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# Constraints on dark matter scenarios from measurements of the galaxy luminosity function at high- $z$

Pier Stefano Corasaniti

*CNRS & Observatoire de Paris*

*Collaborators:*

**Shankar Agarwal (AIMS, Cape Town)**

**Doddy Marsh (Univ. Gottingen )**

**Subinoy Das (IIA, Bangalore)**

**Isabella Carucci (UCL)**

*Publications:*

PSC et al. (2017), arXiv:1611.05892

Carucci & PSC, in preparation

Carucci & PSC (2017), arXiv:1706.09462

*Mainz, "Tensions in the LCDM paradigm", 14-18 May 2018*

*Observatoire de Paris – Meudon Campus*



# Outline

- **Motivations**
- **High-z Galaxy Luminosity Function**
- **N-body Simulations & Numerical Systematics**
- **Statistical Halo-Galaxy Model**
- **DM Model Constraints & Astrophysical Implications**

# Standard DM paradigm

## Dark Matter:

- Foster matter clustering
- Resides in virialized clumps

## Halos:

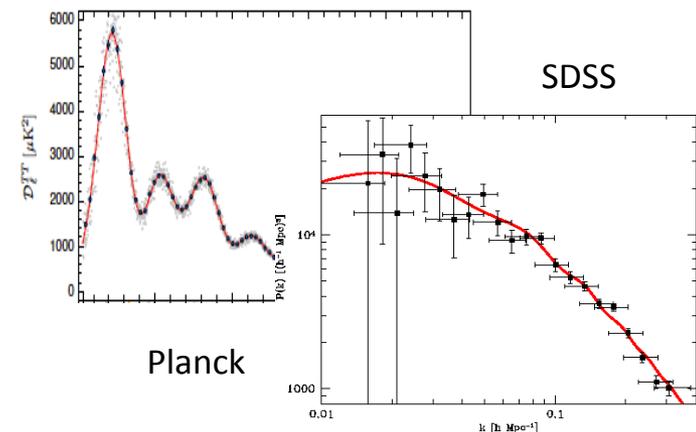
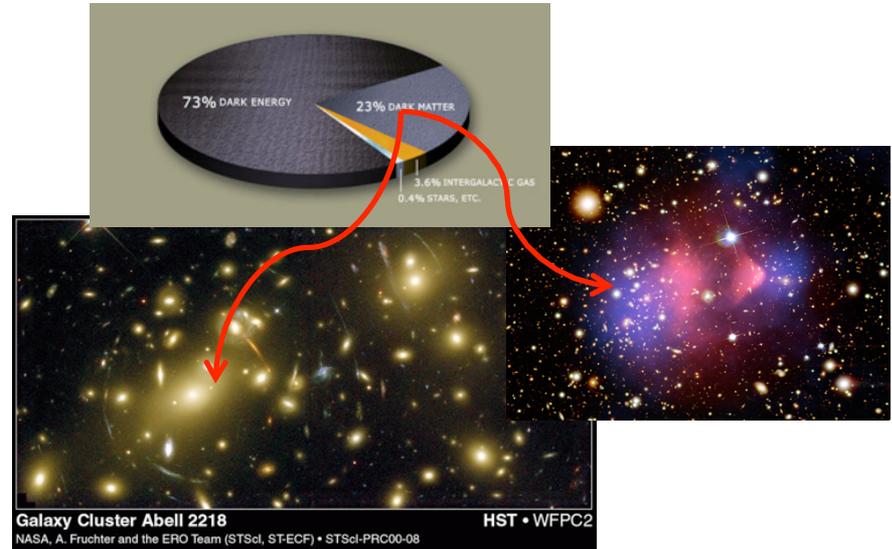
- Building blocks of cosmic structure formation
- Shape baryon distribution

## Large Scales:

- Successful Description CMB spectra
- Clustering of Matter from galaxy surveys
- Bottom-Up Evolution -> Certainly no HDM

## Cold Dark Matter Paradigm

- WIMP miracle



# Small Scales & Beyond CDM

## CDM Anomalies:

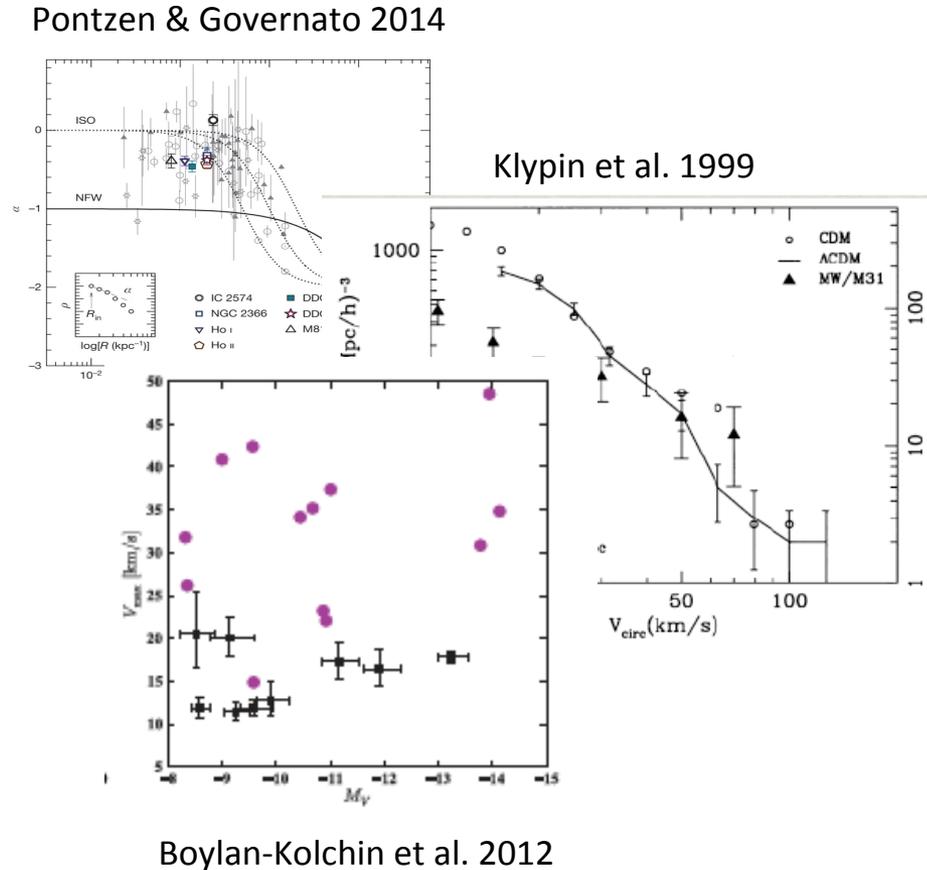
- Core vs Cusp Profiles
- Missing Satellites
- Too-big-too fail

## Non-DM Physics Explanations:

- Baryonic Feedback
- Observational Selection Effects
- Uncertainties of Milky-Way Mass

## DM Direct Searches:

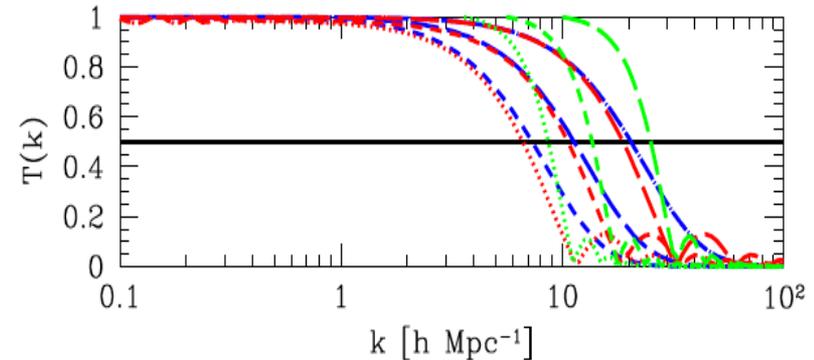
- Negative or Contrasting Results
- No signal at LHC



# Non-CDM Candidates

## Warm Dark Matter

- Thermal Relic ,  $m_{\text{WDM}} \approx \text{keV}$
- Free-Stream  $\leq 100 \text{ kpc}$
- Small-Scale Power Spectrum Cut-off



## Self-Interacting DM Spergel & Steinhardt 2000

- DM scattering cross-section (velocity in/dependent)
- Interaction with radiation

## Ultra-Light Axions and Late-Forming DM

- ULA: Axion field transition from vacuum to matter ( $w=-1 \rightarrow w=0$ )  
see review Marsh (2015)
- LFDM: Before matter/radiation equality, decay of scalar field coupled to radiation ( $w \approx 1/3 \rightarrow w=0$ )  
Fardon, Nelson, Weiner 2006; Das & Weiner 2011

# Non-CDM Observational Consequences

## Warm Dark Matter

- $m_{\text{WDM}} < 0.1$  keV to core the DM profiles e.g. Maccio et al. 2012
- $1.5 < m_{\text{WDM}} [\text{keV}] < 2$  to solve too-big-to-fail Lowell et al. 2012, 2014
- $m_{\text{WDM}} > 3.3$  keV from Lyman- $\alpha$  power spectrum at  $z > 2$  Viel et al. 2012

## Self-Interacting DM

- Lower density sub-halos and core profiles Vogelsberger et al. 2012; Zavala et al. 2013
- Low mass halo abundances unaltered

## Axion DM

- CMB anisotropy and large scale galaxy clustering analysis:  $m > 10^{-24}$  eV

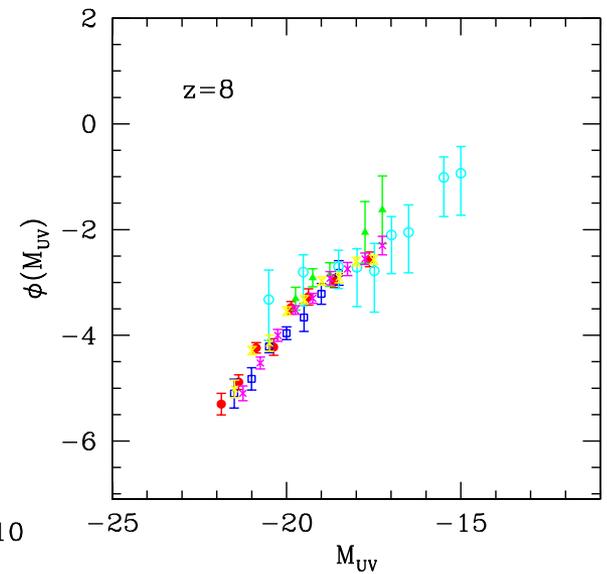
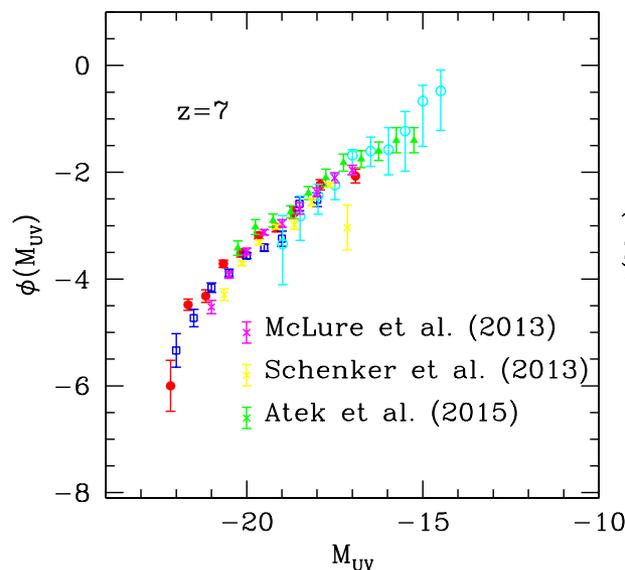
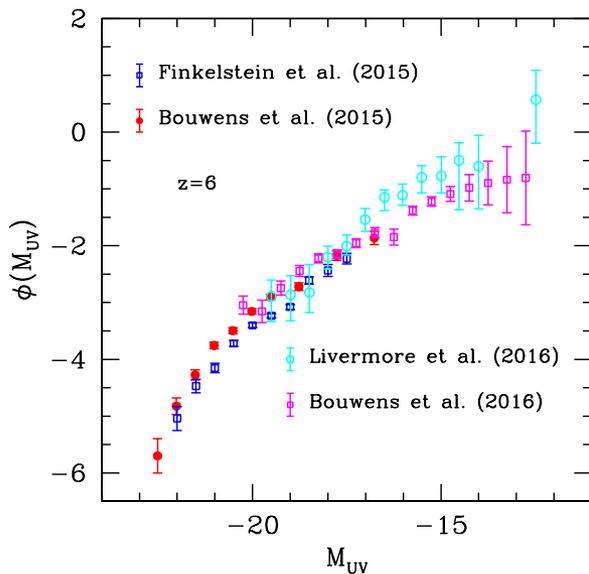
Hlozek et al. 2015

# High-z Galaxy Luminosity Function

- Test for galaxy formation models
- Linked to Cosmic Reionization Scenarios
- Probe low-mass end HMF evolution

## Fast evolving field

- Deep surveys (*Hubble Deep Field*)
- Cluster Gravitational Lensing (*Hubble Frontier Fields*)



# High-z LF vs DM

## Galaxy Number Density:

-  $m_{\text{WDM}} > 1 \text{ keV}$  (CLASH at  $z = 10$ )

