

$\Lambda_b \rightarrow \Lambda_c^{(*)}$ form factors from lattice QCD

Stefan Meinel



Challenges in Semileptonic B Decays, MITP, April 2018

1 Introduction

2 $\Lambda_b \rightarrow \Lambda_c$ form factors from lattice QCD

3 $\Lambda_b \rightarrow \Lambda_c^*$ form factors from lattice QCD

4 Outlook

The $\Lambda_b \rightarrow \Lambda_c^{(*)}$ form factors are needed mainly for:

- $\left| \frac{V_{ub}}{V_{cb}} \right|$ from $\frac{\Gamma(\Lambda_b \rightarrow p \mu^- \bar{\nu}_\mu)}{\Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu)}$
- $|V_{cb}|$ from $\Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu)$
- $R(\Lambda_c^{(*)}) = \frac{\Gamma(\Lambda_b \rightarrow \Lambda_c^{(*)} \tau^- \bar{\nu})}{\Gamma(\Lambda_b \rightarrow \Lambda_c^{(*)} \mu^- \bar{\nu})}$

Name	J^P	Mass [MeV]	Width [MeV]	Strong decay modes
Λ_c	$\frac{1}{2}^+$	2286.46(14)	$3.3(1) \times 10^{-9}$	stable
$\Lambda_c^*(2595)$	$\frac{1}{2}^-$	2592.25(28)	2.6(6)	$\Lambda_c \pi^+ \pi^-$
$\Lambda_c^*(2625)$	$\frac{3}{2}^-$	2628.11(19)	< 0.97	$\Lambda_c \pi^+ \pi^-$

(decays proceed partly through $\Lambda_c^* \rightarrow \Sigma_c^{(*)} (\rightarrow \Lambda_c \pi) \pi$)

[2017 Review of Particle Physics]

In the following, we will treat the Λ_c^* baryons as if they were stable.

Some notation to define the form factors:

$$\langle \Lambda_{c\frac{1}{2}^+}(\mathbf{p}', s') | \bar{c}\Gamma b | \Lambda_b(\mathbf{p}, s) \rangle = \bar{u}(m_{\Lambda_{c\frac{1}{2}^+}}, \mathbf{p}', s') \mathcal{G}^{(\frac{1}{2}^+)}[\Gamma] u(m_{\Lambda_b}, \mathbf{p}, s)$$

$$\langle \Lambda_{c\frac{1}{2}^-}^*(\mathbf{p}', s') | \bar{c}\Gamma b | \Lambda_b(\mathbf{p}, s) \rangle = \bar{u}(m_{\Lambda_{c\frac{1}{2}^-}^*}, \mathbf{p}', s') \gamma_5 \mathcal{G}^{(\frac{1}{2}^-)}[\Gamma] u(m_{\Lambda_b}, \mathbf{p}, s)$$

$$\langle \Lambda_{c\frac{3}{2}^+}^*(\mathbf{p}', s') | \bar{c}\Gamma b | \Lambda_b(\mathbf{p}, s) \rangle = \bar{u}_\lambda(m_{\Lambda_{c\frac{3}{2}^+}^*}, \mathbf{p}', s') \gamma_5 \mathcal{G}^{\lambda(\frac{3}{2}^+)}[\Gamma] u(m_{\Lambda_b}, \mathbf{p}, s)$$

$$\langle \Lambda_{c\frac{3}{2}^-}^*(\mathbf{p}', s') | \bar{c}\Gamma b | \Lambda_b(\mathbf{p}, s) \rangle = \bar{u}_\lambda(m_{\Lambda_{c\frac{3}{2}^-}^*}, \mathbf{p}', s') \mathcal{G}^{\lambda(\frac{3}{2}^-)}[\Gamma] u(m_{\Lambda_b}, \mathbf{p}, s)$$

$$\sum_s u(m, \mathbf{p}, s) \bar{u}(m, \mathbf{p}, s) = m + \not{p}$$

$$\sum_{s'} u_\mu(m', \mathbf{p}', s') \bar{u}_\nu(m', \mathbf{p}', s') =$$

$$-(m' + \not{p}') \left(g_{\mu\nu} - \frac{1}{3} \gamma_\mu \gamma_\nu - \frac{2}{3m'^2} p'_\mu p'_\nu - \frac{1}{3m'} (\gamma_\mu p'_\nu - \gamma_\nu p'_\mu) \right)$$

$\frac{1}{2}^+ \rightarrow \frac{1}{2}^+$ vector and axial vector helicity form factors:

$$\begin{aligned}
 \mathcal{G}^{(\frac{1}{2}^+)}[\gamma^\mu] &= f_0^{(\frac{1}{2}^+)} (m_{\Lambda_b} - m_{\Lambda_c}) \frac{q^\mu}{q^2} \\
 &+ f_+^{(\frac{1}{2}^+)} \frac{m_{\Lambda_b} + m_{\Lambda_c}}{s_+} \left(p^\mu + p'^\mu - (m_{\Lambda_b}^2 - m_{\Lambda_c}^2) \frac{q^\mu}{q^2} \right) \\
 &+ f_\perp^{(\frac{1}{2}^+)} \left(\gamma^\mu - \frac{2m_{\Lambda_c}}{s_+} p^\mu - \frac{2m_{\Lambda_b}}{s_+} p'^\mu \right)
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{G}^{(\frac{1}{2}^+)}[\gamma^\mu \gamma_5] &= -g_0^{(\frac{1}{2}^+)} \gamma_5 (m_{\Lambda_b} + m_{\Lambda_c}) \frac{q^\mu}{q^2} \\
 &- g_+^{(\frac{1}{2}^+)} \gamma_5 \frac{m_{\Lambda_b} - m_{\Lambda_c}}{s_-} \left(p^\mu + p'^\mu - (m_{\Lambda_b}^2 - m_{\Lambda_c}^2) \frac{q^\mu}{q^2} \right) \\
 &- g_\perp^{(\frac{1}{2}^+)} \gamma_5 \left(\gamma^\mu + \frac{2m_{\Lambda_c}}{s_-} p^\mu - \frac{2m_{\Lambda_b}}{s_-} p'^\mu \right)
 \end{aligned}$$

