

# Status of SIMBA

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*Inclusive  $|V_{ub}|$  and New Physics from Inclusive Decays*



Frank Tackmann

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William Sutcliffe

Lu Cao

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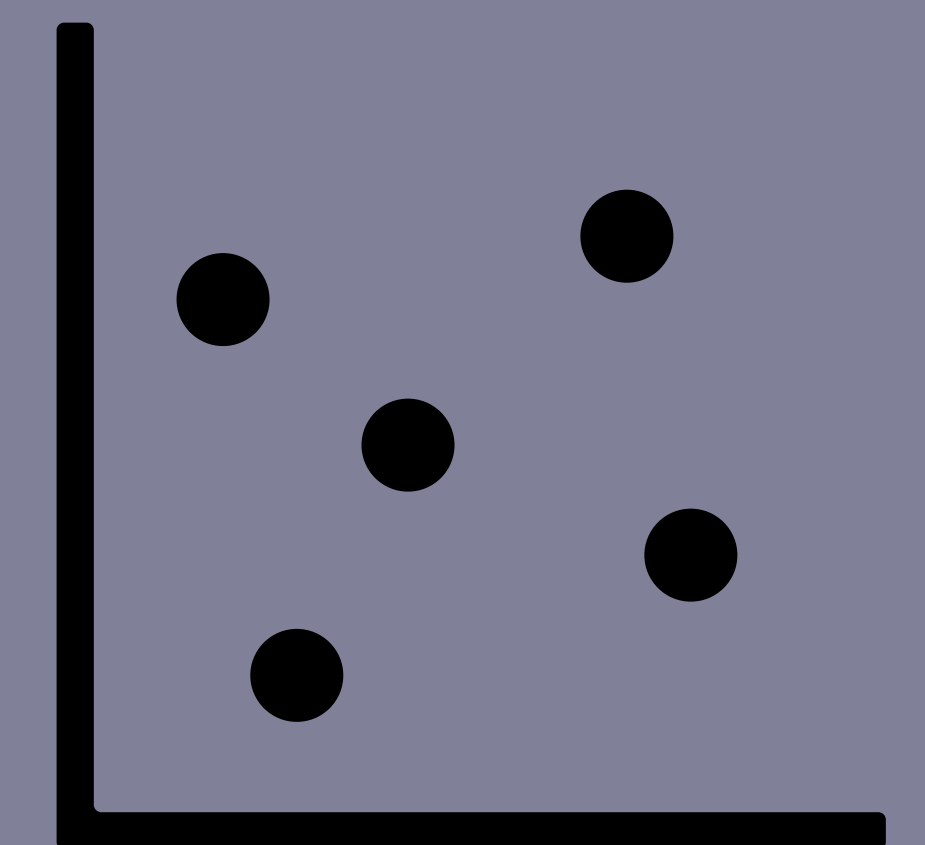
Heiko Lacker



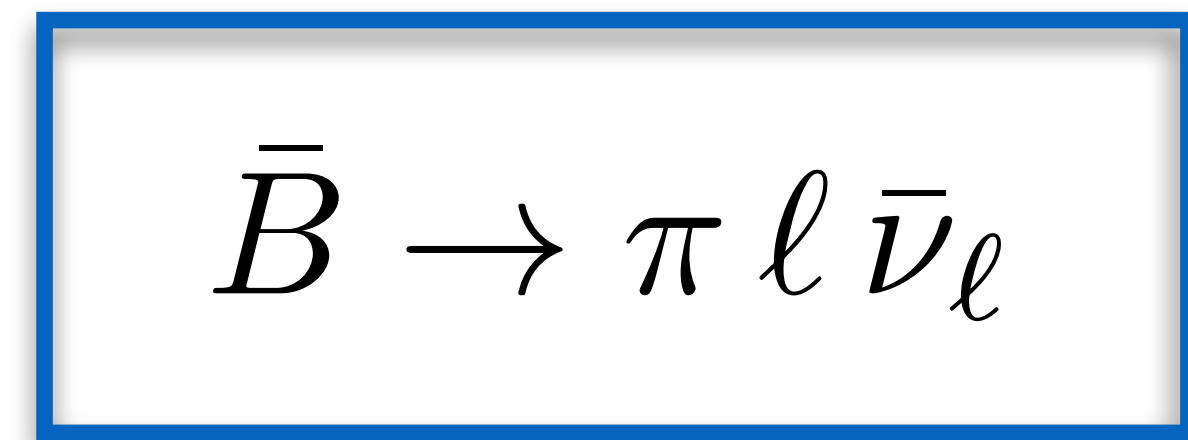
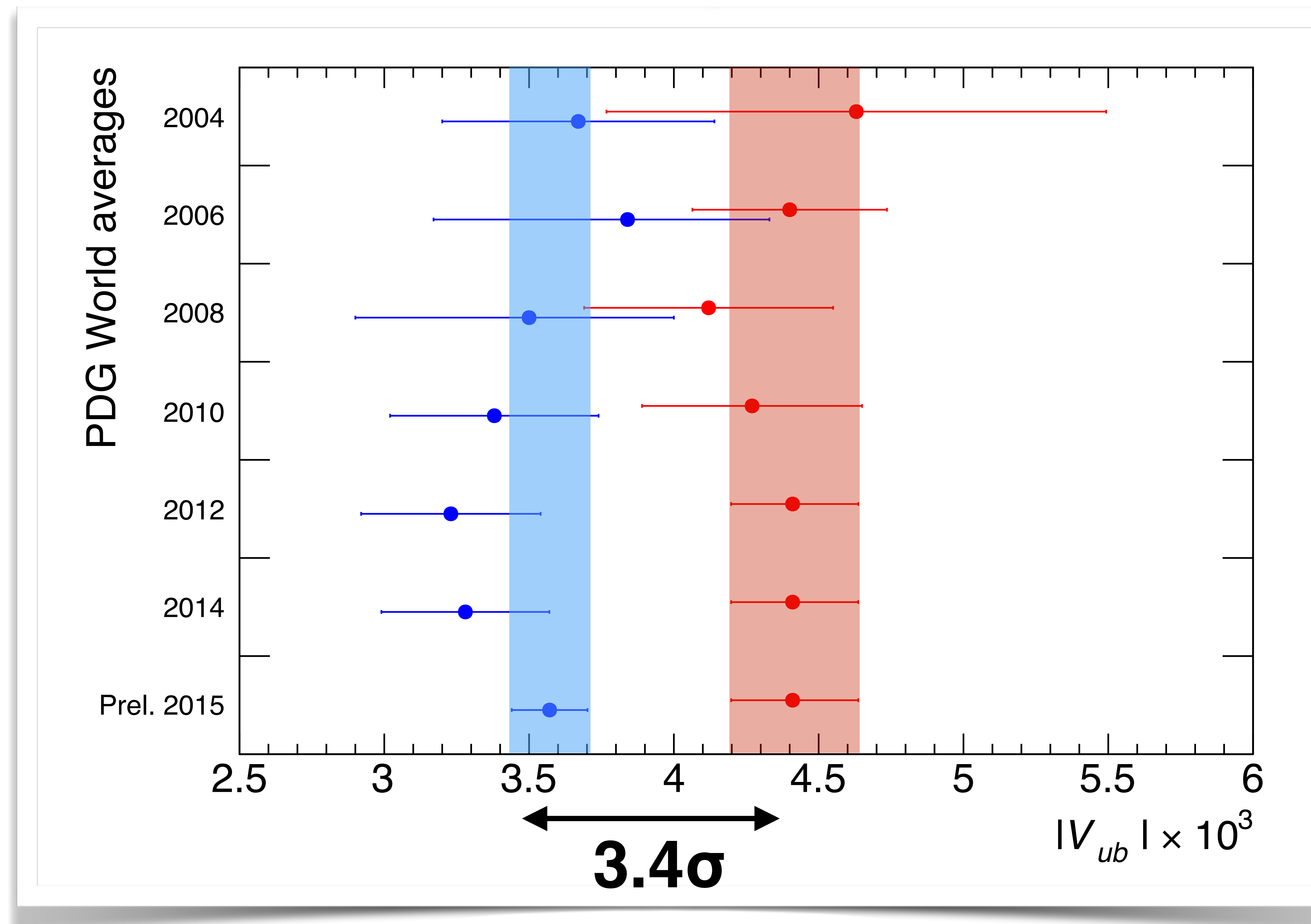
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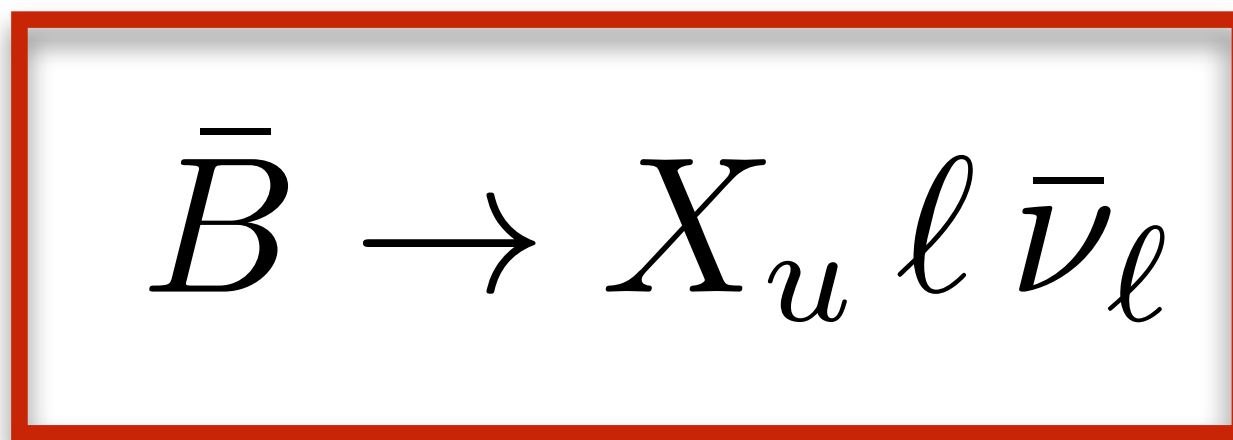
# Introduction to SIMBA



# Inclusive $|V_{ub}|$



Exclusive Ansatz

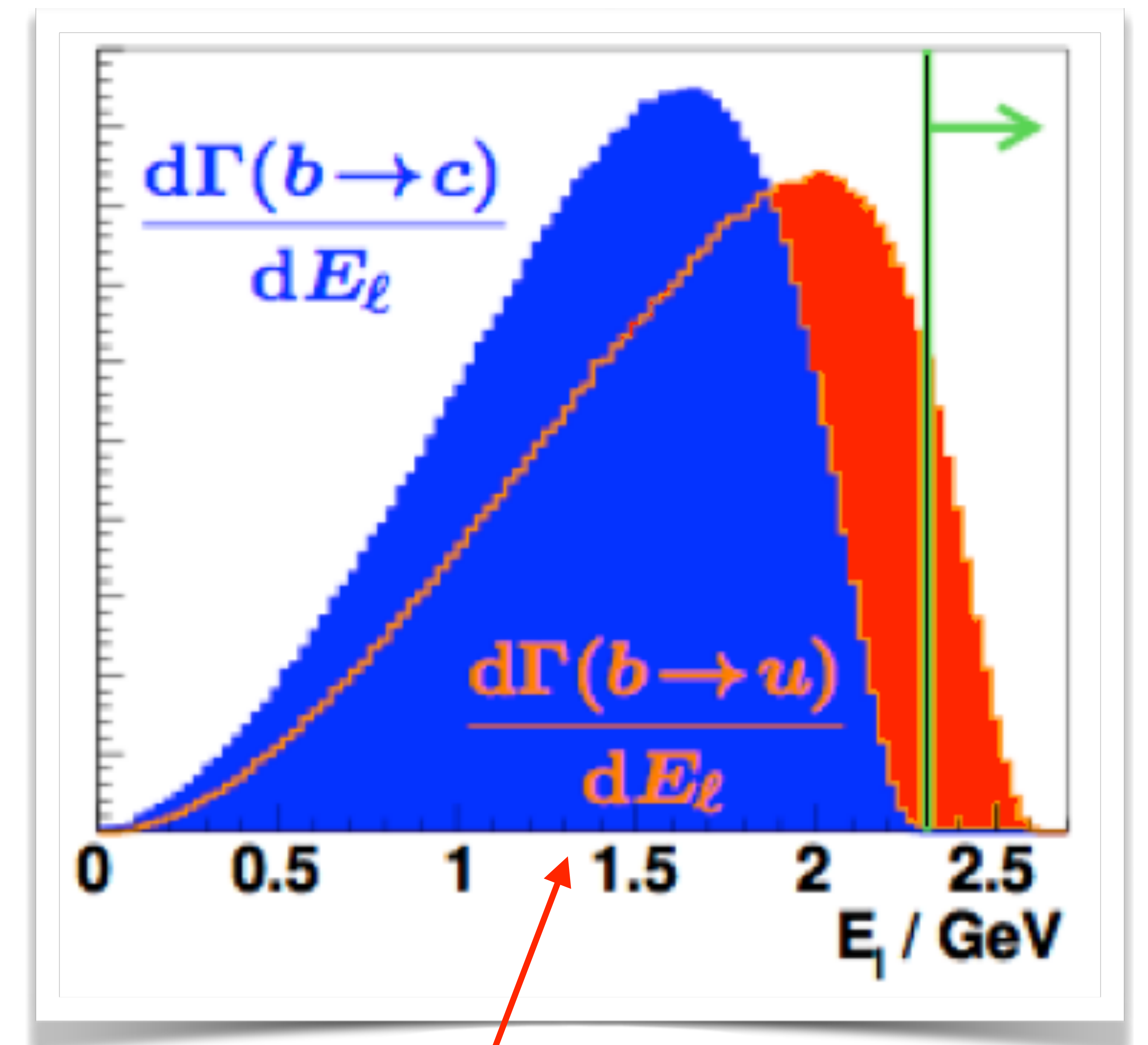
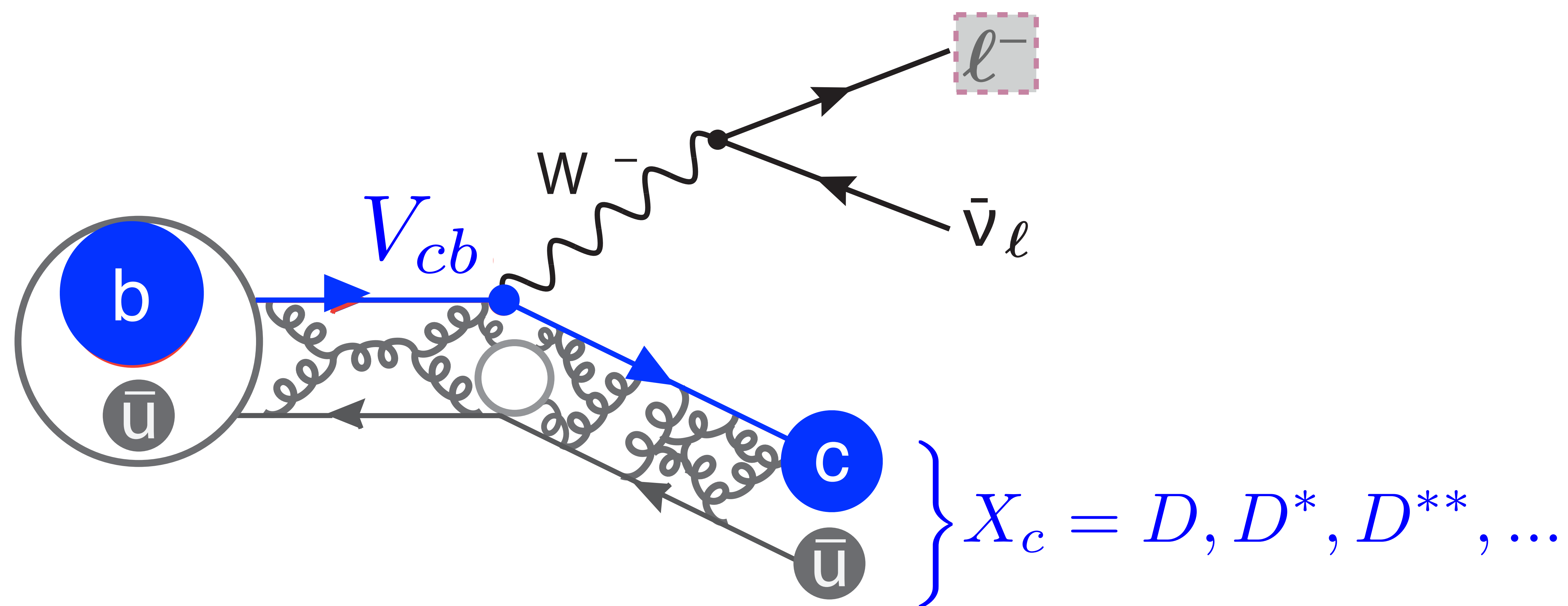


Inclusive Ansatz



# What makes measuring inclusive $|V_{ub}|$ difficult?

- Inclusive  $|V_{ub}|$  determinations are difficult:
  - Large backgrounds from  $\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$



Signal x 100!

