\chi}_{5\sigma}^2$

2.41

1.14

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Dynamical parameters (α , μ)

$$\vec{a}(\vec{x}) = -G_N \int \frac{\rho(\vec{x}')(\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3} \times \left[1 + \alpha - \alpha e^{-\mu|\vec{x} - \vec{x}'|} (1 + \mu|\vec{x} - \vec{x}'|) \right] d^3\vec{x}'$$

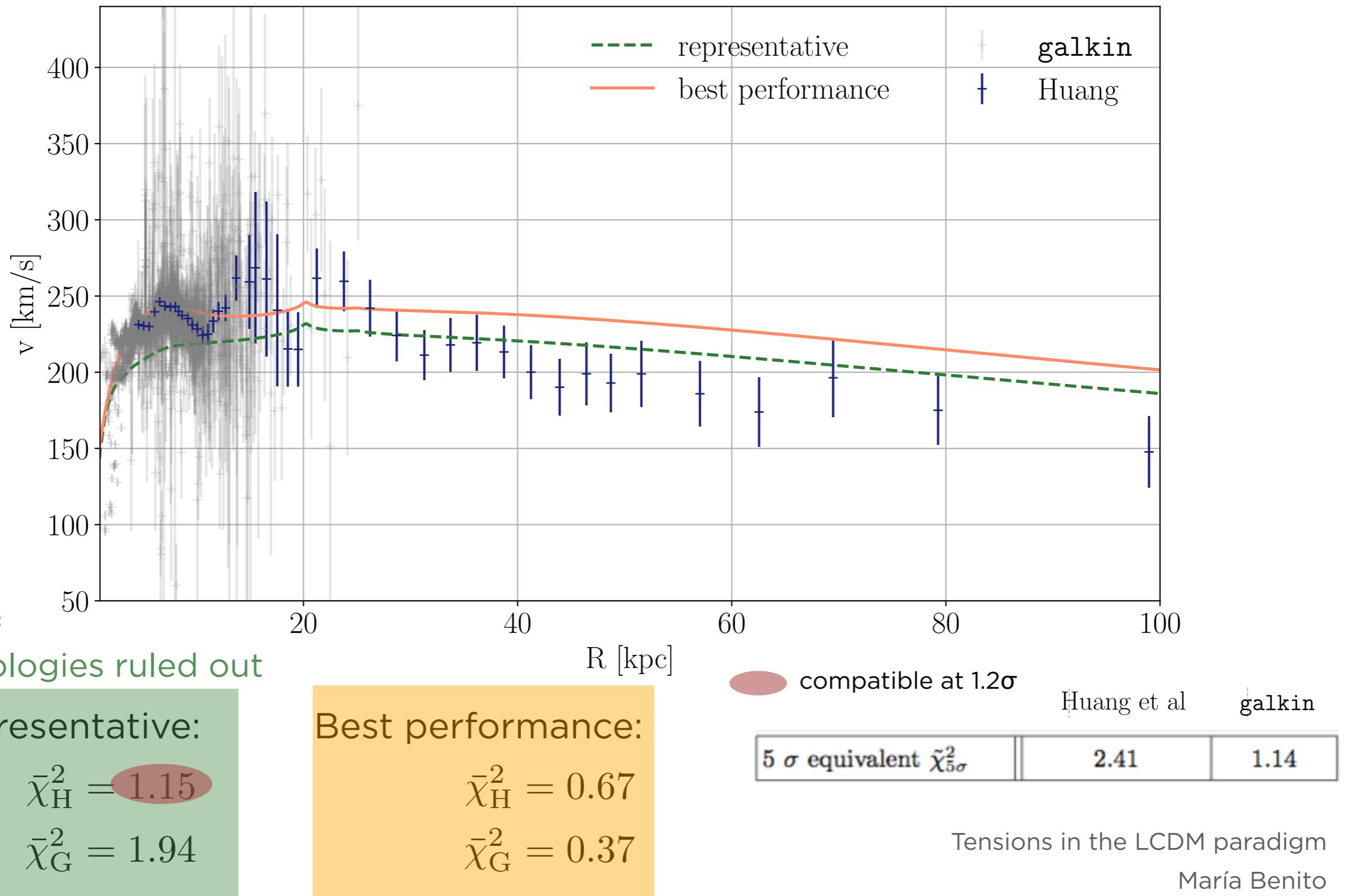
$$\mu = \frac{D}{\sqrt{M}} \quad \alpha = \frac{M}{(\sqrt{M} + E)^2} \left(\frac{G_\infty}{G_N} - 1 \right)$$