$$s_\pm = (m_{\Lambda_b} - m_{\Lambda_c})^2 - q^2$$

$\frac{1}{2}^+ \rightarrow \frac{1}{2}^-$ vector and axial vector helicity form factors:

$$\begin{aligned}
 \mathcal{G}^{(\frac{1}{2}^-)}[\gamma^\mu] &= f_0^{(\frac{1}{2}^-)} (m_{\Lambda_b} + m_{\Lambda_c^*}) \frac{q^\mu}{q^2} \\
 &+ f_+^{(\frac{1}{2}^-)} \frac{m_{\Lambda_b} - m_{\Lambda_c^*}}{s_-} \left(p^\mu + p'^\mu - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2) \frac{q^\mu}{q^2} \right) \\
 &+ f_\perp^{(\frac{1}{2}^-)} \left(\gamma^\mu + \frac{2m_{\Lambda_c^*}}{s_-} p^\mu - \frac{2m_{\Lambda_b}}{s_-} p'^\mu \right),
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{G}^{(\frac{1}{2}^-)}[\gamma^\mu \gamma_5] &= -g_0^{(\frac{1}{2}^-)} \gamma_5 (m_{\Lambda_b} - m_{\Lambda_c^*}) \frac{q^\mu}{q^2} \\
 &- g_+^{(\frac{1}{2}^-)} \gamma_5 \frac{m_{\Lambda_b} + m_{\Lambda_c^*}}{s_+} \left(p^\mu + p'^\mu - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2) \frac{q^\mu}{q^2} \right) \\
 &- g_\perp^{(\frac{1}{2}^-)} \gamma_5 \left(\gamma^\mu - \frac{2m_{\Lambda_c^*}}{s_+} p^\mu - \frac{2m_{\Lambda_b}}{s_+} p'^\mu \right),
 \end{aligned}$$

$$s_\pm = (m_{\Lambda_b} - m_{\Lambda_c^*})^2 - q^2$$

$\frac{1}{2}^+ \rightarrow \frac{3}{2}^+$ vector and axial vector helicity form factors:

$$\begin{aligned}
 \mathcal{G}^{\lambda(\frac{3}{2}^+)}[\gamma^\mu] &= f_0^{\left(\frac{3}{2}^+\right)} \frac{m_{\Lambda_c^*}}{s_-} \frac{(m_{\Lambda_b} + m_{\Lambda_c^*}) p^\lambda q^\mu}{q^2} \\
 &+ f_+^{\left(\frac{3}{2}^+\right)} \frac{m_{\Lambda_c^*}}{s_+} \frac{(m_{\Lambda_b} - m_{\Lambda_c^*}) p^\lambda (q^2(p^\mu + p'^\mu) - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2)q^\mu)}{q^2 s_-} \\
 &+ f_\perp^{\left(\frac{3}{2}^+\right)} \frac{m_{\Lambda_c^*}}{s_+} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu - m_{\Lambda_c^*} p^\mu)}{s_-} \right) \\
 &+ f_{\perp'}^{\left(\frac{3}{2}^+\right)} \frac{m_{\Lambda_c^*}}{s_+} \left(p^\lambda \gamma^\mu + \frac{2 p^\lambda p'^\mu}{m_{\Lambda_c^*}} + \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu - m_{\Lambda_c^*} p^\mu)}{s_-} - \frac{s_+ g^{\lambda\mu}}{m_{\Lambda_c^*}} \right)
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{G}^{\lambda(\frac{3}{2}^+)}[\gamma^\mu \gamma_5] &= -g_0^{\left(\frac{3}{2}^+\right)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_+} \frac{(m_{\Lambda_b} - m_{\Lambda_c^*}) p^\lambda q^\mu}{q^2} \\
 &- g_+^{\left(\frac{3}{2}^+\right)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_-} \frac{(m_{\Lambda_b} + m_{\Lambda_c^*}) p^\lambda (q^2(p^\mu + p'^\mu) - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2)q^\mu)}{q^2 s_+} \\
 &- g_\perp^{\left(\frac{3}{2}^+\right)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_-} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu + m_{\Lambda_c^*} p^\mu)}{s_+} \right) \\
 &- g_{\perp'}^{\left(\frac{3}{2}^+\right)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_-} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda p'^\mu}{m_{\Lambda_c^*}} + \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu + m_{\Lambda_c^*} p^\mu)}{s_+} + \frac{s_- g^{\lambda\mu}}{m_{\Lambda_c^*}} \right)
 \end{aligned}$$

$\frac{1}{2}^+ \rightarrow \frac{3}{2}^-$ vector and axial vector helicity form factors:

$$\begin{aligned}
 \mathcal{G}^{\lambda(\frac{3}{2}^-)}[\gamma^\mu] &= f_0^{(\frac{3}{2}^-)} \frac{m_{\Lambda_c^*}}{s_+} \frac{(m_{\Lambda_b} - m_{\Lambda_c^*}) p^\lambda q^\mu}{q^2} \\
 &+ f_+^{(\frac{3}{2}^-)} \frac{m_{\Lambda_c^*}}{s_-} \frac{(m_{\Lambda_b} + m_{\Lambda_c^*}) p^\lambda (q^2(p^\mu + p'^\mu) - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2) q^\mu)}{q^2 s_+} \\
 &+ f_{\perp}^{(\frac{3}{2}^-)} \frac{m_{\Lambda_c^*}}{s_-} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu + m_{\Lambda_c^*} p^\mu)}{s_+} \right) \\
 &+ f_{\perp'}^{(\frac{3}{2}^-)} \frac{m_{\Lambda_c^*}}{s_-} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda p'^\mu}{m_{\Lambda_c^*}} + \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu + m_{\Lambda_c^*} p^\mu)}{s_+} + \frac{s_- g^{\lambda\mu}}{m_{\Lambda_c^*}} \right)
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$$\begin{aligned}
 \mathcal{G}^{\lambda(\frac{3}{2}^-)}[\gamma^\mu \gamma_5] &= -g_0^{(\frac{3}{2}^-)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_-} \frac{(m_{\Lambda_b} + m_{\Lambda_c^*}) p^\lambda q^\mu}{q^2} \\
 &- g_+^{(\frac{3}{2}^-)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_+} \frac{(m_{\Lambda_b} - m_{\Lambda_c^*}) p^\lambda (q^2(p^\mu + p'^\mu) - (m_{\Lambda_b}^2 - m_{\Lambda_c^*}^2) q^\mu)}{q^2 s_-} \\
 &- g_{\perp}^{(\frac{3}{2}^-)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_+} \left(p^\lambda \gamma^\mu - \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu - m_{\Lambda_c^*} p^\mu)}{s_-} \right) \\
 &- g_{\perp'}^{(\frac{3}{2}^-)} \gamma_5 \frac{m_{\Lambda_c^*}}{s_+} \left(p^\lambda \gamma^\mu + \frac{2 p^\lambda p'^\mu}{m_{\Lambda_c^*}} + \frac{2 p^\lambda (m_{\Lambda_b} p'^\mu - m_{\Lambda_c^*} p^\mu)}{s_-} - \frac{s_+ g^{\lambda\mu}}{m_{\Lambda_c^*}} \right)
 \end{aligned}$$

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Early work (quenched, focused on Isgur-Wise function):

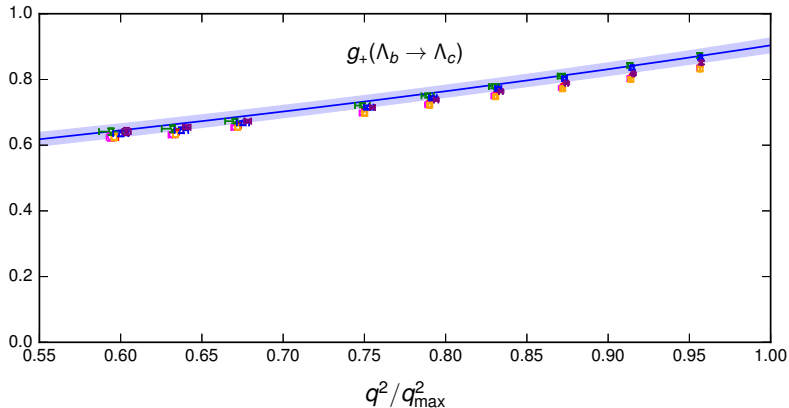
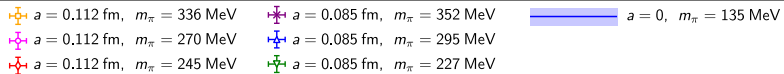
- K. C. Bowler *et al.* (UKQCD Collaboration), hep-lat/9709028/PRD 1998
- S. Gottlieb and S. Tamhankar, hep-lat/0301022/Lattice 2002

Our work:

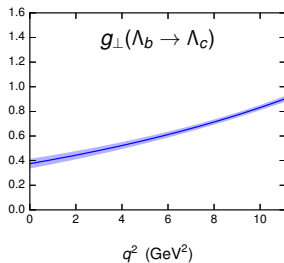
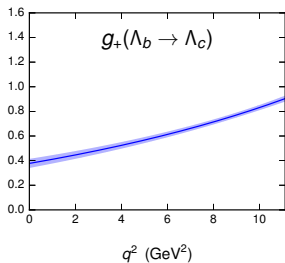
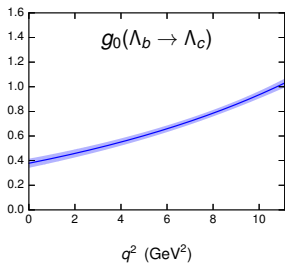
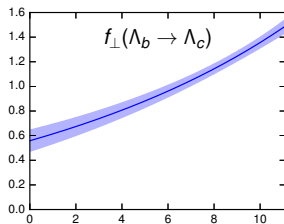
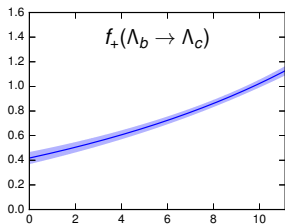
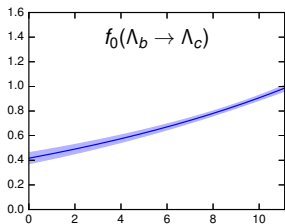
- W. Detmold, C. Lehner, S. Meinel, arXiv:1503.01421/PRD 2015
(vector and axial vector form factors)
- A. Datta, S. Kamali, S. Meinel, A. Rashed, arXiv:1702.02243/JHEP 2017
(tensor form factors)

- Gauge field configurations generated by the RBC and UKQCD collaborations
[Y. Aoki *et al.*, arXiv:1011.0892/PRD 2011]
- u , d , s quarks: domain-wall action
[D. Kaplan, hep-lat/9206013/PLB 1992; V. Furman and Y. Shamir, hep-lat/9303005/NPB 1995]
- c , b quarks: anisotropic clover with two or three parameters tuned nonperturbatively
[A. El-Khadra, A. Kronfeld, P. Mackenzie, hep-lat/9604004/PRD 1997; Y. Aoki *et al.*, arXiv:1206.2554/PRD 2012]
- “Mostly nonperturbative” renormalization
[A. El-Khadra *et al.*, hep-ph/0101023/PRD 2001]
- Three-point functions with 12 source-sink separations
- Combined chiral/continuum/kinematic extrapolation using modified z -expansion [C. Bourrely, I. Caprini, L. Lellouch, arXiv:0807.2722/PRD 2009]