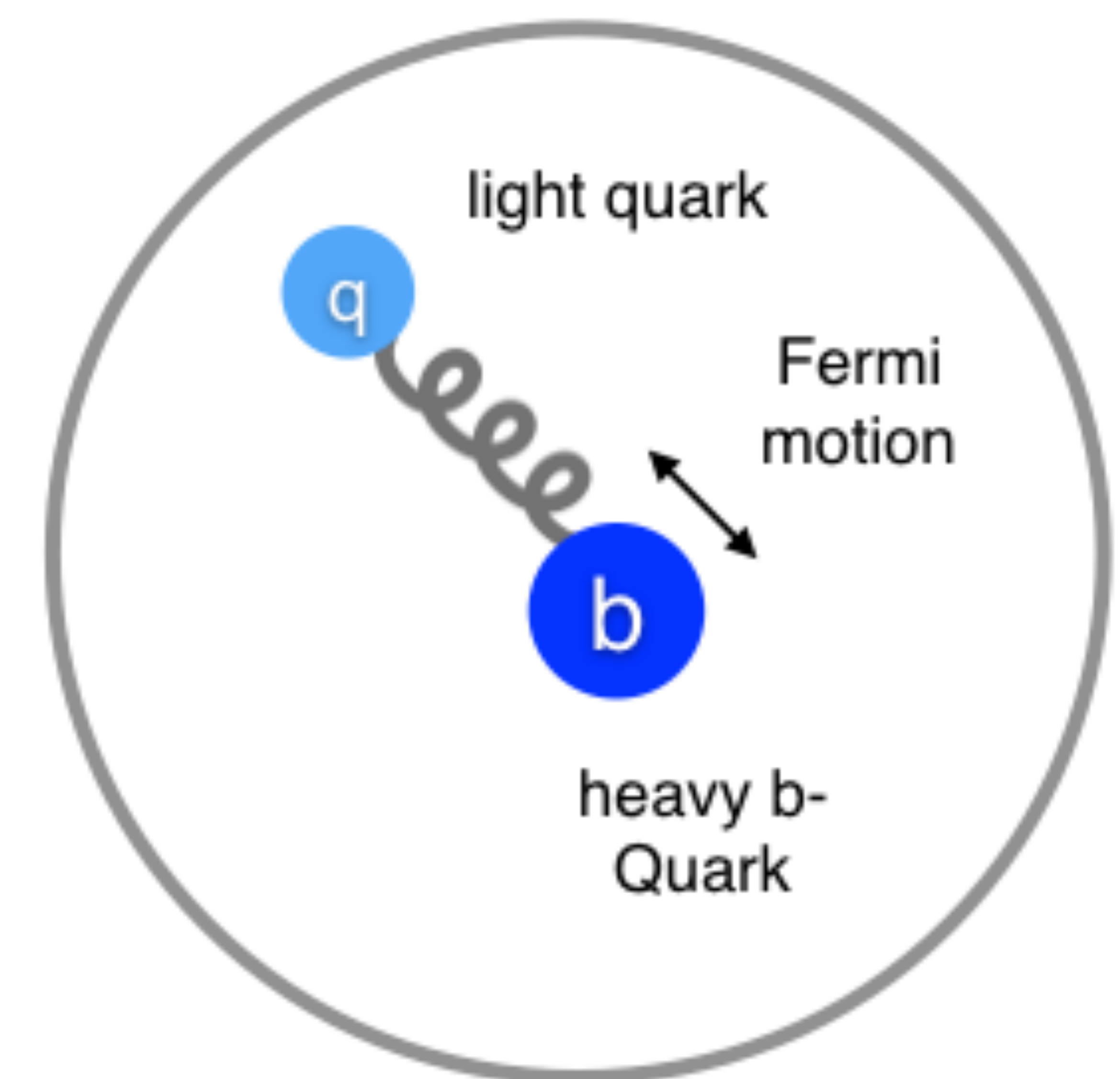
- **O(100)** larger than signal
- Decays involving  $D^{**}$  not well understood
- Clear separation only possible in **corners of phase space**



# What makes measuring inclusive $|V_{ub}|$ difficult?

- Does not help much though, as theory prediction heavily depends on details of shape function

$$|V_{ub}| = \sqrt{\frac{\Delta\mathcal{B}(B \rightarrow X_u \ell \bar{\nu}_\ell)}{\tau_B \Delta\Gamma(B \rightarrow X_u \ell \bar{\nu}_\ell)}}$$



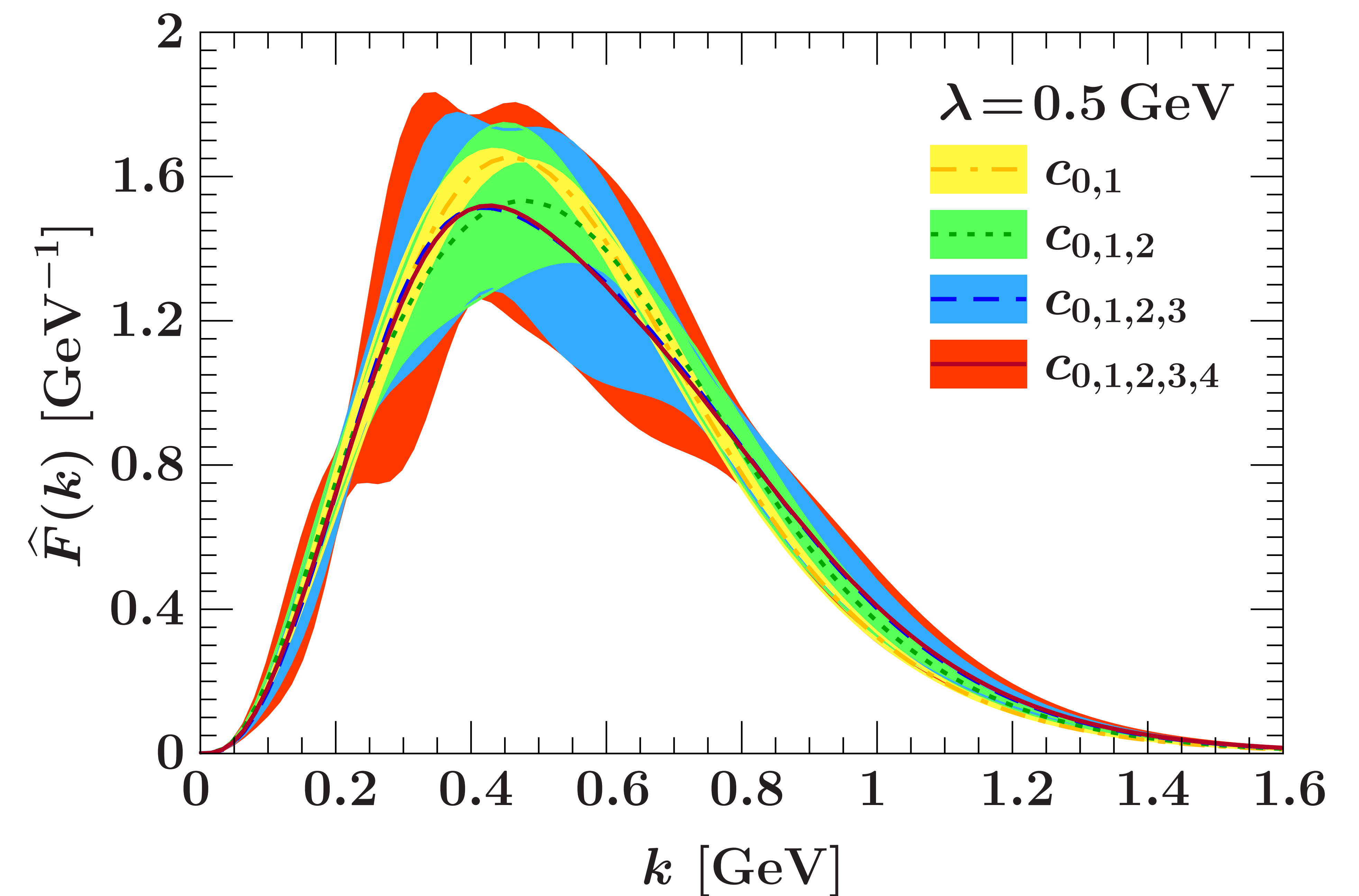
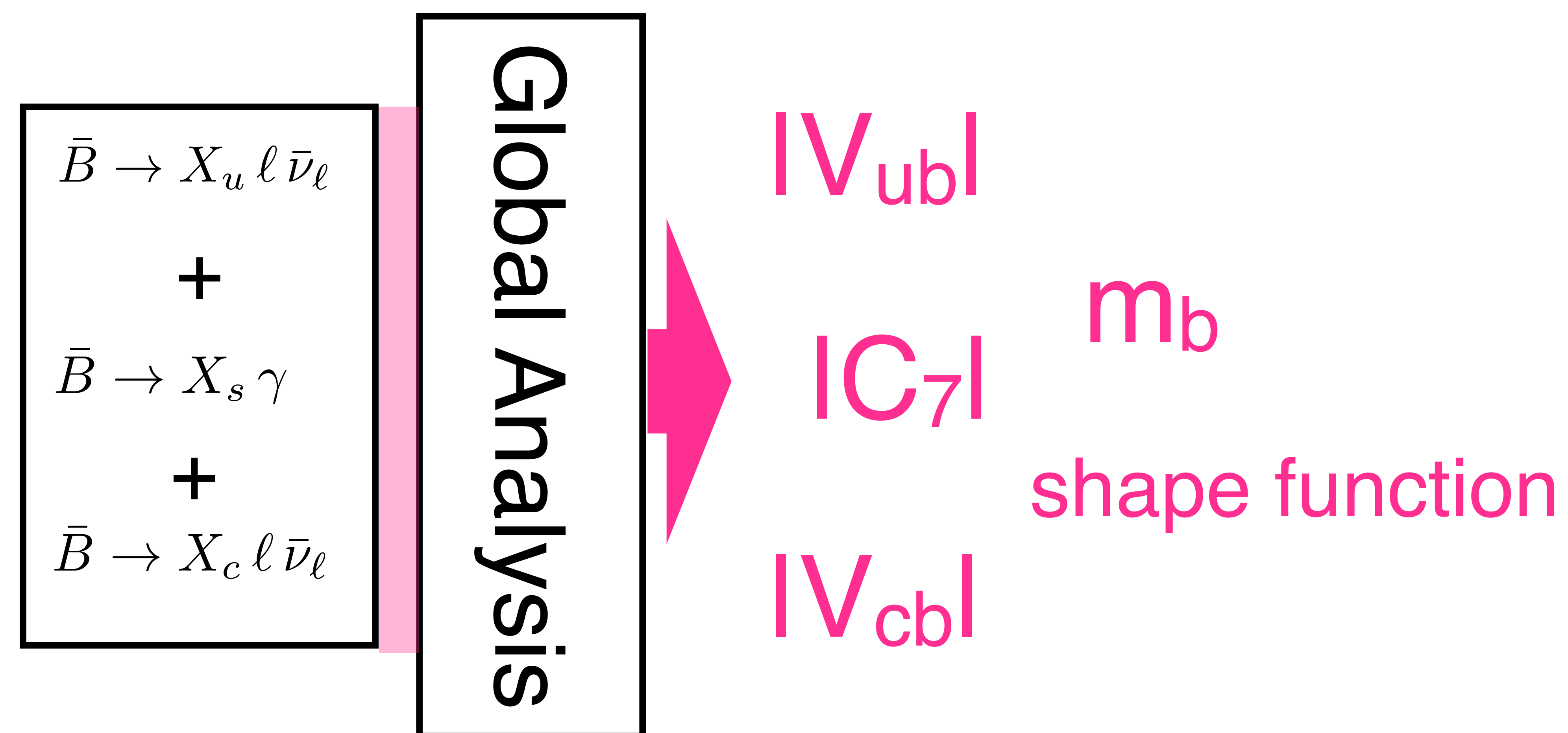
- “Experimentally wonderful region, but completely useless as no theorist can tell you what  $|V_{ub}|$  you measured”



# The Idea in a nutshell



- **The SIMBA idea:** turn this around:
  - **Large model dependence**  $\iff$  **sensitivity to constrain SF**
    - Most information in differential spectra
    - Can use different decay modes (same leading shape function) and carry out global analysis that propagates uncertainties:

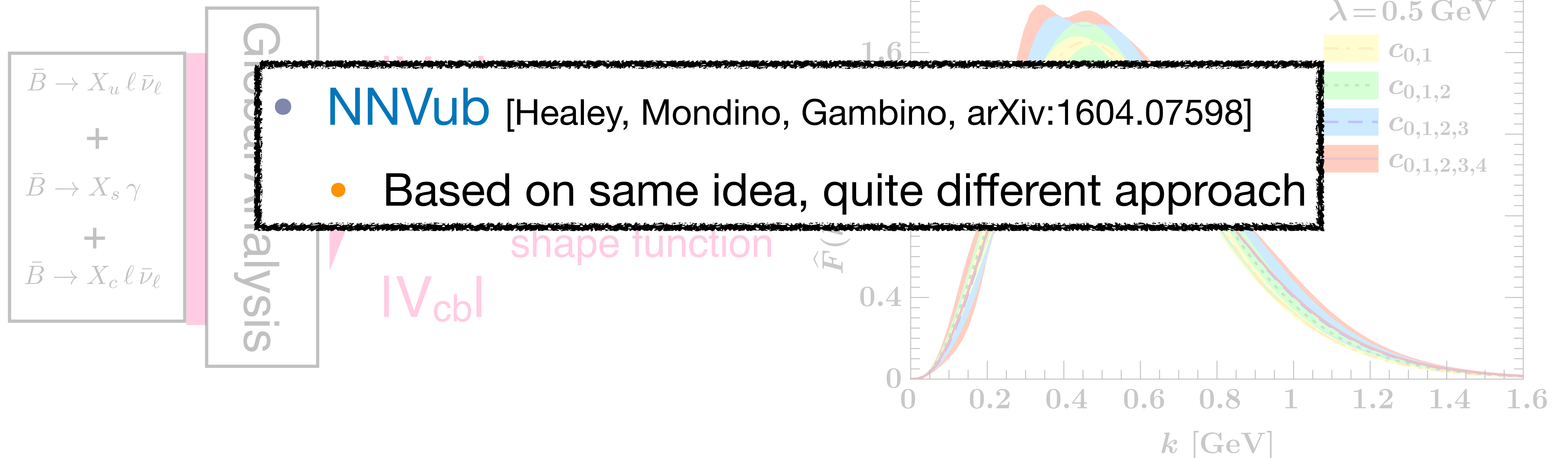




# The Idea in a nutshell



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  - Large model dependence  $\iff$  sensitivity to constrain SF
    - Most information in differential spectra
    - Can use different decay modes (same leading shape function) and carry out global analysis that propagates uncertainties:





# Theory side of Global Fits

- **SIMBA** master formulae:

$$d\Gamma_s = |V_{tb}V_{ts}^*|^2 m_b^2 |C_7^{\text{incl}}|^2 \int dk \widehat{W}_{77}(E_\gamma; k) \widehat{F}(m_B - 2E_\gamma - k) + \dots$$

$$d\Gamma_u = |V_{ub}|^2 \int dk \widehat{W}_u(p_X^-, p_X^+, E_\ell; k) \widehat{F}(p_X^+ - k) + \dots$$

- Fit parameters:  $|V_{tb}V_{ts}^*|^2 m_b^2$ ,  $|V_{ub}|^2$ ,  $\widehat{F}(\lambda x) = \frac{1}{\lambda} [\sum_{n=0}^{\infty} c_n f_n(x)]^2$ 
  - Theory Input:  $\widehat{W}_i(\dots; k)$  computed to (N)NNL'+NNLO in 1S scheme
  - Factorized shape function:

$$S(\omega, \mu_\Lambda) = \int dk \widehat{C}_0(\omega - k, \mu_\Lambda) \widehat{F}(k)$$

$\widehat{F}(k)$  nonperturbative part

- Determines peak region
- Fit from data

$\widehat{C}_0(\omega, \mu_\Lambda)$  perturbative part

- Generates perturbative tail with correct  $\mu_\Lambda$  dependence



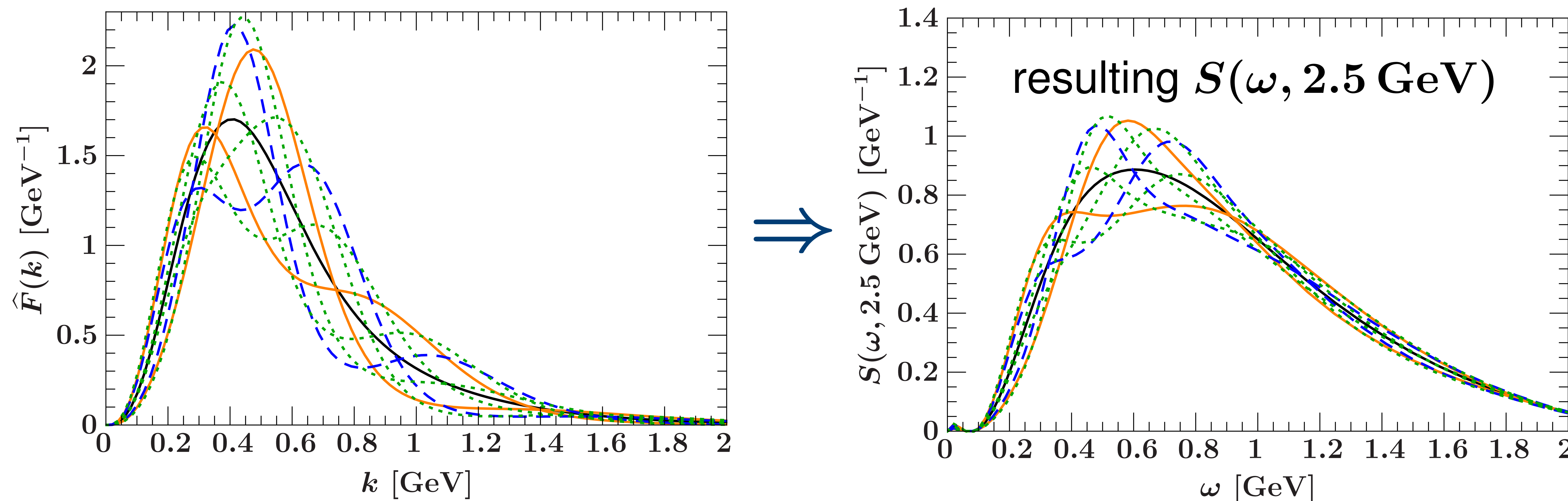
# Theory side of Global Fits

- **SIMBA** master formulae:

$$d\Gamma_s = |V_{tb}V_{ts}^*|^2 m_b^2 |C_7^{\text{incl}}|^2 \int dk \widehat{W}_{77}(E_\gamma; k) \widehat{F}(m_B - 2E_\gamma - k) + \dots$$

$$d\Gamma_u = |V_{ub}|^2 \int dk \widehat{W}_u(p_X^-, p_X^+, E_\ell; k) \widehat{F}(p_X^+ - k) + \dots$$

- Fit parameters:  $|V_{tb}V_{ts}^*|^2 m_b^2$ ,  $|V_{ub}|^2$ ,  $\widehat{F}(\lambda x) = \frac{1}{\lambda} [\sum_{n=0}^{\infty} c_n f_n(x)]^2$ 
  - Theory Input:  $\widehat{W}_i(\dots; k)$  computed to (N)NNL'+NNLO in 1S scheme
  - Factorized shape function:



# Theory side of Global Fits

Expand  $\hat{F}(k)$  into suitable orthonormal basis

$$\hat{F}(\lambda x) = \frac{1}{\lambda} \left[ \sum_{n=0}^{\infty} c_n f_n(x) \right]^2$$

$$\int dk \hat{F}(k) = \sum_{n=0}^{\infty} c_n^2 = 1$$

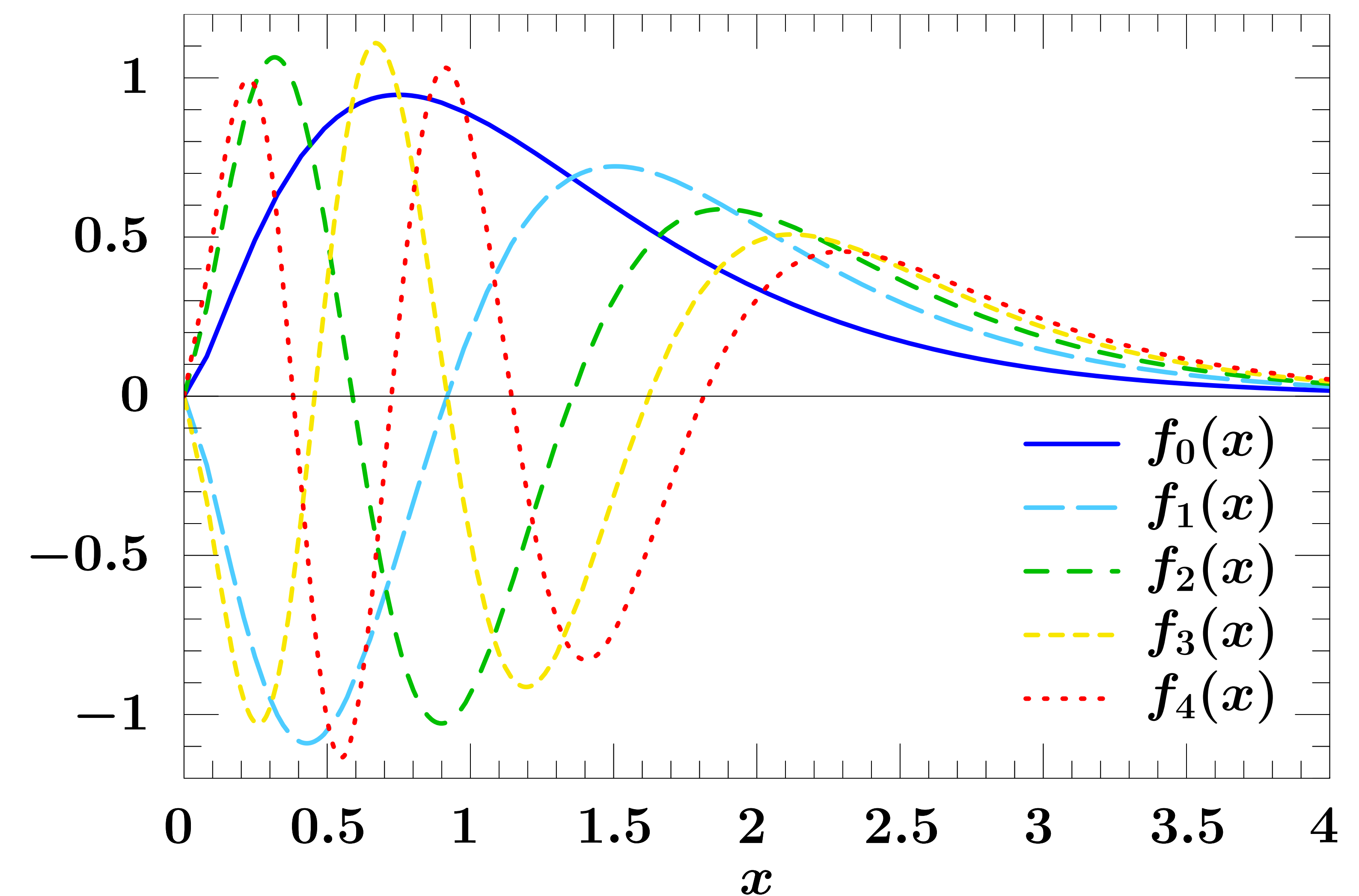
- Provides model-independent description