Pacucci, Mesinger & Haiman (2013)

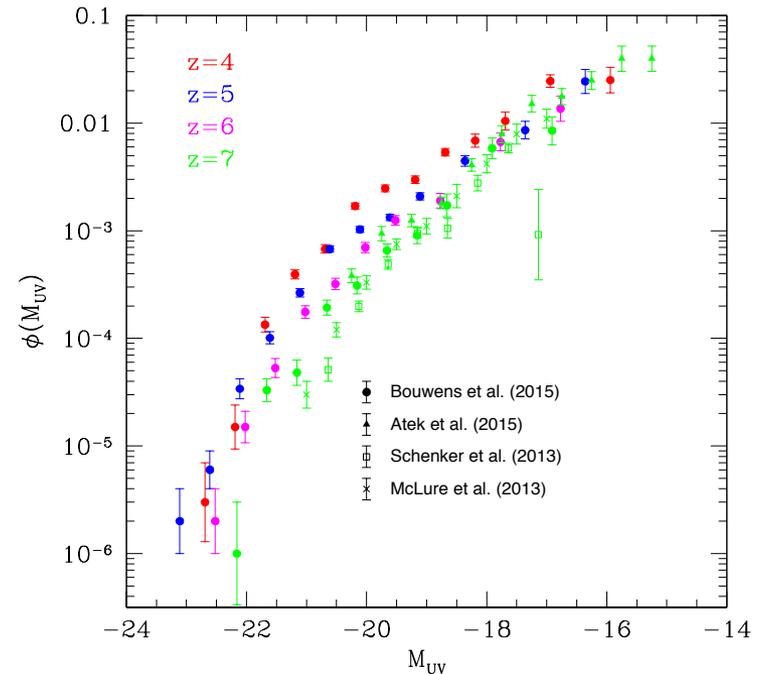
-  $m_{\text{WDM}} > 2.9 \text{ keV}$  (HFF at  $z = 6$ )

Menci et al. (2016)

Does not need to specify  $M_{\text{UV}} - M_h$   
but not free of other caveats

-  $m_{\text{WDM}} > 0.8 \text{ keV}$  (at mag limit  
HUDF data) using HAM of LCDM

Schultz et al. (2014)



## Luminosity Function:

-  $m_{\psi} > 10^{-22} \text{ eV}$  from HAM assuming parametrized  $L_{\text{UV}}(M)$  relation

Schive et al. (2016)

# N-body Simulations

## Models:

- **WDM**:  $m_{\text{WDM}} = 0.7, 1.0, 1.5, 2.0, 2.4$  keV

- **LFDM**:  $z_t = 5, 8, 15 \times 10^5$

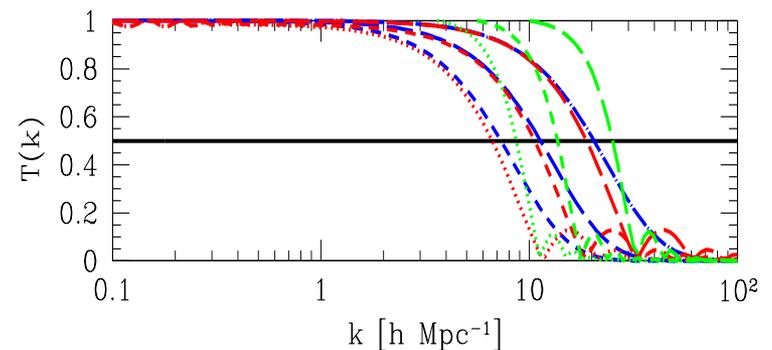
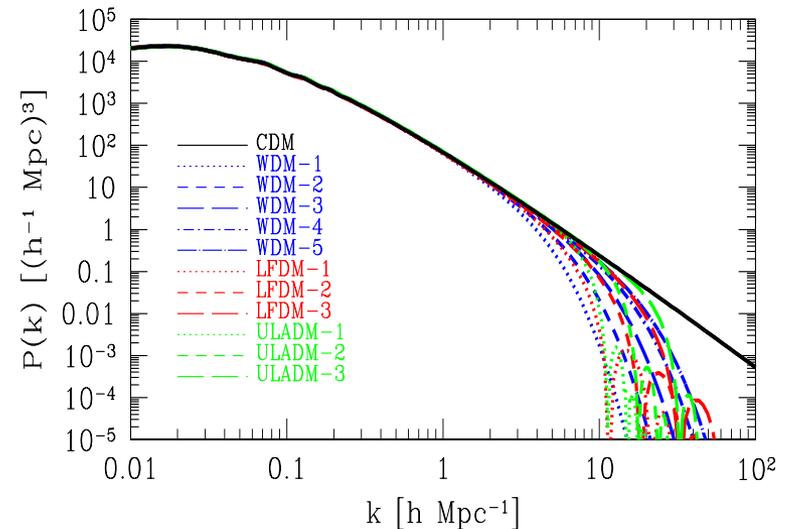
- **ULADM**:  $m_{\text{ULA}} = 1.6, 4.2, 15.4 \times 10^{-22}$  eV

## Runs:

- RAMSES

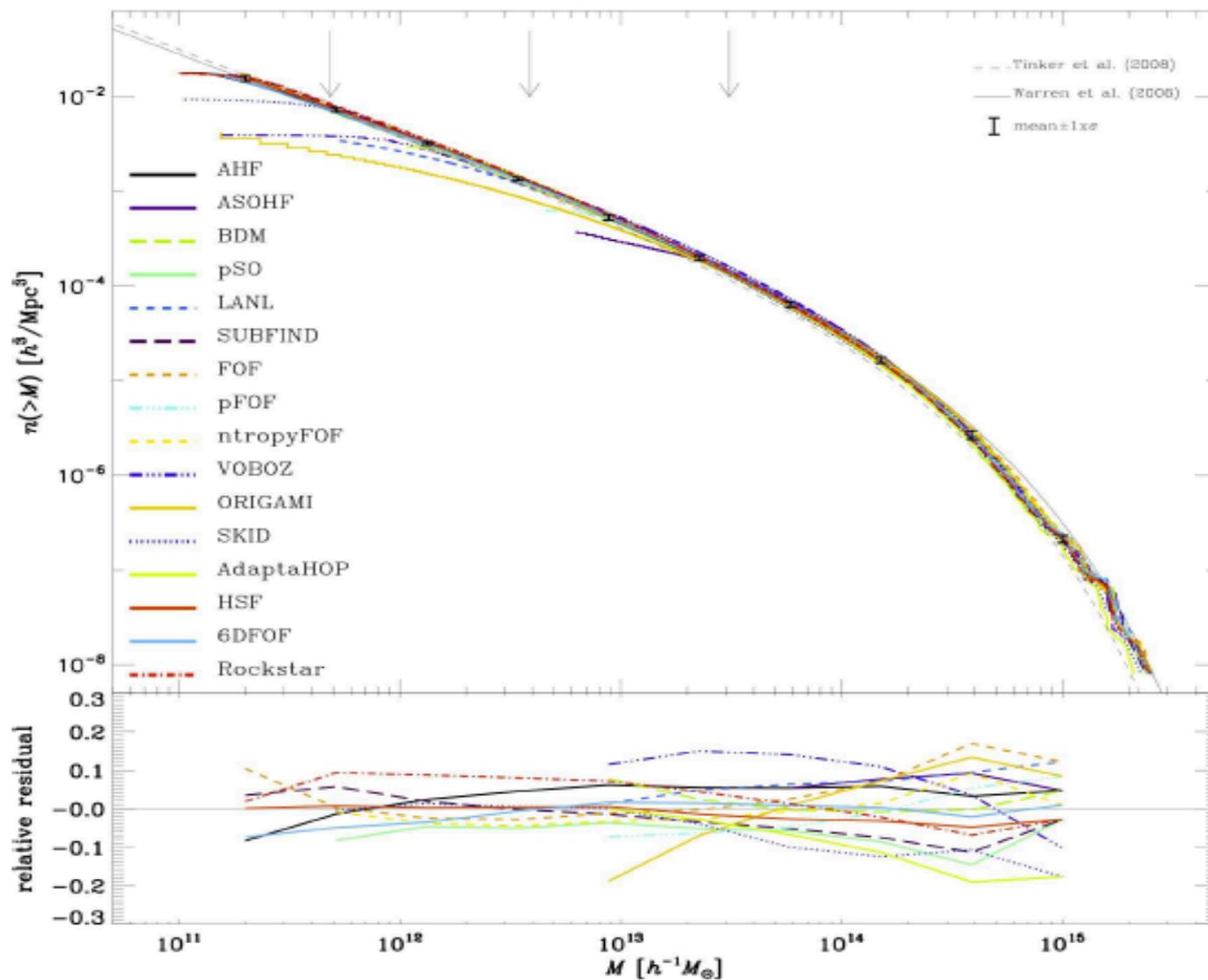
-  $L_{\text{box}} = 27.5$  Mpc/h

-  $N_p = 1024^3$  ( $m_p = 1.6 \times 10^6 M_{\odot}/h$ )



Why can we use N-body?