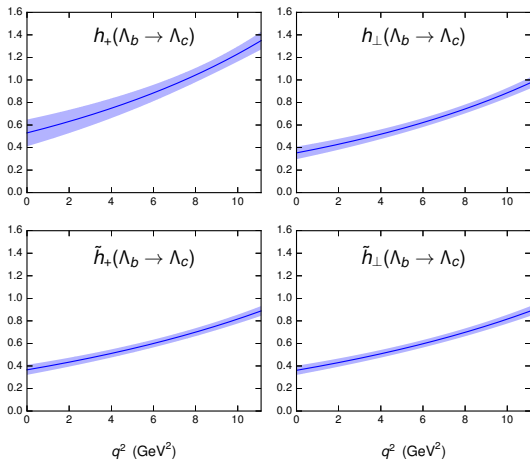
$N_s^3 \times N_t$	β	$am_{u,d}^{(\text{sea})}$	$am_{u,d}^{(\text{val})}$	$am_s^{(\text{sea})}$	a (fm)	$m_\pi^{(\text{sea})}$ (MeV)	$m_\pi^{(\text{val})}$ (MeV)
$24^3 \times 64$	2.13	0.005	0.005	0.04	≈ 0.11	≈ 340	≈ 340
$24^3 \times 64$	2.13	0.005	0.002	0.04	≈ 0.11	≈ 340	≈ 270
$24^3 \times 64$	2.13	0.005	0.001	0.04	≈ 0.11	≈ 340	≈ 250
$32^3 \times 64$	2.25	0.006	0.006	0.03	≈ 0.08	≈ 360	≈ 360
$32^3 \times 64$	2.25	0.004	0.004	0.03	≈ 0.08	≈ 300	≈ 300
$32^3 \times 64$	2.25	0.004	0.002	0.03	≈ 0.08	≈ 300	≈ 230



Vector and axial vector form factors:

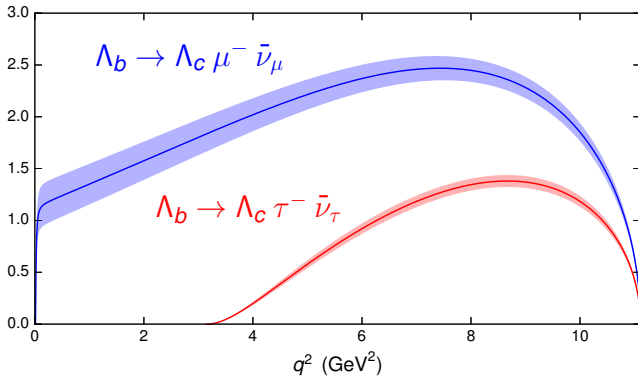


Tensor form factors:



Differential decay rates in the standard model:

$$\frac{d\Gamma/dq^2}{|V_{cb}|^2} \text{ (ps}^{-1} \text{ GeV}^{-2}\text{)}$$



Integrated decay rates in the standard model:

$$\frac{1}{|V_{cb}|^2} \Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu) = (21.5 \pm 0.8_{\text{stat}} \pm 1.1_{\text{syst}}) \text{ ps}^{-1}$$

(6.3% uncertainty \rightarrow 3.2% for $|V_{cb}|$)

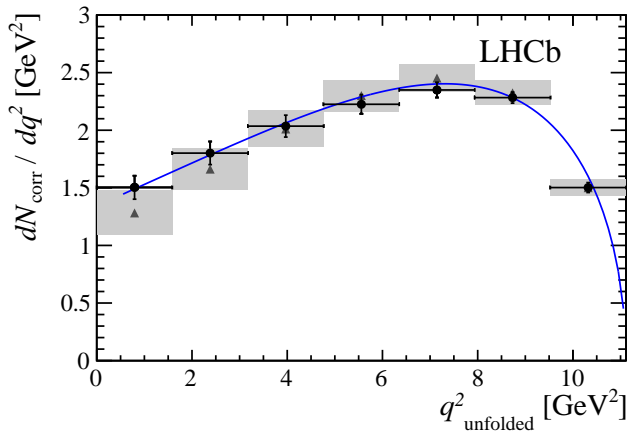
$$\frac{1}{|V_{cb}|^2} \int_7^{\text{GeV}^2} \frac{d\Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu)}{dq^2} dq^2 = (8.37 \pm 0.16_{\text{stat}} \pm 0.34_{\text{syst}}) \text{ ps}^{-1}$$

(4.5% uncertainty \rightarrow 2.3% for $|V_{cb}|$)

$$R(\Lambda_c) = \frac{\Gamma(\Lambda_b \rightarrow \Lambda_c \tau^- \bar{\nu}_\tau)}{\Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu)} = 0.3328 \pm 0.0074_{\text{stat}} \pm 0.0070_{\text{syst}}$$

(3.1% uncertainty)

The shape of the $\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu$ decay rate measured by LHCb (black data points) agrees with our lattice QCD prediction (gray).



[LHCb Collaboration, arXiv:1709.01920/PRD 2017]

BSM phenomenology of $\Lambda_b \rightarrow \Lambda_c \tau^- \bar{\nu}$

[A. Datta, S. Kamali, S. Meinel, A. Rashed, arXiv:1702.02243/JHEP 2017]

- We have shown that a future measurement of $R(\Lambda_c)$ can provide useful constraints on all of the couplings g_L, g_R, g_S, g_P, g_T .
- The paper contains plots showing the correlations between $R(D^{(*)})$ and $R(\Lambda_c)$ for several leptoquark models.