Fit for  $\hat{F}(k)$  by fitting basis coefficients  $c_n$

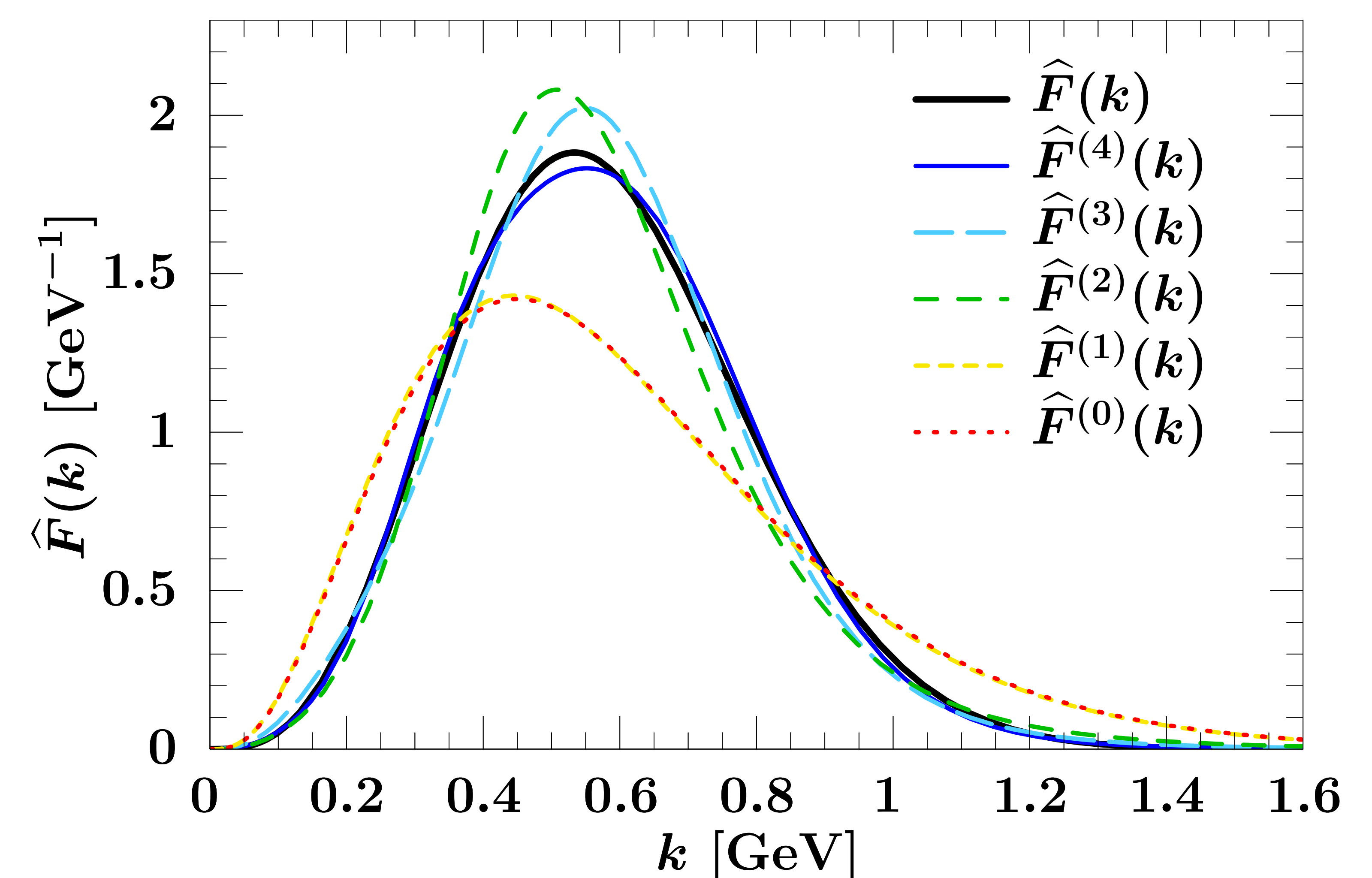
- Experimental uncertainties and correlations can be properly captured in covariance matrix of fitted coefficients  $c_n$

⇒ Allows for *data-driven*, reliable estimation of SF uncertainties

## Basis functions



## Expansion of Gaussian $\hat{F}(k)$





# Theory side of Global Fits

$$\hat{F}(\lambda x) = \frac{1}{\lambda} \left[ \sum_{n=0}^N c_n f_n(x) \right]^2$$

In practice, series must be truncated

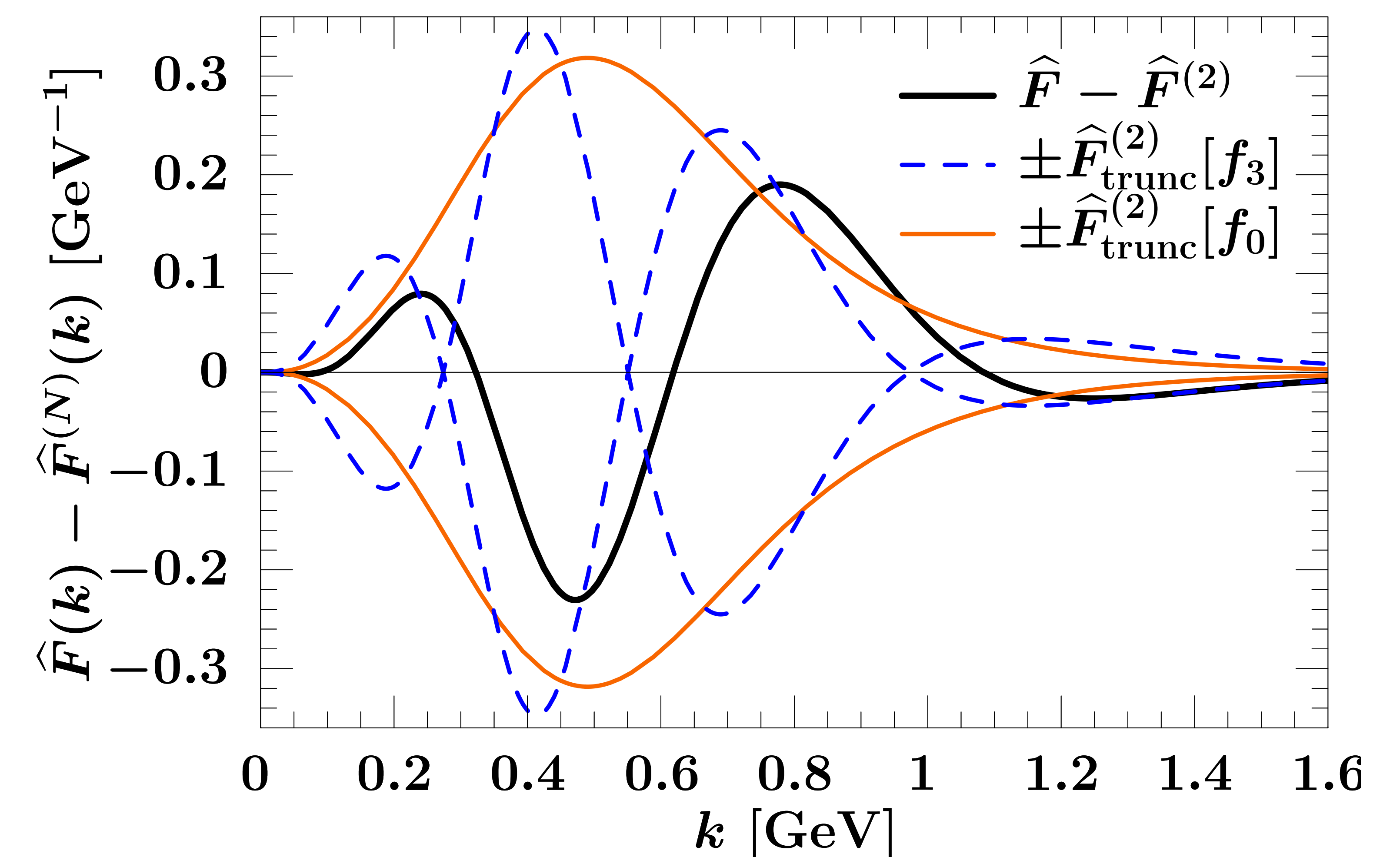
- Induces residual basis (model) dependence

- Truncation error scales as  $1 - \sum_{n=0}^N c_n^2$

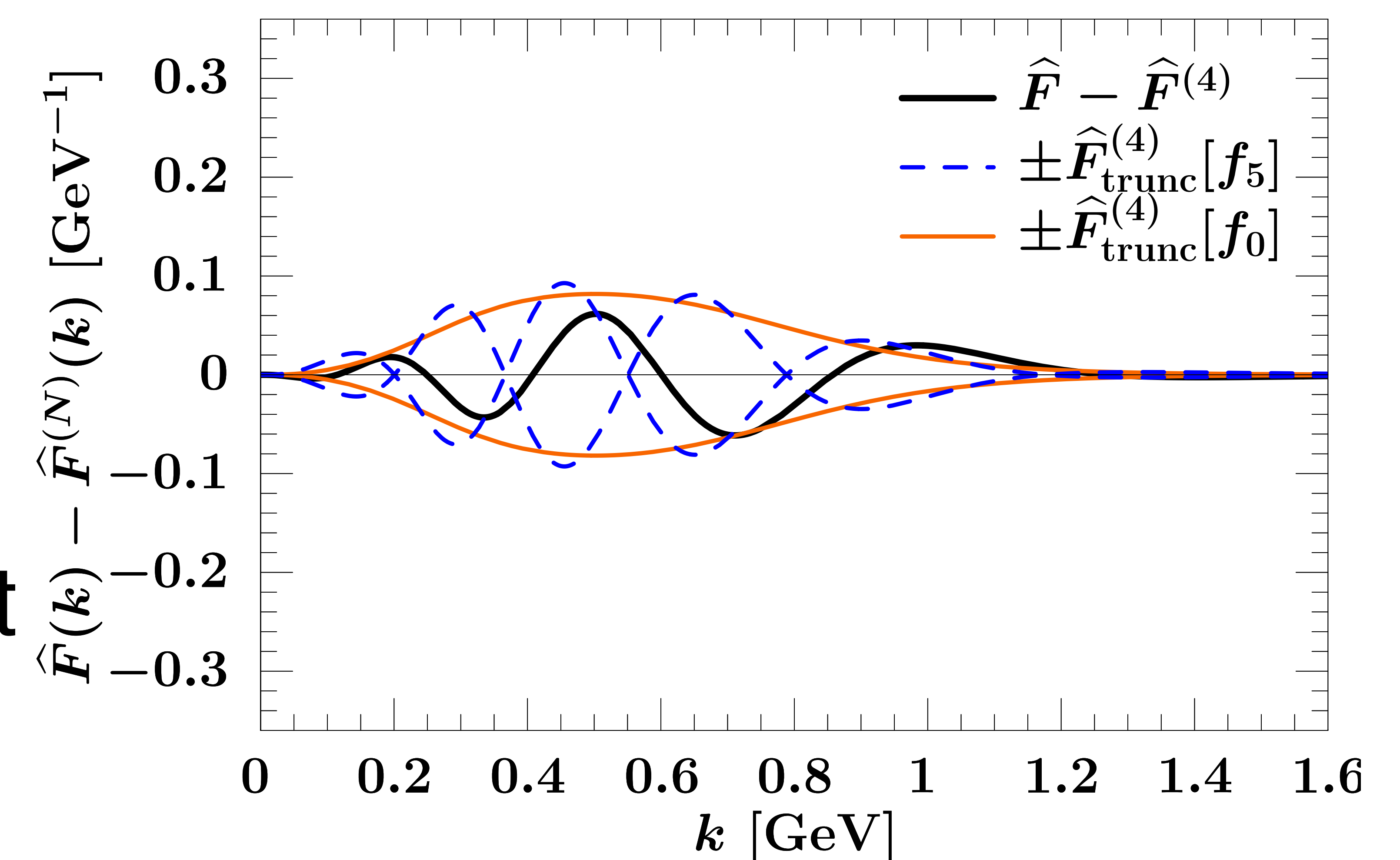
In practice most complications are in choosing good basis ( $\lambda$ ) and  $N$

- Want basis so series converges quickly but still unbiased (e.g. iterate)
- Choose  $N$  large enough so truncation error is smaller than to exp. uncertainties, but small enough to have stable fit and not waste statistical power
- Add coefficients with more precise data

## Truncation error at $N = 2$



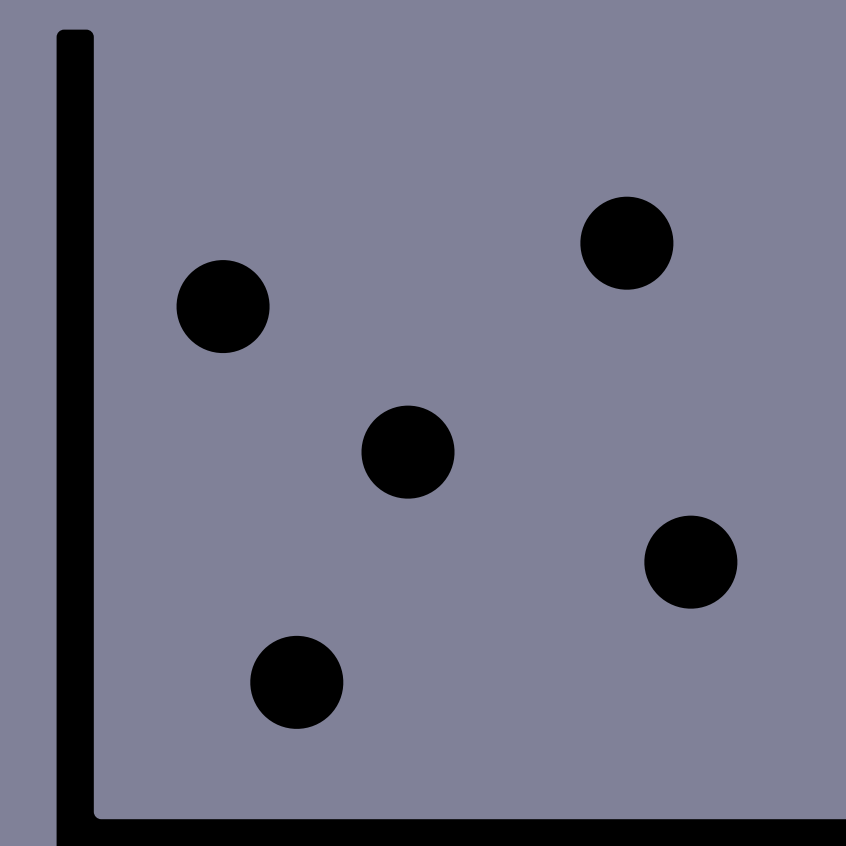
## Truncation error at $N = 4$



$$B \rightarrow X_s \gamma$$



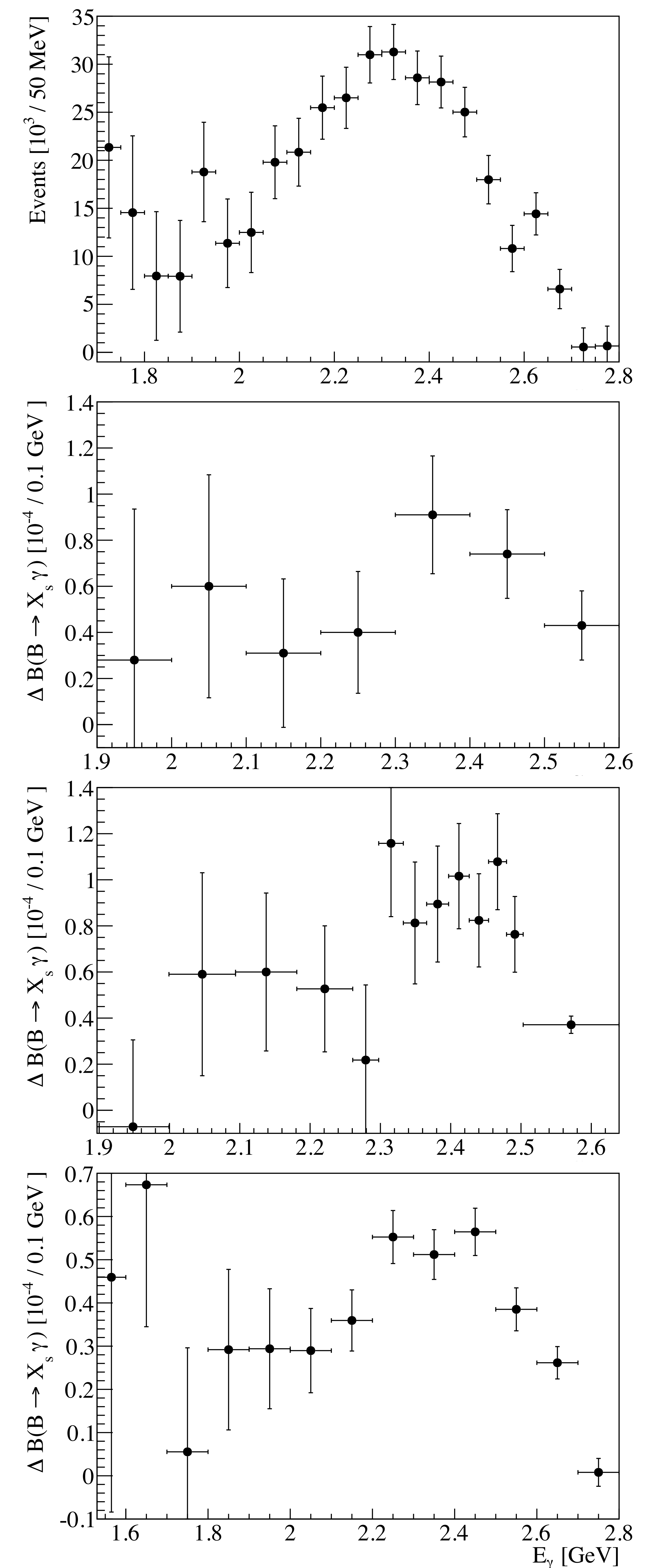
$m_b^{1S}$ ,  $|C_7^{eff}|$  & the Shape Function