(α , μ) calibrated with spiral galaxies

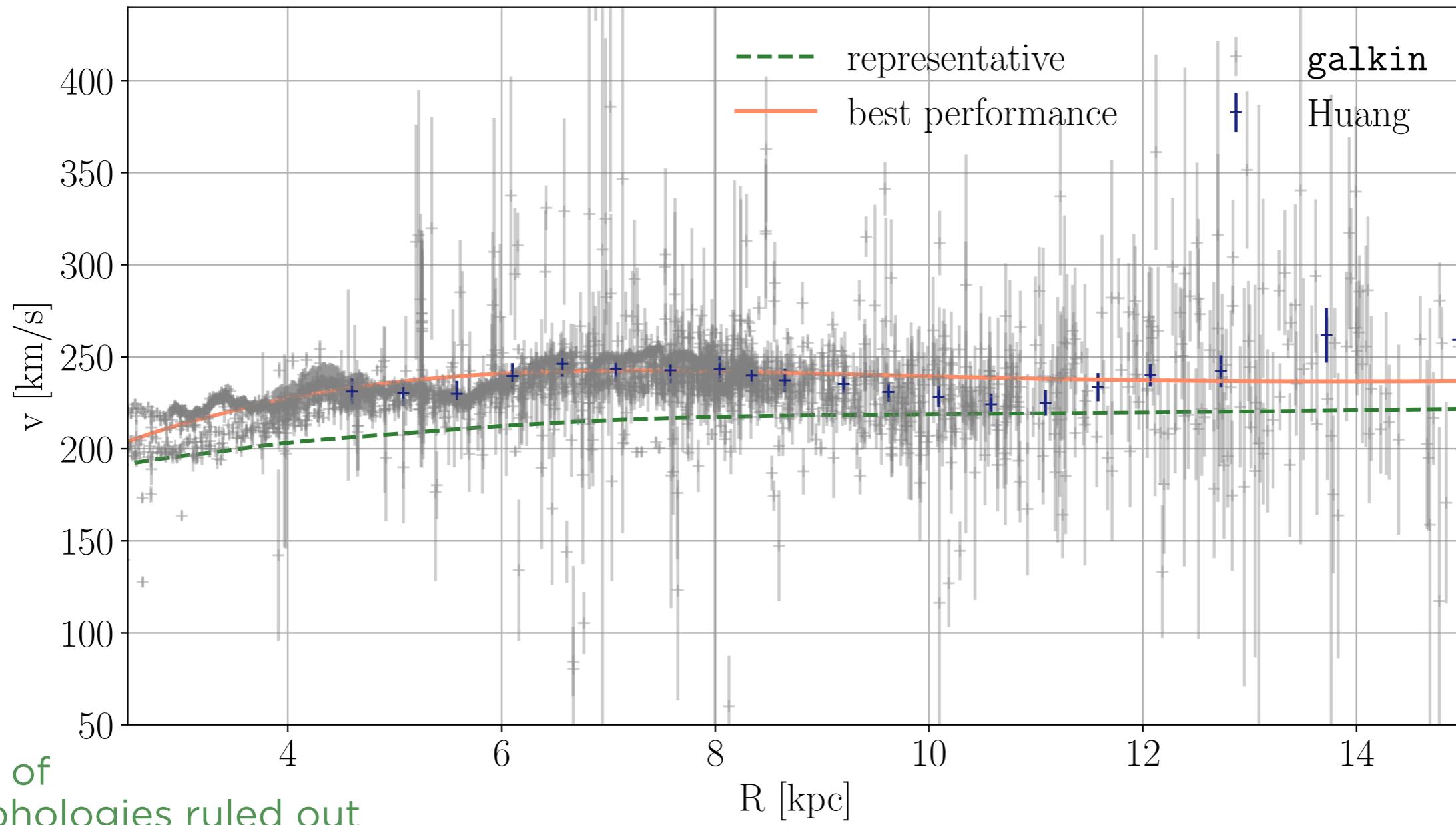
(α , μ) is a function of the enclosed mass:

$$\mu = \mu(R) \propto \int \rho(R) dV \quad \alpha = \alpha(R) \propto \int \rho(R) dV$$

Dynamical parameters (α , μ)



Dynamical parameters (α , μ)