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3 $\Lambda_b \rightarrow \Lambda_c^*$ form factors from lattice QCD

[S. Meinel and G. Rendon, work in progress]

4 Outlook

We work in the Λ_c^* rest frame to allow exact spin-parity projection. We use the interpolating field

$$\begin{aligned}
 (\Lambda_c^*)_{j\gamma} = & \epsilon^{abc} (C\gamma_5)_{\alpha\beta} \left[\tilde{c}_\alpha^a \tilde{d}_\beta^b (\nabla_j \tilde{u})_\gamma^c - \tilde{c}_\alpha^a \tilde{u}_\beta^b (\nabla_j \tilde{d})_\gamma^c \right. \\
 & \left. + \tilde{u}_\alpha^a (\nabla_j \tilde{d})_\beta^b \tilde{c}_\gamma^c - \tilde{d}_\alpha^a (\nabla_j \tilde{u})_\beta^b \tilde{c}_\gamma^c \right]
 \end{aligned}$$

($\tilde{}$ denotes Gaussian smearing)

[S. Meinel and G. Rendon, arXiv:1608.08110/Lattice2016]

This requires light-quark propagators with derivative sources.

We project to $J^P = \frac{1}{2}^-$ and $J^P = \frac{3}{2}^-$ using

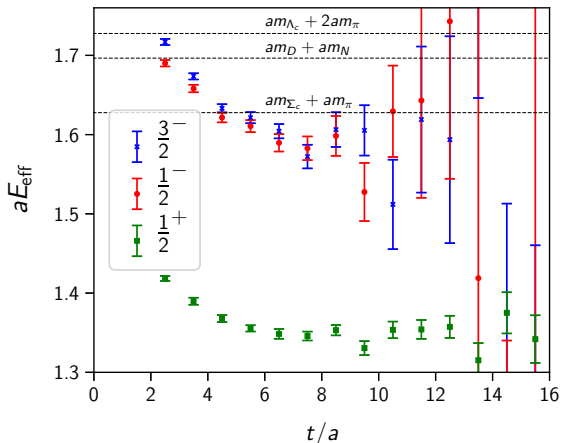
$$\begin{aligned}
 P_{jk}^{(\frac{1}{2}^-)} &= \frac{1}{3} \gamma_j \gamma_k \frac{1 + \gamma_0}{2}, \\
 P_{jk}^{(\frac{3}{2}^-)} &= \left(g_{jk} - \frac{1}{3} \gamma_j \gamma_k \right) \frac{1 + \gamma_0}{2}.
 \end{aligned}$$

- Gauge field configurations generated by the RBC and UKQCD collaborations
[Y. Aoki *et al.*, arXiv:1011.0892/PRD 2011]
- u , d , s quarks: domain-wall action
[D. Kaplan, hep-lat/9206013/PLB 1992; V. Furman and Y. Shamir, hep-lat/9303005/NPB 1995]
- All-mode averaging with 1 exact and 32 sloppy propagators per configuration
[E. Shintani *et al.*, arXiv:1402.0244/PRD 2015]
- c , b quarks: anisotropic clover with three parameters, re-tuned more accurately to $D_s^{(*)}$ and $B_s^{(*)}$ dispersion relation and HFS
- “Mostly nonperturbative” renormalization
[A. El-Khadra *et al.*, hep-ph/0101023/PRD 2001]
- Three-point functions with 7 source-sink separations (plan to add 2 more)

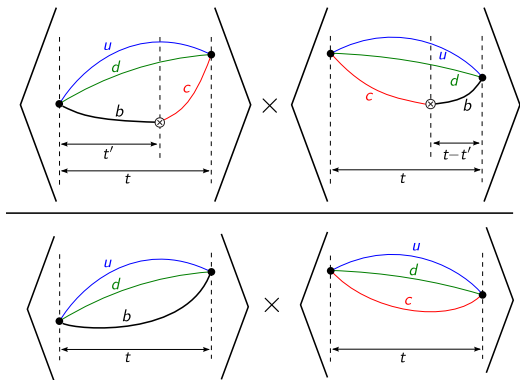
$N_s^3 \times N_t$	β	$am_{u,d}$	am_s	a (fm)	m_π (MeV)	Run status
$24^3 \times 64$	2.13	0.01	0.04	≈ 0.111	≈ 430	1/4 of cfgs done
$24^3 \times 64$	2.13	0.005	0.04	≈ 0.111	≈ 340	1/4 of cfgs done
$32^3 \times 64$	2.25	0.004	0.03	≈ 0.083	≈ 300	1/4 of cfgs done
$48^3 \times 96$	2.31	0.002144	0.02144	≈ 0.071	≈ 230	planned

Results from $24^3 \times 64$, $am_{u,d} = 0.005$ ensemble, 78 configs \times 32 sources

$a^{-1} = 1.785(5)$ GeV



Extracting the form factors from ratios of 3pt and 2pt functions



t = source-sink separation

t' = current insertion time

We have data for two different Λ_b momenta: $\mathbf{p} = (0, 0, 2) \frac{2\pi}{L} \approx 0.9$ GeV and $\mathbf{p} = (0, 0, 3) \frac{2\pi}{L} \approx 1.4$ GeV

Schematically,

$$R_f(\mathbf{p}, t) = \sqrt{(\text{kinematic factors}) \times (\text{polarization vectors}) \times (\text{ratio at } t' = t/2)}$$