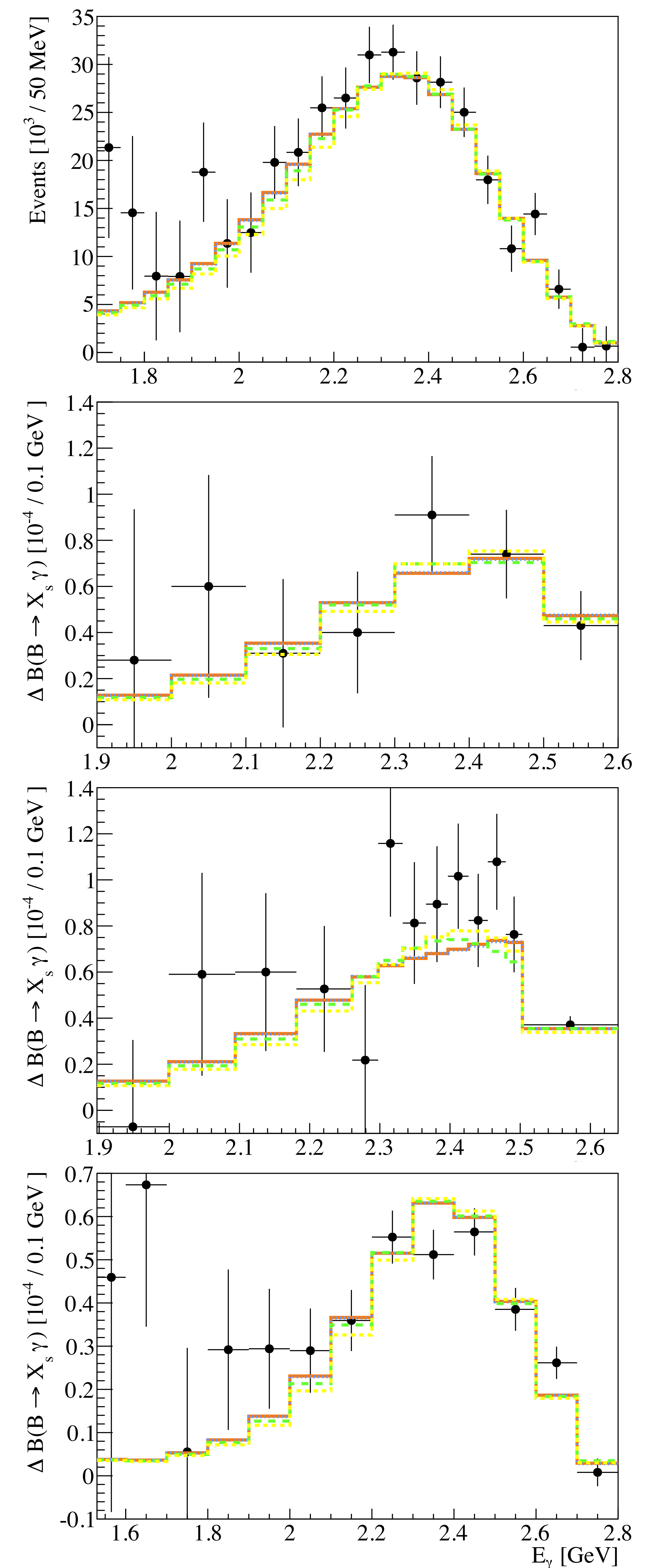
# Global Fit to $B \rightarrow X_s \gamma$

- Theory
  - NNLL' + NNLO
  - non- $C_7$  contributions fixed to SM
- Experimental Inputs
  - Belle Inclusive (in  $Y(4S)$  frame)
    - arXiv:0907.1384
  - BaBar hadronic (in  $B$  frame)
    - arXiv:0711.4889
  - BaBar sum-over-exclusive (in  $B$  frame)
    - hep-ex/0508004
  - BaBar inclusive (in  $Y(4S)$  frame)
    - arXiv:1207.5772



# Global Fit to $B \rightarrow X_s \gamma$

- Too few coefficients lead to clear bias and underestimates uncertainties
- Extracted  $|C_7^{\text{eff}} V_{tb} V_{ts}^*|$  consistent with SM





# Global Fit to $B \rightarrow X_s \gamma$

- Too few coefficients lead to clear bias and underestimates uncertainties
- Extracted  $|C_7^{\text{eff}} V_{tb} V_{ts}^*|$  consistent with SM

# Global Fit to $B \rightarrow X_s \gamma$

- Perturbative uncertainties:
  - Dominant source of uncertainties
  - Important to take into account correlations
  - Evaluated via set of profile scale variations
  - Expected theory uncertainties of comparable size of fit uncertainties

(Illustration only)





# Open issues End of 2017

- Consistent treatment of charm contributions
  - Integrate out charm loops vs keeping charm dynamic
  - Include known massive results
    - In the end small effect, but good to get it right
- Four-quark shape functions
- Sub-leading shape functions
  - irrelevant for fit, but important for interpretation
- Fix fit strategy

# Open issues Now

DESY 14-xxx  
MIT-CTP xxxxx  
January 10, 2018 – 17:06

## Precision Global Analysis of $B \rightarrow X_s \gamma$

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Iain W. Stewart,<sup>4</sup> Frank J. Tackmann,<sup>5</sup> and Kerstin Tackmann<sup>5</sup>

(The SIMBA Collaboration)

<sup>1</sup>Karlsruher Institute of Technology, 76131 Karlsruhe, Germany

<sup>2</sup>Humboldt University of Berlin, D-12489 Berlin, Germany

<sup>3</sup>Lawrence Berkeley National Laboratory, University of California, Berkeley, California 94720, USA

<sup>4</sup>Center for Theoretical Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

<sup>5</sup>Deutsches Elektronen-Synchrotron (DESY), D-22607 Hamburg, Germany

We perform a model independent global fit to all available inclusive  $B \rightarrow X_s \gamma$  data from *BABAR* and *Belle*. We extract the normalization of the  $B \rightarrow X_s \gamma$  decay rate, which is a sensitive probe of physics beyond the standard model, together with the nonperturbative  $b$ -quark distribution function and the  $b$ -quark mass,  $m_b$ . Our theoretical framework consistently combines the correct descriptions

- ~~Consistent treatment of charm con~~
- ~~Integrate out charm loops vs keeping~~
- ~~Include known massive results~~
  - ~~In the end small effect, but good to get it right~~
- ~~Four-quark shape functions~~
- ~~Sub-leading shape functions~~
  - ~~irrelevant for fit, but important for interpretation~~
- ~~Fix fit strategy~~



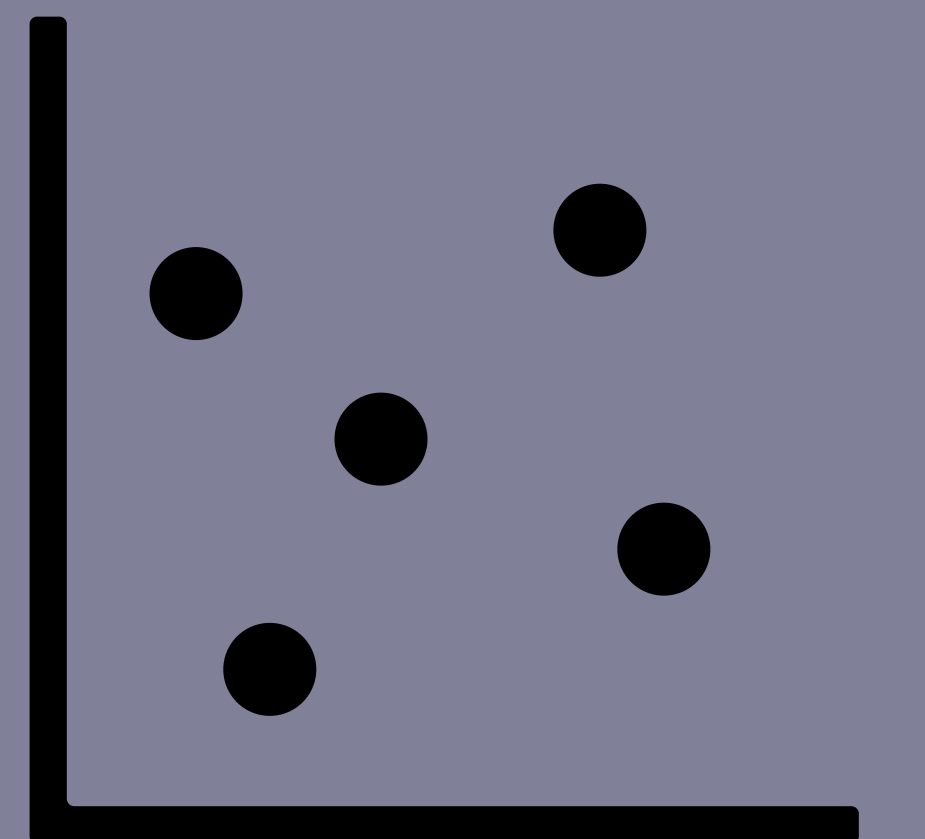
**Meet up in early January at KIT and finished all open items**  
**Will meet up again soon and write the  $B \rightarrow X_s \gamma$  paper and**  
**start working on  $B \rightarrow X_u l \nu$**



$B \rightarrow X_u / \nu$

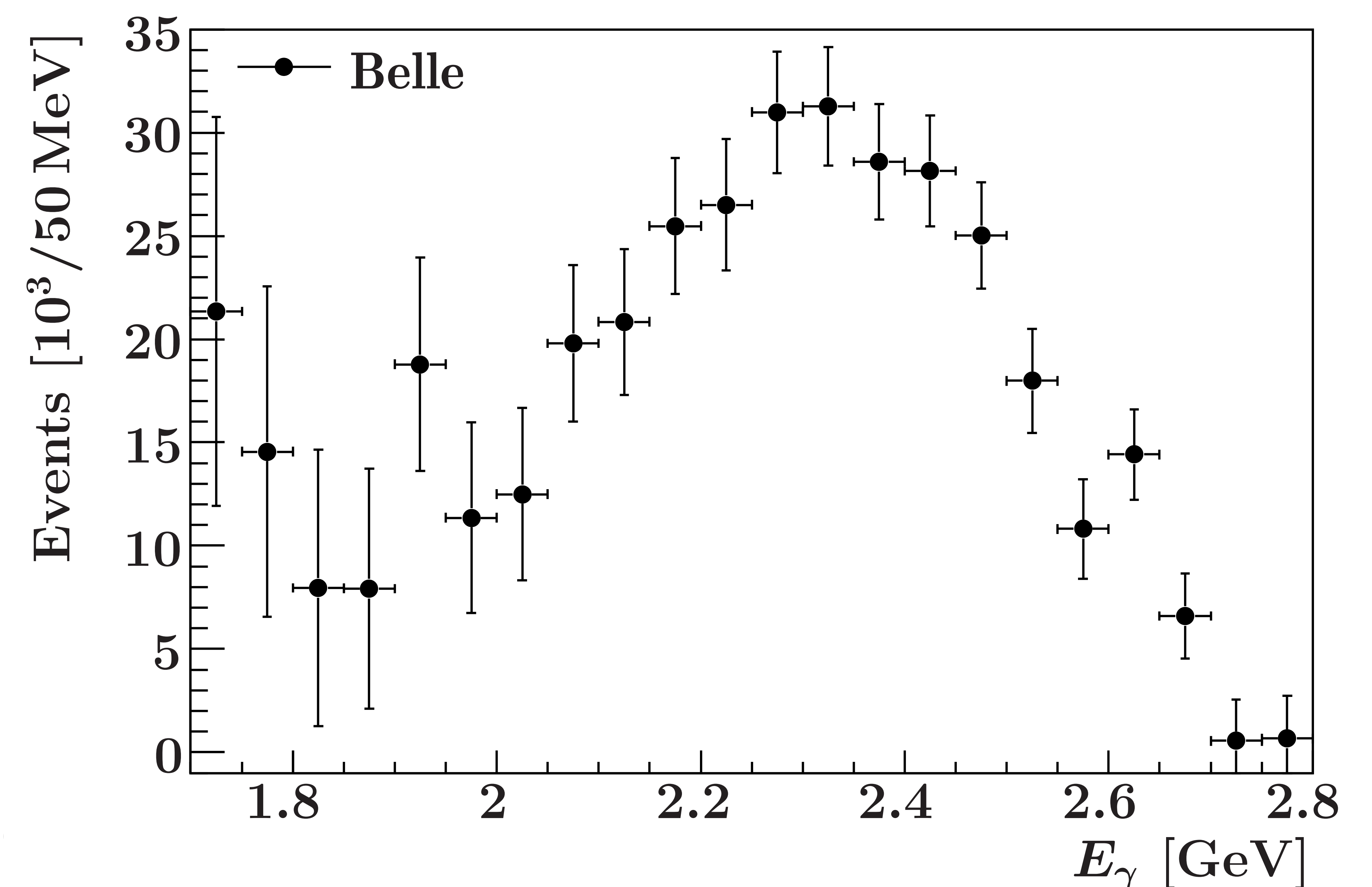
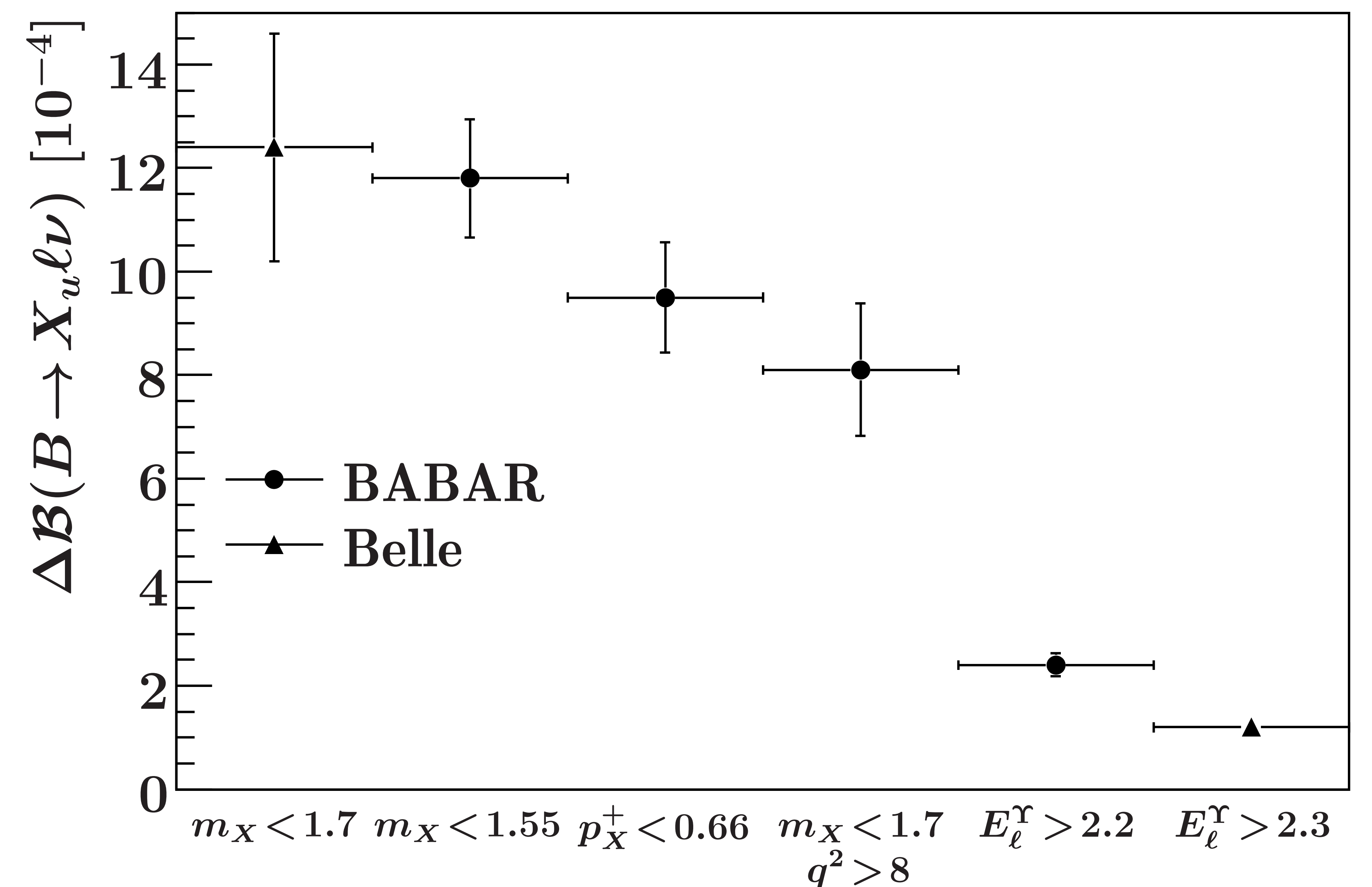


$m_b^{1S}, |C_7^{eff}|$  & the Shape Function



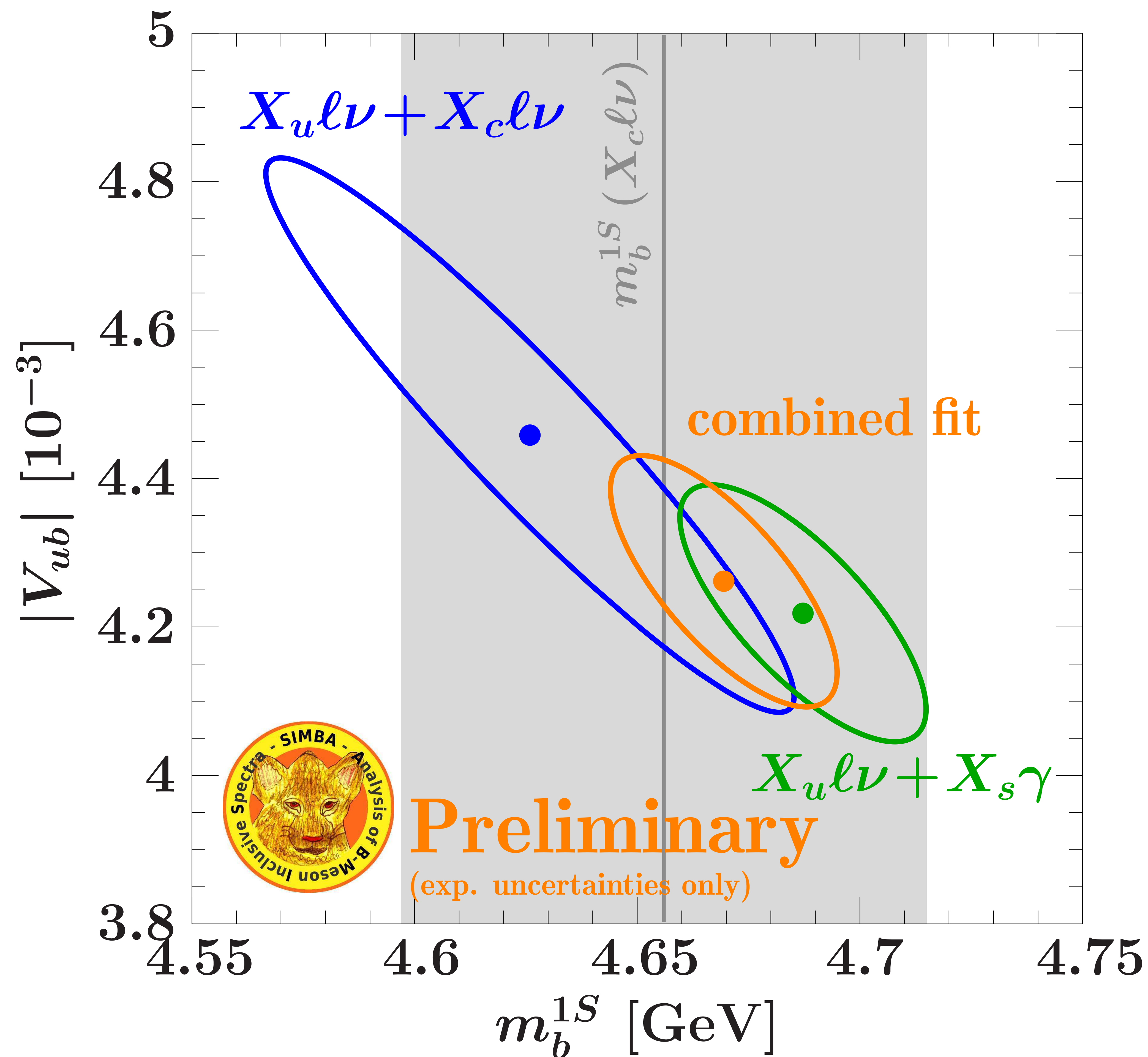
# Global $|V_{ub}|$ fit

- Theory
  - NLL' + NLO
  - ignoring sub-leading shape functions
- Experimental Inputs
  - $B \rightarrow X_u / \nu$  partial branching fractions
    - picked measurement for which we are sure enough that they have negligible (SF) model dependence
    - BaBar & Belle hadronic tag
    - BaBar & Belle lepton endpoint
  - $B \rightarrow X_s \gamma$  partial branching fractions
    - Belle inclusive (shown)
    - (old) BaBar sum-over-exclusive (not shown)
    - BaBar hadronic tag (not shown)
  - $B \rightarrow X_c / \nu$ 
    - $m_b^{1S}$ ,  $\lambda_1$  from moment fits