# pFoF Halo Finder



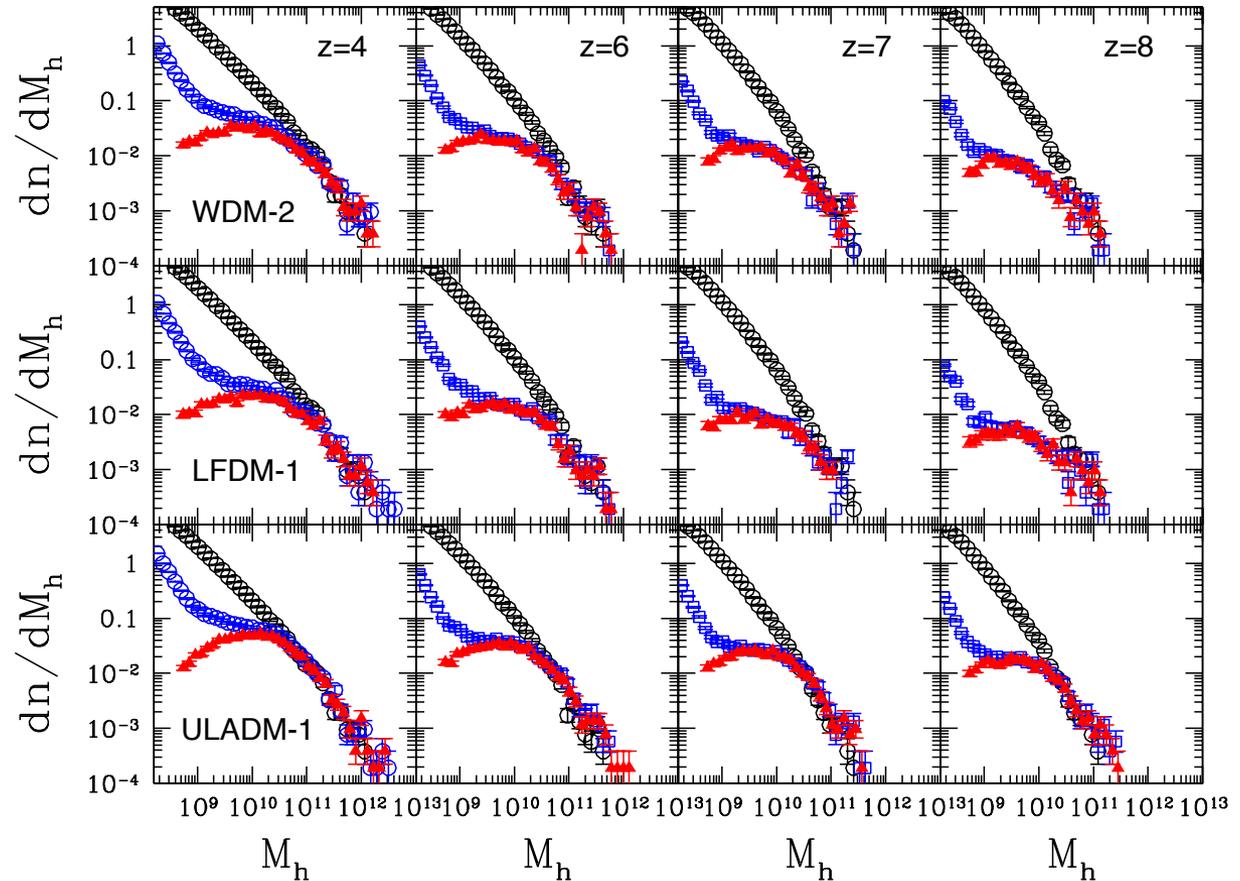
# DM model halo catalogs

## Initial Catalogs:

- FoF halos
- $N_h > 100$
- lots of spurious halos

## Selections:

- $N_h > 300$
- only halos with  $0 < \eta = 2K/E < 1.5$



# Artificial Halos

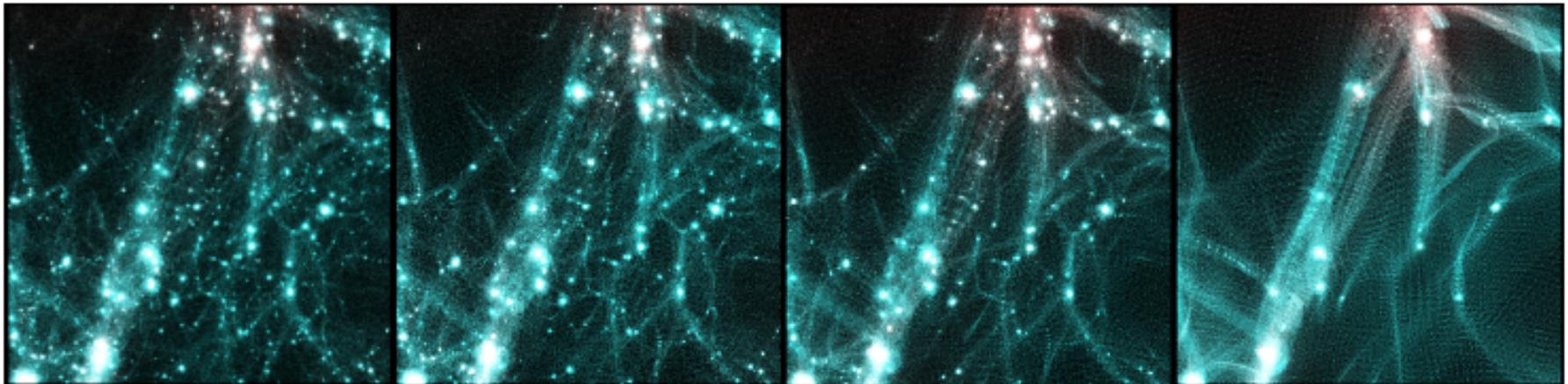
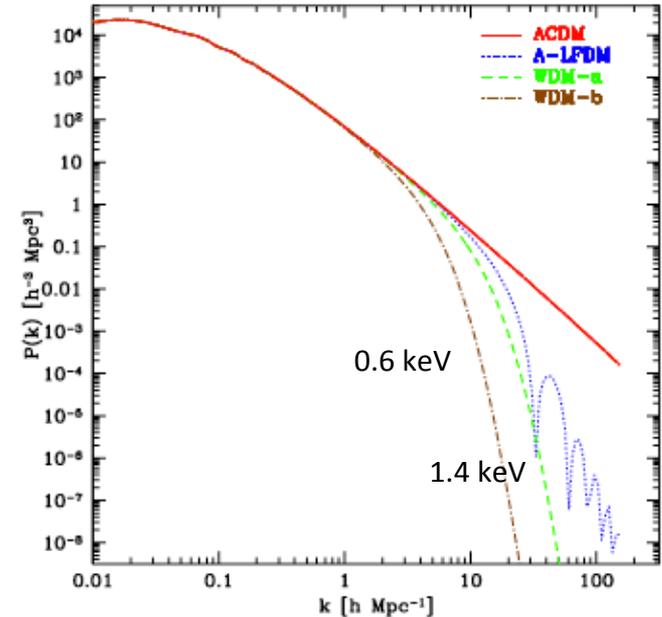
## Discretization Effect

- Sampling Poisson Noise ( $k > k_{\text{cut-off}}$ )
- Spurious Numerical Halos

Gotz & Sommer-Larsen 2002, 2003; Wang & White 2007

## Example

RAMSES -  $N_p = 512^3$  -  $m_p \sim 10^7 M_{\text{sun}} h^{-1}$   
AMR -  $L_{\text{box}} = 27.5 \text{ Mpc } h^{-1}$  -  $dx_{\text{coarse}} \sim 54 \text{ kpc } h^{-1}$



# Spurious Halo Contamination

## Halo Mass Function

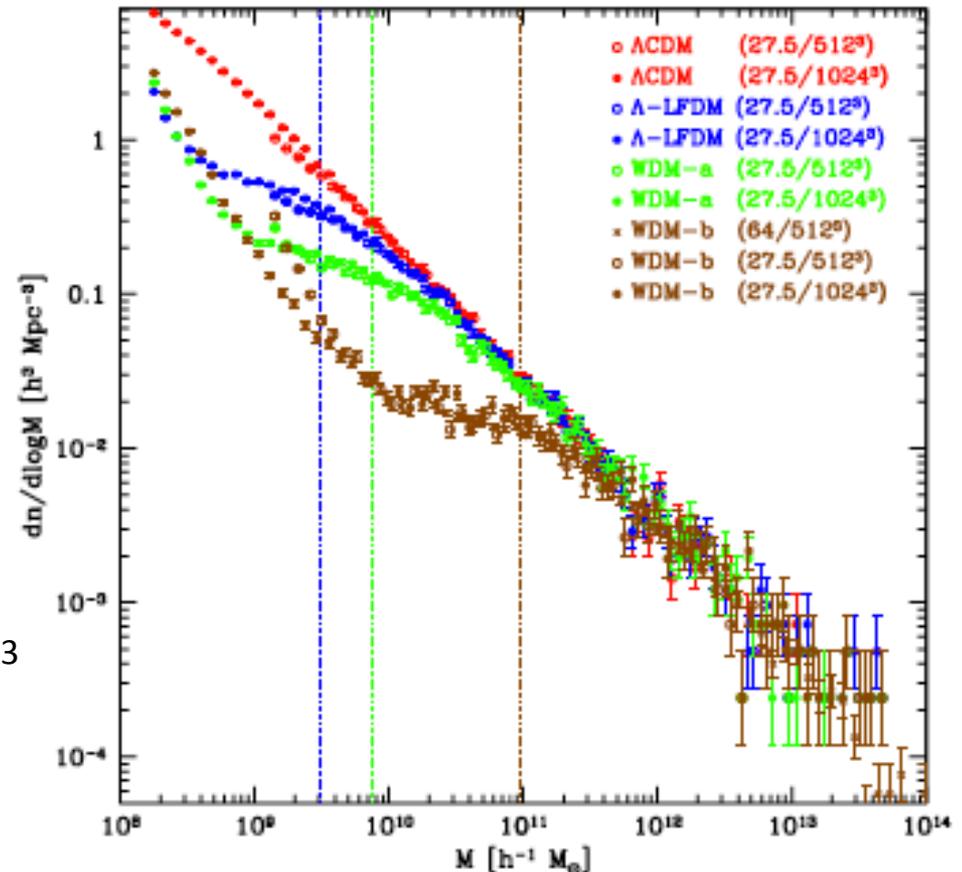
- $N_{\text{h-particles}} > 100$
- Upturn at  $M < M_*$
- Simulation Dependent Slope

## Proposed Cures

- Mass Cut:  $M_{\text{min}} = 10.1 \rho d k_p^{-2}$   
Wang & White 2007
- Select Unflatten Proto-Halos  
in Initial Lagrangian Patch &  
Apply Mass Cut Lowell et al. 2012
- Visual Inspection Angulo, Hahn, Abel 2013
- Tessellation 6-d phase-space  
folding (reduce but doesn't solve)

Hahn, Abel, Kaehler 2013

- $N_p = 1024^3$
- $L_{\text{box}} = 27.5 \text{ Mpc } h^{-1}$
- $m_p \sim 10^6 M_{\text{sun}} h^{-1}$
- $dx_{\text{coarse}} \sim 26 \text{ kpc } h^{-1}$



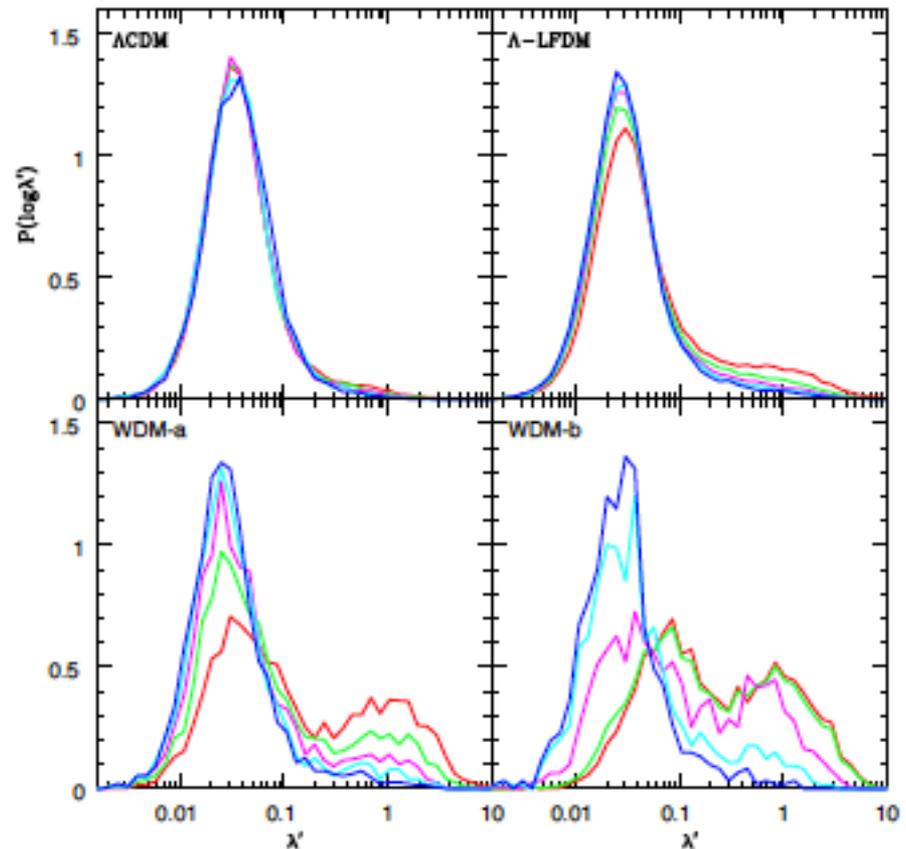
Agarwal & Corasaniti 2015

# Structural Properties of Halos

Agarwal & Corasaniti 2015

## Halo Spin

- Spin parameter  $\lambda' = \frac{J}{\sqrt{2}MVR}$
- $V = v(GM/R)$
- 8 bins:  $4 < M[10^9 M_{\text{sun}} h^{-1}] < 8$
- CDM: log-normal & mass independent
- non-CDM: deviations from lognormality/bimodality and mass dependent
- spurious halos have large spins



# Structural Properties of Halos

## Halo Shape

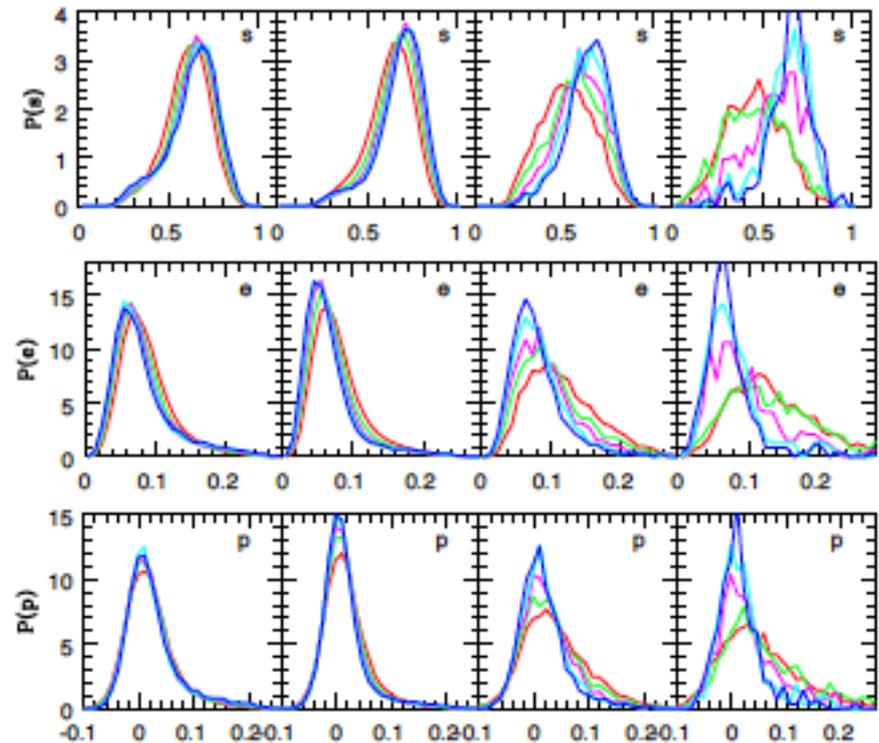
- Symmetric Mass  
Distribution Tensor

$$M_{\alpha\beta} = \frac{m_p}{M} \sum_{i=1}^{N_h} (r_{\alpha,i} - r_{\alpha,c})(r_{\beta,i} - r_{\beta,c})$$