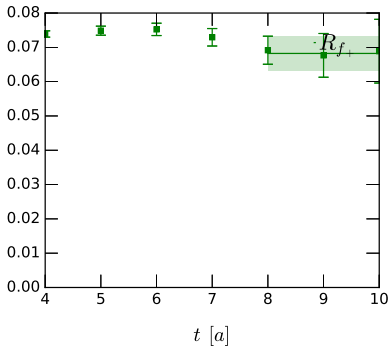
→ $f(\mathbf{p})$ for large t

Example: R_{f_+} for $\Lambda_b \rightarrow \Lambda_c^* (\frac{3}{2}^-)$

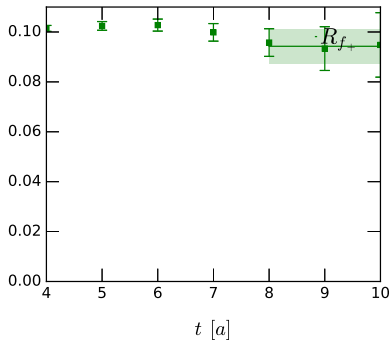
preliminary

Results from $24^3 \times 64$, $am_{u,d} = 0.005$ ensemble, 78 configs \times 32 sources

$$\mathbf{p} = (0, 0, 2) \frac{2\pi}{L}$$

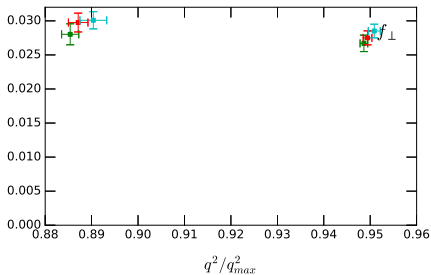
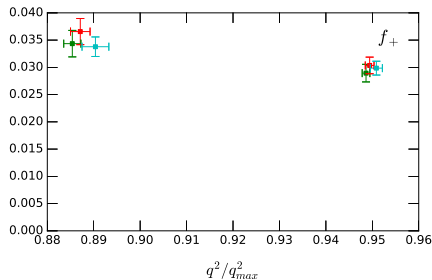
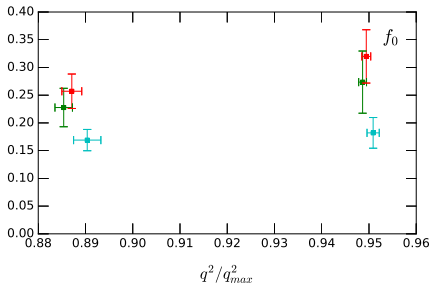


$$\mathbf{p} = (0, 0, 3) \frac{2\pi}{L}$$



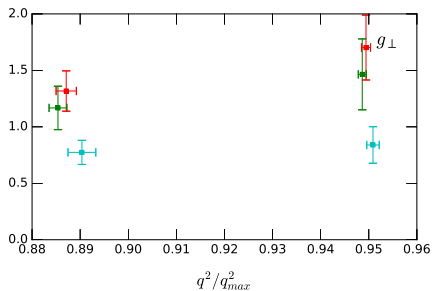
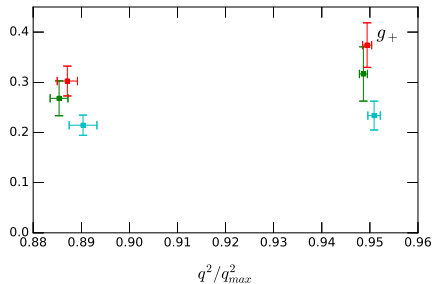
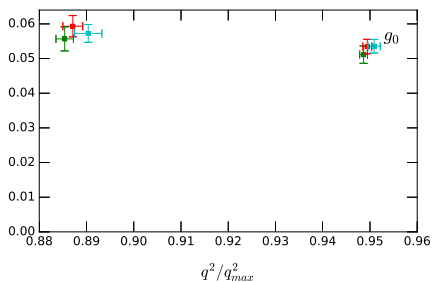
$\Lambda_b \rightarrow \Lambda_c^* (\frac{1}{2}^-)$ vector form factors

preliminary

[red] $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV[green] $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV[cyan] $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

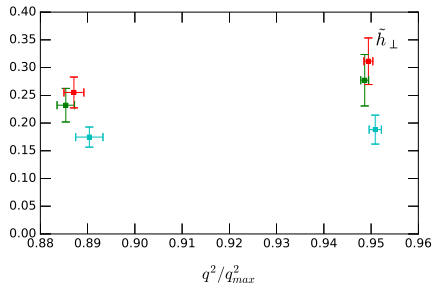
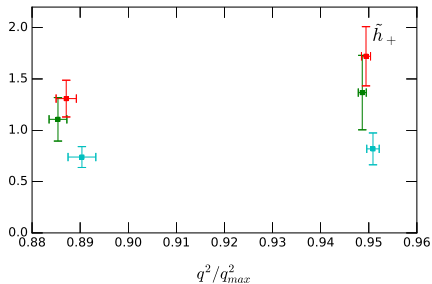
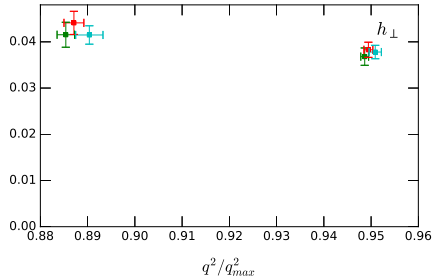
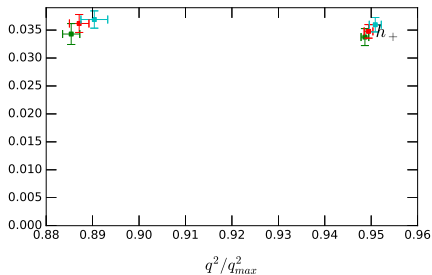
$\Lambda_b \rightarrow \Lambda_c^* (\frac{1}{2}^-)$ axial vector form factors