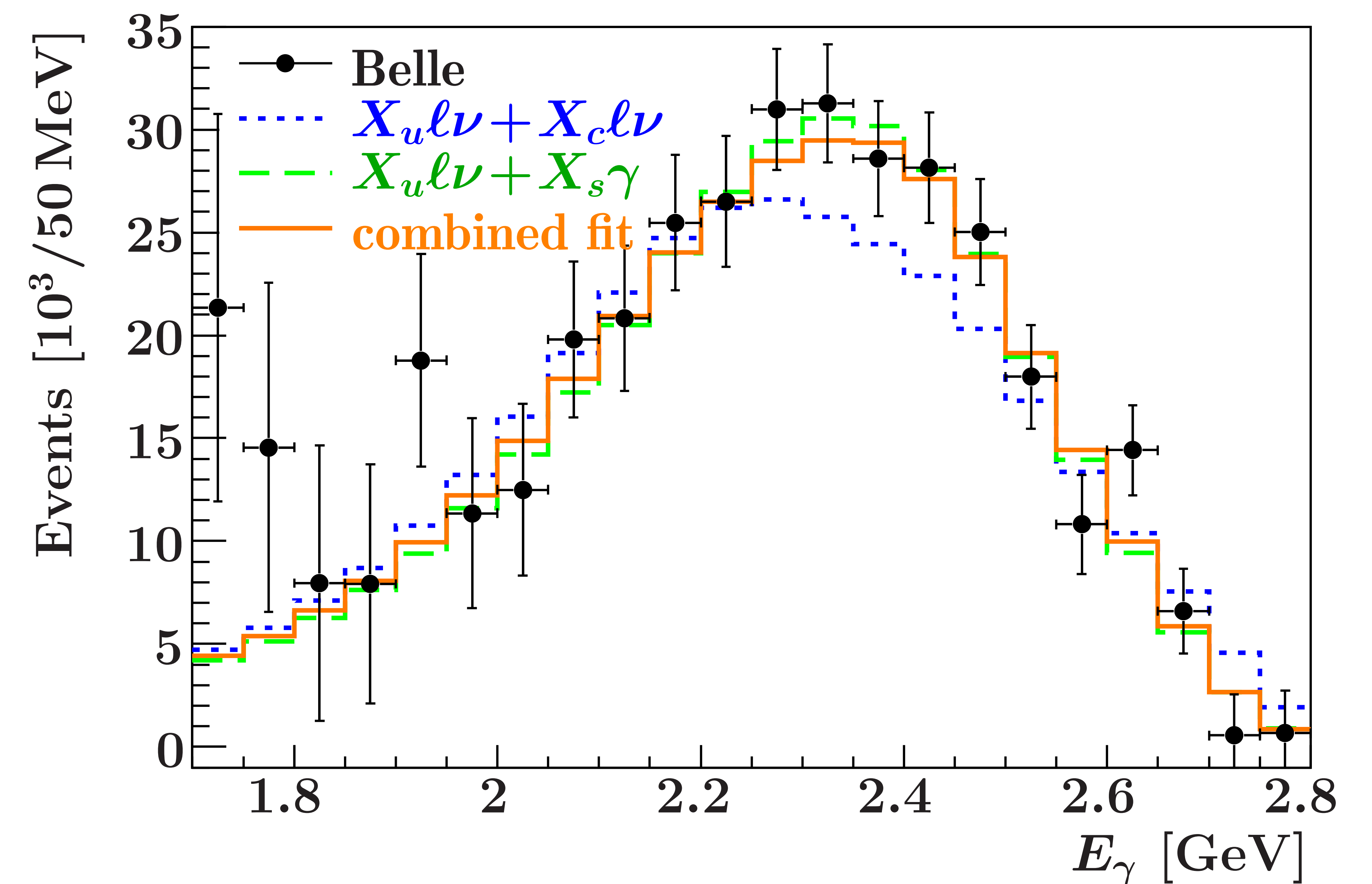
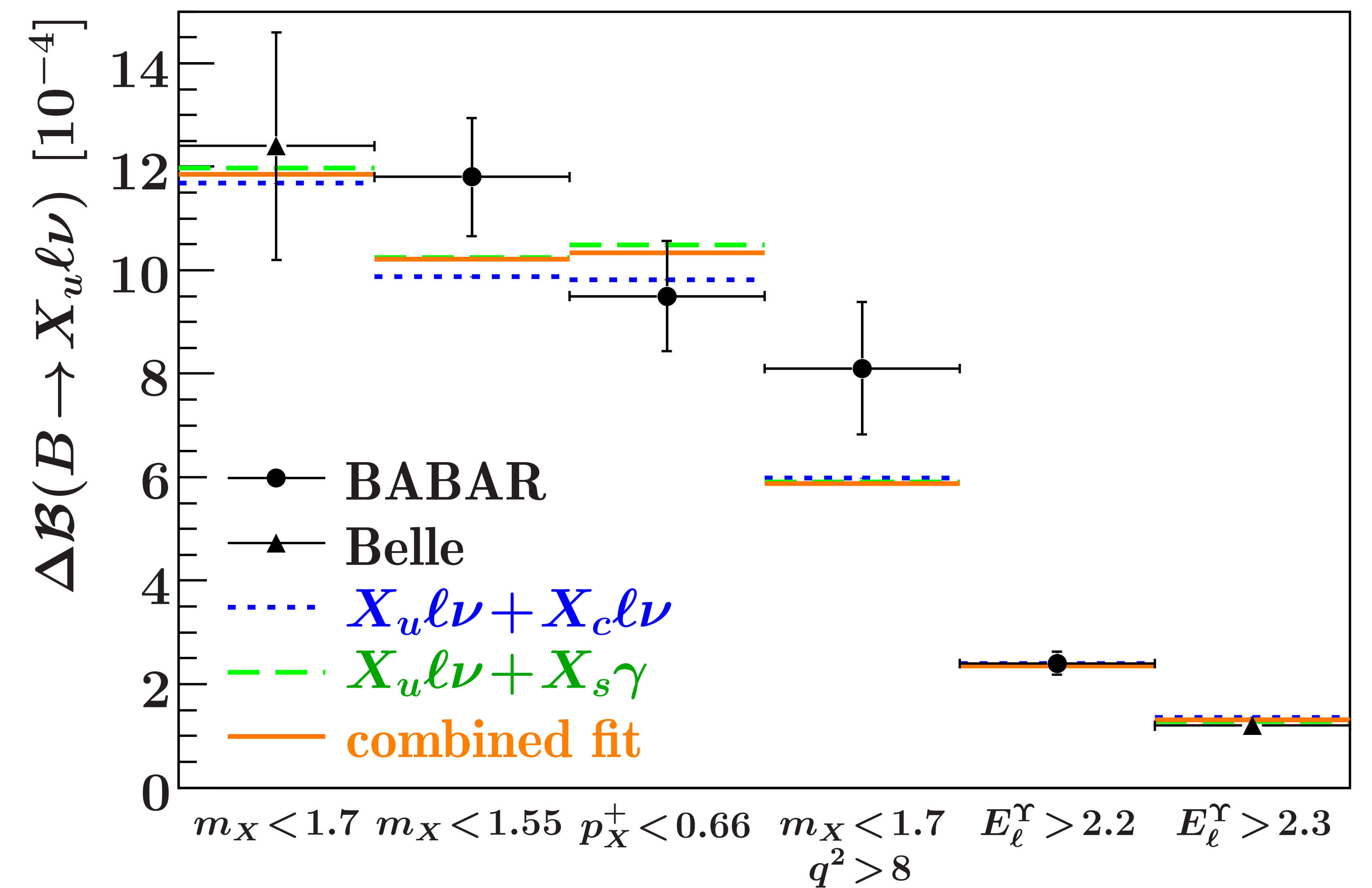




# Global $|V_{ub}|$ fit



- Parametric unc. (SF, mb) are part of fit, no pert. unc. included
- Without  $B \rightarrow X_s \gamma$ :  $\sim 10\%$  uncertainties on  $|V_{ub}|$
- With  $B \rightarrow X_s \gamma$ :  $\sim 5\%$  uncertainty, but also shift in  $|V_{ub}|$

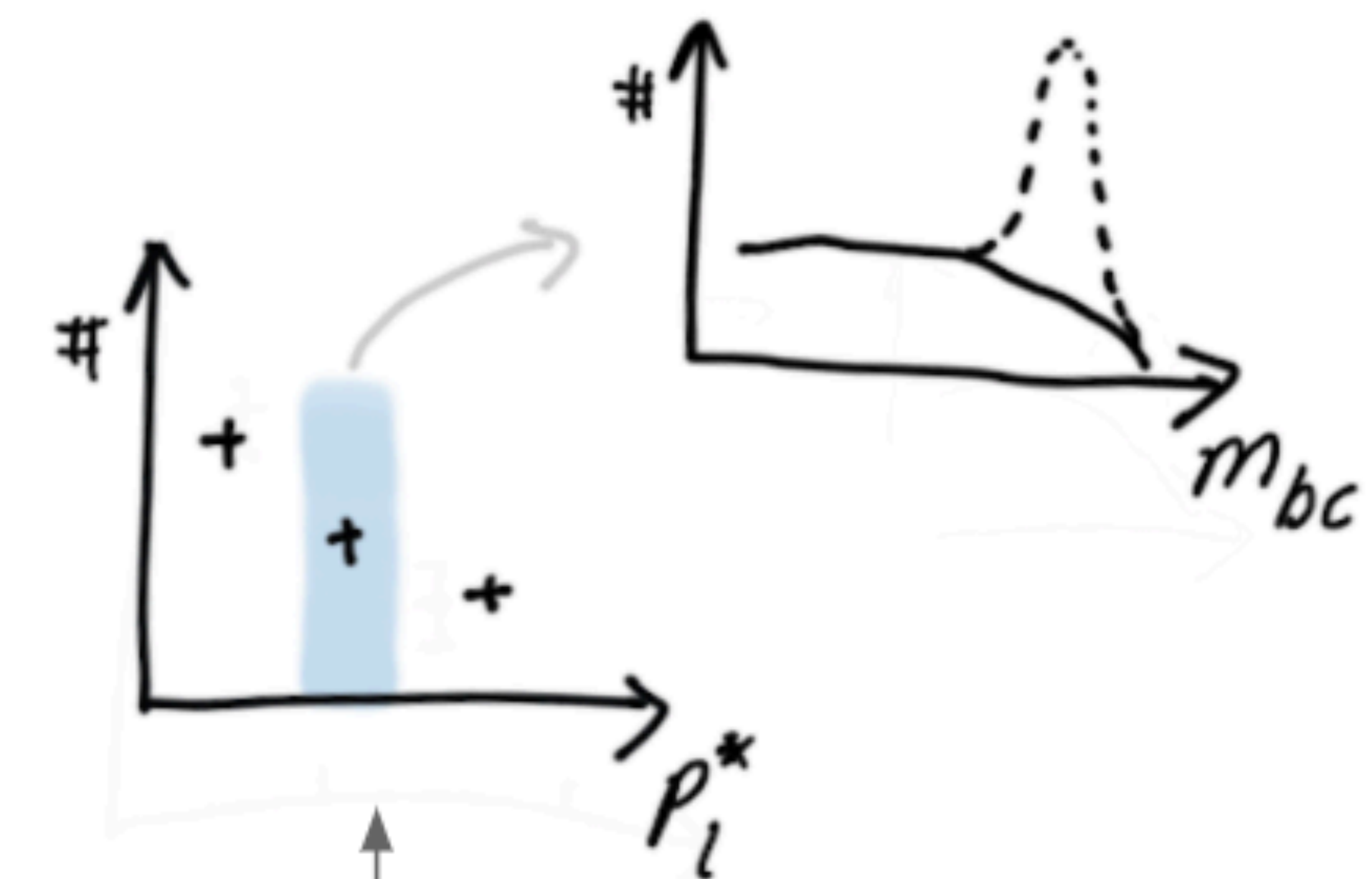
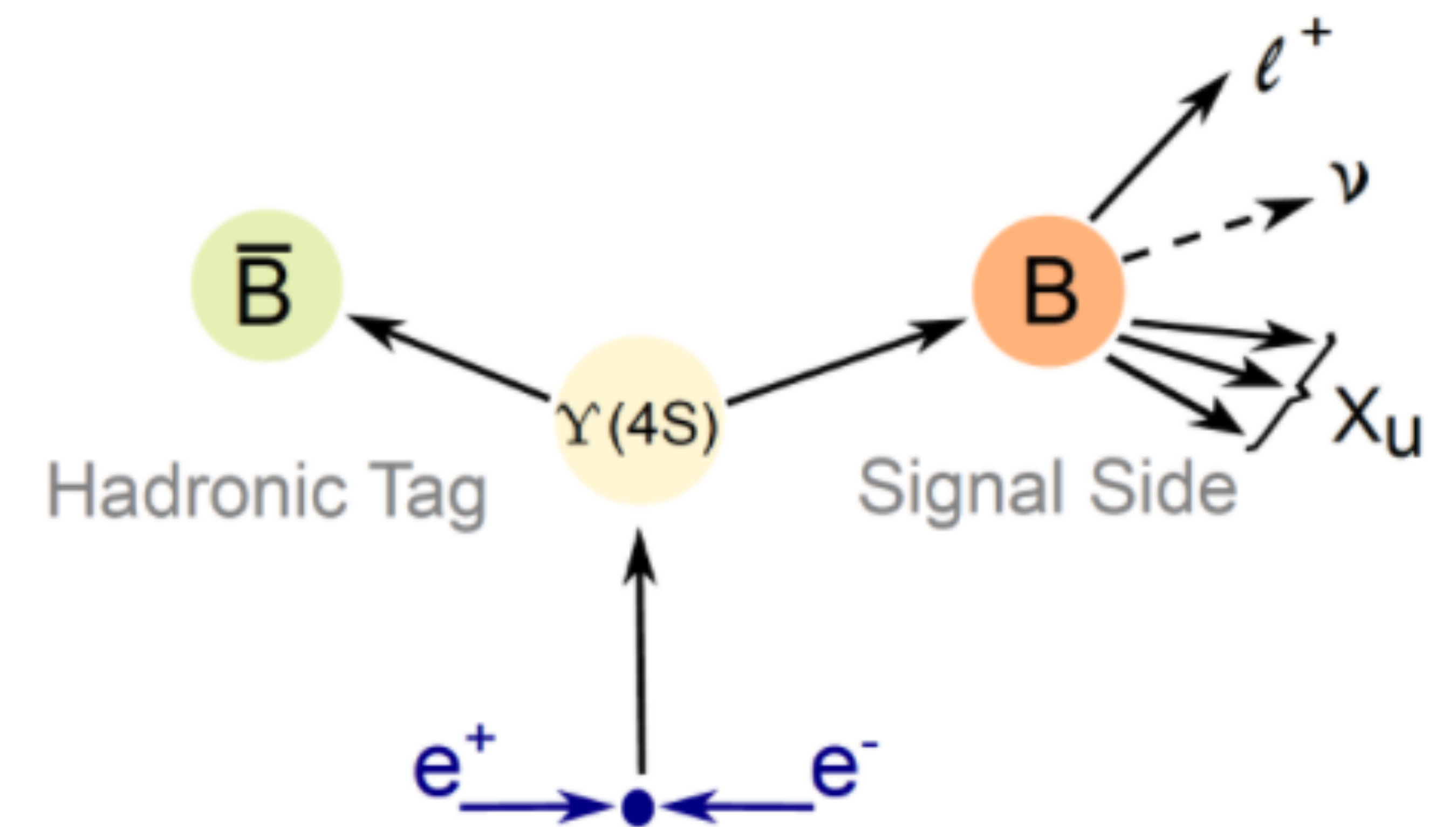




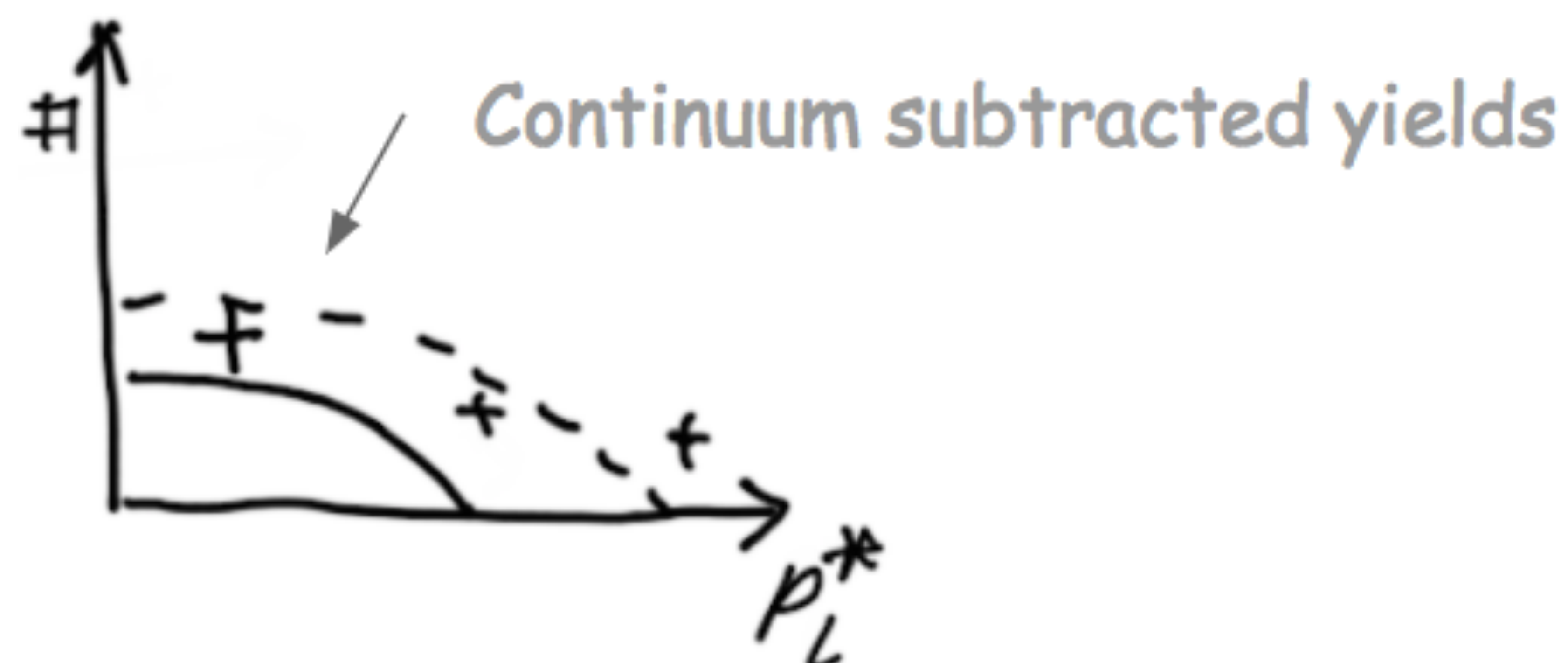
# $B \rightarrow X_u / \nu$ differential Information