Most of
morphologies ruled out

Representative:

$$\bar{\chi}_H^2 = 1.15$$

$$\bar{\chi}_G^2 = 1.94$$

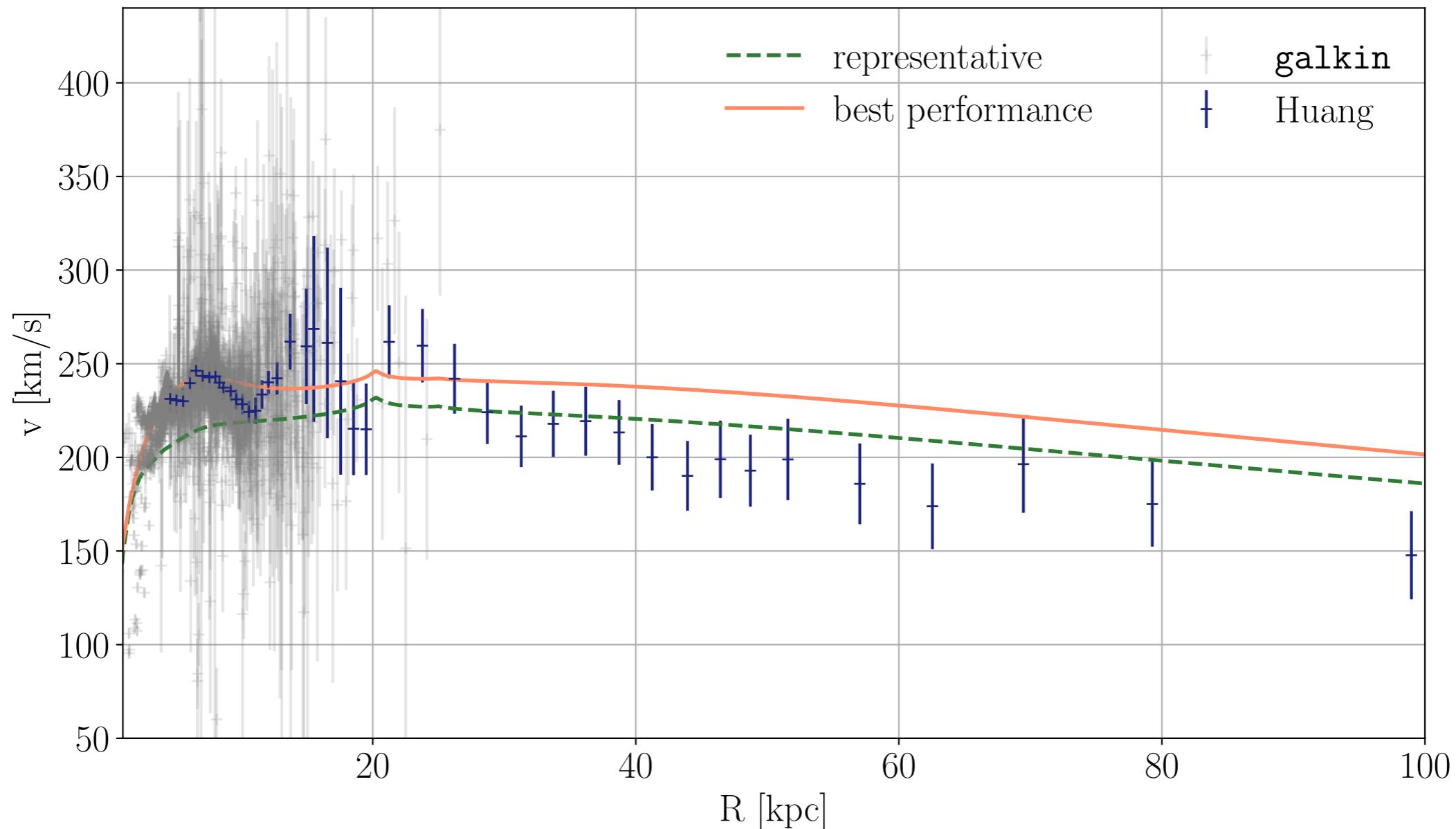
compatible at 1.2σ

5σ equivalent $\tilde{\chi}_{5\sigma}^2$	2.41	1.14
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Dynamical parameters (α , μ)



Best performance:

$$\bar{\chi}_H^2 = 0.67$$

$$\bar{\chi}_G^2 = 0.37$$

5σ equivalent $\tilde{\chi}_{5\sigma}^2$

Huang et al

galkin

2.41

1.14

Conclusions

- Simplified version of MOG ruled out.
- Dynamical (calibrated self-consistently) MOG:
depends on the morphology (ignorance on shape crucial).