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

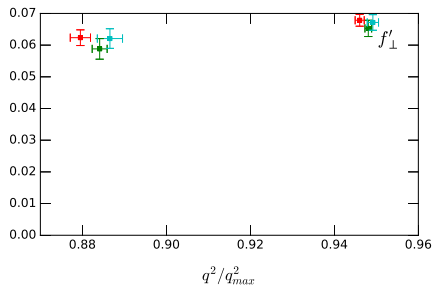
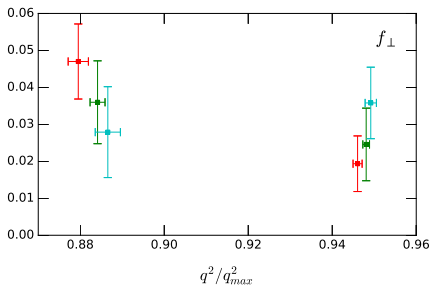
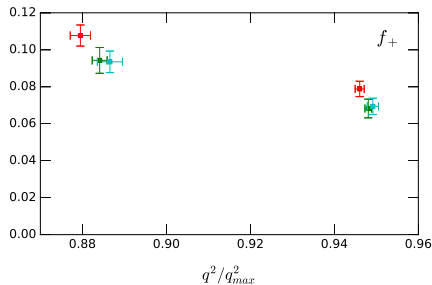
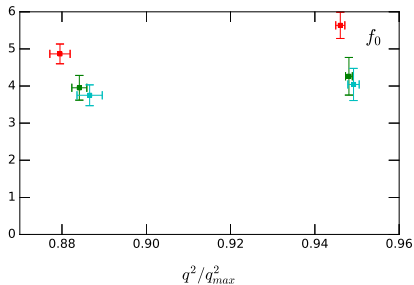
$\Lambda_b \rightarrow \Lambda_c^* (\frac{1}{2}^-)$ tensor form factors

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

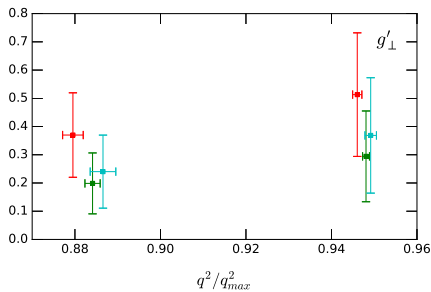
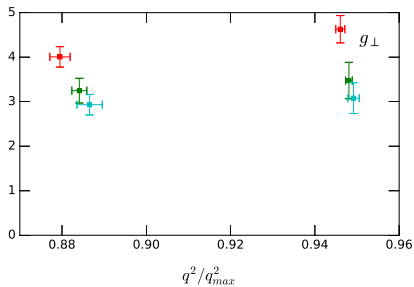
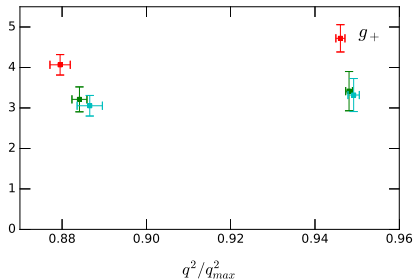
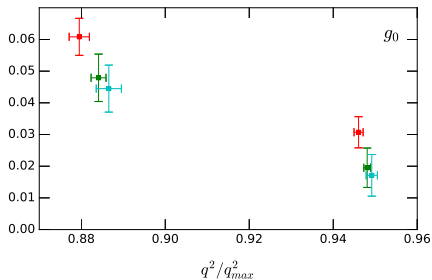
$\Lambda_b \rightarrow \Lambda_c^* \left(\frac{3}{2}^-\right)$ vector form factors

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

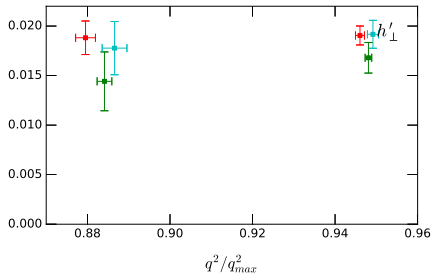
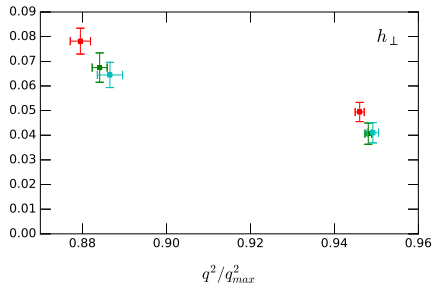
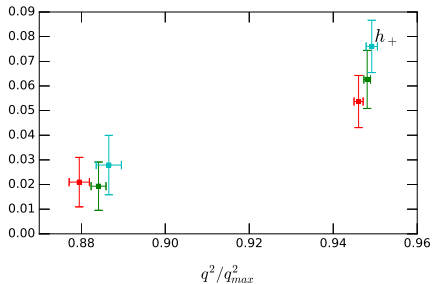
$\Lambda_b \rightarrow \Lambda_c^* \left(\frac{3}{2}^-\right)$ axial vector form factors

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

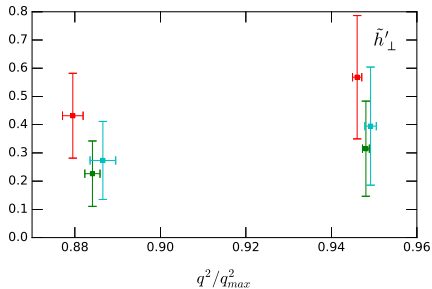
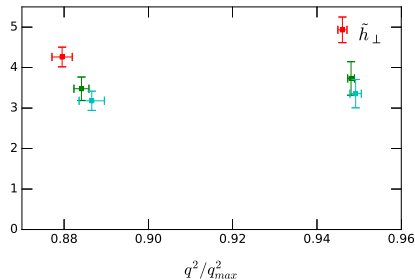
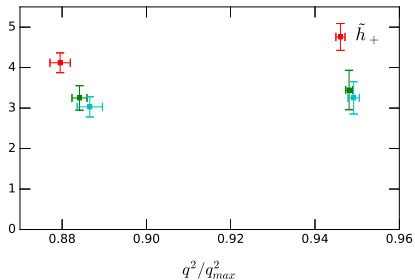
$\Lambda_b \rightarrow \Lambda_c^* \left(\frac{3}{2}^-\right)$ tensor form factors part 1