- sphericity, ellipticity &  
prolatness

- CDM: mass independent &  
elliptical halos

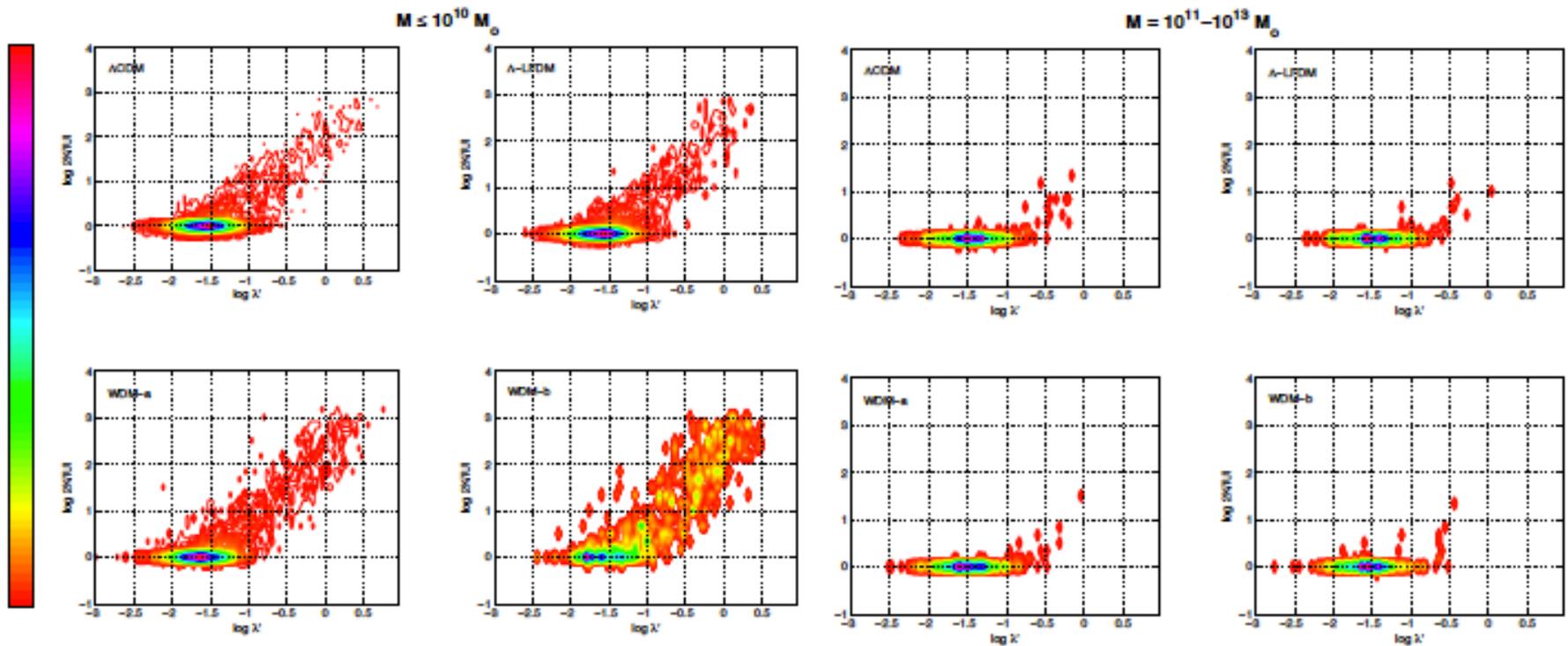
- non-CDM: mass dependent  
& highly non-spherical  
(elliptical & prolate, i.e.  
alignment with filaments)



# Halo Dynamical State

## Virial Condition

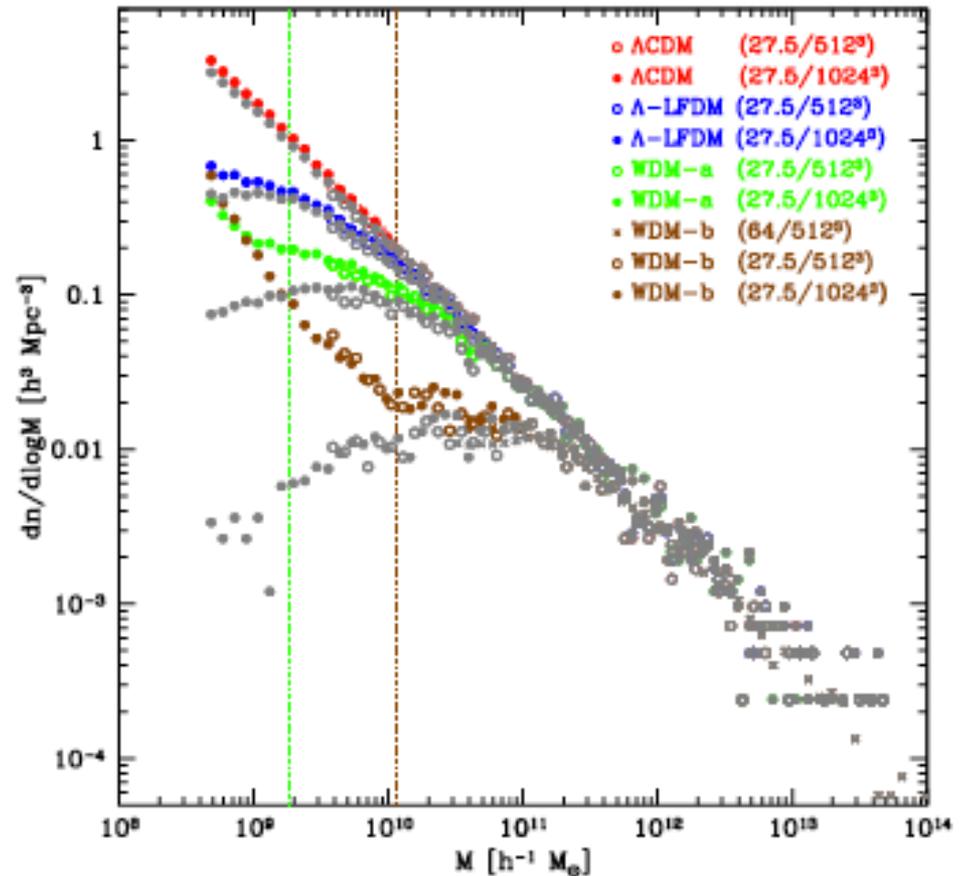
- proxy:  $\eta = 2 K/|E|$
- correlation  $\lambda' - \eta$  for  $\eta > 1$



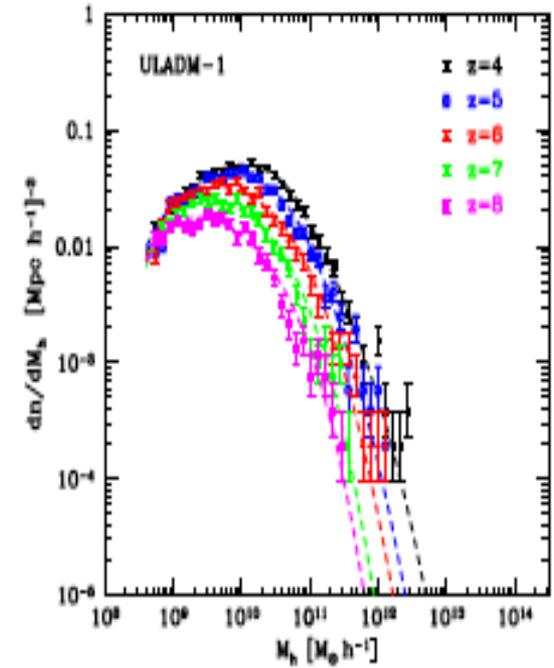
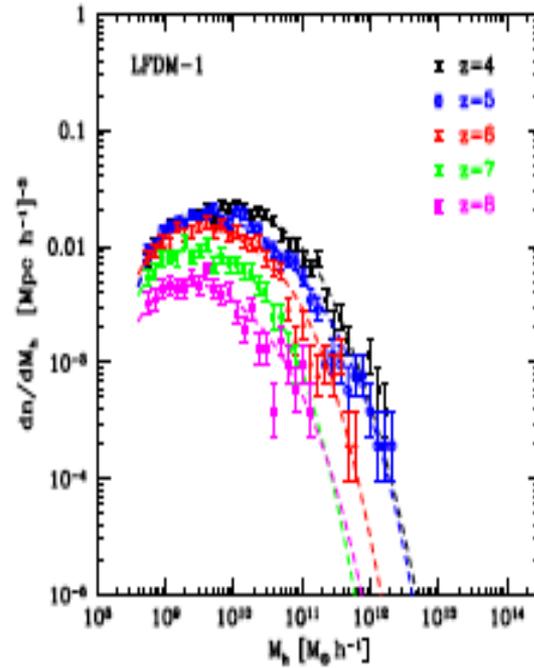
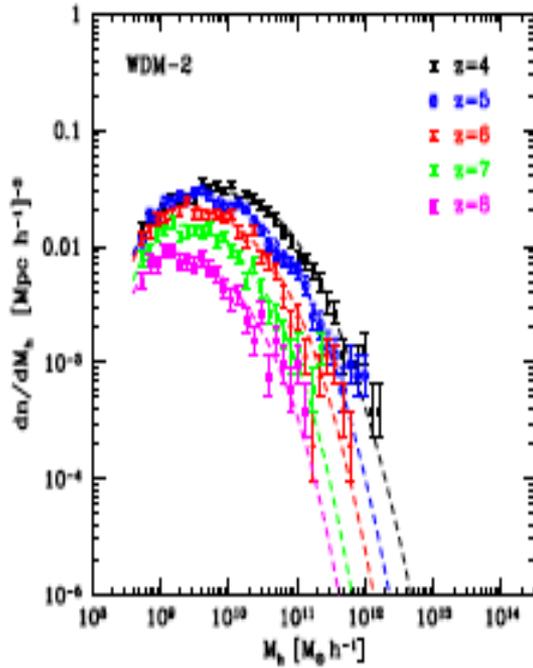
# Virial State Selection

## Removing Spurious Halos

- $0 < \eta = 2K/|E| < 1.5$
- recover halo triaxial distribution
- recover spin log-normality
- recover suppressed mass function at low mass (mass resolution convergence)
- spurious halos still present with simple mass-cut at  $M_{\min}$
- mass range larger than mass cut



# Evolution of HMF in NDM models



## Calibrated Analytical Formula

$$\frac{dn}{dM_h} = 10^{\alpha + \beta \frac{M_*}{M_h}} \left( 1 - e^{-\frac{M_h}{M_*}} \right)^\gamma \frac{dn}{dM_h} \Big|_{\text{CDM}}$$

# Predicting LF at $z > 5$

## From HMF to GLF:

- Must assume:  $M_{UV} - M_h$

- Infer from ensemble averaging using HAM (e.g. Mashian et al. 2016)

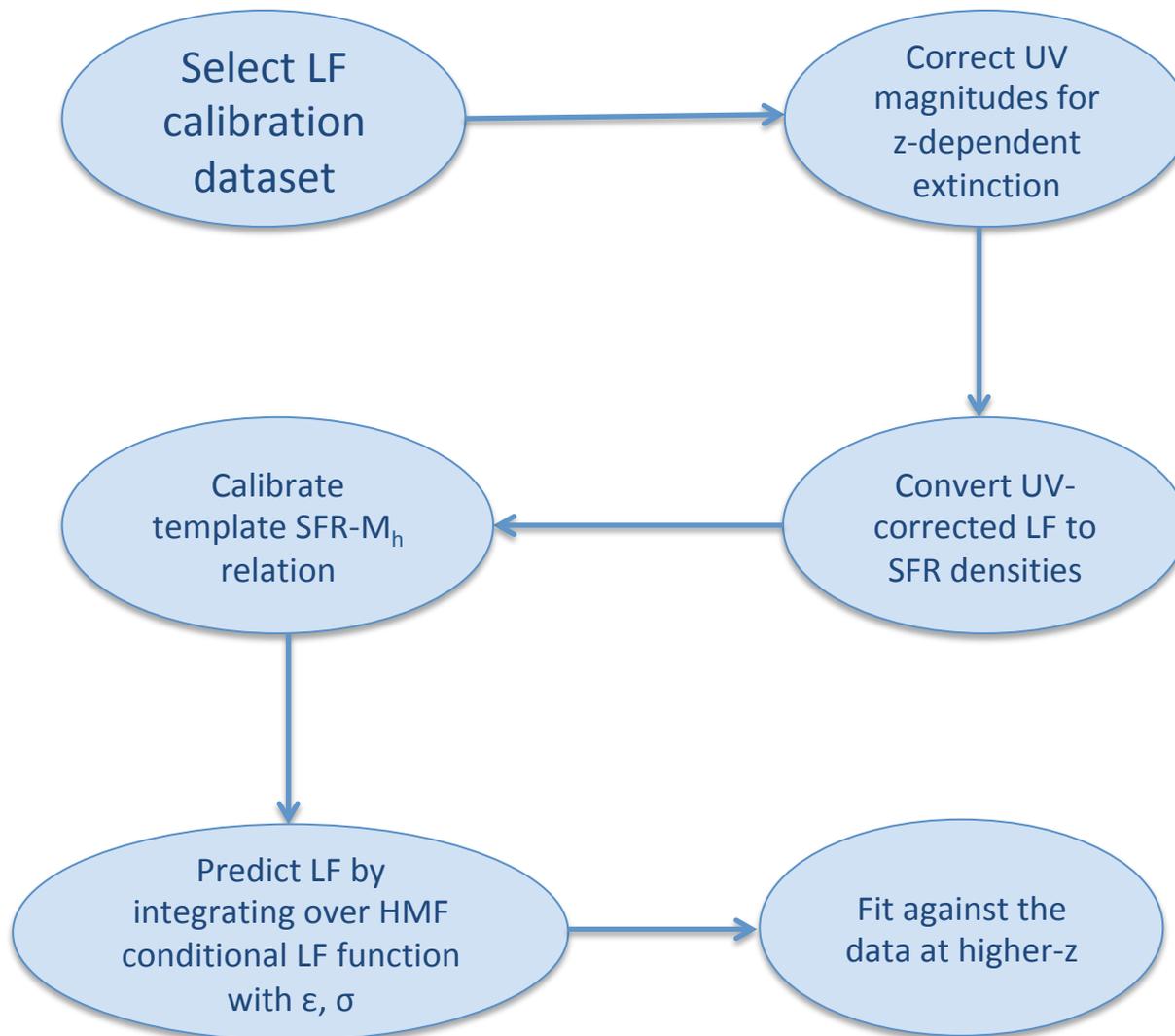
- Calibrate SFR(M) model using single dataset at given  $z$  (e.g. Mason et al. 2015)