- Ongoing measurement with Belle to measure  $E_l, p_{\pm}, m_X, q^2$

- Uses NeuroBayes hadronic tagging
  - Eff. (0.3/0.2% for  $B^+$  and  $B^0$ )
- Pre- and fine-selection finished:
  - 25% Signal efficiency with 2.2% Bkg retention
    - BaBar: 32% / 2.2%
- Analysis strategy:
  - Normalise against  $B \rightarrow X / \nu$
  - Subtract continuum and bad tags via  $m_{bc}$  fit



In each bin of an observable we fit  $m_{bc}$  to subtract continuum and not properly reconstructed B-meson events

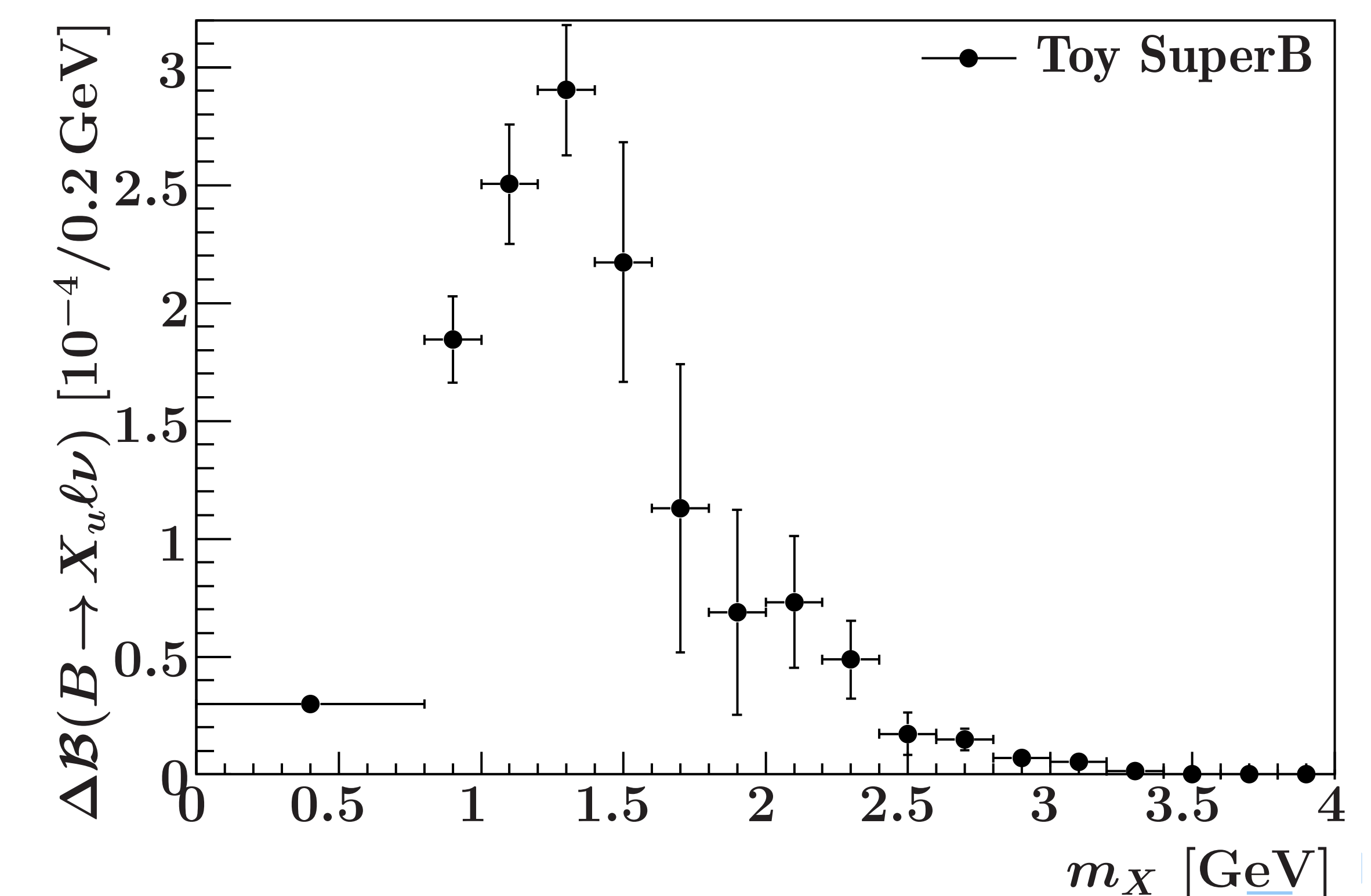
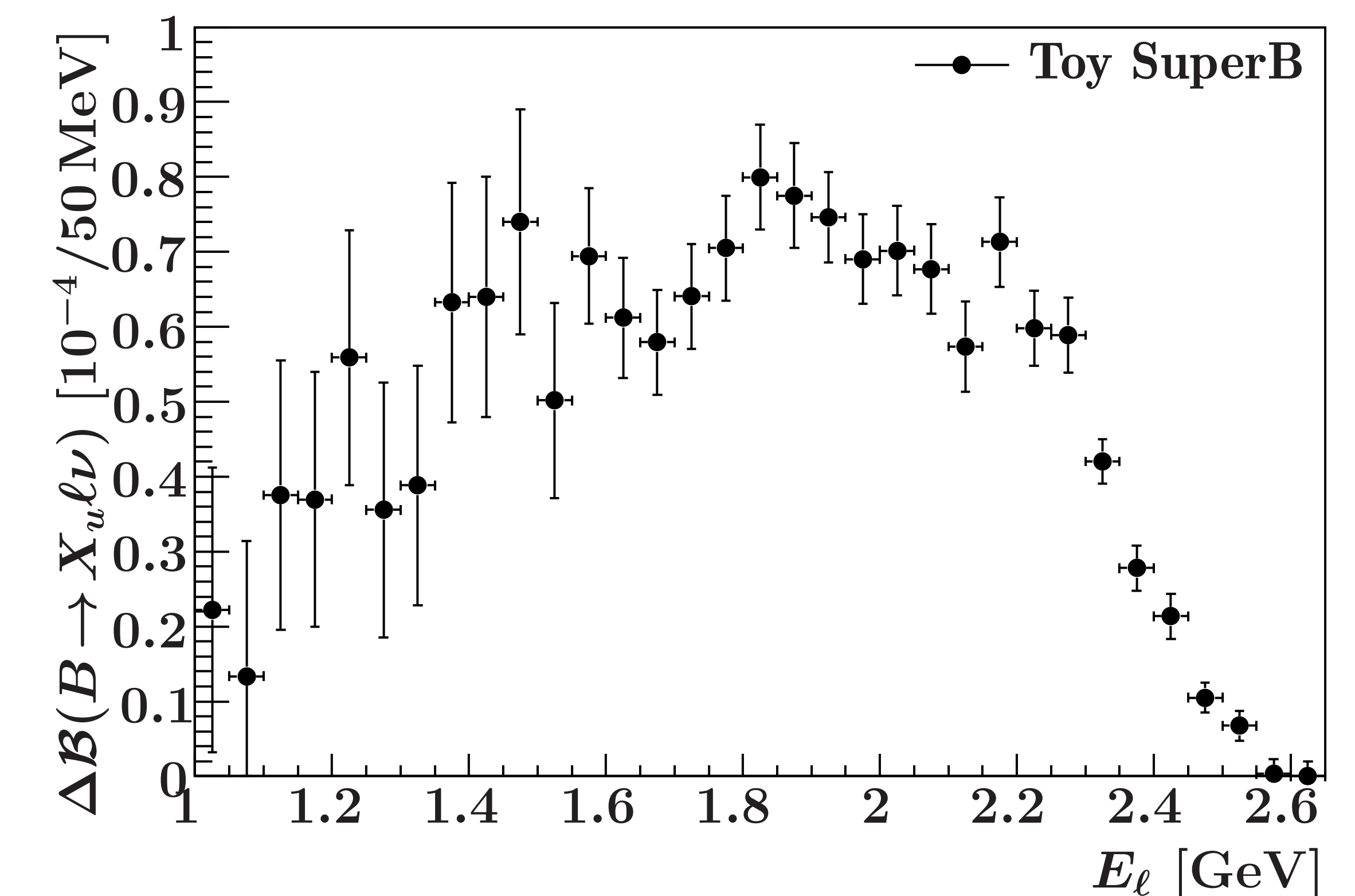
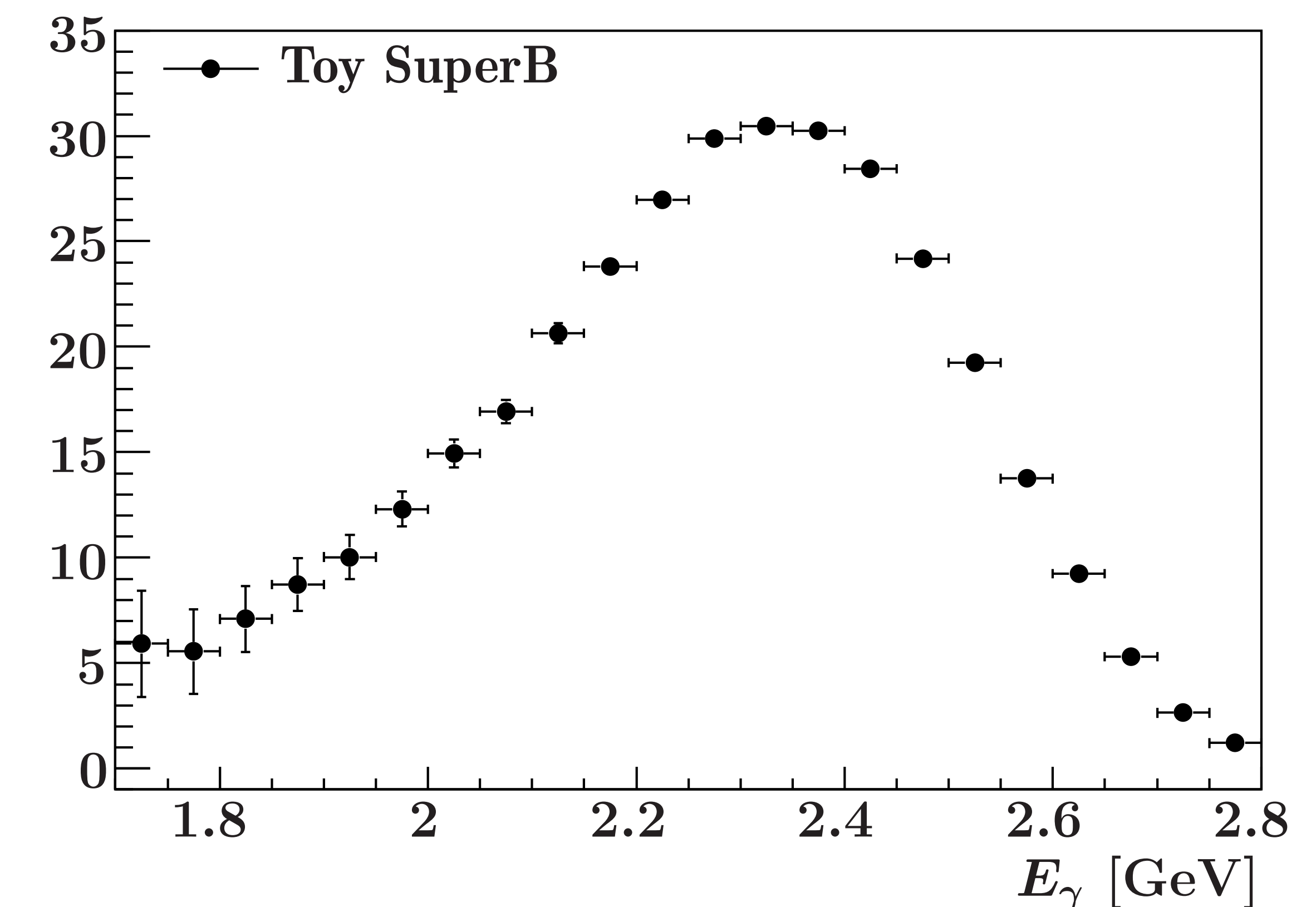