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

$\Lambda_b \rightarrow \Lambda_c^* \left(\frac{3}{2}^-\right)$ tensor form factors part 2

preliminary

■ $a \approx 0.11$ fm, $m_\pi \approx 430$ MeV■ $a \approx 0.11$ fm, $m_\pi \approx 340$ MeV■ $a \approx 0.08$ fm, $m_\pi \approx 300$ MeV

- 1 Introduction
- 2 $\Lambda_b \rightarrow \Lambda_c$ form factors from lattice QCD
- 3 $\Lambda_b \rightarrow \Lambda_c^*$ form factors from lattice QCD
- 4 Outlook

$\Lambda_b \rightarrow \Lambda_c^*$ next steps:

- 4×statistics
- third lattice spacing
- chiral/continuum extrapolations

The lattice form factor results are limited to high q^2 .

To predict $R(\Lambda_c^*)$, it will be helpful to combine the lattice results with experimental data for the shape of the $\Lambda_b \rightarrow \Lambda_c^* \mu \bar{\nu}$ differential decay rates, making use of HQET.

[P. Boer, M. Bordone, E. Graverini, P. Owen, M. Rotondo, and D. Van Dyk, arXiv:1801.08367]

An improved calculation of $\Lambda_b \rightarrow p, \Lambda, \Lambda_c (\frac{1}{2}^+)$ form factors is also underway:

- remove data sets with $m_{u,d}^{(\text{val})} < m_{u,d}^{(\text{sea})}$, add two new ensembles
- for $\Lambda_b \rightarrow \Lambda$: physical $m_s^{(\text{val})}$
- more accurate tuning of charm and bottom actions
- all-mode-averaging for higher statistics
- better source smearing

$N_s^3 \times N_t$	β	$am_{u,d}^{(\text{sea})}$	$am_{u,d}^{(\text{val})}$	$am_s^{(\text{sea})}$	a (fm)	$m_\pi^{(\text{sea})}$ (MeV)	$m_\pi^{(\text{val})}$ (MeV)	Status
$24^3 \times 64$	2.13	0.005	0.005	0.04	≈ 0.111	≈ 340	≈ 340	done
$24^3 \times 64$	2.13	0.005	0.002	0.04	≈ 0.111	≈ 340	≈ 270	
$24^3 \times 64$	2.13	0.005	0.001	0.04	≈ 0.111	≈ 340	≈ 250	
$48^3 \times 96$	2.13	0.00078	0.00078	0.0362	≈ 0.114	≈ 140	≈ 140	done
$32^3 \times 64$	2.25	0.006	0.006	0.03	≈ 0.083	≈ 360	≈ 360	done
$32^3 \times 64$	2.25	0.004	0.004	0.03	≈ 0.083	≈ 300	≈ 300	done
$32^3 \times 64$	2.25	0.004	0.002	0.03	≈ 0.083	≈ 300	≈ 230	
$48^3 \times 96$	2.31	0.002144	0.002144	0.02144	≈ 0.071	≈ 230	≈ 230	planned

Expected completion: 2020. Hope to reduce total uncertainties by factor of 2.

Extra slides

Breakdown of uncertainties in partially integrated $\Lambda_b \rightarrow p\mu\bar{\nu}_\mu$ and $\Lambda_b \rightarrow \Lambda_c\mu\bar{\nu}_\mu$ decay rates (in percent):

	$\zeta_{p\mu\bar{\nu}}(15 \text{ GeV}^2)$	$\zeta_{\Lambda_c\mu\bar{\nu}}(7 \text{ GeV}^2)$	$\frac{\zeta_{p\mu\bar{\nu}}(15 \text{ GeV}^2)}{\zeta_{\Lambda_c\mu\bar{\nu}}(7 \text{ GeV}^2)}$
Statistics	6.2	1.9	6.5
Finite volume	5.0	2.5	4.9
Continuum extrapolation	3.0	1.4	2.8
Chiral extrapolation	2.6	1.8	2.6
RHQ parameters	1.4	1.7	2.3
Matching & improvement	1.7	0.9	2.1
Missing isospin breaking/QED	1.2	1.4	2.0
Scale setting	1.7	0.3	1.8
z expansion	1.2	0.2	1.3
Total	8.8	4.5	9.8

Note: the individual systematic uncertainties are correlated in a complicated way. Use the total uncertainty only.

[W. Detmold, C. Lehner, S. Meinel, arXiv:1503.01421/PRD 2015]

Λ_b and Λ_c decay form factors from lattice QCD: References

	m_b	a [fm]	m_π [MeV]	Reference
$\Lambda_b \rightarrow \Lambda$	∞	0.11, 0.08	230–360	arXiv:1212.4827/PRD 2013
$\Lambda_b \rightarrow p$	∞	0.11, 0.08	230–360	arXiv:1306.0446/PRD 2013
$\Lambda_b \rightarrow p$	phys.	0.11, 0.08	230–360	arXiv:1503.01421/PRD 2015
$\Lambda_b \rightarrow \Lambda_c$	phys.	0.11, 0.08	230–360	arXiv:1503.01421/PRD 2015, arXiv:1702.02243/JHEP 2017
$\Lambda_b \rightarrow \Lambda$	phys.	0.11, 0.08	230–360	arXiv:1602.01399/PRD 2016
$\Lambda_b \rightarrow \Lambda^*(\frac{3}{2}^-)$	phys.	0.11	340	arXiv:1608.08110/Lattice 2016
$\Lambda_c \rightarrow \Lambda$		0.11, 0.08	140 –360	arXiv:1611.09696/PRL 2017
$\Lambda_c \rightarrow p$		0.11, 0.08	230–360	arXiv:1712.05783/PRD 2018