- Parametrize and derive parameters from LF data fitting (e.g. Schive et al. 2016)

## What can we do to..

- Account for  $z$ -evolution of  $M_{UV} - M_h$  due dust extinction
- Infer such relation from the data
- Deduce information on galaxy properties e.g. SFR( $M_h$ )
- Learn about DM model dependencies

# Empirical Approach



# HAM & SFR

## UV-Luminosity to SFR

- Account for dust extinction
- Convert  $M_{UV}$  corrected to SFR (Kennicutt relation)
- Derive SFR-density functions (see Mashian, Oesch, Loeb 2015 for LCDM)

## Extinction Correction

$$\langle A_{UV} \rangle = 4.43 + 0.79 \ln 10 \sigma_{\beta}^2 + 1.99 \langle \beta \rangle$$

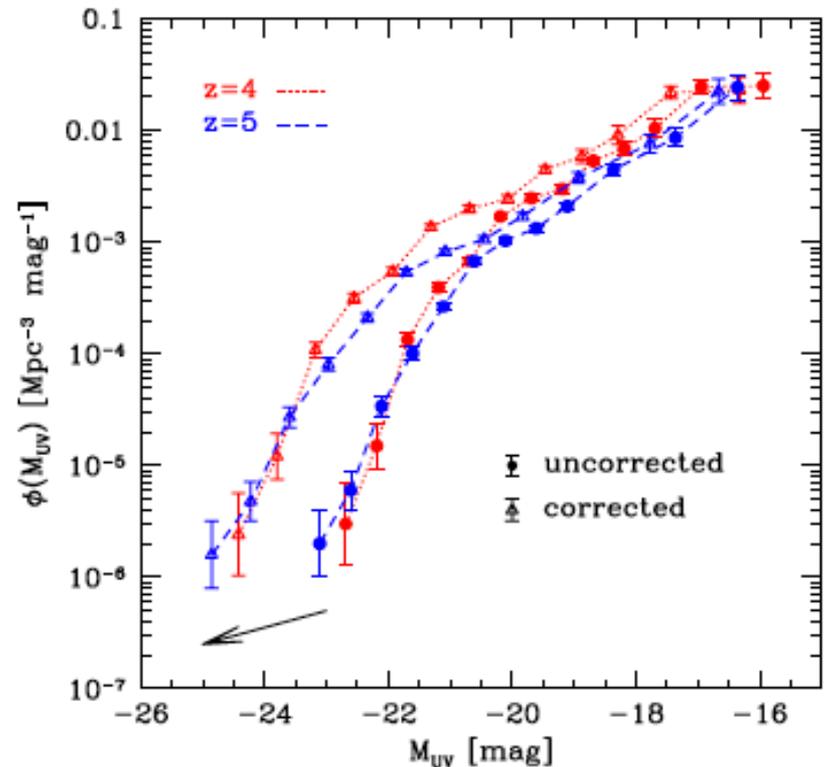
Meurer et al. (1999)

$$\langle \beta(M_{UV}, z) \rangle = \begin{cases} [\beta_{M_0}(z) - C] e^{\beta'(z) \frac{M_{UV} - M_0}{\beta_{M_0}(z) - C}} & M_{UV} \geq M_0 \\ \beta'(z) [M_{UV} - M_0] + \beta_{M_0}(z) & \end{cases}$$

Tacchella et al. (2013), Mason, Trenti & Treu (2015)

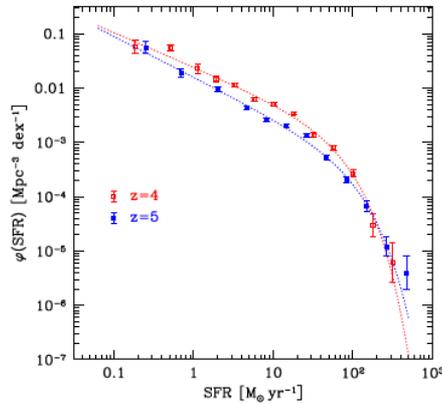
- Changes UV-mag bin size
- Shift toward higher luminosities

Smit et al. (2012)



# Template Function SFR – $M_h$ relation

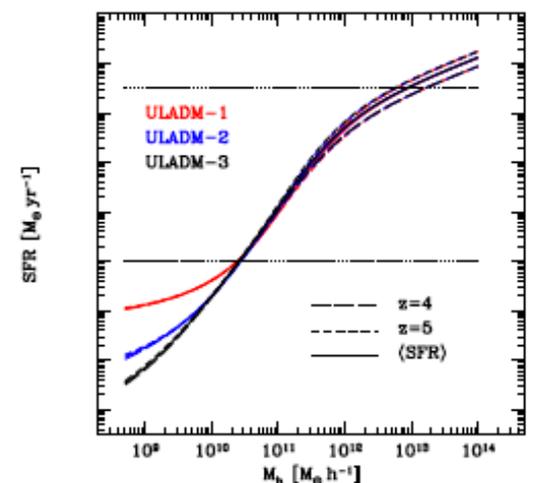
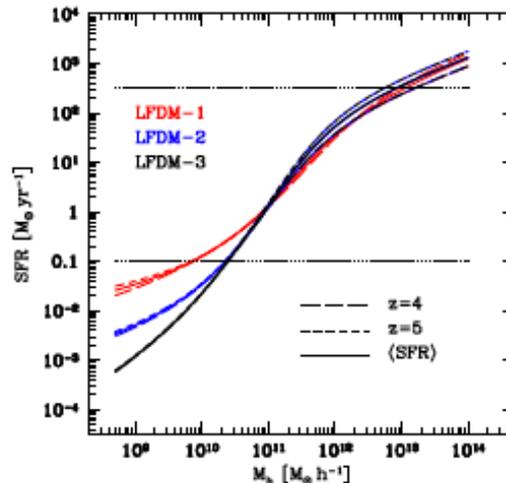
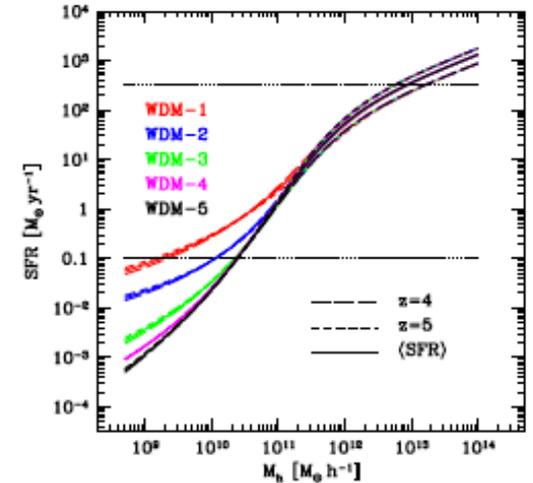
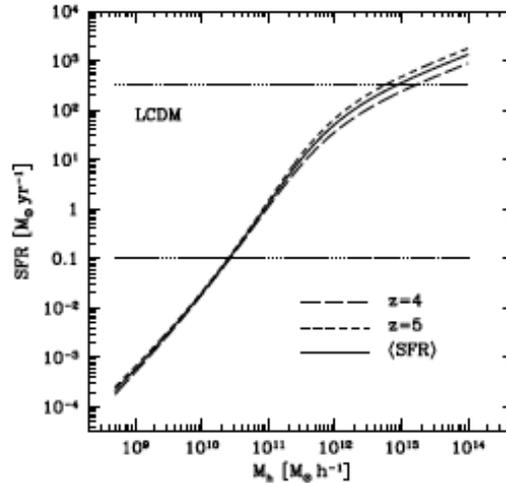
## SFR density function



- Fit  $\Phi(\text{SFR})$  with Press-Schechter function

- HAM

$$\int_M d\tilde{M}_h \left( \frac{dn}{d\tilde{M}_h} \right)_{NDM} = \int_{SFR} dSFR' \phi(SFR')$$



# Modeling Luminosity Function at $z > 5$

## Average amplitude and intrinsic scatter SFR- $M_h$ relation

- Compute

$$\phi(SFR, z) = \frac{1}{\sigma_{\text{int}}^2 \sqrt{2\pi} SFR} \int dM_h \frac{dn}{dM_h}(M_h, z) e^{-\frac{\log_{10}^2[SFR/(\epsilon \langle SFR(M_h) \rangle)]}{2\sigma_{\text{int}}^2}}$$

- Convert to UV luminosities

- Add extinction effect

- Estimate  $\Phi(M_{\text{UV}})$

- Fit against the data  $\epsilon, \sigma_{\text{int}}$

# High-z LF data

## Bright-end side of LF

Bouwens et al. (2015):

- 10,000 galaxies HST data
- subsample for dust model

## Faint-end Slope

HFF observations:

Atek et al. (2015)

Livermore et al. (2016)

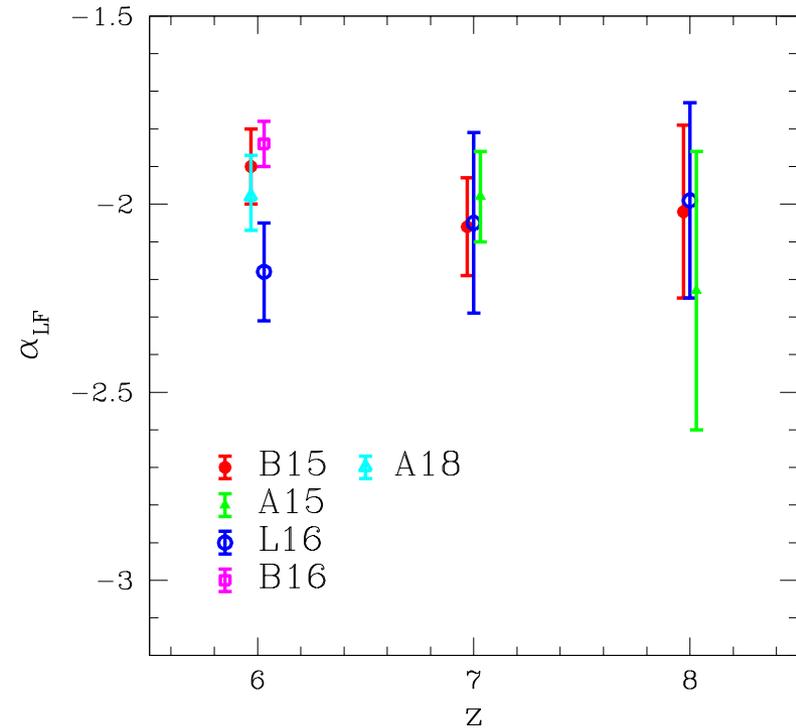
Bouwens et al. (2016)

## Our selected dataset

26 points at  $z=6$  (B15+B16)

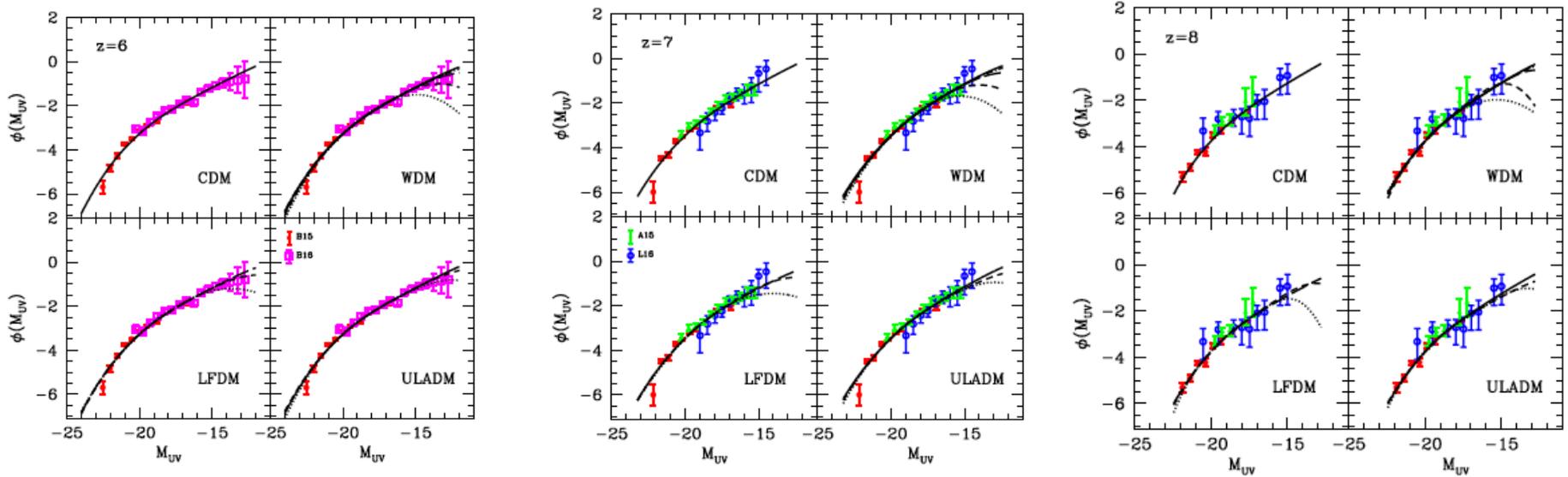
31 points at  $z=7$  (B15+A15+L16)