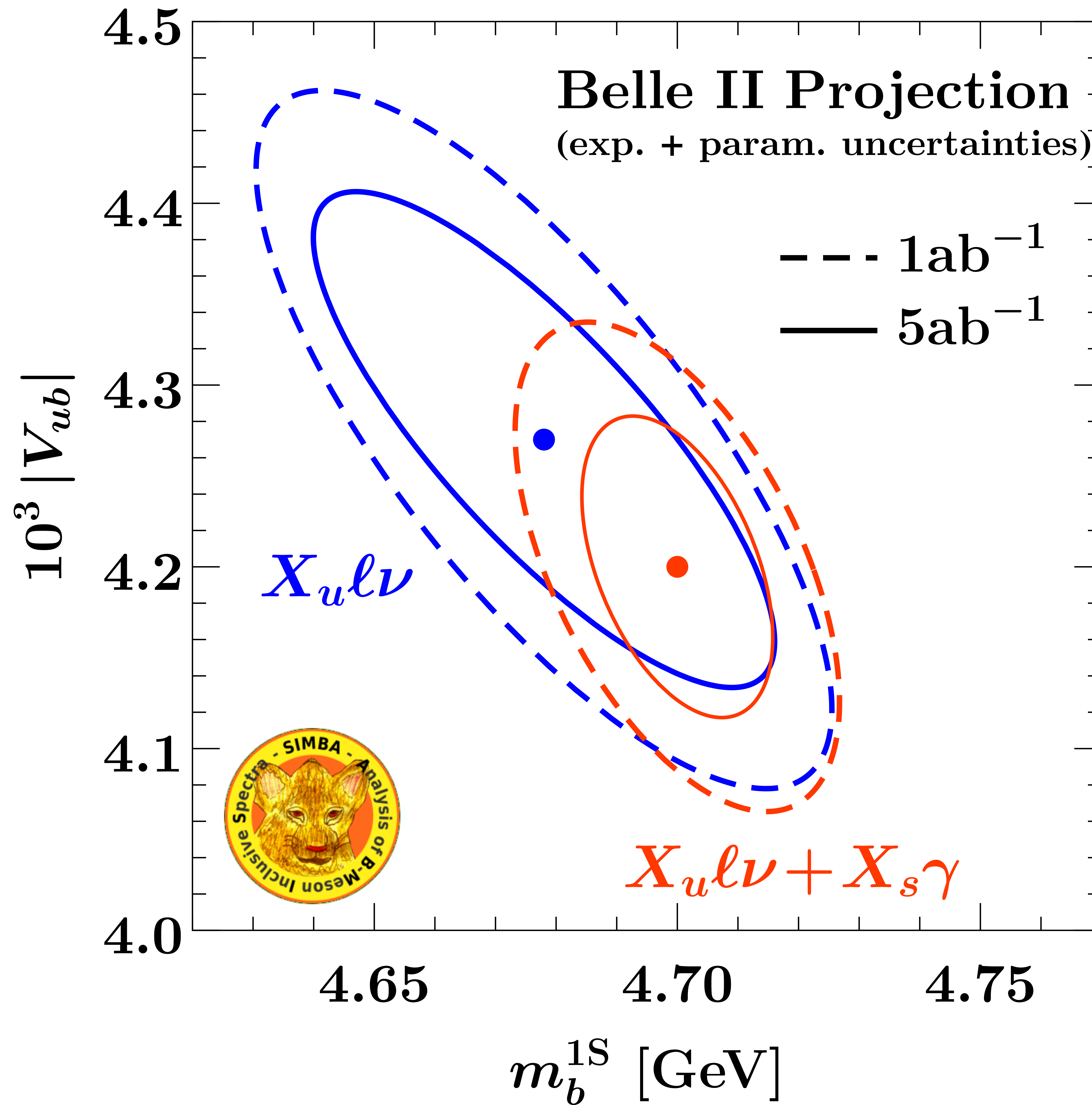
# Projections for Belle II



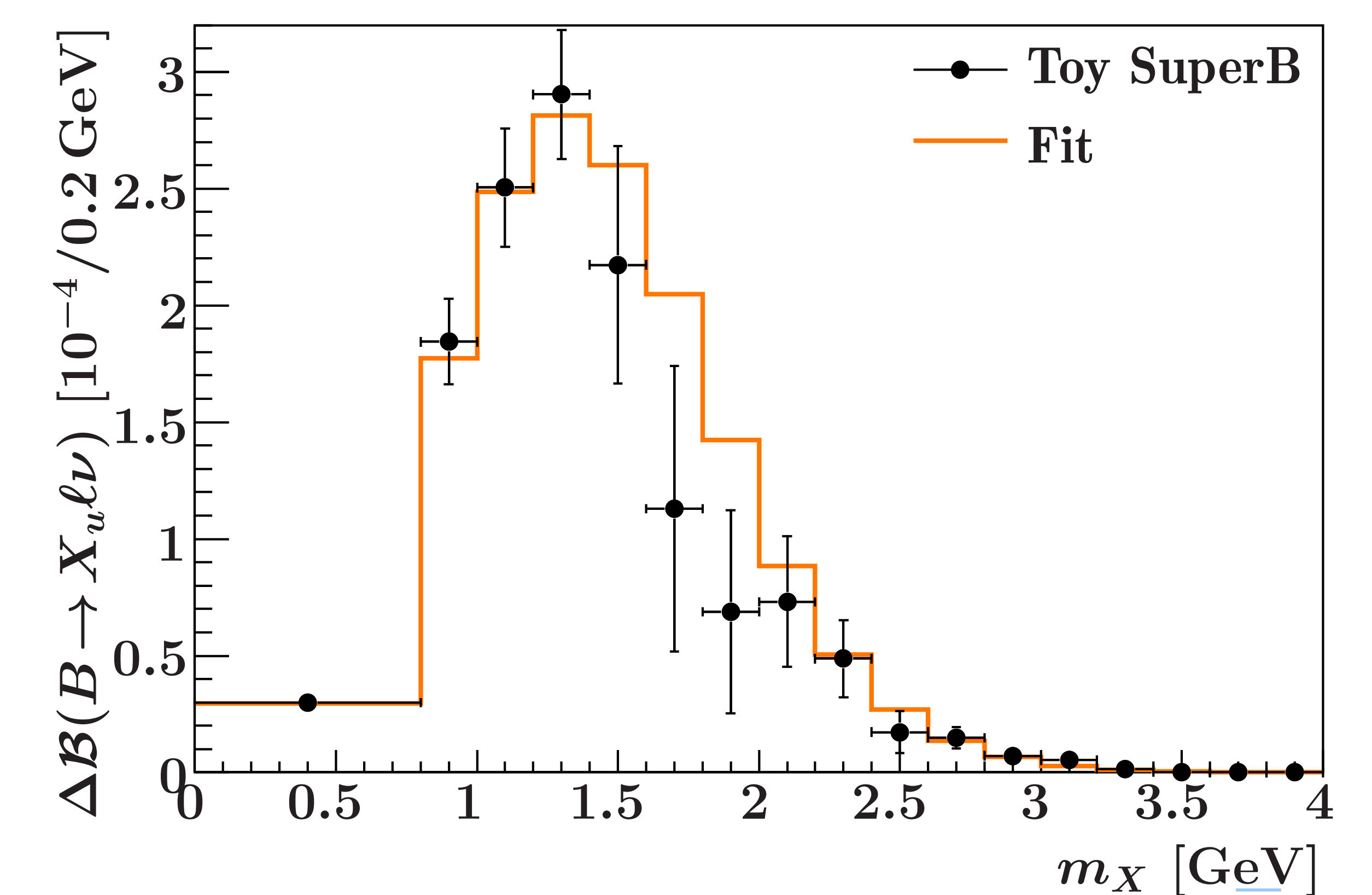
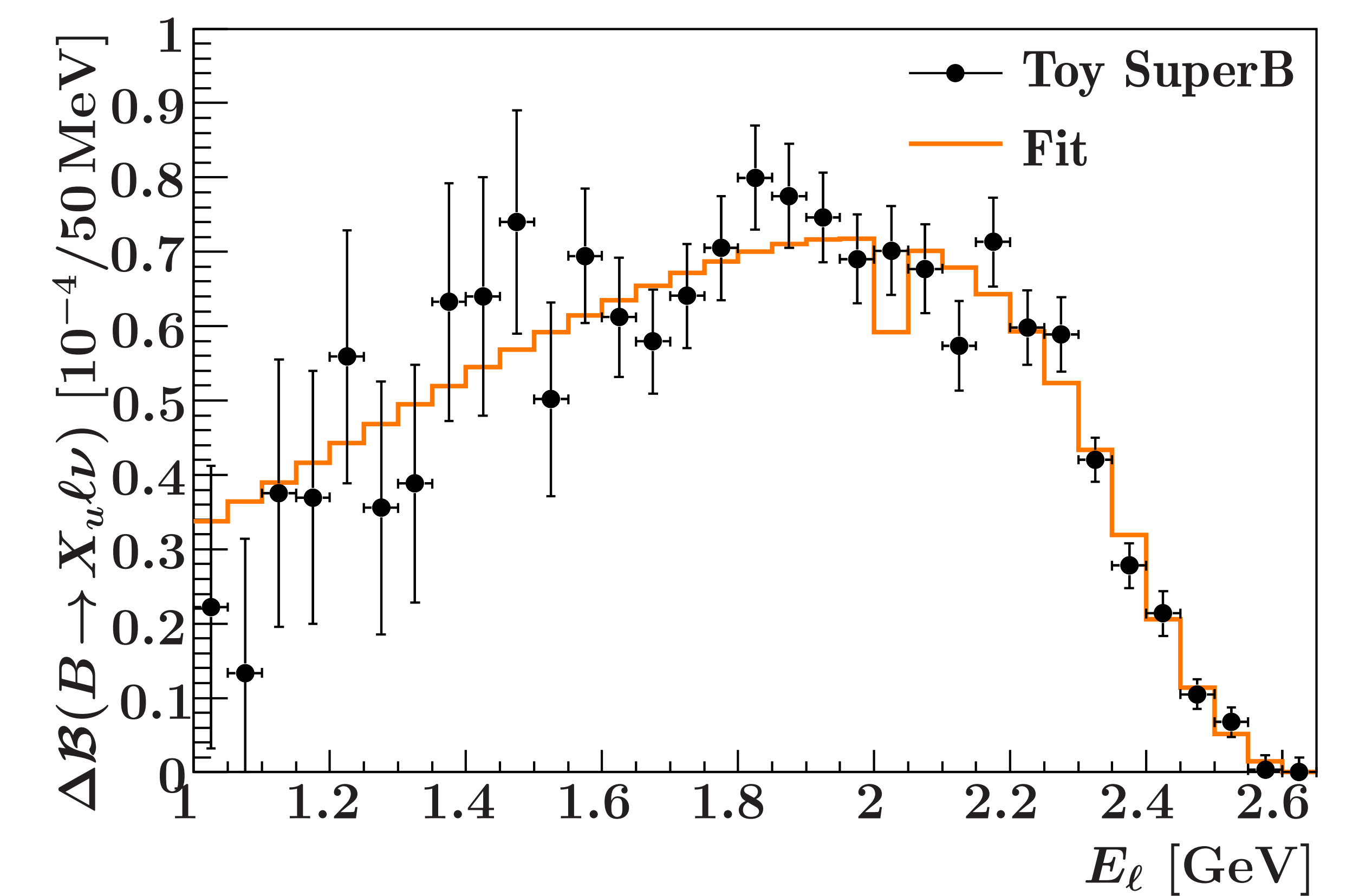
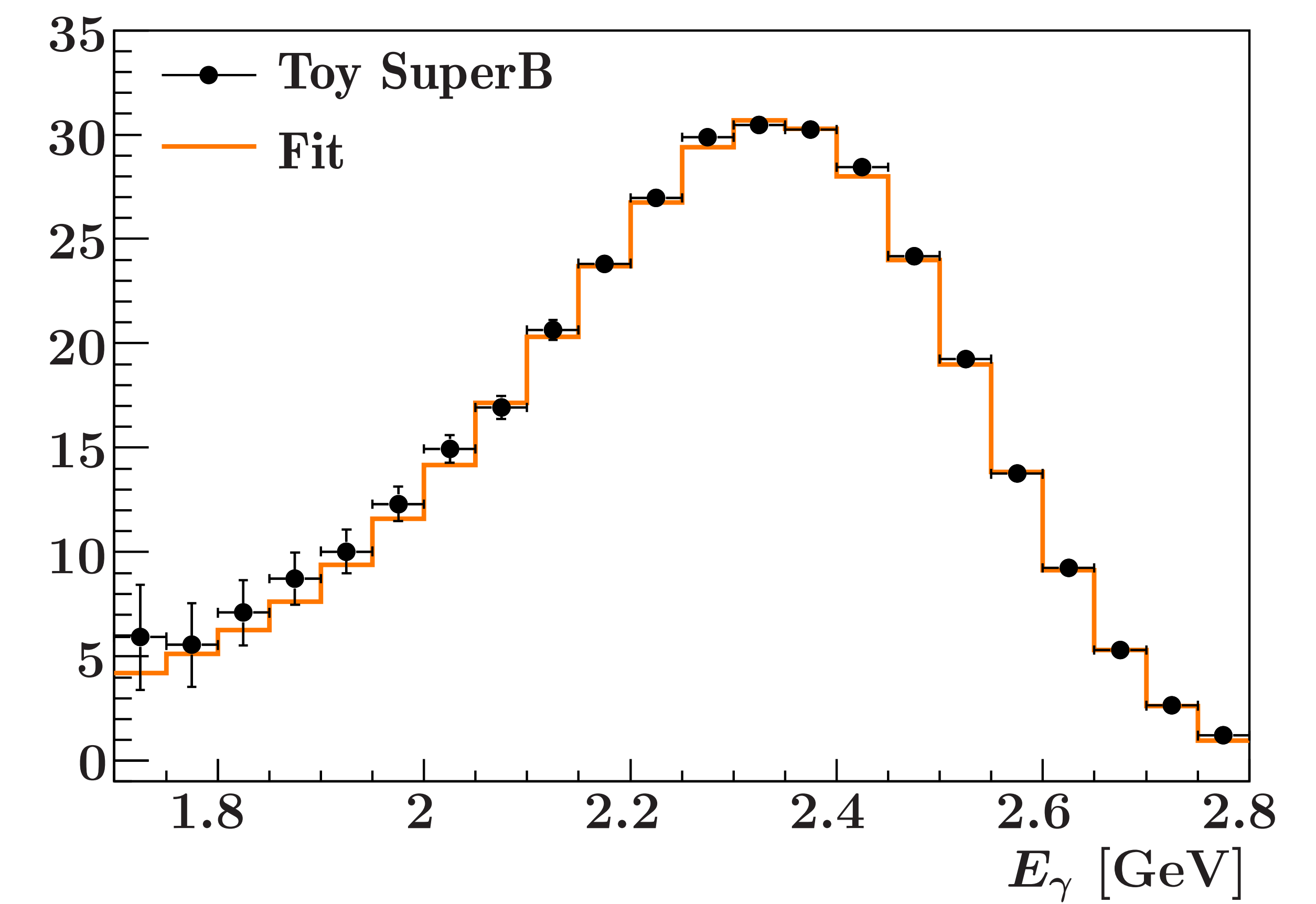
- Theory
  - NLL' + NLO
  - ignore subleading SFs
- Toy study
  - Generate  $m_X$ ,  $E_\ell$ , and  $E_\gamma$  from theory
  - Smearred from uncertainties and correlations inspired by BaBar hadronic tag analysis, Belle II hadronic tagging efficiency is much better by now
  - Target lumi: 1/ab, 5/ab
  - **Caveats:**
    - No resolution effects considered
    - No theory uncertainties included (!)
    - Not done with Belle II MC



# Projections for Belle II




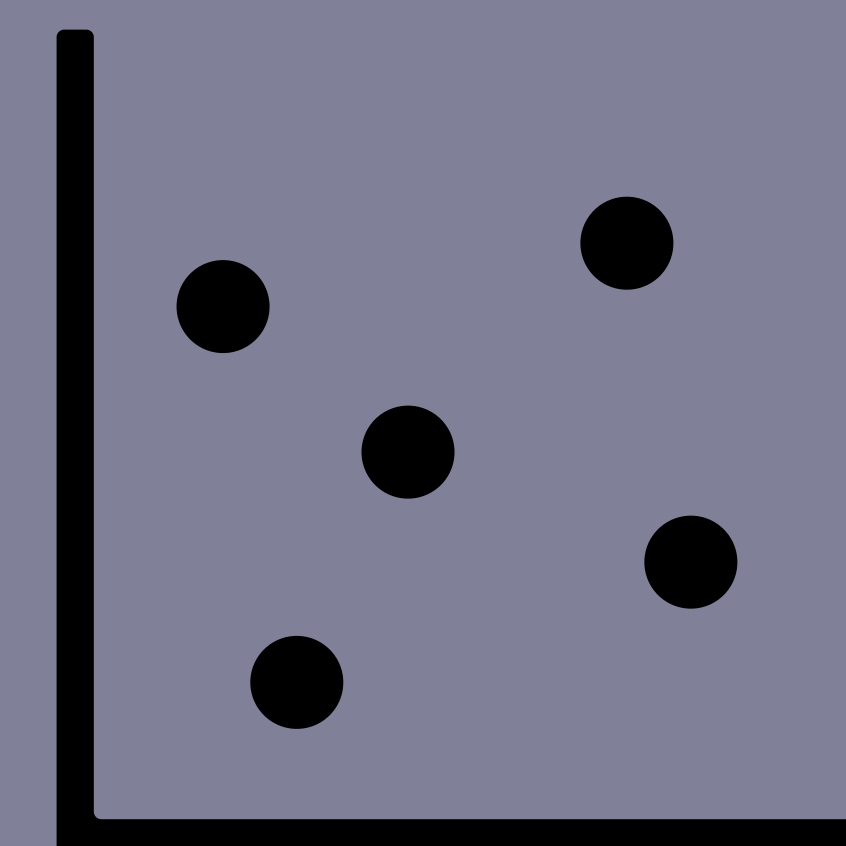
3.5%  
2.2%





$B \rightarrow X_s //$

  $C_9 \& C_{10}$

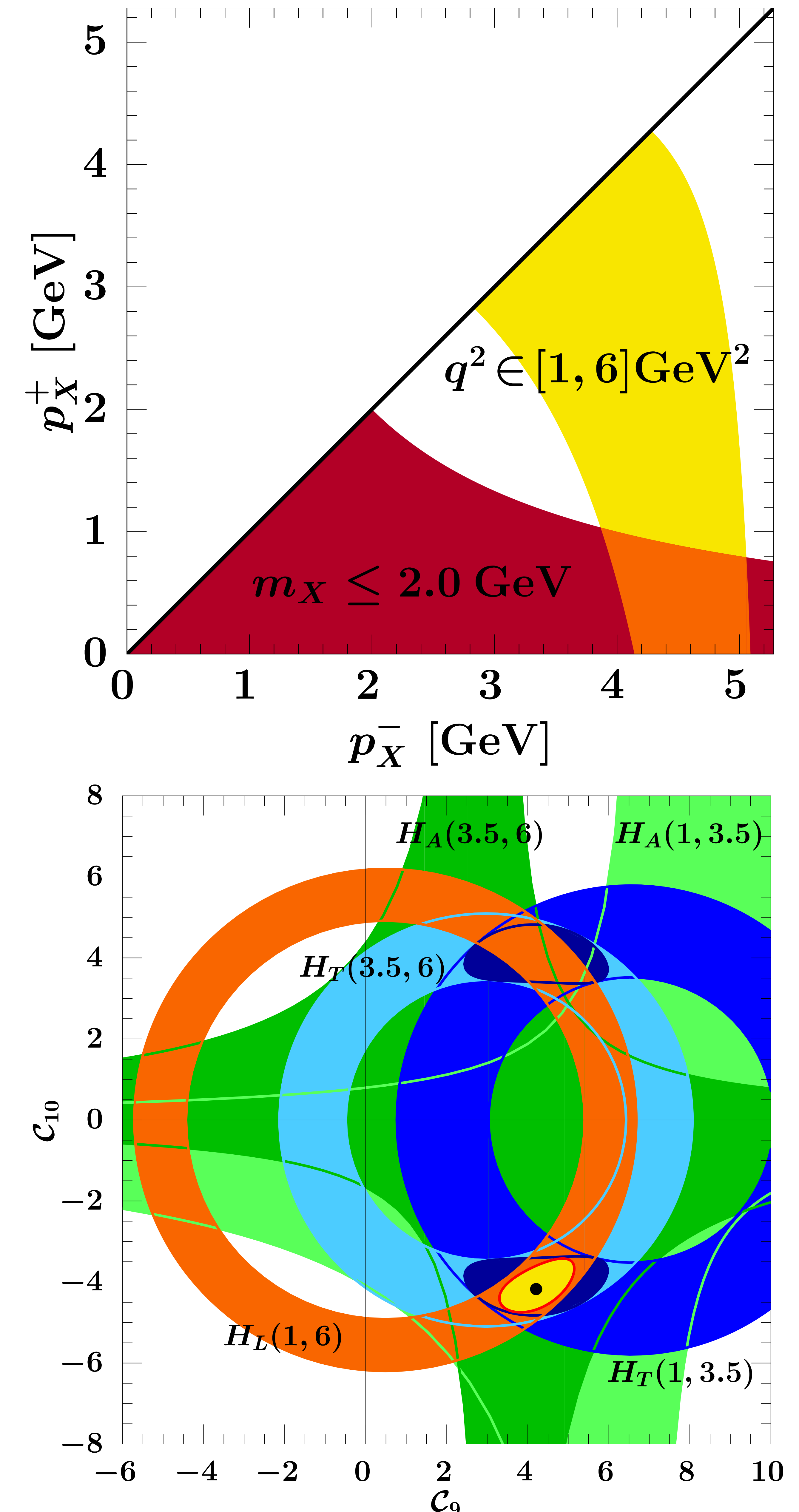


# $B \rightarrow X_s //$

- Experimental kinematic cuts for  $B \rightarrow X_s //$ 
  - $1 < q^2 < 6 \text{ GeV}^2$ ,  $m_X < m_X^{\text{cut}} \sim 2 \text{ GeV}$
  - Unavoidable to suppress huge  $B \rightarrow X_c // \nu \rightarrow X_s // \nu \nu$  background
  - Shape function effects must be taken into account to retain NP sensitivity
- Helicity decomposition for inclusive rate
  - [Lee, Ligeti, Stewart, Tackmann (2008)]

$$\frac{d^3\Gamma}{dp_X^+ dp_X^- dz} = \frac{3}{8} \left[ (1 + z^2) H_T(p_X^\pm) + 2z H_A(p_X^\pm) + 2(1 - z^2) H_L(p_X^\pm) \right]$$

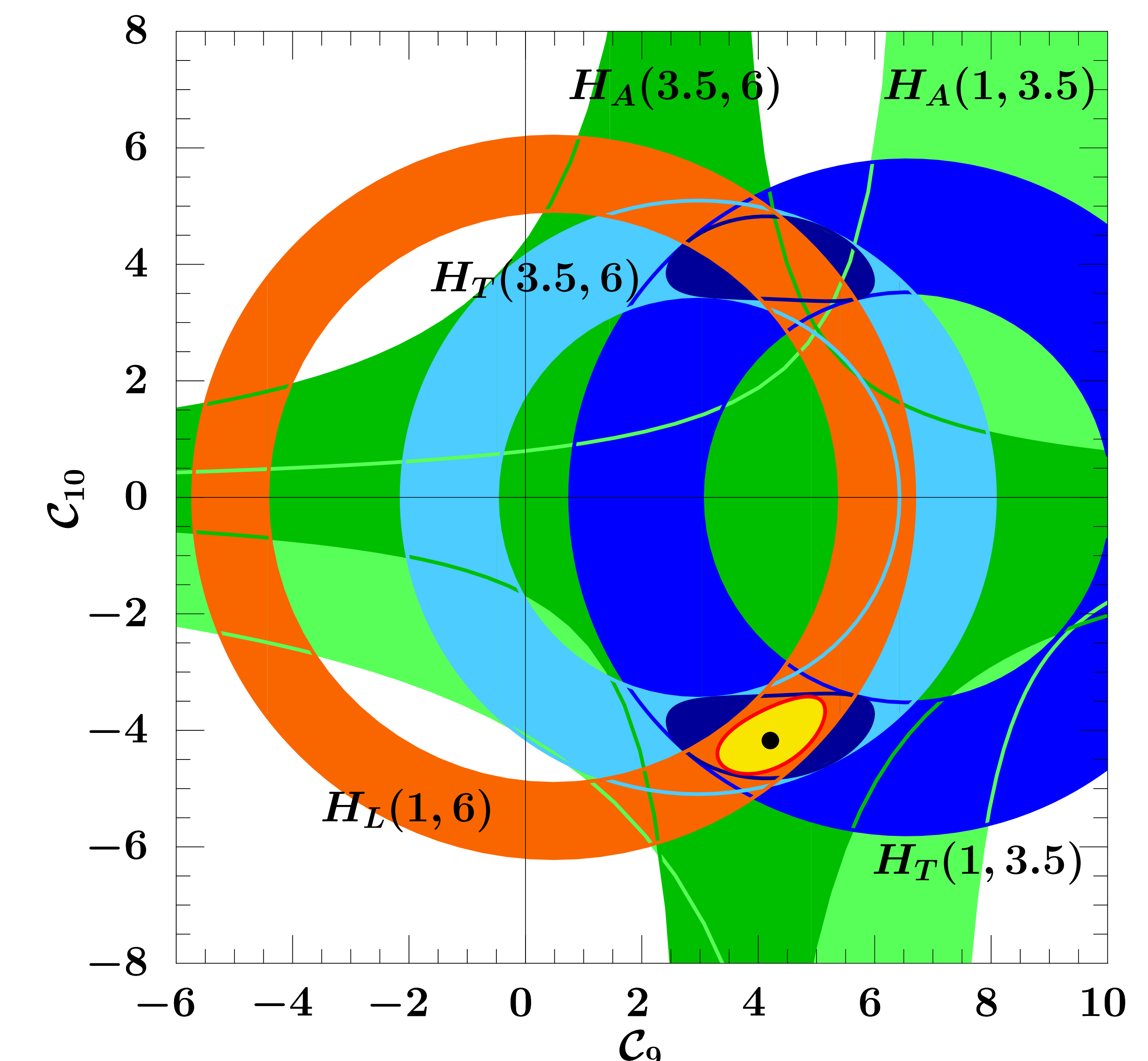
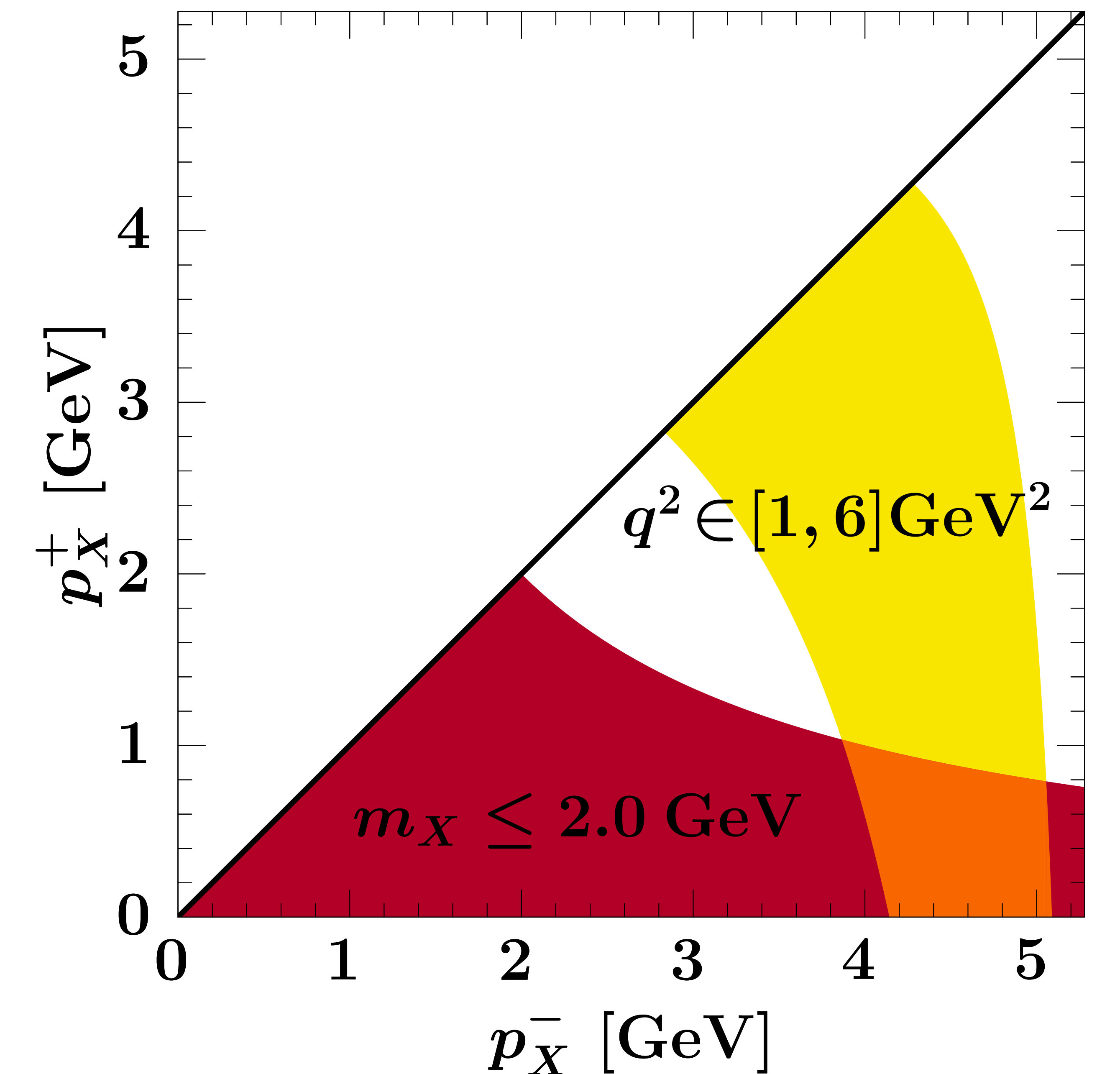
$$z = \cos \theta = 2 \frac{E_\ell - E_{\bar{\ell}}}{p_X^- - p_X^+} \quad : \text{ angle between lepton and B in } W \text{ rest frame}$$