22 points at  $z=8$  (B15+A15+L16)



Sensitivity to lens model  
magnification systematics

# DM models goodness-of-fit



Model	$\log_{10} \varepsilon_{\text{SFR}}^{z=6}$	$\log_{10} \sigma_{\text{SFR}}^{z=6}$	$\chi^2_{z=6}$	$\log_{10} \varepsilon_{\text{SFR}}^{z=7}$	$\log_{10} \sigma_{\text{SFR}}^{z=7}$	$\chi^2_{z=7}$	$\log_{10} \varepsilon_{\text{SFR}}^{z=8}$	$\log_{10} \sigma_{\text{SFR}}^{z=8}$	$\chi^2_{z=8}$	$\chi^2_{\text{tot}}$
<b>CDM</b>	-0.80	-0.23	21.5	-0.52	-0.23	27.6	-0.18	-0.39	15.8	<b>64.9</b>
WDM-1	-0.78	-0.26	87.7	-0.60	-0.25	57.4	-0.58	-0.24	23.8	168.9
WDM-2	-0.79	-0.25	22.4	-0.53	-0.24	31.6	-0.12	-0.47	17.5	71.5
<b>WDM-3</b>	-0.83	-0.22	20.5	-0.54	-0.23	28.1	-0.23	-0.37	15.7	<b>64.3</b>
WDM-4	-0.90	-0.20	21.7	-0.60	-0.21	27.8	-0.28	-0.35	15.9	65.4
WDM-5	-0.85	-0.21	22.0	-0.60	-0.21	27.1	-0.26	-0.36	15.8	64.9
LFDM-1	-0.92	-0.17	37.2	-0.73	-0.14	45.0	-0.29	-0.50	16.6	98.8
<b>LFDM-2</b>	-0.83	-0.23	20.1	-0.53	-0.23	28.3	-0.22	-0.38	15.6	<b>64.0</b>
LFDM-3	-0.85	-0.22	21.7	-0.73	-0.20	30.5	-0.49	-0.29	16.2	68.4
ULADM-1	-0.91	-0.24	21.3	-0.69	-0.24	33.6	-0.48	-0.36	14.9	69.8
ULADM-2	-0.89	-0.20	21.5	-0.78	-0.19	31.4	-0.59	-0.26	16.5	69.4
<b>ULADM-3</b>	-0.81	-0.23	21.9	-0.60	-0.21	27.3	-0.29	-0.34	15.9	<b>65.1</b>

**WDM:**

$$m_{\text{WDM}} \geq 1.5 \text{ keV}$$

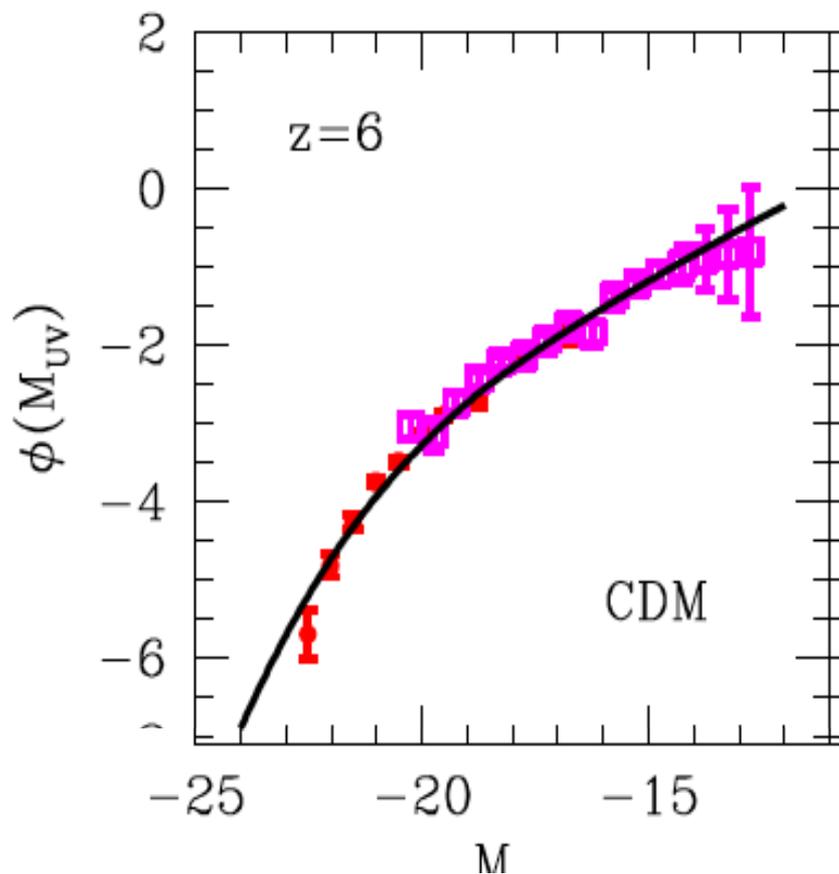
**LFDM:**

$$z_t \geq 8 \times 10^5$$

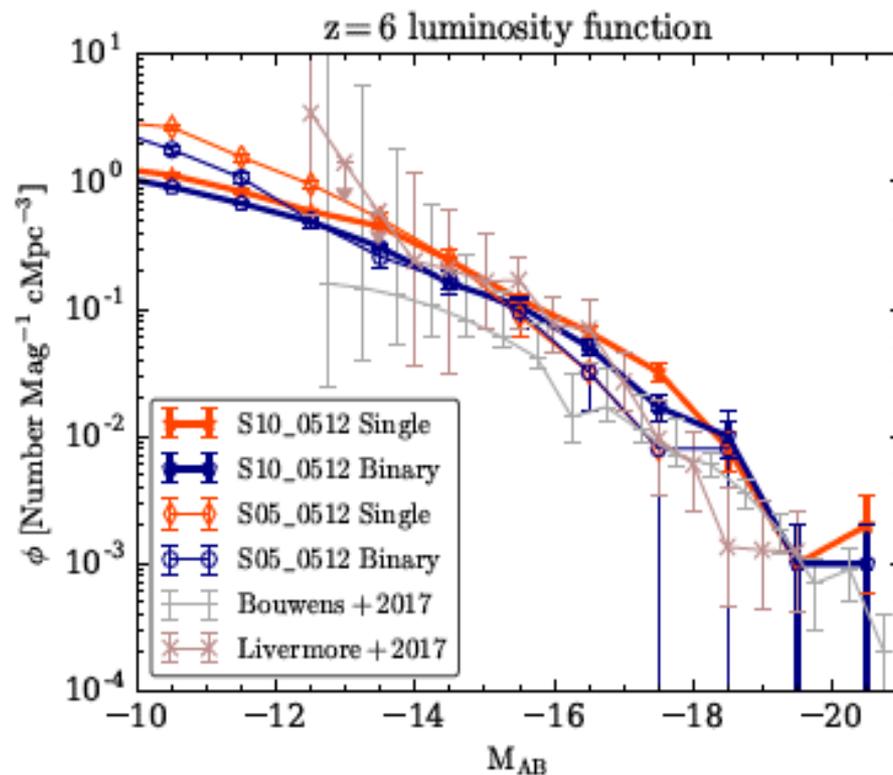
**ULADM:**

$$m_{\text{ULA}} \geq 1.6 \times 10^{-22} \text{ eV}$$

# ΛCDM comparison to Sphinx Simulations



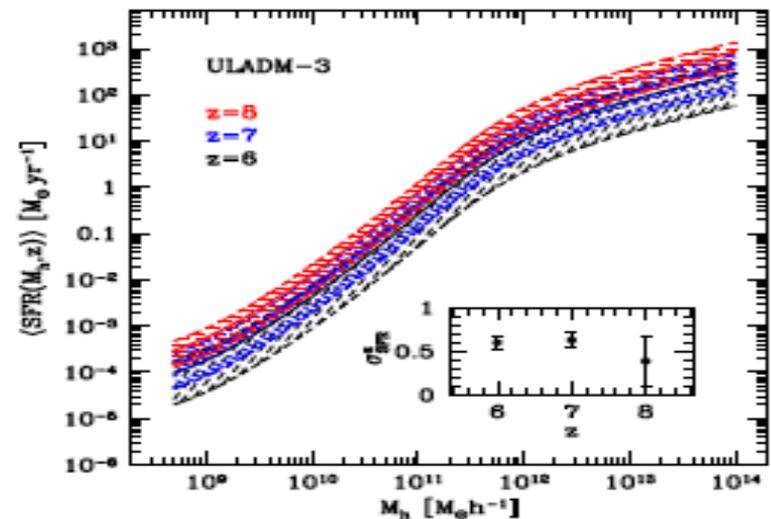
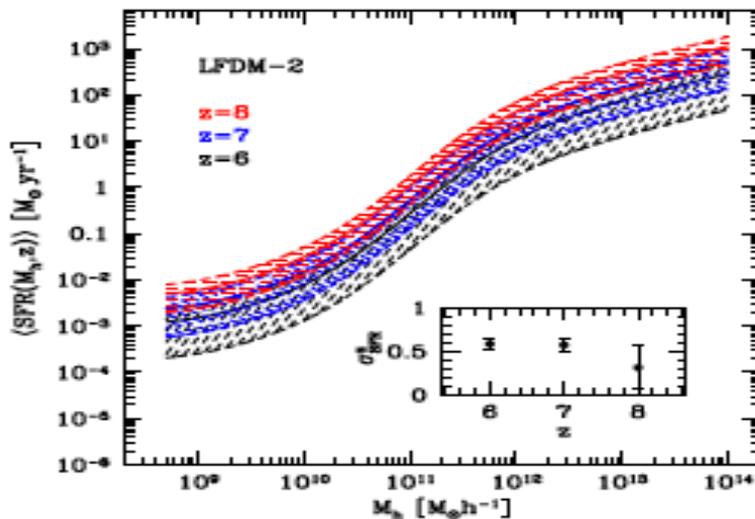
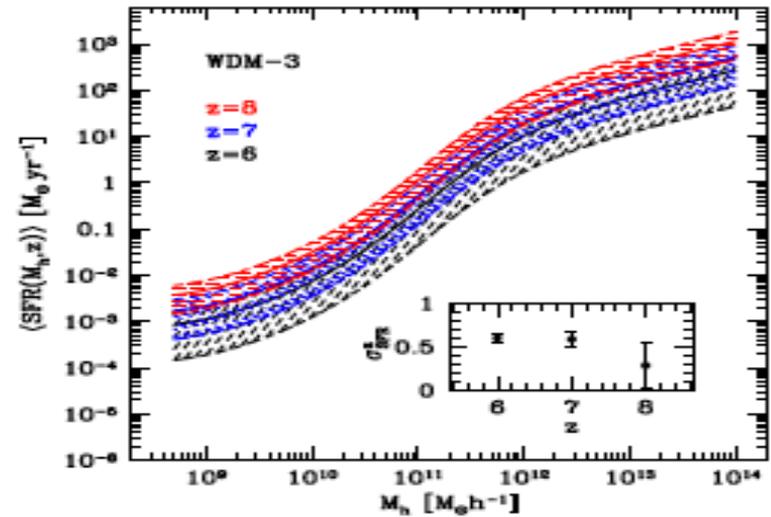
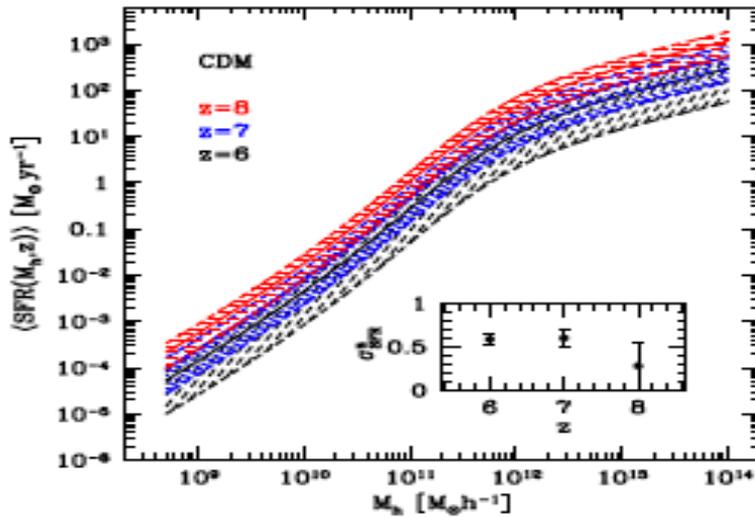
Radiation-Hydrodynamics  
Simulation Suite



Rosdahl et al. (2018)

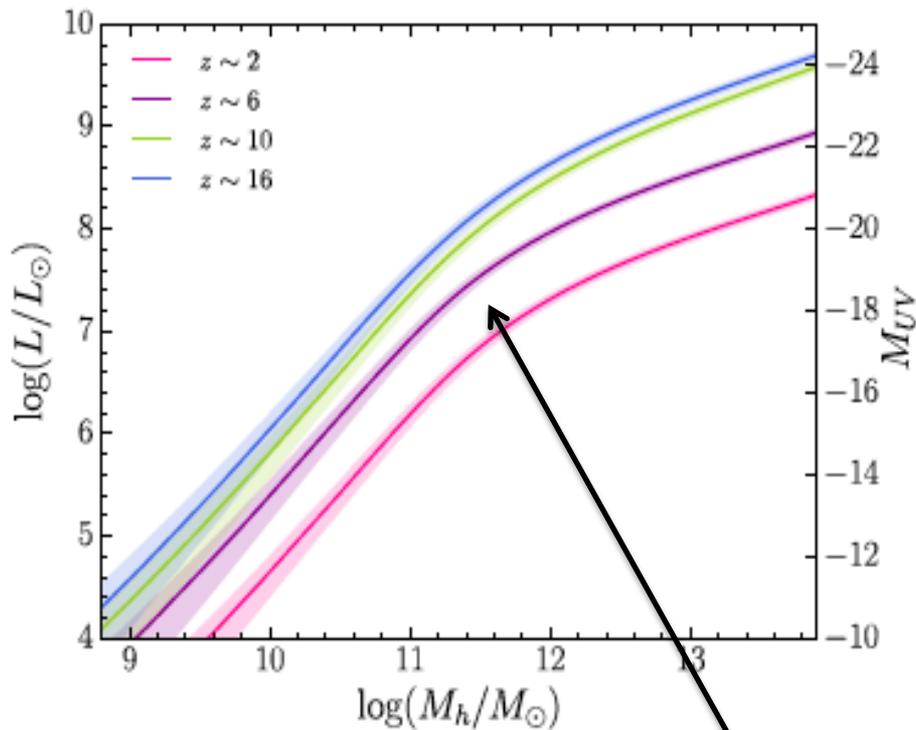
- Consistent LF “predictions”
- Disagreement on the choice of x-axis orientation

# Star-Formation Rate Halo Mass Relation

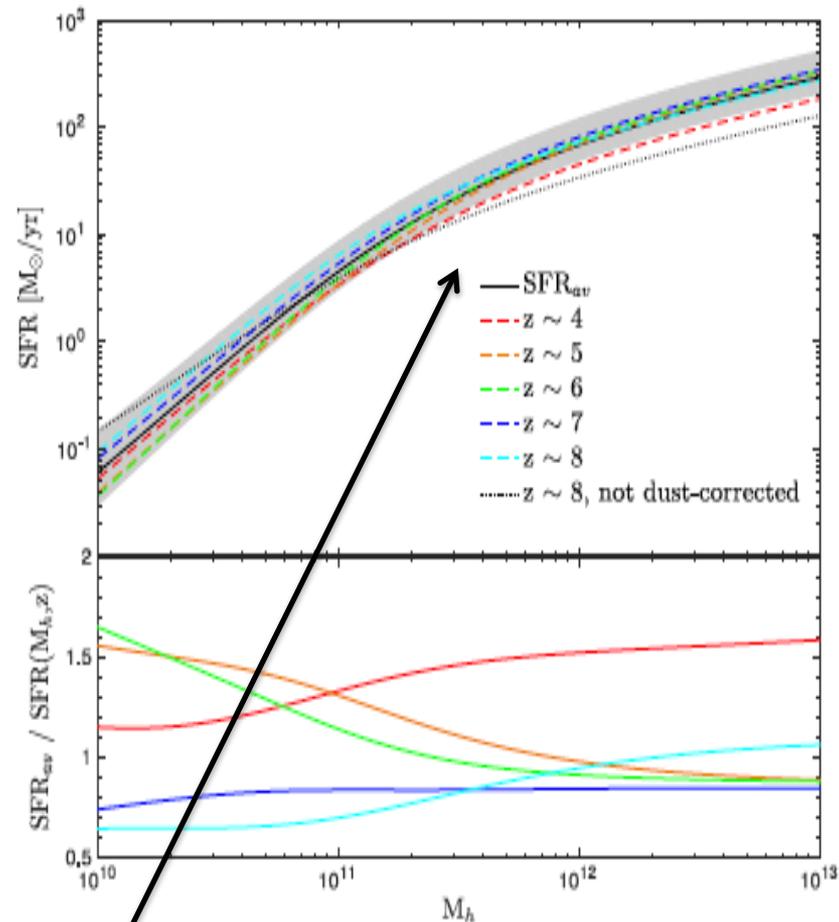


# LCDM “prediction” consistency check

## Different Empirical Models



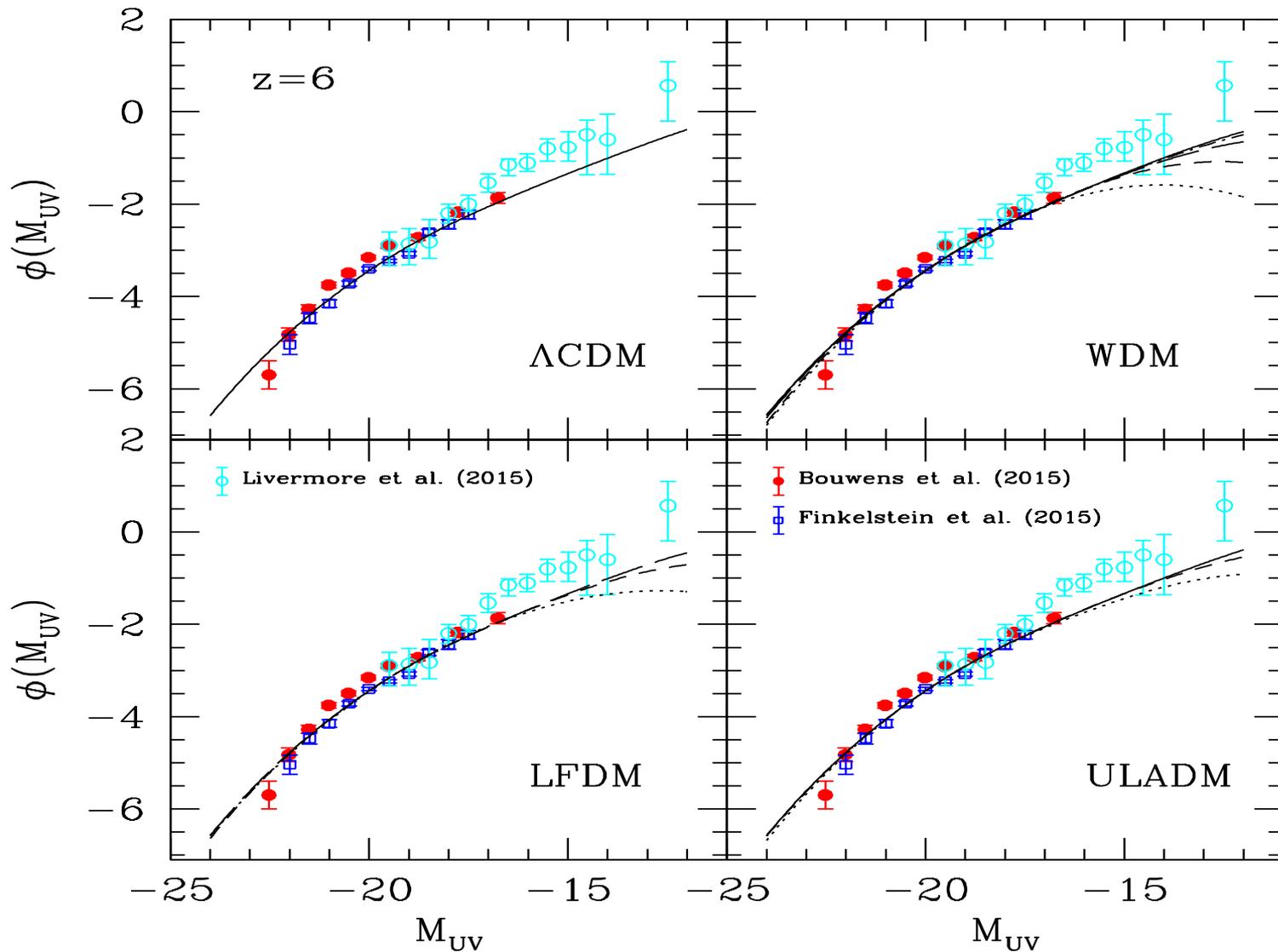
Mason, Trenti & Treu (2015)



Mashian, Oesch, Loeb (2015)

$M_* \sim 2 \cdot 10^{11} M_{\text{Sun}}$

# If you like tensions: results from L16 at $z=6$



# Cosmic Reionization Model

Kuhlen & Faucher-Giguere (2012)

## Ionizing Emissivity Model

$$\dot{n}_{ion}^{com} = f_{esc} \int_{M_{lim}}^{\infty} dM_{UV} \phi(M_{UV}) \gamma_{ion}(M_{UV})$$

## Ionizing Luminosity

$$\gamma_{ion}(M_{UV}) = 2 \cdot 10^{25} \cdot 10^{0.4(51.63 - M_{UV})} \xi_{ion} \quad [s^{-1}]$$

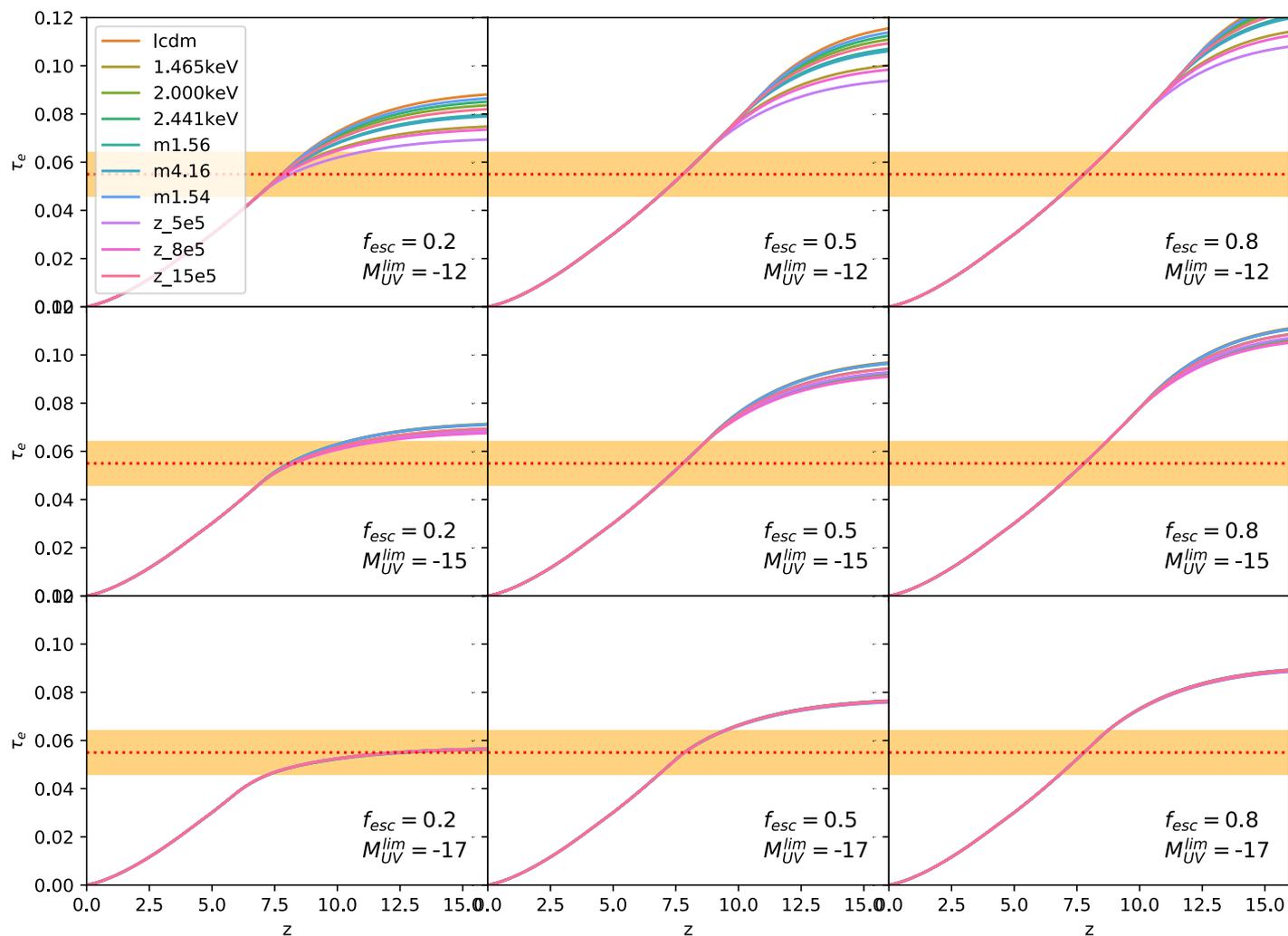
## Volume Average Ionized Hydrogen Fraction

$$\frac{dQ_{HII}}{dt} = \frac{\dot{n}_{ion}^{com}}{\bar{n}_H} - \frac{Q_{HII}}{\bar{t}_{rec}} \quad \bar{t}_{rec} = 0.93 \left( \frac{C_{HII}}{3} \right)^{-1} \left( \frac{T_0}{2 \times 10^4 K} \right)^{0.7} \left( \frac{1+z}{7} \right)^{-3} \quad [\text{Gyr}]$$