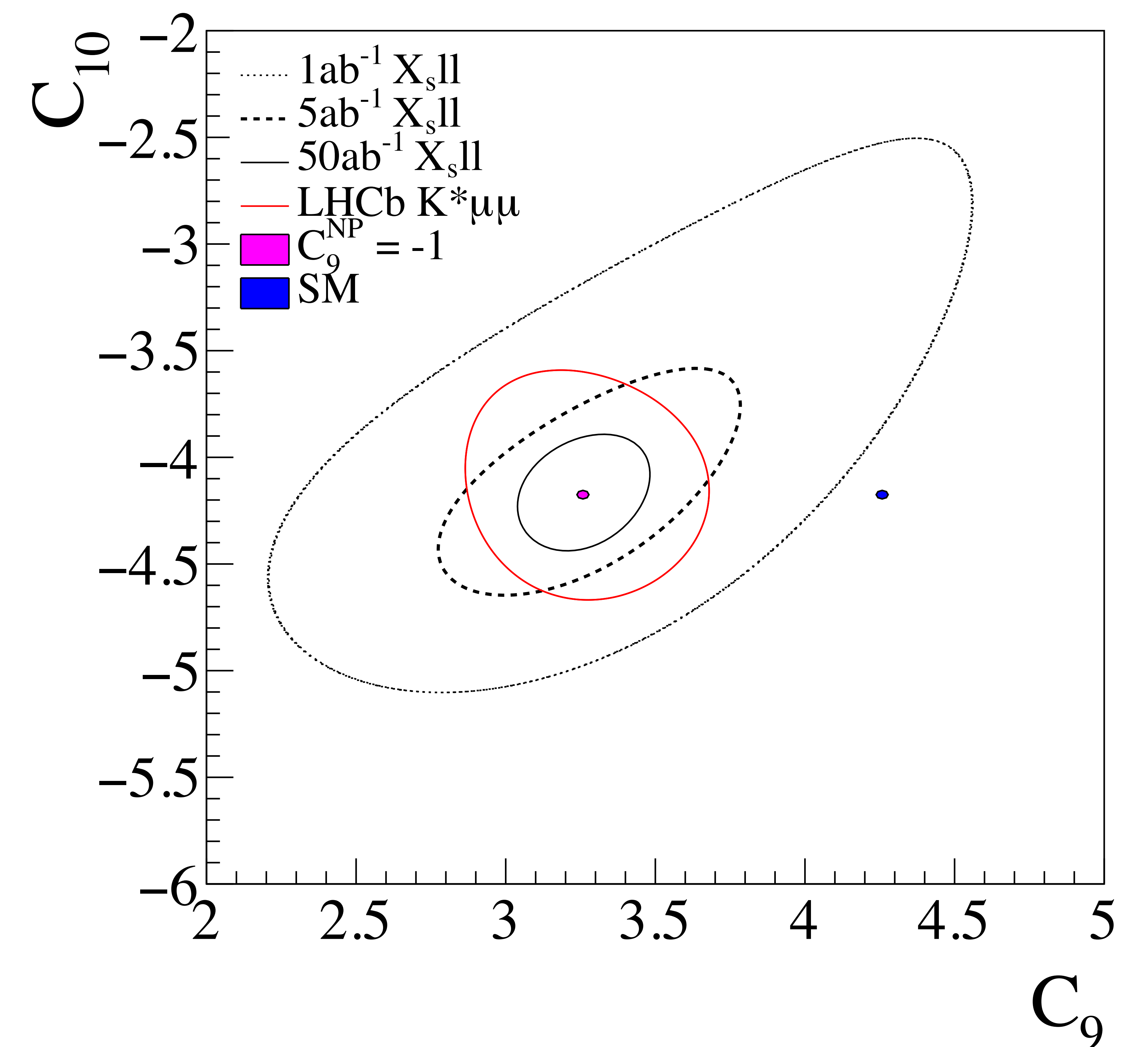
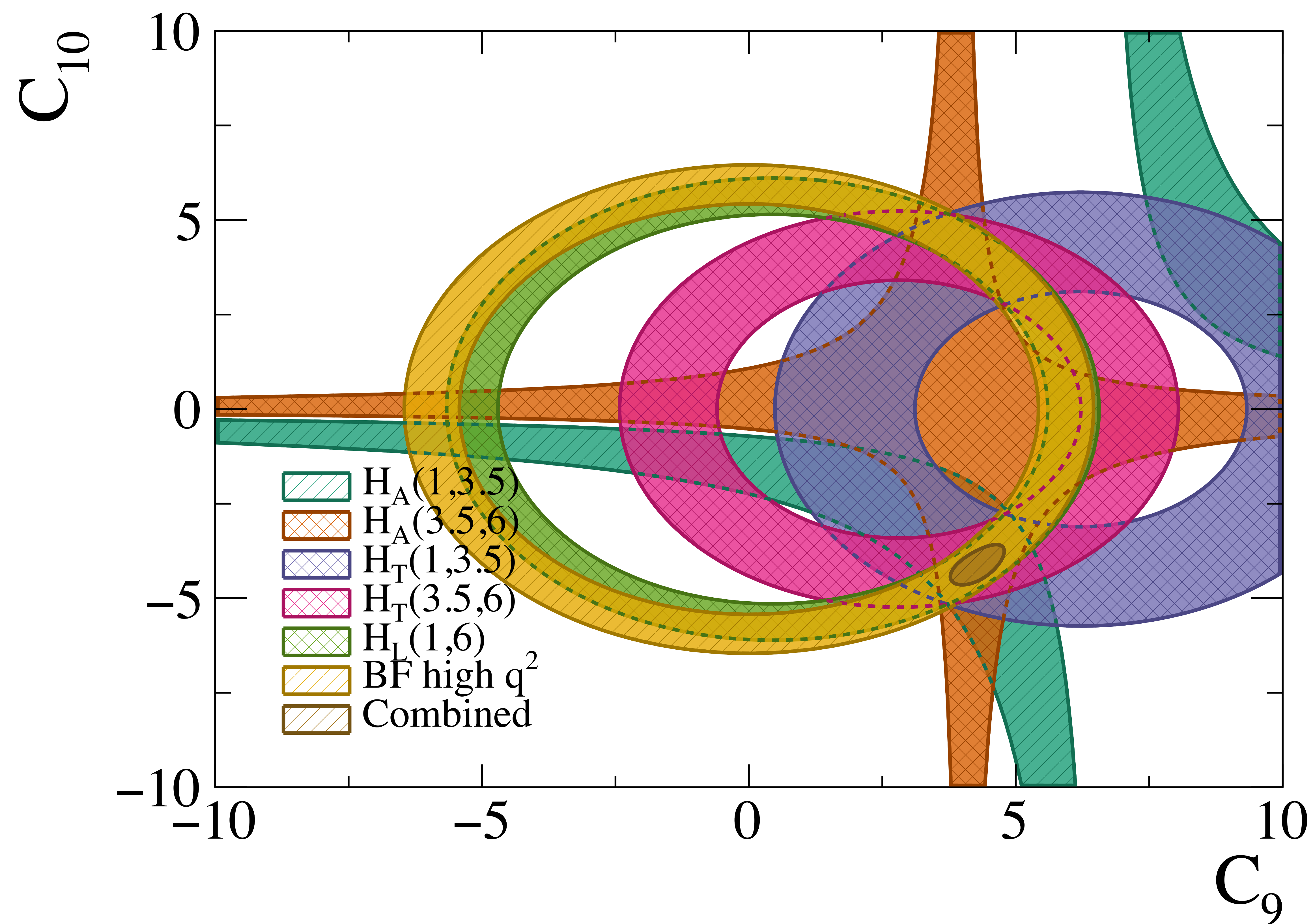
# $B \rightarrow X_s \ell \ell$

- Experimental kinematic cuts for  $B \rightarrow X_s \ell \ell$ 
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  - Unavoidable to suppress huge  $B \rightarrow X_c \ell \nu \rightarrow X_s \ell \ell \nu \nu$  background
  - Shape function effects must be taken into account to retain NP sensitivity
- Helicity decomposition for inclusive rate
  - [Lee, Ligeti, Stewart, Tackmann (2008)]
  - Same basic structure:
 
$$dH_{T,A,L} = \sum_{ij} C_i^{\text{incl}} C_j^{\text{incl}} \int dk \widehat{W}_{ij}^{A,T,L}(p_X^+, E_\ell, E_{\bar{\ell}}; k) \widehat{F}(p_X^+ - k) + \dots$$
  - Combined fit of  $B \rightarrow X_s \ell \ell$  and  $B \rightarrow X \ell \nu$ 
    - Best way to get clean extraction of  $C_9, C_{10}$  with inclusive decays



# $B \rightarrow X_s \ell \ell$

- Belle II sensitivity study from William Sutcliffe
  - Unbinned fit in  $z$  for 2  $q^2$  bins (1-3.5 GeV<sup>2</sup>, 3.5 - 6 GeV<sup>2</sup>)

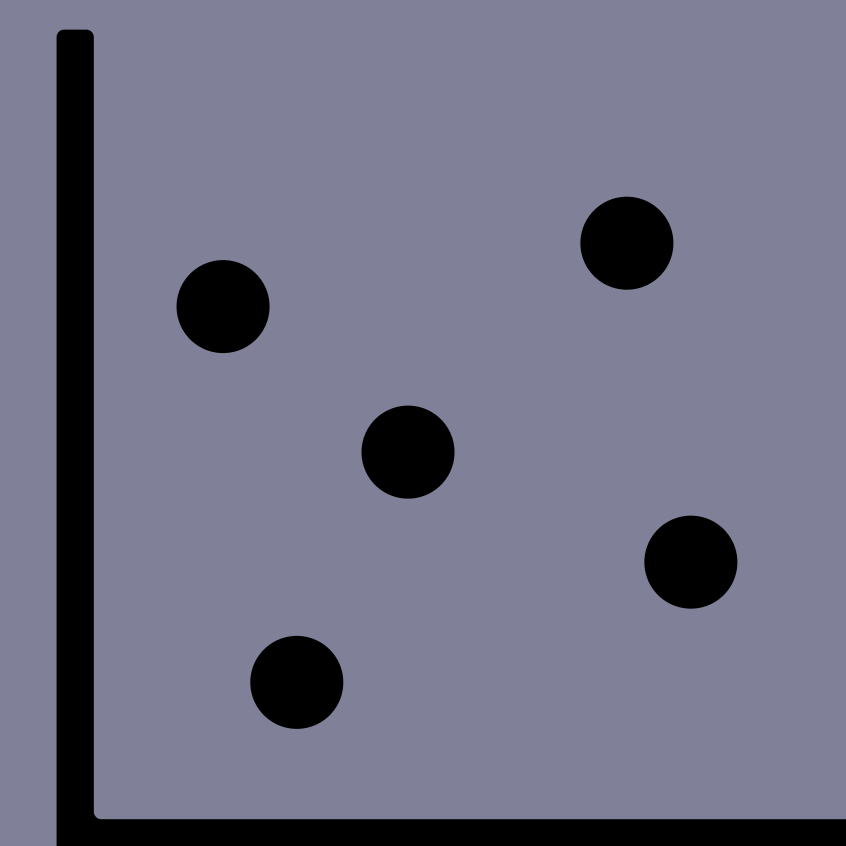




$$B \rightarrow X_c / v$$



$|V_{cb}|$  and  $R(X)$



# $B \rightarrow X_c \ell \nu$ & $B \rightarrow X \tau \nu$

- Combined analysis of  $B \rightarrow X_c \ell \nu$  and  $B \rightarrow X_u \ell \nu$ 
  - Measure very precisely the lepton energy spectrum
  - Allows for fully consistent and correlated treatment of both channels
    - Can constrain leading SF from  $b \rightarrow c$
  - Combined fit to directly extract  $|V_{ub}| / |V_{cb}|$
- $B \rightarrow X \tau \nu$  and  $R(X)$ 
  - Belle II should obviously measure  $R(X)$ 
    - if  $R(D^*)$  is due to NP, it must also manifest itself in  $R(X)$
  - Theory for inclusive decay is as clean
  - Combined analysis of  $B \rightarrow X \ell \nu$  and  $B \rightarrow X \tau \nu$  to measure  $R(X)(q^2)$



# Summary





- **Inclusive  $|V_{ub}|$  and  $|V_{cb}|$  with current approaches are **theory limited**, but not in a way that more calculations alone will help**
- Strategy for Belle II should be to exploit increased data sets **to help** theory by providing maximal amount of information in the form of differential branching fractions measured as model-independent as possible.
- **Global fit to inclusive rare and semileptonic data** with model-independent treatment of shape function will be **key to reach the ultimate** precision for **inclusive  $|V_{ub}|$**
- **Global analysis** will be **essential** to fully exploit **NP** sensitivity of **inclusive  $B \rightarrow X_s \gamma$  and  $B \rightarrow X_s \ell \ell$**





Differential Information

$\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$

$E_\ell$   
 $m_X$   
 $q^2$

+

$\bar{B} \rightarrow X_s \gamma$

$E_\gamma$

+

$\bar{B} \rightarrow X_s \ell \ell$

$q^2$

+

$\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$

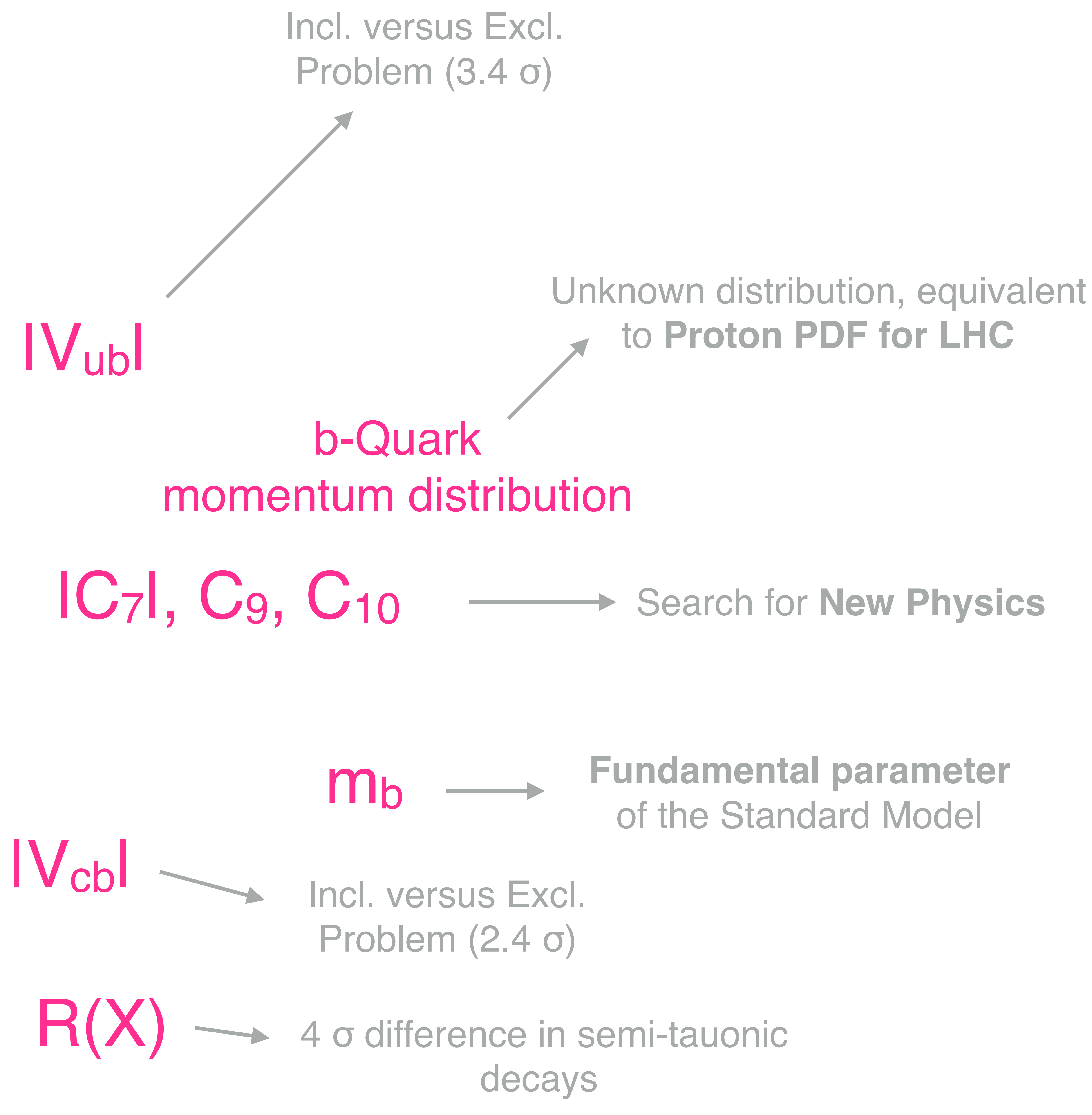
$E_\ell$   
 $q^2$

+

$\bar{B} \rightarrow X \tau \bar{\nu}_\tau$

$q^2$

Global Analysis



$$R(X) = \frac{\mathcal{B}(\bar{B} \rightarrow X \tau \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow X \ell \bar{\nu}_\ell)}$$