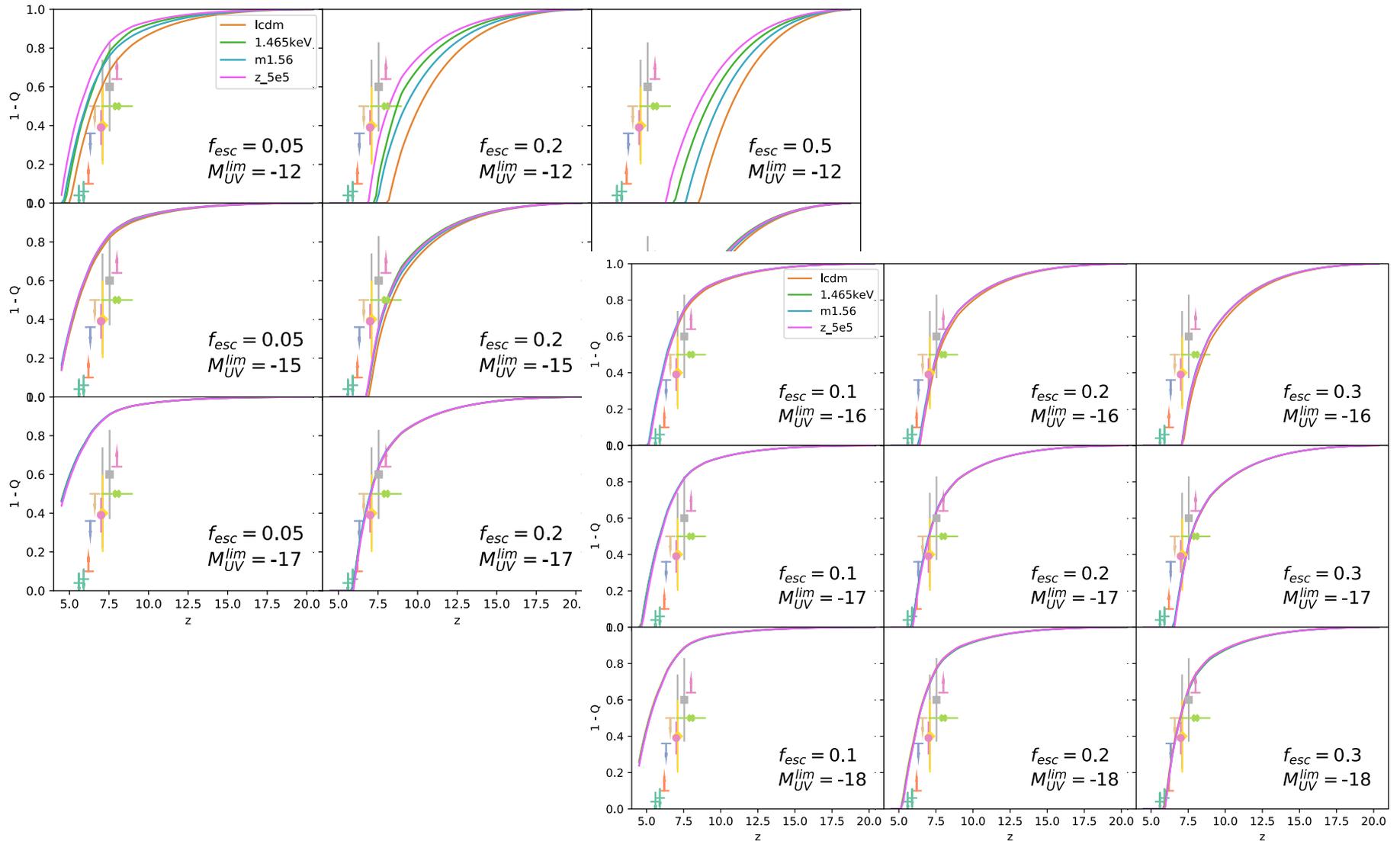
## Optical Depth

$$\tau_e = \int_0^{\infty} dz \frac{c(1+z)^2}{H(z)} Q_{HII}(z) \sigma_T \bar{n}_H (1 + \eta Y/4X),$$

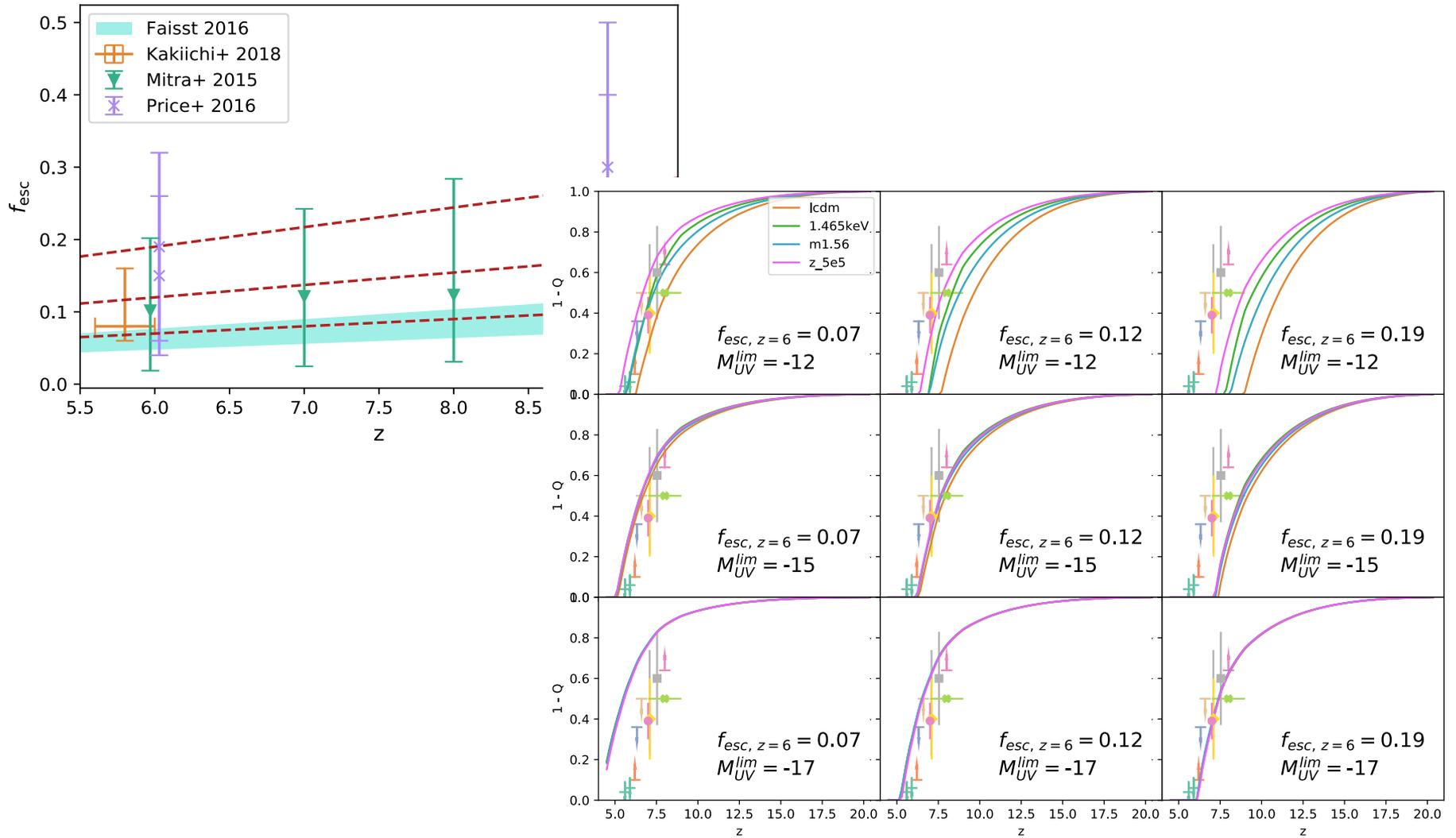
# Planck Limits: $f_{esc}$ vs $M_{lim}^{DM}$ Model Degeneracy



# Ionized Hydrogen Fraction



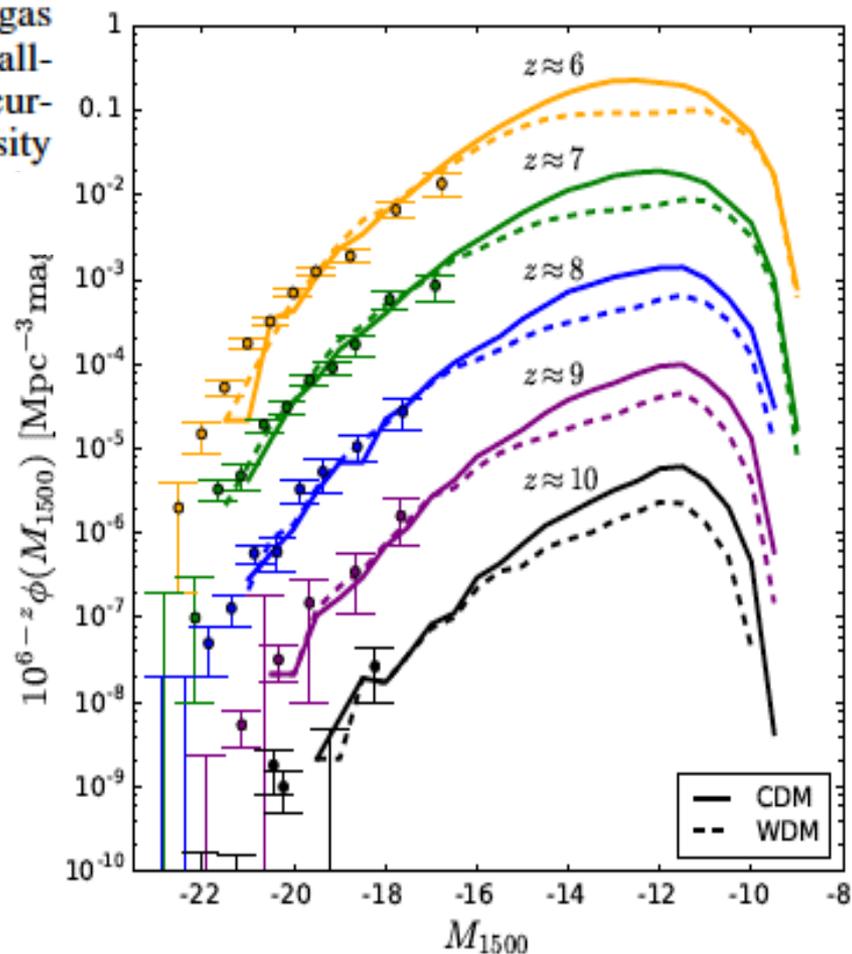
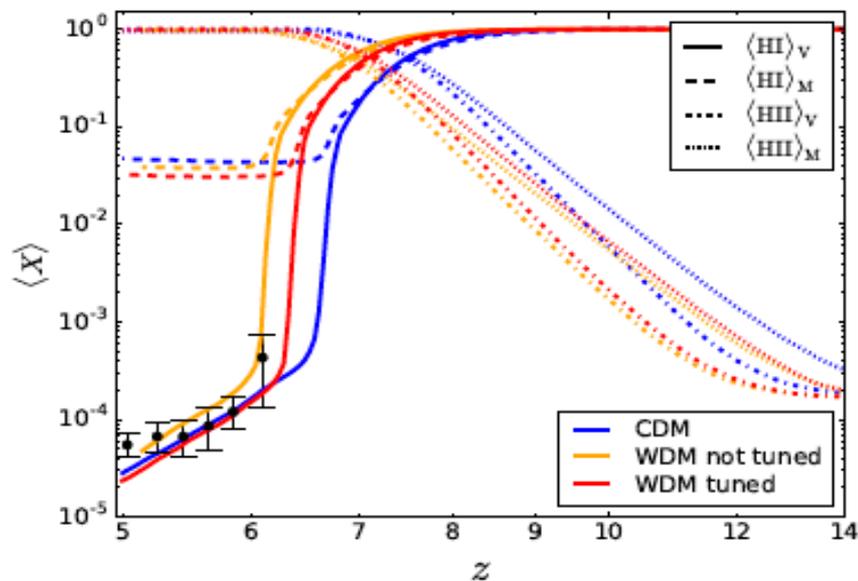
# Impact of $f_{esc}(z)$



# WDM – Reionization Simulations

We have shown that the lack of small scale power in WDM cosmology (relative to the equivalent CDM case) due to the non-negligible free-streaming length of WDM particle considerably delays the reionization processes. However, a higher star formation efficiency (or, equivalently, a lower gas depletion time) compensates for the WDM-suppressed small-scale structure, leading to nearly identical (within the currently observationally constrained range) galaxy luminosity functions in the CDM and WDM cases.

Villanueva-Domingo, Gnedin, Mena (2017)



# Conclusions

- Galaxy formation cannot occur in the same way in CDM and non-standard DM models if they reproduce LF data
- Testing SFR-halo mass relation at low halo masses can provide key insights on DM models
- Faint-end LF sensitive to DM halo abundance, location of turnover signature of DM physics (or minimum mass star forming halo?)
- Implications for reionization models degenerate with astrophysical processes
- What's next?

**Alternative Cosmology Baryon Astrophysics Runs (ACBAR)**